Al Composition Dependence of Band Offsets for SiO₂ on α-(Al_xGa_{1-x})₂O₃

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Abstract

Valence band offsets were measured by X Ray Photoelectron Spectroscopy for SiO₂ deposited by Atomic Layer Deposition on α -(Al_xGa_{1-x})₂O₃ alloys with x= 0.26-0.74 grown with a continuous composition spread to enable investigations of the band alignment as a function of the alloy composition. From measurement of the core levels in the alloys, the bandgaps were determined to range from 5.8 eV (x=0.26) to 7eV (x=0.74). The valence band offsets were -1.2 eV for x=0.26, -0.2 eV for x=0.42, 0.2 eV for x=0.58 and 0.4 eV for x=0.74. Given the bandgap of the SiO₂ was 8.7 eV, this led to conduction band offsets of 4.1 eV (x=0.26) to 1.3 eV (x=0.74). The band alignments were nested for x>0.5, but at lower Al contents the conductions band offsets were negative, with a staggered band alignment. This shows the challenge of finding appropriate dielectrics for this ultra-wide bandgap semiconductor system.

Introduction

 $(Al_xGa_{1-x})_2O_3$ is an ultra-wide bandgap semiconductor that is attracting attention for use in high power transistors and deep ultraviolet (DUV) photodetectors ⁽¹⁻⁴⁾. To this point, most of the work has focused on the β -polytype with a bandgap tunability ranging from 4.8 eV (Ga₂O₃) to 8.7 eV (Al₂O₃) ^(1,2). The advantages of β -(Al_xGa_{1-x})₂O₃ grown on β -Ga₂O₃ include the high channel mobility at the interface and ability to use the heterostructure in field effect transistor structures with improved mobility, as well as Schottky diodes and photodetectors ^(1,2). A drawback of this materials system during growth with techniques such as Metal Organic Chemical Vapor Deposition is the difficulty in maintain phase stability at high Al contents ⁽¹⁾. The α -polytype is a potential solution to extending the stability to higher Al contents and higher temperatures $^{(2-12)}$. This metastable, corundum crystal structure α -polytype, transforms to the β polymorph at temperatures above 750-900°C under equilibrium conditions ⁽⁵⁾. Of all reported polymorphs, this phase has the highest optical bandgap with energies ranging between 5.0 eV and 5.3 eV for the binary. A number of different growth methods have recently been shown to produce high quality rhombohedral α -Al_xGa_{1-x})₂O₃ over a broad composition range ⁽⁶⁻⁴⁸⁾ and in particular, the use of Al₂O₃ cap layers during annealing is able to preserve the α -phase ⁽¹⁶⁾. The α -polytype has an even larger bandgap than the β -polytype and can be grown on low-cost c-plane sapphire substrates which have the same lattice structure, albeit with relatively high dislocation densities $^{(2, 19, 22, 43)}$. Figure 1 shows the crystal structure of rhombohedral α -Ga₂O₃. This rhombohedral phase can also occur for Al₂O₃ and In₂O₃, which allows a potential bandgap engineering over the entire composition range ^(34,45,46), rendering this polymorph interesting for deep UV photodetectors or quantum-well infrared photodetectors ^(1,2), in addition to the power device applications ⁽³⁵⁻³⁷⁾.

It is clearly of interest to establish the band alignments for commonly employed dielectrics on this materials system. SiO₂ is one of the most promising choices, due to its large bandgap and maturity as a dielectric in semiconductor devices. We have previously measured the band alignment and its thermal stability for SiO₂ and Al₂O on β - Al_xGa_{1-x})₂O₃ for 0.2< x <0.65 ^(51,52). In this paper we report on the measurement of valence band offsets in the SiO₂/ α -Al_xGa_{1-x})₂O₃ system using alloys grown by Pulsed Laser Deposition (PLD)^(49,50) over a wide composition range. This was done using circular half-segmented (Ga₂O₃/Al₂O₃) ceramic targets to synthesize large diameter large thin films with a lateral varying cation composition ^(1,7,8,45,46). Since the bandgap of the α -polytype alloys is even larger than those of the β -polytype alloys, this system represents a test in terms of finding dielectrics that exhibit acceptable band offsets (> 1eV is a common rule of thumb) in both the conduction and valence bands.

