



A local Hopf lemma and unique continuation for elliptic equations[☆]



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ABSTRACT

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We prove results on unique continuation at the boundary for the solutions of real analytic elliptic partial differential equations of the form

$$\Delta u + \sum_{i=1}^n b_i(x) \frac{\partial u}{\partial x_i} + c(x)u = 0.$$

The work is motivated by and generalizes the main results of X. Huang et al. in [15] and [16], and M.S. Baouendi and L.P. Rothschild in [5].

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1. Introduction

This paper concerns unique continuation at the boundary and a local Hopf lemma for solutions of

$$Pu = \Delta u + \sum_{i=1}^n b_i(x) \frac{\partial u}{\partial x_i} + c(x)u = 0$$

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where the coefficients of P are real analytic.

The work generalizes the results proved for harmonic functions in [5], [15], [16] and those in [12] for solutions of the Helmholtz's equation $\Delta u + cu = 0$, $c \in \mathbb{R}$. See also [6], [9] and [10].

These uniqueness phenomena extend the classical Hopf lemma about the nonvanishing of the normal derivative at a boundary point where a nonconstant solution attains an extremum: the assumption is local in nature and only imposes conditions at the boundary.

For holomorphic functions of one variable with nonnegative real part on a piece of the boundary, unique continuation and local forms of Hopf's lemma were proved in [7], [15], [16], and [18]. The results were used to prove unique continuation for CR mappings for certain classes of CR manifolds. They were also used to prove a more general Schwarz reflection principle for holomorphic functions mapping the real line into a totally real manifold or a real analytic set. Earlier results along this line appeared in the works [4], [8], [2] and [3]. H. Alexander's paper [3] contains a general local Hopf lemma for holomorphic functions of one variable with applications to unique continuation for CR mappings. See also [11] for an extension of the latter results.

Further extensions of the results of Baouendi and Rothschild were proved by V. Shklover ([22]) and H.S. Shapiro ([21]). In particular, Shklover showed that Theorem 3 in [5] (Theorem 2.1 in this paper) fails in general if the normal direction is replaced with a transverse direction. Shapiro used convolution transforms as discussed in [14] to obtain new proofs and generalizations of the theorems of Baouendi and Rothschild. In the article [23], N. Suzuki established a local Hopf lemma in the spirit of [5] for the one-dimensional heat equation.

The article is organized as follows: Section 2 contains the statements of the results in this work. Section 3 is devoted to the construction of a kernel that effectively (for our purposes here) serves as a Poisson kernel for P on the flat piece of the half ball B_r^+ . The proofs of the theorem and its corollaries are presented in section 4.

2. Statement of the results

We will say that a continuous function u defined on a half ball

$$B_r^+ = \{x = (x', x_n) \in \mathbb{R}^n : |x| < r, x_n > 0\}$$

is flat at 0 if for every positive integer N , there is a constant $C_N > 0$ such that

$$|u(x)| \leq C_N |x|^N.$$

Suppose now $D \subseteq \mathbb{R}^n$ is a smoothly bounded domain, $x_0 \in \partial D$. We will say a function u vanishes to infinite order in a direction v at x_0 , where v is a unit vector pointing inside D and transversal to ∂D if for every N , there is a constant $C_N > 0$ such that:

$$|u(x_0 + tv)| \leq C_N t^N.$$

We also say u vanishes to infinite order on a non-singular smooth curve $S : x = x(t)$, $0 \leq t \leq 1$, in D passing through x_0 and transversal to ∂D if for every N there is $C_N > 0$ such that:

$$|u(x(t))| \leq C_N t^N.$$

It is easy to see that this latter definition is independent of the parametrization.

We recall the main result of [5]:

Theorem 2.1. *Let u be harmonic on the half ball B_r^+ , continuous on the closure. Suppose*

- (1) $u(x', 0) \geq 0$ for $|x'| \leq r$, $x' \in \mathbb{R}^{n-1}$;
- (2) the function $x_n \mapsto u(0', x_n)$ is flat at $x_n = 0$;

Then $u(x', 0) \equiv 0$ for x' near the origin in \mathbb{R}^{n-1} .

Somewhat similar but weaker results under the stronger hypothesis that u is harmonic in the upper half plane and decays exponentially along the y -axis was obtained in [19]. The theorem of Baouendi and Rothschild has the following immediate consequence on boundary unique continuation for harmonic functions:

Corollary 2.2. *Let u be harmonic in B_r^+ , continuous on the closure of B_r^+ . Assume that*

- (1) $u(x', 0) \geq 0$ for $|x'| \leq r$;
- (2) The function u is flat at 0.

Then $u \equiv 0$.

Our generalization is as follows:

Theorem 2.3. *Let u be a solution of*

$$Pu = \Delta u + \sum_{i=1}^n b_i(x) \frac{\partial u}{\partial x_i} + c(x)u(x) = 0$$

in the half ball B_r^+ , C^2 on $\overline{B_r^+}$. Assume that the coefficients of P are real analytic on $\overline{B_r^+}$. Suppose

- (1) $u(x', 0) \geq 0$ for $|x'| \leq r$;
- (2) the function $x_n \mapsto u(0', x_n)$ is flat at $x_n = 0$;

(3) for every positive integer N , the function

$$|x'|^{-N}u(x', 0)$$

is integrable on $|x'| \leq r$.

Then $u(x', 0) \equiv 0$ for x' small.

We remark that by the results in [20], u then extends as a solution in a neighborhood of the origin in \mathbb{R}^n .

Theorem 2.3 has the following consequence on boundary unique continuation:

Corollary 2.4. *Let u be a solution of $Pu = 0$ in B_r^+ , C^2 on the closure of B_r^+ . Assume that*

- (1) $u(x', 0) \geq 0$ for $|x'| \leq r$;
- (2) The function u is flat at 0.

Then $u \equiv 0$.

In [22] the author considered the following refinement of Theorem 2.1 of Baouendi and Rothschild:

Suppose $D \subseteq \mathbb{R}^n$ is a smoothly bounded domain, $x_0 \in \partial D$, $V \subset \partial D$ real analytic, $x_0 \in V$. If u is harmonic in D and continuous on \overline{D} , vanishes to infinite order in a direction v (or a curve) at x_0 , and $u(x) \geq 0$ on V , then $u(x) \equiv 0$ in some neighborhood of x_0 in V .

It was shown in [22] that in the plane, this property holds (for harmonic functions) if and only if V is locally symmetric about the normal to ∂D at x_0 . The author also proved further results for the situation when the normal is replaced by a transversal curve.

These results generalize to the operators P under study:

Corollary 2.5. *Let $n = 2$ and u be a solution of $Pu = 0$ in B_r^+ , C^2 on the closure of B_r^+ . Assume that*

- (1) $u(x) \geq 0$ for $x \in V \subset \mathbb{R}$, $0 \in V$;
- (2) u vanishes to infinite order on an analytic curve S through the origin orthogonal to the x -axis and S is symmetric with respect to the x -axis.

