

# Chemical Response of the Upper Atmosphere due to Lightning-induced Electron Precipitation

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## Key Points:

- We quantify the direct effects via excess odd hydrogen production on the middle atmosphere ozone layer produced by lightning-induced electron precipitation
- The change in ozone concentration (75–85 km altitude) during one thunderstorm via LEP processes is comparable to other types of energetic particle precipitation
- The long-term global chemical effects produced by LEP events are uncertain and need to be properly quantified

14 **Abstract**

15 Terrestrial lightning frequently serves as a loss mechanism for energetic electrons in  
 16 the Van Allen radiation belts, leading to lightning-induced electron precipitation (LEP). Re-  
 17 gardless of the specific causes, energetic electron precipitation from the radiation belts in  
 18 general has a significant influence on the ozone concentration in the stratosphere and meso-  
 19 sphere. The atmospheric chemical effects induced by LEP have been previously investigated  
 20 using subionospheric VLF measurements at Faraday station, Antarctica ( $65.25^{\circ}$ S,  $64.27^{\circ}$ W,  
 21  $L = 2.45$ ). However, there exist large variations in the precipitation flux, ionization produc-  
 22 tion, and occurrence rate of LEP events depending on the peak current of the parent light-  
 23 ning discharge, as well as the season, location, and intensity of the thunderstorm activity.  
 24 These uncertainties motivate us to revisit the calculation of atmospheric chemical changes  
 25 produced by LEP. In this study, we combine a well-validated LEP model and first-principles  
 26 atmospheric chemical simulation, and investigate three intense storms in the year of 2013,  
 27 2015, and 2017 at the magnetic latitude of  $50.9^{\circ}$ ,  $32.1^{\circ}$ , and  $35.7^{\circ}$ , respectively. Modeling  
 28 results show that the LEP events in these storms can cumulatively drive significant changes  
 29 in the  $\text{NO}_x$ ,  $\text{HO}_x$ , and  $\text{O}_x$  concentration in the mesosphere. These changes are as high as  
 30  $\sim 156\%$ ,  $\sim 66\%$ , and  $-5\%$  at 75–85 km altitude, respectively, and comparable to the effects  
 31 typically induced by other types of radiation belt electron precipitation events. Considering  
 32 the high occurrence rate of thunderstorms around the globe, the long-term global chemical  
 33 effects produced by LEP events need to be properly quantified.

34 **1 Introduction**

35 The concept of terrestrial lightning discharge as a loss mechanism for energetic elec-  
 36 trons in the Van Allen radiation belts was first speculated by *Dungey* [1963], and later con-  
 37 firmed by direct measurements from the S81-1 (SEEP) satellite [Voss *et al.*, 1984, 1998].  
 38 This phenomenon is referred to as Lightning-induced Electron Precipitation (LEP), in which  
 39 the Very-Low-Frequency (VLF, 3–30 kiloHertz, kHz) waves emitted from a lightning dis-  
 40 charge propagate in the whistler mode through the Earth's magnetosphere, scatter energetic  
 41 electrons into lower mirroring altitudes in the radiation belts, and ultimately cause the precip-  
 42 itation of some trapped electrons into the upper atmosphere. As a strong coupling between  
 43 the Earth's atmosphere and magnetosphere, LEP has been the main focus of various obser-  
 44 vational studies [e.g., *Clilverd* *et al.*, 2002; *Rodger*, 2003; *Peter and Inan*, 2007; *Inan* *et al.*,  
 45 2010] in the past few decades. Even long before the first spacecraft measurements, LEP has  
 46 been investigated using radio measurements [*Helliwell* *et al.*, 1973; *Rycroft*, 1973; *Lohrey*  
 47 and *Kaiser*, 1979]. The most effective approach of LEP observation is via the indirect mea-  
 48 surements of ionospheric electron density enhancements using subionospheric VLF remote  
 49 sensing. VLF waves radiated from the Naval transmitters are well trapped within the wave-  
 50 guide formed between the ground and the sharp boundary in the lower ionosphere, and thus  
 51 particularly sensitive to the electron density in the D-region ionosphere ( $\sim 60$ – $90$  km), a re-  
 52 gion that is frequently bombarded by LEP fluxes.

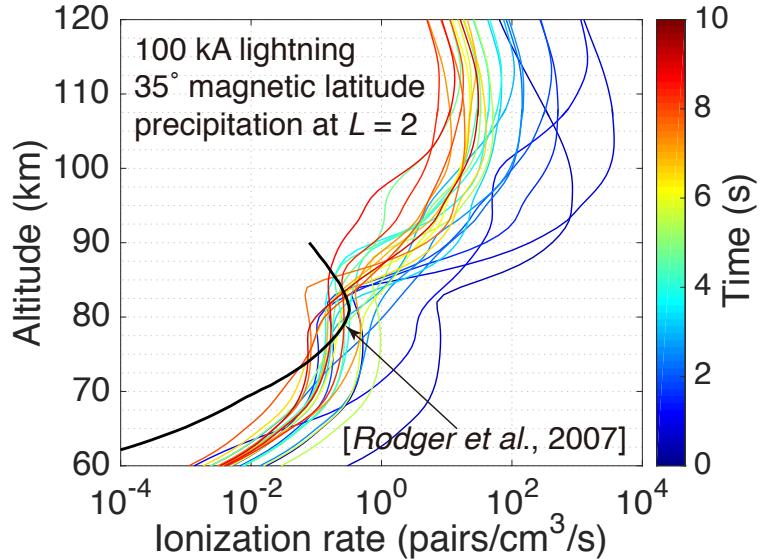
53 A major goal of previous and ongoing VLF observations is to estimate the size, fluxes,  
 54 and spectra of LEP events, and thereby quantify the effects of terrestrial lightning on the ra-  
 55 diation belt fluxes. Many of the pioneer works have been carried out by *Inan* *et al.* [1985];  
 56 *Inan and Carpenter* [1986, 1987]. The authors revealed that the phase and amplitude per-  
 57 turbations of VLF signals associated with LEP events (historically known as Trimpi events  
 58 [*Helliwell* *et al.*, 1973]), in most cases, can be explained using a whistler-induced precipita-  
 59 tion flux ranging from  $10^{-4}$  to  $10^{-2}$  erg/cm $^2$ /s. Using simultaneous measurements from mul-  
 60 tiple ground receivers, the size of LEP events has been estimated, by *Johnson* *et al.* [1999]  
 61 and *Clilverd* *et al.* [2002], to be as large as one thousand kilometers overhead the causative  
 62 lightning discharge. *Clilverd* *et al.* [2004] have further studied the relationship between VLF  
 63 perturbations and the peak current of the source lightning flashes, and pointed out that the  
 64 amplitude change of typical Trimpi events is consistent with lightning peak currents of 70–

250 kA. *Rodger et al.* [2005] have estimated the energy deposition via the LEP process into the middle atmosphere; a mean rate of energy deposition at  $L = 1.9\text{--}3.5$  was found to be  $3 \times 10^{-4}$  erg/cm $^2$ /min, with highs of  $6 \times 10^{-3}$  ergs/cm $^2$ /min above North America. Studies using the Holographic Array for Ionospheric/Lightning Research (HAIL) found that, for a 100-kA lightning discharge, the peak flux of precipitation electrons is on the order of 10 $^{-2}$  ergs/cm $^2$ /s [*Peter and Inan*, 2007]. In general, the spatial scale of lightning-induced electron precipitation ranges from several hundreds to one thousand kilometers, covering several degrees in latitude/longitude [e.g., *Clilverd et al.*, 2002]. As derived from ground VLF measurements, the displacement with respect to the lightning source is largely controlled by the geomagnetic field line, but primarily poleward shifted [e.g., *Peter and Inan*, 2007]. Lightning-generated whistler waves can lead to precipitation of energetic electron from both the inner and outer radiation belts: LEP events have been found to play a significant role in electron losses in the inner radiation belt [e.g., *Bortnik et al.*, 2006a,b; *Claude-pierre et al.*, 2020a,b]; LEP can also lead to electron losses from the outer radiation belt, for example, Trimpi events [*Helliwell et al.*, 1973].

