

## Experimental investigation of high-spin states in $^{90}\text{Zr}$

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High-spin states of  $^{90}\text{Zr}$  have been investigated using the heavy-ion fusion-evaporation reaction  $^{82}\text{Se}(^{13}\text{C}, 5n)$  at a beam energy of 60 MeV. Excited levels of  $^{90}\text{Zr}$  have been observed up to an excitation energy of  $\approx 13$  MeV and a spin of  $\approx 20\hbar$  with the addition of thirty-two new  $\gamma$ -ray transitions to the proposed level scheme. Structures of both the positive- and negative-parity states up to the highest observed spin have been interpreted with shell-model calculations using the GWBXXG interaction and a  $^{68}\text{Ni}$  core. Calculations suggest the role of neutron excitations across the  $N = 50$  shell gap for states with greater than 7 MeV excitation energy. High-spin states in these bands are interpreted to be generated by the recoupling of stretched proton and neutron configurations.

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### I. INTRODUCTION

Nuclei in the vicinity of shell closures are of particular interest in nuclear structure studies because they provide a platform for scrutinizing the validity and details of shell-model theories. It is now computationally feasible to extend shell-model calculations to higher excitations that incur a larger model space and an increased number of valence nucleons, to be used for the purpose. High-spin states in the  $A \approx 90$  mass region, with  $Z \approx 40$  and  $N \approx 50$ , have been subjects of investigation in many studies, both experimentally and theoretically [1–31]. These states have multiquasiparticle configurations, with the  $g_{9/2}$  orbital playing a significant role towards the generation of high-spin states. The contribution of proton  $g_{9/2}$  orbital comes into the picture either due to the (proton) particle occupancy of the orbital, which is for

nuclei with  $Z > 40$ , or due to excitations of protons from the  $fp$  orbitals into the  $g_{9/2}$  orbital across the  $Z = 40$  subshell gap. These excitations dominate the lower-energy part of the level scheme in these nuclei. Similarly, the role of neutron  $g_{9/2}$  orbital towards the high-spin generation is seen to be either due to the already-present holes in it, which is the case for nuclei with  $N < 50$ , or to excitations across the  $N = 50$  shell gap into the  $gd$  orbitals. Due to the larger energy of this gap, the contributions from these excitations underlying the high-spin states in  $N = 50$  nuclei are observed at higher excitation energies. Such excitations of one or more neutrons from the  $1g_{9/2}$  orbital into the  $2d_{5/2}$  and  $1g_{7/2}$  have been observed in the  $^{86}\text{Kr}$  ( $Z = 36$ ) [2],  $^{87}\text{Rb}$  ( $Z = 37$ ) [4],  $^{88}\text{Sr}$  ( $Z = 38$ ) [6],  $^{89}\text{Y}$  ( $Z = 39$ ) [11,12],  $^{91}\text{Nb}$  ( $Z = 41$ ) [21],  $^{92}\text{Mo}$  ( $Z = 42$ ) [23],  $^{93}\text{Tc}$  ( $Z = 43$ ) [27],  $^{94}\text{Ru}$  ( $Z = 44$ ) [29], and  $^{95}\text{Rh}$  ( $Z = 45$ ) [31] isotones, but it has yet to be confirmed in  $^{90}\text{Zr}$  ( $Z = 40$ ) [15]. The shell-model calculations have been used to describe observed excited states based on either proton excitation across  $Z = 40$  or neutron excitation across  $N = 50$ . It will be intriguing to invoke the two excitations simultaneously for  $^{90}\text{Zr}$ , which is the focus of the present study.

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