

Poroelastic near-field inverse scattering

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Abstract

A multiphysics data analytic platform is established for imaging poroelastic interfaces of finite permeability (e.g., hydraulic fractures) from elastic waveforms and/or acoustic pore pressure measurements. This is accomplished through recent advances in design of non-iterative sampling methods to inverse scattering. The direct problem is formulated via the Biot equations in the frequency domain where a network of discontinuities is illuminated by a set of total body forces and fluid volumetric sources, while monitoring the induced (acoustic and elastic) scattered waves in an arbitrary near-field configuration. A thin-layer approximation is deployed to capture the rough and multiphase nature of interfaces whose spatially varying hydro-mechanical properties are a priori unknown. In this setting, the well-posedness analysis of the forward problem yields the admissibility conditions for the contact parameters. In light of which, the poroelastic scattering operator and its first and second factorizations are introduced and their mathematical properties are carefully examined. It is shown that the non-selfadjoint nature of the Biot system leads to an intrinsically asymmetric factorization of the scattering operator which may be symmetrized at certain limits. These results furnish a robust framework for systematic design of regularized and convex cost functionals whose minimizers underpin the multiphysics imaging indicators. The proposed solution is synthetically implemented with application to spatiotemporal reconstruction of hydraulic fracture networks via deep-well measurements.

Keywords: poroelastic waves, waveform tomography, hydraulic fractures, multiphysics sensing

1. Introduction

Coupled physics processes driven by engineered stimulation of the subsurface underlie many emerging technologies germane to advanced energy infrastructure such as unconventional hydrocarbon and geothermal reservoirs [1, 2]. Optimal design and closed-loop control of such systems require real-time feedback on the nature of progressive variations in a target region [3, 4]. 3D in-situ tracking of multiphase developments however is exceptionally challenging [5, 6]. Engineered treatments (such as fluid or gas injection) often occur in a complex background whose structure and hydro-mechanical properties are uncertain [7, 8]. Nevertheless, state-of-the-art solutions to in-situ characterization mostly rely on elastic waves [9, 10, 11, 4]. In particular, micro-seismic waves generated during shear propagation of fractures are commonly used to monitor evolving discontinuities in rock [12, 13]. Such passive schemes are nonetheless agnostic to aseismic evolution [14], and involve significant assumptions on the nature of wave motion in the subsurface which may lead to critical errors [15, 16, 17]. On the other hand, existing approaches to full waveform inversion [18, 19, 11], while offering more comprehensive solutions, are by and large computationally expensive and thus inapplicable for

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real-time sensing. Therefore, there exists a critical need for the next generation of imaging technologies that transcend some of these limitations.

In this vein, recent advances in rigorous design of sampling-based waveform tomography solutions [20, 21, 22, 23, 24] seem particularly relevant owing to their robustness and newfound capabilities pertinent to imaging in complex and unknown domains [21, 23, 25, 26]. So far, these developments mostly reside in the context of self-adjoint and energy preserving systems such as Helmholtz and Navier equations. Wave attenuation, however, is a defining feature of the fractured rock systems with crucial impact on subsurface exploration [27, 28, 29]. In this context, a systematic analysis of sound absorption – e.g., due to the fluid-solid interaction, and its implications on sampling-based data inversion is still lacking. This, in part, may be due to the sparse and single-physics nature of traditional data. Fast-paced progress in sensing instruments such as fiber optic (FBG) sensors – capturing high-resolution multiple physical measurements in time-space [12, 30], may bridge this gap and enable a holistic approach to subsurface characterization. This invites new data analytic tools capable of (a) expedited processing of large datasets, and (b) simultaneous inversion of multiphysics observations. The latter may particularly assist a high-fidelity reconstruction of the sought-for subterranean events.

In this paper, the theoretical foundation of poroelastic inverse scattering is established with application to spatiotemporal tracking of hydraulic fracture networks in unconventional energy systems. This method is nested within the framework of active sensing where the process zone is sequentially illuminated (during stimulation) via deep-well excitations prompting poroelastic scattering whose signature is recorded in the form of seismic and pore pressure waveforms. Thus-obtained multiphase sensory data are deployed for geometric reconstruction of treatment-induced evolution in the target region. Rooted in rigorous foundations of the inverse scattering theory [24], the proposed imaging solution carries a high spatial resolution with carefully controlled sensitivity to noise and illumination frequency.

The forward problem is posed in a suitable dimensional platform catering for simultaneous elastic and acoustic data inversion. A system of arbitrary-shaped hydraulic fractures with non-trivial (i.e., heterogeneous, dissipative, and finitely permeable) interfaces is illuminated by a plausible combination of total body forces and fluid volumetric sources in a medium governed by the coupled Biot equations [31, 32]. Thus-induced multiphase scattered fields are used to define the poroelastic scattering operator mapping the source densities to near-field measurements. In this setting, the reconstruction scheme is based on (i) a custom factorization of the near-field operator, and (b) a sequence of approximate solutions to the scattering equation, seeking (fluid and total-force) source densities whose affiliated waveform pattern matches that of a designated poroelastodynamics solution emanating from a sampling point. The latter is obtained via a sequence of carefully constructed cost functionals whose minimizers can be computed without iterations. Such minimizing solutions are then used to construct a robust multiphysics imaging indicator whose performance is examined through a set of numerical experiments.

In what follows, Section 2 introduces the direct scattering problem and the necessary tools for the inverse analysis. In particular, Section 2.4 investigates the well-posedness of the forward problem which leads to finding the admissibility conditions for the interfacial contact parameters including the elastic stiffness matrix and the permeability, effective stress and Skempton coefficients. This criteria plays a key role later in Section 3 where the essential properties of the poroelastic scattering operator and its components, in the first and second factorizations, are established. Section 4 presents the regularized cost functionals for solving the scattering equations, main theorems, and imaging functionals. Section 5 illustrates a set of numerical experiments implementing the reconstruction of (stationary and evolving) hydraulic fracture networks via the near-field measurements by way of the proposed imaging indicators.

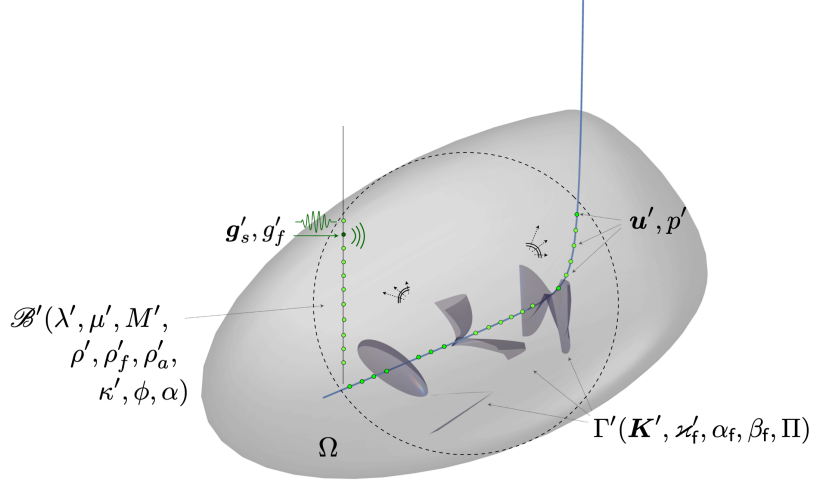


Figure 1: Near-field illumination of a fracture network in the stimulation zone $\mathcal{B}' \subset \Omega$ featuring disjoint multiphasic discontinuities Γ' characterized by heterogeneous elasticity $\mathbf{K}'(\xi')$ and permeability $\kappa'_f(\xi')$. Poroelastic incidents are induced via total body forces and fluid volumetric sources of respective densities $\mathbf{g}'_s(\xi', \omega')$ and $g'_f(\xi', \omega')$ on a designated grid $\xi' \in \mathcal{G}$ at frequency ω' , which give rise to the scattered fields captured on \mathcal{G} in the form of pore pressure p' and solid displacements \mathbf{u}' .

2. Problem statement

Consider the ball \mathcal{B}' of radius $R > 0$, centered at the origin in the subsurface $\Omega \subset \mathbb{R}^3$, encompassing the stimulation zone and sensing grid \mathcal{G} as illustrated in Fig. 1. Let R be sufficiently large so that the probing waves may be assumed negligible in $\Omega \setminus \mathcal{B}'$ due to poroelastic attenuation. It is also assumed that \mathcal{B}' does not meet $\partial\Omega$, which in hydraulic fracturing is a reasonable premise, in particular, since a treatment region is typically four to ten kilometers below the surface while $R = O(10^2)\text{m}$ [33]. The domain $\mathcal{B}' \subset \Omega$ is described by drained Lamé coefficients λ' and μ' , Biot modulus M' , total density ρ' , fluid density ρ'_f , apparent mass density ρ'_a , permeability coefficient κ' , porosity ϕ , and Biot effective stress coefficient α . A system of pre-existing and stimulation-induced fractures and flow networks Γ' is embedded in \mathcal{B}' whose rough and multiphasic interfaces are characterized by a set of heterogeneous parameters, namely: the stiffness matrix $\mathbf{K}'(\xi')$, permeability coefficient $\kappa'_f(\xi')$, effective stress coefficient $\alpha_f(\xi')$, Skempton coefficient $\beta_f(\xi')$, and fluid pressure dissipation factor $\Pi(\xi')$ on $\xi' \in \Gamma'$. The domain is excited by time harmonic total body forces of density $\mathbf{g}'_s(\xi', \omega')$ and fluid volumetric sources of density $g'_f(\xi', \omega')$ on $\xi' \in \mathcal{G}$ at frequency ω' . This gives rise to poroelastic scattering observed on the sensing grid \mathcal{G} in terms of solid displacements $\mathbf{u}'(\xi', \omega')$ and pore pressure $p'(\xi', \omega')$.

2.1. Dimensional platform

To enable multiphasic data inversion, the reference scales

$$\mu_r = \lambda', \quad \rho_r = \rho', \quad \ell_r = 2\pi \sqrt{\frac{\mu'}{(1-\phi)\rho'_s\omega'^2}}, \quad (1)$$

are respectively defined for stress, mass density, and length. Here, ρ'_s denotes the solid-phase density. Note that ℓ_r represents the drained shear wavelength. In this setting, all physical quantities are rendered *dimensionless* [34] as follows

$$(\boldsymbol{\xi}', \omega') \rightarrow (\boldsymbol{\xi}, \omega) := \left(\frac{1}{\ell_r} \boldsymbol{\xi}', \sqrt{\frac{\rho_r}{\mu_r}} \ell_r \omega' \right),$$

$$\begin{aligned} \mathcal{B}'(\lambda', \mu', M', \rho', \rho_f', \rho_a', \kappa', \phi, \alpha) &\rightarrow \\ \mathcal{B}(\lambda, \mu, M, \rho, \rho_f, \rho_a, \kappa, \phi, \alpha) &:= \mathcal{B}\left(\frac{\lambda'}{\mu_r}, \frac{\mu'}{\mu_r}, \frac{M'}{\mu_r}, \frac{\rho'}{\rho_r}, \frac{\rho_f'}{\rho_r}, \frac{\rho_a'}{\rho_r}, \frac{\sqrt{\mu_r \rho_r}}{\ell_r} \kappa', \phi, \alpha\right), \\ \Gamma'(\mathbf{K}', \boldsymbol{\varkappa}'_f, \alpha_f, \beta_f, \Pi) &\rightarrow \Gamma(\mathbf{K}, \boldsymbol{\varkappa}_f, \alpha_f, \beta_f, \Pi) := \Gamma\left(\frac{\ell_r}{\mu_r} \mathbf{K}', \sqrt{\mu_r \rho_r} \boldsymbol{\varkappa}'_f, \alpha_f, \beta_f, \Pi\right), \\ (\mathbf{g}'_s, \mathbf{g}'_f) &\rightarrow (\mathbf{g}_s, \mathbf{g}_f) := \left(\frac{\ell_r}{\mu_r} \mathbf{g}'_s, \mathbf{g}'_f\right), \\ (\mathbf{u}', p') &\rightarrow (\mathbf{u}, p) := \left(\frac{\mathbf{u}'}{\ell_r}, \frac{p'}{\mu_r}\right). \end{aligned} \tag{2}$$

Note that since $\ell_r \propto 1/\omega'$, the dimensionless frequency ω in the first of (2) becomes independent of ω' .

