

# A Novel Sequential Method to Train Physics Informed Neural Networks for Allen Cahn and Cahn Hilliard Equations

Revanth Mattey <sup>a</sup>, Susanta Ghosh  <sup>a,b,\*</sup>

<sup>a</sup>*Department of Mechanical Engineering—Engineering Mechanics, Michigan Technological University, MI, USA*

<sup>b</sup>*The Center for Data Sciences, The Center for Applied Mathematics and Statistics Michigan Technological University, MI, USA*

---

## Abstract

A physics informed neural network (PINN) incorporates the physics of a system by satisfying its boundary value problem through a neural network's loss function. The PINN approach has shown great success in approximating the map between the solution of a partial differential equation (PDE) and its spatio-temporal coordinates. However, the PINN's accuracy suffers significantly for strongly non-linear and higher-order time-varying partial differential equations such as Allen Cahn and Cahn Hilliard equations. To resolve this problem, a novel PINN scheme is proposed that solves the PDE sequentially over successive time segments using a single neural network. The key idea is to re-train the same neural network for solving the PDE over successive time segments while satisfying the already obtained solution for all previous time segments. Thus it is named as backward compatible PINN (bc-PINN). To illustrate the advantages of bc-PINN, the Cahn Hilliard and Allen Cahn equations are solved. These equations are widely used to describe phase separation and reaction-diffusion systems. Additionally, two new techniques have been introduced to improve the proposed bc-PINN scheme. The first technique uses the initial condition of a time-segment to guide the neural network map closer to the true map over that segment. The second technique is a transfer learning approach where the features learned from the previous training are preserved. We have demonstrated that these two techniques improve the accuracy and efficiency of the bc-PINN scheme significantly. It has also been demonstrated that the convergence is improved by using a phase space representation for higher-order PDEs. It is shown that the proposed bc-PINN technique is significantly more accurate and efficient than PINN.

**Keywords:** Physics informed neural networks,, Partial differential equation (PDEs), Allen Cahn equation, Cahn Hilliard equation

---

## 1. Introduction

Traditional physics-based numerical methods for solving partial differential equations (PDEs) have found remarkable success in solving various science and engineering problems. These methods are accurate but computationally expensive for complex problems such as nonlinear PDEs and requires problem-specific techniques. In the last decade data driven methods have gained a lot of attention in almost all areas of science and engineering. Data driven methods for PDEs can help in identifying highly non-linear mappings (between the inputs and outputs) which can substitute or augment expensive physics based simulations. Due to their versatility and fast

---

\*Corresponding author; Email: susantag@mtu.edu

9 evaluation capabilities, machine learning based models can be used as PDE solvers in situations  
10 when there is a requirement of a large number of simulations such as the inverse–problem and  
11 homogenization [1, 2].

12 Several data–driven techniques have been attempted to solve PDEs. For instance, the Gaussian  
13 Process based approaches are described in [3, 4, 5, 6, 7]. Despite the ease of training for Gaussian  
14 Process, this approach did not gain as much popularity as a neural network for solving PDEs  
15 due to its difficulties in handling high dimensional problems.

16

17 Among different data–driven techniques for PDEs the Physics Informed Neural Networks  
18 (PINN) has shown remarkable promise and versatility. PINN is a new class of machine learning  
19 technique where a neural network’s loss function is designed to satisfy the Initial Boundary Value  
20 Problem (IBVP) [8]. A PINN “learns” the non linear map between the spatio–temporal input  
21 and the solution of the PDE in a given domain. Henceforth, throughout this paper the PINN  
22 described in [8] would be referred to as std-PINN. std-PINN utilizes the automatic-differentiation  
23 capability [9] to compute the derivatives of the field variables.

24

25 Different variants of std-PINN are shown to work effectively in solving many forward and  
26 inverse problems [10, 11, 12]. Recently in [13], PINN has been extended to satisfy various  
27 conservation laws while solving the PDEs. This approach is named as cPINNs. cPINNs solve the  
28 problem over several sub-domains and ensure flux continuity at the boundaries of the sub-domains.  
29 Another extension to cPINNs is XPINN known as extended PINN, where the authors propose a  
30 generalized space–time domain decomposition based deep learning framework [14]. The key idea  
31 in XPINN is to decompose the domain into multiple sub-domains and train the sub-domains  
32 using multiple neural networks (sub Net), while ensuring  $C^0$  continuity along the interfaces.  
33 While most of the PINN approaches solve the strong form of a PDE, it can also be used to solve  
34 the weak (variational) form of a PDE. Since the weak form incorporates the natural boundary  
35 conditions, the neural network solution only needs to satisfy the essential boundary conditions *a priori*.  
36 This aspect is used in several numerical methods for PDEs such as the finite element  
37 method. Due to this advantage of weak form over strong form the application of PINN on  
38 the weak form has been investigated in [15]. In this study, the authors have considered the  
39 variational form for stochastic PDEs and applied the idea of PINNs to obtain the solution of  
40 the PDE. Also the corresponding uncertainty propagation through their model is presented in  
41 [15, 16]. Uncertainty quantification provides the variation associated with the prediction of the  
42 model. It is particularly useful for systems where there is a high cost of data acquisition or  
43 lack of high resolution data [17]. The authors in [18] proposed a Bayesian approach for physics  
44 informed neural network to solve forward and inverse problems.

45

46 The promise and versatility of PINN have been demonstrated through its application for a  
47 wide range of problems. PINN has been used in modeling subsurface transport phenomena [19],  
48 approximating Euler equations for high speed flows [1], constitutive modeling of stress–strain  
49 behavior in biological tissues [20], predicting arterial blood pressure from noisy MRI data of flow  
50 velocity [21], and cardiac activation mapping for diagnosing atrial fibrillation [22].

51

52 Recently many other data–driven techniques for solving time varying PDEs have been pro-  
53 posed. In [23] the authors proposed a framework called Neural-PDE which aims to learn the  
54 solution of the PDE by utilizing numerical techniques like FDM (Finite difference method) and  
55 LSTM (Long short term memory) networks. The method aims to predict the solution of the  
56 PDE at  $n$  future time steps by using the meshgrid data (solution) of all the previous time steps.  
57 Artificial neural network based approaches for solving parametric PDEs have been introduced by

58 the authors in [24]. Operator learning is a new and emerging technique for solving PDEs. In  
 59 operator learning a map between initial condition and solution of the PDE is learnt by using  
 60 multiple instantiations [25, 26, 27].

61

62 In the present work, we demonstrate that the accuracy of the std-PINN [8] suffers in the  
 63 presence of (i) Strong Non-linearity, and (ii) Higher order partial differential operators. In order  
 64 to illustrate the above, we chose the Allen Cahn equation having *strong non-linearity* and Cahn  
 65 Hilliard equation having *strong non-linearity and fourth order derivative*. These are the two  
 66 most widely used PDEs to study diffusion separation and multi-phase flows [28, 29, 30, 31]. To  
 67 overcome the drawbacks of the std-PINN, we have proposed an extension, which is named as  
 68 backward compatible PINN (bc-PINN). The proposed bc-PINN solves the PDE over successive  
 69 time segments by re-training the same neural network, where the key idea is:

70 *To ensure that the neural network can reproduce the solution for all the prior time segments  
 71 while solving the PDE for a particular time segment.*

72 Henceforth, this idea is referred to as *backward compatibility*. Some of the main advantages of  
 73 the proposed bc-PINN method are as follows:

- 74 1. It works for higher order and strongly nonlinear PDEs by using fewer iterations and  
 75 collocation points while achieving significantly higher accuracy when compared to std-  
 76 PINN.
- 77 2. A single neural network is used for the entire domain and continuity across the time  
 78 segments is ensured for the predicted solution and its derivatives.

79 The rest of the paper is organized as follows: in section (2) the std-PINN method is briefly  
 80 reviewed; in section (3) the proposed bc-PINN method is presented; in section (4) the Allen  
 81 Cahn and Cahn Hilliard equations are described; in section (5) the bc-PINN method is analyzed  
 82 and compared against the std-PINN and the XPINN method for the one dimensional (1D) Allen  
 83 Cahn and Cahn Hilliard equations; in section (6) two new techniques initial condition guided  
 84 learning and transfer learning based acceleration have been presented along with the results for  
 85 the two dimensional (2D) Allen Cahn and Cahn Hilliard equations. Finally, the conclusions are  
 86 presented in section (7).

87 **2. A brief review of standard physics informed neural network (std-PINN) for  
 88 partial differential equations**

Physics informed neural network (PINN) is a class of machine learning model where the governing PDE is satisfied through the loss function of the neural network [8]. The efficient optimization and prediction capabilities of neural network are exploited in the std-PINN approach. In std-PINN a neural network is trained to predict the solution at any point in the entire spatial-temporal domain. Let's consider the general form of a  $m^{\text{th}}$  order partial differential equation (PDE):

$$h_t = F(h(\mathbf{x}, t), h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), h_{\mathbf{x}}^{(2)}(\mathbf{x}, t), \dots, h_{\mathbf{x}}^{(m)}(\mathbf{x}, t)), \quad \mathbf{x} \in \Omega \subset \mathbb{R}^D, \quad t \in (0, T] \quad (1)$$

Here,  $\Omega$  is an open set of  $\mathbb{R}^D$  ( $D = 1, 2, 3$ ).  $F$  is a non linear function of the solution  $h(\mathbf{x}, t)$  and it's spatial derivatives  $(h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), h_{\mathbf{x}}^{(2)}(\mathbf{x}, t), \dots, h_{\mathbf{x}}^{(m)}(\mathbf{x}, t))$  where  $\mathbf{x}$  and  $t$  are the space and time coordinates respectively. The corresponding boundary conditions and initial conditions are

$$\begin{aligned} h(\mathbf{x}, 0) &= \phi(\mathbf{x}), \quad \mathbf{x} \in \Omega \\ h(-\mathbf{x}, t) &= h(\mathbf{x}, t), \quad (\mathbf{x}, t) \in \Gamma \times (0, T] \\ h_{\mathbf{x}}^{(1)}(-\mathbf{x}, t) &= h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), \quad (\mathbf{x}, t) \in \Gamma \times (0, T] \end{aligned} \quad (2)$$

89 Where,  $\Gamma$  is the boundary of  $\Omega$ . The PDE, the Initial and the Boundary Conditions (given by  
90 equation (1-2)) form a initial-boundary value problem (IBVP) considered in this study. The  
91 boundary conditions are taken as periodic and the initial condition is a real function.

92 std-PINN approximates the map between points in the spatio-temporal domain to the solution  
93 of the PDE. The parameters of the neural network are randomly initialized and iteratively updated  
94 by minimizing the loss function that enforces the PDE. The std-PINN's loss function consists of  
95 three error components, for the prediction of the neural network as in the following (i) Initial  
96 Condition, (ii) Boundary Condition, and (iii) PDE. Let  $\hat{h}(\mathbf{x}, t)$  be the output of neural network.  
97 The three components of the std-PINN's loss function are given below:

- Mean squared error on the Initial Condition

$$\text{MSE}_I = \frac{1}{N_i} \sum_{k=1}^{N_i} \left( \hat{h}(\mathbf{x}_k^i, 0) - h_k^i \right)^2, \quad \mathbf{x}_k^i \in \Omega \quad (3)$$

99 where  $\hat{h}(\mathbf{x}_k^i, 0)$  is the neural network output and  $h_k^i$  is the given initial condition at  $(\mathbf{x}_k^i, 0)$ .  
100 Here, the superscript,  $(\bullet)^i$  stands for initial condition.

- Mean squared error on the Boundary Condition

$$\text{MSE}_B = \frac{1}{N_b} \sum_{k=1}^{N_b} \sum_{d=1}^{n_d} \left( \hat{h}^{(d-1)}(\mathbf{x}_k^b, t_k^b) - \hat{h}^{(d-1)}(-\mathbf{x}_k^b, t_k^b) \right)^2, \quad (\mathbf{x}_k^b, t_k^b) \in \Gamma \times (0, T] \quad (4)$$

101 where  $n_d$  is the highest order of derivative to which the periodicity is enforced on the  
102 boundary,  $\Gamma$ . Here, the superscript,  $(\bullet)^b$  stands for boundary condition.

- The Mean squared error due to Residual of the partial differential equation

$$R := \hat{h}_t - F(\hat{h}, \hat{h}_{\mathbf{x}}^{(1)}, \hat{h}_{\mathbf{x}}^{(2)}, \dots, \hat{h}_{\mathbf{x}}^{(m)})$$

$$\text{MSE}_R = \frac{1}{N_r} \sum_{k=1}^{N_r} (R(\mathbf{x}_k^r, t_k^r))^2, \quad (\mathbf{x}_k^r, t_k^r) \in \Omega \times (0, T] \quad (5)$$

The superscript,  $(\bullet)^r$  stands for residual of the PDE.  $(\mathbf{x}_k^i)$  and  $(\mathbf{x}_k^b, t_k^b)$ , represent the set of  
points where the initial and boundary errors are computed. The residual/collocation error is  
computed at the collocation points  $(\mathbf{x}_k^r, t_k^r)$ . These points on the domain and the boundary are  
obtained using a latin hypercube sampling approach. Therefore, the total loss function of the  
neural network is given by adding all the aforementioned mean squared errors

$$\text{MSE} = \text{MSE}_I + \text{MSE}_B + \text{MSE}_R \quad (6)$$

Once the std-PINN is trained, the accuracy of the predicted solution is computed with respect  
to the reference solution at unknown points (called testing points). Highly accurate solution  
of the initial boundary value problem obtained by the Chebyshev polynomial based numerical  
algorithm [32] and is considered as the reference solution. The relative total error ( $\varepsilon_{total}$ ) of the  
PINN's prediction over the entire domain is obtained by normalizing the error with respect to  
the reference solution as

$$\varepsilon_{total} = \frac{\left[ \frac{1}{N} \sum_{k=1}^N \left( \hat{h}(\mathbf{x}_k, t_k) - h(\mathbf{x}_k, t_k) \right)^2 \right]^{1/2}}{\left[ \frac{1}{N} \sum_{k=1}^N (h(\mathbf{x}_k, t_k))^2 \right]^{1/2}} \quad (7)$$

The relative error ( $\varepsilon$ ) of the PINN's prediction at each point is obtained by normalizing the absolute error with respect to the reference solution as

$$\varepsilon(\mathbf{x}_k, t_k) = \frac{|\hat{h}(\mathbf{x}_k, t_k) - h(\mathbf{x}_k, t_k)|}{\left[\sum_{k=1}^N (h(\mathbf{x}_k, t_k))^2\right]^{1/2}} \quad (8)$$

103 Where  $h(\mathbf{x}_k, t_k)$  is the reference solution and  $\hat{h}(\mathbf{x}_k, t_k)$  is the neural network prediction for a  
104 set of testing points  $\{(\mathbf{x}_k, t_k)\}_{k=1}^N$ ,  $(\mathbf{x}_k, t_k) \in \Omega \times (0, T]$ . For all comparisons between reference  
105 and predicted solutions the relative total error ' $\varepsilon_{total}$ ' and relative error ' $\varepsilon$ ' is used.

### 106 3. The proposed backward compatible sequential PINN method (bc-PINN)

107 In this section, we introduce an extension of the std-PINN technique that solves an initial-  
108 boundary value problem sequentially in time.

#### 109 3.1. bc-PINN

In the proposed method the PDE is solved progressively in time by re-training a single neural network over successive time segments. The limitation of such retraining is that the network can predict only for the latest time segment and cannot predict for previous time segments for those it has been trained earlier. To overcome this limitation, the proposed model is designed to satisfy the solution of all the previous time segments while solving the PDE over a particular time segment. This scheme ensures backward compatibility of the solution by a single network. The proposed method is henceforth referred as backward compatible PINN (bc-PINN). The schematics of bc-PINN for a particular time segment is shown in figure (1) and the sequential scheme of proposed bc-PINN approach is shown in figure (2). In bc-PINN the time domain  $[0, T]$  is discretized into  $n_{max}$  segments as

$$[T_0 = 0, T_1], [T_1, T_2], \dots, [T_{n-1}, T_n], \dots, [T_{n_{max}-1}, T_{n_{max}} = T] \quad (9)$$

110 where the  $n^{\text{th}}$  segment is denoted as  $\Delta T_n = [T_{n-1}, T_n]$ ,  $n = 1, \dots, n_{max}$ .