Due to the indirect relationship between VLF perturbations and the underlying D-region electron density variation, quantification of LEP fluxes using VLF measurements is by nature a nonlinear problem [*Marshall et al.*, 2019b]. The amplitude and phase changes of transmitter VLF signals are controlled not solely by the electron density enhancement, but also by the geometry of the transmitter-receiver path, the ambient ionosphere along the path [e.g., *Xu et al.*, 2019], and the collision frequency profile driven by the background atmosphere [*Marshall*, 2012]. As such, the LEP fluxes as inversely derived from VLF measurements are inherently ambiguous, with large uncertainties in the energy spectrum in particular. Besides VLF technique, LEP fluxes have been directly measured by in situ particle instruments, for example, the Solar Anomalous and Magnetospheric Particle Explorer (SAMPEX) [*Blake et al.*, 2001] and the Detection of Electro-Magnetic Emissions Transmitted from Earthquake Regions (DEMETER) satellite [*Inan et al.*, 2007]. However, nearly all existing space-borne instruments can only resolve part of the loss cone angle [*Marshall et al.*, 2020] and these measurements only provide a coarse estimate of the true precipitating flux.

Regardless of the specific causes, energetic electron precipitation (EEP) into the Earth's atmosphere, in general, has a significant influence on the ozone concentration in the upper atmosphere [e.g., *Thorne*, 1980; *Randall et al.*, 2007; *Sinnhuber et al.*, 2012] through the catalytic cycles of odd nitrogen ( $\text{NO}_x = [\text{N}] + [\text{NO}] + [\text{NO}_2]$ ) [*Rusch et al.*, 1981] and odd hydrogen ( $[\text{HO}_x] = [\text{H}] + [\text{OH}] + [\text{HO}_2]$ ) [*Solomon et al.*, 1981]. Using 60 major EEP events measured during the solar cycle 23, *Andersson et al.* [2014a] revealed that EEP strongly affects the ozone concentration and can cause up to 90% depletion at altitudes of 60–80 km. *Turunen et al.* [2016] have studied the chemical changes during a pulsating aurora event on November 17, 2012, and found a maximum reduction of 14% in ozone concentration at 75 km altitude. As for the chemical effects of LEP, *Rodger et al.* [2007] have performed detailed atmospheric chemistry simulations using the mean LEP energy flux reported in *Rodger et al.* [2005], but rescaled using the Trimpi events observed at Faraday station, Antarctica (65.25°S, 64.27°W,  $L = 2.45$ ) on April 14, 1994. The maximum changes in  $\text{NO}_x$  and  $\text{HO}_x$  concentration were found to be ~0.1% around 80 km altitude, with a reduction of odd oxygen concentration by less than 0.1%, and thus the atmospheric chemistry effects were concluded to be insignificant [*Rodger et al.*, 2007].

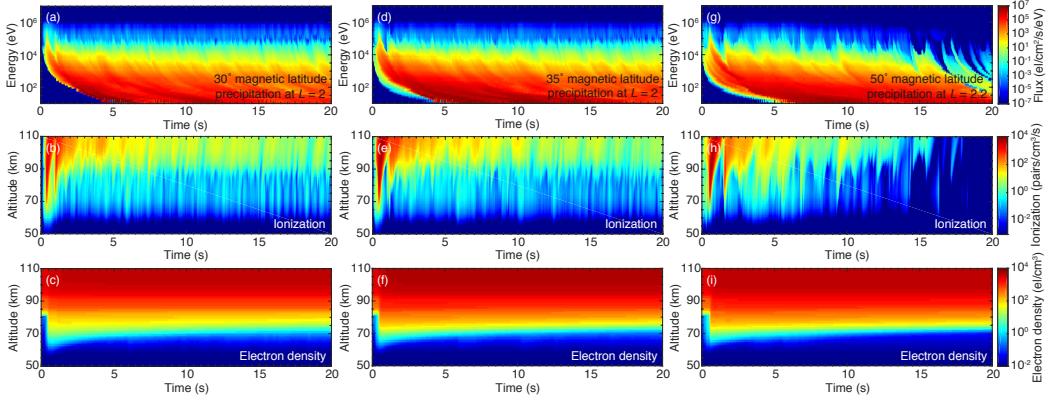
However, there exist large variations in the precipitation flux, ionization production, and occurrence rate of LEP events depending on the peak current of the parent lightning discharge, as well as the season, location, and intensity of thunderstorm activity [*Sousa*, 2018]. For example, Figure 1 shows the comparison of ionization production between the mean LEP ionization production used by *Rodger et al.* [2005]; *Rodger et al.* [2007], and our more recent LEP modeling results [*Marshall et al.*, 2019a]. The colored lines (color-coded using time) show the ionization rate versus altitude produced by an LEP event over the first 10 seconds after a 100-kA lightning discharge at 35° magnetic latitude, as calcu-



94 **Figure 1.** Comparison of the ionization production by lightning-induced electron precipitation between re-  
 95 cent LEP modeling results [Marshall *et al.*, 2019a] using the WIPP code [Lauben *et al.*, 1999; Bortnik, 2004;  
 96 Sousa, 2018] and the mean LEP ionization production [Rodger *et al.*, 2005; Rodger *et al.*, 2007]. The colored  
 97 lines show the altitude profile of ionization production at  $L = 2$  by a lightning source with a peak current of  
 98 100 kA at 35° magnetic latitude.

123 lated using the Stanford Wave-Induced Particle Precipitation (WIPP) code [Lauben *et al.*,  
 124 1999; Bortnik, 2004; Golden *et al.*, 2010; Sousa, 2018]. This code explicitly simulates from  
 125 first principles the entire LEP process from the source lightning discharge to precipitation  
 126 fluxes in the upper atmosphere. This model framework has been extensively used to analyze  
 127 LEP-associated VLF measurements [e.g., Peter and Inan, 2007; Inan *et al.*, 2010], and more  
 128 recently calibrated using X-ray measurements by the Balloon Array for Radiation-belt Rel-  
 129 ativistic Electron Losses (BARREL) during possible LEP events [Marshall *et al.*, 2019a].  
 130 In the first  $\sim$ 5 seconds, the mean LEP ionization production used by Rodger *et al.* [2007]  
 131 (possibly corresponding to a peak current smaller than 100 kA) is on average one order of  
 132 magnitude lower than that produced by this simulated 100-kA lightning discharge [Marshall  
 133 *et al.*, 2019a].

134 Another parameter that is critical for LEP-induced chemical effects is the occurrence  
 135 rate of intense lightning flashes. The flash rate was estimated to be approximately 3.3 events  
 136 per minute for the Trimpi measurements at the Faraday station on April 14, 1994 [Rodger  
 137 *et al.*, 2007]. A reexamination using the U.S. National Lightning Detection Network (NLDN)  
 138 data [Cummins *et al.*, 1998] reveals that this value is not representative of intense thunder-  
 139 storms at lower latitudes. As will be shown in section 3, the flash rate of lightning discharges  
 140 with peak current larger than 50 kA could be as high as  $\sim$ 50 per minute, as observed during  
 141 an intense thunderstorm occurring around 23.5°N, 97.5°W on October 25, 2015 (see Fig-  
 142 ure 3). In addition, the duration of a single LEP event was assumed to be  $\sim$ 0.2 s [Rodger  
 143 *et al.*, 2007] in the ducted case, whereas it can last up to 20 seconds or longer in the non-  
 144 ducted case [Bortnik, 2004; Marshall *et al.*, 2019a] due to multiple magnetospheric reflec-  
 145 tions between the conjugate hemispheres. Considering the high occurrence rate of LEP  
 146 events (nearly once per minute globally) and potential chemical effects, the uncertainties in  
 147 the LEP source (mostly from the uncertainties about the peak current of source lightning  
 148 discharge, energy and pitch angle distribution of precipitation fluxes) motivate us to revisit  
 149 the calculation of atmospheric chemical changes. In this paper, we present first-principles



171 **Figure 2.** (a) Differential flux of precipitation electrons with energies between 10 eV and 10 MeV at  $L = 2$   
 172 due to a 100 kA lightning source at  $30^\circ$  magnetic latitude. (b) Ionization production at altitudes between  
 173 50 and 110 km by these precipitation electrons. The background atmospheric profile used in this ionization  
 174 calculation is obtained using the date, latitude, and longitude of the 2015 storm reported by NLDN (see sec-  
 175 tion 3). (c) Electron density change in the D-region ionosphere produced by this LEP event. Panels (d–f)  
 176 show similar results, but for a lightning source at  $35^\circ$  magnetic latitude and precipitation fluxes at  $L = 2$ .  
 177 Panels (g–i) show the results for a lightning source at  $50^\circ$  magnetic latitude and precipitation fluxes at  $L = 2.2$ .  
 178 The background atmospheric profile used for the ionization calculation of panel (e) and (h) is obtained using  
 179 the date and location of the 2017 and 2013 storm (see section 3), respectively.