2.2. Field equations

The incident field $(\mathbf{u}^t, p^t) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \overline{\mathcal{G}})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \overline{\mathcal{G}})$ generated by $(\mathbf{g}_s, \mathbf{g}_f) \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ in the background is governed by

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}^t)(\boldsymbol{\xi}) - \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla p^t(\boldsymbol{\xi}) + \omega^2 \left(\rho - \frac{\rho_f^2}{\gamma}\right) \mathbf{u}^t(\boldsymbol{\xi}) &= - \int_{\mathcal{G}} \delta(\mathbf{y} - \boldsymbol{\xi}) \mathbf{g}_s(\mathbf{y}) d\mathbf{y}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \overline{\mathcal{G}}, \\ \frac{1}{\gamma \omega^2} \nabla^2 p^t(\boldsymbol{\xi}) + M^{-1} p^t(\boldsymbol{\xi}) + \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla \cdot \mathbf{u}^t(\boldsymbol{\xi}) &= - \int_{\mathcal{G}} \delta(\mathbf{y} - \boldsymbol{\xi}) g_f(\mathbf{y}) d\mathbf{y}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \overline{\mathcal{G}}, \end{aligned} \tag{3}$$

where $\mathbf{C} = \lambda \mathbf{I}_2 \otimes \mathbf{I}_2 + 2\mu \mathbf{I}_4$ is the fourth-order drained elasticity tensor with \mathbf{I}_m ($m = 2, 4$) denoting the m th-order symmetric identity tensor, and

$$\gamma = \frac{\rho_a}{\phi^2} + \frac{\rho_f}{\phi} + \frac{i}{\omega \kappa}. \tag{4}$$

The interaction of (\mathbf{u}^t, p^t) with Γ gives rise to the scattered field $(\mathbf{u}, p) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \overline{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \overline{\Gamma})$ solving

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}) - \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla p + \omega^2 \left(\rho - \frac{\rho_f^2}{\gamma}\right) \mathbf{u} &= \mathbf{0} \quad \text{in } \mathbb{R}^3 \setminus \overline{\Gamma}, \\ \frac{1}{\gamma \omega^2} \nabla^2 p + M^{-1} p + \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla \cdot \mathbf{u} &= 0 \quad \text{in } \mathbb{R}^3 \setminus \overline{\Gamma}, \\ \mathbf{t} + \tilde{\alpha}_f \llbracket p \rrbracket \mathbf{n} = \mathbf{K} \llbracket \mathbf{u} \rrbracket - \mathbf{t}^t - \tilde{\alpha}_f p^t \mathbf{n} &\quad \text{on } \Gamma, \\ \llbracket p \rrbracket + \beta_f \mathbf{t} \cdot \mathbf{n} = - \frac{\kappa_n \beta_f}{\Pi \alpha_f} \llbracket q \rrbracket - p^t - \beta_f \mathbf{t}^t \cdot \mathbf{n} &\quad \text{on } \Gamma, \\ \llbracket q \rrbracket = \frac{\boldsymbol{\varkappa}_f}{i \omega \Pi} \llbracket p \rrbracket - q^t, \quad \llbracket \mathbf{t} \rrbracket = \mathbf{0} &\quad \text{on } \Gamma, \end{aligned} \tag{5}$$

where $(\mathbf{t}, q) \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma)$ and $(\mathbf{t}^t, q^t) \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma)$ are defined by

$$\begin{aligned} [\mathbf{t}, q](\mathbf{u}, p) &:= [\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u} - \alpha p \mathbf{n}, \frac{1}{\gamma \omega^2} (\nabla p \cdot \mathbf{n} - \rho_f \omega^2 \mathbf{u} \cdot \mathbf{n})] \quad \text{on } \Gamma, \\ [\mathbf{t}^t, q^t](\mathbf{u}^t, p^t) &:= [\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^t - \alpha p^t \mathbf{n}, \frac{1}{\gamma \omega^2} (\nabla p^t \cdot \mathbf{n} - \rho_f \omega^2 \mathbf{u}^t \cdot \mathbf{n})] \quad \text{on } \Gamma; \end{aligned} \tag{6}$$

the unit normal $\mathbf{n} = \mathbf{n}^-$ on Γ is explicitly identified in Section 2.3;

$$\tilde{\alpha}_f := \frac{\alpha_f \Pi}{1 - \alpha_f \beta_f (1 - \Pi)};$$

$(\llbracket \mathbf{u} \rrbracket, \llbracket p \rrbracket, \llbracket q \rrbracket) = (\mathbf{u}^+ - \mathbf{u}^-, p^+ - p^-, q^+ - q^-)$ signifies the jump in (\mathbf{u}, p, q) across Γ , while

$$(\langle \mathbf{u} \rangle, \langle p \rangle, \langle q \rangle) = \frac{1}{2}(\mathbf{u}^+ + \mathbf{u}^-, p^+ + p^-, q^+ + q^-)$$

is the respective mean fields on Γ ; $\mathbf{K} = \mathbf{K}(\boldsymbol{\xi})$, according to [35], is a *symmetric* and possibly *complex-valued* matrix of specific stiffnesses which in the fracture's local coordinates $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{n})(\boldsymbol{\xi})$ may be described as the following

$$\mathbf{K} := k_t \mathbf{e}_1 \otimes \mathbf{e}_1 + k_t \mathbf{e}_2 \otimes \mathbf{e}_2 + \frac{\tilde{\alpha}_f k_n}{\alpha_f \Pi} \mathbf{n} \otimes \mathbf{n} \quad \text{on } \Gamma, \quad (7)$$

wherein $k_t(\boldsymbol{\xi})$ and $k_n(\boldsymbol{\xi})$ are known respectively as the tangential and normal drained elastic stiffnesses. It should be mentioned that the contact condition across Γ makes use of the generalized Schoenberg's model for a poroelastic interface of finite permeability [35].

Remark 2.1. *The contact law at high-permeability interfaces [35] is reduced to the following form*

$$\begin{aligned} \mathbf{t} + \alpha_f p \mathbf{n} &= \mathbf{K}_\infty \llbracket \mathbf{u} \rrbracket - \mathbf{t}^\iota - \alpha_f p^\iota \mathbf{n} && \text{on } \Gamma, \\ p + \beta_f \mathbf{t} \cdot \mathbf{n} &= -k_n \frac{\beta_f}{\alpha_f} \llbracket q \rrbracket - p^\iota - \beta_f \mathbf{t}^\iota \cdot \mathbf{n} && \text{on } \Gamma, \\ \llbracket p \rrbracket &= 0, \quad \llbracket \mathbf{t} \rrbracket = \mathbf{0} && \text{on } \Gamma. \end{aligned}$$

where $\mathbf{K}_\infty := k_t \mathbf{e}_1 \otimes \mathbf{e}_1 + k_t \mathbf{e}_2 \otimes \mathbf{e}_2 + k_n \mathbf{n} \otimes \mathbf{n}$ on Γ .

The formulation of the direct scattering problems may now be completed by requiring that (\mathbf{u}, p) and $(\mathbf{u}^\iota, p^\iota)$ satisfy the radiation condition as $|\boldsymbol{\xi}| \rightarrow \infty$ [36]. On uniquely decomposing the incident and scattered fields into two irrotational parts (p_1, p_2) and a solenoidal part (s) as

$$(\mathbf{u}^\iota, p^\iota) = (\mathbf{u}_{p_1}^\iota, p_{p_1}^\iota) + (\mathbf{u}_{p_2}^\iota, p_{p_2}^\iota) + (\mathbf{u}_s^\iota, p_s^\iota), \quad (\mathbf{u}, p) = (\mathbf{u}_{p_1}, p_{p_1}) + (\mathbf{u}_{p_2}, p_{p_2}) + (\mathbf{u}_s, p_s), \quad (8)$$

the radiation condition can be stated as

$$\begin{aligned} \frac{\partial \mathbf{u}_\varsigma}{\partial r} - i k_\varsigma \mathbf{u}_\varsigma &= o(r^{-1} e^{-a_\varsigma r}) \quad \text{as } r := |\boldsymbol{\xi}| \rightarrow \infty, \quad a_\varsigma := \Im(k_\varsigma), \quad \varsigma = p_1, p_2, s, \quad \mathbf{u} = \mathbf{u}^\iota, \mathbf{u}, \\ \frac{\partial p_\varsigma}{\partial r} - i k_\varsigma p_\varsigma &= o(r^{-1} e^{-a_\varsigma r}) \quad \text{as } r \rightarrow \infty, \quad \varsigma = p_1, p_2, \quad p = p^\iota, p, \end{aligned} \quad (9)$$

uniformly with respect to $\hat{\boldsymbol{\xi}} := \boldsymbol{\xi}/r$ where $k_\varsigma, \varsigma = p_1, p_2, s$, is the complex-valued wavenumber associated with the slow and fast p-waves and the transverse s-wave [37].

Remark 2.2. *It should be noted that (3)-(9) are pertinent to the low-frequency regime of wave motion which, on defining the Biot's characteristic frequency $\omega_c := \phi/(\kappa \rho_f \alpha_\infty)$ [38], imply that $\omega/\omega_c \ll 1$ wherein α_∞ is the real-valued tortuosity. The latter may be justified by noting that the nominal value of ω_c in shale gas, granite, and many sandstones (e.g., Berea, Pecos, Ohio, and Ruhr) is of $O(\text{MHz})$, while the reported range of probing frequencies in the field is generally within $O(\text{Hz})$ - $O(\text{kHz})$ [16].*

2.3. Function spaces

For clarity of discussion, it should be mentioned that the support of Γ may be decomposed into N smooth open subsets $\Gamma_n \subset \Gamma$, $n = 1, \dots, N$, such that $\Gamma = \bigcup_{n=1}^N \Gamma_n$. The support of Γ_n may be arbitrarily extended to a closed Lipschitz surface ∂D_n of a bounded simply connected domain D_n . In this setting, the unit normal vector \mathbf{n} to Γ_n coincides with the outward normal vector to ∂D_n . Let $D = \bigcup_{n=1}^N D_n$ be a multiply connected Lipschitz domain of bounded support such that $\Gamma \subset \partial D$, then Γ is assumed to be an open set relative to ∂D with a positive surface measure, and the closure of Γ is denoted by $\bar{\Gamma} := \Gamma \cup \partial\Gamma$. Following [39], we define

$$\begin{aligned} H^{\pm\frac{1}{2}}(\Gamma) &:= \{f|_{\Gamma} : f \in H^{\pm\frac{1}{2}}(\partial D)\}, \\ \tilde{H}^{\pm\frac{1}{2}}(\Gamma) &:= \{f \in H^{\pm\frac{1}{2}}(\partial D) : \text{supp}(f) \subset \bar{\Gamma}\}, \end{aligned} \quad (10)$$

and recall that $H^{-1/2}(\Gamma)$ and $\tilde{H}^{-1/2}(\Gamma)$ are respectively the dual spaces of $\tilde{H}^{1/2}(\Gamma)$ and $H^{1/2}(\Gamma)$. Accordingly, the following embeddings hold

$$\tilde{H}^{\frac{1}{2}}(\Gamma) \subset H^{\frac{1}{2}}(\Gamma) \subset L^2(\Gamma) \subset \tilde{H}^{-\frac{1}{2}}(\Gamma) \subset H^{-\frac{1}{2}}(\Gamma). \quad (11)$$

Note that since $(\mathbf{u}, p) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$, then by trace theorems $([\![\mathbf{u}]\!], [\![p]\!]) \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma)$, $(\mathbf{t}, q) \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma)$, and $[\![q]\!] \in \tilde{H}^{-1/2}(\Gamma)$.

Remark 2.3 (Wellposedness of the forward scattering problem). *In light of (5), let us define the interface operator $\mathfrak{P} : \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma) \rightarrow H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ such that*

$$\mathfrak{P}([\![\mathbf{u}]\!], [\![p]\!], -[\![q]\!]) := (\mathbf{t} + \mathbf{t}^t, \langle\!\langle q \rangle\!\rangle + q^t, \langle\!\langle p \rangle\!\rangle + p^t) \quad \text{on } \Gamma. \quad (12)$$

Then, (5)-(9) is wellposed provided that \mathfrak{P} is bounded and

$$\Im \langle \mathfrak{P}\varphi, \varphi \rangle_{\Gamma} \leq 0, \quad \forall \varphi \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma) : \varphi \neq \mathbf{0}, \quad (13)$$

given the duality product

$$\langle \cdot, \cdot \rangle_{\Gamma} = \langle H^{-\frac{1}{2}}(\Gamma)^3 \times H^{-\frac{1}{2}}(\Gamma) \times H^{\frac{1}{2}}(\Gamma), \tilde{H}^{\frac{1}{2}}(\Gamma)^3 \times \tilde{H}^{\frac{1}{2}}(\Gamma) \times \tilde{H}^{-\frac{1}{2}}(\Gamma) \rangle.$$

Detailed analysis is included in Appendix B.

2.4. Poroelastic scattering

In this section, the poroelastic scattering operator is introduced and its first and second factorizations are formulated. In what follows, the Einstein's summation convention applies over the repeated indexes.

2.4.1. Incident wave function

Observe that for a given density $L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \ni \mathbf{g} := (\mathbf{g}_s, g_f)$, the poroelastic incident field $(\mathbf{u}^t, p^t) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})$ solving (3) may be recast in the integral form

$$\begin{bmatrix} u_i^t \\ p^t \end{bmatrix}(\boldsymbol{\xi}) = \int_{\mathcal{G}} \begin{bmatrix} U_{ij}^s & u_i^f \\ p_j^s & p^f \end{bmatrix}(\mathbf{y}, \boldsymbol{\xi}) \cdot \begin{bmatrix} g_s^j \\ g_f \end{bmatrix}(\mathbf{y}) d\mathbf{y}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \bar{\mathcal{G}}, \quad i, j = 1, 2, 3, \quad (14)$$

whose kernel is the poroelastodynamics fundamental solution tensor provided in Appendix A. Here, $U_{ij}^s(\mathbf{y}, \boldsymbol{\xi})$ ($i, j = 1, 2, 3$) signifies the i^{th} component of the solid displacement at $\boldsymbol{\xi}$ due to a unit total body force applied

at \mathbf{y} along the coordinate direction j , while p_j^s is the associated pore pressure. Likewise, $u_i^f(\mathbf{y}, \boldsymbol{\xi})$ ($i = 1, 2, 3$) stands for the solid displacement at $\boldsymbol{\xi}$ in the i^{th} direction due to a unit fluid source i.e., volumetric injection at \mathbf{y} , and p^f is the induced pore pressure.

2.4.2. Poroelastic scattered field

By way of the reciprocal theorem of poroelastodynamics [40, 41], one may show that the following integral representation holds for the scattered field $(\mathbf{u}, p) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ satisfying (5)-(9),

$$\begin{bmatrix} u_i \\ p \end{bmatrix}(\boldsymbol{\xi}) = \int_{\Gamma} \begin{bmatrix} \mathbf{T}_{ij}^s & \mathbf{q}_i^s & p_i^s \\ \mathbf{t}_j^f & \mathbf{q}^f & p^f \end{bmatrix}(\mathbf{y}, \boldsymbol{\xi}) \cdot \begin{bmatrix} \llbracket u_j \rrbracket \\ \llbracket p \rrbracket \\ -\llbracket q \rrbracket \end{bmatrix}(\mathbf{y}) dS_{\mathbf{y}}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \bar{\Gamma}, \quad i, j = 1, 2, 3, \quad (15)$$

whose kernel is derived from the fundamental solution tensor and provided in Appendix A. More specifically, \mathbf{T}^s and \mathbf{t}^f indicate the fundamental tractions on Γ ; \mathbf{q}^s and \mathbf{q}^f signify the associated relative fluid-solid displacements across Γ ; p^s and p^f are the fundamental pore pressure solutions.