For the first time segment  $\Delta T_1$  the solution of the PDE is sought through the std-PINN by minimizing the following loss function

$$\begin{aligned} \text{MSE}_{\Delta T_1} = & w_i \text{MSE}_I(\mathbf{x}_k^i, 0) + w_b \text{MSE}_B(\mathbf{x}_k^b, t_k^b) + w_r \text{MSE}_R(\mathbf{x}_k^r, t_k^r) \\ & \mathbf{x}_k^i \in \Omega, \quad (\mathbf{x}_k^b, t_k^b) \in \Gamma \times (0, T_1] \quad (\mathbf{x}_k^r, t_k^r) \in \Omega \times (0, T_1] \end{aligned} \quad (10)$$

111 Here,  $(\mathbf{x}_k^i, t_k^i)$  represent the set of points where the error on initial condition is computed and  
112  $(\mathbf{x}_k^b, t_k^b)$  represent the set of points where the error on boundary condition is computed within the  
113 time segment  $\Delta T_1 = (0, T_1]$ . For all of the subsequent time segments (i.e.  $\Delta T_n$ ,  $n = 2, \dots, n_{max}$ )  
114 we propose a novel loss function, which satisfies the solution of all previous time segments. The  
115 solution of all previous time segments is enforced by penalizing the departure from the already  
116 obtained solutions from the previous training, as given by

$$\begin{aligned} \text{MSE}_{\Delta T_n} = & w_i \text{MSE}_I(\mathbf{x}_k^i, T_{n-1}) + w_b \text{MSE}_B(\mathbf{x}_k^b, t_k^b) + w_r \text{MSE}_R(\mathbf{x}_k^r, t_k^r) \\ & + w_s \text{MSE}_S(\mathbf{x}_k^s, t_k^s), \quad n = 2, \dots, n_{max} \\ & \mathbf{x}_k^i \in \Omega, \quad (\mathbf{x}_k^b, t_k^b) \in \Gamma \times (T_{n-1}, T_n] \\ & (\mathbf{x}_k^r, t_k^r) \in \Omega \times (T_{n-1}, T_n], \quad (\mathbf{x}_k^s, t_k^s) \in \Omega \times [0, T_{n-1}] \end{aligned} \quad (11)$$

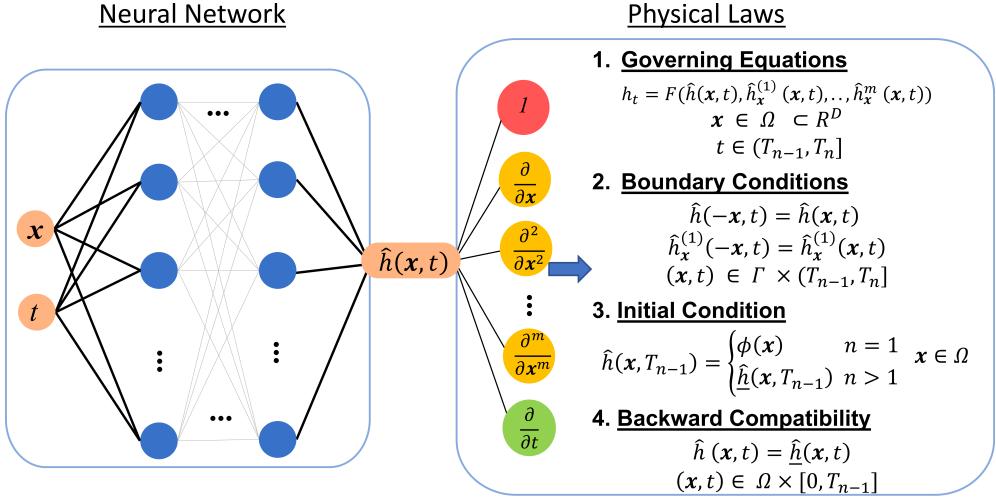


Figure 1: The schematics of the proposed backward compatible PINN (bc-PINN) approach for a time segment  $((T_{n-1}, T_n])$ . The neural network re-trains the PDE over  $(T_{n-1}, T_n]$  while satisfying the solution for all previous time segments. The error in the Initial Condition is computed at time  $t = 0$  for the first time segment and at time  $t = (T_{n-1})$  for the  $n$ th time segment.

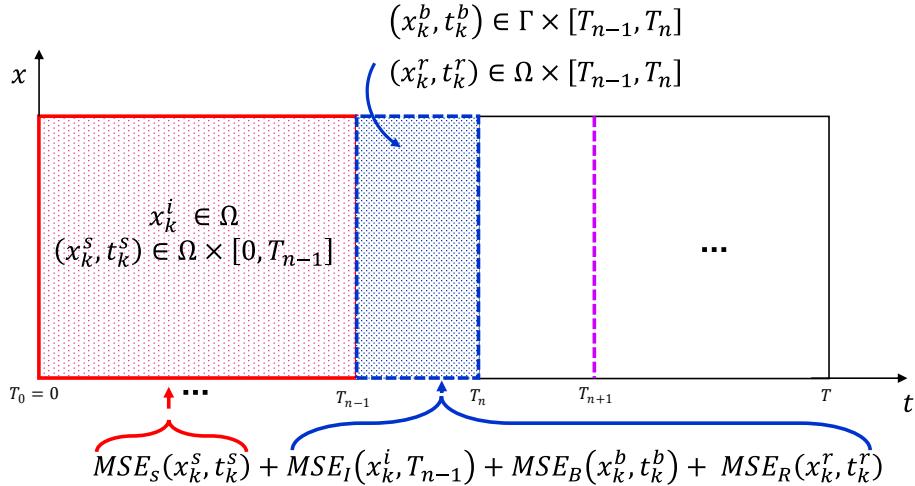


Figure 2: Illustration of the proposed backward compatibility scheme (over a 1D domain) that satisfies the solutions obtained on all previous time segments  $([0, T_{n-1}])$  while satisfying the PDE on the current time segment  $((T_{n-1}, T_n])$ . The error in the Initial Condition is computed at time  $t = 0$  for the first time segment and at time  $t = (T_{n-1})$  for the  $n$ th time segment.

117 Here,  $(x_k^i, T_{n-1})$  represent the set of points where the error on initial condition is com-  
 118 puted and  $(x_k^b, t_k^b)$  represent the set of points where the error on boundary conditions is  
 119 computed within the time segment  $(T_{n-1}, T_n]$ . The residual/collocation error as given  
 120 in equation (5) is computed at the collocation points  $(x_k^r, t_k^r)$ . We also minimize the  
 121 departure from the already obtained solution that were stored at the grid points  $(x_k^s, t_k^s)$ .  
 122 The solution obtained (on  $(0, T_n]$ ) at the  $n$ th segment is stored for using it in the  $(n+1)$ th segment.  
 123  
 124 The weights  $(w_i, w_b, w_r, w_s)$  given in equation (11) help in faster convergence to the true solution.

125 These weights are utilized to scale the difference in the magnitude of the errors. In the present  
 126 work, the difference in the order of derivatives (spatial) and the number of spatial dimensions  
 127 are used for scaling. From a mathematical perspective, the weighting of the loss function can be  
 128 seen as a mechanism that forces the learning process to focus on the terms where high prediction  
 129 accuracy is required.

130

131 In sections (3.2)3.3, two techniques have been proposed to further improve the accuracy and  
 132 efficiency of the bc-PINN scheme.

133 *3.2. Initial condition guided learning (ICGL)*

In initial condition guided learning, the key idea is to perform the training in two stages for each time segment. In the first stage the neural network is trained to match only the initial condition of that time segment, using a small fraction of the total iterations. Therefore, the loss function for ICGL ( $MSE_{SI}$ ) is as follows:

$$MSE_{SI} = \frac{1}{N_{SI}} \sum_{k=1}^{N_{SI}} \left( \hat{h}(\mathbf{x}_k^{SI}, t_k^{SI}) - \underline{h}_k^{SI}(\mathbf{x}_k^{SI}, T_{n-1}) \right)^2, \quad (\mathbf{x}_k^{SI}, t_k^{SI}) \in \Omega \times (T_{n-1}, T_n] \quad (12)$$

134 Here,  $\hat{h}(\mathbf{x}, t)$  is the neural network prediction and  $\underline{h}(\mathbf{x}, t)$  is the known solution through the neural  
 135 network at the previous time step,  $T_{n-1}$ . Whereas, in the second stage the weights obtained in  
 136 the first stage are taken as the initial weights and the bc-PINN as described earlier is trained.  
 137 This is motivated by the fact that the solution of a PDE in a time segment is expected to be  
 138 close to the initial condition if the segment is small. Thus matching the initial condition brings  
 139 the neural network map closer to the true map. Since in the first step there are no derivative  
 140 calculations involved, it accelerates the training.

141 *3.3. Transfer learning based acceleration (TL)*

142 A transfer learning approach is implemented that uses the weights and biases from a bc-PINN  
 143 that has been trained on a different initial condition or the previous segment. The transfer  
 144 learning approach preserves the features from a previous training as reported in [33]. This helps  
 145 in faster convergence for a new initial condition. A notable advantage of this technique is that  
 146 the training time can be significantly reduced as the number of trainable parameters decreases.

147

148 The two proposed techniques proposed (*Initial condition guided learning* and *Transfer learning*  
 149 *based acceleration*) shows significant improvement in the accuracy of the bc-PINN - solution for  
 150 the Allen Cahn and Cahn Hilliard equations. These two techniques are independent of each  
 151 other and can be used either in tandem or individually.

152 *3.4. Details of the neural network of bc-PINN*

153 We have used a standard (deep) neural network with two input neurons consisting of the  
 154 spatial variables ( $\mathbf{x}$ ) and temporal variable ( $t$ ). The output of the neural network ( $\hat{h}(\mathbf{x}, t)$ )  
 155 approximates the solution of the PDE ( $h(\mathbf{x}, t)$ ). To avoid model bias due to input features of  
 156 different scales we have performed “min-max” normalization to scale the data uniformly. For  
 157 solving 1D Allen Cahn and Cahn Hilliard equations using bc-PINN the architecture of the neural  
 158 network chosen has 6 hidden layers with 128 neurons in each layer. Whereas, while solving 2D  
 159 Allen Cahn and Cahn Hilliard equations using the aforementioned architecture didn’t yield a  
 160 good performance. Thus, in order to choose the neural network architecture for solving 2D PDEs  
 161 (section (6)) the following approach has been used. Preference has been given to increasing

162 the number of neurons of a hidden layer rather than increasing the number of hidden layers.  
 163 This is backed by the fact that increasing hidden layers is more computationally intensive than  
 164 increasing the number of neurons. For both std-PINN and bc-PINN, *tanh* is chosen as the  
 165 activation function. Even though it's well known that *tanh* activation function has a problem of  
 166 vanishing gradient in very deep networks [34]. The advantages of *tanh* activation function are  
 167 that its continuous (range [-1,1] ) and differentiable. Since, *tanh* is a non-linear function it gives  
 168 neural network the capability to learn non-linear maps [35]. The neural network has more than  
 169 100,000 learning parameters which have been initialized using the “xavier initialization” [36]  
 170 technique. The optimization of the loss function and updating the learning parameters (weights  
 171 and biases of the neural network) is performed using the ADAM and LBFGS optimizers. The  
 172 learning rate for ADAM optimizer is considered as 0.001 with all other parameters as suggested  
 173 in [37]. Following std-PINN, after training the neural network using the ADAM optimizer we  
 174 again train it using the L-BFGS optimizer until one of the following stopping criteria is met:  
 175 (i) Maximum iterations are equal to 50,000, (ii) Maximum number of function evaluations are  
 176 equal to 50,000, (iii) Maximum number of line search steps (per iteration) equal to 50, (iv) The  
 177 maximum number of variable metric corrections used to define the limited memory matrix are  
 178 equal to 50, (v) The iteration stops when  $\frac{f^k - f^{k+1}}{\max(|f^k|, |f^{k+1}|, 1)} \leq 2.22044604925e - 16$ , where  $f$  is  
 179 the neural network objective function and  $k$  is the iteration number.

### 180 3.5. Details of the Computational Platform

181 All the neural networks are trained on Nvidia Tesla P100 (3584 CUDA cores and 16GB of  
 182 HBM2 vRAM) and Nvidia Volta V100 GPU (5120 CUDA cores, 640 Tensor cores and 16GB of  
 183 HBM2 vRAM). For inferencing and generating the reference solutions via chebfun, we have used  
 184 Dell precision 3630 workstation with Intel core i7-9700k 8 core (4.9 GHz Turbo) and 32 GB  
 185 RAM. The software packages used for all the computations are Tensorflow 1.15 and MATLAB  
 186 R2020a. All the variables defined for computations in tensorflow are of float32 data type.  
 187

### 188 3.6. The reference solution

189 Accurate numerical solutions for the Allen Cahn and Cahn Hilliard equations are obtained  
 190 using the chebfun package [32]. The chebfun approach provides a polynomial interpolant  
 191 for smooth functions in Chebyshev points. To solve time varying PDEs an exponential time  
 192 differencing with Runge–Kutta time stepping scheme [38] has been implemented in chebfun,  
 193 which is used in the present work. Henceforth, these solutions are considered as the reference  
 194 solutions. We have taken 512 points for spatial discretization and 201 points for discretization in  
 195 time scale. A fourth order Runge–Kutta time integrator with time step  $\Delta t = 10^{-5}$  is used.  
 196

197 The bc-PINN approach is applied to solve Allen Cahn equation and Cahn Hilliard equation  
 198 in sections (5 and 6) to demonstrate its advantages for nonlinear and higher order PDEs in  
 199 comparison to std-PINN method [8].

## 200 4. bc-PINN for Allen Cahn and Cahn Hilliard equations

201 The Allen Cahn and Cahn Hilliard<sup>1</sup> equations are two of the widely used partial differential  
 202 equations for studying the phenomena of phase separation [39]. There are innumerable practical

---

<sup>1</sup>The Cahn Hilliard equation plays an essential role in the field of material science for describing the qualitative features in a phase separation process for two phase systems (assuming isotropy and constant temperature). The process of phase separation can be observed when a binary alloy is cooled down adequately. This leads to a state

203 applications of these equations in various fields such as material science [40, 41, 42, 43, 44, 45],  
 204 biological systems [46, 47, 48, 49, 50, 51, 52], electro-chemical systems etc. [53, 54, 55].

205 *4.1. Allen Cahn Equation*

206 For every  $\mathbf{x} \in \Omega$ , ( $\Omega$  is an open set of  $\mathbb{R}^D$ ) the Allen Cahn equation can be written as:

$$h_t - c_1^2 \nabla^2 h + f(h) = 0, \quad t \in (0, T], \quad \mathbf{x} \in \Omega \subset \mathbb{R}^D$$

$$f(h) = c_2(h^3 - h) \quad (13)$$

207 For a phase separation problem, the parameter  $h$ , represents the concentration of the individual  
 208 component and the parameter ‘ $c_1$ ’ represents the interfacial thickness. The solution progressively  
 209 develops interfaces separating different phases. For a given initial condition,  $h_0(\mathbf{x}) \in L^2(\Omega)$   
 210 and  $T > 0$ , we seek a function  $h : \Omega \times (0, T] \rightarrow \mathbb{R}$  which satisfies the above equation. In order  
 211 to implement the bc-PINN scheme for Allen Cahn equation, the following loss function has been  
 212 used.

- Mean squared error on the initial Condition

$$\text{MSE}_I = \frac{1}{N_i} \sum_{k=1}^{N_i} \left( \hat{h}(\mathbf{x}_k^i, T_{n-1}) - h_k^i \right)^2, \quad \mathbf{x}_k^i \in \Omega \quad (14)$$

- Mean squared error on the boundary Condition

$$\text{MSE}_B = \frac{1}{N_b} \sum_{k=1}^{N_b} \sum_{d=1}^{n_d} \left( \hat{h}^{(d-1)}(\mathbf{x}_k^b, t_k^b) - \hat{h}^{(d-1)}(-\mathbf{x}_k^b, t_k^b) \right)^2 \quad (\mathbf{x}_k^b, t_k^b) \in \Gamma \times (T_{n-1}, T_n] \quad (15)$$

213 where  $n_d$  is the order to which periodicity is enforced on the boundary  $\Gamma$ . Here, the  
 214 superscript,  $(\bullet)^b$  stands for boundary condition.

- The Mean squared error due to residual of the partial differential equation

$$R := \hat{h}_t - c_1^2 \nabla^2 \hat{h} + f(\hat{h})$$

$$\text{MSE}_R = \frac{1}{N_r} \sum_{k=1}^{N_r} (R(\mathbf{x}_k^r, t_k^r))^2, \quad (\mathbf{x}_k^r, t_k^r) \in \Omega \times (T_{n-1}, T_n] \quad (16)$$

215 The superscript,  $(\bullet)^r$  stands for residual of the PDE.

- Mean squared error for backward compatibility

$$\text{MSE}_S = \frac{1}{N_s} \sum_{k=1}^{N_s} \left( \hat{h}(\mathbf{x}_k^s, t_k^s) - \underline{h}(\mathbf{x}_k^s, t_k^s) \right)^2, \quad (\mathbf{x}_k^s, t_k^s) \in \Omega \times [0, T_{n-1}] \quad (17)$$

216 where,  $\hat{h}(\mathbf{x}, t)$  is the neural network prediction and  $\underline{h}(\mathbf{x}, t)$  is the known solution through  
 217 the neural network from the previous time steps  $\Omega \times [0, T_{n-1}]$ . The superscript,  $(\bullet)^s$  stands  
 218 for the backward compatible solution.

219 • The total mean squared error is the same as given in equation (11)

---

of total nucleation which is mainly referred to as spinodal decomposition. In the subsequent stage coarsening occurs in the nucleated microstructure at a much slower rate. This whole phase separation phenomena affects the mechanical properties (eg. strength, hardness and fracture toughness) of the material.

220 4.2. Cahn Hilliard Equation

221 For every  $\mathbf{x} \in \Omega$ , ( $\Omega$  is an open set of  $\mathbb{R}^D$ ) the Cahn Hilliard equation can be written as:

$$h_t - \nabla^2 (\kappa f(h) - (\alpha \kappa) \nabla^2 h(\mathbf{x}, t)) = 0, \quad t \in (0, T], \quad \mathbf{x} \in \Omega \subset \mathbb{R}^D \quad (18)$$

222 To simplify the derivative calculation, a phase space representation of the Cahn Hilliard equation  
223 has been adopted. The phase space representation is widely used to represent a high order PDE  
224 into coupled multiple lower order PDEs. The phase space representation of the Cahn Hilliard  
225 equation (a fourth order PDE, equation (18)) yields two coupled second order PDEs.

$$\begin{aligned} h_t - \nabla^2 (-(\alpha \kappa) \mu + \kappa f(h)) &= 0, \quad \mu = \nabla^2 h \quad t \in (0, T], \quad \mathbf{x} \in \Omega \subset \mathbb{R}^D \\ f(h) &= h^3 - h \end{aligned} \quad (19)$$

226 Since the entire process is governed by the Cahn Hilliard equation it is essential to understand  
227 the physical significance of each individual variable. The order parameter ( $h$ ) in equation (18),  
228 refers to the rescaled density or concentration of one of the material components in the system  
229 and it takes values between (-1 and 1, which corresponds to their pure states). The density of  
230 second component is  $(1 - h)$ , and this ensures that the total density over the simulation domain  
231 is a conserved quantity. The parameter  $\kappa$  is the mobility parameter and the parameter  $\alpha$  is  
232 related to the surface tension at the interface. In order to implement the bc-PINN scheme for  
233 the Cahn Hilliard equation, the following loss function has been used.

