

Ultrahigh-rectification near-field radiative thermal diode using infrared-transparent film backsided phase-transition metasurface

Cite as: Appl. Phys. Lett. **119**, 123101 (2021); <https://doi.org/10.1063/5.0058779>

Submitted: 02 June 2021 • Accepted: 04 September 2021 • Published Online: 21 September 2021

Yang Liu, Yanpei Tian, Fangqi Chen, et al.



View Online



Export Citation



CrossMark

ARTICLES YOU MAY BE INTERESTED IN

[Near-field photonic thermal diode based on hBN and InSb films](#)

Applied Physics Letters **119**, 181111 (2021); <https://doi.org/10.1063/5.0068775>

[Dynamic tuning of near-field radiative transport between an overlapping pair of movable comb-like metamaterials](#)

Applied Physics Letters **118**, 153106 (2021); <https://doi.org/10.1063/5.0045563>

[Tailoring broadband Kerr soliton microcombs via post-fabrication tuning of the geometric dispersion](#)

Applied Physics Letters **119**, 121103 (2021); <https://doi.org/10.1063/5.0061238>

Lock-in Amplifiers up to 600 MHz



Zurich
Instruments



Ultrahigh-rectification near-field radiative thermal diode using infrared-transparent film backside phase-transition metasurface

Cite as: Appl. Phys. Lett. **119**, 123101 (2021); doi: [10.1063/5.0058779](https://doi.org/10.1063/5.0058779)

Submitted: 2 June 2021 · Accepted: 4 September 2021 ·

Published Online: 21 September 2021



View Online



Export Citation



CrossMark

Yang Liu,¹ Yanpei Tian,¹ Fangqi Chen,¹ Andrew Caratenuto,¹ Xiaojie Liu,¹ Mauro Antezza,^{2,3}  and Yi Zheng^{1,a)} 

AFFILIATIONS

¹Department of Mechanical and Industrial Engineering, Northeastern University, Boston, Massachusetts 02115, USA

²Laboratoire Charles Coulomb (L2C), UMR 5221 CNRS-Université de Montpellier, F-34095 Montpellier, France

³Institut Universitaire de France, 1 rue Descartes, F-75231 Paris Cedex 05, France

^{a)}Author to whom correspondence should be addressed: y.zheng@northeastern.edu

ABSTRACT

We present a theoretical study of near-field radiative thermal rectification combining phase-transition and high-infrared-transmittance materials. The phase-transition material vanadium dioxide (VO₂), with a metal-insulator transition near 341 K, is utilized under a reasonable temperature. Four types of high-infrared-transmittance materials, including potassium bromide, sodium chloride, polyethylene, and magnesium fluoride, are introduced as thin film substrates under a VO₂ grating on one side of the near-field rectifier. We explore the effects of various high-infrared-transmittance thin-film substrates and relevant geometric parameters on the thermal rectification of the device. The results show that thermal rectification can be greatly enhanced by using a one-dimensional VO₂ grating backed with a high-infrared-transmittance thin-film substrate. With the introduction of a high-infrared-transmittance substrate, the rectification ratio is dramatically boosted due to the enhancement of the substrate transmittance. This work predicts a remarkable rectification ratio as high as 161—greater than the recently reported peak values for comparable near-field radiative thermal rectification. The results outlined herein will shed light on the rapidly expanding fields of nanoscale thermal harvesting, conversion, and management.

Published under an exclusive license by AIP Publishing. <https://doi.org/10.1063/5.0058779>

The study of near-field radiative thermal rectification, which can preferentially facilitate radiative heat transfer in one direction through a micro/nanoscale gap, has recently garnered significant attention.^{1–7} Due to their contactless structures, radiative thermal rectification devices do not suffer the same limitations as conduction-based thermal rectifiers, which are restrained by both the speed of phonons and the presence of Kapitza resistances.^{8–14} Meanwhile, the near-field radiative heat transfer (NFRHT) between closely spaced objects can exceed the blackbody radiation limit by several orders of magnitude due to the tunneling of evanescent waves and coupling of surface phonons or plasmon polaritons, which provides near-field radiative thermal rectifiers with excellent thermal rectification performance.^{15–29} By taking advantage of the near-field radiative thermal rectification, devices like thermal diodes,⁴ transistors,⁹ switches,³⁰ and thermal memories³¹ can be leveraged to attain highly efficient functionality in nanoscale heat control and thermal modulation.

