

Secondary structure of peptides mimicking the Gly-rich regions of major ampullate spidroin protein 1 and 2

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Abstract

Spider dragline silk has highly desirable material properties, possessing high extensibility, strength, and biocompatibility. Before it is spun, the constituent proteins are stored in a concentrated dope that is void of fibrils. To investigate the structural properties of the amorphous fiber regions in the dope, computer simulations were performed on model peptides representing the *N. clavipes* Gly-rich regions. Analysis of the secondary structure found predominantly turns, bends and coils; a small 3_1 -helical population decreased with increasing concentration. Interestingly, the population of 3_1 -helices saw a large increase in octanol. These results indicate that the unusual 3_1 -helical secondary structure of the Gly-rich region of the fiber is a consequence of the spinning process, and that the low dielectric environment of the fiber may assist in favoring this structure.

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Introduction

Spider dragline silk is remarkably tough due to a combination of high extensibility and strength [1-3]. Dragline silk fibers are predominantly composed of two large proteins, major ampullate spidroin 1 (MaSp1) and major ampullate spidroin 2 (MaSp2) [4, 5]. The sequence of these proteins consist of non-repetitive termini, which are separated by a highly repetitive core that makes up the majority (~90%) of the sequence. The repetitive core consists of blocks of 40 to 200 residues that are repeated between 20 and 100 times; each block consists of a 4 to 10 residue poly-Ala or poly-Gly-Ala region that is adjacent to a Gly-rich region [2]. In the fiber, the poly-Ala and poly-Gly-Ala regions form $2 \times 5 \times 7$ nm crystallites consisting of antiparallel β -sheets [6, 7]. These crystallites are aligned to the fiber axis and embedded in

an amorphous matrix [1]. The crystallites are responsible for the fiber's strength, while the amorphous Gly-rich regions provide its elasticity [8]. In MaSp1 the Gly-rich regions consist of GGX motifs [4], and in MaSp2 they consist of GPGXX motifs [5] (where X=Ala, Gly, Leu, Gln, Tyr). In silk fibers, the GGX motifs adopt 3_1 -helical structures, also known as polyproline II (PPII) helices [9-19], while the GPGXX motifs adopt type II β -turn structures [20] that collectively enables formation of β -spirals [21]. 3_1 -helices were first identified in fibrous collagen [22], but are relatively rare in folded proteins [23]. They are present in the glycine-rich anaplastic lymphoma kinase [24] and snow-flea antifreeze protein [25], for example, but not all glycine-rich protein motifs form 3_1 -helices [26-31]. In contrast to the more familiar α - and 3_{10} -helices, 3_1 -helices do not form intramolecular hydrogen bonds. They possess 3-fold symmetry, with three residues per turn, and their carbonyl oxygen atoms point away from the helical axis.

Before spider silk is spun into a fiber by a specialized organ (spinneret), the constituent proteins are stored in a highly concentrated solution in an abdominal gland [32]. Depending on the species, this gland consists of two or three different zones [33]. The A zone at the start of the gland is neutral or slightly alkaline and secretes and stores MaSp1 and MaSp2, while the other zones towards the spinnerets are acidic, secrete protein coating, and have a role in preprocessing of the fiber [34]. At concentrations that range between 30-50% [35-37], the proteins are believed to form micellar structures in the A zone, that congregate to form liquid crystal intermediates [38-40]. This viscous, gel-like solution is void of fibrils, through an unknown mechanism that is controlled by the pH [34, 39].

Insights into the protein microstructure in the A zone dope would aid the understanding of this mechanism, but so far, experimental data is somewhat conflicting. Early NMR, FTIR and CD measurements indicated that the dope lacks β -sheets, and consists of a dynamical ensemble of loosely formed α and 3_1 -helices [37], while other CD measurements indicated a structure poor in α -helices and β -sheets [41]. CD indicated that the proteins are in a random coil conformation with some β -sheet structure in the gland [34], and a random coil in aqueous solution [42, 43]. NMR indicated that the dope is dynamically disordered without well-defined secondary structure [44-46] with fast backbone dynamics on the sub-nanosecond time scale [45]. Raman confocal spectromicroscopy and vibrational CD indicated a mixture of random coil and 3_1 -helical segments, with some α -helical structures [47, 48]. Other Raman spectromicroscopy experiments indicated that the GGX and GPGXX motifs have indistinguishable conformations in the dope, which both form a disordered random coil with some α -helical contribution [49].