- Mean squared error on the initial condition for  $h(\mathbf{x}, t)$  and  $\mu(\mathbf{x}, t)$

$$\text{MSE}_I = \frac{1}{N_i} \left\{ \sum_{k=1}^{N_i} \left( \hat{h}(\mathbf{x}_k^i, T_{n-1}) - h_k^i \right)^2 + \sum_{k=1}^{N_i} \left( \hat{\mu}(\mathbf{x}_k^i, T_{n-1}) - \mu_k^i \right)^2 \right\}, \quad \mathbf{x}_k^i \in \Omega \quad (20)$$

- Mean squared error on the boundary Condition

$$\begin{aligned} \text{MSE}_{B_h} &= \frac{1}{N_b} \left\{ \sum_{k=1}^{N_b} \sum_{d=1}^{n_d} \left( \hat{h}^{(d-1)}(x_k^b, t_k^b) - \hat{h}^{(d-1)}(-x_k^b, t_k^b) \right)^2 \right\} \\ \text{MSE}_{B_\mu} &= \frac{1}{N_b} \left\{ \sum_{k=1}^{N_b} \sum_{d=1}^{n_d} \left( \hat{\mu}^{(d-1)}(x_k^b, t_k^b) - \hat{\mu}^{(d-1)}(-x_k^b, t_k^b) \right)^2 \right\} \\ \text{MSE}_B &= \text{MSE}_{B_h} + \text{MSE}_{B_\mu}, \quad (x_k^b, t_k^b) \in \Gamma \times (T_{n-1}, T_n] \end{aligned} \quad (21)$$

234 Here, the superscript,  $(\bullet)^b$  stands for boundary condition.

- The Mean squared error due to residual of the partial differential equation

$$\begin{aligned} R_1 &:= \hat{h}_t - \nabla^2 \left( -(\alpha \kappa) \mu + \kappa f(\hat{h}) \right) \\ R_2 &:= \hat{\mu} - \nabla^2 \hat{h} \\ \text{MSE}_R &= \frac{1}{N_r} \left\{ \sum_{k=1}^{N_r} R_1(\mathbf{x}_k^r, t_k^r)^2 + \sum_{k=1}^{N_r} R_2(\mathbf{x}_k^r, t_k^r)^2 \right\}, \quad (\mathbf{x}_k^r, t_k^r) \in \Omega \times (T_{n-1}, T_n] \end{aligned} \quad (22)$$

235 The superscript,  $(\bullet)^r$  stands for residual of the PDE.

- Mean squared error for backward compatibility

$$\text{MSE}_S = \frac{1}{N_s} \sum_{k=1}^{N_s} \left( \hat{h}(\mathbf{x}_k^s, t_k^s) - \underline{h}(\mathbf{x}_k^s, t_k^s) \right)^2, \quad (\mathbf{x}_k^s, t_k^s) \in \Omega \times [0, T_{n-1}] \quad (23)$$

236 where,  $\hat{h}(\mathbf{x}, t)$  is the neural network prediction and  $\underline{h}(\mathbf{x}, t)$  is the known solution through  
 237 the neural network from the previous time steps  $\Omega \times [0, T_{n-1}]$ . The superscript,  $(\bullet)^s$  stands  
 238 for the backward compatible solution.

239 • The total mean squared error is the same as given in equation (11)

240 The boundary loss (equation (21) is applied for  $n_d = 1$  on  $\hat{h}$  and  $\hat{\mu}$ , to represent periodic boundary  
 241 conditions. Equation (22) describes two components of residual for two PDEs in the phase space  
 242 form of the Cahn Hilliard equation (equation (19)). In the following sections (5 and 6), the  
 243 results of Allen Cahn and Cahn Hilliard equations obtained using bc-PINN are presented..

244 **5. Results of Allen Cahn and Cahn Hilliard equations using bc-PINN in 1D**

245 *5.1. Allen Cahn equation in one dimension*

In this section, a 1D time varying Allen Cahn equation has been considered. The PDE for Allen Cahn equation remains the same as described in section (4.1). Since, the domain considered is one dimensional,  $\Omega \in [-1, 1]$  and  $t \in (0, 1]$ . The values of the parameters considered in equation (13) are,  $c_1^2 = 0.0001$  and  $c_2 = 5$ .

$$\begin{aligned} h(x, 0) &= x^2 \cos(\pi x) \\ h(x, t) &= h(-x, t), \quad (x, t) \in \Gamma \times (0, T] \\ h_x^{(1)}(x, t) &= h_x^{(1)}(-x, t), \quad (x, t) \in \Gamma \times (0, T] \end{aligned} \quad (24)$$

246 The above equation (24) gives details about the initial and boundary conditions (where  $\Gamma$   
 247 describes the boundary).

248 *5.1.1. std-PINN for Allen Cahn equation*

249 At first, the aforementioned Allen Cahn equation is solved using the std-PINN to demonstrate  
 250 the challenge associated with non-linearity. The number of collocation points considered for  
 251 training the std-PINN are 512 points for initial condition, 201 points for the boundary condition  
 252 and 20,000 spatio-temporal points for computing the residual. The neural network architecture  
 253 used has 4 hidden layers with 200 neurons in each layer and which is the same as given in (8).  
 254 For training the std-PINN, we have used both ADAM and L-BFGS optimizers. Training is  
 255 performed using 100,000 ADAM iterations and the subsequent training has been performed  
 256 using the L-BFGS optimization method until one of the stopping criteria is met: (i) Maximum  
 257 iterations are equal to 50,000 (ii) Maximum number of function evaluations are equal to 50,000  
 258 (iii) Maximum number of line search steps (per iteration) equal to 50 (iv) The maximum number  
 259 of variable metric corrections used to define the limited memory matrix are equal to 50 (v) The  
 260 iteration stops when  $\frac{f^k - f^{k+1}}{\max(|f^k|, |f^{k+1}|, 1)} \leq 2.22044604925e-16$ , where  $f$  is the neural network  
 261 objective function and  $k$  is the iteration number. The loss function for std-PINN is described in  
 262 equation (3), (4) and (5). The solution of std-PINN is quite erroneous as shown in figure (5).  
 263 In order to understand the reason for failure of the std-PINN, we analyze its prediction for the  
 264 individual terms of the Allen Cahn equation. Figure (4) shows the individual terms of the Allen

265 Cahn equation obtained through the Chebfun method and the std-PINN. We observe that the  
 266 std-PINN fail to predict the non-linear term ( $5(h^3 - h)$ ) of the Allen Cahn equation. Therefore  
 267 we have shown that the std-PINN [8] does not work for the Allen Cahn equation that consist of  
 268 a strongly non-linear term.

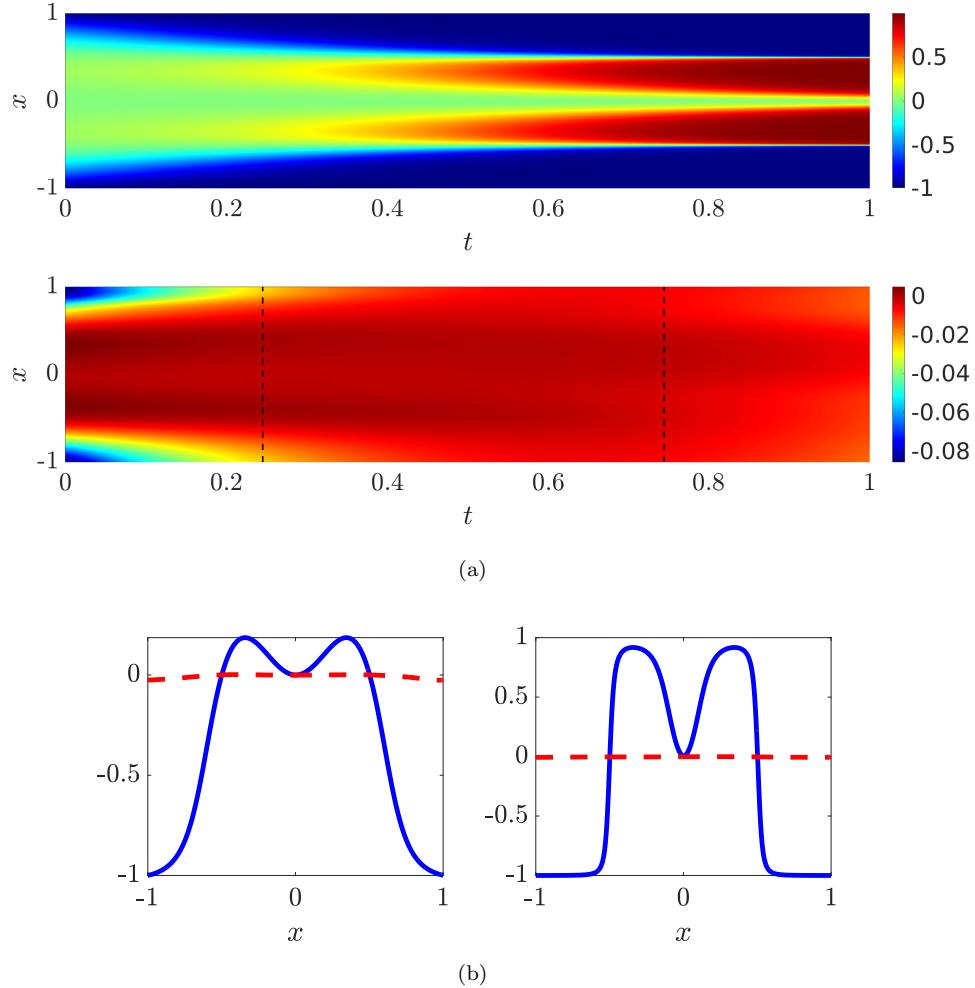


Figure 3: (a): The reference solution (Top) and the std-PINN solution (Bottom) of the Allen Cahn equation for the entire spatio-temporal domain. (b): Time snapshots for the reference solution (—) and the std-PINN solution (---) at  $t = 0.25$  and  $t = 0.75$ .

269 *5.1.2. bc-PINN for Allen Cahn equation*

270 To overcome this limitation of std-PINN, the backward compatible PINN approach has  
 271 been used along with ICGL and TL techniques. Therefore for the proposed bc-PINN approach,  
 272 the loss function given in equations (11,14,17) is used for any given time segment  $\Delta T_n$ . The  
 273 hyper-parameters associated with training the bc-PINN are number of ADAM iterations ( $N_{iter}$ ),  
 274 time steps per segment and number of residual collocation points ( $N_r$ ) per segment.

275

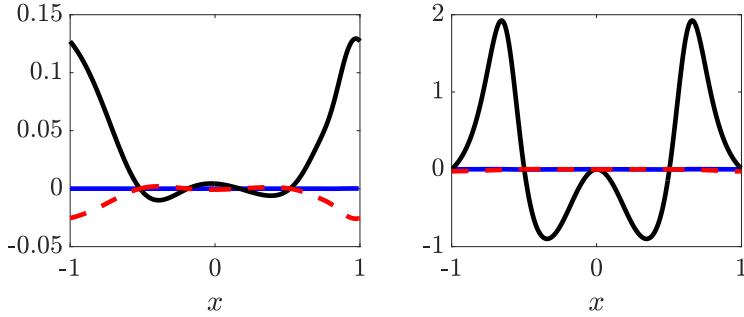


Figure 4: Individual terms of the Allen Cahn Equation obtained through the (Left): std-PINN method (Right): Chebfun method.  $h(x, t)$ , (---),  $0.0001\nabla^2 h$  (—) and  $5(h^3 - h)$  (—) at  $t = 0.25$ .

Variable	Description	Number
$N_i$	Initial collocation points	128
$N_b$	Boundary collocation points	50/segment
$N_r$	Residual collocation points	20000/segment
$N_{iter}$	Number of ADAM iterations	10000/segment

Table 1: Description of Training Data for Allen Cahn Equation. The segment considered here consists of 50 time steps.

276 In table 1, initial collocation points ( $N_i$ ) refer to the input spatio-temporal points where  
 277 the initial condition is prescribed. ( $N_b$ ) and ( $N_r$ ) refer to the number of boundary and residual  
 278 collocation points respectively and the output of the neural network at these points is used to  
 279 compute the loss function. For computing the backward compatibility loss the solution predicted  
 280 by the neural network in all the previous time steps is utilized. To reiterate, the reference  
 281 solution is only available at the initial condition whereas for all other points in the entire domain  
 282 the solution of the PDE is obtained through minimizing the bc-PINN loss function.

283  
 284 The reference and predicted solution at time  $t = 0.25$  obtained by the std-PINN and bc-PINN  
 285 are shown in figure 5. While the std-PINN fails, the proposed bc-PINN predicts the solution  
 286 quite accurately. The relative total errors ( $\varepsilon_{total}$ ) and the total training time for std-PINN,  
 287 XPINN, bc-PINN and bc-PINN with ICGL and TL approaches are shown in table 2. By using  
 288 techniques like ICGL and TL we have been able to observe a significant improvement in accuracy  
 289 and also reduction in training time. In order to implement the XPINN 14 for the time-varying  
 290 Allen Cahn equation the entire domain has been decomposed into 5 sub-domains sequentially  
 291 across time and each sub-domain is trained by a sub-net. Interfacial solution continuity across  
 292 different sub-domains has also been implemented and trained using 50,000 ADAM iterations  
 293 and L-BFGS optimizer with the same stopping criteria as bc-PINN. But the main drawback  
 294 of XPINN in solving forward problems is that while training subnet-1 all other subsequent  
 295 subnets are also trained which increases the computational cost. For example, before even the  
 296 solution converges in subnet-1, subnet-2 is simultaneously being trained which searches for the  
 297 solution of the PDE in an infinite-dimensional space in other words subnet-2 is trying to predict  
 298 the solution with an incorrect initial condition, which has not yet converged to the correct solution.

299  
 300 The comparison between the predicted solution using bc-PINN and the reference solution  
 301 is shown in figure 6. This shows that the bc-PINN can accurately predict the solution for  
 302 the entire domain. The solutions and errors by the std-PINN and bc-PINN are compared in

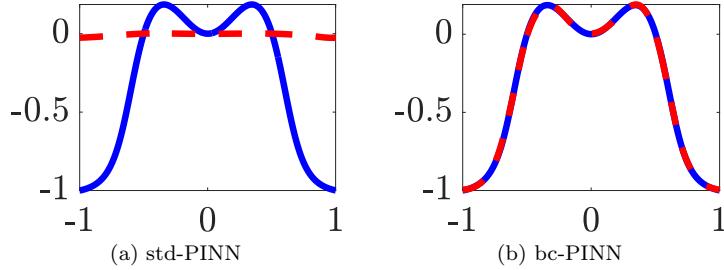


Figure 5: Reference (—) and Predicted (---) solution at time  $t = 0.25$

Method	Error( $\varepsilon_{\text{total}}$ )	Training Time
std-PINN	0.9919	4.5hrs
XPINN	0.9612	4 hrs
bc-PINN	0.0701	2 hrs
bc-PINN with ICGL and TL	0.0168	0.75 hrs

Table 2: Relative total errors (equation (7)) over the entire domain with respect to Chebfun solution for different methods.

303 figure (7), showing much higher accuracy by the bc-PINN. The error plots confirms high accuracy  
 304 of bc-PINN. The error increases with time very slowly. This is due to two reasons: (i) the solution  
 305 becomes progressively phase-separated (between zero and one) yielding greater curvatures and  
 306 sharp phase-boundaries that are difficult to capture, and (ii) due to the sequential nature of  
 307 the bc-PINN approach the error accumulates with time progression, which is similar to the  
 308 time-integrators. To illustrate the high accuracy of the bc-PINN approach, solutions and errors  
 309 for different values of the interfacial thickness ( $c_1$ ) is plotted in figure (8). As we decrease the  
 310 parameter  $c_1$ , it can be seen that the error in the prediction decreases. The parameter  $c_1$  controls  
 311 the effect of the double derivative of the solution ( $\nabla^2 h$ ). Therefore, as we decrease  $c_1$  the error  
 312 due to the approximation in derivative reduces and thus the accuracy of the bc-PINN solution  
 313 increases. In Appendix B, a new loss function including a logarithmic residual for the Allen  
 314 Cahn equation is discussed. This new logarithmic residual bc-PINN approach and its results are  
 315 presented in comparison with the simple bc-PINN approach without a logarithmic residual.