To gauge the rectification efficiency of a thermal diode, we utilize the definition of the rectification ratio $R = (Q_F - Q_R)/Q_R$, where Q_F

and Q_R refer to the forward and reverse heat fluxes, respectively.³² Otey *et al.* proposed photon-mediated thermal rectification through vacuum, calculated near-field radiative heat transfer between SiC-3C and SiC-6H bulks, and predicted a maximum thermal rectification ratio of 0.41.³ Basu and Francoeur designed a near-field thermal rectifier using a doped silicon film and a doped silicon bulk and investigated the effects of the vacuum gap and the film thickness on rectification, realizing a rectification greater than 0.5 as a result.¹ Yang *et al.* theoretically calculated that a thermal rectification ratio of almost 2 could be achieved utilizing bulk vanadium dioxide (VO₂) and silicon dioxide (SiO₂) separated by a 10 nm gap. The same work postulated that a rectification ratio of 3 can be attained with a 100 nm gap when bulk SiO₂ is replaced by a thin film.³³ The enhancing effects of the VO₂ thin film and the SiO₂ substrate on a near-field vacuum thermal switch have also been studied.³⁴ Huang *et al.* designed a near-field thermal rectification structure consisting of VO₂ and La_{0.7}Ca_{0.15}Sr_{0.15}MnO₃ (LCSMO), which can achieve a rectification ratio of 7.7 with a vacuum gap of 10 nm.³⁵ Ghanekar *et al.* theoretically

demonstrated that the thermal rectification ratio could be greatly enhanced to 16 in such cases by using one-dimensional (1D) rectangular and triangular VO₂ surface gratings.⁴ Recently, Chen *et al.* studied an active-tuning near-field thermal rectifier using a stretchable structure composed of a polydimethylsiloxane (PDMS) thin film and a VO₂ grating, which can achieve an ultrahigh rectification ratio of 23.7.⁵ Moreover, extensive research has achieved and further enhanced thermal rectification by utilizing various materials (both polar and phase-transition) and structures (plate-plate structures, micro/nano-structures, and surface gratings).^{9,36–47} Recently, Li *et al.* theoretically achieved rectification ratios over 140 with nanofilms of vanadium dioxide and cubic boron nitride (cBN) in the parallel plane geometry at the gap of 100 nm by leveraging the distinct scaling behaviors of the local density of states with film thicknesses for metals and insulators.⁴⁸ Though many studies on the near-field radiative thermal rectification aim to further enhance the thermal rectification performance, the thermal rectification ratio has not been greatly improved.

In this Letter, we investigate the thermal rectification in a near-field radiative thermal diode utilizing the phase-transition material VO₂. To increase the thermal rectification ratio, different high-infrared-transmittance materials are employed, such as potassium bromide (KBr), sodium chloride (NaCl), polyethylene (PE), and magnesium fluoride (MgF₂) as a thin-film substrate under the VO₂ grating. Differing from previous studies that focus on the structure of the phase-transition material VO₂, this study concentrates on the materials and structures of the substrate film. Our work indicates that the rectification ratio can be raised significantly by using a high-infrared-transmittance thin-film substrate (HITTS).

Here, we propose a near-field radiative thermal diode composed of a VO₂ 1D grating and a HITTS. A typical design is shown in Fig. 1, incorporating two planar structures separated by a nanoscale distance less than the thermal wavelength ($1.8 \mu\text{m} \leq \lambda \leq 14 \mu\text{m}$). The passive

side consists of a 1 μm cubic boron nitride (cBN) layer on top of a 1 μm gold (Au) layer, which is kept constant throughout these analyses. The active counterpart is composed of a VO₂ grating deposited on top of the HITTS. Four high-infrared-transmittance materials are considered within this work: KBr, NaCl, PE, and MgF₂.

To obtain radiative heat fluxes across the near-field thermal diode, we utilize the expression for NFRHT based on the dyadic Green's function formalism.^{49–52} The NFRHT expression is as follows:

$$Q_{1 \rightarrow 2}(T_1, T_2, L) = \int_0^\infty \frac{d\omega}{2\pi} [\Theta(\omega, T_1) - \Theta(\omega, T_2)] \int_0^\infty \frac{k_\rho dk_\rho}{2\pi} \xi(\omega, k_\rho), \quad (1)$$

where $\Theta(\omega, T) = (\omega/2) \coth(\omega/2k_B T)$ is the energy of the harmonic oscillator. The function $\int_0^\infty \frac{k_\rho dk_\rho}{2\pi} \xi(\omega, k_\rho)$ is the spectral transmissivity of the radiative transport between media 1 and 2 separated by a distance L , where $\xi(\omega, k_\rho)$ is the energy transmission coefficient. Additionally, we use the effective medium approximation to obtain the effective dielectric properties of our proposed 1D grating structure.^{53–57}

To evaluate the effects of the HITTS on the thermal rectification performance of the diode described in Fig. 1, the heat flux is calculated based on the temperature difference between the active and passive sides of the diode ($2\Delta T$). Heat flux as a function of $2\Delta T$ is plotted for the four high-infrared-transmittance material cases studied in Fig. 2. To evaluate the specific effects of the HITTS on performance, our study calculates the heat flux for a PDMS film-backed diode as a comparison. This provides a useful juxtaposition as PDMS has lower infrared transmittance than the other materials considered and has been recently studied in similar works.⁵ The temperature of the active side is $T_1 = 341 \text{ K} + \Delta T$, while that of the passive side is set as $T_2 = 341 \text{ K} - \Delta T$. This is chosen based on the metal-insulator-transition temperature of VO₂,