Here, we present molecular dynamics (MD) simulations to further assess the structural propensities of several peptide sequence motifs of the Gly-rich regions of *Nephila clavipes* MaSp1 and MaSp2. Five peptide sequences with GGX and GPGXX motifs from the amorphous region of *N. clavipes* were studied at differing concentrations to assess the inherent propensity for 3_1 -helix formation under dope-like conditions. In order to fully sample the conformational space, simulations were performed with the multiple scaling replica exchange (MREST) method [50, 51]. This method combines replica exchange [52] with a temperature-scaled potential for the solvent-solvent and solvent-protein interactions [53] and Tsallis biasing [54] of all degrees of freedom for fast convergence. Simulations were also performed in octanol to mimic a low dielectric environment, as would be encountered in the fiber.

Methods

A total of five peptide sequences representative of the Gly-rich regions in MaSp1 [4] and MaSp2 [5] of *Nephila clavipes* were studied (Table 1). The MaSp1 sequences will be indicated by 1a-c, the MaSp2 sequences by 2a-b. Systems consisting of one strand and three strands per box were simulated in explicit water [55] using a solvent layer of 20 Å. In order to fully sample conformational space in reasonable time scales, sampling was performed with the MREST enhanced sampling method [50, 51]; details of this method can be found in the original paper [50]. The MREST simulations were performed with the CHARMM program [56], using six replicas with temperature and Tsallis q factors of $(T, q) = (300 \text{ K}, 1)$, $(332 \text{ K}, 1.00001)$, $(368 \text{ K}, 1.00002)$, $(407 \text{ K}, 1.00003)$, $(451 \text{ K}, 1.00004)$, and $(500 \text{ K}, 1.00005)$ for the single stranded systems, and 8 replicas with $(T, q) = (300 \text{ K}, 1)$, $(312 \text{ K}, 1.000010)$, $(325 \text{ K}, 1.000020)$, $(339 \text{ K}, 1.000030)$, $(353 \text{ K}, 1.000040)$, $(368 \text{ K}, 1.000050)$, $(383 \text{ K}, 1.000055)$, and $(400 \text{ K}, 1.000060)$ for the triple stranded systems. Systems were first heated in the NPT ensemble from 150 to 300 K in 50 K increments over a period of 400 ps, further heated (using 10 K increments of 100 ps) to the desired temperature, and equilibrated in the NVT ensemble for 2 ns. After a further equilibration using MREST for 1 ns, production runs of 40 ns per replica were performed. Swapping attempts were made every 5 ps. Simulations were run until the distribution of measured variables (backbone root mean square deviation (RMSD), secondary structure, and Ramachandran plots) did not change with additional simulation time; convergence was further tested by assessing swapping rates, round-trips of the entire (T, q) ladder for each replica, transitions between different structural clusters, and correlation times.

- Table 1 here -

The low dielectric environment inside a fiber was mimicked by simulations of a single strand in octanol [57]. The setup of this system was similar to the water solvated systems, except that 10 replicas were used with $(T, q) = (300 \text{ K}, 1), (317 \text{ K}, 1.000025), (336 \text{ K}, 1.000030), (355 \text{ K}, 1.000040), (376 \text{ K}, 1.000045), (398 \text{ K}, 1.000050), (421 \text{ K}, 1.000055), (446 \text{ K}, 1.000065), (472 \text{ K}, 1.000075), \text{ and } (500 \text{ K}, 1.000080)$.

The CHARMM 36 force field [58] was used for all simulations. To test whether this force field correctly handled 3_1 -helices, a simulation of the 3_1 -rich snow-flea antifreeze protein in the NPT ensemble was performed. The starting structure was taken from PDB ID 2PNE [25]. A solvent layer of 20 Å of TIP3 [55] water was added beyond the protein. The system was then minimized and heated with backbone restraints of 5 kcal/mol from 150 K to 300 K in 50 K increments over 1 ns each. The restraints were gradually released from 5 kcal/mol, 2.5 kcal/mol, 1 kcal/mol, 0.5 kcal/mol, 0.1 kcal/mol over 2 ns simulations each, followed by a 2 ns unrestrained equilibration and a 100 ns production run using OpenMM [59].