316 *5.2. Cahn Hilliard equation in one dimension*

317 In this section, a 1D time varying Cahn Hilliard equation has been considered. The PDE  
 318 for Cahn Hilliard equation remains the same as described in section (4.2). Since, the domain  
 319 considered is one dimensional,  $\Omega \in [-1, 1]$  and  $t \in (0, 1]$ . The values of the parameters considered  
 320 in equation (18) are,  $\alpha = 0.02$  and  $\kappa = 1$ .

$$\begin{aligned}
 h(x, 0) &= \cos(\pi x) - \exp(-4(\pi x)^2) \\
 h(-x, t) &= h(x, t), \quad (x, t) \in \Gamma \times (0, T] \\
 h_x^{(1)}(-x, t) &= h_x^{(1)}(x, t), \quad (x, t) \in \Gamma \times (0, T] \\
 \mu(-x, t) &= \mu(x, t), \quad (x, t) \in \Gamma \times (0, T] \\
 \mu_x^{(1)}(-x, t) &= \mu_x^{(1)}(x, t), \quad (x, t) \in \Gamma \times (0, T]
 \end{aligned} \tag{25}$$

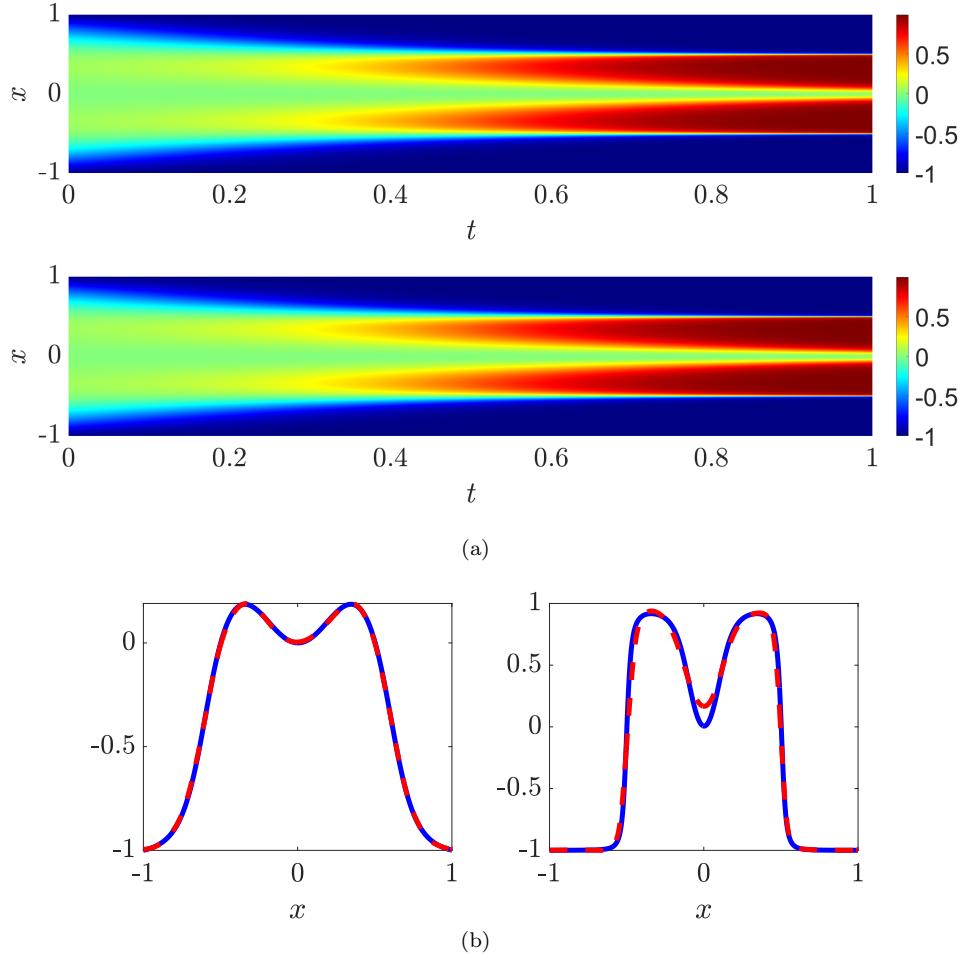
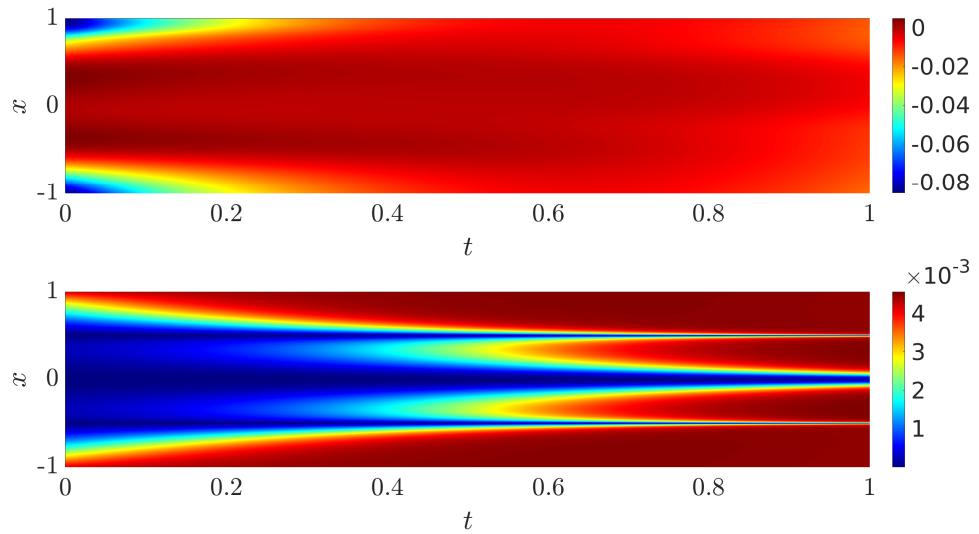


Figure 6: (a): Reference (Top) and bc-PINN (Bottom) solutions of the Allen Cahn equation for the entire spatio-temporal domain. (b): The Reference (—) and the bc-PINN (---) solutions at time  $t = 0.25$  and  $t = 0.75$ .

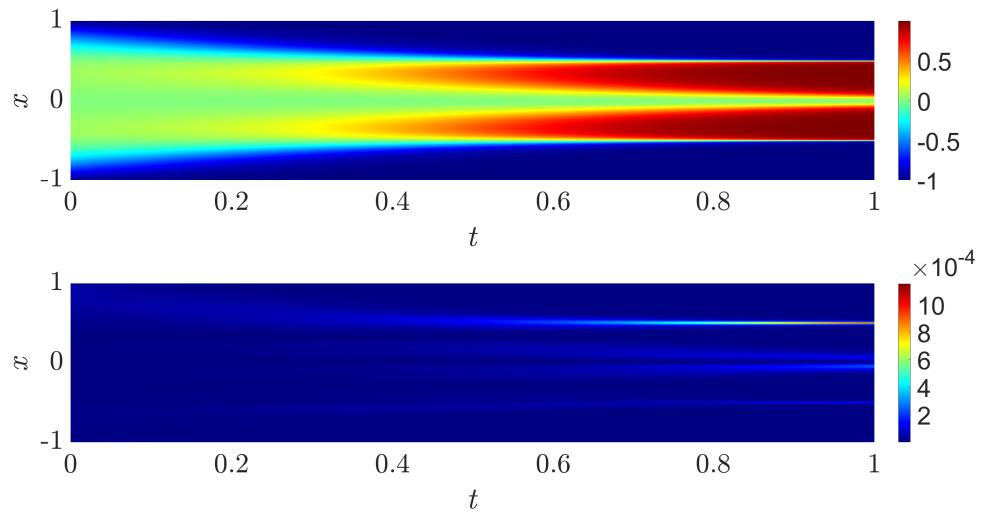
321 The above equation (25) gives details about the initial and boundary conditions (where  $\Gamma$   
322 describes the boundary).

### 323 5.2.1. std-PINN for Cahn Hilliard equation

324 Initially, the Cahn Hilliard equation (18) (without phase space) is solved using the std-PINN  
325 to demonstrate the challenge associated with high order. The neural network architecture  
326 used has 4 hidden layers with 200 neurons in each layer and which is the same as given in  
327 [8]. For training the std-PINN, we have used 20,000 collocation points and the loss function is  
328 minimized by using 100,000 ADAM iterations and subsequently L-BFGS optimizer until one  
329 of the stopping criteria are met: (i) Maximum iterations are equal to 50,000 (ii) Maximum  
330 number of function evaluations are equal to 50,000 (iii) Maximum number of line search  
331 steps (per iteration) equal to 50 (iv) The maximum number of variable metric corrections  
332 used to define the limited memory matrix are equal to 50 (v) The iteration stops when  
333  $\frac{f^k - f^{k+1}}{\max(|f^k|, |f^{k+1}|, 1)} <= 2.22044604925e - 16$ , where  $f$  is the neural network objective function and  $k$

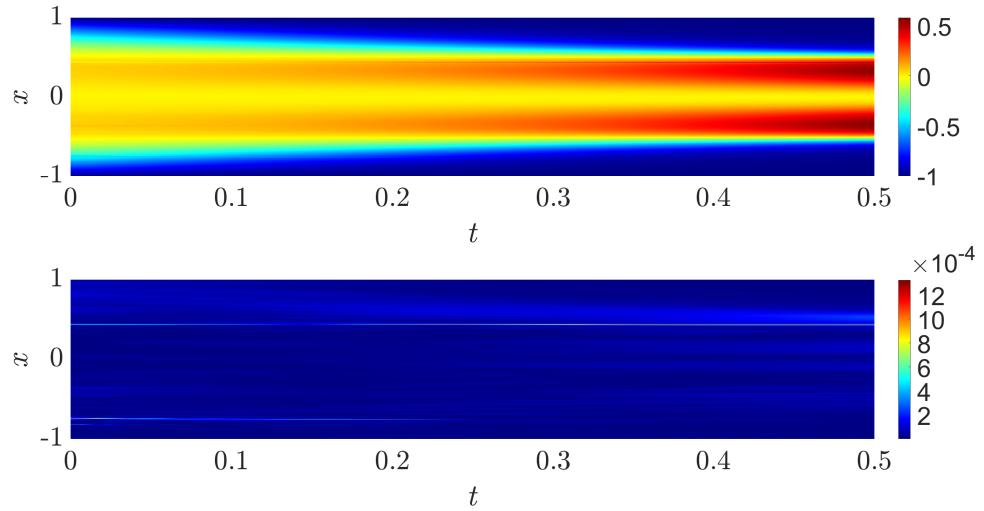


(a) Predicted solution (top) and relative error (bottom) obtained using std-PINN

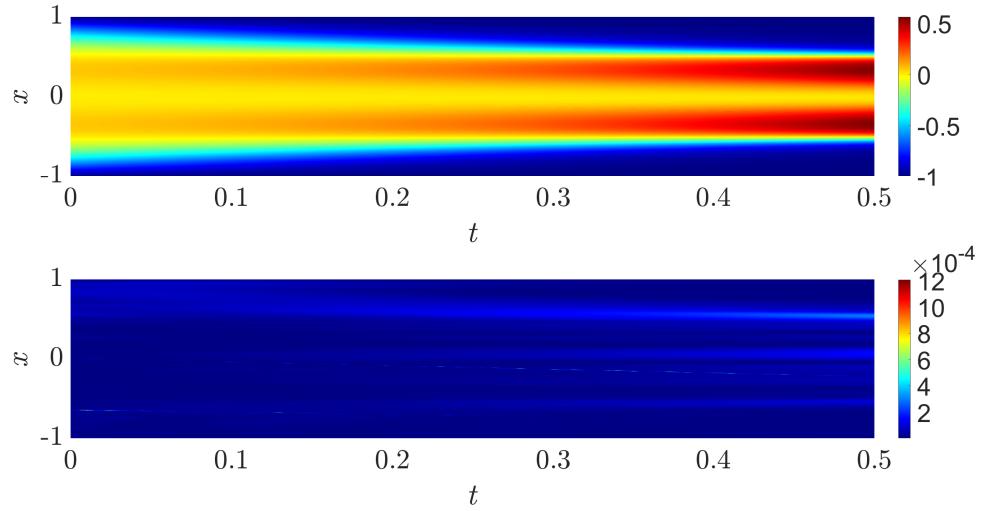


(b) Predicted solution (top) and relative error (bottom) obtained using bc-PINN

Figure 7: Solution and relative error associated with respect to the reference solution for Allen-Cahn equation.



(a) Predicted solution (top) and relative error (bottom) obtained using bc-PINN for  $c_1^2 = 0.00001$



(b) Predicted solution (top) and relative error (bottom) obtained using bc-PINN for  $c_1^2 = 0.00005$

Figure 8: Solutions and relative errors of the Allen Cahn equation for different  $c_1^2$  (of equation (13)) obtained by the bc-PINN method.

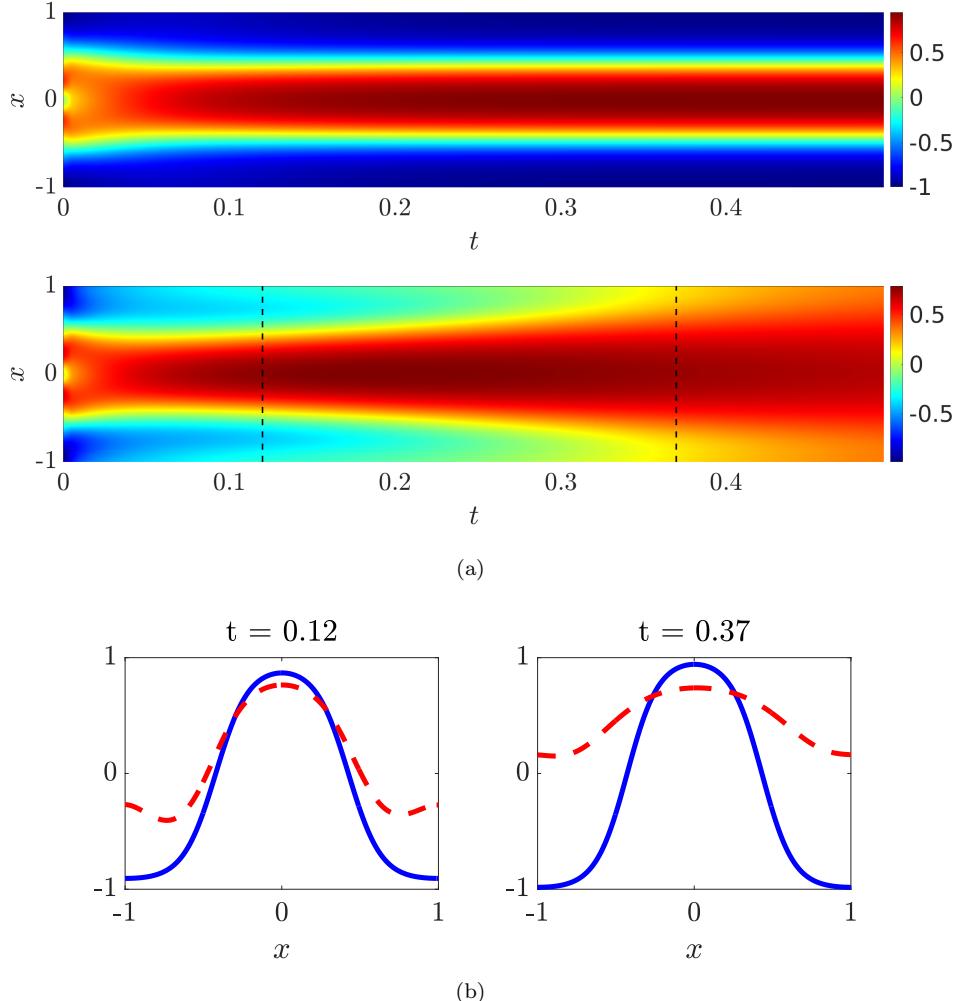


Figure 9: (a): The reference solution (Top) and the std-PINN solution (Bottom) of the Cahn Hilliard equation for the entire spatio-temporal domain. (b): Time snapshots for the reference solution (—) and the std-PINN solution (---) at  $t = 0.12$  and  $t = 0.37$ .

334 is the iteration number. The loss function for std-PINN is described in equation (3), (4) and (5).  
 335 The solution predicted after training is shown in figure (9) and it can be observed that there is  
 336 significant mismatch between the std-PINN prediction and the reference solution.

337  
 338 The two possible reasons for the inaccurate solution are strong non-linearity and the high  
 339 order derivative terms (fourth order). In PINN the derivatives are approximated using automatic  
 340 differentiation. It has been shown that as the order of the derivative increases the complexity in  
 341 automatic differentiation increases and it becomes computationally expensive [9]. In order to  
 342 overcome the difficulty in approximating the higher order derivative via automatic differentiation,  
 343 we adopt the phase space representation in the proposed bc-PINN.

344    5.2.2. *bc-PINN for Cahn Hilliard equation*

345    In this section, we introduce the bc-PINN approach with a phase space representation for  
 346 solving the Cahn Hilliard Equation (equation (19)). Additionally we have used the ICGL and TL  
 347 techniques as described in sections (3.2, 3.3). Therefore, there are two outputs of the neural  
 348 network  $\hat{h}(x, t)$  and  $\hat{\mu}(x, t)$  in the present method. The input features are the spatio-temporal  
 349 variables  $(x, t)$ . The modified loss function for the coupled phase space system includes an  
 350 error on initial condition, error on the boundary conditions and error on the residual. In  
 351 addition it will have the error for the backward compatibility. Therefore, the total loss function  
 352 (equation (11)) for any time segment  $\Delta T_n$  is sum of all the aforementioned errors given in  
 353 equation (20-23).

354

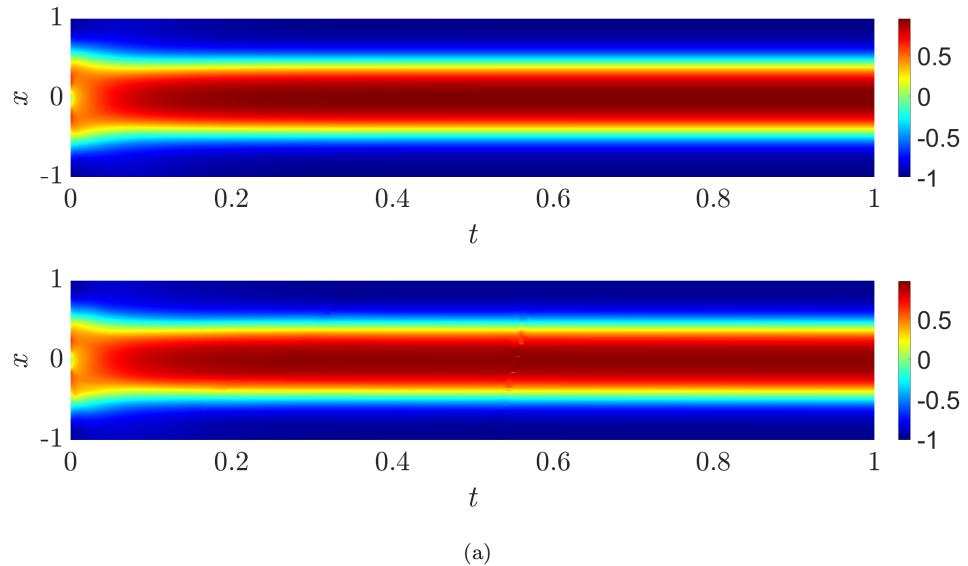
Variable	Description	Number
$N_i$	Initial collocation points	128
$N_b$	Boundary collocation points	20/segment
$N_r$	Residual collocation points	10000/segment
$N_{iter}$	Number of ADAM iterations	10000/segment

Table 3: Description of training data for Cahn Hilliard equation. 20 time steps/segment have been considered and the amount of collocation points generated remains same and doesn't increase as we progress through time.