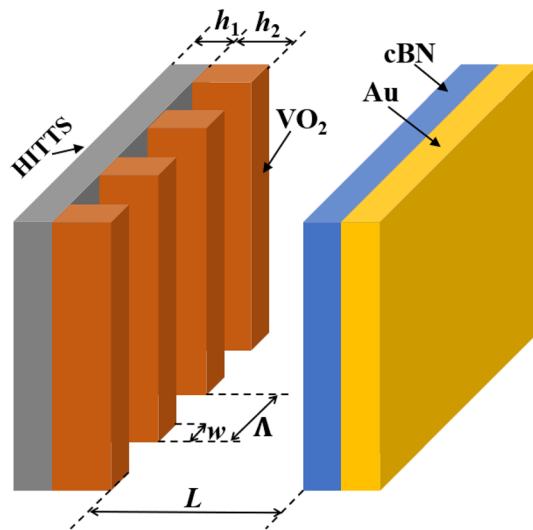


FIG. 1. Schematic of near-field thermal diodes using an infrared transparent film backside phase-transition metasurface. The active side (left) top layer is a 1D rectangular grating made of VO₂ with height h_2 , width w , period Λ , and filling ratio $\phi = w/\Lambda$ and is deposited on a HITTS with thickness h_1 . The passive counterpart (right) consists of a 1 μm cBN layer atop a 1 μm Au layer.

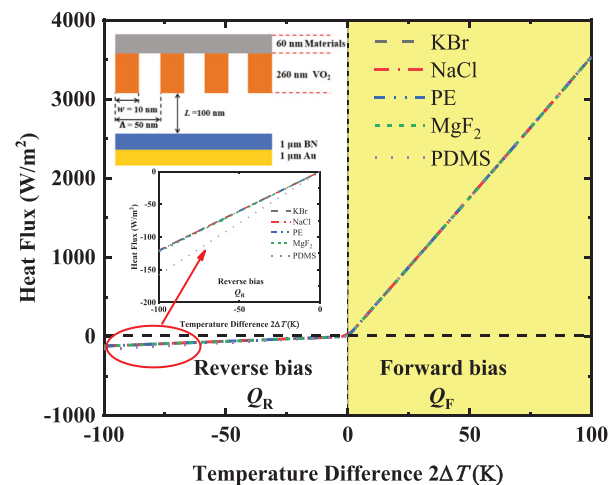


FIG. 2. Forward and reverse heat fluxes (Q_F and Q_R) as functions of the temperature difference between active and passive sides with a 100 nm gap for five different substrate materials, all with $h_1 = 60 \text{ nm}$, $h_2 = 260 \text{ nm}$, $\Lambda = 50 \text{ nm}$, $\phi = 0.2$, and $L = 100 \text{ nm}$.

which exists near 341 K. As such, when $T_1 > T_2$ (forward bias), VO_2 is in the metallic phase; when $T_1 < T_2$ (reverse bias), the VO_2 layer is in the insulator phase.³³ It can be clearly seen that the slopes of Q_F are much larger than those of Q_R , indicating apparent diode-like characteristics. To illustrate why there is a strong asymmetry in the heat transport across the interfaces of the thermal diode, we plot the energy transmission coefficient $\xi(\omega, k_p)$ across the interfaces of the thermal diode with the KBr substrate at a gap of 100 nm shown in Fig. 3. Here, $k_p c/\omega$ is the normalized parallel wavevector. Transmission is high at two prominent frequencies that occur near the characteristic wavelengths of cBN (7.6 and 9.8 μm). Since metallic VO_2 does not support surface phonon polariton in the infrared wavelength region, the NFRHT is mainly due to the symmetric and antisymmetric surface phonon polaritons supported by the cBN layer. Aside from that, there is a weak contribution based on the Fabry-Pérot modes and frustrated modes ($k_p c/\omega \approx 1$) in the near-field regime. High energy transmission is attributed to the tunneling of surface waves across interfaces when the diode is in the forward bias. However, the NFRHT is dominated by non-resonant surface waves when the diode is in the reverse bias because cBN and insulator VO_2 support the non-overlapping surface phonon modes. Meanwhile, the surface phonons of insulator VO_2 appear in the frequency range, where the extinction coefficient of cBN is very low ($\kappa \approx 0$). Therefore, tunneling between cBN and insulator VO_2 is much weaker than that of cBN and metallic VO_2 , which contributes to the strong asymmetry in the heat transport across the interfaces of the thermal diode. The diodes backed with high-infrared-transmittance materials (KBr, NaCl, PE, or MgF_2) on the active side realize improved thermal rectification performance with respect to the PDMS-backed diode. This is because HITTS cannot absorb more infrared spectrum radiation emitted from the passive side when the diode is in the reverse bias shown in the inset of Fig. 2. These data prove that the transmission properties of the thin film substrate play an important role in the radiative heat transfer across the diode.