150 modeling results of LEP events, including the precipitation fluxes, ionization production,  
 151 and chemical changes. We use three NLDN-reported intense storms as extreme examples to  
 152 quantify the chemical effects produced by LEP.

## 153 2 Numerical Simulations

154 In this study, we combine the WIPP-LEP simulations of lightning-induced electron  
 155 precipitation [Lauben *et al.*, 1999; Bortnik, 2004; Sousa, 2018], the Boulder Electron Radia-  
 156 tion to Ionization (BERI) model [Xu *et al.*, 2020], and the Sodankylä Ion and Neutral Chem-  
 157 istry (SIC) model [Turunen *et al.*, 1996; Verronen *et al.*, 2005], specifically in three steps.  
 158 First, following the framework formulated by Lauben *et al.* [1999]; Bortnik [2004], the WIPP  
 159 model is employed to simulate LEP events produced by source lightning discharges at differ-  
 160 ent magnetic latitudes and calculate the resultant precipitation fluxes at different observation  
 161 locations ( $L$  values). Second, three intense storms are picked from the NLDN database for  
 162 the years of 2013–2017 at the magnetic latitudes of  $30^\circ$ – $50^\circ$ . Using the WIPP results ob-  
 163 tained in the first step, we calculate the total ionization production by the precipitation fluxes  
 164 induced by all lightning flashes in these storms (denoted as the cumulative ionization produc-  
 165 tion hereafter). Finally, the cumulative ionization production is utilized as an external forcing  
 166 in SIC simulations in order to quantify the atmospheric changes to constituents of interest.  
 167 Similar to previous EEP studies [Turunen *et al.*, 1996, 2009], the main focus of this study is  
 168 the relative change in the molecular concentration of odd hydrogen, odd nitrogen, and odd  
 169 oxygen ( $[O_x] = [O] + [O_3]$ ). In the following, we introduce the numerical models and the  
 170 initial parameters used in these simulations.

180 The WIPP code was built upon the modeling work of Inan [1977], and has been re-  
 181 fined through the past three decades of LEP modeling work at Stanford by Lauben *et al.*  
 182 [1999]; Bortnik [2004]; Golden *et al.* [2010]; Cotts [2011]; Sousa [2018]. The details of  
 183 this code, as well as the most recent updates, can be found in Sousa [2018]. In short, a stan-

184 dard WIPP-LEP simulation includes four steps [Bortnik, 2004; Sousa, 2018; Marshall *et al.*,  
 185 2019a]: i) The electromagnetic pulse (EMP) energy emitted by the return stroke current of  
 186 a lightning discharge is calculated and mapped to the base of the ionosphere at 100 km alti-  
 187 tude, ii) We calculate the attenuation of lightning-emitted VLF waves during their propaga-  
 188 tion through the lossy ionosphere (100–1000 km altitude) using the VLF attenuation curves  
 189 [Helliwell, 1965; Graf *et al.*, 2013a], iii) Starting from 1000 km altitude, the propagation of  
 190 each frequency component in the plasmasphere is simulated using the improved Stanford ray-  
 191 tracing program [Bortnik *et al.*, 2006a; Golden *et al.*, 2010]. Each ray is tracked through the  
 192 plasmasphere for 20 seconds due to multiple magnetospheric reflections, iv) We calculate  
 193 the resonant wave-particle interactions between these waves and the radiation belt fluxes, and  
 194 mainly focus on the modification to the pitch angles of trapped electrons near the loss cone  
 195 edge.

196 The fluxes and spectra of radiation belt electrons are calculated using the AE8 model  
 197 [Vette, 1991] and their pitch angles are assumed to be sine-distributed between 0° and 90°  
 198 [Sousa, 2018; Marshall *et al.*, 2019a]. Note that precipitation fluxes are strongly dependent  
 199 on the assumption of the background fluxes and pitch angle distributions; the most important  
 200 part of pitch angle distribution is the region near the loss cone angle; the uncertainty brought  
 201 by the assumption of pitch angle distribution has been previously discussed in Marshall *et al.*  
 202 [2019a]; our results were validated using Van Allen Probes data, as reported in Marshall  
 203 *et al.* [2019a]. WIPP simulations are performed for lightning discharges at magnetic lati-  
 204 tudes between 15° and 55° with 5° steps. For each magnetic latitude, the differential fluxes  
 205 of precipitation electrons with energies between 10 eV and 10 MeV are calculated for differ-  
 206 ent observation locations ( $L$  values). To quantify the maximum chemical effects, the location  
 207 with the highest precipitation flux (as denoted in the upper panels of Figure 2) is used to cal-  
 208 culate the ionization production and resultant chemical changes. It is important to emphasize  
 209 that in the WIPP simulations, the precipitation fluxes scale linearly with the total input EMP  
 210 energy, which is proportional to the square of lightning peak current. Following this relation,  
 211 WIPP-calculated LEP fluxes can be readily rescaled to the NLDN-reported peak current.  
 212 This relation is close to that empirically determined by Clilverd *et al.* [2004], wherein the  
 213 precipitation flux scales as the 2.3 power of lightning peak current.

214 Knowing the precipitation fluxes, we calculate the ionization production at altitudes  
 215 below 150 km altitude (the upper boundary of the SIC model) using the BERI model [Xu  
 216 *et al.*, 2020]. This model is largely based on a lookup table of atmospheric ionization pro-  
 217 duction by monoenergetic electrons with energies between 3 keV and 33 MeV, and pitch an-  
 218 gles between 0° and 90°. This lookup table was developed using physics-based Monte Carlo  
 219 simulations [e.g., Lehtinen *et al.*, 1999; Xu and Marshall, 2019], and allows rapid and accu-  
 220 rate specification of ionization production by arbitrary precipitation energy and pitch angle  
 221 distribution in any atmospheric condition. In this study, it is assumed that the precipitation  
 222 electrons of each LEP burst at 500 km altitude are isotropically distributed in pitch angles  
 223 between 0° and 90°. The mass density profile of background atmosphere is calculated using  
 224 the NRLMSISE-00 model [Tobiska and Bouwer, 2006] for the date, latitude, and longitude  
 225 of the storms reported by NLDN (see section 3).

226 Figure 2a shows WIPP modeling results of precipitation fluxes at  $L = 2$  produced by  
 227 a 100-kA lightning discharge at the magnetic latitude of 30°. The two peaks at  $\sim 0.4$  s and  
 228  $\sim 1$  s, as typical of satellite measurements of LEP events [Voss *et al.*, 1998], are caused by  
 229 the interaction between radiation belt electrons and the initial upward-going whistler waves  
 230 and the reflected whistler waves, respectively. The ionization production by these precipi-  
 231 tation electrons at altitudes between 50 and 110 km is shown in Figure 2b. The background  
 232 atmospheric profile used in this ionization calculation is obtained using the date, latitude,  
 233 and longitude of the 2015 storm reported by NLDN (see Figure 3). Using a 5-species chem-  
 234 istry model [Glukhov *et al.*, 1992; Lehtinen and Inan, 2007], we have further calculated the  
 235 electron density change in the D-region ionosphere, as shown in Figure 2c. Figures 2d–2f  
 236 show similar results, but for a lightning source at 35° magnetic latitude and precipitation