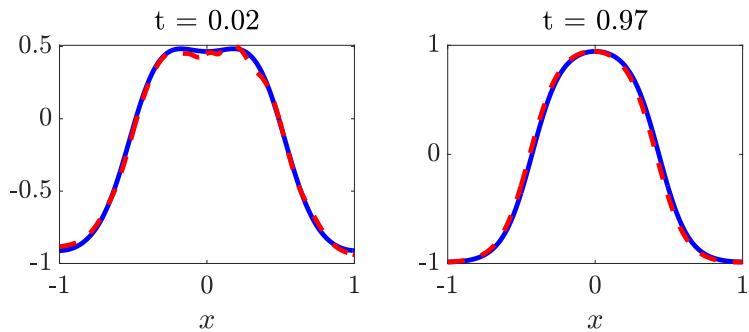
355    Table 3 describes the values of the hyper-parameters used in bc-PINN.  $N_i$  and  $N_b$  refers to  
 356 the number of points considered to enforce the initial and boundary condition respectively.  $N_r$   
 357 is the number of residual collocation points per time segment and  $N_{iter}$  is the number of ADAM  
 358 iterations used to train the neural network per time segment. As described in section 5.1.2,  
 359 only the reference solution at the initial points ( $N_i$ ) is used to compute the initial loss and for  
 360 computing the remaining terms in loss function (equation (20-23)) the output of the neural  
 361 network (predicted solution of the PDE) is used.

362

363    The accuracy of the proposed bc-PINN approach is shown by comparing it against the  
 364 reference solution obtained by the chebfun method in figure (10). This shows that the phase  
 365 space representation with bc-PINN can closely match the reference solution for the Cahn Hilliard  
 366 equation. The relative total error ( $\varepsilon_{total}$ ) obtained for the bc-PINN solution is 0.0186 whereas for  
 367 the std-PINN solution the error is 0.8594. It is evident from the error plots given in figure (11)  
 368 that a more accurate solution is obtained by using the bc-PINN compared to std-PINN. The  
 369 higher accuracy can be accredited to the fact that approximating lower order derivatives using  
 370 automatic differentiation is much simpler. One key observation to note is that the solution in  
 371 the  $n^{\text{th}}$  time segment takes the solution at time  $T_{n-1}$  from the  $(n-1)^{\text{th}}$  time segment as initial  
 372 condition. Thus, only the error at the end point in a time segment is propagated to the next time  
 373 segment. For instance, only the error at the time  $T_{n-1}$  in  $(n-1)^{\text{th}}$  time segment is propagated  
 374 to the next time segment. Errors at all other time steps in  $(n-1)^{\text{th}}$  time segment does not  
 375 propagate to the  $n^{\text{th}}$  time segment. This can be observed in figure (11b), even though the error  
 376 at time 0.01 is quite high but since this is not the end point of the time segment  $[0, 0.05]$  it does  
 377 not propagate with time. The error in the first time segment can be further reduced by using  
 378 more iterations and the accuracy of the total solution can be improved. To further demonstrate  
 379 the effectiveness of the current phase space backward compatible training approach, we have  
 380 taken different values of the parameter ( $\alpha\kappa$ ) and compared the predicted solutions with the  
 381 reference solutions generated using chebfun which is shown in figure (12). The proposed phase  
 382 space representation with bc-PINN approach can be extended to any partial differential equation  
 383 consisting higher order derivatives and non-linearity.

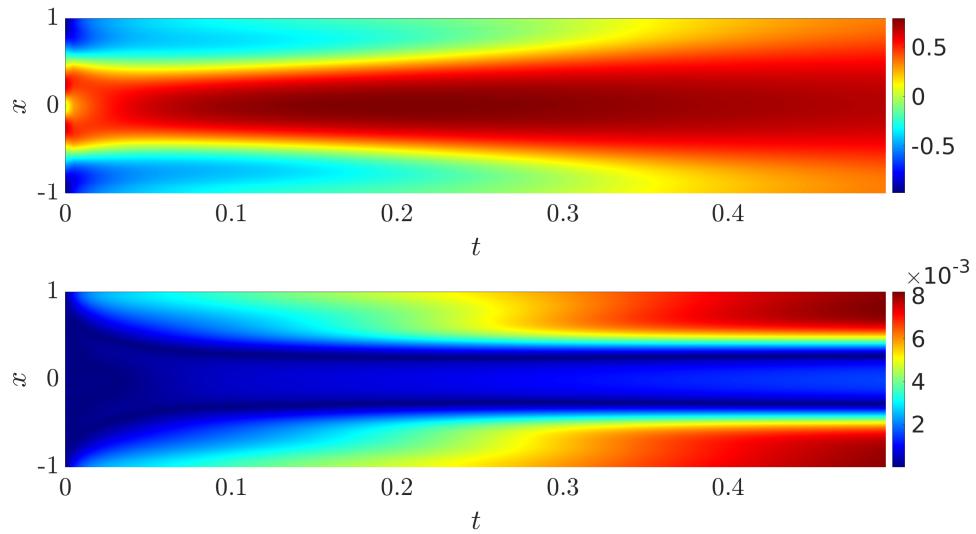


(a)

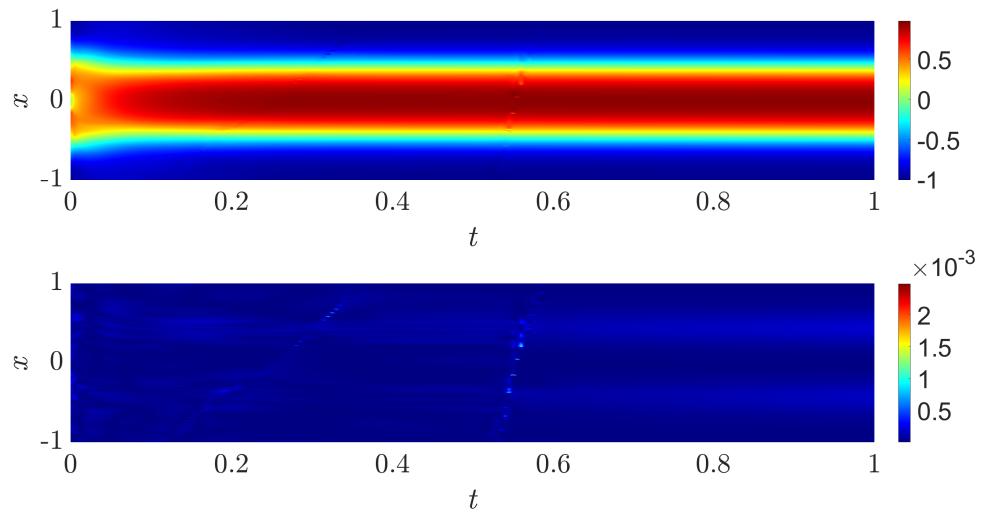


(b)

Figure 10: (a): Reference (Top) and bc-PINN (Bottom) solutions of the Cahn Hilliard equation for the entire spatio-temporal domain. (b): The Reference (—) and the bc-PINN (---) solutions at time  $t = 0.02$  and  $t = 0.97$ .

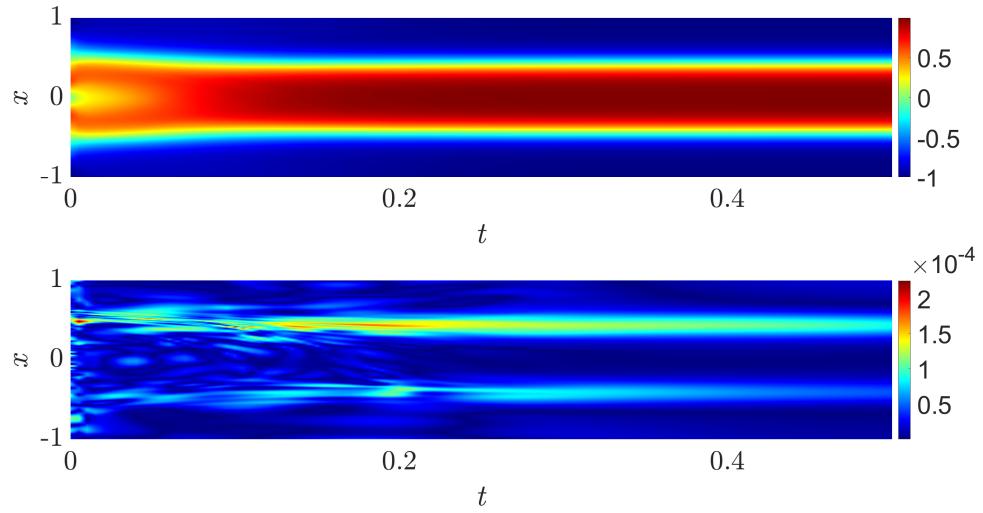


(a) Solution (top) and relative error (bottom) via std-PINN

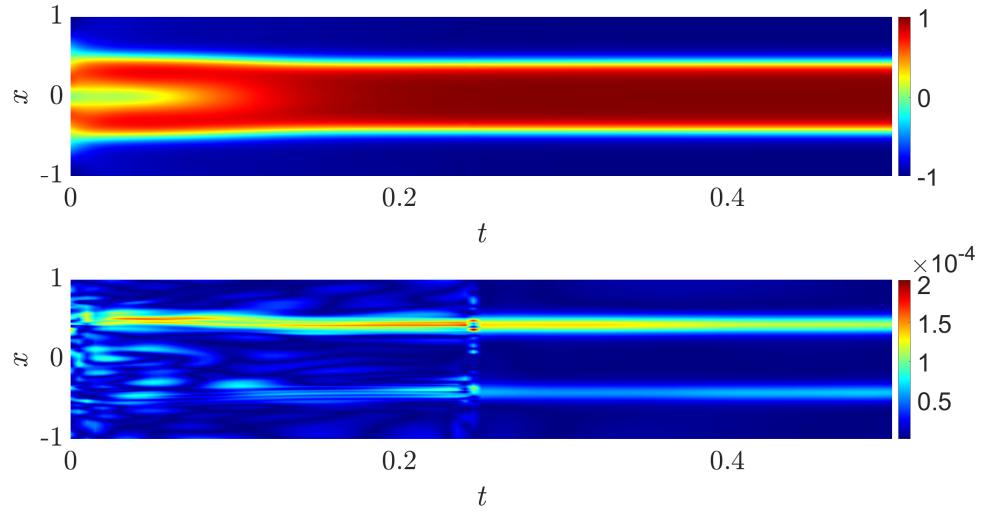


(b) Solution (top) and relative error (bottom) via bc-PINN

Figure 11: Solution and error associated with respect to the reference solution for Cahn-Hilliard equation.



(a) Solution (top) and error (bottom) for  $\alpha \kappa = 0.01$



(b) Solution (top) and error (bottom)  $\alpha \kappa = 0.005$

Figure 12: Solution and relative errors of the Cahn Hilliard equation for different parameters ( $\alpha \kappa$  of equation (18)) obtained by the bc-PINN method.

384 **6. Results of Allen Cahn and Cahn Hilliard equations using bc-PINN in 2D**

385 In this section, the Allen Cahn and Cahn Hilliard equations are solved in two dimensions  
 386 using the proposed bc-PINN scheme. Solving the Allen Cahn and Cahn Hilliard equations in 2D  
 387 is more computationally intensive than in 1D. Thus in order to solve the 2D equations, the  
 388 bc-PINN technique has been implemented along with the ICGL and TL techniques as proposed  
 389 in section (3).  
 390

391 *6.1. Allen Cahn equation in two dimensions*

392 In this section, a two dimensional time varying Allen Cahn equation has been considered.  
 393 The PDE for Allen Cahn equation remains the same as described in section (4.1). The values of  
 394 the parameters considered in equation (13) are,  $c_1^2 = 0.0001$  and  $c_2 = 1$ . To demonstrate the  
 395 proposed method following two IBVPs have been considered :

396 IBVP-1: The PDE used is the same as that given in equation (13). The domain for the IBVP  
 397 is taken as  $\mathbf{x} \times t \in [0, 1]^2 \times (0, 1]$ . The initial condition chosen is,  $\sin(4\pi x_1) \cos(4\pi x_2)$ , where,  
 398  $(x_1, x_2)$  are points in the domain  $[0, 1]^2$ . The boundary conditions have been considered to  
 399 be periodic,  $h^{(d-1)}(\mathbf{x}, t) = h^{(d-1)}(-\mathbf{x}, t)$ , for  $d = 1, 2$

400 IBVP-2: The PDE used is the same as that given in equation (13). The domain for the  
 401 IBVP is taken as  $\mathbf{x} \times t \in [0, 1]^2 \times (0, 2]$ . The initial condition chosen is a random doubly  
 402 periodic function where the maximum amplitude is 0.3. The boundary conditions have  
 403 been considered to be periodic,  $h^{(d-1)}(\mathbf{x}, t) = h^{(d-1)}(-\mathbf{x}, t)$ , for  $d = 1, 2$

404 *6.1.1. bc-PINN for Allen Cahn equation in two dimensions*

405 First in order to solve IBVP-1, the ICGL technique has been using along with bc-PINN. The  
 406 loss function for ICGL and bc-PINN remains the same as described in equations (12, 14, 17).  
 407 Furthermore, as described in section (3.6), Chebfun is used to obtain the reference solution.  
 408 Here, the spatial domain has been discretized into a grid containing 64 points along each  
 409 axis and the temporal domain has been discretized into 101 grid points. The neural network  
 410 architecture has 3 input neurons and consists of 6 hidden layers with 128 neurons in each layer.  
 411 The output layer contains only one neuron for the output/solution of the PDE ( $h$ ). Following, to  
 412 solve the IBVP-1, bc-PINN is used along with the ICGL technique. The loss function has  
 413 been optimized using both ADAM and L-BFGS optimizers. 15,000 ADAM iterations (20% for  
 414  $MSE_{SI}$  and 80% for  $MSE_{\Delta T_n}$ ) and 10,000 LBFGS iterations are used for every time segment.  
 415 Apart from the maximum iterations other stopping criteria for L-BFGS remain the same as  
 416 those mentioned in section (3.4) (criteria ii - v). The total number of collocation points used to  
 417 train the bc-PINN for IBVP-1 are given in table 4. The total error in prediction ( $\varepsilon_{total}$ ), using  
 418 bc-PINN with ICGL is 2.5%. Whereas, the total error in prediction using bc-PINN without  
 419 ICGL is more than 95%. Figure (13) shows the evolution of solution as time increases.  
 420

$N_i$	Initial collocation points	4096
$N_b$	Boundary collocation points	640 per segment
$N_r$	Residual collocation points	20,000 per segment

Table 4: Description of training data for 2D Allen Cahn equation (IBVP - 1). 20 time steps/segment have been considered and the amount of collocation points generated remains same and doesn't increase as we progress through time.

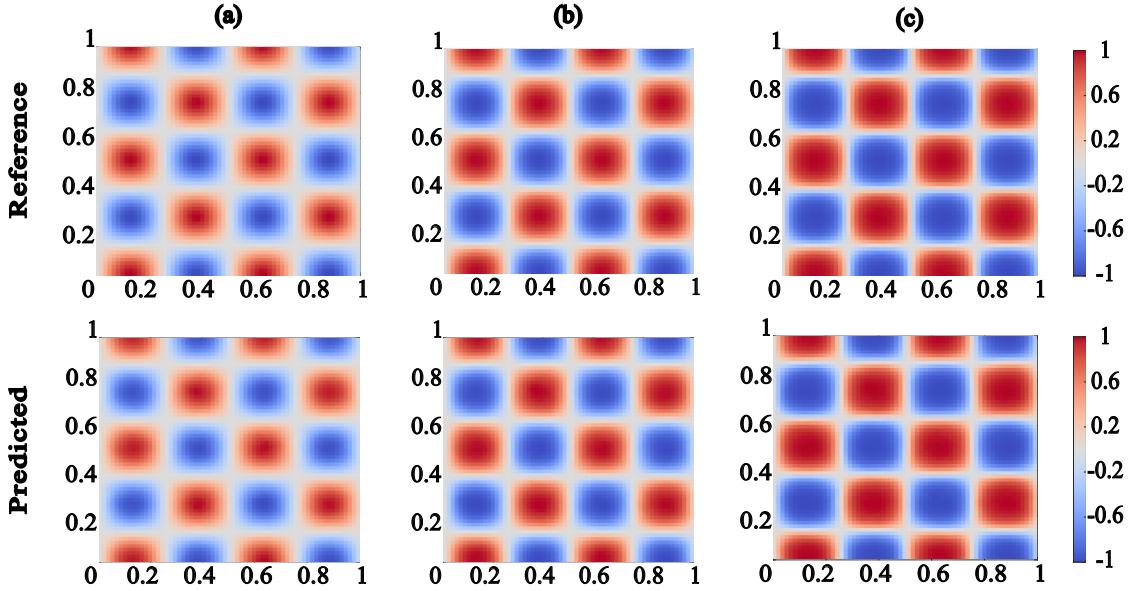


Figure 13: Reference and bc-PINN predicted solution of the 2D Allen Cahn equation (IBVP-1) at different time snapshots (a)  $t = 0$  , (b)  $t = 0.5$  , (c)  $t = 1$

421 To demonstrate bc-PINN with (ICGL and TL), IBVP-2 has been considered. The loss function for  
 422 ICGL, initial condition, and boundary condition remains the same as described in equations (12),  
 423 (14), (17). The total loss function also remains the same as given in equation (11). The optimization  
 424 of the loss function is performed using both ADAM and L-BFGS optimizers. For minimizing  
 425 the loss function in the first segment 30,000 ADAM iterations and 10,000 LBFGS iterations are  
 426 used. From the second segment 40,000 ADAM iterations and 10,000 LBFGS iterations are used  
 427 in optimizing the bc-PINN loss function. Apart from the maximum iterations other stopping  
 428 criteria for L-BFGS are the same as mentioned in section (3.4) (criteria ii - v). Also, during  
 429 the minimization process in each time segment, 20% of the total ADAM iterations are used  
 430 to implement the ICGL technique. The neural network architecture remains the same as that  
 431 used for solving IBVP - 1 (3 input neurons, 6 hidden layers with 128 neurons, and 1 output).  
 432 Following, the TL technique has been implemented by freezing the parameters obtained after  
 433 solving the IBVP-1. Initially, only a single layer is fixed. As the training progresses through  
 434 different time segments the number of layers frozen has been slowly incremented. Table (5) gives  
 435 details about the number of collocation points and the time segment details. From figure (14) it  
 436 can be observed that the bc-PINN prediction matches well with the reference solution, whereas  
 437 the std-PINN cannot solve this using the same number of collocation points.