In order to further illustrate the effects of different high-infrared-transmittance materials as the substrate on the thermal rectification performance of the diode, we analyze the impact all the related geometric parameters on the rectification ratio of the near-field radiative thermal diode, including the thickness of the substrate film h_1 , the height of the VO_2 grating h_2 , the filling ratio ϕ , and the gap L . The rectification ratios are calculated when the temperature of the active side is set as 351 K for the forward bias and 331 K for the reverse bias, and that of the passive side is kept constant at 341 K. Figure 4(a) shows the rectification ratio R vs the filling ratio of the VO_2 grating for different HITTSs. It is evident that when the value of the filling ratio ϕ decreases from 1.0 to 0.1, these rectification ratios trend upwards with varying slopes. A sharp decrease in the rectification ratio is noted when the value of filling ratio ϕ increases from 0.1 to 0.3. This trend proves that the filling ratio greatly affects the rectification ratio of the diode, because the filling ratio directly affects the optical characteristics of the VO_2 grating structure, resulting in changes in the surface waves across the interfaces. It is also easy to understand that the VO_2 grating behaves as a thin film with a relatively weaker thermal rectification when the filling ratio approaches to 1, which has been proved in Ref. 4. The rectification ratio will continue to grow when the filling ratio further decreases ($\phi < 0.1$), but it is impractical as the grating period is $\Lambda = 50$ nm. Meanwhile, it is apparent that the rectification ratio curves for HITTS cases are above those of PDMS cases for all values of the filling ratio, deviating most notably at low values of ϕ . Therefore, it can be concluded that the HITTS can enhance the thermal rectification performance with respect to other substrates.

To illustrate how other geometric parameters affect the thermal rectification of the diode, Fig. 4(b) provides results for the rectification ratio R simulated with different VO_2 grating heights (h_2). For the four HITTSs, the rectification ratio sharply decreases as the grating height h_2 increases from 10 to 100 nm. Conversely, the PDMS film rectification ratio trends upwards with gradually decreasing slopes. Here, we plot spectral heat flux $dq/d\lambda$ in forward and reverse biases for KBr and

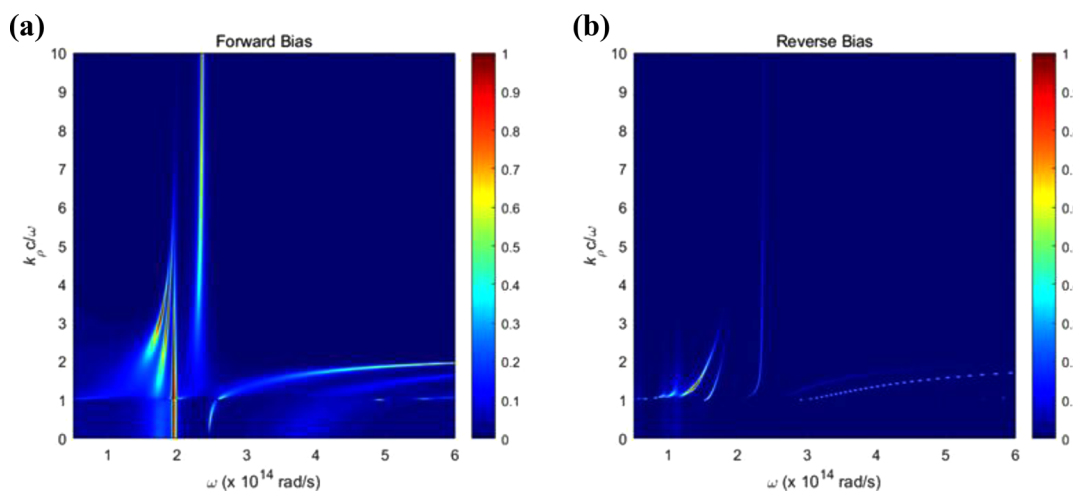


FIG. 3. Coefficient of energy transmission $\xi(\omega, k_p)$ across the two interfaces of an ultrahigh-rectification thermal diode with the KBr substrate against angular frequency ω and normalized parallel wavevector $k_p c/\omega$ for the (a) forward bias and (b) reverse bias.