Then u vanishes on some subinterval of V about the origin.

In the case of a general domain D we have:

Corollary 2.6. *Let $n = 2$ and u be a solution of $Pu = 0$ in a domain D , C^2 on the closure of D . Suppose $V \subseteq \partial D$ is real analytic, and is tangent to the real axis at the origin. Assume that*

- (1) $u(x) \geq 0$ for $x \in V$;
- (2) the function $y \mapsto u(0, y)$ vanishes to infinite order at 0;
- (3) V is locally symmetric about the imaginary axis. Then u vanishes on some subinterval of V about the origin.

Corollary 2.7. *Let $n = 2$ and u be a solution of $Pu = 0$ in a domain D , C^2 on the closure of D . Suppose $V \subseteq \partial D$ is real analytic and $x_0 \in V$. Assume that*

- (1) $u(x) \geq 0$ for $x \in V$;
- (2) The function u is flat at x_0 .

Then u vanishes on D .

3. Construction of a Poisson kernel for P

For elliptic differential operators of any order with constant coefficients, Poisson kernels for the upper half space \mathbb{R}^n were constructed by Agmon-Douglis-Nirenberg in the work [1]. For elliptic operators with real analytic coefficients, the existence of a local Poisson kernel was proved in [17]. However, this latter kernel is not explicit, and it doesn't serve our purpose since we will need precise estimates on arbitrarily high order derivatives of the kernel. Let

$$P = P(D_x) = \Delta + \sum_{j=1}^n b_j(x) \frac{\partial}{\partial x_j} + c(x)$$

be as in Theorem 2.3. In this section we will construct what will essentially serve as a Poisson kernel $K = K(x, y')$ for P using the ideas and methods of Hadamard ([13]). In [13] it is shown that P has a fundamental solution $E(x, y)$ ($x, y \in \mathbb{R}^n$) of a form that depends on the parity of n .

Case 1: Assume n is odd. In that case, E has the form

$$E(x, y) = \frac{1}{d(x, y)^{n-2}} \sum_{k=0}^{\infty} E_k(x, y) d(x, y)^{2k}, \quad d(x, y) = |x - y|$$

for x, y in a ball B centered at the origin in \mathbb{R}^n .

We will use the same idea to construct first a solution

$$F(x, y) = \frac{1}{\bar{d}(x, y)^{n-2}} \sum_{k=0}^{\infty} F_k(x, y) \bar{d}(x, y)^{2k}$$

of $P(D_x)F(x, y) = 0$ for $x \neq \bar{y}$, where $\bar{d}(x, y) = |x - \bar{y}|$ and for any $y = (y_1, \dots, y_n)$, $\bar{y} = (y_1, \dots, y_{n-1}, -y_n)$. In addition, the F_k will be required to satisfy the initial condition

$$F_k(x', 0, y) = E_k(x', 0, y).$$

Consider

$$PF = \sum_{k=0}^{\infty} P\left(F_k \bar{d}^{2k-n+2}\right).$$

We have

$$\begin{aligned} & P\left(F_k \bar{d}^{2k-n+2}\right)(x, y) \\ &= (PF_k)\bar{d}^{2k-n+2} + 2\langle \nabla F_k, \nabla(\bar{d}^{2k-n+2}) \rangle \\ &+ F_k \Delta(\bar{d}^{2k-n+2}) + F_k \left(\sum_{j=1}^n b_j \frac{\partial}{\partial x_j} \bar{d}^{2k-n+2} \right) \\ &= (PF_k)\bar{d}^{2k-n+2} + 2(2k-n+2) \left[\sum_{j=1}^{n-1} (x_j - y_j) \frac{\partial F_k}{\partial x_j} + (x_n + y_n) \frac{\partial F_k}{\partial x_n} \right] \bar{d}^{2k-n} \\ &+ (2k-n+2) \left(2kF_k \bar{d}^{2k-n} + \sum_{j=1}^{n-1} b_j(x)(x_j - y_j) + b_n(x)(x_n + y_n) \right) F_k \bar{d}^{2k-n} \\ &= (PF_k) \bar{d}^{2k-n+2} + 2(2k-n+2) \\ & \left[\sum_{j=1}^{n-1} (x_j - y_j) \frac{\partial F_k}{\partial x_j} + (x_n + y_n) \frac{\partial F_k}{\partial x_n} + (\theta(x, y) + k - m - 1) F_k \right] \bar{d}^{2k-n} \end{aligned}$$

where $m = \frac{n-2}{2}$ and

$$\theta(x, y) = \frac{1}{4} \left(2n + 2 \sum_{i=1}^{n-1} b_i(x)(x_i - y_i) + 2b_n(x)(x_n + y_n) \right).$$

Let

$$x(t) = tx + (1-t)\bar{y}, \quad 0 < t \leq 1.$$

Setting the coefficients of the powers of $\bar{d}(x, y)$ equal to 0, we are led to the initial value problems

$$t \frac{d}{dt} F_0(x(t), y) + (\theta(x(t), y) - m - 1) F_0(x(t), y) = 0, \quad (3.1)$$

$$F_0(x', 0, y) = E_0(x', 0, y) \quad (3.2)$$

and for any $k \geq 1$,

$$t \frac{d}{dt} F_k(x(t), y) + (\theta(x(t), y) + k - m - 1) F_k(x(t), y) = -\frac{1}{4(k-m)} P(F_{k-1})(x(t), y), \quad (3.3)$$

$$F_k(x', 0, y) = E_k(x', 0, y). \quad (3.4)$$

For $0 < t \leq 1$, define F_0 by

$$F_0(x(t), y) = E_0(x', 0, y) \exp\left\{2 \int_0^1 b(\tilde{x}(r)) \cdot (x - \bar{y}) dr\right\} \cdot \exp\left\{-2 \int_0^t b(x(r)) \cdot (x - \bar{y}) dr\right\}$$

where $\tilde{x}(r) = r(x', 0) + (1-r)\bar{y}$.