All simulations were performed with periodic boundary conditions using Langevin dynamics [60] and a Monte Carlo barostat for the NPT simulations [61]. A 12 Å cutoff was used for nonbonded interaction with a switching potential. Long-range electrostatics were treated using the particle-mesh Ewald method [62]. Bonds with hydrogen atoms were constrained using SHAKE [63], allowing a timestep of 0.002 ps. Snapshots were saved every 5 ps. Secondary structure was calculated using DSSP [64]. Since 3_1 -helices are not recognized as a secondary structure by this method, an internal script was used to calculate the angle between adjacent C=O backbone vectors. Angles of $120^\circ \pm 20^\circ$ were then classified as 3_1 -helices. Clustering based on backbone RMSD was performed on all 300 K snapshots using the Art-2 algorithm in CHARMM with a radius of 1.5 Å and a maximum error of 10^{-4} [65].

Results

Because of the relative scarcity of 3_1 -helices in known protein structures, and issues with simulating 3_1 -helices in implicit solvent [13], MD simulations in explicit water were performed to test the ability of the CHARMM 36 force field to replicate this structural motif. For this purpose, the snow flea antifreeze protein, a protein with a high content of 3_1 -helices, was simulated (Fig. 1). The protein was stable throughout the simulation and retained its secondary structure without experiencing significant structural changes. A small spike in RMSD was observed ~ 75 ns, which corresponded to a twist in the structure. The average fraction of residues in 3_1 -helical conformation was 27%, which agreed well with the fraction

in the crystal structure (25%). These results show that the CHARMM 36 force field is able to replicate this structural motif with high fidelity.

- Figure 1 here -

Convergence of the MaSp1/2 Gly-rich motifs simulations was tested in a variety of ways. Each replica visited all (T, q) values within 1 ns of sampling, and round-trip traveling across the entire (T, q) ladder occurred for each replica in all simulations. Structural clustering of the trajectories showed frequent transitions between clusters and that each cluster was visited multiple times. Structures transitioned between clusters continuously throughout the simulation. Autocorrelation times for the backbone RMSD using the first frame as a reference were 1 ns for the single stranded systems in water, 1.5 ns in octanol, and 2 ns for the triple stranded systems in water.

Ramachandran plots are shown in Fig. 2. Right-handed 3_1 -helices are typically associated with the region between $(\phi, \psi) = (-60^\circ, 150^\circ)$ to $(-90^\circ, 135^\circ)$ [9, 12, 14-18]; however, this region of the Ramachandran plot overlaps with the region for β -turns [13]. The single stranded 1a-1c peptides had similar Ramachandran plots in water, with major populations centered around $(-60^\circ, 150^\circ)$. Small populations were observed around the helical regions (corresponding to α -helices and 3_{10} -helices) of the Ramachandran plots as well, centered at $(-60^\circ, -150^\circ)$ and $(60^\circ, \pm 150^\circ)$. The single stranded 2a-2b peptides also had the largest population centered at $(-60^\circ, 150^\circ)$ in water. However, the shape of this population differed from the 1a-1c systems, with additional density observed around $(-50^\circ, 180^\circ)$. Populations at $(60^\circ, \pm 150^\circ)$ were greatly reduced in 2a and reduced in 2b, which is likely due to the presence of Pro residues.

- Figure 2 here -

All triple stranded systems showed peaks around $(-60^\circ, 150^\circ)$, as well as small populations in the $\alpha/3_{10}$ -helical region. 1a had a sizeable population around $(60^\circ, 150^\circ)$, which was significantly reduced in the other systems. Peaks around $(-60^\circ, -180^\circ)$ were also observed, and these populations were connected to the $(-60^\circ, 150^\circ)$ populations. Several differences between the Ramachandran plots for the single and triple stranded systems in water were observed. For 1a-1c, the increase in concentration reduced the population of left-handed structures. This was evidenced by the loss of density at $(60^\circ, 150^\circ)$. This loss was also observed for 2b. Increasing concentration also reduced the population in the $\alpha/3_{10}$ -helical regions.

Since the dihedral angles of 3_1 -helices and β -turns overlap [13], the presence of 3_1 -helices and other structural elements was further quantified by analysis of the carbonyl angles and DSSP (Fig. 3). This revealed high propensities for turns/bends and coils for the single stranded systems in water, with relatively small populations of 3_1 -helices. Very few β -sheet and α -helices were found. Overall, 1a-1c had ~45% turns/bends, 45% coils and ~7% 3_1 -helices. Typically, structural formation was local and did not result in the formation of large structural motifs. The distribution of these structural motifs was peaked around 2 residues for all structures, with tails extending to as many as 8 residues. Distributions for bends were the broadest, followed by turns and 3_1 -helices. Distributions for 3_1 -helices were very similar between sequences and ~90% of these helices were 2 residues in length. A similar percentage of turns/bends was seen for 2a-2b (45%), while its overall coil was slightly reduced (42%) and its 3_1 -helical contents was slightly larger (11%) than 1a-1c. These differences are likely due to the presence of proline in peptides 2a-b, which is the only amino acid with a strong 3_1 -helical propensity [23]. Of note was a large 3_1 -helical propensity for the QQ motif of 2a (~45%) and the GPG motif of 2b (~40%).

