

1 **Solitary magnetic structures developed from gyro-resonance with solar wind**
2 **ions at Mars and Earth**

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26 **Key points**

27 • Solitary magnetic structures with density enhancements, plasma heating, and ion
28 reflection are observed in the Martian foreshock

29 • The structures resemble those developed from foreshock waves gyro-resonant with
30 solar wind ions at Earth

31 • The structures and ion distributions are reproduced by simulations, and shown to be
32 self-induced turbulence

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37 Abstract

38 We investigate solitary magnetic structures in the foreshock of Mars and Earth. The
39 structures exhibit pulse-like magnetic field and density enhancements along with plasma
40 heating and local solar wind ion reflection. The structures at Mars resemble the foreshock
41 structures developed from Ultra-Low-Frequency electromagnetic waves gyro-resonant
42 with solar wind ions at Earth, and propagate toward the magnetosphere. We perform fully
43 kinetic simulations to reproduce the solitary structures and the nonlinear evolution of ion
44 distribution functions, illustrating their resonance with solar wind ions. The structures
45 present self-induced foreshock turbulence that can have space weather effects. Our results
46 advance the fundamental understanding of how solar wind interacts with planetary
47 magnetospheres, and have potential impact on the current picture of planet-origin ion
48 escape.

49 1. Introduction

50 Mars has an induced magnetosphere [Nagy et al., 2004; Ramstad et al., 2020], and a
51 highly active environment upstream of the bow shock (e.g., review by Mazelle et al.,
52 2004). The region upstream of the bow shock is permeated by a number of plasma waves,
53 ranging from Alfvén waves [Dubinin et al., 2000; Halekas et al., 2017], Ultra-Low-
54 Frequency (ULF) magnetosonic waves [Collinson et al., 2018; Ruhunusiri et al., 2016;
55 Shan et al., 2020a] including those near the proton gyro-frequency [Brain et al., 2002;
56 Russell et al., 1990; Romanelli et al., 2016], whistler-mode waves at frequencies of a few
57 Hz [Brain et al., 2002; Halekas et al., 2020], and other higher frequency waves [Grard et
58 al., 1989; Sagdeev et al., 1990; Skalsky et al., 1992]. Upstream particles, including
59 reflected solar wind [Dubinin et al., 2006; Yamauchi et al., 2011] and freshly ionized
60 Martian [Yamauchi et al., 2015a] protons as well as heavier ions [Yamauchi et al.,
61 2015b] have been reported. A variety of structures have been observed upstream of the
62 bow shock, including those identified as fast-mode small-scale shocks developed from
63 ULF magnetosonic waves [Shan et al., 2020b], diamagnetic cavities [Øieroset et al.,
64 2001], hot flow anomalies [Collinson et al., 2015], spontaneous hot flow anomalies
65 [Collinson et al., 2017], foreshock cavities [Collinson et al., 2020], and quasi-periodic
66 compressive structures reminiscent of Short Large Amplitude Magnetic Structures
67 (SLAMS [Schwartz et al., 1992]) [Fowler et al., 2019; Halekas et al., 2017]. The
68 ionospheric impact of ULF magnetosonic foreshock structures are documented
69 [Collinson et al., 2018; Fowler et al., 2019]. Compressive foreshock structures have been
70 shown to occur most frequently during times when the interplanetary magnetic field

71 (IMF) is nearly Sun-planet aligned (quasi-radial), but can occur under a broad range of
72 upstream Mach number and dynamic pressure [Halekas et al., 2017].