Experimental

The α -Al_xGa_{1-x})₂O₃ alloys were grown by combinatorial PLD using a KrF excimer laser (Coherent LPX Pro 305, 248 nm, 2.6 J.cm⁻²) directed through a quartz window onto a segmented ceramic target (Ga₂O₃/Al₂O₃) under an O₂ partial pressure ^(1,49,50). The substrate and target were mounted rotatable opposite each other at a distance of 10 cm and with a small lateral offset to each other ⁽⁴⁶⁾. PLD has been utilized previously to obtain AGO alloys over a broad composition range for both α and β -polytypes ^(8, 20, 40) The resulting cation ratio was determined by energy dispersive X-ray spectroscopy (EDX) measurements performed using a FEI Nova Nanolab 200 equipped with an Ametek EDAX detector and the crystal structure by X-ray diffraction (XRD). measured by a PANalytical X'pert PRO MRD diffractometer equipped with a PIXcel3D detector operating in 1D scanning line mode with 255 channels (for 2Θ- ω scans), receiving slit mode (for Φ -scans) and fast 2D frame based mode (for reciprocal space maps). The cation gradient ranged between x= 0.13 and x=0.84 as measured by EDX and marked as black dots in Figure 2(a). The composition between these data points was interpolated and the resulting cation ratio represented as false color map. Along the cation gradient direction, marked with a black arrow, additional EDX measurements were performed in 1mm steps to obtain the spatial Al dependence with high lateral resolution as shown in Figure 2(b) ⁽²⁰⁾. The Al incorporation covers a range of $0.13 \le x \le 0.84$. The rhombohedral crystal structure of the thin films was confirmed by the 2Θ - ω XRD scans. The resulting false color map is presented in Figure 3 and reveals the crystallization in the rhombohedral phase in the entire investigated Al range. We examined four compositions, x=0.26, 0.42, 0.58 and 0.74. More details and discussion on the characterization of these samples can be found in Hassa et al. ⁽²⁰⁾.

The SiO₂ layers were deposited by Atomic Layer Deposition (ALD) at 200°C using the thermal mode in a Cambridge Nano Fiji 200. The precursors were Tris (dimethylamino) silane and an inductively coupled plasma (ICP) of O₂ at 300 W ^(51,52). After solvent cleaning, the substrates were exposed to Ozone for 15 min to form a protective oxide that is easily thermally desorbed during the ALD step. Both thick (200 nm) and thin (1.5 nm) layers of SiO₂ were deposited for measuring bandgaps and core levels on the α -(Al_xGa_{1-x})₂O₃.

XPS survey scans were performed to measure the chemical state of the SiO₂ and (Al_xGa₁. _x)₂O₃ layers in a Physical Instruments ULVAC PHI, with an Al x-ray source (energy 1486.6 eV, source power 300W)⁽⁵¹⁻⁵³⁾. The analysis size was 100 μ m diameter, with a take-off angle of 50° and acceptance angle of ±7 degrees. The electron pass energy was 23.5 eV for high-resolution scans and 93.5 eV for survey scans. The total energy resolution of this XPS system is 0.5 eV, and the accuracy of the observed binding energy is ~ 0.03 eV. Charge compensation was performed using an electron flood gun and simultaneous ion beam. The C 1*s* core levels of the surface adsorbate (284.8 eV) were used to calibrate the binding energy ⁽⁵³⁾. Only the relative energy position is needed to determine the valence band offsets, so the absolute energy calibration has no effect on that offset. The samples were electrically insulated from the chuck to avoid uneven charge dispersion along the sample. All electron analyzers and equipment were grounded. Differential charging was not observed in any of the samples with the use of the electron gun ^(51,52).