2.4.3. Poroelastic scattering operator

The linearity of (5) implies that the scattered wave function may be expressed as a linear integral operator. To observe this, let $\mathbb{R}^4 \ni \mathbf{d} := (\mathbf{d}_s, d_p)$ be a constant vector characterizing an incident point source wherein \mathbf{d}_s indicates the amplitude and direction of the total body force, while d_p represents the magnitude of fluid volumetric source both applied at $\mathbf{y} \in \mathcal{G}$. The resulting scattered field is observed in terms of seismic displacements $\mathbf{u}_{\mathbf{d}}$ and pore pressure $p_{\mathbf{d}}$ at $\boldsymbol{\xi} \in \mathcal{G}$. Now, let us define the scattering kernel $\boldsymbol{\mathfrak{S}}(\mathbf{y}, \boldsymbol{\xi}) \in \mathbb{C}^{4 \times 4}$ so that

$$\boldsymbol{\mathfrak{S}}(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{d} := (\mathbf{u}_{\mathbf{d}}, p_{\mathbf{d}})(\boldsymbol{\xi}), \quad \mathbf{y}, \boldsymbol{\xi} \in \mathcal{G}. \quad (16)$$

Then, given $\mathbf{g} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$, one may easily verify that

$$(\mathbf{u}, p)(\boldsymbol{\xi}) = \int_{\mathcal{G}} \boldsymbol{\mathfrak{S}}(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{g}(\mathbf{y}) dS_{\mathbf{y}}, \quad \mathbf{y}, \boldsymbol{\xi} \in \mathcal{G}, \quad (17)$$

where $(\mathbf{u}, p) \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ satisfies (5)-(9).

Accordingly, one may define the poroelastic scattering operator $\Lambda : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ by

$$\Lambda(\mathbf{g}) = (\mathbf{u}, p)|_{\mathcal{G}}. \quad (18)$$

2.4.4. Factorization of Λ

With reference to (3) and (6), let us define the operator $\mathcal{S} : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ given by

$$\mathcal{S}(\mathbf{g}) := (\mathbf{t}^t, q^t, p^t) \quad \text{on } \Gamma, \quad (19)$$

where \mathbf{t}^t specifies the incident field traction, q^t is the specific relative fluid-solid displacement across Γ , and p^t is the incident pore pressure on Γ . Next, define $\mathcal{P} : H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ as the map taking the incident traction, relative flow and pressure (\mathbf{t}^t, q^t, p^t) over Γ to the multiphase scattered data $(\mathbf{u}, p) \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ satisfying (5)-(9). Then, the scattering operator (18) may be factorized as

$$\Lambda = \mathcal{P}\mathcal{S}. \quad (20)$$

Lemma 2.1. *The adjoint operator $\mathcal{S}^*: \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ takes the form*

$$\mathcal{S}^*(\mathbf{a}) := (\mathbf{u}_a, p_a)(\boldsymbol{\xi}), \quad \boldsymbol{\xi} \in \mathcal{G}, \quad (21)$$

where $(\mathbf{u}_a, p_a) \in H_{loc}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{loc}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ solves

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}_a) - (\alpha - \frac{\rho_f}{\bar{\gamma}}) \nabla p_a + \omega^2 (\rho - \frac{\rho_f^2}{\bar{\gamma}}) \mathbf{u}_a &= \mathbf{0} & \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ \frac{1}{\bar{\gamma} \omega^2} \nabla^2 p_a + M^{-1} p_a + (\alpha - \frac{\rho_f}{\bar{\gamma}}) \nabla \cdot \mathbf{u}_a &= 0 & \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ ([\mathbf{u}_a], [p_a], -[q_a]) &= \mathbf{a}, \quad [\mathbf{t}_a] = \mathbf{0} & \text{on } \Gamma. \end{aligned} \quad (22)$$

Here, the ‘bar’ indicates complex conjugate, and

$$[\mathbf{t}_a, q_a](\mathbf{u}_a, p_a) := [\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}_a - \alpha p_a \mathbf{n}, \frac{1}{\bar{\gamma} \omega^2} (\nabla p_a \cdot \mathbf{n} - \rho_f \omega^2 \mathbf{u}_a \cdot \mathbf{n})] \quad \text{on } \Gamma.$$

Note that (\mathbf{u}_a, p_a) satisfy the radiation condition as $|\boldsymbol{\xi}| \rightarrow \infty$, similar to the complex conjugate of (9).

Proof. see Appendix C. ■

On the basis of (15) and (21), the operator \mathcal{P} can be further decomposed as $\mathcal{P} = \bar{\mathcal{S}}^* T$ where the middle operator $T: H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \rightarrow \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)^3$ is given by

$$T(\mathbf{t}^t, q^t, p^t) := ([\mathbf{u}], [p], -[q]) \quad \text{on } \Gamma, \quad (23)$$

such that $(\mathbf{u}, p) \in H_{loc}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{loc}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ satisfies (15) and q is computed from (6). Thanks to this new decomposition of \mathcal{P} , a second factorization of $\Lambda: L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is obtained as

$$\Lambda = \bar{\mathcal{S}}^* T \mathcal{S}. \quad (24)$$

It is worth noting that the asymmetry of the second factorization (24) stems from the non-selfadjoint nature of the Biot system which is evident from (22).

3. Key properties of the poroelastodynamic operators

Assumption 1. *We assume that Γ contains $M \geq 1$ (possibly disjoint) analytic surfaces $\Gamma_m \subset \Gamma$, $m = 1, \dots, M$, whose (unique) analytic continuation ∂D_m identifies a bounded interior domain $D_m \subset \mathcal{B}$. In this setting, ω is not a “Neumann” eigenfrequency of the Navier system*

$$\begin{aligned} \nabla \cdot (\mathbf{C}_u : \nabla \mathbf{u}^m) + \rho \omega^2 \mathbf{u}^m &= \mathbf{0} & \text{in } D_m, \\ \mathbf{n} \cdot \mathbf{C}_u : \nabla \mathbf{u}^m &= \mathbf{0} & \text{on } \partial D_m, \end{aligned} \quad (25)$$

where $\mathbf{C}_u = (\lambda + \alpha^2 M) \mathbf{I}_2 \otimes \mathbf{I}_2 + 2\mu \mathbf{I}_4$.

Lemma 3.1. *Operator $\mathcal{S}^*: \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ in Lemma 2.1 is compact, injective, and has a dense range.*

Proof. Observe that \mathcal{S}^* of Lemma 2.1 possesses the integral form

$$\mathcal{S}^*(\mathbf{a}) = \int_{\Gamma} \bar{\Sigma}(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{a}(\mathbf{y}) dS_{\mathbf{y}}, \quad \Sigma = \begin{bmatrix} \mathbf{T}^s & \mathbf{q}^s & \mathbf{p}^s \\ \mathbf{t}^f & \mathbf{q}^f & \mathbf{p}^f \end{bmatrix}, \quad \boldsymbol{\xi} \in \mathcal{G}. \quad (26)$$

Note that the kernel $\Sigma(\mathbf{y}, \boldsymbol{\xi})$ is derived from the poroelastodynamics fundamental solution of Appendix A which also appeared in the integral representation (15). In this setting, the smoothness of $\bar{\Sigma}$ substantiates the compactness of \mathcal{S}^* as an application from $\tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$ into $L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$. Next, suppose that there exists $\mathbf{a} \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$ such that $\mathcal{S}^*(\mathbf{a}) = \mathbf{0}$. Then, the vanishing trace of $(\mathbf{u}_{\mathbf{a}}, p_{\mathbf{a}})$ on \mathcal{G} implies, by the unique continuation principle, that $(\mathbf{u}_{\mathbf{a}}, p_{\mathbf{a}}) = \mathbf{0}$ in $\mathbb{R}^3 \setminus \bar{\Gamma}$, and whereby $([\![\mathbf{u}_{\mathbf{a}}]\!], [\![p_{\mathbf{a}}]\!], -[\![q_{\mathbf{a}}]\!]) = \mathbf{a} = \mathbf{0}$ which guarantees the injectivity of \mathcal{S}^* . The denseness of the range of \mathcal{S}^* is established via the injectivity of \mathcal{S} . To observe the latter, let $(\mathbf{g}_s, g_f) \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ such that $\mathcal{S}(\mathbf{g}_s, g_f) = \mathbf{0}$. The definition (19) then reads $(\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^t, p^t) = \mathbf{0}$ on Γ with (\mathbf{u}^t, p^t) satisfying (3). On recalling Assumption 1, let us consider $\Gamma_m \subset \Gamma$, $m = 1, \dots, M$, and its analytic continuation ∂D_m . Observe that $(\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^t, p^t) = \mathbf{0}$ also on Γ_m which thanks to the analyticity of $\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^t$ with respect to the surface coordinates on ∂D_m leads to $(\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^t, p^t) = \mathbf{0}$ on ∂D_m . One may show that $\omega \in \mathbb{R}$ may not be an eigenfrequency associated with

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}^m) - (\alpha - \frac{\rho_f}{\gamma}) \nabla p^m + \omega^2 (\rho - \frac{\rho_f^2}{\gamma}) \mathbf{u}^m &= \mathbf{0} & \text{in } D_m, \\ \frac{1}{\gamma \omega^2} \nabla^2 p^m + M^{-1} p^m + (\alpha - \frac{\rho_f}{\gamma}) \nabla \cdot \mathbf{u}^m &= \mathbf{0} & \text{in } D_m, \\ (\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^m, p^m) &= \mathbf{0} & \text{on } \partial D_m, \end{aligned} \quad (27)$$

when the relative pore flow

$$\mathbf{q}^m := \frac{1}{\gamma \omega^2} (\nabla p^m - \rho_f \omega^2 \mathbf{u}^m) \quad \text{in } D_m, \quad (28)$$

is nontrivial. This may be observed by (i) premultiplying the first of (27) by $\bar{\mathbf{u}}^m \in H^1(D_m)^3$, and (ii) multiplying the conjugate of its second by $-p^m \in H^1(D_m)$. Summation of the results and integration by parts over D_m yields

$$\omega^2 = \frac{\int_{D_m} \nabla \bar{\mathbf{u}}^m : \mathbf{C} : \nabla \mathbf{u}^m dV_{\boldsymbol{\xi}} + M^{-1} \|p^m\|_{L^2(D_m)}^2}{\int_{D_m} 2\rho_f \Re(\bar{\mathbf{u}}^m \cdot \mathbf{q}^m) dV_{\boldsymbol{\xi}} + \rho \|\mathbf{u}^m\|_{L^2(D_m)^3}^2 + \gamma \|\mathbf{q}^m\|_{L^2(D_m)^3}^2}. \quad (29)$$

Observe that the numerator and all terms, except $\gamma \|\mathbf{q}^m\|_{L^2(D_m)^3}^2$, in the denominator of (29) are real-valued. Therefore, $\omega \in \mathbb{C}$ unless $\|\mathbf{q}^m\|_{L^2(D_m)^3}^2 = 0$. In the latter case, signifying a purely elastic response in a bounded obstacle, one should further assume that ω is not an eigenfrequency affiliated with the Navier system (25) per Assumption 1. It should be mentioned that according to [42, Chapter 7, Theorem 1.4], the set of real eigenfrequencies characterizing (25) is discrete for a bounded D_m . In this setting, one may conclude that $(\mathbf{u}^t, p^t) = \mathbf{0}$ in D_m . The unique continuation principle then implies that $(\mathbf{u}^t, p^t) = \mathbf{0}$ in \mathbb{R}^3 and thus $(\mathbf{g}_s, g_f) = \mathbf{0}$ according to (3) which completes the proof. \blacksquare

Lemma 3.2. *Operator $T: H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \rightarrow \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)^3$ in (23) is bounded and satisfies*

$$\Im \langle \boldsymbol{\psi}, T\boldsymbol{\psi} \rangle_{\Gamma} < 0, \quad \forall \boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) : \boldsymbol{\psi} \neq \mathbf{0}. \quad (30)$$

Proof. The boundedness of T stems from the well-posedness of the forward problem (5) and trace theorems. To

observe (30), let $\boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ and consider (\mathbf{u}, p) satisfying (5) with $(\mathbf{t}^t, q^t, p^t) = \boldsymbol{\psi}$. Following similar steps used to obtain (B.5), one finds $\forall \boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$,

$$\Im \langle \boldsymbol{\psi}, T\boldsymbol{\psi} \rangle_\Gamma = -\frac{\omega}{\kappa} \|\mathbf{q}\|_{L^2(\mathbb{R}^3 \setminus \Gamma)}^2 + \Im \langle \mathfrak{P}(\llbracket \mathbf{u} \rrbracket, \llbracket p \rrbracket, -\llbracket q \rrbracket), (\llbracket \mathbf{u} \rrbracket, \llbracket p \rrbracket, -\llbracket q \rrbracket) \rangle_\Gamma. \quad (31)$$

Then, the Lemma's claim follows immediately from (13). \blacksquare

Lemma 3.3. *Operator $T: H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \rightarrow \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)^3$: (i) has a bounded (and thus continuous) inverse, and (ii) is coercive, i.e., there exists constant $c > 0$ independent of $\boldsymbol{\psi}$ such that*

$$|\langle \boldsymbol{\psi}, T(\boldsymbol{\psi}) \rangle| \geq c \|\boldsymbol{\psi}\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}^2, \quad \forall \boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma). \quad (32)$$

Proof. 1. First, observe that T given by (23) is Fredholm with index zero. The argument directly follows the proof of [43, Lemma 2.3] and makes use of the integral representation theorems and asymptotic behavior of the poroelastodynamics fundamental solution on Γ [41]. One then further demonstrate the injectivity of T by assuming that there exists $\boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ so that $T(\boldsymbol{\psi}) = \mathbf{0}$. Consequently, (15) reads that $(\mathbf{u}, p) = \mathbf{0}$ in $\mathbb{R}^3 \setminus \bar{\Gamma}$, which implies that $(\mathbf{t}, \langle q \rangle) = \mathbf{0}$ on Γ according to (6). Then, the third to fifth of (5) indicate that $\boldsymbol{\psi} = \mathbf{0}$, and thus the Lemma's claim follows.