73 In this paper, we analyze pulse-like magnetic field and density enhancements along
74 with plasma distribution functions at Mars, compare them with structures from the
75 Earth's foreshock with similar parameters, and elucidate the key defining physics with
76 particle-in-cell (PIC) simulations. We term the pulses solitary magnetic structures
77 because they exhibit properties (including polarization, spatial-temporal scales,
78 capabilities to heat the plasma and reflect solar wind ions) similar to those developed
79 from foreshock ULF waves gyro-resonant with solar wind ions at Earth [Chen et al.,
80 2020; Paper I hereafter], in contrast with the ULF waves [Eastwood et al., 2005] gyro-
81 resonant with backstreaming ions [Gary, 1991; Akimoto et al., 1993]. Intense solitary
82 magnetic structures are shown by a global simulation of the terrestrial magnetosphere to
83 bombard the magnetopause and induce reconnection as well as Earth-sized indents,
84 opening potential dayside escape channels for planet-origin ions [Chen et al., 2021]. The
85 induced magnetosphere of Mars, created by its ionosphere and crustal magnetic field,
86 stands off the solar wind and produces a bow shock and magnetosheath that are much
87 closer to the planet than those at Earth. We envision that the processes discussed in this
88 paper have powerful capacities to facilitate dayside loss of Martian ions.

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90 2. Spacecraft Measurements

91 In this section, we present examples of solitary magnetic structures from the
92 foreshock of Mars and Earth, identify their similarities, and extract key parameters for
93 simulations and linear instability analysis. The analyzed spacecraft measurements are

94 from Mars Atmosphere and Volatile EvolutioN (MAVEN) [Jakosky et al., 2015] and
95 Magnetospheric Multiscale (MMS) [Burch et al., 2016] missions. Two key parameters
96 determining the dominant wave mode are: ratio of the density of ions streaming back
97 toward the Sun (backstreaming, viewed in the plasma center-of-mass frame) to the total
98 plasma density (n_b/n_0) and relative drift velocity (V_d/V_A) between the two ion populations
99 [e.g., Gary, 1984; Akimoto et al., 1993; Weidl et al., 2019]. The selected MMS foreshock
100 crossing has both parameters similar to those in the MAVEN event. MMS provides
101 higher-resolution (~50 times higher cadence ion distribution functions than those from
102 MAVEN, for example) plasma measurements to enable the ion phase space structure to
103 be fully resolved and compared with simulation results.

104 The MAVEN measurements employed in this study are: magnetic field data (32
105 samples/s; 1 sample/s) from the Magnetometer [Connerney et al., 2015]; ion and electron
106 energy flux (4 s per sample) and distribution functions (8 s per sample) from the Solar
107 Wind Ion Analyzer (SWIA; Halekas et al., 2015) and the Solar Wind Electron Analyzer
108 (SWEA; Mitchell et al., 2016), respectively; the plasma density data from the SWIA
109 onboard moments (4 s per sample). Determination of the density and velocity of
110 backstreaming protons in the foreshock region is based on data from both SWIA and the
111 Supra-Thermal and Thermal Ion Composition (STATIC) instrument [McFadden et al.,
112 2015]. All vector are shown in the Mars-Solar-Orbital (MSO) coordinate system in which
113 the x-axis points from Mars to the Sun, the z-axis normal to the orbital plane (positive to
114 the north), and the y-axis completes the right-handed coordinate system.

115 For the presented MMS data, the magnetic fields (128 samples/s) are measured by
116 the Flux Gate Magnetometer [Russell et al., 2016], and electron (30 ms/sample) as well

117 as ion (150 ms/sample for burst mode; 4.5 s/sample for fast survey mode) data from the
118 Fast Plasma Investigation [Pollock *et al.*, 2016]. Vectors are shown in the Geocentric
119 solar ecliptic (GSE) coordinate system which is defined in an analogous manner as MSO.