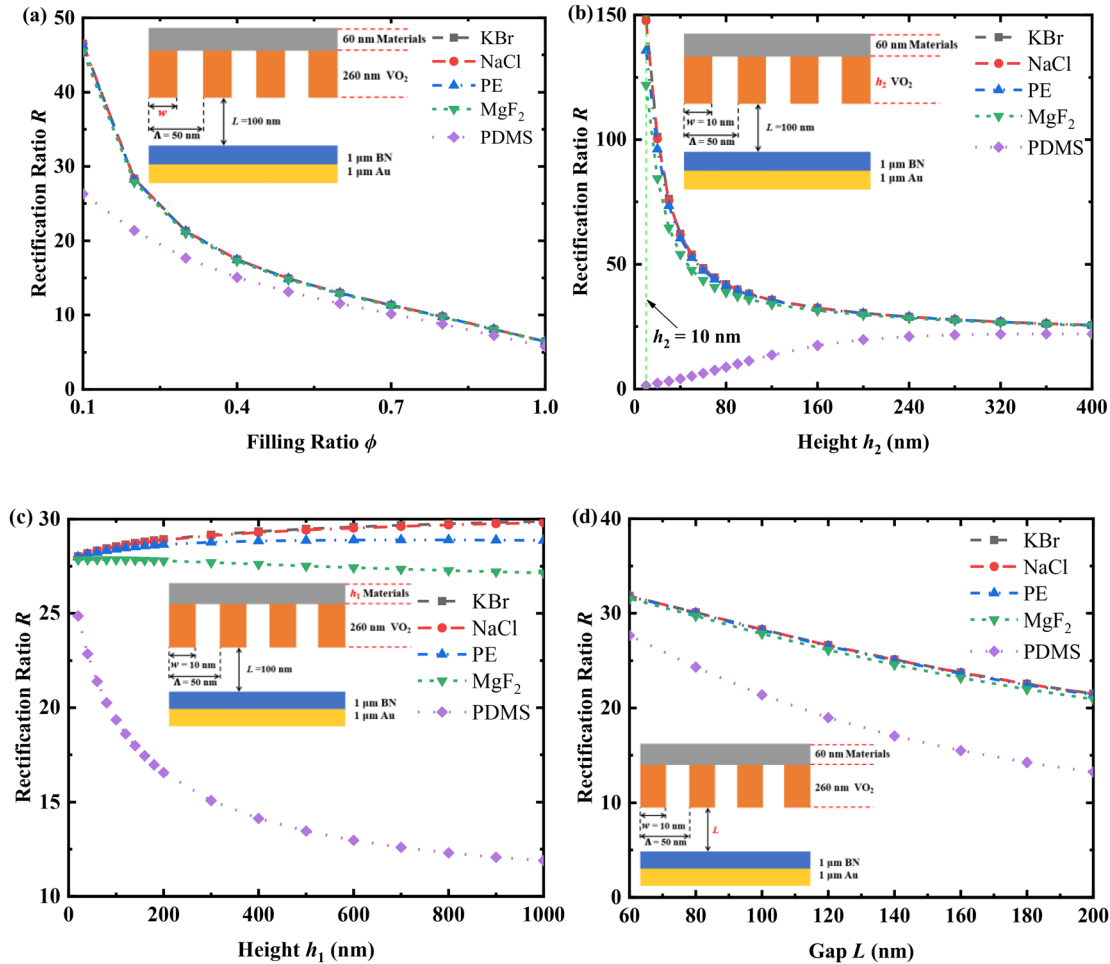


FIG. 4. Rectification ratio R as a function of the following geometric parameters: (a) filling ratio ϕ , (b) height of VO_2 grating h_2 , (c) thickness of the substrate film h_1 , and (d) gap L . For all parameters not varied in each figure, the following constant values are used: $h_1 = 60$ nm, $h_2 = 260$ nm, $\Lambda = 50$ nm, $\phi = 0.2$, and $L = 100$ nm. The temperature of the active side is set as 351 K for the forward bias and 331 K for the reverse bias, and that of the passive side is kept constant at 341 K.

PDMS cases with different VO_2 grating heights, including 20, 40, 60, 80, and 100 nm, at a gap of 100 nm shown in Fig. 5. The reason for a pair of opposite trends is that the increase in the grating height can improve the radiation absorption emitted from the passive side when the diode is in the reverse bias for HITTS cases shown in Fig. 5(a), while the optical properties of the PDMS film can cause the two prominent absorption peaks around the 7.6 and 9.8 μm to descend with the increase in the height of VO_2 grating shown in Fig. 5(b), resulting in the decrease in the overall heat flux in the reverse bias. When h_2 reaches a certain threshold (approximately greater than 200 nm), the rectification ratios become quite similar for all different substrate materials. This indicates that the HITTS affects the thermal rectification of diodes to a much greater degree when the VO_2 grating height is significantly smaller than 100 nm.

Next, Fig. 4(c) shows the thermal rectification ratios calculated vs the thickness of the thin film substrate h_1 for the five different materials mentioned above. For HITTSs made by KBr and NaCl, respectively, an increase in film thickness h_1 leads to a slow increase in the

rectification ratio. In the case of the PE or MgF_2 backing, the slope remains nearly the same for all values of the film thickness. On the other hand, the rectification ratio for the PDMS-backed diode trends downwards with different slopes as h_1 increases. Because the increased thickness of the PDMS thin film substrate can enhance the radiation absorption of the active side when the diode is in the reverse bias, while HITTSs cannot achieve it as mentioned above. Thus, we can conclude that compared with the PDMS film backing, the change in the thickness of HITTS h_1 has little impact on the thermal rectification of the diode.

Figure 4(d) displays the rectification ratio as a function of the diode gap L for five cases with different substrate films. It can be seen clearly that R is inversely proportional to the gap L . This is because surface waves gradually become more dominant at smaller gaps, and vice versa.

Based on the above analysis of design parameters, the highest performing values for the four parameters are determined: $h_1 = 1$ μm , $\Lambda = 50$ nm, $\phi = 0.2$, and $L = 100$ nm. With these values, and by