$N_i$	Initial collocation points	4096
$N_b$	Boundary collocation points	1600 per segment
$N_r$	Residual collocation points	30,000 per segment

Table 5: Description of training data for 2D Allen Cahn equation (IBVP - 2). 50 time steps/segment have been considered and the amount of collocation points generated remains same and doesn't increase as we progress through time.

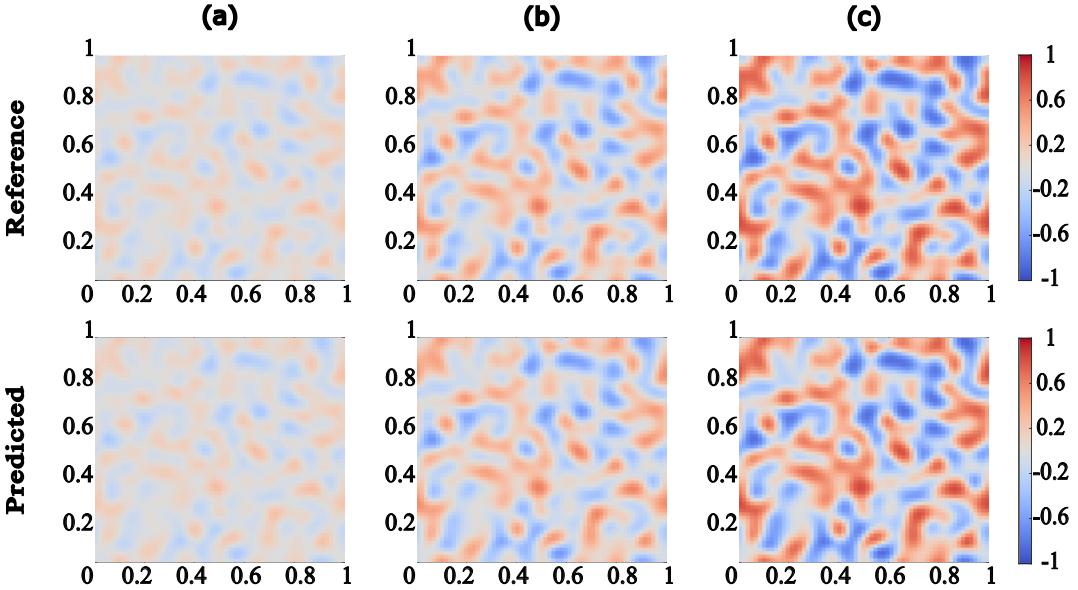


Figure 14: Reference and bc-PINN predicted solution of the 2D Allen Cahn equation (IBVP-2) at different time snapshots (a)  $t = 0$  , (b)  $t = 1.2$  , (c)  $t = 2$

438 *6.2. Cahn Hilliard equation in two dimensions*

439 In this section, a two dimensional time varying Cahn Hilliard equation has been considered.  
440 The PDE for the Cahn Hilliard equation remains the same as described in section (4.2). The  
441 values of the parameters considered in equation (19) are,  $(\alpha, \kappa) = (0.02, 1)$ . To demonstrate the  
442 applicability of bc-PINN for 2D Cahn Hilliard, the following two IBVP's has been considered:

443 **IBVP-1:** The PDE used is the same as that given in equation (19). The domain for the IBVP  
444 is taken as  $\mathbf{x} \times t \in [0, 1]^2 \times (0, 0.005]$ . The initial condition chosen is,  $0.4 \cos(3\pi x_1) \cos(3\pi x_2)$ ,  
445 where,  $(x_1, x_2)$  are points in the domain  $[0, 1]^2$ . Homogeneous Neumann boundary conditions  
446 have been considered.

447 **IBVP-2:** The PDE used is the same as that given in equation (19). The domain for the IBVP  
448 is taken as  $\mathbf{x} \times t \in [0, 1]^2 \times (0, 0.005]$ . The initial condition chosen is,  
449  $0.4 \cos(4\pi x_1) \cos(4\pi(x_1 + x_2))$ , where,  $(x_1, x_2)$  are points in the domain  $[0, 1]^2$ . Periodic  
450 boundary conditions have been considered.

451 **IBVP-3:** The PDE used is the same as that given in equation (19). The domain for  
452 the IBVP is taken as  $\mathbf{x} \times t \in [0, 1]^2 \times (0, 0.00375]$ . The initial condition chosen is a  
453 random doubly periodic function where the maximum amplitude is 0.5. Periodic boundary  
454 conditions have been considered.

455 *6.2.1. bc-PINN for Cahn Hilliard equation in two dimensions*

456 As described in section (3), the ICGL technique has been used along with bc-PINN to solve  
457 IBVP-1. A key point to be noted here is that the loss function for ICGL remains the same as  
458 that in equation (12). Therefore, the total loss function (equation (11)) for any time segment  
459  $\Delta T_n$  is a sum of all the aforementioned errors given in equation (20-23).

460

461 Further in order to generate the reference solution for IBVP-1, a cosine transform has been  
462 utilized as given in [56]. Here, the spatial domain is discretized into a grid containing 64 points  
463 along each axis. The temporal domain is discretized into 201 grid points. The neural network  
464 architecture has 3 input neurons and consists of 5 hidden layers with 128 neurons in each layer.  
465 Two output neurons have been taken to represent the output/solution of the PDE and the phase  
466 space term ( $h, \mu$ ). The minimization of bc-PINN loss function has been performed using ADAM  
467 and L-BFGS optimizers. The number of ADAM iterations are 30,000 for the first segment and  
468 for all the later segments 50,000 iterations have been used. The maximum number of L-BFGS  
469 iterations are 15,000 per segment and the other stopping criteria are the same as those mentioned  
470 in section (3.4) (criteria ii - v). The details of collocation points used to train the bc-PINN for  
471 IBVP1 are given in table (6). Figure (15), shows the evolution of  $h(\mathbf{x}, t)$  with time and there's a  
472 good match between the predicted and reference solution.

$N_i$	Initial collocation points	4096
$N_b$	Boundary collocation points	1600 per segment
$N_r$	Residual collocation points	50,000 per segment

Table 6: Description of training data for 2D Cahn Hilliard equation (IBVP-1). 50 time steps/segment have been considered and the amount of collocation points generated remains same and doesn't increase as we progress through time.

472

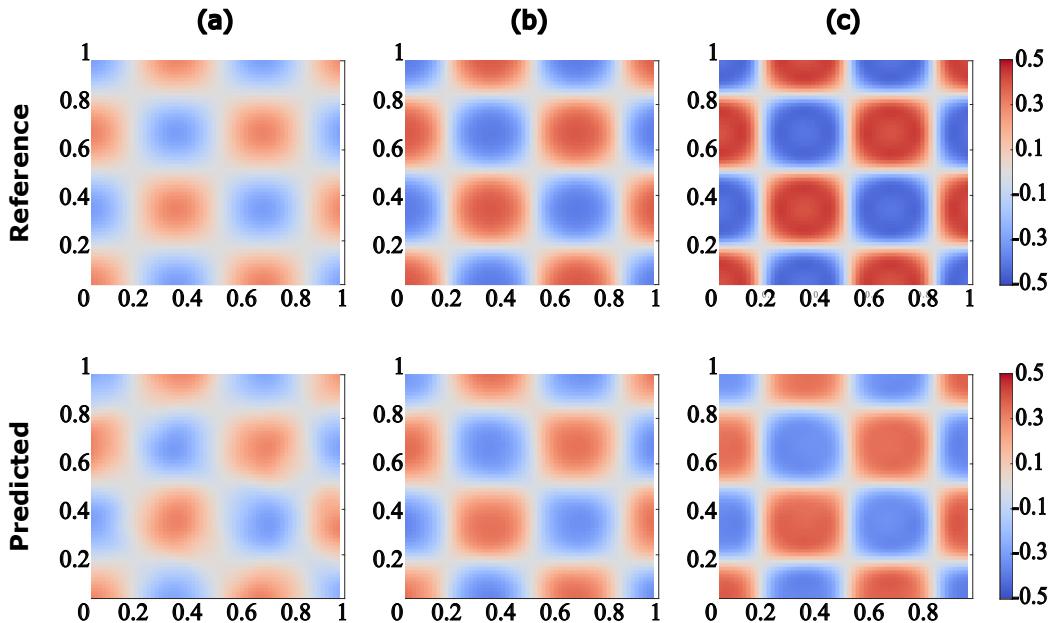


Figure 15: Reference and bc-PINN predicted solution of the 2D Cahn Hilliard equation (IBVP-1) at different time snapshots (a)  $t = 0$  , (b)  $t = 0.003$  , (c)  $t = 0.005$

473 In order to demonstrate bc-PINN with (ICGL and TL), IBVP-2 has been considered. To solve  
474 the IBVP-2, the loss function for ICGL, initial condition and boundary condition remain the  
475 same as described in equations (12), [20], [23], [11]. To generate the reference solution for IBVP-2  
476 and IBVP-3, an explicit time stepping scheme is used for time integration and a 9-stencil finite  
477 difference method for computing the spatial derivatives. Here, the spatial domain is discretized

$N_i$	Initial collocation points	4096
$N_b$	Boundary collocation points	1600 per segment
$N_r$	Residual collocation points	50,000 per segment

Table 7: Description of training data for 2D Cahn Hilliard equation (IBVP-2,IBVP-3). 25 time steps/segment have been considered and the amount of collocation points generated remains same and doesn't increase as we progress through time.

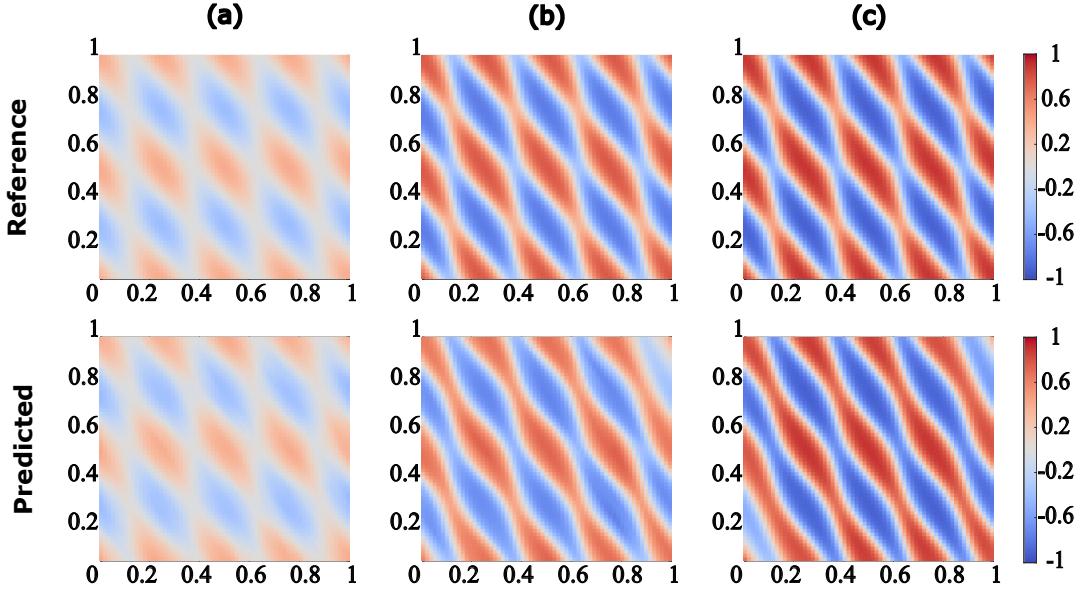


Figure 16: Reference and bc-PINN predicted solution of the 2D Cahn Hilliard equation (IBVP-2) at different time snapshots (a)  $t = 0$  , (b)  $t = 0.003$  , (c)  $t = 0.005$

478 into a grid containing 64 points along each axis, whereas, the temporal domain is discretized  
479 into 101 grid points. The neural network architecture has 3 input neurons and consists 5 hidden  
480 layers with 128 neurons in each layer. Two output neurons have been taken to represent the  
481 output/solution of the PDE and the phase space term  $(h, \mu)$ . The neural network architecture  
482 remains the same as that used for solving IBVP - 1 (3 input neurons, 5 hidden layers with 128  
483 neurons and 2 outputs). Following, the TL technique has been implemented for IBVP2 by  
484 freezing the parameters obtained after solving the first segment of IBVP-1. Initially, only a  
485 single layer is fixed and as the training progresses through different time segments the number  
486 of layers frozen has been incremented by one after every time segment. The minimization of  
487 bc-PINN loss function for IBVP-2 has been performed using ADAM and L-BFGS optimizers.  
488 The number of ADAM iterations are 50,000 for all the segments. The maximum number of  
489 L-BFGS iterations are 75,000 per segment and the other stopping criteria are same as those  
490 mentioned in in section (3.4) (criteria ii - v). The details of collocation points used to train the  
491 bc-PINN for IBVP-2 are given in table (7). Figure (16) shows the evolution of  $h(\mathbf{x}, t)$  with time.  
492

Despite using all the aforementioned techniques (ICGL, TL and weighting of the loss function) the solution doesn't converge efficiently for IBVP-3. This is because the initial condition in this case is random and the evolution is complex. Therefore, we have used a sum of  $\mathcal{L}^1$  norm and  $\mathcal{L}^2$  norm of the individual error terms in the total loss function. The reason behind using such a

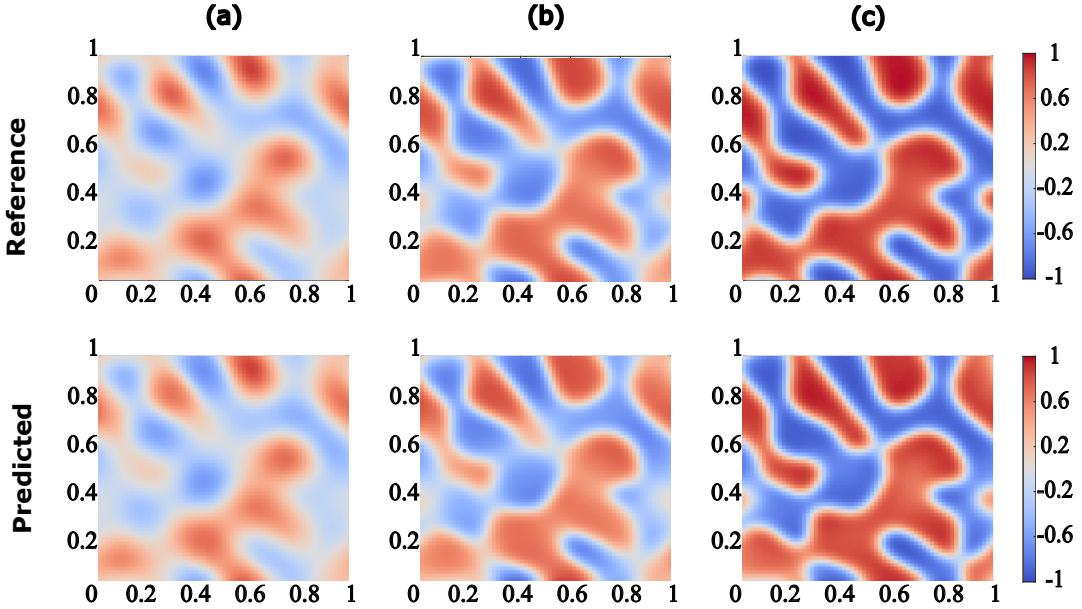


Figure 17: Reference and bc-PINN predicted solution of the 2D Cahn Hilliard equation (IBVP-3) at different time snapshots (a)  $t = 0$  , (b)  $t = 0.002$  , (c)  $t = 0.005$

loss function is that, when the error is small the penalization due to  $\mathcal{L}^1$  norm is more than  $\mathcal{L}^2$  norm. Thus, as given in equation (26) a combination of  $\mathcal{L}^1$  norm +  $\mathcal{L}^2$  norm has been used to solve the IBVP-3.

$$\text{MSE}_{\Delta T_n} = w_i (\text{MSE}_I + \text{MAE}_I) + w_b (\text{MSE}_B + \text{MAE}_B) + w_r (\text{MSE}_R + \text{MAE}_R) + w_s (\text{MSE}_S + \text{MAE}_S) \quad (26)$$

493 Here,  $\text{MAE}_I$ , is the mean absolute error on the initial condition ;  $\text{MAE}_B$ , is the mean absolute  
 494 error on the boundary condition;  $\text{MAE}_R$ , is the mean absolute error on the residual ( $R_1, R_2$ ) as  
 495 given in equations (22);  $\text{MSE}_S$ ; is the mean absolute error on backward compatibility;

496 Other hyper-parameters like neural network architecture, number of training iterations, number  
 497 of collocation points etc. remain the same as those used for solving IBVP-2. The only difference  
 498 is that the TL technique has been implemented for IBVP-3, by using the parameters obtained  
 500 after solving the first segment of IBVP-2. Considering the complexity of the Cahn Hilliard  
 501 equation from figure (17) it can be observed that the bc-PINN prediction matches the reference  
 502 solution remarkably well.

503 **7. Conclusions**

504 A new PINN approach (named as bc-PINN) has been proposed for solving the Allen Cahn  
 505 and Cahn Hilliard equations, however, the methodology in general should be applicable to  
 506 any PDE. The bc-PINN re-trains the neural network over successive time segments while  
 507 satisfying the solution for all previous time segments. Additionally, bc-PINN incorporates  
 508 two new techniques to improve the accuracy and efficiency of training. Firstly, while solving  
 509 a boundary value problem using bc-PINN, the initial condition of that segment is used to  
 510 bring the neural network map closer to the true map. Secondly, a transfer learning approach

511 is implemented to accelerate training, where the parameters learned from previous training  
512 are used to train a subsequent segment or a new initial condition. In addition by using  
513 a phase space representation for the Cahn Hilliard equation, better convergence has been achieved.

514

515 The key advantages of bc-PINN are summarized below. The proposed bc-PINN method can  
516 provide an accurate solution for nonlinear or higher-order PDEs such as the Cahn Hilliard and  
517 Allen Cahn equations in 1D and 2D. Moreover, the proposed method can achieve high accuracy  
518 by using fewer collocation points compared to std-PINN. Despite the segmentation of the time  
519 domain, it uses only one neural network and provides a continuous solution for the entire  
520 spatio-temporal domain. The proposed backward compatibility scheme may enhance many other  
521 machine learning approaches applied to complex systems represented by time dependent PDEs.