Clearly,

$$F_0(x', 0, y) = E_0(x', 0, y)$$

and equation (3.1) is also satisfied. Thus

$$F_0(x, y) = E_0(x', 0, y) \exp\left\{2 \int_0^1 b(\tilde{x}(r)) \cdot (x - \bar{y}) dr\right\} \cdot \exp\left\{-2 \int_0^1 b(x(r)) \cdot (x - \bar{y}) dr\right\}.$$

For $k \geq 1$, define F_k by

$$\begin{aligned} \frac{t^k F_k(x(t), y)}{F_0(x(t), y)} &= \frac{-1}{4(k-m)} \int_0^t \frac{r^{k-1} P F_{k-1}(x(r), y)}{F_0(x(r), y)} dr \\ &+ \frac{1}{4(k-m)} \int_0^1 \frac{r^{k-1} P F_{k-1}(\tilde{x}(r), y)}{F_0(\tilde{x}(r), y)} dr + \frac{E_k(x', 0, y)}{E_0(x', 0, y)}. \end{aligned}$$

It is easy to see that $F_k(x', 0, y) = E_k(x', 0, y)$. To show that equation (3.3) is satisfied, note that since

$$\begin{aligned} \frac{d}{dt} \left\{ \frac{t^k F_k(x(t), y)}{F_0(x(t), y)} \right\} &= -\frac{t^{k-1} P F_{k-1}(x(t), y)}{4(k-m) F_0(x(t), y)}, \\ t^k \frac{d}{dt} \left(\frac{F_k}{F_0} \right) + k t^{k-1} \frac{F_k}{F_0} &= -\frac{t^{k-1} P F_{k-1}(x(t), y)}{4(k-m) F_0(x(t), y)} \end{aligned}$$

and hence

$$t \frac{dF_k}{dt} - \frac{t}{F_0} \frac{dF_k}{dt} F_k + k F_k = -\frac{P F_{k-1}}{4(k-m)}.$$

The latter together with equation (3.1) leads to (3.3).

Thus for $k \geq 1$,

$$\begin{aligned} \frac{F_k(x, y)}{F_0(x, y)} &= \frac{-1}{4(k-m)} \int_0^1 r^{k-1} P F_{k-1}(x(r), y) dr \\ &+ \frac{1}{4(k-m)} \int_0^1 r^{k-1} P F_{k-1}(\tilde{x}(r), y) dr + \frac{E_k(x', 0, y)}{E_0(x', 0, y)}. \end{aligned}$$

We will next show that there is $\delta > 0$ such that the series

$$\sum_{k=0}^{\infty} F_k(x, y) \bar{d}(x, y)^{2k} \quad (3.5)$$

converges uniformly on $\{(x, y) : x \in B_{\delta}(0), |x - \bar{y}| < \delta\}$.

As in [13], this is accomplished by the method of majorants analogous to the Cauchy-Kowalewski technique. For the purpose of the convergence proof, we may assume that $F_0 \equiv \text{constant}$, since the substitution $u_1 = \frac{u}{F_0}$ reduces $Pu = 0$ to a new partial differential equation

$$P_1(u_1) = P(F_0 u_1) = 0 \quad (3.6)$$

of the same type as P for which such an assumption is valid.

With $F_0 \equiv \text{const}$, for $k \geq 1$,

$$\begin{aligned} F_k(x, y) &= \frac{-1}{4(k-m)} \int_0^1 r^{k-1} P F_{k-1}(x(r), y) dr \\ &+ \frac{1}{4(k-m)} \int_0^1 r^{k-1} P F_{k-1}(\tilde{x}(r), y) dr + E_k(x', 0, y). \end{aligned} \quad (3.7)$$

Assume $y = 0$ first. Let

$$\frac{K\epsilon}{\epsilon - |x_1| - \dots - |x_n|} = \sum_{j=0}^{\infty} \frac{K}{\epsilon^j} \left(\sum_{i=1}^k |x_j| \right)^j$$

be a majorant for the Taylor expansions in powers of x of all of the coefficients of P . In [13] it is shown that if

$$M\{E_k\} = \frac{M_k}{(1 - \frac{|x_1| + \dots + |x_n|}{\epsilon})^{2k}}$$

denotes a majorant of E_k , then

$$M\{E_{k+1}\} = \frac{M_{k+1}}{(1 - \frac{|x_1| + \dots + |x_n|}{\epsilon})^{2k+2}}$$

with

$$M_{k+1} = \frac{k(2k+1)}{2(k+1)(k-m+1)} \left(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2}\right) KM_k$$

(note that [13] has $\frac{n^2}{\epsilon^2}$, but in our particular case, it can be replaced with $\frac{2n}{\epsilon^2}$). Suppose now

$$M\{F_k\} = \frac{\tilde{M}_k}{(1 - \frac{|x_1| + \dots + |x_n|}{\epsilon})^{2k}}$$

for some constant \tilde{M}_k with $\tilde{M}_k \geq M_k$ is a majorant of F_k .

Then

$$M\{PF_k\} = \frac{2k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K\tilde{M}_k}{(1 - \frac{\sum_{j=1}^n |x_j|}{\epsilon})^{2k+3}}. \quad (3.8)$$

We will show that for some constants C_k ,

$$\tilde{M}_{k+1} = 2C_k \tilde{M}_k + M_{k+1}.$$

Recall that $y = 0$. Write

$$F_{k+1}(x, 0) = I_1 + I_2 + I_3,$$

where

$$I_1 = \frac{-1}{4(k+1-m)} \int_0^1 r^k PF_k(x(r), 0) dr,$$

$$I_2 = \frac{1}{4(k+1-m)} \int_0^1 r^k PF_k(\tilde{x}(r), 0) dr,$$

and

$$I_3 = E_{k+1}(x', 0, 0).$$

Using (3.8), with $A_k = 2k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K\tilde{M}_k$, we get:

$$\begin{aligned}
M\{I_1\} &= \frac{A_k}{4(k+1-m)} \int_0^1 [1 + (\frac{|rx_1| + \dots + |rx_n|}{\epsilon})] \frac{r^k}{(1 - \frac{\sum_i r|x_i|}{\epsilon})^{2k+3}} dr \\
&= \frac{A_k}{4(k+1-m)(k+1)} \int_0^1 \frac{d}{dr} \left[\frac{r^{k+1}}{(1 - \frac{\sum_i r|x_i|}{\epsilon})^{2k+2}} \right] dr \\
&= \frac{A_k}{4(k+1-m)(k+1)} \left(\frac{1}{(1 - \frac{\sum_{i=1}^n |x_i|}{\epsilon})^{2k+2}} \right).
\end{aligned}$$

Thus

$$M\{I_1\} = \frac{k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K\tilde{M}_k}{2(k+1)(k-m+1)} \cdot \frac{1}{(1 - \frac{\sum_{i=1}^n |x_i|}{\epsilon})^{2k+2}}.$$

Likewise,

$$M\{I_2\} = \frac{k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K\tilde{M}_k}{2(k+1)(k-m+1)} \cdot \frac{1}{(1 - \frac{\sum_{i=1}^n |x_i|}{\epsilon})^{2k+2}}$$

and from [13], we recall that

$$M\{I_3\} = \frac{M_{k+1}}{(1 - \frac{\sum_{i=1}^n |x_i|}{\epsilon})^{2k+2}}.$$

It follows that when $y = 0$,

$$M\{F_{k+1}\} = \frac{\tilde{M}_{k+1}}{(1 - \frac{\sum_{i=1}^n |x_i|}{\epsilon})^{2k+2}}$$

with

$$\tilde{M}_{k+1} = \frac{k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K\tilde{M}_k}{(k+1)(k-m+1)} + M_{k+1}.$$

We have

$$\begin{aligned}
\frac{\tilde{M}_{k+1}}{\tilde{M}_k} &= \frac{k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K}{(k+1)(k-m+1)} + \frac{M_{k+1}}{\tilde{M}_k} \\
&\leq \frac{k(2k+1)(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K}{(k+1)(k-m+1)} + \frac{M_{k+1}}{M_k}.
\end{aligned}$$

Since

$$\lim_{k \rightarrow \infty} \frac{M_{k+1}}{M_k} = (1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K,$$

given $\delta > 0$, we can get N such that for $k \geq N$,

$$\frac{\tilde{M}_{k+1}}{\tilde{M}_k} < \alpha + \delta,$$

where $\alpha = 3(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K$.