120 Intense magnetic pulses (Figure 1c) disrupting the solar wind (Figure 1a) and
121 heating electrons (Figure 1b) are observed upstream of the Martian bow shock (shown by
122 the green shaded bar and orbit above Figure 1a and in Figure 1f, respectively) at ~1725-
123 1830 UT on July 11, 2019. The IMF is dominantly Sun-Mars aligned with B_x much
124 larger than $B_{y,z}$ as seen in the beginning of the shown interval (Figure 1e) and on average
125 throughout the interval with magnetic pulses. The pulses reach magnitudes 10-18 nT ~5-
126 9 times the background IMF strength (Figure 1c). Associated with the magnetic pulses
127 are the plasma density enhancements to $4-5 \text{ cm}^{-3}$ (the upstream solar wind density ~1.2
128 cm^{-3} , averaged from SWIA data at 1710-1720 UT). Solar wind protons exhibit strong
129 deceleration at the magnetic field and density spikes, as seen by the intense flux
130 extending from 1 keV to ~100 eV (Figures 1a and 2c). The nature of this deceleration
131 will be further discussed in Figure 3.

132 The sudden disappearance of magnetic pulses at ~1830UT suggests a sharp
133 boundary (at approximately $Z \sim 1.2 R_M$) beyond which the pulses are absent, rendering
134 the primary impact region of the solitary structures to be the quasi-parallel foreshock
135 within $\sim 1 R_M$ from the bow shock nose. Note that the IMF is still dominated by B_x and the
136 ions are solar wind-like. The location of MAVEN at 1800 UT in MSO coordinates is
137 $[2.0, 0.4, 0.9] R_M$, estimated to be $\sim 0.5 R_M$ ($\sim 1695 \text{ km} \sim 8 d_i$, where d_i is the ion skin
138 depth based on the upstream solar wind density) from the Martian bow shock encountered
139 after 1840 UT (Figure 1a).

140 Similar to that which observed in the foreshock of Mars, disruption of the solar wind
141 (Figure 1g) and heating of electrons (Figure 1h) are observed at the magnetic pulses in
142 the selected MMS event. The magnetic field amplitude (Figure 1i) and plasma density
143 (Figure 1j) are enhanced by about an order of magnitude compared with the upstream
144 values (IMF \sim 4.1 nT; density \sim 3.7 cm $^{-3}$). The intense B_y and B_z (Figures 1e and 1k)
145 exhibit both positive and negative excursions, and can enlarge magnetic shear angles
146 across the magnetopause [Chen et al., 2021]. The MMS four spacecraft configuration in
147 the selected foreshock crossing is colinear with an inter-spacecraft separation a few d_i ,
148 enabling definitive determination that the solitary magnetic structures evolve from the
149 ULF waves gyro-resonant with solar wind ions (Figure 3 in Paper I).

150 For the wave gyro-resonant with solar wind ions, the inter-pulse separation results
151 from the nature of the solitary pulse formation: growing from the maxima of beat-wave-
152 like envelopes [Paper I], in contrast to the quasi-periodic shocks whose inter-pulse
153 separation is dictated by the proton cyclotron periodicity [Shan et al., 2020b]. The
154 separation in time between two consecutive pulses varies substantially, ranging from 5 to
155 25 s in the zoom-in interval (Figure 2f). For comparison, the proton cyclotron period for
156 the background magnetic field (2 nT, taken from 1720 UT; MAVEN MAG) is \sim 33 s. The
157 inter-pulse durations observed by MMS are also not organized by the cyclotron period
158 which is \sim 17s for the MMS foreshock event (Figure 2n).

159 The backstreaming ion population exhibits a net sunward velocity in the spacecraft
160 frame. The v_x reduced distribution function (summed over v_y and v_z ; Figures 2b, 2d, and
161 2l) shows the colder solar wind population at $v_x \sim -440$ km/s, and the backstreaming
162 hotter population (predominantly at $v_x \sim 100$ -200 km/s in Figures 2b and 2l) which

163 exhibits variations in both the phase-space-density (PSD) and velocity as does the
164 backstreaming population in the simulation (Figures 3aeh).