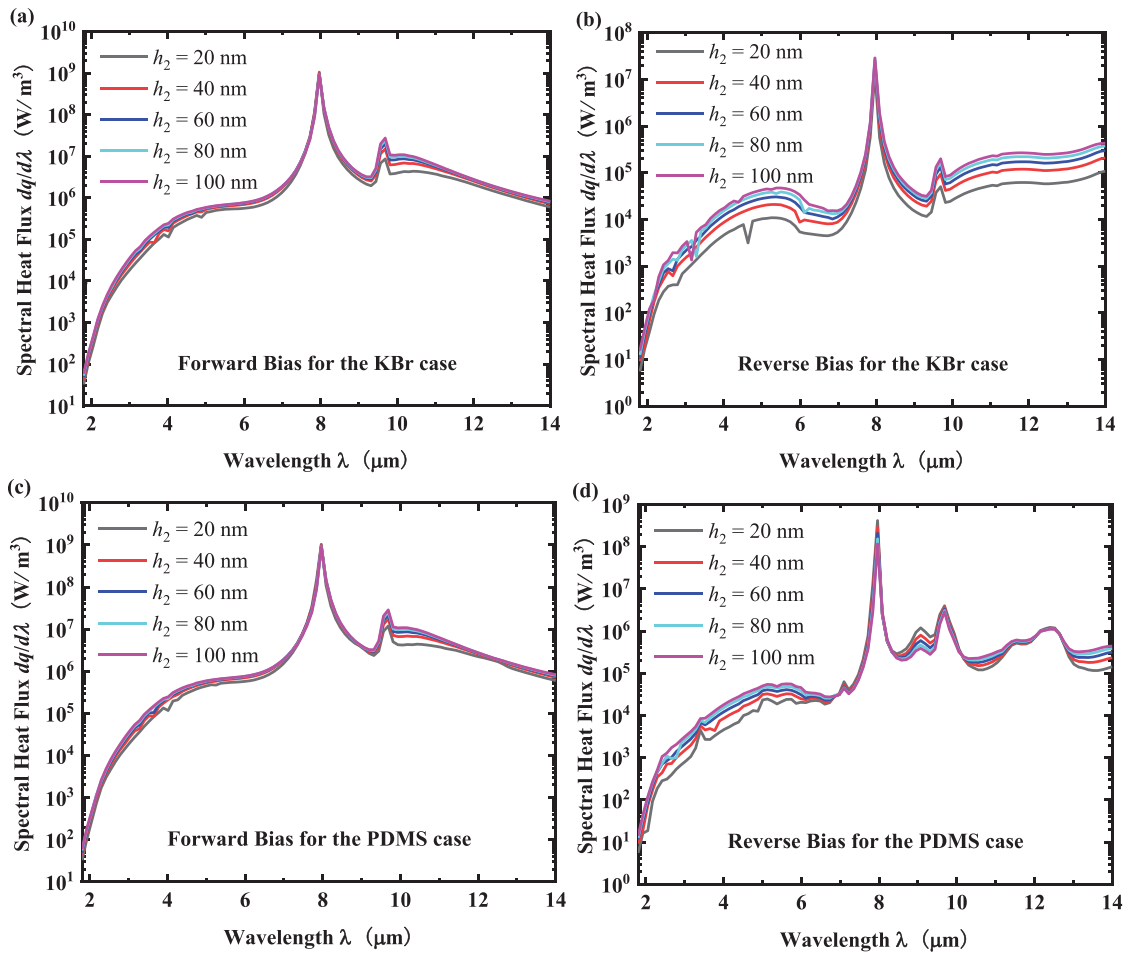


FIG. 5. Comparison of spectral heat fluxes in forward and reverse biases for KBr and PDMS thin-film substrates with different VO₂ grating heights, including 20, 40, 60, 80, and 100 nm. (a) KBr case: forward bias, (b) KBr case: reverse bias, (c) PDMS case: forward bias, and (d) PDMS case: reverse bias.

optimizing the height of the VO₂ grating h_2 , a greatly increased diode rectification ratio can be achieved, as shown in Fig. 6. It is apparent that, for the HITTS, the rectification ratio increases with a decrease in the height of the VO₂ grating. Conversely, for the PDMS case, the rectification ratio decreases with the decrease in h_2 . The rectification ratios for the HITTS cases are also much larger than those of the PDMS case. The rectification ratio reaches a maximum value of $R = 161$ when the VO₂ grating height is reduced to $h_2 = 10$ nm for the KBr-backed diode. While the value of rectification ratio R can reach a higher value when $h_2 < 10$ nm, the VO₂ grating height $h_2 < 10$ nm is impractical for the nanofabrication process. Therefore, considering the above comprehensive analysis of design parameters and the feasibility of nanofabrication, a set of optimal parameters is determined: $h_1 = 1$ μm, $h_2 = 10$ nm, $\Lambda = 50$ nm, $\phi = 0.2$, and $L = 100$ nm. By utilizing a 1D grating made of the phase-transition material VO₂ and a thin-film substrate made of high-infrared-transmittance material KBr, an ultrahigh rectification ratio of 161 can be obtained—a substantial increase compared with previous studies.

In summary, this theoretical investigation proves that highly efficient thermal rectification can be achieved using the 1D VO₂ grating metasurface backed with a large variety of high-infrared-transmittance thin-film substrates. This study shows that the design parameters of the thermal diode play an incredibly important role in enhancing rectification, such as the filling ratio, the height of the VO₂ grating, the thickness of the substrate film, and the selection of substrate materials. Through the study of four materials (KBr, NaCl, PE, and MgF₂), we conclude that thin film substrates composed of high-infrared-transmittance materials such as these can greatly promote the thermal rectification rate of a diode. With optimized geometric parameters and a thin film substrate made of KBr, an ultrahigh rectification ratio of 161 can be obtained at a gap of 100 nm. This work verifies that employing a highly infrared transmissive thin film as the active side substrate plays a significant role in the enhancement of the near-field thermal rectification of a diode. The data and analyses shed light on high performance thermal rectification devices, which can be leveraged for small-scale thermal management in the future.