Therefore, for some $c > 0$,

$$\tilde{M}_k \leq c(\alpha + \delta)^k \quad \forall k \geq 1.$$

If $y \neq 0$, we replace x with $x - \bar{y}$ and still arrive at the same estimate for the \tilde{M}_k .

It follows that the series

$$\sum_{k=0}^{\infty} F_k(x, y) \bar{d}(x, y)^{2k}$$

converges on the open set

$$\left\{ (x, y) : |x - \bar{y}|^2 < \frac{(1 - \sum_{i=1}^n \frac{|x_i - \bar{y}_i|^2}{\epsilon})^2}{3(1 + \frac{n}{\epsilon} + \frac{2n}{\epsilon^2})K} \right\}$$

where we used the notation

$$\bar{y}_i = y_i \text{ for } 1 \leq i \leq n-1, \quad \bar{y}_n = -y_n.$$

Case 2: Suppose n is even. This time $m = \frac{n-2}{2}$ is an integer and so the approach used under case 1 breaks down because the formula for F_k involves division by $k - m$. We recall from [13] that the fundamental solution S for P has the form

$$S = \sum_{j=0}^{m-1} U_j(x, y) d(x, y)^{2j-n+2} + V(x, y) \log d(x, y)^2 + W(x, y)$$

with the U_j, V , and W real analytic. Motivated by this, we seek a solution $H = H(x, y)$ of the form

$$H(x, y) = \sum_{j=0}^{m-1} A_j(x, y) \bar{d}(x, y)^{2j-n+2} + B(x, y) \log \bar{d}(x, y)^2 + C(x, y)$$

where A_j, B, C are real analytic,

$$P(D_x)H(x, y) = 0 \quad \text{for } x \neq \bar{y}, \tag{3.9}$$

$$A_j(x', 0, y) = U_j(x', 0, y), \quad 0 \leq j \leq m-1, \tag{3.10}$$

$$B(x', 0, y) = V(x', 0, y), \tag{3.11}$$

and

$$C(x', 0, y) = W(x', 0, y). \quad (3.12)$$

At a point $(x(t), y)$, $x(t) = tx + (1-t)\bar{y}$, we have

$$\begin{aligned} P(A_j \bar{d}^{2j-n+2}) \\ = P(A_j) \bar{d}^{2j-n+2} + 2(2j-n+2)t \frac{dA_j}{dt} \bar{d}^{2j-n} \\ + A_j(j-m)4j\bar{d}^{2j-n} + A_j(j-m) \left(\sum_{j=1}^n b_j \frac{\partial \bar{d}^2}{\partial x_j} \right) \bar{d}^{2j-n} \\ = P(A_j) \bar{d}^{2j-n+2} + 2(2j-n+2)[t \frac{d}{dt} A_j + (\theta(x(t), y) + j-m-1)A_j] \bar{d}^{2j-n} \end{aligned}$$

where

$$\theta(x, y) = \frac{1}{4} \left(2n + 2 \sum_{i=1}^n b_i(x)(x_i - \bar{y}_i) \right),$$

$$\bar{y}_i = y_i \text{ for } 1 \leq i \leq n-1, \bar{y}_n = -y_n.$$

$$\begin{aligned} P(B \log \bar{d}^2) &= (\Delta B) \log \bar{d}^2 + \frac{2(n-2)B}{\bar{d}^2} + \frac{4}{\bar{d}^2} \sum_{i=1}^n (x_i - \bar{y}_i) \frac{\partial B}{\partial x_i} \\ &+ \left(\sum_{j=1}^n b_j \frac{\partial B}{\partial x_j} \right) \log \bar{d}^2 + \frac{2B}{\bar{d}^2} \sum_{j=1}^n b_j (x_j - \bar{y}_j) + cB \log \bar{d}^2 \\ &= (PB) \log \bar{d}^2 + \frac{2(n-2)}{\bar{d}^2} B + \frac{2B}{\bar{d}^2} \sum_{j=1}^n b_j (x_j - \bar{y}_j) + \frac{4}{\bar{d}^2} \sum_{j=1}^n (x_j - \bar{y}_j) \frac{\partial B}{\partial x_j} \\ &= (PB) \log \bar{d}^2 + \frac{4}{\bar{d}^2} \left\{ (\theta - 1)B + \sum_{i=1}^n (x_i - \bar{y}_i) \frac{\partial B}{\partial x_i} \right\}. \end{aligned}$$

We choose the A_j like the F_j , that is,

$$A_0(x, y) = U_0(x', 0, y) \exp \left\{ 2 \int_0^1 b(\tilde{x}(r)) \cdot (x - \bar{y}) dr \right\} \cdot \exp \left\{ -2 \int_0^1 b(x(r)) \cdot (x - \bar{y}) dr \right\}$$

and for $1 \leq j \leq m-1$,

$$\begin{aligned} \frac{A_j(x, y)}{A_0(x, y)} &= \frac{-1}{4(j-m)} \int_0^1 r^{j-1} P A_{j-1}(x(r), y) dr \\ &+ \frac{1}{4(j-m)} \int_0^1 r^{j-1} P A_{j-1}(\bar{x}(r), y) dr + \frac{U_j(x', 0, y)}{U_0(x', 0, y)}. \end{aligned}$$

Thus

$$PH = 0 \text{ for } x \neq \bar{y}$$

if

$$\frac{1}{\bar{d}^2} \{P(A_{m-1}) + 4[t \frac{dB}{dt} + (\theta - 1)B]\} + P(B) \log \bar{d}^2 + P(W) = 0.$$

The logarithmic term has to cancel out and so we will also need B to satisfy

$$PB = 0 \quad (3.13)$$

Moreover, we need B to satisfy the equation

$$t \frac{dB}{dt} + (\theta - 1)B = \frac{-P(A_{m-1})}{4} \quad (3.14)$$

on the set where $\bar{d}(x, y) = 0$.