165 We note the following additional features in the side-by-side zoom-in views of
166 solitary magnetic structures from Mars and Earth shown in Figure 2: (1) At the solitary
167 magnetic structures, the solar wind (SW) ion population is decelerated to below 100 eV
168 as well as accelerated to \sim 4 keV (Figures 2c and 2k; the acceleration at Mars is further
169 supported by H^+ data from STATIC). SW ion acceleration also occurs in the simulation
170 (below $-4 V_A$ in Figures 3h and 3n). (2) The most intense flux of decelerated SW ions
171 occurs at the upstream edge of the structure (Figures 2cdg and 2klo; examples: 120131,
172 120220, and 120245 UT), corresponding to the density peak which tends to be displaced
173 upstream from the magnetic field peak (Figure 2o; the structure at 120130 UT for
174 example). (3) The duration of a single pulse is \sim 3-10 s (Figures 2f, 2j, and 2n; a 10s-
175 duration pulse is at 183005-183015UT, not shown) and the pulse polarization is
176 dominantly right-handed in the spacecraft frame (pulse moving toward $-x$, B_y leads B_z
177 variations by \sim 90 degrees and $B_x > 0$ in Figures 2f and 2n; magnetic wave power shown in
178 Figures 2hi and 2pq), consistent with the polarization of ULF waves gyro-resonant with
179 solar wind ions. (4) Electrons remain one population throughout the interval of interest
180 (Figures 2e and 2m). The electron thermal spread in v_x enhances in intervals when
181 thermalized ions are observed (Figures 2cde and 2klm; the interval at \sim 1821 UT for
182 Mars, and 120120-30 UT for MMS) and the associated densities are relatively low.

183 We separate the solar wind and backstreaming populations in the 3D ion distribution
184 data to compute their moments to obtain the key parameters needed for PIC simulations
185 and linear instability analysis. For Mars, $n_b/n_0 \sim 0.15-0.17$, $V_d/V_A \sim 15$ based on

186 MAVEN data (SWIA and STATIC H+) from 20190711/1732-1734 UT. For Earth, n_b/n_0
187 ~ 0.17 , $V_d/V_A \sim 15$ based on MMS3 data from 20190222/115819-20 UT. Solitary
188 magnetic structures have been shown by PIC simulations to develop from ULF
189 electromagnetic waves gyro-resonant with solar wind ions using a backstreaming ion
190 density ratio $n_b/n_0=0.5$ extracted from an Earth foreshock event where N and $|B|$ peaks
191 coincide with each other without a relative phase shift [Paper I]. Using the same setup as
192 in Paper I and the Mars/MMS parameters, the simulations reproduce structure properties
193 including the displacement of the density peak slightly upstream of the $|B|$ maximum as
194 discussed below.

195

196 3. Particle-in-cell simulations

197 We carry out PIC simulations to demonstrate that given the conditions at the
198 foreshock of Mars and Earth, intense solitary magnetic structures develop from waves
199 gyro-resonant with solar wind ions under a constant IMF. The initial conditions are: two
200 counterstreaming ion and one electron populations (as seen in Figures 2be and 2lm) with
201 zero net total current. We perform simulations and corresponding linear instability
202 analyses with $n_b/n_0 = 0.1, 0.2$ and $V_d/V_A = 10, 15, 20$. The mode gyro-resonant with solar
203 wind ions dominates in all combinations of parameters in the above range. We show
204 results from the run with $n_b/n_0 = 0.2$ and $V_d/V_A = 20$ to illustrate the key physics. In the
205 plasma center-of-mass frame (plasma frame), the initial conditions are: one cold ion
206 population moving at $V_{sw} = -4V_A$ (corresponding to $V_{sw} = -9V_A \sim -440 \text{ km/s}$ for Mars and
207 Earth in the spacecraft frame), where V_A is the Alfvén speed based on the background
208 magnetic field B_0 and the total number density (n_0); one hot backstreaming ion

209 population at $V_b = 16V_A$; the temperature ratios $T_b/T_{sw} = 25$ and $T_e/T_{sw} = 2$; a uniform
210 $B_0 = (0.067, 0, 0)m_e c/|e|d_e$, where m_e is the electron mass, c is the speed of light, e is the
211 electron charge, and d_e is the electron inertia length. Other parameters are: the mass ratio
212 $m_i/m_e = 100$, electron plasma to cyclotron frequency ratio $\omega_{pe}/\omega_{ce} = 15$, and solar wind
213 ion beta $\beta = 1$. The simulation has one spatial and three velocity dimensions, performed
214 using the particle-in-cell code VPIC, which solves a system of relativistic Vlasov-
215 Maxwell equations [Bowers et al 2008].