For the aqueous triple stranded systems the 3_1 -helical population was reduced. For 1a-c this reduction was modest, to 5%, but for 2a-b the reduction was severe (to 3%). These reductions are likely due to interstrand interactions. Coils comprised 35%, while turns/bends were 55%. While $\alpha/3_{10}$ -helices did not form, there was a slight increase in β -sheets. This was more pronounced for 2a-2b (7%), than for 1a-1c (5%). Compared to the single stranded systems, the triple stranded systems were relatively more compact and had less structural variability. Their RMSD distributions were shifted by 0.15 Å to lower values, and their radius of gyration were generally reduced as well.

- Figure 3 here -

Instead of maintaining their secondary structure throughout the simulations, the peptides frequently transitioned between different structural motifs. This observation indicated that the model peptides formed a highly dynamic ensemble, and is in agreement with NMR measurements of the dope [44-46]. High disorder, low α -helical and β -sheet propensity of the aqueous peptides also agreed with experimental findings of the dope [34, 37, 41-49], but the lack of experimental high resolution structural data complicated a detailed comparison.

In octanol the Ramachandran plot for systems 1a-1b showed high density in the $(-60^\circ, -45^\circ)$ range, which corresponds to α -helices (Fig. 2). Both also had populations around $(40^\circ, 20^\circ)$, suggesting the presence of left-handed 3_{10} -helices. Small populations around $(-60^\circ, 150^\circ)$ were observed for 1b. 1c and 2a had

shifted helical populations closer to $(-40^\circ, 0^\circ)$, indicating greater 3_{10} -helical content. 1c also had a left-handed helical population, and a large population at $(-60^\circ, 150^\circ)$ and $(-90^\circ, 135^\circ)$, indicative of right-handed 3_1 -helices. Similar populations were observed in 2a, while a population at $(-60^\circ, 150^\circ)$ was observed in 2b. 2b had small populations in the $\alpha/3_{10}$ -helical regions of the plot, with larger populations at $(-60^\circ, 180^\circ)$ and $(60^\circ, 150^\circ)$. Further secondary structure analysis confirmed that the peptides were more α -helical in octanol (Fig. 3). 1a and 1b had $\sim 12\%$ α -helix content, while 1c had $\sim 5\%$. 1c had similar coil, turns/bends and β -sheet content as in water. 1a-c had dramatic increases in 3_1 -helices: a four-fold increase over water for 1a-b to $\sim 25\%$, and a two-fold for 1c to 12% . Nearly all residues of 1a showed significant 3_1 -helical propensity, while the propensity was especially large for the GAGQ motif of 1b and GLG of 1c. The overall 3_1 -helical propensity was similar in octanol and water for 2a, but double in octanol for 2b (20%). Especially the PGG motif of 2b showed large propensity (over 50%). The GPG motif of 2b was mostly in the bend conformation (80%). The RMSD distributions of 2a and 2b were shifted to slightly lower values than the single stranded water systems, and the distributions widths were reduced. The radii of gyration were also reduced with lower widths compared to the single stranded systems in water, while the distributions had higher widths than for the triple stranded systems in water. More hydrogen bonding was observed in octanol than in water, with broader distributions centered around 0.3-0.5 hydrogen bonds per residue.

Conclusion

Simulations of the snow-flea antifreeze protein demonstrated that the force field used in this study was able to accurately replicate 3_1 -helical secondary structures. Simulations of peptides corresponding to the Gly-rich regions of MaSp1 and MaSp2 showed that these peptides predominantly form coils, turns and bends in water. 3_1 -helices were found at a small fraction, and this fraction decreased with increasing peptide concentration. This data indicates that the Gly-rich regions of MaSp1 and MaSp2 mostly form coils, bends and turns in the aqueous dope. These regions have little inherent propensity toward the 3_1 -helix, and this propensity does not increase with concentration.