We seek B of the form

$$B(x, y) = \sum_{k=0}^{\infty} B_k(x, y) \bar{d}(x, y)^{2k}$$

that satisfy (3.11), (3.13) and (3.14). This will be achieved if B_0 solves

$$\begin{cases} t \frac{dB_0}{dt} + (\theta - 1)B_0 = \frac{-P(A_{m-1})}{4} \\ B_0(x', 0, y) = V_0(x', 0, y) \end{cases} \quad (3.15)$$

and for $k \geq 1$

$$\begin{cases} t \frac{dB_k}{dt} + (\theta + k - 1)B_k = \frac{-1}{4k} PB_{k-1} \\ B_k(x', 0, y) = V_k(x', 0, y). \end{cases} \quad (3.16)$$

Define B_0 by

$$\begin{aligned}
t^m \frac{B_0(x(t), y)}{A_0(x(t), y)} &= -\frac{1}{4} \int_0^t \frac{PA_{m-1}(x(r), y)}{A_0(x(r), y)} r^{m-1} dr \\
&\quad + \frac{1}{4} \int_0^1 \frac{PA_{m-1}(\tilde{x}(r), y)}{A_0(\tilde{x}(r), y)} r^{m-1} dr + \frac{V_0(x', 0, y)}{A_0(x', 0, y)}
\end{aligned} \tag{3.17}$$

and for $k \geq 1$, B_k is defined by

$$\begin{aligned}
t^{k+m} \frac{B_k(x(t), y)}{A_0(x(t), y)} &= -\frac{1}{4k} \int_0^t \frac{PB_{k-1}(x(r), y)}{A_0(x(r), y)} r^{k+m-1} dr \\
&\quad + \frac{1}{4k} \int_0^1 \frac{PB_{k-1}(\tilde{x}(r), y)}{A_0(\tilde{x}(r), y)} dr + \frac{V_k(x', 0, y)}{A_0(x', 0, y)}
\end{aligned} \tag{3.18}$$

To see that B_0 defined by (3.17) solves (3.15) observe that since

$$\begin{aligned}
\frac{d}{dt} \left(\frac{t^m B_0}{A_0} \right) &= -\frac{1}{4} \frac{PA_{m-1}}{A_0} t^{m-1}, \\
t \frac{d}{dt} \left(\frac{B_0}{A_0} \right) + \frac{mB_0}{A_0} &= \frac{-PA_{m-1}}{4A_0}
\end{aligned}$$

which implies that

$$t \frac{dB_0}{dt} - \frac{t \frac{dA_0}{dt} B_0}{A_0} + mB_0 = \frac{-PA_{m-1}}{A_0}.$$

Using

$$t \frac{dA_0}{dt} + (\theta - m - 1)A_0 = 0$$

in the latter equation we conclude that

$$t \frac{dB_0}{dt} + (\theta - 1)B_0 = \frac{-PA_{m-1}}{4}.$$

From the formula for $B_0(x, y) = B_0(x(1), y)$, we see that

$$B_0(x', 0, y) = V_0(x', 0, y).$$

When $k \geq 1$, the equation

$$\frac{d}{dt} \left(\frac{t^{k+m} B_k}{A_0} \right) = \frac{-1}{4k} \frac{PB_{k-1}}{A_0} t^{k+m-1}$$

leads to

$$t \frac{dB_k}{dt} - \frac{t \frac{dA_0}{dt} B_k}{A_0} + (k+m)B_k = \frac{-PB_{k-1}}{4k}$$

which together with

$$t \frac{dA_0}{dt} + (\theta - m - 1)A_0 = 0$$

imply that

$$t \frac{dB_k}{dt} + (\theta + k - 1)B_k = \frac{-PB_{k-1}}{4k}.$$

The formula for $B_k(x, y)$ shows that

$$B_k(x', 0, y) = V_k(x', 0, y).$$

The convergence of the series

$$\sum_{k=0}^{\infty} B_k(x, y) \bar{d}(x, y)^{2k}$$

is proved using the majorant method as in case 1.

The functions $A_j(x, y)$ ($0 \leq j \leq m - 1$) and $B(x, y)$ were constructed so that

$$Q(x, y) = P(D_x) \left(\sum_{j=0}^{m-1} A_j \bar{d}^{2j-n+2} + B(x, y) \log \bar{d}^2 \right)$$

is real analytic. We choose $C(x, y)$ that satisfies

$$\begin{cases} P(D_x)C(x, y) = -Q(x, y) \\ C(x', 0, y) = W(x', 0, y). \end{cases}$$

It follows that $H = H(x, y)$ satisfies all the requirements.

By taking $E(x, y) - F(x, y)$ when n is odd and $\underline{S(x, y)} - H(x, y)$ when n is even, we have found a function $G(x, y)$ defined for $(x, y) \in B_r^+ \times B_r^+$ (after decreasing r), $x \neq y$, such that

$$\begin{cases} P(D_x)G(x, y) = \delta(x - y) \text{ on } B_r^+ \times B_r^+ \\ G(x', 0, y) = 0 \text{ when } y_n > 0. \end{cases} \quad (3.19)$$

The transpose of P denoted by ${}^t P$ has the same form as P and so we also have G^t that solves

$$\begin{cases} {}^t P(x, D_x) G^t(x, y) = \delta(x - y) \text{ on } B_r^+ \times B_r^+ \\ G^t(x', 0, y) = 0 \text{ when } y_n > 0. \end{cases} \quad (3.20)$$

We will next show that the function $G^t(x, y) - G(y, x)$ extends as a real analytic function in a neighborhood of $(0, 0)$ in $\mathbb{R}^n \times \mathbb{R}^n$.

Given a C^1 domain $D \subseteq B_r^+$, for $w_1, w_2 \in C^2(\overline{D})$, by Green's identity,

$$\begin{aligned} \int_D \left(w_1(x) P w_2(x) - w_2(x) {}^t P w_1(x) \right) dx &= \int_{\partial D} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &\quad + \int_{\partial D} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma \end{aligned} \quad (3.21)$$

where $b(x) = (b_1(x), \dots, b_n(x))$ and $n(x)$ is unit outer normal vector.

Fix $p, q \in B_r^+$, $p \neq q$. Let $w_1(x) = G^t(x, q)$ and $w_2(x) = G(x, p)$. For $\epsilon > 0$ small, let $D_\epsilon = B_r^+ \setminus (B_\epsilon(p) \cup B_\epsilon(q))$.

Since ${}^t P w_1 = 0$ and $P w_2 = 0$ in D_ϵ , (3.21) becomes

$$\int_{\partial D_\epsilon} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma = - \int_{\partial D_\epsilon} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma \quad (3.22)$$

Write $\partial B_r^+ = \Sigma \cup \Sigma'$ where $\Sigma = \{x \in \partial B_r^+ : x_n = 0\}$.