216 In the simulation, the magnetic field $B_{y,z}$ grows from right-hand polarized
217 electromagnetic waves (electric field variations are similar to those presented in Figures 2
218 and 4 in Paper I) in which B_y leads B_z by 90 degrees, consistent with the observed pulse
219 forms shown in Figure 2. The magnetic field first develops as quasi-sinusoidal waves like
220 the profiles from t-2 (Figure 3i; the even earlier stage is similar to that shown in Figure 4,
221 Paper I). Later the field grows further at local maxima of $|B|$ to form solitary structures at
222 time t and $t+2.5$.

223 Simulation results show ion reflection, thermalization, and electron heating at sites
224 of magnetic field and density enhancements, consistent with the observations (Figure 2)
225 at the foreshock of Mars and Earth. Majorities of the incoming solar wind ions are
226 reflected slightly upstream of the $|B|$ peak, forming an arc structure in the $[x, v_x]$ phase-
227 space (Figure 3a). Such arc structures are seen in the high-cadence MMS ion v_x
228 distribution (120130 and 120243 UT in Figure 2l). Electron heating occurs at the solitary
229 structure, and tends to be more intense in regions where ions are more thermalized (e.g.,
230 $x \sim 580-585$ in Figure 3b), similar to that observed by MAVEN and MMS (Figure 2).

231 The density (N) maximum is displaced slightly upstream (to the right) of the $|B|$
232 peak (Figure 3d), a feature resolved by MMS (Figure 2o). At later time (for example, 2.5
233 ω_{ci}^{-1} later, where ω_{ci} is the ion cyclotron frequency), further ion thermalization occurs
234 (Figure 3e), the magnetic field enhancements become more isolated (Figure 3f), and the
235 relationship between the $|B|$ and N peaks becomes more variable across different solitary
236 structures (Figure 3g). The less correlated $|B|$ and N profiles observed by MAVEN and
237 MMS in the more thermalized regions (e.g., ~182050-182110 UT in Figures 2c and 2g
238 for Mars; ~120120-120130 UT in Figures 2k and 2o for Earth) can thus be attributed to
239 temporal evolution of the structures.

240 Solitary magnetic structures grow from waves gyro-resonant with solar wind ions.
241 The mode is right-hand polarized and non-resonant with the backstreaming ions. Before
242 the formation of solitary structures at time t , the wave propagates at $V_{ph}=-2V_A$ with a
243 dominant wave number of $kd_i \sim 0.53$, consistent with linear instability analysis predicting
244 the maximum growth rate to be $0.81 \omega_{ci}$ at $kd_i \sim 0.51$ for the non-resonant mode. The
245 cyclotron resonant velocity is $V_{res} = V_{ph} + \omega_{ci}/k = -3.9V_A$, within one solar-wind ion
246 thermal speed ($1.1 V_A$) from $V_{sw} = -4V_A$, and hence satisfies the solar wind gyro-
247 resonance condition [Gary, 1991; Weidl et al., 2019].

248 The capacity of solitary structures to reflect SW ions like a collisionless shock is
249 further demonstrated by v_x - v_y ion velocity distribution functions (summed over v_z) from
250 PIC (Figure 3j), Mars (Figure 3k), and Earth (Figure 3l). All three distribution functions
251 (in the spacecraft frame) show the co-existence of the slightly slowed down incoming
252 SW and the reflected SW population at the upstream edge just before the $|B|$ peak. The

253 v_x - v_y distributions are 2D velocity-space views of a slice (at a fixed x in PIC or fixed time
254 in MMS) of the arc structure in the v_{ix} reduced distribution (Figures 2l and 3a).