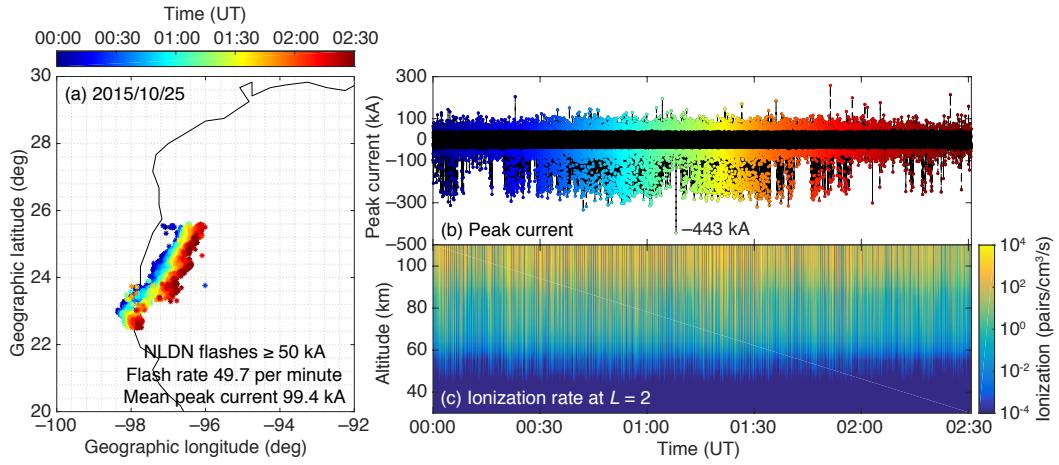
237 fluxes at  $L = 2$ , while Figures 2g–2i show those for a lightning source at  $50^\circ$  magnetic latitude  
 238 and precipitation fluxes at  $L = 2.2$ . The background atmospheric profile used for the  
 239 calculation of Figure 2e and Figure 2h is obtained using the date and location of the 2017  
 240 and 2013 storm (see section 3), respectively. Of note, the electron density variations shown  
 241 in the bottom panels of Figure 2 are capable of reproducing the typical amplitude changes  
 242 ( $\sim 0.5$ – $2$  dB) of transmitter VLF signal during LEP events [Peter and Inan, 2007]. Moreover,  
 243 the peak precipitation flux produced by the lightning discharge at  $30^\circ$ ,  $35^\circ$ , and  $50^\circ$  latitude  
 244 is  $0.06$ ,  $0.14$ , and  $0.15$  ergs/cm $^2$ /s, somewhat higher but not unreasonably different from the  
 245 values suggested by Peter and Inan [2007].

246 Because of the dense electron density in the D- and E-region ionosphere, the lightning  
 247 EMP energy is severely attenuated during the daytime ionospheric conditions [Graf *et al.*,  
 248 2013b] and the resultant precipitation fluxes are considerably lower than nighttime [Sousa,  
 249 2018]. Therefore, in this study, we mainly focus on those thunderstorms with high flash rates  
 250 and large peak currents, occurring during local nighttime conditions. Given these criteria,  
 251 three intense storms were chosen from NLDN data for the years of 2013, 2015, and 2017 at  
 252 the magnetic latitudes of  $50.9^\circ$ ,  $32.1^\circ$ , and  $35.7^\circ$ , respectively. The geolocation, temporal  
 253 evolution, and peak current of lightning flashes in these storms are presented in section 3.

254 The total ionization production during a thunderstorm is a key parameter in chemi-  
 255 cal simulations, and is calculated using all the flashes with peak current larger than  $50$  kA  
 256 in the present study. A minimum value of  $50$  kA is used since it is close to what is needed  
 257 to trigger Trimpie events ( $70$  kA) with detectable ionosphere enhancements [Clilverd *et al.*,  
 258 2004]. Clilverd *et al.* [2002] also noted that, if the lightning peak current is less than  $45$  kA,  
 259 the chance of observing any Trimpie events is almost zero. We have checked that, if a lower  
 260 threshold value is instead used, the cumulative ionization production would not change sig-  
 261 nificantly since the LEP flux scales linearly with the square of lightning peak current. Specif-  
 262 ically, for a given storm, we use the WIPP results at the corresponding magnetic latitude (up-  
 263 per panels of Figure 2) and rescale the ionization results (middle panels of Figure 2) using  
 264 the peak current of all NLDN-reported lightning flashes ( $\geq 50$  kA). The rescaled ionization  
 265 production is then sorted using the NLDN-tagged time of each flash (see Figure 3b). Finally,  
 266 these ionization results are summed together and we calculate the cumulative ionization pro-  
 267 duction versus altitude and time for each storm (see Figure 3c).

268 The cumulative ionization production is then used as an input into SIC chemistry sim-  
 269 ulations. SIC is a 1-D atmospheric model that dynamically solves for the concentration of  
 270 16 minor neutral species and 72 ionic species in the altitude range between  $20$  and  $150$  km  
 271 with  $1$  km resolution [Turunen *et al.*, 1996; Verronen *et al.*, 2005; Verronen, 2006]. Ver-  
 272 tical motion of species is included as molecular and eddy diffusion, neglecting transport  
 273 by prevailing neutral winds. The latest version of this model takes into account 389 ion-  
 274 neutral and neutral-neutral reactions, and 2523 ion-ion and electron-ion recombination re-  
 275 actions. The background profile of neutral density used in SIC modeling is obtained from the  
 276 NRLMSISE-00 model [Tobiska and Bouwer, 2006] using the daily average values of solar  
 277 radio flux ( $F_{10.7}$ ) and the geomagnetic activity index ( $A_p$ ). Note that horizontal mixing is not  
 278 included in the 1-D SIC model; this effect will be investigated in our next-step study using  
 279 global circulation simulations.

280 In this study, the neutral density profile is calculated using the specific date and loca-  
 281 tion of each storm as reported by NLDN. Solar proton precipitation is provided as an op-  
 282 tional external force in the SIC model, but not included in present simulations. The initial  
 283 profiles of  $\text{HO}_x$ ,  $\text{NO}_x$ ,  $\text{O}_x$  are obtained by running the SIC model at the thunderstorm loca-  
 284 tion for 5 days, prior to the LEP forcing with photoionization only. Chemical changes are  
 285 simulated for a period of 24 hours starting from the first LEP event and stored every 1 minute  
 286 of simulation. An average ionization rate of the event is used. Two sets of SIC simulations  
 287 are performed for each storm, one with and another without the LEP-induced ionization pro-  
 288 duction. The simulation results obtained without applying LEP ionization are regarded as



**Figure 3.** NLDN measurements of lightning flashes between  $\sim$ 00:00 and  $\sim$ 02:30 UT on October 25, 2015 near  $23.5^{\circ}$ N,  $97.5^{\circ}$ W. (a) Longitude and latitude of all lightning flashes with peak current magnitude larger than 50 kA. The rate of lightning flashes with peak current magnitude larger than 50 kA is  $\sim$ 49.7 flashes per minute and the average value of peak current for these flashes is  $\sim$ 99.4 kA. (b) Peak current of lightning flashes versus the occurrence time. The largest peak current recorded by NLDN was  $-443$  kA at 01:08:08 UT. (c) Altitude profile of ionization production by the lightning flashes with peak current magnitude larger than 50 kA shown in panel (b).