2. We proceed using a contradiction argument. Assume that estimation (32) does not hold true, then one can find a sequence $(\boldsymbol{\psi}_n)_{n \in \mathbb{N}^*} \subset H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ such that for all $n \in \mathbb{N}^*$,

$$|\langle \boldsymbol{\psi}_n, T(\boldsymbol{\psi}_n) \rangle| < \frac{1}{n}, \quad \|\boldsymbol{\psi}_n\| = 1. \quad (33)$$

Since $(\boldsymbol{\psi}_n)$ is bounded, one may assume up to extracting a subsequence that $(\boldsymbol{\psi}_n)$ converges weakly to some $\boldsymbol{\psi} \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$. Now, let $(\mathbf{u}_n, p_n) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ solve the poroelastic scattering problem (5) for $(\mathbf{t}^t, q^t, p^t) = \boldsymbol{\psi}_n$. In this setting, (31) reads that for $\boldsymbol{\psi} = \boldsymbol{\psi}_n$,

$$|\langle \boldsymbol{\psi}_n, T(\boldsymbol{\psi}_n) \rangle| \geq \frac{\omega}{\kappa} \|\mathbf{q}_n\|_{L^2(\mathbb{R}^3 \setminus \bar{\Gamma})}^2 - \Im \langle \mathfrak{P}(\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket), (\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket) \rangle_\Gamma, \quad (34)$$

wherein

$$\begin{aligned} \mathbf{q}_n &:= \frac{1}{\gamma \omega^2} (\nabla p_n - \rho_f \omega^2 \mathbf{u}_n) && \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ q_n &:= \frac{1}{\gamma \omega^2} (\nabla p_n \cdot \mathbf{n} - \rho_f \omega^2 \mathbf{u}_n \cdot \mathbf{n}) && \text{on } \Gamma. \end{aligned} \quad (35)$$

In light of (13), (33) and (34), observe that \mathbf{q}_n converges strongly to zero in $L^2(\mathbb{R}^3 \setminus \bar{\Gamma})^3$. Next, from the wellposedness of the scattering problem and the continuity of the trace operator, one finds that

$$\begin{aligned} \|(\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket)\|_{\tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)} &\leq c \|(\mathbf{u}_n, p_n)\|_{H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})} \\ &\leq c' \|\boldsymbol{\psi}_n\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)} = c', \end{aligned} \quad (36)$$

for a constant $c > 0$ (resp. $c' > 0$) independent of (\mathbf{u}_n, p_n) (resp. $\boldsymbol{\psi}_n$). The boundedness of sequences $\llbracket \mathbf{u}_n \rrbracket$, $\llbracket p_n \rrbracket$ and $\llbracket q_n \rrbracket$ indicates, up to extracting a subsequence, that they respectively converge weakly to some $\llbracket \mathbf{u} \rrbracket$, $\llbracket p \rrbracket$ and $\llbracket q \rrbracket$. Then, the compactness of \mathcal{S}^* , mapping $(\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket)$ to (\mathbf{u}_n, p_n) according to (15), implies that (\mathbf{u}_n, p_n) converges strongly to some $(\mathbf{u}, p) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$. From (5)

and (9), one may show that \mathbf{u}_n satisfies

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}_n) + (\rho - \alpha \rho_f) \omega^2 \mathbf{u}_n &= (\alpha \gamma - \rho_f) \omega^2 \mathbf{q}_n && \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ \frac{\partial \mathbf{u}_\varsigma}{\partial r} - i k_\varsigma \mathbf{u}_\varsigma &= o(r^{-1} e^{-\Im(k_\varsigma) r}), \quad \varsigma = p_1, p_2, s, \quad \mathbf{u} = \mathbf{u}_n && \text{as } r := |\boldsymbol{\xi}| \rightarrow \infty. \end{aligned} \quad (37)$$

As $n \rightarrow \infty$, one obtains that \mathbf{u} satisfies

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}) + (\rho - \alpha \rho_f) \omega^2 \mathbf{u} &= \mathbf{0} && \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ \frac{\partial \mathbf{u}_\varsigma}{\partial r} - i k_\varsigma \mathbf{u}_\varsigma &= o(r^{-1} e^{-\Im(k_\varsigma) r}), \quad \varsigma = p_1, p_2, s, \quad \text{as } r := |\boldsymbol{\xi}| \rightarrow \infty, \end{aligned} \quad (38)$$

implying that $\mathbf{u} = \mathbf{0}$. Note that the conservative Navier system in first of (38) does not generate the exponentially decaying radiation pattern in the second of (38) (which is associated with the lossy Biot system). In this setting, (35) indicates that p is constant in $\mathbb{R}^3 \setminus \bar{\Gamma}$ and since p is subject to a radiation condition similar to the second of (38), one deduces that $p = 0$. Now, from the wellposedness of the forward scattering problem, one finds that $\boldsymbol{\psi} = 0$. On the other hand, observe from (12) that

$$\|\boldsymbol{\psi}_n\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}^2 = \langle \boldsymbol{\psi}_n, \mathfrak{P}(\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket) \rangle_\Gamma - \langle \boldsymbol{\psi}_n, (\mathbf{t}_n, \langle q_n \rangle, \langle p_n \rangle) \rangle_\Gamma. \quad (39)$$

Then, the regularity of the trace operator implies that $(\llbracket \mathbf{u}_n \rrbracket, \llbracket p_n \rrbracket, -\llbracket q_n \rrbracket)$ and $(\mathbf{t}_n, \langle q_n \rangle, \langle p_n \rangle)$ converges strongly to zero. Consequently, (39) reads that $\|\boldsymbol{\psi}_n\|^2$ goes to zero as n goes to infinity. This result contradicts (33) and establishes the lemma's statement. ■

Lemma 3.4. *Operator $\mathcal{P} = \bar{\mathcal{S}}^* T : H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is compact over $H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$.*

Proof. The claim follows immediately from Lemmas 3.1 and 3.2 establishing, respectively, the compactness of \mathcal{S}^* and the boundedness of T . ■

Lemma 3.5. *The poroelastic scattering operator $\Lambda : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is injective, compact and has a dense range.*

Proof. The injectivity of Λ is implied by the injectivity of operators \mathcal{S}^* , T , and \mathcal{S} in the factorization $\Lambda = \bar{\mathcal{S}}^* T \mathcal{S}$ (see Lemmas 3.1 and 3.3). The compactness of Λ follows immediately from the compactness of \mathcal{S}^* – and thus that of \mathcal{S} (Lemma 3.1), and the boundedness of T (Lemma 3.2). The last claim may be verified by establishing the injectivity of $\Lambda^* = \mathcal{S}^* \mathcal{P}^*$ which in light of Lemma 3.1, results from the injectivity of \mathcal{P}^* . To demonstrate the latter, one first finds from the definition of \mathcal{P} that the adjoint operator $\mathcal{P}^* : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$ is given by $\mathcal{P}^*(\mathbf{g}) = (\llbracket \mathbf{u}^* \rrbracket, \llbracket p^* \rrbracket, -\llbracket q^* \rrbracket)$ on Γ where $(\mathbf{u}^*, p^*) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\bar{\Gamma} \cup \bar{\mathcal{G}}\})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\bar{\Gamma} \cup \bar{\mathcal{G}}\})$ solves

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}^*)(\boldsymbol{\xi}) - (\alpha - \frac{\rho_f}{\bar{\gamma}}) \nabla p^*(\boldsymbol{\xi}) + \omega^2 (\rho - \frac{\rho_f^2}{\bar{\gamma}}) \mathbf{u}^*(\boldsymbol{\xi}) &= - \int_{\mathcal{G}} \delta(\mathbf{y} - \boldsymbol{\xi}) \mathbf{g}_s(\mathbf{y}) dS_{\mathbf{y}} && \text{in } \mathbb{R}^3 \setminus \{\bar{\Gamma} \cup \bar{\mathcal{G}}\}, \\ \frac{1}{\bar{\gamma} \omega^2} \nabla^2 p^*(\boldsymbol{\xi}) + M^{-1} p^*(\boldsymbol{\xi}) + (\alpha - \frac{\rho_f}{\bar{\gamma}}) \nabla \cdot \mathbf{u}^*(\boldsymbol{\xi}) &= - \int_{\mathcal{G}} \delta(\mathbf{y} - \boldsymbol{\xi}) g_f(\mathbf{y}) dS_{\mathbf{y}} && \text{in } \mathbb{R}^3 \setminus \{\bar{\Gamma} \cup \bar{\mathcal{G}}\}, \end{aligned} \quad (40)$$

$$(\mathbf{t}^*, \langle\langle q^* \rangle\rangle, \langle\langle p^* \rangle\rangle) = \bar{\mathfrak{P}}^T(\llbracket \mathbf{u}^* \rrbracket, \llbracket p^* \rrbracket, -\llbracket q^* \rrbracket), \quad \llbracket \mathbf{t}^* \rrbracket = \mathbf{0} \quad \text{on } \Gamma,$$

wherein ‘T’ indicates transpose, and

$$[\mathbf{t}^*, q^*](\mathbf{u}^*, p^*) := [\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^* - \alpha p^* \mathbf{n}, \frac{1}{\bar{\gamma} \omega^2} (\nabla p^* \cdot \mathbf{n} - \rho_f \omega^2 \mathbf{u}^* \cdot \mathbf{n})] \quad \text{on } \Gamma.$$

Note that (\mathbf{u}^*, p^*) satisfy the radiation condition as $|\boldsymbol{\xi}| \rightarrow \infty$, similar to the complex conjugate of (9). It should be mentioned that \mathcal{P}^* is obtained by applying the poroelastodynamics reciprocal theorem between (\mathbf{u}^*, p^*) of (40) and (\mathbf{u}, p) satisfying (5) similar to the argument in proof of Lemma 2.1. Now, let $\mathbf{g} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$, and assume that $\mathcal{P}^*(\mathbf{g}) = \mathbf{0}$, i.e., $(\llbracket \mathbf{u}^* \rrbracket, \llbracket p^* \rrbracket, -\llbracket q^* \rrbracket) = \mathbf{0}$. In this case, (\mathbf{u}^*, p^*) and q^* are continuous across Γ independent of $\bar{\mathfrak{P}}^T$. Therefore, $(\mathbf{u}^*, p^*) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})$ and $(\mathbf{t}^*, q^*, p^*) = \mathbf{0}$. Thanks to this result, one can show as in the proof of Lemma 3.1 that $(\mathbf{u}^*, p^*) = \mathbf{0}$, and consequently $\mathbf{g} = \mathbf{0}$ which concludes the proof. \blacksquare

Remark 3.1. *It is worth noting that almost all of the analysis in this section is conducted via the variational formulation. Related arguments in the case of elasticity [44] mostly rely on the boundary integral formulations – to take advantage of symmetries in the far-field patterns of elastodynamic fundamental solution. The latter is a convenient feature which no longer exists in the case of poroelastic scattering.*

4. Inverse poroelastic scattering

This section presents an adaptation of the key results from sampling approaches to inverse scattering for the problem of poroelastic-wave imaging of finitely permeable interfaces. The fundamental idea stems from the nature of solution $\mathbf{g} = (\mathbf{g}_s, g_f) \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ to the poroelastic scattering equation

$$\Lambda \mathbf{g} = \Phi_L, \quad \Lambda = \mathcal{P}\mathcal{S} = \bar{\mathcal{S}}^* T \mathcal{S}, \quad (41)$$

where Φ_L is the near-field pattern of a trial poroelastodynamic field specified by Definition 1.

Definition 1. *With reference to (15), for every admissible density $\mathbf{a} \in \tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L)$ specified over a smooth, non-intersecting trial interface $L \subset \mathcal{B} \subset \mathbb{R}^3$, the induced near-field pattern $\Phi_L: \tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is given by*

$$\Phi_L(\mathbf{a})(\boldsymbol{\xi}) := (\mathbf{u}_L, p_L)(\boldsymbol{\xi}) = \int_L \Sigma(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{a}(\mathbf{y}) \, dS_{\mathbf{y}}, \quad \Sigma = \begin{bmatrix} \mathbf{T}^s & \mathbf{q}^s & \mathbf{p}^s \\ \mathbf{t}^f & \mathbf{q}^f & \mathbf{p}^f \end{bmatrix}, \quad \boldsymbol{\xi} \in \mathcal{G}. \quad (42)$$

Remark 4.1. *Let $L \subset \mathcal{B}$ (containing the origin) denote a poroelastic discontinuity that is much smaller than ℓ_r in (1), and let $L = \mathbf{z} + \mathbf{R}L$ where $\mathbf{z} \in \mathcal{B}$ and $\mathbf{R} \in U(3)$ is a unitary rotation matrix. Given an admissible distribution $\mathbf{a} \in \tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L)$, solving (41) over (\mathbf{z}, \mathbf{R}) is simply an effort to probe the range of operator Λ (or that of $\bar{\mathcal{S}}^*$), through synthetic reshaping of the illuminating wavefront, for fingerprints in terms of Φ_L . As shown by Theorems 4.2, the latter is possible if and only if $L \subset \Gamma$. Otherwise, the norm of any approximate solution to (41) can be made arbitrarily large, which provides a criterion for reconstructing Γ .*

Theorem 4.1. *Under the assumptions of Remark 2.3 for the wellposedness of the forward scattering problem, for every smooth and non-intersecting trial dislocation $L \subset \mathcal{B} \subset \mathbb{R}^3$ and density profile $\mathbf{a}(\boldsymbol{\xi}) \in \tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L)$, one has*

$$\Phi_L \in \text{Range}(\bar{\mathcal{S}}^*) \iff L \subset \Gamma.$$

Proof. Consider the following:

- If $L \subset \Gamma$, then $\tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L) \subset \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$. By extending the domain of $\mathbf{a}(\boldsymbol{\xi})$ from L to Γ through zero padding, one immediately obtains $\Phi_L \in \text{Range}(\bar{\mathcal{S}}^*)$ thanks to (26) and (42).
- Assume that $L \not\subset \Gamma$ and that $\Phi_L \in \text{Range}(\bar{\mathcal{S}}^*)$. Then there exists $\mathbf{b} \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$ such that

$$\Phi_L(\boldsymbol{\xi}) = \bar{\mathcal{S}}^*(\mathbf{b})(\boldsymbol{\xi}), \quad \boldsymbol{\xi} \in \mathcal{G}. \quad (43)$$

According to (26), the right-hand side of (43) is associated with the layer potential

$$\Phi_\Gamma(\boldsymbol{\xi}) = \int_\Gamma \Sigma(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{b}(\mathbf{y}) \, dS_{\mathbf{y}}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \bar{\Gamma}. \quad (44)$$

which, by definition (42), should coincide with

$$\Psi_L(\mathbf{a})(\boldsymbol{\xi}) = \int_L \Sigma(\mathbf{y}, \boldsymbol{\xi}) \cdot \mathbf{a}(\mathbf{y}) \, dS_{\mathbf{y}}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \overline{\Gamma \cup L}. \quad (45)$$

Now, let $\Gamma \not\ni \boldsymbol{\xi}^\circ \in L$ and let \mathcal{B}_ϵ be a small ball centered at $\boldsymbol{\xi}^\circ$ such that $\mathcal{B}_\epsilon \cap \Gamma = \emptyset$. In this case, Φ_Γ is analytic in \mathcal{B}_ϵ , while Ψ_L has a discontinuity across $\mathcal{B}_\epsilon \cap L$, which by contradiction completes the proof. ■

On the basis of Theorem 4.1, one arrives at the following statement:

Theorem 4.2. *Under the assumptions of Theorem 4.1,*

- If $L \subset \Gamma$, there exists a density vector $\mathbf{g}_\epsilon^L \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ such that $\|\Lambda \mathbf{g}_\epsilon^L - \Phi_L\|_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})} \leq \epsilon$ and $\limsup_{\epsilon \rightarrow 0} \|\mathcal{S} \mathbf{g}_\epsilon^L\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)} < \infty$.
- If $L \not\subset \Gamma$, then $\forall \mathbf{g}_\epsilon^L \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ such that $\|\Lambda \mathbf{g}_\epsilon^L - \Phi_L\|_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})} \leq \epsilon$, one has $\lim_{\epsilon \rightarrow 0} \|\mathcal{S} \mathbf{g}_\epsilon^L\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)} = \infty$.

Proof. Let $L \subset \Gamma$, i.e., $\Phi_L \in \text{Range}(\bar{\mathcal{S}}^*)$ by Theorem 4.1, then there exists $\mathbf{a}^L \in \overline{\text{Range}(T)}$ such that $\bar{\mathcal{S}}^* \mathbf{a}^L = \Phi_L$. On invoking (a) continuity of T^{-1} as per Lemma 3.3, and (b) denseness of the range \mathcal{S} by the injectivity of $\bar{\mathcal{S}}^*$ according to Lemma 3.1, one finds that $\forall \epsilon > 0$, $\exists \mathbf{g}_\epsilon^L \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ such that $\|T^{-1} \mathbf{a}^L - \mathcal{S} \mathbf{g}_\epsilon^L\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)} \leq \epsilon$. This establishes the first part of the claim owing to the continuity of $\bar{\mathcal{S}}^* T$ per Lemma 3.4, and the fact that $\mathbf{a}^L \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$.

Next, consider the case where $L \not\subset \Gamma$, and consequently, $\Phi_L \notin \text{Range}(\bar{\mathcal{S}}^*)$ by Theorem 4.1. Then, Lemma 3.1 implies that for every $\epsilon > 0$, and some regularization parameter $\alpha = \alpha(\epsilon)$ ($\lim_{\epsilon \rightarrow 0} \alpha = 0$), a nearby solution $\mathbf{a}_\epsilon^L \in \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)$ can be built, e.g., via the Tikhonov regularization [45], such that $\|\Phi_L - \bar{\mathcal{S}}^* \mathbf{a}_\epsilon^L\|_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})} \leq \epsilon$, and $\lim_{\epsilon \rightarrow 0} \|\mathbf{a}_\epsilon^L\|_{\tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma)} = \infty$ due to the compactness of $\bar{\mathcal{S}}^*$. At this point, the same argument as in the first part of the proof – by deploying the continuity of T^{-1} and the range denseness of \mathcal{S} – can be used to establish the second claim. ■

4.1. The multiphysics \mathfrak{L} indicator

Theorem 4.2 poses two challenges in that: (i) the featured indicator $\|\mathcal{S}\mathbf{g}_\epsilon^L\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}$ inherently depends on the unknown support Γ of hidden scatterers, and (ii) construction of the wavefront density $\mathbf{g}_\epsilon^L \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is implicit in the theorem [20, 26]. These are conventionally addressed by replacing $\|\mathcal{S}\mathbf{g}_\epsilon^L\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}$ with $\|\mathbf{g}_\epsilon^L\|_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})}$ which in turn is computed by way of Tikhonov regularization [45] as follows.

First, note that when the measurements are contaminated by noise (e.g., sensing errors and fluctuations in the medium properties), one has to deal with the noisy operator Λ^δ satisfying

$$\|\Lambda^\delta - \Lambda\| \leq \delta, \quad (46)$$

where $\delta > 0$ is a measure of perturbation in data independent of Λ . Assuming that Λ^δ is compact, the Tikhonov-regularized solution \mathbf{g}_η to (41) is obtained by non-iteratively minimizing the cost functional $J_\eta(\Phi_L; \cdot) : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow \mathbb{R}$ defined by

$$J_\eta(\Phi_L; \mathbf{g}) := \|\Lambda^\delta \mathbf{g} - \Phi_L\|_{L^2}^2 + \eta \|\mathbf{g}\|_{L^2}^2, \quad \mathbf{g} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G}), \quad (47)$$

where the regularization parameter $\eta = \eta(\delta, L)$ is obtained by way of the Morozov discrepancy principle [45]. On the basis of (47), the well-known linear sampling indicator \mathfrak{L} is constructed as

$$\mathfrak{L} := \frac{1}{\|\mathbf{g}_\eta\|_{L^2}}, \quad \mathbf{g}_\eta := \min_{\mathbf{g} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})} J_\eta(\Phi_L; \mathbf{g}). \quad (48)$$

By design, \mathfrak{L} achieves its highest values at the loci of hidden scatterers Γ . More specifically, the behavior of \mathfrak{L} within \mathcal{B} may be characterized as the following,

$$\begin{aligned} \text{if } L \subset \Gamma & \iff \liminf_{\eta \rightarrow 0} \mathfrak{L}(\mathbf{g}_\eta) > 0, \\ \text{if } L \subset \mathcal{B} \setminus \bar{\Gamma} & \iff \lim_{\eta \rightarrow 0} \mathfrak{L}(\mathbf{g}_\eta) = 0. \end{aligned} \quad (49)$$

4.2. The modified \mathfrak{G} indicator

In cases where the illumination frequency and/or the background's permeability is sufficiently large so that $\Im\gamma \rightarrow 0$ according to (4), one may deduce from (22)-(24) that the second factorization of the poroelastic scattering operator is symmetrized as follows,

$$\Lambda = \mathcal{S}^* T \mathcal{S}. \quad (50)$$

In this setting, one may deploy the more rigorous \mathfrak{G} indicator [44] to dispense with the approximations underlying the \mathfrak{L} imaging functional. This results in a more robust reconstruction especially with noisy data [26]. The \mathfrak{G} indicator takes advantage of the positive and self-adjoint operator $\Lambda_\sharp : L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ defined on the basis of the scattering operator Λ by

$$\Lambda_\sharp := \frac{1}{2} |\Lambda + \Lambda^*| + \left| \frac{1}{2i} (\Lambda - \Lambda^*) \right|, \quad (51)$$

with the affiliated factorization [46]

$$\Lambda_\sharp = \mathcal{S}^* T_\sharp \mathcal{S}, \quad (52)$$

where in light of Lemma 3.3, the middle operator T_\sharp is coercive with reference to [44, Lemma 5.7] i.e., there exists a constant $c > 0$ independent of $\Psi = \mathcal{S}\mathbf{g}$ such that $\forall \Psi \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$,

$$(\mathbf{g}, \Lambda_{\#} \mathbf{g})_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})} = \langle \Psi, T_{\#} \Psi \rangle_{\Gamma} \geq c \|\Psi\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}^2. \quad (53)$$

Thanks to (53), the term $\|\mathcal{S} \mathbf{g}_{\epsilon}\|_{H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)}^2$ in Theorem 4.2 may be safely replaced by $(\mathbf{g}_{\epsilon}, \Lambda_{\#} \mathbf{g}_{\epsilon})_{L^2(\mathcal{G})^3 \times L^2(\mathcal{G})}$ which is computable without prior knowledge of Γ .

In presence of noise, the perturbed operator $\Lambda_{\#}^{\delta}$ is deployed satisfying

$$\|\Lambda_{\#}^{\delta} - \Lambda_{\#}\| \leq \delta, \quad (54)$$

which is assumed to be compact similar to Λ^{δ} in (46). Then, the regularized cost functional $J_{\alpha, \delta}(\Phi_L; \cdot): L^2(\mathcal{G})^3 \times L^2(\mathcal{G}) \rightarrow \mathbb{R}$ is constructed according to [23, Theorems 4.3],

$$J_{\alpha, \delta}(\Phi_L; \mathbf{g}) := \|\Lambda^{\delta} \mathbf{g} - \Phi_L\|_{L^2}^2 + \alpha(\mathbf{g}, \Lambda_{\#}^{\delta} \mathbf{g})_{L^2} + \alpha \delta \|\mathbf{g}\|_{L^2}^2, \quad \mathbf{g} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G}). \quad (55)$$

Here, $\alpha > 0$ represents the regularization parameter specified in terms of η from (47) as

$$\alpha(\delta, L) := \frac{\eta(\delta, L)}{\|\Lambda\|_{L^2} + \delta}. \quad (56)$$

In addition, $\delta > 0$ signifies both a measure of perturbation in data and a regularization parameter that, along with α , is designed to create a robust imaging indicator as per Theorems 4.3. It should be mentioned that $J_{\alpha, \delta}$ is convex and that its minimizer $\mathbf{g}_{\alpha, \delta} \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ solves the linear system

$$\Lambda^{\delta*}(\Lambda^{\delta} \mathbf{g}_{\alpha, \delta} - \Phi_L) + \alpha((\Lambda_{\#}^{\delta})^{\frac{1}{2}*}(\Lambda_{\#}^{\delta})^{\frac{1}{2}} \mathbf{g}_{\alpha, \delta} + \delta \mathbf{g}_{\alpha, \delta}) = \mathbf{0}, \quad (57)$$

which can be computed without iterations [23]. Within this framework, the following theorem rigorously establishes the relation between the range of operator \mathcal{S}^* and the norm of penalty term in (55).

Theorem 4.3. *Under the assumptions of Theorem 4.1 and additional hypothesis that Λ^{δ} and $\Lambda_{\#}^{\delta}$ are compact,*

$$\begin{aligned} \Phi_L \in \text{Range}(\mathcal{S}^*) &\iff \left\{ \limsup_{\alpha \rightarrow 0} \limsup_{\delta \rightarrow 0} (|\langle \mathbf{g}_{\alpha, \delta}, \Lambda_{\#}^{\delta} \mathbf{g}_{\alpha, \delta} \rangle| + \delta \|\mathbf{g}_{\alpha, \delta}\|^2) < \infty \right. \\ &\iff \left. \liminf_{\alpha \rightarrow 0} \liminf_{\delta \rightarrow 0} (|\langle \mathbf{g}_{\alpha, \delta}, \Lambda_{\#}^{\delta} \mathbf{g}_{\alpha, \delta} \rangle| + \delta \|\mathbf{g}_{\alpha, \delta}\|^2) < \infty \right\}, \end{aligned}$$

where $\mathbf{g}_{\alpha, \delta}$ is a minimizer of the perturbed cost functional (55) in the sense of [23, Lemma 6.8].

In this setting, the imaging indicator \mathfrak{G} takes the form

$$\mathfrak{G} = \frac{1}{\sqrt{\|(\Lambda_{\#}^{\delta})^{\frac{1}{2}} \mathbf{g}_{\alpha, \delta}\|^2 + \delta \|\mathbf{g}_{\alpha, \delta}\|^2}}, \quad (58)$$

with similar characteristics to \mathfrak{L} as in (49) yet more robustness against noise [26].

Remark 4.2. *It is worth noting that the sampling-based approaches to inverse scattering exhibit minor sensitivity to the interfacial condition – described according to (5) by the heterogeneous distribution of: the stiffness matrix $\mathbf{K}(\boldsymbol{\xi})$, permeability coefficient $\kappa_f(\boldsymbol{\xi})$, effective stress coefficient $\alpha_f(\boldsymbol{\xi})$, Skempton coefficient $\beta_f(\boldsymbol{\xi})$, and fluid pressure dissipation factor $\Pi(\boldsymbol{\xi})$ on $\boldsymbol{\xi} \in \Gamma$. This attribute may be traced back to Remark 4.1 where the density \mathbf{a} is deemed an arbitrary distribution in $\tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L)$. This attribute makes the proposed method particularly attractive in situations where the interfacial parameters are unknown a priori, which calls for hybrid inverse solutions for geometrical reconstruction and interfacial characterization of such features e.g., see [47, 48].*

Remark 4.3. *It should be mentioned that there are recent efforts to systematically adapt the \mathfrak{G} indicator for application to asymmetric scattering operators [21]. These developments, however, do not lend themselves to poroelastic inverse scattering due to the non-selfadjoint nature of Λ . A fundamental treatment of the poroelastic \mathfrak{G} indicator in the general case where $\gamma \in \mathbb{C}$ could be an interesting subject for a future study.*

5. Computational treatment and results

This section examines the performance of multiphasic indicators (48) and (58) through a set of numerical experiments. In what follows, the synthetic data are generated within the FreeFem++ computational platform [49].