255 The solitary-structure velocity lies between the incoming SW and the locally
256 reflected SW populations, as verified by the PIC and MMS distributions. Based on MMS
257 multi-spacecraft timing analysis, the propagation velocities of terrestrial solitary
258 structures (such as the ones shown in Figure 2n) are estimated to be 200-300 km/s anti-
259 sunward in the spacecraft frame, consistent with the average velocity of the incoming and
260 reflected SW components ($\sim 5-6V_A \sim 250-300$ km/s). For Mars, the observed velocities
261 of the two populations indicate that the solitary structure moves toward the planet with a
262 speed approximately 250-300 km/s.

263 Solitary magnetic structures are kinetic in nature. The solar wind ion population
264 having a finite thermal spread in the velocity space enables gyro-resonance of a subset of
265 SW ions with the wave and a rich spectrum of deceleration/acceleration (including
266 reflections) throughout the solitary structure as seen in the x - v_x structure (Figure 3). To
267 further illustrate this kinetic nature, we trace all SW ions within a spatial bin (size 0.02 d_i ;
268 marked by a magenta vertical line in Figure 3m) forward in time from $t-2$ (Figure 3m) to
269 t , and display all traced ions (as magenta open circles) in the x - v_x phase-space at t (Figure
270 3n). The SW ions spread out from a 0.02 d_i bin to an x range of 13 d_i ($x \sim 600-613$),
271 occupy a much larger v_x range, and populate key x - v_x phase-space structures such as the
272 arc ($x \sim 611-613$; Figure 3n). Examination of traced solar wind ion orbits (one example
273 shown as the cyan curve in Figure 3n) further confirms that reflection occurs at different
274 locations within the solitary structure, and the most probable turning points correspond to
275 the density peak. The histogram (Figure 3o) of the number of traced SW ions exhibits a

276 profile qualitatively similar to the density profile (black curves in Figures 3m-n
277 represents 3N). The differing acceleration and deceleration features of the SW ions and
278 their local reflections throughout the magnetic structure indicate that the solar wind
279 population cannot be viewed as a fluid element, but need to be treated as particles
280 interacting through their mean fields.

281

282 4. Summary, discussion, and conclusion

283 In summary, we report solitary magnetic structures developed from gyro-resonance
284 with solar wind ions in the foreshock of Mars, based on comparing MAVEN
285 measurements with the high-cadence MMS data from an Earth foreshock crossing that
286 exhibits similar parameters and comparing the observations with PIC simulations. The
287 solitary structures in both observations and simulations decelerate/accelerate solar wind
288 ions, locally reflect portions of the incoming solar wind population, and heat electrons,
289 behaving like kinetic shocks moving toward the planet. The dominant instability is
290 consistent with the right-hand non-resonant mode in the literature [Gary, 1991; Paper I],
291 and requires higher backstreaming ion density and/or velocity than does the resonant
292 mode. The evidence for the non-resonant mode reported in the paper indicates that a
293 quasi-radial IMF may stimulate the foreshock/shock/magnetosphere environment of Mars
294 to provide more/faster backstreaming ions through reflecting the incoming SW or
295 through enabling ionized martian ions to flow sunward. The exact process remains an
296 open question.

297 Strictly speaking, given the single spacecraft measurements from MAVEN, one
298 cannot entirely rule out the possibility that the solitary structures are from the upstream

299 solar wind and propagate to the foreshock of Mars. However, the presence of
300 backstreaming ions, the similarity of structure properties with those observed in Earth's
301 foreshock where the evolution of ULF waves into solitary structures are directly
302 observed, and the similarity with solitary structures generated in PIC simulations all
303 combine to support that the structures in the Martian foreshock are generated by the
304 kinetic interaction of the solar wind and backstreaming ions.