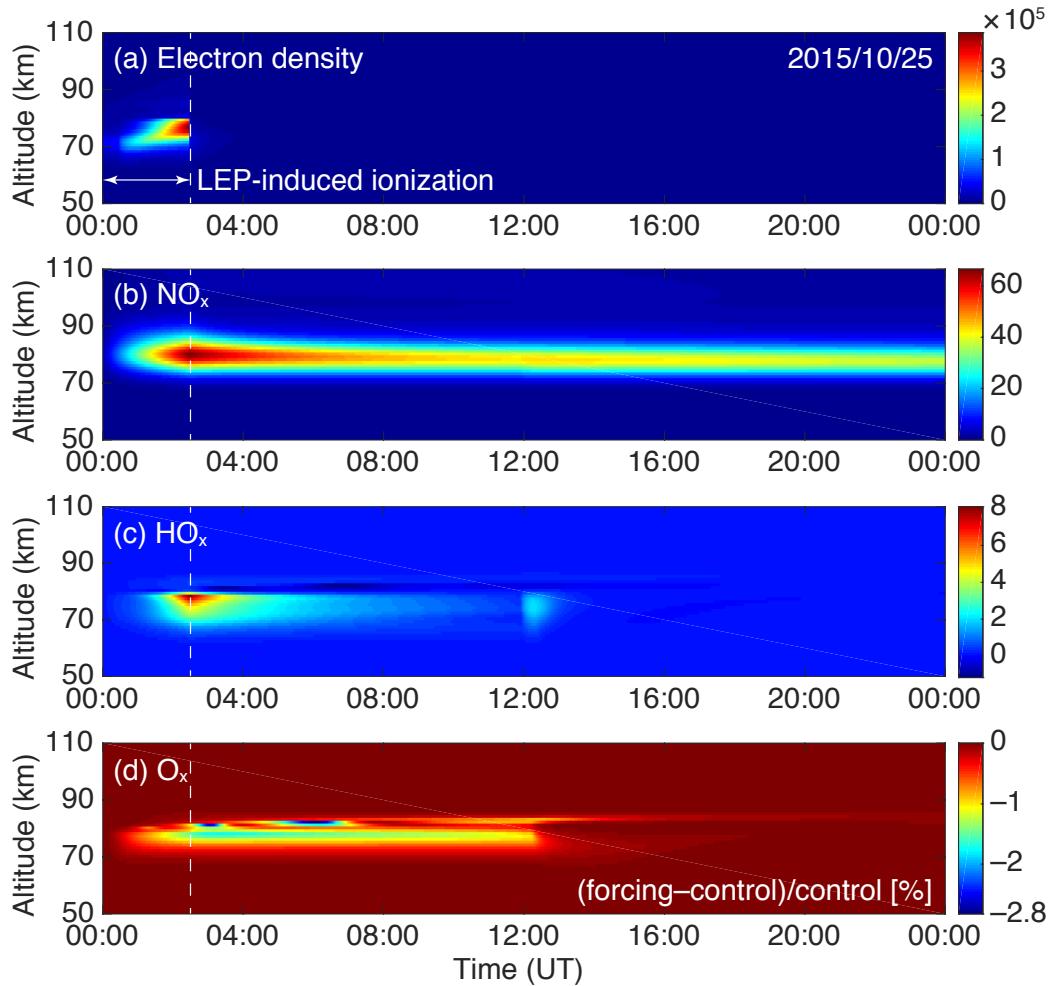
the baseline, against which we compute the relative changes in neutral species and therefore quantify the chemical effects.

We emphasize that all above-mentioned numerical models have been well validated in previous studies. The WIPP code has been extensively used to interpret LEP-modulated VLF signals [e.g., Peter and Inan, 2007], and lately the X-ray fluxes recorded by BARREL during possible LEP events [Marshall *et al.*, 2019a]. Different from VLF measurements, X-ray measurements at balloon altitudes are directly linked to the precipitation fluxes and energy spectra of LEP bursts; WIPP results can fully explain the X-ray fluxes, temporal signature, and energy spectra measured by BARREL [Marshall *et al.*, 2019a]. As for the BERI model [Xu *et al.*, 2020], it shows good agreements with the parameterization method of Fang *et al.* [2010] in terms of the peak ionization rate and altitude, with a maximum difference of  $\sim$ 20% among tests using different precipitation energy and pitch angle distributions. The SIC model has been employed for the estimation of atmospheric chemical effects due to a wide variety of external sources, including radiation belt electron precipitation [e.g., Turunen *et al.*, 2016; Xu *et al.*, 2018], solar eclipse [e.g., Xu *et al.*, 2019], and solar proton events [e.g., Clilverd *et al.*, 2005].

### 3 Results

#### 3.1 Storm 1: October 25, 2015

The first storm in this case study occurred on October 25, 2015 at geographic latitudes between  $22.5^{\circ}$ N and  $25.6^{\circ}$ N, and geographic longitudes between  $96.0^{\circ}$ W and  $98.5^{\circ}$ W, along the Caribbean coast of Texas / Mexico. Figure 3a shows the longitudes and latitudes of all lightning flashes with peak current magnitude larger than 50 kA recorded by the NLDN network between  $\sim$ 00:00 and  $\sim$ 02:30 UT, with color progressing in time from blue to red. The peak current and occurrence time of these lightning flashes are shown separately in Fig-



313 **Figure 4.** SIC modeling results of the relative changes in (a) electron, (b)  $\text{NO}_x$ , (c)  $\text{HO}_x$ , and (d)  $\text{O}_x$   
 314 concentration produced by all the LEP events in the 2015 storm. Sunrise time: 12:30:27 UT. This figure  
 315 specifically shows the percentage difference between simulation results obtained with (denoted as forcing)  
 316 and without (denoted as control) applying the cumulative ionization production (Figure 3c). The cumulative  
 317 ionization production between 00:00 and 02:30 UT is marked using white dashed lines in these panels.

ure 3b. The black and colored dots show the lightning flashes with peak current magnitude smaller and larger than 50 kA, respectively.

NLDN reported a total of 33,504 flashes from this region between 00:00 and 02:30 UT, ~22.3% of which had peak current larger than 50 kA (7,453 flashes). The flash rate of intense lightning discharges ( $\geq 50$  kA) was approximately 49.7 per minute. The average value of the peak current for these 7,453 flashes was ~99.4 kA and the majority were negative cloud-to-ground discharges. Out of the 33,504 flashes, the fraction of lightning flashes with peak current larger than 70 kA and 100 kA was approximately 11.2% (3,748 flashes) and 6.4% (2,157 flashes), respectively; the largest peak current was -443 kA at 01:08:08 UT, likely associated with large-scale high-altitude luminous events, for example, elves [e.g., Marshall *et al.*, 2010] or sprites [e.g., Pasko *et al.*, 1997].

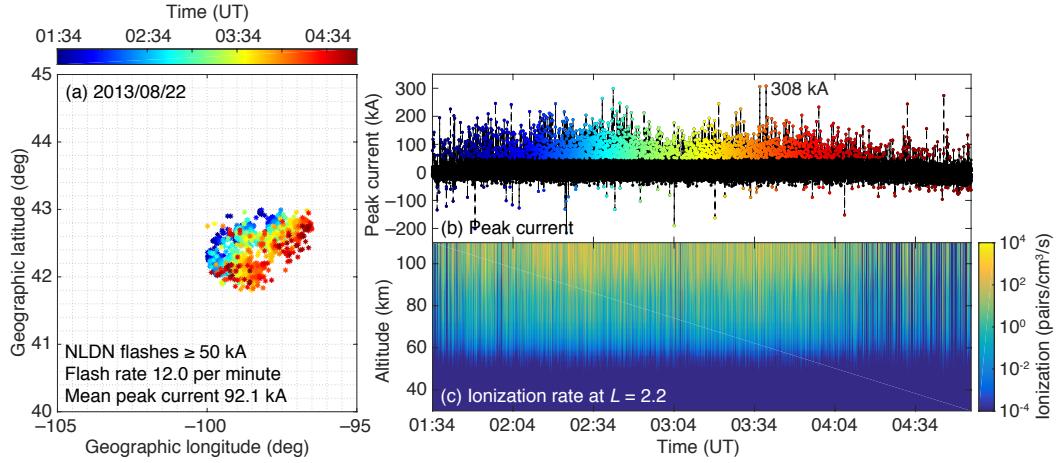
The magnetic latitude corresponding to the center of this storm is approximately 32.1°N, as calculated using the International Geomagnetic Reference Field (IGRF) model [Thébault *et al.*, 2015]. Thus, the WIPP-calculated precipitation fluxes produced by the lightning discharge at 30° magnetic latitude (Figures 2a–2b) are used for the ionization calculation. The LEP-produced ionization profile in Figure 2b is rescaled using the peak current of the lightning flashes (colored dots with peak current  $\geq 50$  kA) shown in Figure 3b, sorted using the NLDN-tagged occurrence time (Figure 3b), and then summed together. Figure 3c shows the cumulative ionization production via all LEP processes during this 150-minute storm.