5.1. Testing configuration

Two test setups shown in Figs. 2 and 3 are considered where the background is a poroelastic slab P of dimensions 22.3×22.3 endowed with (evolving and stationary) hydraulic fracture networks. Following [32, 50], the properties of Pecos sandstone are used to characterize P . On setting the reference scales $\mu_r = 5.85$ GPA, $\rho_r = 10^3$ kg/m³, and $\ell_r = 0.14$ m for stress, mass density, and length, respectively, the non-dimensionalized quantities of Table 1 are obtained and used for simulations. Accordingly, the complex-valued wave speeds [16] affiliated with the modal decomposition in (8) read $c_s = 0.66 + 8.8 \times 10^{-6}i$, $c_{p_1} = 1.26 + 3 \times 10^{-7}i$, and $c_{p_2} = 5.8 \times 10^{-3} + 5.8 \times 10^{-3}i$. The boundary condition on ∂P is such that the total traction and pore pressure vanish for both incident and scattered fields i.e., $(\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}^i, p^i) = (\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}, p) = \mathbf{0}$. In *Setup I*, a hydraulic fracture network $\Gamma_1 - \Gamma_9$ is induced in four steps as shown in Fig. 2. A detailed description of scatterers including the center (x_c, y_c) , length ℓ , and orientation ϕ of each crack Γ_κ , $\kappa = \{1, 2, \dots, 9\}$ is provided in Table 2. All fractures in this configuration are highly permeable as per Remark 8 and characterized by the interfacial stiffness $\mathbf{K}(\boldsymbol{\xi}) = \mathbf{0}$, effective stress coefficient $\alpha_f(\boldsymbol{\xi}) = 0.85$, and Skempton coefficient $\beta_f(\boldsymbol{\xi}) = 0.3$ on $\boldsymbol{\xi} \in \bigcup_{\kappa=1}^9 \Gamma_\kappa$. The latter quantities are taken from [35]. Note that in *Setup I*, the excitation and sensing grid \mathcal{G} straddles two (vertical) monitoring wells and the horizontal section of the injection well. Depicted in Figs. 3 and 4, *Setup II* features hydraulic fractures $\Gamma_{10} - \Gamma_{15}$ of distinct length scales as described in Table 3. The discontinuity interfaces in this configuration are modeled as thin poroelastic inclusions characterized by $\lambda_f = 0.1$, $\mu_f = 0.2$, $M_f = 0.33$, $\kappa_f = 5 \times 10^{-7}$, and $\phi_f = 0.35$, while the remaining material parameters are similar to their counterparts in the background as reported in Table 1. In *Setup II*, the excitation and measurements are solely conducted in the treatment well shown in Fig. 3.

Table 1: Poroelastic properties of the background.

| property | value | dimensionless value |
|-----------------------------------|------------------------------------|--------------------------------|
| first Lamé parameter (drained) | $\lambda' = 2.74$ GPA | $\lambda = 0.47$ |
| drained shear modulus | $\mu' = 5.85$ GPA | $\mu = 1$ |
| Biot modulus | $M' = 9.71$ GPA | $M = 1.66$ |
| total density | $\rho' = 2270$ kg/m ³ | $\rho = 2.27$ |
| fluid density | $\rho'_f = 1000$ kg/m ³ | $\rho_f = 1$ |
| apparent mass density | $\rho'_a = 117$ kg/m ³ | $\rho_a = 0.117$ |
| permeability coefficient | $\kappa' = 0.8$ mm ⁴ /N | $\kappa = 24.5 \times 10^{-7}$ |
| porosity | $\phi = 0.195$ | $\phi = 0.195$ |
| Biot effective stress coefficient | $\alpha = 0.83$ | $\alpha = 0.83$ |
| frequency | $\omega' = 12$ kHz | $\omega = 3.91$ |

Table 2: Process zone configuration illustrated in Fig. 2: center (x_c, y_c) , length ℓ , and orientation ϕ (with respect to x axis) of cracks Γ_κ , $\kappa = \{1, 2, \dots, 9\}$.

| κ | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
|-----------------------|-----------|----------|-----------|-----------|-----------|----------|-----------|----------|-----------|
| $x_c(\Gamma_\kappa)$ | -5.5 | -0.25 | 4.3 | -3.3 | 1.6 | -4.3 | -3.4 | -2 | 2.9 |
| $y_c(\Gamma_\kappa)$ | 0 | 0 | -1 | 0 | 0.5 | 0.5 | -1.07 | 0 | 0 |
| $\ell(\Gamma_\kappa)$ | 3 | 2.2 | 3 | 1 | 2.75 | 1.8 | 1.25 | 2.4 | 2.2 |
| $\phi(\Gamma_\kappa)$ | 0.47π | 0.6π | 0.56π | 0.56π | 0.42π | 0.5π | 0.37π | 0.5π | 0.56π |

Table 3: Process zone configuration illustrated in Figs. 3, 4: center (x_c, y_c) , length ℓ , and orientation ϕ (with respect to x axis) of cracks Γ_κ , $\kappa = \{10, 11, \dots, 15\}$.

| κ | 10 | 11 | 12 | 13 | 14 | 15 |
|-----------------------|----------|----------|-----------|----------|----------|-----------|
| $x_c(\Gamma_\kappa)$ | -4.76 | -1.13 | 3.46 | -3.06 | -1.13 | 1.16 |
| $y_c(\Gamma_\kappa)$ | 0 | 0 | 0 | 0 | 0 | 0 |
| $\ell(\Gamma_\kappa)$ | 2.35 | 0.74 | 7.03 | 2.35 | 0.74 | 7.03 |
| $\phi(\Gamma_\kappa)$ | 0.3π | 0.5π | 0.46π | 0.3π | 0.5π | 0.46π |

5.2. Forward scattering simulations

Every sensing step entails in-plane harmonic excitation via total body forces \mathbf{g}_s and fluid volumetric sources g_f at a set of points on \mathcal{G} . The excitation frequency $\omega = 3.91$ is set such that the induced shear wavelength λ_s in the specimen is approximately one, serving as a reference length scale. The incident waves interact with the hydraulic fracture network in each setup giving rise to the scattered field (\mathbf{u}, p) , governed by (5), whose pattern $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$ over the observation grid \mathcal{G} is then computed. For both illumination and sensing purposes, \mathcal{G} is sampled by a uniform grid of N excitation and observation points. In *Setup I*, the H-shaped incident/observation grid is comprised of $N = 330$ source/receiver points, while in *Setup II*, the L-shaped support of injection well is uniformly sampled at $N = 130$ points for excitation and sensing.

5.3. Data Inversion

With the preceding data, one may generate the multiphase indicator maps affiliated with (48) and (58) in three steps, namely by: (1) constructing the discrete scattering operators $\mathbf{\Lambda}$ and $\mathbf{\Lambda}^\delta$ from synthetic data $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$, (2) computing the trial signature patterns Φ_L pertinent to a poroelastic host domain, and (3) evaluating the imaging indicators (\mathfrak{L} or \mathfrak{G}) in the sampling area through careful minimization of the discretized cost functionals (47) and (55) as elucidated in the sequel.

Step 1: construction of the multiphase scattering operator

Given the in-plane nature of wave motion, i.e., that the polarization amplitude of excitation \mathbf{g}_s and the nontrivial components of associated scattered fields \mathbf{u}^{obs} lay in the $x - y$ plane of orthonormal bases $(\mathbf{e}_1, \mathbf{e}_2)$, the discretized scattering operator $\mathbf{\Lambda}$ may be represented by a $3N \times 3N$ matrix with components

$$\mathbf{\Lambda}(4i+1:4i+4, 4j+1:4j+4) = \begin{bmatrix} u_{11} & u_{12} & | & \mathbf{u}_1 \\ u_{21} & u_{22} & | & \mathbf{u}_2 \\ \hline p_1 & p_2 & | & \mathbf{p} \end{bmatrix} (\boldsymbol{\xi}_i, \mathbf{y}_j), \quad i, j = 0, \dots, N-1, \quad (59)$$

where $\Lambda_{rs}(\boldsymbol{\xi}_i, \mathbf{y}_j)$ ($r, s=1, 2$) is the r^{th} component of the solid displacement measured at $\boldsymbol{\xi}_i$ due to a unit force applied at \mathbf{y}_j along the coordinate direction s and p_s is the associated pore pressure measurement; also, \mathbf{u} signifies the 2×1 solid displacement at $\boldsymbol{\xi}_i$ due to a unit injection at \mathbf{y}_j and \mathbf{p} is the affiliated pore pressure.

Note that here the general case is presented where excitations and measurements are conducted along all dimensions. $\mathbf{\Lambda}$ may be downscaled as appropriate according to the testing setup.

Noisy operator. To account for the presence of noise in measurements, we consider the perturbed operator

$$\mathbf{\Lambda}^\delta := (\mathbf{I} + \mathbf{N}_\epsilon) \mathbf{\Lambda}, \quad (60)$$

where \mathbf{I} is the $3N \times 3N$ identity matrix, and \mathbf{N}_ϵ is the noise matrix of commensurate dimension whose components are uniformly-distributed (complex) random variables in $[-\epsilon, \epsilon]^2$. In what follows, the measure of noise in data with reference to definition (46) is $\delta = \|\mathbf{N}_\epsilon \mathbf{\Lambda}\| = 0.05 \|\mathbf{\Lambda}\|$.

Remark 5.1. *In practice, $\delta > 0$ may be specified via sensor calibration and/or signal processing. In [25], for instance, $\delta = \|\mathbf{\Lambda}^n - \mathbf{\Lambda}^p\|_{L^2}$ wherein $\mathbf{\Lambda}^n$ is the noisy scattering operator constructed using “raw” i.e., unprocessed measurement data, whereas $\mathbf{\Lambda}^p$ is computed from the denoised i.e., processed signals.*

Step 2: generation of a multiphysics library of scattering patterns

This step aims to construct a suitable right hand side for the discretized scattering equation in unbounded and bounded domains pertinent to the analytical developments of Section 4 and numerical experiments of this section, respectively.

Unbounded domain in \mathbb{R}^3 . In this case, the poroelastic trial pattern $\Phi_L \in L^2(\mathcal{G})^3 \times L^2(\mathcal{G})$ is given by Definition 1 indicating that (a) the right hand side is not only a function of the dislocation geometry L but also a function of the trial density $\mathbf{a} \in \tilde{H}^{1/2}(L)^3 \times \tilde{H}^{1/2}(L) \times \tilde{H}^{-1/2}(L)$, and (b) computing Φ_L generally requires an integration process at every sampling point \mathbf{x}_o . Conventionally, one may dispense with the integration process by considering a sufficiently localized (trial) density function e.g., see [44, 23].

Bounded domain. This case corresponds to the numerical experiments of this section where the background is a poroelastic plate P of finite dimensions. In this setting, following [51], it is straightforward to show that the trial patterns $\Phi_L = (\mathbf{u}_L, p_L)$ for a finite domain is governed by

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}_L) - \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla p_L + \omega^2 \left(\rho - \frac{\rho_f^2}{\gamma}\right) \mathbf{u}_L &= \mathbf{0} & \text{in } P \setminus \bar{L}, \\ \frac{1}{\gamma \omega^2} \nabla^2 p_L + M^{-1} p_L + \left(\alpha - \frac{\rho_f}{\gamma}\right) \nabla \cdot \mathbf{u}_L &= 0 & \text{in } P \setminus \bar{L} \\ (\mathbf{n} \cdot \mathbf{C} : \nabla \mathbf{u}_L, p_L) &= \mathbf{0} & \text{on } \partial P, \\ ([\mathbf{u}_L], [p_L], -[q_L]) &= \mathbf{a}, \quad [\mathbf{t}_L] = \mathbf{0} & \text{on } L. \end{aligned} \quad (61)$$

In what follows, the trial signatures $\Phi_{\mathbf{x}_o}^{\mathbf{n}, \mathbf{t}} = (\mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \mathbf{t}}, p_{\mathbf{x}_o}^{\mathbf{n}, \mathbf{t}})(\xi_i)$ over the observation grid $\xi_i \in \mathcal{G}$ are computed separately for every sampling point $\mathbf{x}_o \in L \subset P$ by solving

$$\begin{aligned}
\nabla \cdot (C : \nabla \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota}) - (\alpha - \frac{\rho_f}{\gamma}) \nabla p_{\mathbf{x}_o}^{\mathbf{n}, \iota} + \omega^2 (\rho - \frac{\rho_f^2}{\gamma}) \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota} &= \mathbf{0} & \text{in } P \setminus \bar{L}, \\
\frac{1}{\gamma \omega^2} \nabla^2 p_{\mathbf{x}_o}^{\mathbf{n}, \iota} + M^{-1} p_{\mathbf{x}_o}^{\mathbf{n}, \iota} + (\alpha - \frac{\rho_f}{\gamma}) \nabla \cdot \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota} &= -(1 - \iota) \delta(\mathbf{x}_o - \boldsymbol{\xi}) & \text{in } P \setminus \bar{L} \\
(\mathbf{n} \cdot C : \nabla \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota}, p_{\mathbf{x}_o}^{\mathbf{n}, \iota}) &= \mathbf{0} & \text{on } \partial P, \\
\iota (\mathbf{n} \cdot C : \nabla \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota} - |\mathbf{L}|^{-1} \delta(\mathbf{x}_o - \boldsymbol{\xi}) \mathbf{n}, p_{\mathbf{x}_o}^{\mathbf{n}, \iota}) &= \mathbf{0} & \text{on } \mathbf{x}_o + \mathbf{RL}, \\
[[p_{\mathbf{x}_o}^{\mathbf{n}, \iota}]] = 0, \quad [[\mathbf{n} \cdot C : \nabla \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota}]] &= \mathbf{0} & \text{on } \mathbf{x}_o + \mathbf{RL}.
\end{aligned} \tag{62}$$

where $\iota = 0, 1$, and the trial dislocation L is described by an infinitesimal crack $L = \mathbf{x}_o + \mathbf{RL}$ wherein \mathbf{R} is a unitary rotation matrix, and \mathbf{L} represents a vanishing penny-shaped crack of unit normal $\mathbf{n}_o := \{1, 0, 0\}$, so that $\mathbf{n} = \mathbf{R}\mathbf{n}_o$. On recalling (59), the three non-trivial components of $(\mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota}, p_{\mathbf{x}_o}^{\mathbf{n}, \iota})(\boldsymbol{\xi}_i)$ in the $x - y$ plane, with orthonormal bases $(\mathbf{e}_1, \mathbf{e}_2)$, are arranged into a $3N \times 1$ vector for $\boldsymbol{\xi}_i \in \mathcal{G}$ as the following

$$\Phi_{\mathbf{x}_o}^{\mathbf{n}, \iota}(3i + 1 : 3i + 3) = \begin{bmatrix} \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota} \cdot \mathbf{e}_1 \\ \mathbf{u}_{\mathbf{x}_o}^{\mathbf{n}, \iota} \cdot \mathbf{e}_2 \\ p_{\mathbf{x}_o}^{\mathbf{n}, \iota} \end{bmatrix}(\boldsymbol{\xi}_i), \quad i = 0, \dots, N - 1. \tag{63}$$