305 The solitary magnetic structures discussed in this paper develop from the ULF
306 electromagnetic waves gyro-resonant with solar wind protons (coded as "SWr structures"
307 in this paragraph), to be distinguished from the periodic shocks that evolve from gyro-
308 resonance with freshly ionized Martian hydrogen [Shan et al., 2020b]. The key observed
309 differences are: 1) Polarization in the spacecraft frame: SWr structures are dominantly
310 right-handed. The quasi-periodic shocks are primarily left-handed. 2) Time scales: for
311 SWr structures, the inter-pulse duration is not organized by the cyclotron period, but as
312 emphasized in Paper I, the solitary nature stems from a beat-like magnetic field envelope
313 where the $|B|$ local maxima set the initial locations for fastest nonlinear growth; the
314 individual pulse span is 3-10 s, consistent with a time scale dominated by the Doppler
315 shift $k V_{sw} / 2 \pi \sim 0.1\text{-}0.3$ Hz (f_{ci} is only 0.03 Hz), where k is the most unstable wave
316 number ($kd \sim 0.4\text{-}0.7$) for the parameter range $n_b/n_0 = 0.1\text{-}0.2$ and $V_d/V_A = 10\text{-}20$. The
317 quasi-periodic shocks are governed by the proton cyclotron periodicity. 3) Backstreaming
318 ion velocity in the spacecraft frame: SWr structures are associated with ions streaming
319 with a net sunward velocity. The backstreaming ion population for the quasi-periodic
320 shocks are newborn protons that have a net anti-sunward velocity or nearly at rest [Shan
321 et al., 2020b]. 4) Pulse shapes: SWr structures have bipolar B_{xyz} (the bipolar B_x with

322 respect to the background field is most likely due to the wave vector not along x)
323 dominated by B_{yz} . The B_{xyz} profiles of the quasi-periodic shocks are primarily unipolar (B_x
324 never crosses zero, in particular; Figure 1b in Shan et al. [2020b]). Note that the above
325 contrast is based on two case studies. Further investigations are required to assess the
326 generality for both types of structures.

327 We emphasize that the physics leading to the density increase at the structures
328 reported here is primarily local reflections of solar wind ions. Our simulation predicts that
329 the most probable site of reflection corresponds to the density maximum, and occurs
330 slightly upstream of the magnetic field peak given the parameters from the reported
331 foreshock crossings at Mars and Earth. This displacement is observed by MMS in the
332 high cadence measurements.

333 Solitary magnetic structures have the capacity to lead to dayside ion escape. The
334 structures have been demonstrated to further intensify in the magnetosheath and lead to
335 magnetic reconnection as well as indents at the magnetopause by a global hybrid
336 simulation of the terrestrial magnetosphere [Chen et al., 2021]. The induced
337 magnetosphere of Mars is much smaller and weaker than Earth's, and hence more
338 susceptible to the impact of these planet-sized structures. For example, the foreshock
339 generated and amplified bipolar B_y and B_z (Figures 2) can increase the shear angle
340 between the magnetosheath field and the crustal magnetic field at the dayside, and hence
341 may increase the likelihood of dayside reconnection such as that observed by MAVEN
342 [Harada et al., 2018], and open escape pathways from the dayside ionosphere, a major
343 reservoir of Martian ions. Even though the wave modes differ, the structures reported in
344 this paper exhibit magnetic field and density enhancements similar in strengths as the

345 ULF magnetosonic waves (both types of structures can be viewed as subclasses of
346 SLAMS [Schwartz et al., 1992], as they meet the criterion of $|B| > \sim 2B_0$), and hence may
347 impact the ionosphere of Mars in ways similar to those documented [Collinson et al.,
348 2018; Fowler et al., 2019] and contribute to the atmospheric loss.

349 In conclusion, the solitary magnetic structures in the foreshock of Mars are
350 consistent with the terrestrial solitary structures developed from the ULF waves gyro-
351 resonant with solar wind ions under a constant quasi-radial IMF. The structures are
352 planet-sized and kinetic in nature, behaving like mini-shocks moving toward the
353 magnetosphere. They have the potential to lead to dayside escape of planetary ions. Our
354 results suggest that planet-sized space weather events can be self-induced under steady
355 IMF and solar wind conditions.