These ionization results are then used as the forcing input to the SIC model in order to calculate the atmospheric chemistry response. Figure 4, from top to bottom, presents the relative change in the concentration of electrons,  $\text{NO}_x$ ,  $\text{HO}_x$ , and  $\text{O}_x$  produced by all the LEP events in storm 1. This figure specifically shows the percentage difference between the simulation results obtained with (denoted as forcing in Figure 4) and without (denoted as control in Figure 4) applying the cumulative ionization production during this storm ((forcing – control) / control  $\times 100$  [%]). The cumulative ionization production between 00:00 and 02:30 UT is marked using white dashed lines in this figure. The electron density is dramatically enhanced, by three orders of magnitude in the D-region ionosphere, as seen in Figure 4a. This level of electron density change is close to previously reported results [e.g., Peter, 2007, Figure 5.2]. The maximum change of  $\text{NO}_x$  concentration, due to the LEP-induced ionization, is approximately 67% at ~80 km altitude.  $\text{NO}_x$  is relatively stable and these changes can persist for a long time; about 37% of the excess  $\text{NO}_x$  production remains at the end of this 24-hour simulation period (see Figure 4b).

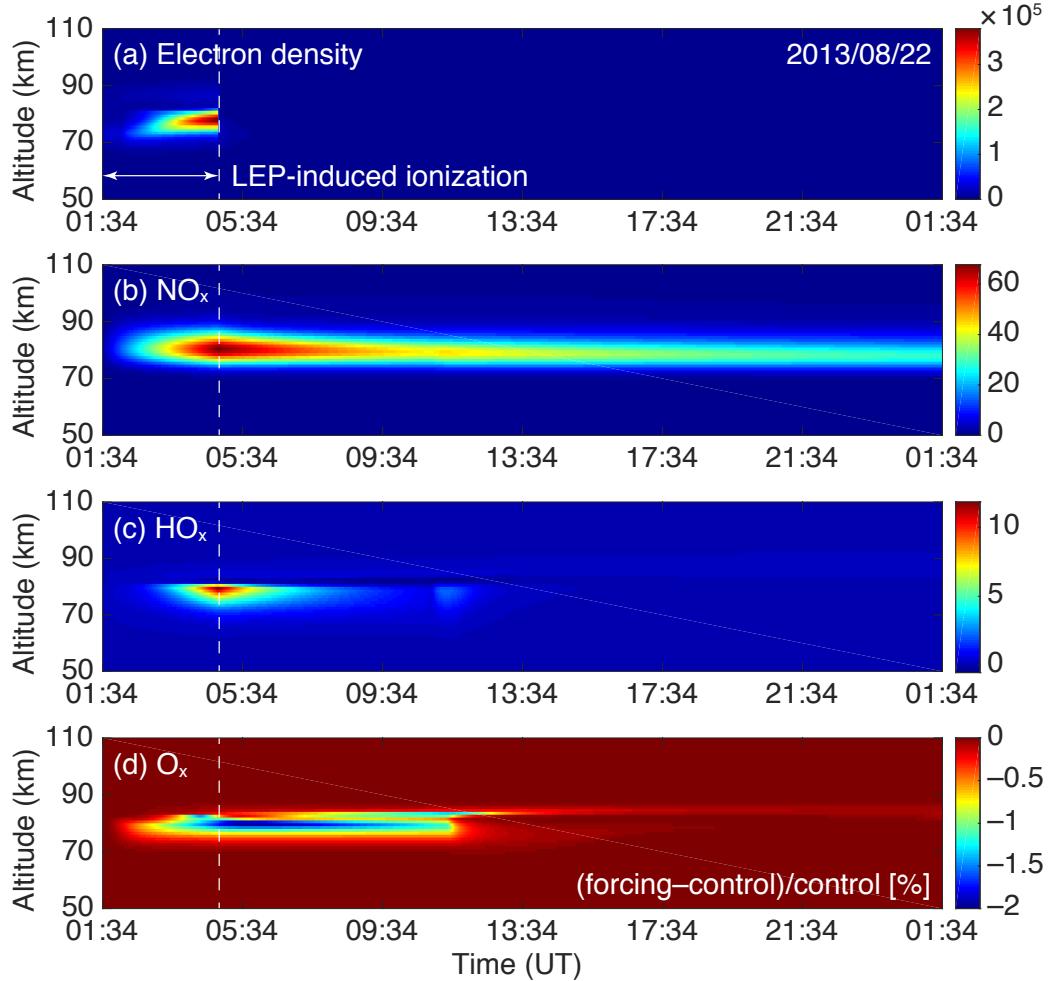
Different from  $\text{NO}_x$ , the  $\text{HO}_x$  change is considerably smaller and non-monotonic.  $\text{HO}_x$  concentration increases by ~8.1% at the altitudes near 78 km, and decreases by ~1.2% at 82 km altitude after the LEP forcing at ~06:52 UT. Above 82 km, the  $\text{HO}_x$  change is very small because of the limited abundance of water vapor in this altitude range [Turunen *et al.*, 2016]. A reduction of 2.8% is predicted for the  $\text{O}_x$  density around the local minimum in the mesospheric ozone profile, ~82 km altitude. However, because of the solar radiation and enhanced photochemistry during the sunrise [Verronen, 2006], the  $\text{HO}_x$  concentration is shortly enhanced (see Figure 4c) and the  $\text{O}_x$  concentration returns to the background value around ~12:00 UT (see Figure 4d).

### 3.2 Storm 2: August 22, 2013

Lightning at higher magnetic latitudes projects to higher  $L$ -shells in the radiation belts, and thus has the potential to impact fluxes in the heart of the radiation belts and produce more intense precipitation signatures. To investigate the LEP effects at higher latitudes, the second case study is performed for the thunderstorm occurring on August 22, 2013 at geographic latitudes between 41.8°N and 43.0°N, and longitudes between 96.5°W and 100.0°W, in northeastern Nebraska. Nebraska is well known for producing uncommonly intense positive lightning discharges [e.g., Stolzenburg, 1994]. The magnetic latitude of this storm is ~50.9°N and the NLDN lightning data from 01:34 to 04:54 UT are used. Similar to Fig-



367 **Figure 5.** Similar to Figure 3, but for the NLDN-reported lightning flashes near  $42.4^{\circ}\text{N}$ ,  $98.7^{\circ}\text{W}$  between  
368 01:34 and 04:54 UT on August 22, 2013.



369 **Figure 6.** Similar to Figure 4, but for the NLDN-reported lightning flashes near  $42.4^{\circ}\text{N}$ ,  $98.7^{\circ}\text{W}$  between  
370 01:34 and 04:54 UT on August 22, 2013. Sunrise time: 11:51:14 UT.

380 ure 3, Figure 5a shows the temporal evolution of lightning flashes with peak current mag-  
 381 nitude larger than 50 kA and Figure 5b shows the peak current versus occurrence time for  
 382 these flashes. The WIPP results corresponding to a lightning source at 50° magnetic latitude  
 383 (Figure 2h) are utilized for the calculation of cumulative ionization production, as shown in  
 384 Figure 5c.

385 Compared to the baseline runs, the relative change in  $\text{NO}_x$ ,  $\text{HO}_x$ , and  $\text{O}_x$  concentration  
 386 is  $\sim 67\%$ ,  $\sim 12\%$ , and  $-2\%$ , respectively. The  $\text{O}_x$  change in this case is smaller than that of  
 387 the first storm, the  $\text{NO}_x$  change is comparable, and the  $\text{HO}_x$  change is slightly higher. Note  
 388 that the relative change of these neutral species is somewhat sensitive to the baseline condi-  
 389 tions, for example, the season, location, and background atmospheric condition of baseline  
 390 simulations. On average, the cumulative ionization production in the second storm is notably  
 391 less than that of the first storm since the peak current and flash rate are lower. In the second  
 392 storm, 2,390 lightning flashes with peak current larger than 50 kA were detected by NLDN  
 393 and the flash rate was  $\sim 12.0$  events per minute, a quarter of that for the first storm. The av-  
 394 erage value of peak current for these flashes ( $\geq 50$  kA) was  $\sim 92.1$  kA, which is also 7.3%  
 395 lower than the first storm (99.4 kA). Out of these 2,390 flashes, the number of flashes with  
 396 peak current larger than 70 kA and 100 kA was 1,585 and 775, respectively, and the largest  
 397 peak current reported by NLDN was 308 kA.