The SiO₂ bandgap was obtained from Reflection Electron Energy Loss Spectroscopy (REELS) using a 1 kV electron beam and hemispherical electron analyzer. The bandgaps of the $(Al_xGa_{1-x})_2O_3$ for each composition were obtained from XPS energy loss measurements of the O1S peak. This is done simultaneously with the band alignment measurements ^(52,53).

Results and Discussion

Figure 4 shows the high resolution XPS spectra for the α - (Al_xGa_{1-x})₂O₃ to SiO₂ vacuumcore delta regions of the four compositions, while the data for the SiO₂ is shown in Figure 5. The valence band offsets were extracted from the shift of the core levels for the heterostructure samples with the thin dielectric on top of the four different compositions of the alloy. A compilation of the valence band maxima (VBM) are collected in Table 1, with values of (2.7±0.15) eV for x=0.26, (3.3±0.15) eV for x=0.42, (4.4±0.15)eV for x=0.58 and (4.5±0.15) eV for x=0.74. These were obtained by fitting of the leading edge of the valence band. The error bars in the different binding energies were combined in a root sum square relationship to determine the overall error bars in the valence band offsets (ΔE_V) ⁽⁵¹⁾. These band offsets are then obtained by measuring the shift of the core levels in the α - (Al_xGa_{1-x})₂O₃ when SiO₂ was deposited. These are also tabulated in Table 1, with values of -1.2 eV for x=0.26, -0.2 eV for x=0.42, 0.2 eV for x=0.58 and 0.4 eV for x=0.74. Once the valence band offsets are established, to determine the conduction band offset, it is necessary to measure the band gap of each composition. We measured the bandgaps of the four $(Al_xGa_{1-x})_2O_3$ compositions, as shown in Figure 6 (a), from the separation between the core level peak energy and the onset of inelastic (plasmon) losses in each O 1s photoemission spectra ⁽³⁶⁾. To find the band-gap energy, a linear fit is made to the measured loss spectra curve near the approximate location of onset of inelastic losses. Next, by subtracting the background, the "zero" level is determined. The energy corresponding to the onset of inelastic losses is found by extrapolating the linear-fit line and calculating its intersection with the "zero" level ^(33,40). The bandgap energy is equal to the difference between the core-level peak energy and the onset of inelastic losses.

The bandgap energies were determined as 5.8 eV for x=0.26, 6.1 eV for x=0.42, 6.4 eV for x=0.58 and 7 eV for x= 0.74, respectively. Figure 7 shows the relationship between bandgap of the alloys and Al composition for our samples. As previously reported, the bandgap as a function of composition, x, is given by (28,54)

$$E_G(x) = (1 - x)E_G^{GaO} + xE_G^{AlO} - bx (1 - x)$$

where E_G^{GaO} and E_G^{AIO} are the bandgaps of the α -polytype binary endpoints and b is the bowing parameter. This has been previously reported in the range 0.84-2.13 ^(9, 34). We get a value closer to 3 for our samples in the composition range measured. It was noted by Varley ⁽⁵³⁾ that the *b* values are strongly influenced by the selective occupation of different cation sublattice by Al and that a single bowing parameter is not accurate over the entire composition range

The SiO₂ bandgap was obtained from Reflection Electron Energy Loss Spectroscopy (REELS) and showed a value of 8.7 eV, similar to previous reports $^{(51,52)}$.

The conduction band offsets ΔE_C were obtained from the relation

$$\Delta E_C = E_G^{SiO2} - E_G^{GaO} - \Delta E_V$$

The corresponding conduction band offsets were 4.1 eV (x=0.26), 2.8 eV (x=0.46), 2.1 eV (x=0.58) and 1.30 eV (x=0.74).