Since $w_1(x', 0) = 0 = w_2(x', 0)$,

$$\begin{aligned} &\int_{\partial D_\epsilon} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &= \int_{\Sigma'} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &\quad - \int_{\partial B_\epsilon(p)} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &\quad - \int_{\partial B_\epsilon(q)} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \end{aligned} \quad (3.23)$$

Consider the integral

$$\int_{\partial B_\epsilon(p)} w_2(x) \nabla w_1(x) \cdot n(x) d\sigma.$$

The function $\nabla w_1(x) \cdot n(x)$ is continuous on $B_\epsilon(p)$. When n is odd, $w_2(x) = \frac{e(x)}{|x-p|^{n-2}}$ with $e(x)$ continuous on $B_\epsilon(p)$. When n is even, the principal singularity is of the same form except when $n = 2$ in which case it is a logarithmic term. In both cases,

$$\lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(p)} w_2(x) \nabla w_1(x) \cdot n(x) d\sigma = 0.$$

Likewise,

$$\lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(q)} w_1(x) \nabla w_2(x) \cdot n(x) d\sigma = 0.$$

Therefore,

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(p)} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &= \lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(p)} w_1(x) \nabla w_2(x) \cdot n(x) d\sigma \\ &= w_1(p) = G^t(p, q) \end{aligned} \tag{3.24}$$

and

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(q)} \left(w_1(x) \nabla w_2(x) - w_2(x) \nabla w_1(x) \right) \cdot n(x) d\sigma \\ &= - \lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(q)} w_2(x) \nabla w_1(x) \cdot n(x) d\sigma \\ &= -w_2(q) = -G(q, p) \end{aligned} \tag{3.25}$$

$$\begin{aligned} \int_{\partial D_\epsilon} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma &= \int_{\Sigma'} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma \\ &\quad - \int_{\partial B_\epsilon(p) \cup \partial B_\epsilon(q)} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma \end{aligned}$$

and

$$\lim_{\epsilon \rightarrow 0^+} \int_{\partial B_\epsilon(p) \cup \partial B_\epsilon(q)} w_1(x) w_2(x) b(x) \cdot n(x) d\sigma = 0. \tag{3.26}$$

From (3.23)-(3.26) we conclude that

$$\begin{aligned}
G^t(p, q) - G(q, p) &= - \int_{\Sigma'} \left(G^t(x, q) \nabla G(x, p) - G(x, p) \nabla G^t(x, q) \right) \cdot n(x) d\sigma \\
&\quad - \int_{\Sigma'} G^t(x, q) G(x, p) b(x) \cdot n(x) d\sigma \\
&\doteq \varphi(q, p)
\end{aligned} \tag{3.27}$$

which is real analytic on a neighborhood of $(0, 0)$ in $\mathbb{R}^n \times \mathbb{R}^n$, say $B_r(0) \times B_r(0)$ (after decreasing r).

Let $g(x, y)$ be a real analytic function on $B_r(0) \times B_r(0)$ (r may have to be decreased) that is a solution of

$$\begin{cases} P(D_x)g(x, y) = P(D_x)(\varphi_{y_n}(x, y)) \\ g(x', 0, y) = 0 \end{cases}$$

Define

$$K(x, y') = \frac{\partial G^t}{\partial y_n}(y', 0, x) - g(x, y', 0).$$

We will show that $K(x, y')$ behaves like a local Poisson Kernel for P for the upper half space.

Indeed, first note that for $x, y \in B_r^+, x \neq y$,

$$\begin{aligned}
P(D_x) \left(\frac{\partial}{\partial y_n} G^t(y, x) - g(x, y) \right) &= P(D_x) \left(\frac{\partial}{\partial y_n} G^t(y, x) - \varphi_{y_n}(x, y) \right) \\
&= \frac{\partial}{\partial y_n} P(D_x) G(x, y) \equiv 0.
\end{aligned} \tag{3.28}$$

Second, in Section 4 we will show that if we define the function v on B_r^+ by

$$v(x) = - \int_{\mathbb{R}^{n-1}} K(x, y') \psi(y') u(y', 0) dy',$$

where ψ is a smooth cutoff function identically equal to 1 near the origin, and supported in $\{x' : |x'| < r\}$, then v differs from u by a function that actually is real analytic in a neighborhood of the origin in \mathbb{R}^n . In particular, the trace $u(x', 0) - v(x', 0)$ is real analytic in a neighborhood of the origin in \mathbb{R}^{n-1} .

4. Proofs of Theorem 2.3 and its corollaries

We begin by estimating the derivatives

$$\partial_{x_n}^k \left[\frac{\partial G^t}{\partial y_n}(y', 0, x) \right] \text{ at } x = 0, y' \neq 0.$$

Assume first that n is odd. Recall that in that case

$$G^t(x, y) = E(x, y) - F(x, y)$$

with

$$E(x, y) = \sum_{k=0}^{\infty} E_k(x, y) d(x, y)^{2k-2m}$$

and

$$F(x, y) = \sum_{k=0}^{\infty} F_k(x, y) \bar{d}(x, y)^{2k-2m}.$$

Using the identity

$$E_k(y', 0, x) = F_k(y', 0, x) \quad \forall k,$$

we get

$$\begin{aligned} \frac{\partial G^t}{\partial y_n}(y', 0, x) &= \sum_{k=0}^{\infty} \left(\frac{\partial E_k}{\partial y_n}(y', 0, x) - \frac{\partial F_k}{\partial y_n}(y', 0, x) \right) d(y', 0, x)^{2k-2m} \\ &\quad - 4 \left[\sum_{k=0}^{\infty} (k-m) E_k(y', 0, x) d(y', 0, x)^{2k-2m-2} \right] x_n \\ &= \frac{p(x, y')}{d(y', 0, x)^{2m}} - \frac{q(x, y') x_n}{d(y', 0, x)^{2m+2}} \end{aligned}$$

with $p(x, y')$ and $q(x, y')$ real analytic.

For an odd integer M , we will estimate first

$$\partial_{x_n}^M \left(\frac{p}{d^{2m}} \right) = \sum_{N=0}^M \binom{M}{N} \partial_{x_n}^N d^{-2m} \partial_{x_n}^{M-N} p$$

at $x = 0, y' \neq 0$. To compute higher order derivatives, we use the formula of Faá di Bruno according to which

$$\frac{d^N}{dt^N} Q(f(t)) = \sum \frac{N!}{N_1! \dots N_N!} Q^{(N_1+...+N_N)}(f(t)) \prod_{j=1}^N \left(\frac{f^{(j)}(t)}{j!} \right)^{N_j}$$

where the sum is over all N -tuples of nonnegative integers (N_1, \dots, N_N) that satisfy the constraint

$$N_1 + 2N_2 + \dots + NN_N = N.$$

Let

$$f(t) = |y'|^2 + t^2 \text{ and } Q(s) = s^{-m}.$$

At $t = 0$ all the terms in the sum above equal zero except when $2N_2 = N$ in which case we get

$$\frac{N!}{N_2!} Q^{(N_2)}(f(0)).$$

Hence, at $x = 0, y' \neq 0$, if $N = 2N_2$,

$$\partial_{x_n}^N d^{-2m} = (-1)^{N_2} \frac{N!}{N_2!} \frac{m(m+1)\dots(m+N_2-1)}{|y'|^{N+2m}}.$$