398 Atmospheric chemical changes, in essence, are positively related with the cumula-  
 399 tive ionization production during a thunderstorm for a given atmospheric condition, which  
 400 is largely controlled by the LEP fluxes and lightning flash rate if the dependence on the pre-  
 401 cipitation energy spectrum is not considered. As explained in section 2, the precipitation flux  
 402 of a single LEP event is linearly proportional to the square of lightning peak current. The  
 403 lightning flash rate can enhance or diminish the cumulative effects of ionization production  
 404 during a thunderstorm. Thus, these two parameters can be roughly considered as a proxy for  
 405 the extent of ionization and chemical effects produced by thunderstorm activity via LEP pro-  
 406 cesses. It is important to note that the fluxes of LEP bursts are also dependent on the  $L$ -shell  
 407 from which lightning whistler induces electron precipitation, i.e., the availability of energetic  
 408 electrons in the radiation belts, as well as their pitch angle distribution.

### 409 3.3 Storm 3: May 29, 2017

410 For completeness of this case study, a third storm is chosen at a magnetic latitude be-  
 411 tween the first two storms, at  $\sim 35.7^\circ\text{N}$ . This storm took place on May 29, 2017 at geographic  
 412 latitudes between  $25.5^\circ\text{N}$  and  $28.2^\circ\text{N}$ , and longitudes between  $96.0^\circ\text{W}$  and  $100.5^\circ\text{W}$ , near  
 413 the U.S. and Mexico border along the Caribbean coast. The NLDN data between 04:29 and  
 414 09:59 UT are used for the chemistry simulation. The geolocation, peak current, and ion-  
 415 ization production by the lightning flashes in this storm are shown in Figure 7. A total of  
 416 440,266 lightning flashes were identified by NLDN to originate from this storm, and the  
 417 fraction of lightning discharges with peak current greater than 50 kA, 70 kA, and 100 kA  
 418 was 2.5% (11,148 flashes), 0.9% (3,971 flashes), and 0.4% (1,753 flashes), respectively.  
 419 These flashes were mostly negative cloud-to-ground discharges. The rate of lightning flashes  
 420 with peak current larger than 50 kA was 33.8 flashes per minute, with the average value of  
 421 the peak current (flashes  $\geq 50$  kA) being 77.6 kA and the largest value being 292 kA.

422 Given the NLDN-reported flash rate and peak current, it is not unexpected that this  
 423 storm leads to the largest chemical changes among all cases, as evidenced in Figure 8. SIC  
 424 modeling results show that LEP-induced ionization results in notable  $\text{NO}_x$  and  $\text{HO}_x$  changes:  
 425 the  $\text{NO}_x$  concentration is more than doubled at altitudes between  $\sim 76$  and  $\sim 84$  km, with  
 426 a maximum enhancement of 156% at  $\sim 80$  km; the  $\text{HO}_x$  concentration increases by  $\sim 66\%$   
 427 around 78 km altitude compared to the control runs. The  $\text{O}_x$  change in this case closely fol-  
 428 lows that of  $\text{HO}_x$ : at altitudes between  $\sim 77$  and  $\sim 83$  km, the  $\text{O}_x$  concentration reduces by  
 429 more than  $\sim 3\%$  and the maximum reduction is approximately 5% at 79 km altitude.

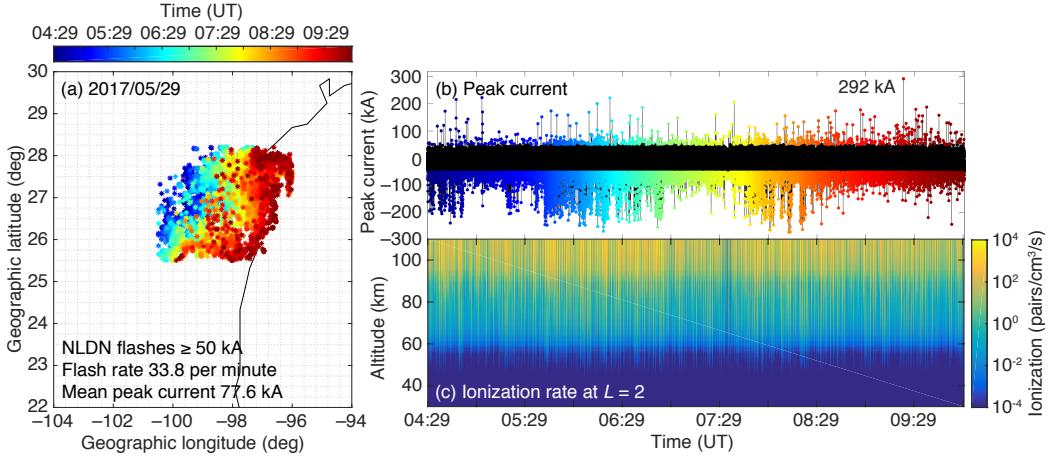


Figure 7. Similar to Figure 3, but for the NLDN-reported lightning flashes near  $27.3^{\circ}\text{N}$ ,  $98.3^{\circ}\text{W}$  between 04:29 and 09:59 UT on May 29, 2017.

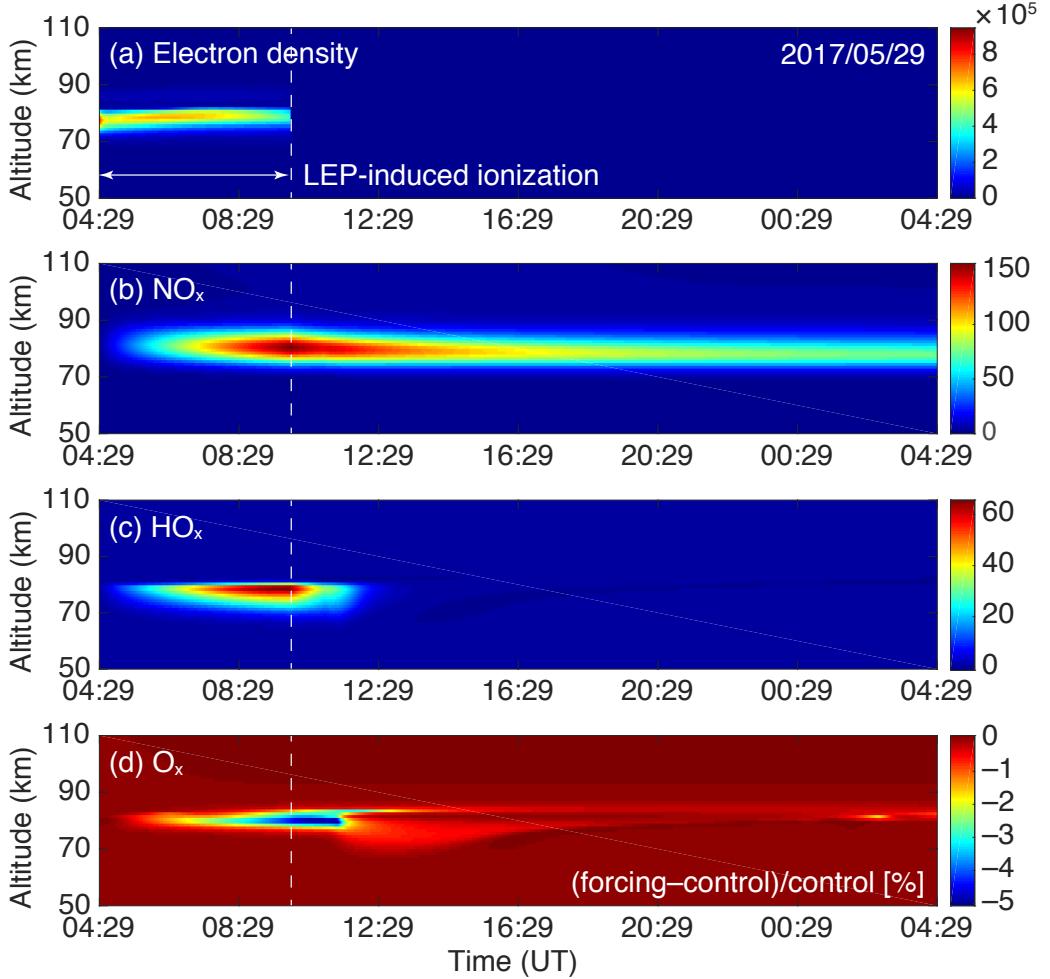


Figure 8. Similar to Figure 4, but for the NLDN-reported lightning flashes near  $27.3^{\circ}\text{N}$ ,  $98.3^{\circ}\text{W}$  between 04:29 and 09:59 UT on May 29, 2017. Sunrise time: 11:42:58 UT.

