concentrations were all relaxed (geometrically optimized), causing the CdTe($\bar{1}00$) facet to adopt a (2 × 1) reconstruction while the CdTe($\overline{111}$) facet remained a (1 × 1) unreconstructed surface (from hereon, the CdTe facets under investigation will have their reconstructions written in subscript form). The CdTe$(\bar{1}00)_{(2\times1)}$ model retained a Te surface coverage of 1.0 and has been verified to exist via Scanning Tunneling Microscopy [16]. All DFT computations are performed using a Linear Combination of Atomic Orbitals (LCAO) basis scheme. Each CdTe surface modeling domain consists of Cd and Te atoms using the FHI pseudopotential (PP) with a double-zeta polarized basis set and the OMX [17, 18] PP with a low(s2p2d1) basis set, respectively. The Perdew-Zunger [19] form of the LDA exchange-correlation functional was employed in all calculations. A Hubbard-U correction was applied to Cd (Cd-4d = 4.6 eV) and Te (Te-5p = 2.55 eV) to establish the proper CdTe bulk band gap value E_g of 1.49 eV as documented in the previous DFT+SGF Cd-terminated CdTe surface study [15]. The Brillouin zone is sampled by 3 × 3 × 150 (4.38 Å × 4.38 Å × 154.7 Å) and 4 × 4 × 100 (5.05 Å × 5.05 Å × 178.63 Å) for the CdTe($\bar{1}00$) and CdTe($\overline{111}$) atomistic models, respectively.

Each of the Te-terminated CdTe models maintained similar domain lengths and vacuum regions for easier comparison. A representative modeling domain for the CdTe surfaces is presented in Fig. 1a, which consists of a bulk electrode, surface region and vacuum region. The semi-infinite CdTe bulk electrode ensures that the electronic properties far away from the termination layer accurately correspond to a periodic CdTe system. The CdTe surface region is coupled to the bulk electrode by means of the Green's function formalism [20] with a large enough thickness to account for possible screening effects. Electronic features along the CdTe surface presented in the current work were determined to not change with longer modeling domains (please see the Supplementary Material for more details). The vacuum region of each CdTe surface model does not contribute to the LCAO-based calculation of the Hamiltonian matrix and thus is maintained at 5 nm simply for visual consistency. The Dirichlet and Neumann boundary conditions are implemented between the respective bulk/surface and surface/vacuum interfaces following the same procedure as previous work [15]. Only the relaxed modeling domains are presented here as they are more representative of the effects seen on epitaxial surfaces during real-world processing conditions. The respective number of atoms used in the bulk electrode and surface region of the CdTe$(\bar{1}00)_{(2\times1)}$ (CdTe$(\overline{111})_{(1\times1)}$) facet is 16 (24) and 528 (480). The bulk electrode size was set as 6.48 Å(11.22 Å). The eight relaxed CdTe surface models used in the study are illustrated in Figs. 1b thru 1i. A force threshold of 0.01 eV/Å was applied to the last 5 nm of the surface region while all other atoms were fixed in each CdTe surface model. All investigated surface supercell sizes after relaxation are provided in Table 1.

3. Results and discussion

The resulting electronic structures of CdTe surfaces with varied concentrations of Cl passivation are evaluated according to their respective CdTe$(\bar{1}00)_{(2\times1)}$ and CdTe$(\overline{111})_{(1\times1)}$ facets. A detailed description of energy band alignment and surface electronic states associated with passivation and relaxation effects is presented below.

Table 1
CdTe one-probe modeling dimensions of the relaxed surface regions with respect to Cl passivation on the surface for the CdTe$(\bar{1}00)_{(2\times1)}$ and CdTe$(\overline{111})_{(1\times1)}$ facets.