Since p and q are real analytic, for some $C > 0$, and for all $k \geq 1$,

$$|\partial_{x_n}^k p| + |\partial_{x_n}^k q| \leq C^{k+1} k!.$$

It follows that at $x = 0, y' \neq 0$,

$$\begin{aligned} \left| \partial_{x_n}^{2k+1} \left(\frac{p}{d^{2m}} \right) \right| &\leq \sum_{N=0, N \text{ even}}^{2k} \frac{(2k+1)!}{N_2!} C^{2k-N+2} \frac{m(m+1)\dots(m+N_2-1)}{|y'|^{N+2m}} \\ &= \frac{(2k+1)! C^{2k+2}}{|y'|^{2m}} \sum_{N=0, N \text{ even}}^{2k} \frac{m(m+1)\dots(m+N_2-1)}{N_2! (C|y'|)^N}. \end{aligned}$$

Using

$$\frac{m(m+1)\dots(m+N_2-1)}{N_2!} \leq \frac{(m+1)\dots(m+k)}{k!},$$

the sum above

$$\leq \frac{C^2 (2k+1)! (m+1)\dots(m+k)}{k! |y'|^{2m+2k}} \sum_{N=0}^{2k} (C|y'|)^{2k-N}$$

and hence choosing $r < \frac{1}{C}$ so that $|y'| < \frac{1}{C}$, for some $C_1 > 0$,

$$\left| \partial_{x_n}^{2k+1} \left(\frac{p}{d^{2m}} \right) \right| \leq C_1 \frac{(2k+1)! (m+1)\dots(m+k)}{k! |y'|^{2m+2k}}. \quad (4.1)$$

For $x = 0, y' \neq 0$,

$$\begin{aligned}
& \partial_{x_n}^{2k+1} \left(\frac{qx_n}{d^{2m+2}} \right) \\
&= (2k+1) \partial_{x_n}^{2k} \left(\frac{q}{d^{2m+2}} \right) \\
&= (2k+1) \sum_{N=0}^{2k-1} \partial_{x_n}^N (d^{-2m-2}) \partial_{x_n}^{2k-N} q + (2k+1) \partial_{x_n}^{2k} (d^{-2m-2}) q \\
&\geq - \frac{(2k+1)! C^{2k+1}}{|y'|^{2m+2}} \sum_{\substack{N=0, \\ \text{even}}}^{2k-1} \frac{(m+1) \cdots (m+N_2)}{N_2!} \frac{1}{(C|y'|)^N} \\
&\quad + (2k+1) \partial_{x_n}^{2k} (d^{-2m-2}) q \\
&\geq - \frac{(2k+1)! (m+1) \cdots (m+k)}{k! |y'|^{2m+2}} \sum_{N=0}^{2k-1} \frac{C^{2k+1}}{(C|y'|)^N} + (2k+1) \partial_{x_n}^{2k} (d^{-2m-2}) q \\
&= - \frac{C(2k+1)! (m+1) \cdots (m+k)}{k! |y'|^{2k+2m+2}} \sum_{N=0}^{2k-1} (C|y'|)^{2k+2-N} \\
&\quad + \frac{(-1)^k (2k+1)! (m+1) \cdots (m+k) q}{k! |y'|^{2k+2m+2}} \\
&\geq - \frac{A(2k+1)! (m+1) \cdots (m+k)}{k! |y'|^{2m+2k}} + \frac{(-1)^k (2k+1)! (m+1) \cdots (m+k) q}{k! |y'|^{2k+2m+2}} \quad (4.2)
\end{aligned}$$

for some $A > 0$. In the second inequality above, we have used the fact that for any $N_2 \leq k$,

$$\frac{(m+1) \cdots (m+N_2)}{N_2!} \leq \frac{(m+1) \cdots (m+k)}{k!}.$$

From (4.1) and (4.2), for k even, since $q(0, 0') = E_0(0, 0) \neq 0$, we get at $x = 0, y' \neq 0$

$$\left| \partial_{x_n}^{2k+1} \frac{\partial G^t}{\partial y_n}(y', 0, x) \right| \geq \frac{C(2k+1)!}{|y'|^{2k+2m+2}}. \quad (4.3)$$

Suppose now n is even. In that case recall that $G^t = S - H$ Where

$$S(x, y) = \sum_{j=0}^{m-1} U_j(x, y) d(x, y)^{2j-n+2} + V(x, y) \log d(x, y)^2 + W(x, y)$$

and

$$H(x, y) = \sum_{j=0}^{m-1} A_j(x, y) \bar{d}(x, y)^{2j-n+2} + B(x, y) \log \bar{d}(x, y)^2 + C(x, y),$$

$$A_j(x', 0, y) = U_j(x', 0, y), B(x', 0, y) = V(x', 0, y) \text{ and } C(x', 0, y) = W(x', 0, y).$$

Since

$$U_j(x', 0, y) = A_j(x', 0, y) \quad (0 \leq j \leq m-1),$$

the arguments for the odd case show that for some $C_3 > 0$, at $x = 0, y' \neq 0$,

$$\partial_{x_n}^{2k+1} \partial_{y_n} \left(\sum_{j=0}^{m-1} U_j d^{2j-n+2} - \sum_{j=0}^{m-1} A_j \bar{d}^{2j-n+2} \right) \geq C_3 \frac{(2k+1)!}{|y'|^{2k+2m+2}}. \quad (4.4)$$

We have:

$$\partial_{x_n}^{2k+1} \partial_{y_n} (V \log d^2) = \partial_{x_n}^{2k+1} (V_{y_n} \log d^2) + 2 \partial_{x_n}^{2k+1} \left(\frac{V x_n}{d^2} \right).$$

The term

$$\partial_{x_n}^{2k+1} (V_{y_n} \log d^2) = \sum_{N=0}^{2k+1} \binom{2k+1}{N} \partial_{x_n}^N \log d^2 \partial_{x_n}^{2k+1-N} V_{y_n}.$$

Faá di Bruno's formula shows that

$$\partial_{x_n}^N \log d^2 = 0 \quad \text{if } N \text{ is odd}$$

and when $N = 2N_2$,

$$\partial_{x_n}^N \log d^2 = \frac{(-1)^{N_2+1} N!}{N_2 |y'|^N}.$$

Choose $C > 0$ such that

$$\left| \partial_{x_n}^{2k+1-N} V_{y_n} \right| \leq C^{2k+2-N} (2k+1-N)!.$$

It then follows that

$$\begin{aligned} |\partial_{x_n}^{2k+1} (V_{y_n} \log d^2)| &\leq \sum_{N=0}^{2k} \frac{(2k+1)! C^{2k+2-N}}{N_2 |y'|^N} \\ &\leq \frac{(2k+1)!}{|y'|^{2k}} C^2 \sum_{N=0}^{2k} (C |y'|)^{2k-N} \\ &\leq \frac{C_4 (2k+1)!}{|y'|^{2k}} \end{aligned} \quad (4.5)$$

where we have assumed that $C |y'| < 1$.