Sampling. With reference to Figs. 2 and 3, the search area i.e., the sampling region is a square $[-5, 5]^2 \subset P$ probed by a uniform 100×100 grid of sampling points \mathbf{x}_o where the featured indicator functionals are evaluated, while the unit circle spanning possible orientations for the trial dislocation L is sampled by a grid of 16 trial normal directions $\mathbf{n} = \mathbf{R}\mathbf{n}_o$, and the excitation form – as a fluid source or an elastic force, is selected via ι . Accordingly, the multiphase indicator maps \mathfrak{L} and \mathfrak{G} are constructed through minimizing the cost functionals (47) and (55) for a total of $M = 10000 \times 8 \times 2$ trial triplets $(\mathbf{x}_o, \mathbf{n}, \iota)$.

Remark 5.2. *It is worth mentioning that in the discretized equation*

$$\Lambda^\delta \mathbf{g}_{\mathbf{x}_o}^{\mathbf{n}, \iota} = \Phi_{\mathbf{x}_o}^{\mathbf{n}, \iota}, \tag{64}$$

the scattering operators Λ^δ is independent of any particular choice of $L(\mathbf{x}_o, \mathbf{n})$ or ι . Therefore, one may replace $\Phi_{\mathbf{x}_o}^{\mathbf{n}, \iota}$ in (64) with an assembled $3N \times M$ matrix, encompassing all variations of $L(\mathbf{x}_o, \mathbf{n})$ and ι , to solve only one equation for the reconstruction of hidden fractures which is computationally more efficient.

Step 3: computation of the multiphysics imaging functionals

\mathfrak{L} map. With reference to Theorem 4.2, (64) is solved by minimizing the regularized cost functional (47). It is well-known, however, that the Tikhonov functional J_η is convex and its minimizer can be obtained without iteration [45]. Thus-obtained solution $\mathbf{g}_{\mathbf{x}_o, \mathfrak{L}}^{\mathbf{n}, \iota}$ to (64) is a $3N \times 1$ vector (or $3N \times M$ matrix for all right-hand sides) identifying the structure of multiphase wavefront densities on \mathcal{G} . On the basis of (48) and (49), the multiphysics \mathfrak{L} indicator is then computed as

$$\mathfrak{L}(\mathbf{x}_o) := \frac{1}{\|\mathbf{g}_{\mathbf{x}_o, \mathfrak{L}}^{\mathfrak{L}}\|_{L^2}}, \quad \mathbf{g}_{\mathbf{x}_o}^{\mathfrak{L}} := \underset{\mathbf{g}_{\mathbf{x}_o, \mathfrak{L}}^{\mathbf{n}, \iota}}{\operatorname{argmin}} \|\mathbf{g}_{\mathbf{x}_o, \mathfrak{L}}^{\mathbf{n}, \iota}\|_{L^2}^2. \tag{65}$$

\mathfrak{G} map. In the case where $\Im \gamma \rightarrow 0$, one may construct a more robust approximate solution to (64) by minimizing $J_{\alpha, \delta}$ in (55). As indicated in Section 4, $J_{\alpha, \delta}$ is also convex and its minimizer $\mathbf{g}_{\mathbf{x}_o, \mathfrak{G}}^{\mathbf{n}, \iota}$ can be computed by solving the discretized linear system (57). Then, with reference to (58), the multiphase indicator

\mathfrak{G} is obtained as

$$\mathfrak{G}(\mathbf{x}_o) := \frac{1}{\sqrt{\|(\Lambda_{\sharp}^{\delta})^{\frac{1}{2}} \mathbf{g}_{\mathbf{x}_o}^{\mathfrak{G}}\|^2 + \delta \|\mathbf{g}_{\mathbf{x}_o}^{\mathfrak{G}}\|^2}}, \quad \mathbf{g}_{\mathbf{x}_o}^{\mathfrak{G}} := \underset{\mathbf{g}_{\mathbf{x}_o, \mathfrak{G}}^{\mathbf{n}, \mathbf{t}}}{\operatorname{argmin}} \|\mathbf{g}_{\mathbf{x}_o, \mathfrak{G}}^{\mathbf{n}, \mathbf{t}}\|_{L^2}^2. \quad (66)$$

5.4. Simulation results

Near-field tomography of an evolving network. With reference to Fig. 2, a hydraulic fracture network $\bigcup_{\kappa=1}^9 \Gamma_{\kappa}$ is induced in P in four steps. Following each treatment, the numerical experiments are conducted as described earlier where the multiphase scattering patterns $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$ are obtained over an H-shaped sensing grid \mathcal{G} resulting in a $3N \times 3N = 990 \times 990$ scattering matrix. The latter is then used to compute the distribution of \mathfrak{L} indicator in the sampling region, at every sensing step, to recover the sequential evolution of the network as shown in the figure.

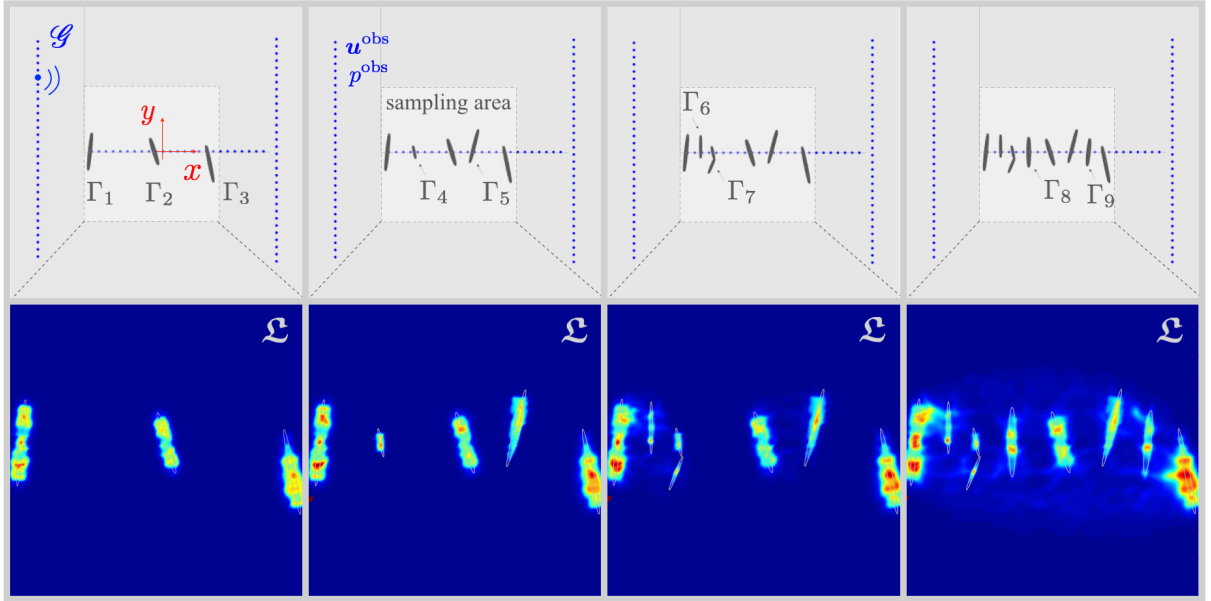


Figure 2: Reconstruction of an evolving hydraulic fracture network via the multiphysics \mathfrak{L} indicator: (top row) sensing and network configurations at each treatment stage, and (bottom row) \mathfrak{L} maps computed via (65) on the basis of solid displacements and pore pressures $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$ measured on an H-shaped grid $\xi_i \in \mathcal{G}, i = \{1, 2, \dots, 330\}$ in the injection and two monitoring wells.

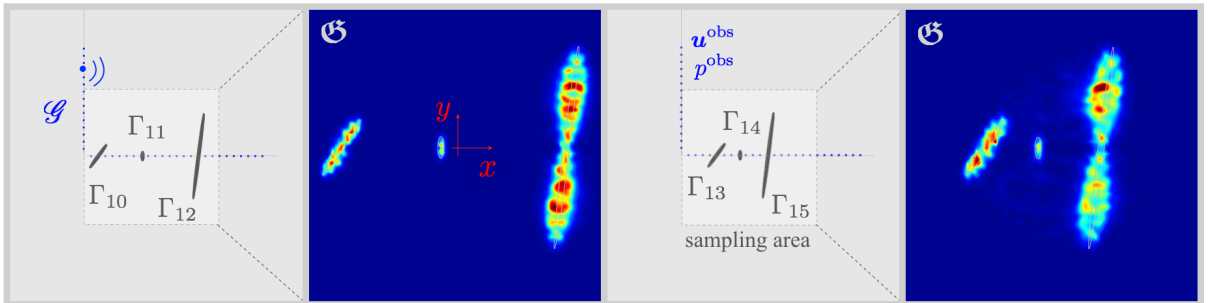


Figure 3: Limited-aperture reconstruction via the multiphysics \mathfrak{G} indicator: (panels 1 and 3) sensing and network configurations, and (panels 2 and 4) the affiliated \mathfrak{G} maps constructed via (66) from the complete dataset i.e., solid displacements and pore pressures $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$ observed on an L-shaped grid $\xi_i \in \mathcal{G}, i = \{1, 2, \dots, 130\}$ in the injection well.

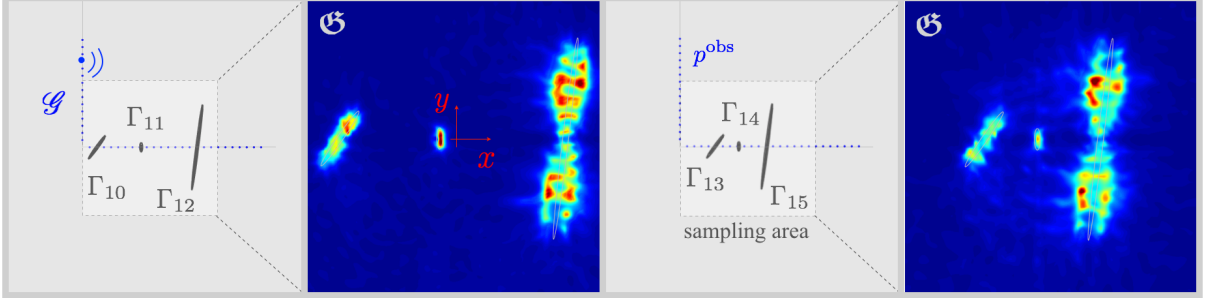


Figure 4: Near-field imaging via the pore fluid: (panels 1 and 3) sensing and network configurations similar to Fig. 3, and (panels 2 and 4) the associated \mathfrak{G} maps computed via (66) when the excitation at every point $\xi_i \in \mathcal{G}, i = \{1, 2, \dots, 130\}$ is a fluid volumetric source g_f and the measurements are in the form of pore pressure data p^{obs} .

Limited-aperture imaging. Fig. 3 shows the reconstructed \mathfrak{G} maps via (66) from the complete dataset, including both solid displacements and pore pressure measurements $(\mathbf{u}^{\text{obs}}, p^{\text{obs}})$, on an L-shaped grid \mathcal{G} associated with the injection (or treatment) well, which results in a $3N \times 3N = 390 \times 390$ poroelastic scattering matrix. Here, both networks $\bigcup_{\kappa=10}^{12} \Gamma_{\kappa}$ and $\bigcup_{\kappa=13}^{15} \Gamma_{\kappa}$ involve hydraulic fractures of different length scales e.g., Γ_{11} is of $O(1)$ while Γ_{12} is of $O(10)$. One may also observe the interaction between scatterers in the \mathfrak{G} maps when the pair-wise distances between fractures is less than a few wavelengths. It should be mentioned that similar maps are obtained by deploying the \mathfrak{L} imaging functional. However, it is interesting to note that the \mathfrak{G} indicator is successful in reconstructing the fracture network even though $\gamma \in \mathbb{C}$. This may be due to the particular sensor arrangement germane to the hydraulic fracturing application where a subset of the sensing grid (within the injection well) is in fact in close proximity of scatterers $\Gamma_{\kappa}, \kappa = 1, 2, \dots, 15$.

Acoustic imaging via the pore fluid. In hydrofracking, it may be convenient to generate distributed fluid volumetric excitation within the treatment well and simultaneously measure thus-induced pore pressures on the same grid. Data inversion on the basis of acoustic excitation and measurements may expedite reconstruction of the induced fractures. In this spirit, Fig. 4 illustrates the reconstruction results for the network and sensing configurations similar to Fig. 3 where the excitation takes only the form of a fluid volumetric source, and the measurements are pore pressure data p^{obs} . The \mathfrak{L} indicator generates very similar maps.