522

523 The code accompanying this manuscript would be available on Github repository:

524 <https://github.com/vmattey/bc-PINN>

525

526 **Acknowledgments:** SG acknowledges the financial support by NSF (CMMI MoMS) under  
527 grant number 1937983. We acknowledge Superior, a high-performance computing facility at  
528 MTU and Google Colab, a cloud service hosted by Google. This work used the Extreme Science  
529 and Engineering Discovery Environment (XSEDE) (allocation number MSS200004), which is  
530 supported by the NSF grant number ACI-1548562. RM acknowledges Ponkrshnan Thiagarajan  
531 and Shashank Pathrudkar for their valuable discussions and insights on the topic.

532 **Appendix A. Hyper-parameter selection for bc-PINN**

533 As discussed in section 5 and 6, the proposed method has a number of hyper-parameters like  
 534 number of ADAM iterations ( $N_{iter}$  per segment), time steps per segment, number of collocation  
 535 points ( $N_r$ ) etc. The accuracy of the bc-PINN's solution depends on proper choice of these  
 536 hyper-parameters. To optimize each of the hyper-parameters, various cases and metrics like  
 537 computational time and accuracy have been considered. In the current section, we chose the  
 538 Cahn Hilliard equation as the canonical example for all the analysis performed. Table (A.8)  
 539 describes the optimum parameters required to train 100 steps. The optimum parameters are  
 540 chosen to achieve an accurate solution while balancing the computational cost as shown in  
 541 figure (A.18). It can be also seen that as the number of collocation points and number of  
 542 iterations are increased the accuracy increases. Therefore, a segment size of 10 steps with 5000  
 543 collocation points and 10000 iterations per time segment has been chosen.

544 In general the bc-PINN framework is robust enough to handle any arbitrary length of time  
 545 segments and any number of time segments. One effective way for choosing a minimum time  
 546 segment is to try training the bc-PINN initially for a specific length of time segment where the  
 547 accuracy doesn't get affected. Based on this minimum length of time segment the total number  
 548 of time segments can be chosen depending on the length of the total time domain.

Model	Time steps/segment	$N_r$	$N_{iter}$
A	10	5000	10000
B	10	5000	20000
C	10	10000	10000
D	10	10000	20000
E	25	5000	10000
F	25	5000	20000
G	25	10000	10000
H	25	10000	20000

Table A.8: Parameter combinations for choosing the optimum segment size, collocation points and number of ADAM iterations to apply the bc-PINN technique for Cahn Hilliard equation.

550 **Appendix B. bc-PINN with a logarithmic residual for Allen Cahn Equation**

551 In this section we show how the bc-PINN with logarithmic residual compares against the  
 552 std-PINN and bc-PINN without the logarithmic residual. The loss function for the bc-PINN  
 553 with a logarithmic residual is same as the bc-PINN except the equation(16) is replaced by the  
 554 following loss term for the residual of the PDE:

$$R := \hat{h}_t - c_1^2 \nabla^2 \hat{h} + f(\hat{h})$$

$$\text{MSE}_R^{(\ln)} = \frac{1}{N_r} \sum_{k=1}^{N_r} \ln (1 + (R(x_k^r, t_k^r))^2), \quad (x_k^r, t_k^r) \in \Omega \times (T_{n-1}, T_n] \quad (\text{B.1})$$

555 It turns out that the bc-PINN with a logarithmic residual is more accurate than the bc-PINN  
 556 without the logarithmic residual. A possible explanation is that the logarithmic function reduces  
 557 the relative weight on the  $\text{MSE}_R$ , which would have larger inaccuracy due to its derivative and  
 558 nonlinear terms. Thus the initial and boundary terms are satisfied more accurately, which in

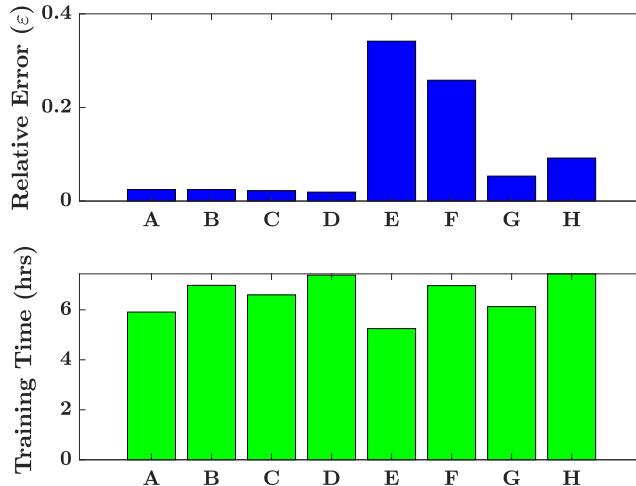


Figure A.18: Relative error ( $\varepsilon$ ) and time taken for various models given in table (A.8).

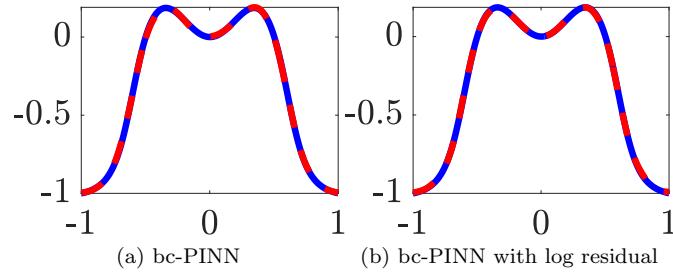


Figure B.19: Reference (—) and Predicted (---) solution at time  $t = 0.25$

Method	Error( $\varepsilon$ )
std-PINN	0.9919
bc-PINN	0.0701
bc-PINN with logresidual	0.03

Table B.9: Relative errors (equation (7)) over the entire domain with respect to Chebfun solution for different methods.

turn yields a more accurate solution. This explanation is substantiated by the fact that when the logarithmic function is used on all of the four terms in the loss function then the accuracy decreases.

### Appendix C. Minimization of the bc-PINN loss function

In section 3.4 we have described the learning rates and stopping criteria for the ADAM and LBFGS optimizer that are utilized to train the bc-PINN. Figure (C.20) shows the minimization of the loss function (equation (11)) while training the bc-PINN for 1D Cahn Hilliard equation in

566 time segment  $[0.45, 0.5]$ .

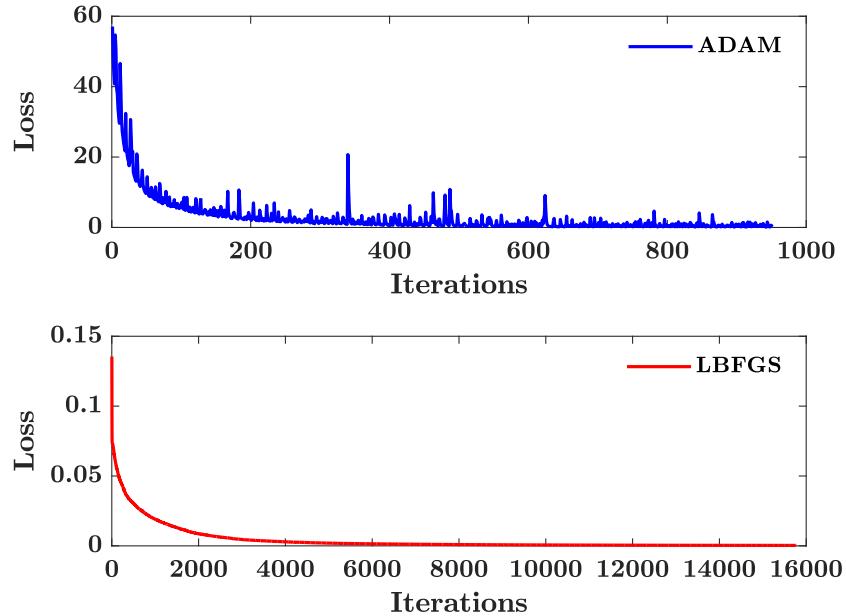


Figure C.20: Loss vs Iterations (Top): using ADAM optimizer (Bottom): using the LBFGS optimizer for training the time segment  $[0.45, 0.5]$  of the Cahn Hilliard equation.

567 **Appendix D. Data driven identification of parameters for PDEs using bc-PINN**

568 The framework of bc-PINN is versatile and can be applied to both forward and inverse  
 569 problems. Data driven methods for inverse problems have been shown to be extremely effective  
 570 [57, 24, 58, 59]. Especially in situations where only partial physics and noisy data are present  
 571 physics informed neural networks perform well. In this section, the performance of std-PINN [8]  
 572 and bc-PINN for non-parametric PDEs and parametric PDEs has been showcased. For the  
 573 non-parametric PDE case, Allen Cahn equation which has been described in section 4 will be  
 574 considered as the canonical example. Finally for the parametric PDE, burgers equation with a  
 575 time-varying parameter has been considered.

576 Consider a general,  $m^{th}$  order partial differential equation with parameters  $\lambda_1, \lambda_2$  etc. as given  
 577 in equation (D.1).

$$h_t = F(h(\mathbf{x}, t), h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), h_{\mathbf{x}}^{(2)}(\mathbf{x}, t), \dots, h_{\mathbf{x}}^{(m)}(\mathbf{x}, t), \lambda_1, \lambda_2, \dots, \lambda_n), \quad \mathbf{x} \in \Omega \subset \mathbb{R}^D, \quad t \in (0, T] \quad (\text{D.1})$$

577 In equation (D.1), we are interested in finding the parameters  $\lambda_1, \lambda_2 \dots \lambda_n$ . These parameters  
 578 are learnt by defining them as trainable parameters to the physics informed neural networks.  
 579 Therefore the loss function can be written as

- Mean squared error on the observed data

$$\text{MSE}_o = \frac{1}{N_o} \sum_{k=1}^{N_o} \left( \hat{h}(\mathbf{x}_k^o, t_k^o) - h_k^o \right)^2, \quad (\mathbf{x}_k^o, t_k^o) \in \Omega \times (0, T] \quad (\text{D.2})$$

- Mean squared error of the residual at observed data points

$$\begin{aligned} R &:= h_t - F(h(\mathbf{x}, t), h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), h_{\mathbf{x}}^{(2)}(\mathbf{x}, t), \dots, h_{\mathbf{x}}^{(m)}(\mathbf{x}, t), \lambda_1, \lambda_2, \dots, \lambda_n) \\ \text{MSE}_{oR} &= \frac{1}{N_o} \sum_{k=1}^{N_o} (R(\mathbf{x}_k^o, t_k^o))^2, \quad (\mathbf{x}_k^o, t_k^o) \in \Omega \times (0, T] \end{aligned} \quad (\text{D.3})$$

580 where  $\hat{h}(\mathbf{x}_k^o, t_k^o)$  is the neural network output and  $h_k^o$  is the observed data at  $(\mathbf{x}_k^o, t_k^o)$ . Here,  
581 the superscript,  $(\bullet)^o$  stands for observed data.

- The total mean squared error for the inverse std-PINN is given as

$$\text{MSE} = \text{MSE}_o + \text{MSE}_{oR} \quad (\text{D.4})$$

582 The equations (D.2) (D.3) (D.4) represent the total loss function of std-PINN for inverse identification  
583 of parameters. In order to implement the bc-PINN scheme for inverse problems, the domain  
584 has been divided into multiple segments as shown in figure (D.21). To identify the parameters  
585 in a particular time segment, the residual has been minimized only on the observed data in  
586 that particular segment. Additionally, to implement the backward compatibility scheme the  
587 solution of all the previous segments is satisfied simultaneously. Moreover, the parameters for  
588  $n^{th}$  time segment are initialized from the parameters learned in the  $(n-1)^{th}$  time segment.  
589 This framework helps in identifying the parameters of the PDE in a given segment and also  
590 learning the solution in the entire domain. The loss function for parameter identification using  
591 the bc-PINN framework is as follows:

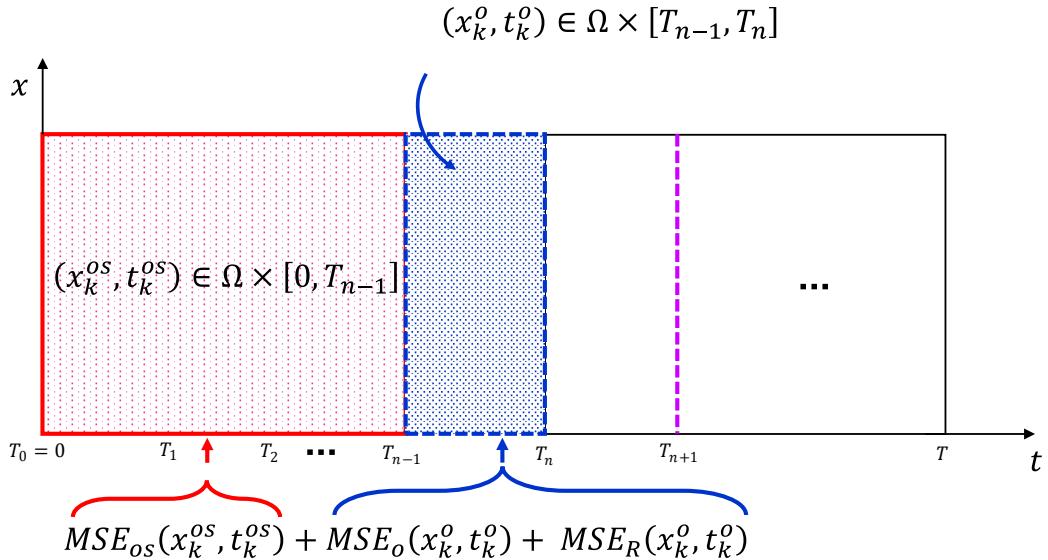


Figure D.21: Illustration of the proposed backward compatibility scheme that satisfies the data observed in the previous time segments at the same time learning the observed data and parameters in a particular segment.

- Mean squared error on the observed data

$$\text{MSE}_o = \frac{1}{N_o} \sum_{k=1}^{N_o} \left( \hat{h}(\mathbf{x}_k^o, t_k^o) - h_k^o \right)^2, \quad (\mathbf{x}_k^o, t_k^o) \in \Omega \times [T_{n-1}, T_n] \quad (\text{D.5})$$

- Mean squared error of the residual at observed data points

$$\begin{aligned} R &:= h_t - F(h(\mathbf{x}, t), h_{\mathbf{x}}^{(1)}(\mathbf{x}, t), h_{\mathbf{x}}^{(2)}(\mathbf{x}, t), \dots, h_{\mathbf{x}}^{(m)}(\mathbf{x}, t), \lambda_1, \lambda_2, \dots, \lambda_n) \\ \text{MSE}_{oR} &= \frac{1}{N_o} \sum_{k=1}^{N_o} (R(\mathbf{x}_k^o, t_k^o))^2, \quad (\mathbf{x}_k^o, t_k^o) \in \Omega \times [T_{n-1}, T_n] \end{aligned} \quad (\text{D.6})$$

- Mean squared error on the previous segments observed data

$$\text{MSE}_{os} = \frac{1}{N_{os}} \sum_{k=1}^{N_{os}} \left( \hat{h}(\mathbf{x}_k^{os}, t_k^{os}) - h_k^{os} \right)^2, \quad (\mathbf{x}_k^{os}, t_k^{os}) \in \Omega \times [0, T_{n-1}] \quad (\text{D.7})$$

592 where  $\hat{h}(\mathbf{x}_k^{os}, t_k^{os})$  is the neural network output and  $h_k^{os}$  is the observed data at  $(\mathbf{x}_k^{os}, t_k^{os})$ .  
593 Here, the superscript,  $(\bullet)^{os}$  stands for observed data in the previous time segments.

- Total mean squared error for inverse bc-PINN is given as

$$\text{MSE}_{\Delta T_n} = \text{MSE}_o + \text{MSE}_{oR} + \text{MSE}_{os} \quad (\text{D.8})$$

594 *Appendix D.1. Parameter identification of Allen Cahn equation 1D*

595 The Allen Cahn as described in earlier sections is a very widely used PDE in material science  
596 for studying the diffusion separation process. Therefore inverse identification of the Allen Cahn  
597 equation is essential to understand the governing physics of a process. Let us consider the explicit  
598 form of Allen Cahn equation

$$h_t = \lambda_1 \nabla^2 h + \lambda_2 (h^3 - h), \quad t \in (0, T], \quad x \in \Omega \subset \mathbb{R} \quad (\text{D.9})$$

599 Using the proposed framework of bc-PINN for inverse problems, the parameters learned are  
600 compared against std-PINN. To generate the required data set, random points have been sampled  
601 from the reference solution of Allen Cahn equation. The reference solution is computed as given  
602 in section (3.6). Also to test the effectiveness of both the methods noise has been added to  
603 the reference solution. Table (D.10) shows the parameter values obtained using std-PINN and  
604 bc-PINN.

	Predicted $\lambda_1$ (True value : 1e-04)		Predicted $\lambda_2$ (True value : 5)		Total relative error in the predicted solution	
Noise (%)	bc-PINN	PINN	bc-PINN	PINN	bc-PINN	PINN
0	1.504e-04	1.522e-04	5.03237	5.01081	0.0068	0.0062
2	1.375e-04	7.592e-05	4.98174	4.98214	0.0078	0.0070

Table D.10: True and Predicted parameters using bc-PINN and std-PINN for the 1D Allen Cahn equation.