356 **Captions**

357 **Figure 1.** Overview of solitary magnetic structures at the foreshock of Mars (left) and
358 Earth (lower right). (a) Ion energy flux. (b) Electron energy flux. (c) Magnetic field
359 amplitude enhanced with the density (d). (e) magnetic field components B_{xyz} . (f)
360 MAVEN Orbit (color coded with shades of green along with the horizontal bar above
361 Figure 1a to aid visualization). (g-k) Measurements from Earth's foreshock presented in
362 the same format as Figures 1a-e.

363 **Figure 2.** Zoom-in view of example solitary structures and corresponding plasma
364 distribution functions at Mars and Earth. (a, f, n) Magnetic field components. (b, d, l) Ion
365 phase-space density in the v_x space (summed over v_y and v_z). (c, k) Ion energy flux. (e,
366 m) Electron phase-space density in v_x . (g, o) pulse-like enhancements of the magnetic
367 field amplitude and density. Ion thermalization (d, l) and electron heating (c, m) occur at
368 solitary structures. (h-i, p-q) Magnetic wave power showing that the fluctuations at \sim
369 0.1Hz and slightly above are dominantly right-hand (RH) polarized, in contrast with the
370 much weaker wave power in the LH panels. The vertical blue dashed lines mark the
371 upstream side of example solitary structures and the center time of the distribution
372 functions presented in Figure 3, magenta arrows example backstreaming ion populations
373 when the SW is approximately un-disturbed, blue arrows example accelerated SW ion
374 features, and black arrows the arc structures signifying local SW ion reflection.

375 **Figure 3.** PIC simulation results showing ion reflection, thermalization, and electron
376 heating at sites of magnetic field and density enhancements, consistent with the
377 observations at the foreshock of Mars and Earth. Local reflection of solar wind ions at the
378 solitary structure is featured. (a-b, e, h) Ion and electron $x-v_x$ phase space (summed over

379 v_y and v_z) presented in the plasma frame. (c, f, i) Magnetic field components showing $B_{y,z}$
380 develop from right-hand polarized waves. (d, g, m, o) Magnetic field amplitude ($|B|$) and
381 density (N) profiles. Ion v_x - v_y velocity distribution functions from PIC (j), Mars (k), and
382 Earth (l), presented in the same normalized unit, same velocity axes and velocity range in
383 the spacecraft frame, showing that the incoming SW (denoted as “IncSW” and indicated
384 by green arrows) co-exists with the locally reflected SW (“LocRefSW,” magenta arrows).
385 Zoom-in views of the SW ion x - v_x phase space at $t-2$ (m) and t (n) shows how all the solar
386 wind ions in the 3D velocity distribution from $x=615$ (marked by a vertical magenta
387 dashed line) at $t-2$ are distributed at t (magenta open circles). One example orbit of
388 reflected SW ions is displayed as the cyan curve (n). The highest number of traced SW
389 ions at time t (o) occurs at the location of the most probable reflection, corresponding to
390 the bend-over location of the arc structure in x - v_x (n). (a-d) are from $t = 13$, (e-g) from
391 $t + 2.5$, and (h-i) from $t - 2$. Time is in unit of ω_{ci}^{-1} . The horizontal green bar marks the x
392 range shown in Figures 3m-o. The vertical line in (a-d) indicates the location from which
393 the v_x - v_y distribution function (j) is taken.

394

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401

402 Open Research

403 MMS data are available at <https://lasp.colorado.edu/mms/sdc/public/about/browse-wrapper/>. MAVEN data are available at
404 <https://lasp.colorado.edu/maven/sdc/public/data/sci/>. The study uses L2 FPI and FGM
405 data from MMS, and L2 swi, swe, sta, and mag data from MAVEN.

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