We also have

$$\begin{aligned} \left| \partial_{x_n}^{2k+1} \left(\frac{V_{x_n}}{d^2} \right) \right| &= (2k+1) \left| \partial_{x_n}^{2k} \left(\frac{V}{d^2} \right) \right| \\ &\leq C \frac{(2k+1)!}{|y'|^{2k+2}} \end{aligned} \quad (4.6)$$

for some $C > 0$. Thus for some $C > 0$,

$$\left| \partial_{x_n}^{2k+1} \partial_{y_n} (V \log d^2) \right| \leq C \frac{(2k+1)!}{|y'|^{2k+2}}. \quad (4.7)$$

Likewise,

$$\left| \partial_{x_n}^{2k+1} \partial_{y_n} (V \log \bar{d}^2) \right| \leq C \frac{(2k+1)!}{|y'|^{2k+2}}. \quad (4.8)$$

Since $W(x, y)$ and $C(x, y)$ are real analytic, from (4.4), (4.7) and (4.8), we conclude that for some $C > 0$, at $x = 0, y' \neq 0$,

$$\partial_{x_n}^{2k+1} \partial_{y_n} G^t(y', 0, x) \geq C \frac{(2k+1)!}{|y'|^{2k+2m+2}}. \quad (4.9)$$

It can easily be checked that (4.9) also holds when $n = 2$. Thus from (4.3) and (4.9), we see that for any n , at $x = 0, y' \neq 0$, when k is even,

$$\partial_{x_n}^{2k+1} \partial_{y_n} G^t(y', 0, x) \geq C \frac{(2k+1)!}{|y'|^{2k+2m+2}}. \quad (4.10)$$

Recall that

$$K(x, y') = \frac{\partial G^t}{\partial y_n}(y', 0, x) - g(x, y', 0).$$

Let $\psi = \psi(x') \in C_0^\infty(\mathbb{R}^{n-1})$ supported in $|x'| < r$, $\psi(x') \equiv 1$ on $|x'| \leq \frac{r}{2}$ and $0 \leq \psi \leq 1$.

Define

$$v(x) = - \int_{\mathbb{R}^{n-1}} K(x, y') \psi(y') u(y', 0) dy', \quad x \in B_r^+.$$

Let $w(x) = u(x) - v(x)$ for $x \in B_r^+$. The function w extends as a real analytic function to a neighborhood of the origin. To see this, recall from (3.21) that since u is a solution on B_r^+ and $G^t(y', 0, x) \equiv 0$, for $x \in B_r^+$, $\Sigma = \{(x', 0) \in \partial B_r^+\}$, we have the following representation formula:

$$\begin{aligned}
u(x) &= \int_{\partial B_r^+} u(y) \frac{\partial G^t}{\partial \eta}(y, x) d\sigma - \int_{\partial B_r^+} \frac{\partial u}{\partial \eta}(y) G^t(y, x) d\sigma \\
&\quad - \int_{\partial B_r^+} u(y) G^t(y, x) b(y) \cdot n(y) d\sigma \\
&= - \int_{\Sigma} u(y', 0) \frac{\partial G^t}{\partial y_n}(y', 0, x) dy' - \int_{\partial B_r^+ \setminus \Sigma} u(y) \frac{\partial G^t}{\partial \eta}(y, x) d\sigma \\
&\quad - \int_{\partial B_r^+ \setminus \Sigma} \frac{\partial u}{\partial \eta}(y) G^t(y, x) d\sigma - \int_{\partial B_r^+ \setminus \Sigma} u(y) G^t(y, x) b(y) \cdot n(y) d\sigma \\
&= - \int_{\Sigma} u(y', 0) K(x, y') dy' + f(x) \\
&= v(x) + w(x)
\end{aligned}$$

where clearly $w(x)$ is real analytic on some ball $B_\delta(0)$. Note also that by (3.28) $Pw = 0$ in B_r^+ and so because of analyticity, $Pw = 0$ in $B_\delta(0)$.

The integrability of $|x'|^{-N} u(x', 0)$ for all N and the nature of the singularity of $K(x, y')$ imply that the function

$$v(0', x_n) = - \int K(0', x_n, y') \psi(y') u(y', 0) dy'$$

and hence

$$u(0', x_n) = v(0', x_n) + w(0', x_n)$$

are C^∞ up to $x_n = 0$. Since $u(0, x_n)$ is flat at $x_n = 0$, and $u - v = w$ is real analytic on $B_\delta(0)$, we can find a constant $D > 0$ such that $\forall k$

$$|\partial_{x_n}^{2k+1} v(0)| = |\partial_{x_n}^{2k+1} (u - v)(0)| \leq D^{2k+2} (2k+1)! \quad (4.11)$$

On the other hand, since $u(y', 0) \geq 0$, for k even, by (4.10), we have, for ϵ small,

$$\begin{aligned}
|\partial_{x_n}^{2k+1} v(0)| &= \int_{\Sigma} \partial_{x_n}^{2k+1} K(0', x_n, y')|_{x_n=0} \psi(y') u(y', 0) dy' \\
&\geq C(2k+1)! \int_{\Sigma} \frac{\psi(y') u(y', 0)}{|y'|^{2k+2m+2}} dy' \\
&\geq C(2k+1)! \int_{|y'|<\epsilon} \frac{u(y', 0)}{|y'|^{2k+2m+2}} dy'
\end{aligned}$$

$$\geq \frac{C(2k+1)!}{\epsilon^{2k+2m+2}} \int_{|y'|<\epsilon} u(y', 0) dy' \quad (4.12)$$

Since (4.11) and (4.12) hold for any even integer k , by choosing ϵ small enough (depending only on D), taking the $(2k+2m+2)$ th root and letting $k \rightarrow \infty$, we conclude that $u(x', 0) \equiv 0$ for x' in a neighborhood of the origin in \mathbb{R}^{n-1} .

Proof of Corollary 2.4. The hypotheses of Theorem 2.3 are satisfied and so $u(x', 0) \equiv 0$ for x' small. By the boundary analyticity result of [20], u extends as a real analytic function to some neighborhood of the origin. Since it is flat at an interior point, $u \equiv 0$ in that neighborhood. But since u is real analytic on B_r^+ , u vanishes everywhere.

Proofs of Corollaries 2.5 and 2.6. Observe that a conformal map preserves the form of the operator P . Therefore, the proofs of both corollaries follow from Theorem 2.3 and the following lemma known in connection with Poincare's local problem of conformal geometry (see [22] for the proof).

Lemma 4.1. *Given an arc V (or S), we can find a conformal map sending V into the x -axis (S into the y -axis), $z = 0$ to $z = 0$ and the y -axis (x -axis) into itself if and only if V is locally symmetric about the y -axis (S is locally symmetric about the x -axis).*

Proof of Corollary 2.7. The proof follows from an application of Corollary 2.2 and the boundary regularity of the Riemann mapping.

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