605 *Appendix D.2. Parameter identification of Burgers equation with a time-varying parameter*

606 This section highlights the ability of the proposed bc-PINN framework for solving parametric  
 607 PDEs. Most of the real world systems are governed by parametric PDEs. The explicit parametric  
 608 form of burgers equation is as follows:

$$h_t = \lambda_1(t) h h_x + \lambda_2 h_{xx} \quad (x, t) \in [-8, 8] \times (0, 10] \quad (\text{D.10})$$

609 where,  $\lambda_1(t)$  is a time varying parameter and,  $\lambda_2$  is a constant value equal to 0.1. The reference  
 610 solution has been computed as given in [57].

$$\lambda_1(t) = \begin{cases} 1.00 & 0 \leq t < 2 \\ 0.75 & 2 \leq t < 4 \\ 0.50 & 4 \leq t < 6 \\ 0.75 & 6 \leq t < 8 \\ 1.00 & 8 \leq t < 10 \end{cases} \quad (\text{D.11})$$

611 Using the proposed framework of bc-PINN for inverse problems, the parameters of the parametric  
 612 burgers equation are learnt using an arbitrary length of time segment. In the present example  
 613 the total time domain has been discretized into 256 time steps. Therefore, in order to identify  
 614 the time varying parameter ( $\lambda_1$ ) an arbitrary time segment of 10 steps has been chosen. In  
 615 order to identify the parameter  $\lambda_1$ , all the observed data within the 10 steps segment has been  
 616 considered. A similar procedure has also been adopted for identifying the parameter  $\lambda_2$ .

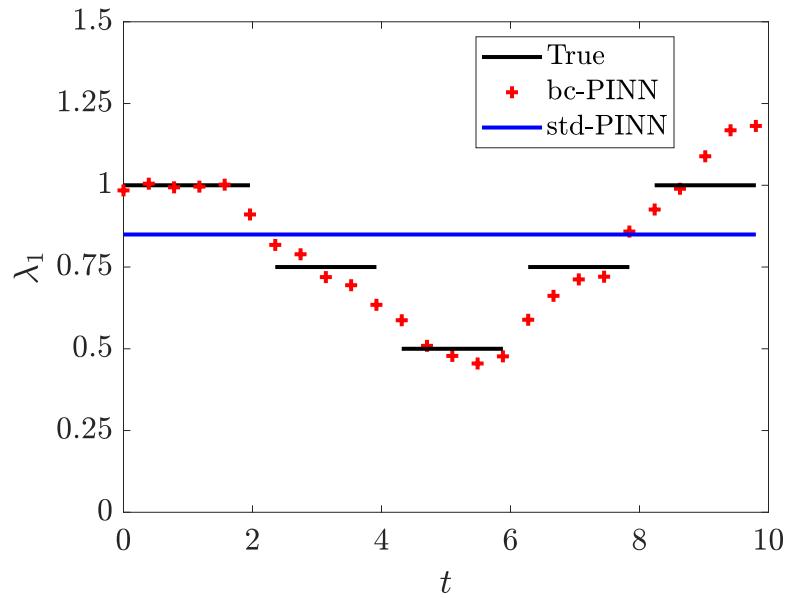


Figure D.22: ‘ $\lambda_1$ ’ learned using the bc-PINN and std-PINN approach for the parametric burgers equation

617 From figure (D.22|D.23), it can be observed that the predicted parameters follow a trend  
 618 similar to the true values. The main advantage of inverse bc-PINN scheme is that without any  
 619 prior knowledge, the nature of a PDE (constant or time-varying) can be identified. Moreover, the  
 620 values of the time varying parameters can be obtained without any prior information about the  
 621 time segments over which its constant.

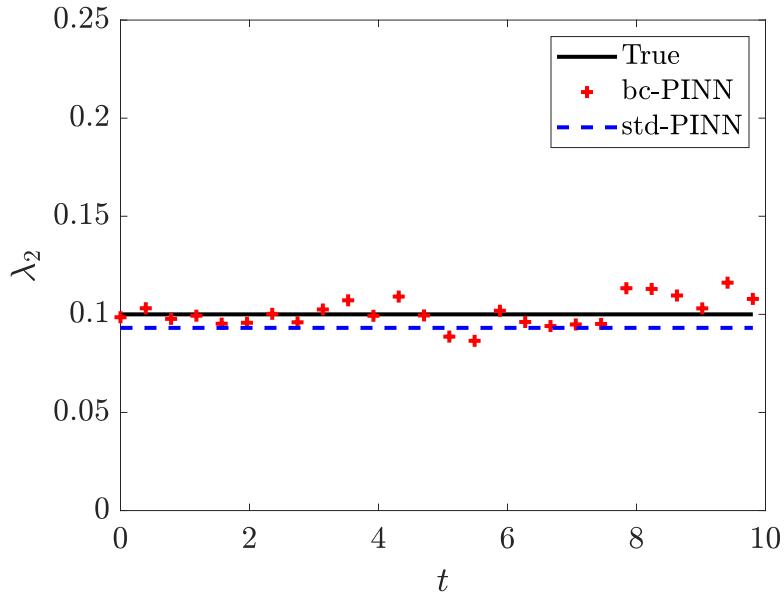


Figure D.23: ‘ $\lambda^2$ ’ learned using the bc-PINN and std-PINN approach for the parametric burgers equation

## 622 References

- 623 [1] Z. Mao, A. D. Jagtap, G. E. Karniadakis, Physics-informed neural networks for high-speed  
624 flows, Computer Methods in Applied Mechanics and Engineering 360 (2020) 112789.
- 625 [2] H. Arbabi, J. E. Bunder, G. Samaey, A. J. Roberts, I. G. Kevrekidis, Linking machine  
626 learning with multiscale numerics: Data-driven discovery of homogenized equations, JOM  
627 72 (12) (2020) 4444–4457.
- 628 [3] M. Raissi, P. Perdikaris, G. E. Karniadakis, Inferring solutions of differential equations using  
629 noisy multi-fidelity data, Journal of Computational Physics 335 (2017) 736–746.
- 630 [4] M. Raissi, P. Perdikaris, G. E. Karniadakis, Machine learning of linear differential equations  
631 using Gaussian processes, Journal of Computational Physics 348 (2017) 683–693.
- 632 [5] M. Raissi, G. E. Karniadakis, Hidden physics models: Machine learning of nonlinear  
633 partial differential equations, Journal of Computational Physics 357 (2018) 125–141. [arXiv:1708.00588](https://arxiv.org/abs/1708.00588).
- 635 [6] S. Atkinson, N. Zabaras, Structured bayesian gaussian process latent variable model:  
636 Applications to data-driven dimensionality reduction and high-dimensional inversion, Journal of  
637 Computational Physics 383 (2019) 166 – 195.
- 638 [7] I. Bilionis, N. Zabaras, B. A. Konomi, G. Lin, Multi-output separable gaussian process:  
639 Towards an efficient, fully bayesian paradigm for uncertainty quantification, Journal of  
640 Computational Physics 241 (2013) 212 – 239.
- 641 [8] M. Raissi, P. Perdikaris, G. E. Karniadakis, Physics-informed neural networks: A deep  
642 learning framework for solving forward and inverse problems involving nonlinear partial  
643 differential equations, Journal of Computational Physics 378 (2019) 686–707.

644 [9] A. G. Baydin, B. A. Pearlmutter, A. A. Radul, J. M. Siskind, Automatic differentiation in  
 645 machine learning: a survey, *Journal of machine learning research* 18 (2018).

646 [10] D. Zhang, L. Lu, L. Guo, G. E. Karniadakis, Quantifying total uncertainty in physics-  
 647 informed neural networks for solving forward and inverse stochastic problems, *Journal of*  
 648 *Computational Physics* 397 (2019) 108850.

649 [11] L. Yang, D. Zhang, G. E. Karniadakis, Physics-informed generative adversarial networks  
 650 for stochastic differential equations, *SIAM Journal on Scientific Computing* 42 (1) (2020)  
 651 A292–A317.

652 [12] X. Meng, G. E. Karniadakis, A composite neural network that learns from multi-fidelity data:  
 653 Application to function approximation and inverse pde problems, *Journal of Computational*  
 654 *Physics* 401 (2020) 109020.

655 [13] A. Jagtap, E. Kharazmi, G. Karniadakis, Conservative physics-informed neural networks  
 656 on discrete domains for conservation laws: Applications to forward and inverse problems,  
 657 *Computer Methods in Applied Mechanics and Engineering* 365 (2020) 113028.

658 [14] A. D. Jagtap, G. Em Karniadakis, Extended physics-informed neural networks (xpinnns): A  
 659 generalized space-time domain decomposition based deep learning framework for nonlinear  
 660 partial differential equations, *Communications in Computational Physics* 28 (5) (2020)  
 661 2002–2041.

662 [15] S. Karumuri, R. Tripathy, I. Bilionis, J. Panchal, Simulator-free solution of high-dimensional  
 663 stochastic elliptic partial differential equations using deep neural networks, *Journal of*  
 664 *Computational Physics* 404 (2020) 109120. [arXiv:1902.05200](https://arxiv.org/abs/1902.05200).

665 [16] R. K. Tripathy, I. Bilionis, Deep UQ: Learning deep neural network surrogate models for  
 666 high dimensional uncertainty quantification, *Journal of Computational Physics* 375 (2018)  
 667 565–588.

668 [17] Y. Yang, P. Perdikaris, Adversarial uncertainty quantification in physics-informed neural  
 669 networks, *Journal of Computational Physics* 394 (2019) 136–152.

670 [18] L. Yang, X. Meng, G. E. Karniadakis, B-pinns: Bayesian physics-informed neural networks  
 671 for forward and inverse pde problems with noisy data (2020). [arXiv:2003.06097](https://arxiv.org/abs/2003.06097).

672 [19] Q. He, D. Barajas-Solano, G. Tartakovsky, A. M. Tartakovsky, Physics-informed neural net-  
 673 works for multiphysics data assimilation with application to subsurface transport, *Advances*  
 674 *in Water Resources* 141 (2020) 103610.

675 [20] M. Liu, L. Liang, W. Sun, A generic physics-informed neural network-based constitutive  
 676 model for soft biological tissues, *Computer Methods in Applied Mechanics and Engineering*  
 677 372 (2020) 113402.

678 [21] G. Kissas, Y. Yang, E. Hwuang, W. R. Witschey, J. A. Detre, P. Perdikaris, Machine learning  
 679 in cardiovascular flows modeling: Predicting arterial blood pressure from non-invasive 4d flow  
 680 mri data using physics-informed neural networks, *Computer Methods in Applied Mechanics and*  
 681 *Engineering* 358 (2020) 112623.

682 [22] F. Sahli Costabal, Y. Yang, P. Perdikaris, D. E. Hurtado, E. Kuhl, Physics-informed neural  
 683 networks for cardiac activation mapping, *Frontiers in Physics* 8 (2020) 42.

684 [23] Y. Hu, T. Zhao, Z. Xu, L. Lin, Neural time-dependent partial differential equation (2020).  
685 [arXiv:2009.03892](https://arxiv.org/abs/2009.03892).

686 [24] Y. Khoo, J. Lu, L. Ying, Solving parametric pde problems with artificial neural networks,  
687 European Journal of Applied Mathematics 32 (3) (2020) 421–435.

688 [25] Z. Li, N. Kovachki, K. Azizzadenesheli, B. Liu, K. Bhattacharya, A. Stuart, A. Anandkumar,  
689 Neural operator: Graph kernel network for partial differential equations (2020). [arXiv:2003.03485](https://arxiv.org/abs/2003.03485)

690 [26] Z. Li, N. Kovachki, K. Azizzadenesheli, B. Liu, K. Bhattacharya, A. Stuart, A. Anandkumar,  
691 Fourier neural operator for parametric partial differential equations (2021). [arXiv:2010.08895](https://arxiv.org/abs/2010.08895).

692 [27] L. Lu, P. Jin, G. Pang, Z. Zhang, G. E. Karniadakis, Learning nonlinear operators via  
693 deeponet based on the universal approximation theorem of operators, Nature Machine  
694 Intelligence 3 (3) (2021) 218–229.

695 [28] H. Abels, H. Garcke, G. Grün, Thermodynamically consistent, frame indifferent diffuse  
696 interface models for incompressible two-phase flows with different densities, Mathematical  
697 Models and Methods in Applied Sciences 22 (03) (2012) 1150013.

698 [29] K. Deckelnick, G. Dziuk, C. M. Elliott, Computation of geometric partial differential  
699 equations and mean curvature flow, Acta Numerica 14 (2005) 139–232.

700 [30] J. Lowengrub, L. Truskinovsky, Quasi incompressible cahn hilliard fluids and topological  
701 transitions, Proceedings of the Royal Society of London. Series A: Mathematical, Physical  
702 and Engineering Sciences 454 (1998) 2617–2654.

703 [31] B. Li, J. Lowengrub, A. Ratz, A. Voigt, Geometric evolution laws for thin crystalline films:  
704 modeling and numerics, Communications in Computational Physics 6 (3) (2009) 433.

705 [32] L. N. Trefethen, N. Hale, T. A. Driscoll, Chebfun Guide, Pafnuty Publications, Oxford,  
706 2014.

707 [33] S. J. Pan, Q. Yang, A survey on transfer learning, IEEE Transactions on Knowledge and  
708 Data Engineering 22 (10) (2010) 1345–1359. [doi:10.1109/TKDE.2009.191](https://doi.org/10.1109/TKDE.2009.191).

709 [34] Y. Bengio, P. Simard, P. Frasconi, Learning long-term dependencies with gradient descent  
710 is difficult, IEEE Transactions on Neural Networks 5 (2) (1994) 157–166.

711 [35] Y. Lecun, L. Bottou, G. Orr, K.-R. Müller, Efficient backprop (08 2000).

712 [36] X. Glorot, Y. Bengio, Understanding the difficulty of training deep feedforward neural  
713 networks, in: Proceedings of the thirteenth international conference on artificial intelligence  
714 and statistics, JMLR Workshop and Conference Proceedings, 2010, pp. 249–256.

715 [37] D. P. Kingma, J. Ba, Adam: A method for stochastic optimization (2017). [arXiv:1412.6980](https://arxiv.org/abs/1412.6980).

716 [38] S. Cox, P. Matthews, Exponential time differencing for stiff systems, Journal of Computational Physics 176 (2) (2002) 430–455.

717 [39] M. Z. Bazant, Thermodynamic stability of driven open systems and control of phase  
718 separation by electro-autocatalysis, Faraday discussions 199 (2017) 423–463.

722 [40] S. M. Allen, J. W. Cahn, A microscopic theory for antiphase boundary motion and its  
723 application to antiphase domain coarsening, *Acta Metallurgica* 27 (6) (1979) 1085–1095.

724 [41] S. Bartels, Numerical methods for nonlinear partial differential equations, Vol. 47, Springer,  
725 2015.

726 [42] J. Shen, X. Yang, Numerical approximations of allen-cahn and cahn-hilliard equations,  
727 *Discrete & Continuous Dynamical Systems-A* 28 (4) (2010) 1669.

728 [43] J. W. Cahn, J. E. Hilliard, Free energy of a nonuniform system. i. interfacial free energy,  
729 *The Journal of chemical physics* 28 (2) (1958) 258–267.

730 [44] A. Miranville, The cahn-hilliard equation and some of its variants, *AIMS Mathematics* 2  
731 (2017) 479–544.

732 [45] J. Kim, S. Lee, Y. Choi, S.-M. Lee, D. Jeong, Basic principles and practical applications of  
733 the cahn–hilliard equation, *Mathematical Problems in Engineering* 2016 (2016).

734 [46] S. C. Takatori, J. F. Brady, Towards a thermodynamics of active matter, *Phys. Rev. E* 91  
735 (2015) 032117.

736 [47] T. Speck, J. Bialké, A. M. Menzel, H. Löwen, Effective cahn-hilliard equation for the phase  
737 separation of active brownian particles, *Phys. Rev. Lett.* 112 (2014) 218304.

738 [48] S. C. Takatori, W. Yan, J. F. Brady, Swim pressure: Stress generation in active matter,  
739 *Phys. Rev. Lett.* 113 (2014) 028103.

740 [49] A. A. Hyman, C. A. Weber, F. Jülicher, Liquid-liquid phase separation in biology, *Annual  
741 Review of Cell and Developmental Biology* 30 (1) (2014) 39–58.

742 [50] D. Zwicker, A. A. Hyman, F. Jülicher, Suppression of oswald ripening in active emulsions,  
743 *Physical Review E* 92 (1) (2015) 012317.

744 [51] C. P. Brangwynne, C. R. Eckmann, D. S. Courson, A. Rybarska, C. Hoege, J. Gharakhani,  
745 F. Jülicher, A. A. Hyman, Germline p granules are liquid droplets that localize by controlled  
746 dissolution/condensation, *Science* 324 (5935) (2009) 1729–1732.

747 [52] C. P. Brangwynne, P. Tompa, R. V. Pappu, Polymer physics of intracellular phase transitions,  
748 *Nature Physics* 11 (11) (2015) 899–904.

749 [53] B. Horstmann, T. Danner, W. G. Bessler, Precipitation in aqueous lithium–oxygen batteries:  
750 a model-based analysis, *Energy & Environmental Science* 6 (4) (2013) 1299–1314.

751 [54] J. Erlebacher, M. J. Aziz, A. Karma, N. Dimitrov, K. Sieradzki, Evolution of nanoporosity  
752 in dealloying, *Nature* 410 (6827) (2001) 450–453.

753 [55] W. Tian, X. Mao, P. Brown, G. C. Rutledge, T. A. Hatton, Electrochemically nanostructured  
754 polyvinylferrocene/polypyrrole hybrids with synergy for energy storage, *Advanced Functional  
755 Materials* 25 (30) (2015) 4803–4813.

756 [56] D. Lee, J.-Y. Huh, D. Jeong, J. Shin, A. Yun, J. Kim, Physical, mathematical, and  
757 numerical derivations of the cahn–hilliard equation, *Computational Materials Science* 81  
758 (2014) 216–225.

759 [57] S. Rudy, A. Alla, S. L. Brunton, J. N. Kutz, Data-driven identification of parametric partial  
760 differential equations (2018). [arXiv:1806.00732](https://arxiv.org/abs/1806.00732).

761 [58] V. Dwivedi, N. Parashar, B. Srinivasan, Distributed learning machines for solving forward  
762 and inverse problems in partial differential equations, Neurocomputing 420 (2021) 299–316.

763 [59] M. Cheng, T. Y. Hou, M. Yan, Z. Zhang, A data-driven stochastic method for elliptic pdes  
764 with random coefficients, SIAM/ASA Journal on Uncertainty Quantification 1 (1) (2013)  
765 452–493.