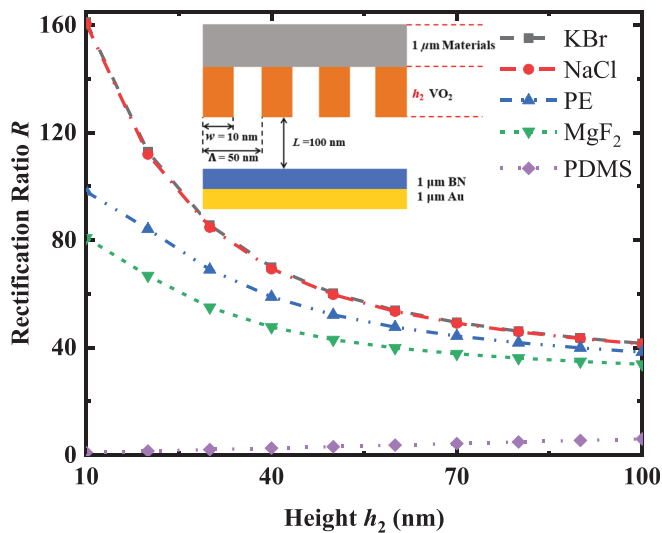


FIG. 6. Effect of VO_2 grating height h_2 on the rectification ratio R when $h_1 = 1 \mu\text{m}$, $\Lambda = 50 \text{ nm}$, $\phi = 0.2$, and $L = 100 \text{ nm}$.

This work was supported by the National Science Foundation (No. CBET-1941743).