These observations therefore suggest that formation of 3_1 -helices in the fiber is a result of the spinning process. While spinning involves a complex and poorly understood interplay of chemical and mechanical factors, our simulations identified a contribution that may aid the formation of 3_1 -helices. We observed that 3_1 -helices were formed significantly more in octanol than in water, and in octanol the 3_1 -helices contributed to a significant fraction of the overall structure. These findings indicate that a lower dielectric environment

favors 3_1 -helices, and imply that the low dielectric environment of the spider silk fiber may directly contribute to the formation and stabilization of these unusual helices.

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Figures

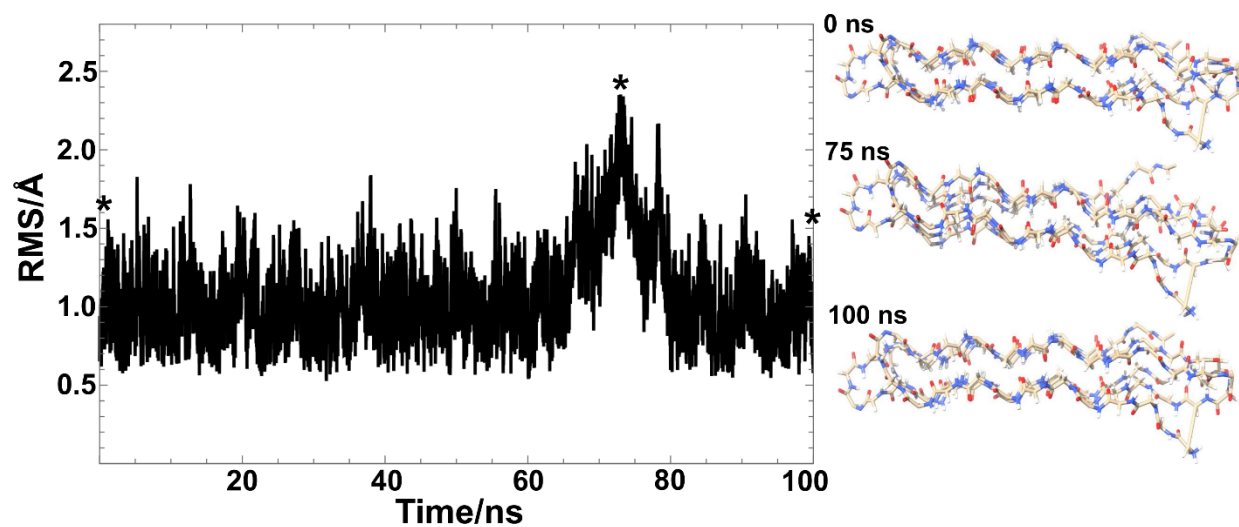


Figure 1. RMSD for the snow-flea antifreeze protein relative to the crystal structure. Snapshots at 0, 75 and 100 ns are shown on the right and indicated with stars on the RMSD graph.