#### 434 4 Conclusion and Discussion

435 In this study, using a suite of well-validated LEP and atmospheric ionization models,  
 436 we have calculated the precipitation fluxes and ionization production by lightning flashes at  
 437 different magnetic latitudes. Case studies for three intense storms in the years from 2013 to  
 438 2017 at magnetic latitudes of  $30^{\circ}$ – $50^{\circ}$  are performed using NLDN-reported lightning data.  
 439 Using SIC modeling of atmospheric changes, we have further quantified the relative changes  
 440 in the electron,  $\text{NO}_x$ ,  $\text{HO}_x$ , and  $\text{O}_x$  concentration due to the LEP events in these storms.

441 Because of LEP-induced ionization, the  $\text{NO}_x$  and  $\text{HO}_x$  concentration at altitudes be-  
 442 tween 75 and 85 km is enhanced by up to  $\sim 156\%$  and  $\sim 66\%$ , respectively, during these storms.  
 443 The maximum reduction in ozone concentration is approximately 5%, as driven by the cat-  
 444 alytic reaction cycles of  $\text{HO}_x$ . These atmospheric changes are one order of magnitude larger  
 445 than those suggested by Rodger *et al.* [2007], mainly because of the variation in lightning  
 446 flash rate and thunderstorm intensity. The mean LEP ionization production used by Rodger  
 447 *et al.* [2007] was calculated using a mean precipitation energy flux of  $2 \times 10^{-3}$  ergs/cm $^2$ /s  
 448 [Rodger *et al.*, 2005], as required to explain Faraday Trimpi measurements. The chemical  
 449 simulation of Rodger *et al.* [2007] was conducted using the Trimpi events observed at the  
 450 Faraday station on April 14, 1994, corresponding to a Trimpi rate of 3.3 events per minute.  
 451 However, as shown in section 3, the NLDN-reported lightning peak current and flash rate for  
 452 the three storms reported herein are orders of magnitude higher.

453 We emphasize that the three storms investigated in this study do not represent the most  
 454 intense cases on a global scale, albeit stronger than majority of the thunderstorm activity in  
 455 North America. According to space-borne measurements of lightning activity, north and  
 456 central Argentina is the region that hosts the most intense convective storms on the Earth  
 457 [e.g., Houze *et al.*, 2015]. As outlined above, the atmospheric effects are positively related  
 458 with the intensity of thunderstorm activity, and it is conceivable that the chemical effects of  
 459 Argentinian storms could be even more dramatic. On the other hand, Argentinian storms  
 460 occur at low magnetic latitudes, corresponding to the inner radiation belt, where the available  
 461 fluxes of electrons for precipitation may be lower.

462 The main findings of our study are not contradictory to those of Rodger *et al.* [2007],  
 463 but more complementary. The main difference between the ionization calculation of Rodger  
 464 *et al.* [2007] and Marshall *et al.* [2019a] is that Rodger *et al.* [2007] modeled LEP events in  
 465 the ducted case, while Marshall *et al.* [2019a] modeled the LEP process in the nonducted  
 466 case due to multiple magnetospheric reflections; the energy and pitch angle distribution of  
 467 LEP fluxes are different. The main focus of Rodger *et al.* [2007] is the average LEP effects  
 468 produced by lightning discharge at high magnetic latitude, while this study mainly focuses  
 469 on the LEP effects at relatively lower latitudes and in extreme cases, which have not been  
 470 previously investigated. This study represents the first step of a series of studies towards bet-  
 471 ter understanding on the atmospheric chemical effects brought by LEP. The main goal is to  
 472 evaluate the immediate effects produced by LEP against other known ionization sources. The  
 473 next-step study is to quantify the indirect effects produced by LEP events using a 3D global  
 474 circulation model. Future studies can also aim at comparing the chemical effects reported in  
 475 this study with ground- and/or space-based measurements.

476 A 5% ozone depletion at 75–85 km altitude is comparable to that brought by other EEP  
 477 processes, for example, microburst precipitation [Seppälä *et al.*, 2018], EMIC-driven elec-  
 478 tron precipitation [Hendry *et al.*, 2021], as well as pulsating auroras [Turunen *et al.*, 2016].  
 479 Compared to other types of energetic particle precipitation (EPP), a single LEP event is  
 480 considerably shorter in duration and lower in precipitation fluxes. LEP produced by a 100-  
 481 kA lightning discharge has a peak energy flux ranging from  $10^{-2}$  [Peter and Inan, 2007] to  
 482  $10^{-1}$  ergs/cm $^2$ /s [Marshall *et al.*, 2019a], while typical values for the precipitation flux as-  
 483 sociated with visible aurora are 0.1–10 ergs/cm $^2$ /s [Meng, 1976; Rees, 1992]. Nevertheless,  
 484 with thousands of lightning flashes repetitively occurring within a short time window of a

485 few hours, i.e., an intense thunderstorm, the cumulative effects are pronounced, and the ion-  
486 ization production and chemical changes become accordingly amplified.

487 An  $O_x$  change of several percent is more significant than it appears since the occur-  
488 rence rate of LEP events globally is overwhelmingly higher than other EPP processes. The  
489 global lightning flash rate ranges from several tens to one hundred per second [Rakov, 2016],  
490 although not all flashes are sufficiently charged to give rise to radiation belt precipitation.  
491 Using Trimpi measurements, a representative value of the mean LEP rate at the Faraday sta-  
492 tion was found to be 0.79 per minute [Rodger *et al.*, 2004]. Rodger *et al.* [2003] have fur-  
493 ther estimated the global LEP rate using lightning observation data. An average value was  
494 suggested to be 0.18, 0.29, and 0.35 per minute at the  $L$  value of 2.4, 2, and 1.7, respec-  
495 tively. These values however, as noted by the authors, should only be considered as the lower  
496 bounds [Rodger *et al.*, 2003] since not all LEP events cause significant VLF changes, due to  
497 the interference of waveguide modes and the nonlinear relation between the electron density  
498 enhancement and VLF perturbation.

499 As in the 1D chemical simulation, the ozone reduction is mainly due to the immedi-  
500 ate effects of  $HO_x$  variation. Because of self-dissociation,  $HO_x$  has a relatively short lifetime  
501 and its effects on the ozone concentration are highly localized in space and time [Turunen  
502 *et al.*, 2016]. Andersson *et al.* [2014b] have shown that the global distribution of nighttime  
503 OH is mostly influenced by EEP events at high latitudes, whereas lightning discharge and  
504 associated LEP event occur more frequently at low latitudes. As such, the ozone reduction  
505 due to  $HO_x$  changes produced by LEP events could be insignificant. Present results show  
506 that, due to LEP ionization, the  $NO_x$  concentration could be enhanced by as high as  $\sim 156\%$   
507 in the mesosphere, but their effects on the ozone layer are not captured by the present chem-  
508 istry simulation. In the context of LEP events, the  $NO_x$  effects could become even greater  
509 than what is predicted in the present study for a single thunderstorm if we take the global oc-  
510 currence rate into account, although lightning activity tends to be more intense and frequent  
511 during summer times at low- and mid-latitude regions [Sousa, 2018]. From this considera-  
512 tion, the long-term global chemical effects of LEP events may be potentially important, but  
513 have been largely overlooked in previous studies.

514 To quantify these effects, the ionization results presented in Figure 2 can be rescaled  
515 using the lightning peak current reported by real time lightning-monitoring network, for ex-  
516 ample, the World Wide Lightning Location Network (WWLLN) [Dowden *et al.*, 2002]. The  
517 lightning data can be converted into altitude profiles of ionization production by LEP events,  
518 and then incorporated into global atmospheric chemistry and transport models such as the  
519 Whole Atmosphere Community Climate Model (WACCM) [Verronen *et al.*, 2016]. Future  
520 studies can thereby take the latitudinal and seasonal variation of thunderstorm activity into  
521 account, and aim at assessing the long-term global chemical effects produced by LEPs.

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527 paper are available at <https://doi.org/10.5281/zenodo.4599480>.

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