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

REFERENCES

- S. Basu and M. Francoeur, *Appl. Phys. Lett.* **98**, 113106 (2011).
- T. Ruokola, T. Ojanen, and A. P. Jauho, *Phys. Rev. B* **79**, 144306 (2009).
- C. R. Otey, W. T. Lau, and S. Fan, *Phys. Rev. Lett.* **104**, 154301 (2010).
- A. Ghanekar, J. Ji, and Y. Zheng, *Appl. Phys. Lett.* **109**, 123106 (2016).
- F. Q. Chen, X. J. Liu, Y. P. Tian, and Y. Zheng, *Adv. Eng. Mater.* **23**, 2000825 (2020).
- P. Ben-Abdallah and S. A. Biehs, *Appl. Phys. Lett.* **103**, 191907 (2013).
- Z. Chen, C. Wong, S. Lubner, S. Yee, J. Miller, W. Jang, C. Hardin, A. Fong, J. E. Garay, and C. Dames, *Nat. Commun.* **5**, 5446 (2014).
- P. Ben-Abdallah and S. A. Biehs, *AIP Adv.* **5**, 053502 (2015).
- P. Ben-Abdallah and S. A. Biehs, *Phys. Rev. Lett.* **112**, 044301 (2014).
- W. Kobayashi, Y. Teraoka, and I. Terasaki, *Appl. Phys. Lett.* **95**, 171905 (2009).
- T. Takeuchi, H. Goto, R. S. Nakayama, Y. I. Terazawa, K. Ogawa, A. Yamamoto, T. Itoh, and M. Mikami, *J. Appl. Phys.* **111**, 093517 (2012).
- D. Sawaki, W. Kobayashi, Y. Moritomo, and I. Terasaki, *Appl. Phys. Lett.* **98**, 081915 (2011).
- W. Kobayashi, D. Sawaki, T. Omura, T. Katsufuji, Y. Moritomo, and I. Terasaki, *Appl. Phys. Express* **5**, 027302 (2012).
- M. Martínez-Pérez, A. Fornieri, and F. Giazotto, *Nat. Nanotechnol.* **10**, 303 (2015).
- D. Polder and M. V. Hove, *Phys. Rev. B* **4**, 3303 (1971).
- A. I. Volokitin and B. N. J. Persson, *Rev. Mod. Phys.* **79**, 1291 (2007).
- J. L. Song, Q. Cheng, B. Zhang, L. Lu, and R. Hu, *Rep. Prog. Phys.* **84**, 036501 (2021).
- Y. M. Xuan, *Photonics Nanostruct.-Fundam. Appl.* **12**, 93 (2014).
- B. Song, D. Thompson, A. Fiorino, Y. Ganjeh, P. Reddy, and E. Meyhofer, *Nat. Nanotechnol.* **11**, 509 (2016).
- A. Fiorino, D. Thompson, L. Zhu, B. Song, P. Reddy, and E. Meyhofer, *Nano Lett.* **18**, 3711 (2018).
- M. Lim, J. Song, S. S. Lee, and B. J. Lee, *Nat. Commun.* **9**, 4302 (2018).
- J. DeSutter, L. Tang, and M. Francoeur, *Nat. Nanotechnol.* **14**, 751 (2019).
- S. Basu, Z. M. Zhang, and C. J. Fu, *Int. J. Energy Res.* **33**, 1203 (2009).
- A. Narayanaswamy and G. Chen, *Phys. Rev. B* **77**, 075125 (2008).
- S. Shen, A. Narayanaswamy, and G. Chen, *Nano Lett.* **9**, 2909 (2009).
- L. Tang, J. Desutter, and M. Francoeur, *ACS Photonics* **7**, 1304 (2020).
- B. Song, Y. Ganjeh, S. Sadat, D. Thompson, A. Fiorino, V. Fernández-Hurtado, J. Feist, F. J. García-Vidal, J. C. Cuevas, and P. Reddy, *Nat. Nanotechnol.* **10**, 253 (2015).
- K. Kim, B. Song, V. Fernández-Hurtado, W. Lee, W. Jeong, L. Cui, D. Thompson, J. Feist, M. T. H. Reid, F. J. García-Vidal, J. C. Cuevas, E. Meyhofer, and P. Reddy, *Nature* **528**, 387 (2015).
- F. Yang and B. Song, *Phys. Rev. B* **103**, 235415 (2021).
- W. Gu, G. H. Tang, and W. Q. Tao, *Int. J. Heat Mass Transfer* **82**, 429 (2015).
- V. Kubyskyi, S. A. Biehs, and P. Ben-Abdallah, *Phys. Rev. Lett.* **113**, 074301 (2014).
- B. Song, A. Fiorino, E. Meyhofer, and P. Reddy, *AIP Adv.* **5**, 053503 (2015).
- Y. Yang, S. Basu, and L. Wang, *Appl. Phys. Lett.* **103**, 163101 (2013).
- Y. Yang, S. Basu, and L. Wang, *J. Quant. Spectrosc. Radiat. Transfer* **158**, 69 (2015).
- J. Huang, L. Qiang, Z. Zheng, and Y. Xuan, *Int. J. Heat Mass Transfer* **67**, 575 (2013).
- R. Audhkhasi and M. L. Povinelli, *J. Appl. Phys.* **126**, 063106 (2019).
- G. Xu, J. Sun, H. Mao, and T. Pan, *J. Appl. Phys.* **124**, 183104 (2018).
- A. Didari, E. B. Eliolu, T. Okutucu-Özyurt, and M. P. Mengüç, *J. Quant. Spectrosc. Radiat. Transfer* **212**, 120 (2018).
- J. Shen, X. Liu, H. He, W. Wu, and B. Liu, *J. Quant. Spectrosc. Radiat. Transfer* **211**, 1–8 (2018).
- L. P. Wang and Z. M. Zhang, *Nanoscale Microscale Thermophys. Eng.* **17**, 337 (2013).
- K. Ito, K. Nishikawa, H. Iizuka, and H. Toshiyoshi, *Appl. Phys. Lett.* **105**, 253503 (2014).
- A. M. Morsy, M. T. Barako, V. Jankovic, V. D. Wheeler, M. W. Knight, G. T. Papadakis, L. A. Sweatlock, P. Hon, and M. L. Povinelli, *Sci. Rep.* **10**, 13964 (2020).
- F. Q. Chen, X. J. Liu, Y. P. Tian, D. Y. Wang, and Y. Zheng, *J. Quant. Spectrosc. Radiat. Transfer* **259**, 107414 (2021).
- A. Ghanekar, Y. Tian, M. Ricci, S. Zhang, and Z. Yi, *Opt. Express* **26**, A209 (2018).
- Y. Zhang, C. L. Zhou, H. L. Yi, and H. P. Tan, *Phys. Rev. Appl.* **13**, 034021 (2020).
- C. L. Zhou, L. Qu, Y. Zhang, and H. L. Yi, *J. Quant. Spectrosc. Radiat. Transfer* **251**, 107023 (2020).
- P. Zhang, P. Yang, Z. Zheng, and W. Yu, *Int. J. Heat Mass Transfer* **155**, 119707 (2020).
- Q. Z. Li, H. Y. He, Q. Chen, and B. Song, *Phys. Rev. Appl.* **16**, 014069 (2021).
- X. L. Liu and Z. M. Zhang, *Appl. Phys. Lett.* **104**, 251911 (2014).
- X. L. Liu, B. Zhao, and Z. M. Zhang, *Phys. Rev. A* **91**, 062510 (2015).
- A. Narayanaswamy and Y. Zheng, *J. Quant. Spectrosc. Radiat. Transfer* **132**, 12 (2014).
- J. Lussange, R. Guérout, F. S. S. Rosa, J. J. Greffet, A. Lambrecht, and S. Reynaud, *Phys. Rev. B* **86**, 085432 (2012).
- E. N. Glytsis and T. K. Gaylord, *Appl. Opt.* **31**, 4459 (1992).
- D. H. Raguin and G. M. Morris, *Appl. Opt.* **32**, 1154 (1993).
- P. Lalanne and D. Lemerrier-Lalanne, *J. Mod. Opt.* **43**, 2063 (1996).
- Y. B. Chen, Z. M. Zhang, and P. J. Timans, *J. Heat Transfer* **129**, 79 (2007).
- X. L. Liu, T. J. Bright, and Z. M. Zhang, *J. Heat Transfer* **136**, 092703 (2014).