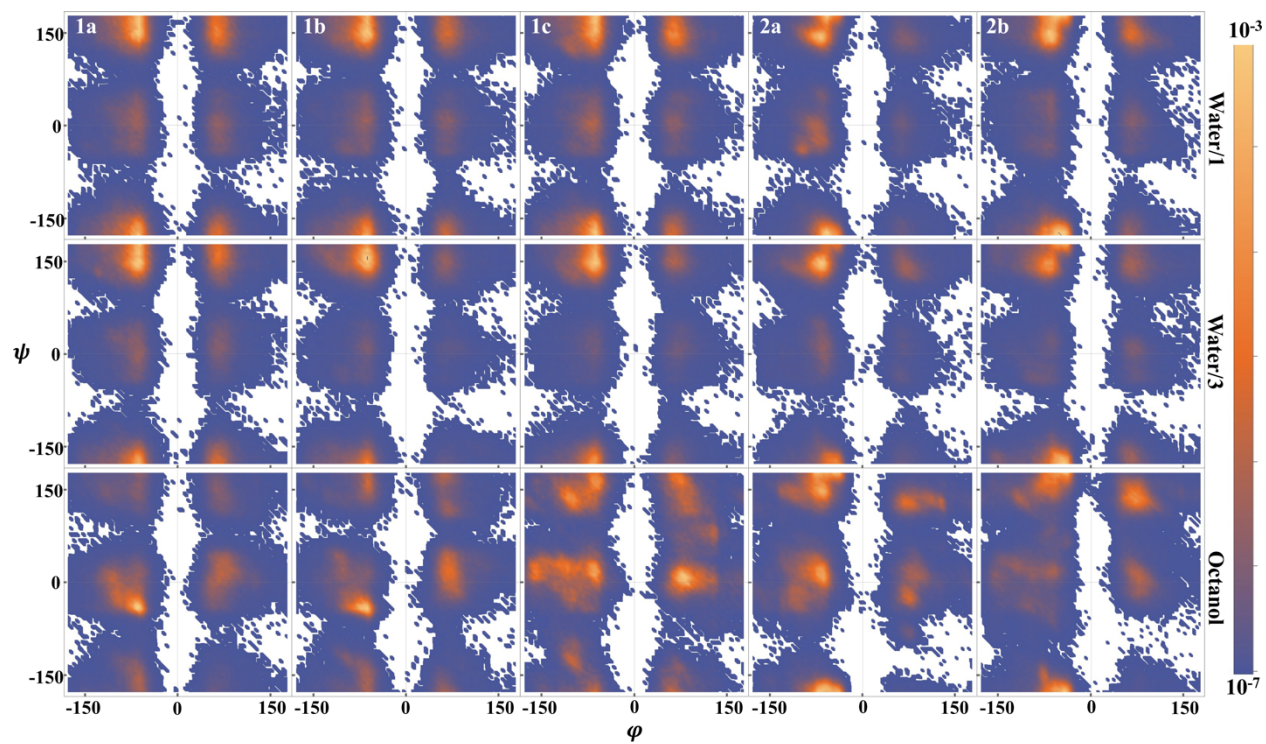


Figure 2. Ramachandran plot for Gly-rich peptides. Colors indicate the probabilities; white regions are unsampled and sterically forbidden.

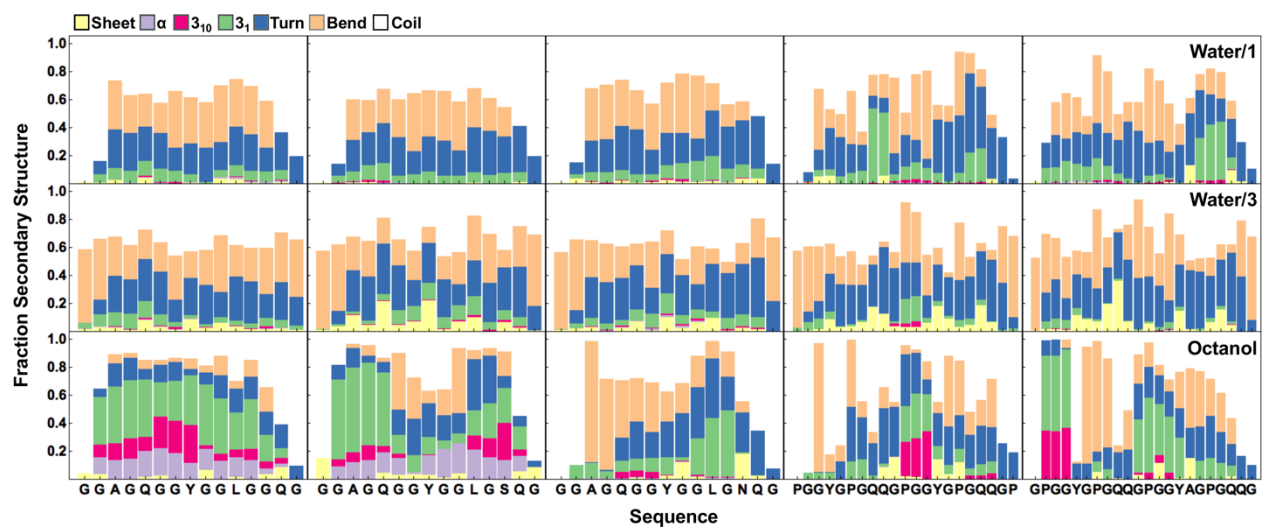


Figure 3. Secondary structure of Gly-rich peptides in water and octanol.

Tables

Table 1. Studied peptide sequences.

Name	Source	Sequence
1a	MaSp1	GGAGQGGYGGLGSQG
1b	MaSp1	GGAGQGGYGGLGSGG
1c	MaSp1	GGAGQGGYGGLGNNG
2a	MaSp2	PGGYGPGQQGPGGYGPGQQGP
2b	MaSp2	GPGGYGPGQQGPGGYAPGQQGP