Figure 7 shows the band diagrams for the SiO₂/ α (Al_xGa_{1-x})₂O₃ heterostructure. The SiO₂ has adequate conduction band offsets, but the valence band offsets are small and actually negative at x<0.5. These would be inadequate confinement of holes in SiO₂/ α - (Al_xGa_{1-x})₂O₃ samples over the entire composition range. The conduction band offsets are high enough to provide effective electron confinement. The SiO₂/ (Al_xGa_{1-x})₂O₃ band alignment is type I for x>0.5 and staggered for x<0.5. It will be interesting to measure the band alignments for Al₂O₃ deposited on the α -(Al_xGa_{1-x})₂O₃ alloys, since that is one of the few dielectrics with a high enough bandgap to provide effective carrier confinement for electrons and holes in these alloys.

Summary and Conclusions

The valence band offsets of SiO₂/ α -(Al_xGa_{1-x})₂O₃ heterojunctions were measured over a range of Al contents (x=0.26 - 0.74). The band alignments are staggered type II for x<0.5 and straddling type I for x>0.5, with conduction band offsets >1.3 eV across the composition range examined. The thermal stability of this system is of interest for future work, as the α -polytype I not the most energetically favorable state, although kinetically limited process such as rapid thermal annealing will be able to be used in device processing steps without major issues.

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Table 1. Summary of the measured reference and heterostructure peaks for SiO₂ on (Al_xGa₁₋

x)2O3 (eV)

	Reference (Al _x Ga _{1-x}) ₂ O ₃			Reference SiO ₂			Thin SiO ₂ on $(Al_xGa_{1-x})_2O_3$	
Aluminum Concentration	Core Level Peak (Ga 2p _{3/2})	VBM	Core - VBM	Core Level Peak (Si 2p)	VBM	Core - VBM	Δ Core Level (Ga 2p _{3/2} - Si 2p)	Valence Band Offset
$(Al_{0.26}Ga_{0.74})_2O_3$	1117.7	2.7±0.15	1115	103.40	4.80	98.60	1017.6	-1.2
$(Al_{0.42}Ga_{0.58})_2O_3$	1118.2	3.3±0.15	1114.9	-	-	-	1016.5	-0.2
(Al _{0.58} Ga _{0.42}) ₂ O ₃	1119.2	4.4±0.15	1114.7	-	-	-	1015.9	0.2
(Al _{0.74} Ga _{0.26}) ₂ O ₃	1118.9	4.5±0.15	1114.5	-	-	-	1015.5	0.4

Figure Captions

Figure 1. Crystal structure of rhombohedral α -Ga₂O₃. The oxygen atoms are displayed as red and the Ga atoms as green spheres. The graphic was created with the software Crystal Maker X, http://www.crystalmaker.com/about/index.html .

Figure 2. (a) Al content x at 49 points across the thin film surface, marked on the map as black dots and determined by EDX. The data between the measurement points was interpolated, the black arrow represents the direction of the gradient. (b) Al ratio x acquired along the indicated gradient in (a) by EDX and XPS, respectively.

Figure 3. False color map of θ - ω scans of α -(Al_xGa_{1-x})₂O₃ recorded along the composition gradient indicated in Figure 2(b).

Figure 4. XPS spectra of core levels to valence band maximum (VBM) for reference (Al_xGa_{1-x})₂O₃ with x= 0.26, 0.42, 0.58 or 0.74 (top to bottom) Aluminum, and (c) ALD thick film Al_2O_3 and SiO₂. The intensity is in arbitrary units (a.u.).

Figure 5. XPS spectra of core levels to valence band maximum (VBM) for ALD thick film SiO₂.The intensity is in arbitrary units (a.u.).

Figure 6. Bandgap of α -(Al_xGa_{1-x})₂O₃ determined using the onset of the plasmon loss feature in O 1s photoemission spectrum. The intensities are in arbitrary units (a.u.).

Figure 7. Bandgap of α -(Al_xGa_{1-x})₂O₃ for 4 Al compositions (experimental points) fitted to a bowing parameter of 3eV, along with a theoretical curve with bowing parameter 1.37 eV.

Figure 8. Band diagrams for the SiO₂/ α -(Al_xGa_{1-x})₂O₃ heterostructure in which the SiO₂ was deposited by ALD.