CdTe Facet	$Cl_{Te,surf}$	$L_{z,surf}$ (Å)
$(\bar{1}00)_{(2\times1)}$	0%	213.52
	25%	213.65
	50%	213.52
	100%	213.3
$(\overline{111})_{(1\times1)}$	0%	223.88
	25%	223.85
	50%	223.86
	100%	223.6

3.1. *Cl passivation along CdTe$(\bar{1}00)_{(2\times1)}$ facet*

Fig. 2 depicts the energy band alignments at different Cl surface concentrations for the CdTe$(\bar{1}00)_{(2\times1)}$ surface configuration. The magnitudes of electronic features associated with the band alignments are provided in Table 2. In general, the calculated band gaps ($E_{g,CdTe}$) after applying the Hubbard U correction are in good agreement with the expected value of 1.5 eV since for single crystal CdTe the band gap can vary down to as low as 1.41 eV with temperature changes [21]. Establishing an appropriate CdTe bulk band gap enables a more quantitative representation of the measured energy levels of each band alignment plot.

The original CdTe surface without Cl (0% Cl_{Te}, Fig. 2a) contains minimal band bending effects and thus retains flattened conduction and valence bands across the entire length of the modeling domain. The energetically favorable atomic (2 × 1) reconstruction that involves Te atoms forming dimers as seen in Fig. 1b is responsible for the flat band condition and satisfies the electron counting rule [22]. The relaxed CdTe $(\bar{1}00)_{(2\times1)}$ surface is found to be one of the experimentally stable reconstructed surfaces forming within an excess Te flux and evaporation temperatures exceeding 300–340 °C [23–25]. The stable configuration predicted by the DFT+SGF method after geometry optimization for CdTe($\bar{1}00$) is indeed the (2 × 1) reconstruction and has been reported in past DFT calculations [8].

As the Cl concentration increases to 25% and 50% (Fig. 2b and Fig. 2c), a noticeable upward band bending effect of at least 0.5 eV occurs in both the valence and conduction bands of the CdTe surface. The band bending is due to a negative charging effect that results from charge rearrangement at the surface (please see the Supplementary Material for verification using the Mulliken population analysis). The negative charge is induced by both the surface relaxation and passivation by Cl, which due to its lower valence than Te would cause a surface charge imbalance. As the CdTe($\bar{1}00$) facet contains only two bonds per surface atom, substituting Te with Cl at fractional amounts disrupts the charge balance that was initially maintained as seen for the (2 × 1) reconstruction. The surface cross-sectional views of the 25% and 50% Cl_{Te} passivated CdTe($\bar{1}00$) facets as depicted in Figs. 1c and 1d do not form a dimer to neutralize the dangling bonds, which as a result causes a negative charge along the CdTe($\bar{1}00$) surface. In turn, the negatively charged surfaces due to fractional amounts for Cl passivation force the bands to bend upward with a large screening length. It is noted that no immediate cusp feature develops within the valence band edge. Instead, only a visible plateau forms for the 50% case that is most likely due to the atomic relaxation that induces strain-dependent changes to the valence band within that region.

When the CdTe($\bar{1}00$) surface experiences complete Cl passivation (Fig. 2d), the band bending effect is removed from the band alignment and now includes a small additional cusp potential seen 63.45 Å away from the surface ($L_{z,\mathrm{surf}}$ - $\delta_{\mathrm{v,cusp}}$ = 63.45 Å). A similar cusp feature has been reported in the earlier study for Cd-terminated CdTe surfaces [15] and is directly attributed to the relaxation applied to the model. The cusp may offer an additional driving force for hole charge carriers to move toward the CdTe surface but is not expected to make significant improvements to device performance. In a similar manner as the bare (0% Cl_{Te}) CdTe($\bar{1}00$) (2 × 1) facet, the 100% Cl_{Te} covered CdTe($\bar{1}00$) surface forms dimers that mitigate the presence of dangling bonds. As a result, both the 0% and 100% Cl_{Te} CdTe$(\bar{1}00)_{(2\times1)}$ surfaces demonstrate minimal surface electronic states and flattened valence/conduction bands throughout the modeling domain.

Further details of how the surface atoms contribute to the band

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