6. Conclusions

A multiphysics data inversion framework is developed for near-field tomography of hydraulic fracture networks in poroelastic backgrounds. This imaging solution is capable of spatiotemporal reconstruction of discontinuous interfaces of arbitrary (and possibly fragmented) support whose poro-mechanical properties are unknown. The data is processed within an appropriate dimensional framework to allow for simultaneous inversion of elastic and acoustic (i.e., pore pressure) measurements, which are distinct in physics and scale. The proposed imaging indicator is (a) inherently non-iterative enabling fast inversion of dense data in support of real-time sensing, (b) flexible with respect to the sensing configuration and illumination frequency, and (c) robust against noise in data. Limited-aperture and partial data inversion – e.g., deep-well tomography as well as imaging based solely on pore pressure excitation and measurements – are explored. Both \mathfrak{L} and \mathfrak{G} indicators showed success in the numerical experiments germane to the limiting case of high-frequency illumination where the probing wavelength is significantly smaller than the characteristic length scale of the sought-for scatterers. It should be mentioned that this imaging modality can be naturally extended to more complex e.g., highly heterogeneous backgrounds. Given the multiphysics nature of data, a remarkable potential would be to deploy

this approach within a hybrid data analytic platform to enable not only geometric reconstruction, but also interfacial characterization of discontinuity surfaces, e.g., involving the retrieval of permeability profile of hydraulic fractures from poroelastic data.

7. Acknowledgements

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Appendix A. Poroelastodynamics fundamental solution

The poroelastodynamics fundamental solution [41, 52] takes the form

$$\mathbf{G}(\mathbf{y}, \boldsymbol{\xi}) := \begin{bmatrix} \mathbf{U}_{ij}^s & \mathbf{u}_i^f \\ \mathbf{p}_j^s & \mathbf{p}^f \end{bmatrix} (\mathbf{y}, \boldsymbol{\xi}), \quad \mathbf{y}, \boldsymbol{\xi} \in \mathbb{R}^3, \quad \boldsymbol{\xi} \neq \mathbf{y}, \quad i, j = 1, 2, 3, \quad (\text{A.1})$$

where $(\mathbf{U}^s, \mathbf{p}^s) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\mathbf{y}\})^{3 \times 3} \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\mathbf{y}\})^3$ satisfies

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{U}^s)(\mathbf{y}, \boldsymbol{\xi}) - (\alpha - \frac{\rho_f}{\gamma}) \nabla \mathbf{p}^s(\mathbf{y}, \boldsymbol{\xi}) + \omega^2 (\rho - \frac{\rho_f^2}{\gamma}) \mathbf{U}^s(\mathbf{y}, \boldsymbol{\xi}) &= -\delta(\mathbf{y} - \boldsymbol{\xi}) \mathbf{I}_{3 \times 3}, \\ \frac{1}{\gamma \omega^2} \nabla^2 \mathbf{p}^s(\mathbf{y}, \boldsymbol{\xi}) + M^{-1} \mathbf{p}^s(\mathbf{y}, \boldsymbol{\xi}) + (\alpha - \frac{\rho_f}{\gamma}) \nabla \cdot \mathbf{U}^s(\mathbf{y}, \boldsymbol{\xi}) &= \mathbf{0}, \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \{\mathbf{y}\}, \boldsymbol{\xi} \neq \mathbf{y} \in \mathbb{R}^3, \end{aligned} \quad (\text{A.2})$$

while $(\mathbf{u}^f, \mathbf{p}^f) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\mathbf{y}\})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \{\mathbf{y}\})$ solves

$$\begin{aligned} \nabla \cdot (\mathbf{C} : \nabla \mathbf{u}^f)(\mathbf{y}, \boldsymbol{\xi}) - (\alpha - \frac{\rho_f}{\gamma}) \nabla \mathbf{p}^f(\mathbf{y}, \boldsymbol{\xi}) + \omega^2 (\rho - \frac{\rho_f^2}{\gamma}) \mathbf{u}^f(\mathbf{y}, \boldsymbol{\xi}) &= \mathbf{0}, \\ \frac{1}{\gamma \omega^2} \nabla^2 \mathbf{p}^f(\mathbf{y}, \boldsymbol{\xi}) + M^{-1} \mathbf{p}^f(\mathbf{y}, \boldsymbol{\xi}) + (\alpha - \frac{\rho_f}{\gamma}) \nabla \cdot \mathbf{u}^f(\mathbf{y}, \boldsymbol{\xi}) &= -\delta(\mathbf{y} - \boldsymbol{\xi}), \quad \boldsymbol{\xi} \in \mathbb{R}^3 \setminus \{\mathbf{y}\}, \boldsymbol{\xi} \neq \mathbf{y} \in \mathbb{R}^3, \end{aligned} \quad (\text{A.3})$$

both subject to the relevant radiation conditions similar to (9). On setting

$$\begin{aligned} r &:= |\mathbf{y} - \boldsymbol{\xi}|, \quad G(k_\varsigma, r) := \frac{e^{ik_\varsigma r}}{4\pi r}, \quad \varsigma = s, p_1, p_2, \quad [f]_{,i} := \frac{\partial f}{\partial \xi_i}, \quad i = 1, 2, 3, \\ A_1 &:= \frac{k_{p_2}^2 (k_{p_1}^2 M - \gamma \omega^2)}{\gamma \omega^2 (k_{p_1}^2 - k_{p_2}^2)}, \quad A_2 := \frac{k_{p_1}^2 (k_{p_2}^2 M - \gamma \omega^2)}{\gamma \omega^2 (k_{p_2}^2 - k_{p_1}^2)}, \end{aligned}$$

$(\mathbf{U}^s, \mathbf{p}^s)$ permits the explicit form

$$\begin{aligned} \mathbf{U}_{ij}^s &= \frac{\gamma}{\omega^2 (\rho \gamma - \rho_f^2)} \left[\left(G(k_s, r) - A_1 G(k_{p_1}, r) - A_2 G(k_{p_2}, r) \right)_{,ij} + \delta_{ij} k_s^2 G(k_s, r) \right], \\ \mathbf{p}_j^s &= \frac{\omega^2 (\alpha \gamma - \rho_f)}{(\lambda + 2\mu) (k_{p_1}^2 - k_{p_2}^2)} \left(G(k_{p_1}, r) - G(k_{p_2}, r) \right)_{,j}, \quad i, j = 1, 2, 3, \end{aligned} \quad (\text{A.4})$$

and $(\mathbf{u}^f, \mathbf{p}^f)$ reads

$$\mathbf{u}^f = -\mathbf{p}^s, \quad \mathbf{p}^f = -\frac{\gamma\omega^2}{k_{p_2}^2 - k_{p_1}^2} \left[\left(k_{p_1}^2 - \frac{\omega^2(\rho\gamma - \rho_f^2)}{\gamma(\lambda + 2\mu)} \right) G(k_{p_1}, r) - \left(k_{p_2}^2 - \frac{\omega^2(\rho\gamma - \rho_f^2)}{\gamma(\lambda + 2\mu)} \right) G(k_{p_2}, r) \right], \quad (\text{A.5})$$

In this setting, the associated fundamental tractions \mathbf{T}^s and \mathbf{t}^f on a surface $\Gamma \subset \mathbb{R}^3 \setminus \{\mathbf{y}\}$ characterized by the unit normal vector \mathbf{n} may be specified as

$$\begin{aligned} \mathbf{T}_{ij}^s(\mathbf{y}, \boldsymbol{\xi}) &= n_k(\boldsymbol{\xi}) C_{ik\ell\varsigma} \mathbf{U}_{\ell j, \varsigma}^s(\mathbf{y}, \boldsymbol{\xi}) - \alpha n_i(\boldsymbol{\xi}) \mathbf{p}_j^s(\mathbf{y}, \boldsymbol{\xi}), \\ \mathbf{t}_i^f(\mathbf{y}, \boldsymbol{\xi}) &= n_k(\boldsymbol{\xi}) C_{ik\ell\varsigma} \mathbf{u}_{\ell, \varsigma}^f(\mathbf{y}, \boldsymbol{\xi}) - \alpha n_i(\boldsymbol{\xi}) \mathbf{p}^f(\mathbf{y}, \boldsymbol{\xi}), \quad \boldsymbol{\xi} \in \Gamma, \mathbf{y} \in \mathbb{R}^3 \setminus \bar{\Gamma}, \end{aligned} \quad (\text{A.6})$$

where Einstein's summation convention applies over the repeated indexes. In addition, the normal relative fluid-solid displacements \mathbf{q}^s and \mathbf{q}^f across Γ may be expressed as

$$\begin{aligned} \mathbf{q}_j^s(\mathbf{y}, \boldsymbol{\xi}) &= \frac{1}{\gamma\omega^2} \left(n_i(\boldsymbol{\xi}) \mathbf{p}_{j,i}^s(\mathbf{y}, \boldsymbol{\xi}) - \rho_f \omega^2 n_i(\boldsymbol{\xi}) \mathbf{U}_{ij}^s(\mathbf{y}, \boldsymbol{\xi}) \right), \\ \mathbf{q}_i^f(\mathbf{y}, \boldsymbol{\xi}) &= \frac{1}{\gamma\omega^2} \left(n_i(\boldsymbol{\xi}) \mathbf{p}_i^f(\mathbf{y}, \boldsymbol{\xi}) - \rho_f \omega^2 n_i(\boldsymbol{\xi}) \mathbf{u}_i^f(\mathbf{y}, \boldsymbol{\xi}) \right), \quad \boldsymbol{\xi} \in \Gamma, \mathbf{y} \in \mathbb{R}^3 \setminus \bar{\Gamma}. \end{aligned} \quad (\text{A.7})$$

Appendix B. Wellposedness of the poroelastic scattering problem

Given $(\mathbf{t}', q^t, p') \in H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ and the interface operator $\mathfrak{P} : \tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times \tilde{H}^{-1/2}(\Gamma) \rightarrow H^{-1/2}(\Gamma)^3 \times H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ satisfying (13), then the direct scattering problem (5)-(9) has a unique solution that continuously depends on (\mathbf{t}', q^t, p') . Note that q^t is continuously dependent on $(\mathbf{u}' \cdot \mathbf{n}, \nabla p' \cdot \mathbf{n})$ according to (6). First, observe that (5)-(9) may be expressed on $\mathcal{B} \setminus \bar{\Gamma}$ in the variational form in terms of $(\mathbf{u}, p) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ so that $\forall (\mathbf{u}', \mathbf{p}') \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ satisfying the radiation condition (9),

$$\begin{aligned} &\left\langle (\mathbf{t}' + \tilde{\alpha}_f p' \mathbf{n}, q^t, p'), (\llbracket \mathbf{u}' \rrbracket, \llbracket \mathbf{p}' \rrbracket, \vartheta_f \langle \mathbf{p}' \rangle) \right\rangle_{\Gamma} = \\ &\int_{\mathcal{B} \setminus \Gamma} \left\{ \nabla \bar{\mathbf{u}}' : \mathbf{C} : \nabla \mathbf{u} + \frac{1}{\gamma\omega^2} \nabla \bar{\mathbf{p}}' \cdot \nabla p - \left(\alpha - \frac{\rho_f}{\gamma} \right) (\nabla \cdot \bar{\mathbf{u}}' p + \bar{\mathbf{p}}' \nabla \cdot \mathbf{u}) \right\} dV_{\boldsymbol{\xi}} - \\ &\int_{\mathcal{B} \setminus \Gamma} \left\{ \omega^2 \left(\rho - \frac{\rho_f^2}{\gamma} \right) \bar{\mathbf{u}}' \cdot \mathbf{u} + M^{-1} \bar{\mathbf{p}}' p \right\} dV_{\boldsymbol{\xi}} + \frac{\rho_f}{\gamma} \int_{\Gamma} \left(\langle p \rangle \llbracket \bar{\mathbf{u}}' \rrbracket + \llbracket p \rrbracket \langle \bar{\mathbf{u}}' \rangle + \langle \bar{\mathbf{p}}' \rrbracket \llbracket \mathbf{u} \rrbracket + \llbracket \bar{\mathbf{p}}' \rrbracket \langle \mathbf{u} \rangle \right) \cdot \mathbf{n} dS_{\boldsymbol{\xi}} - \quad (\text{B.1}) \\ &\left\langle \left(-\llbracket \mathbf{u} \rrbracket \cdot \mathbf{K} + \tilde{\alpha}_f \langle p \rangle \mathbf{n}, -\frac{\mathcal{K}_f}{i\omega\Pi} \llbracket p \rrbracket, \langle p \rangle + \frac{\beta_f}{1 - \tilde{\alpha}_f \beta_f} \mathbf{n} \cdot \mathbf{K} \llbracket \mathbf{u} \rrbracket \right), (\llbracket \mathbf{u}' \rrbracket, \llbracket \mathbf{p}' \rrbracket, \vartheta_f \langle \mathbf{p}' \rangle) \right\rangle_{\Gamma} - \\ &\int_{\partial \mathcal{B}} \left\{ \bar{\mathbf{u}}' \cdot \mathbf{t} + \frac{1}{\gamma\omega^2} \bar{\mathbf{p}}' \nabla p \cdot \mathbf{n} + \frac{\rho_f}{\gamma} \bar{\mathbf{u}}' \cdot \mathbf{n} p \right\} dS_{\boldsymbol{\xi}}, \end{aligned}$$

where $\vartheta_f \in L^\infty(\Gamma)$ is defined by

$$\vartheta_f := \overline{\frac{\Pi \alpha_f}{k_n \beta_f} (1 - \tilde{\alpha}_f \beta_f)} \quad \text{on } \Gamma,$$

with the 'bar' indicating complex conjugate. Note that in light of (9), one has

$$\lim_{R \rightarrow \infty} \int_{\partial \mathcal{B}} \left\{ \bar{\mathbf{u}}' \cdot \mathbf{t} + \frac{1}{\gamma\omega^2} \bar{\mathbf{p}}' \nabla p \cdot \mathbf{n} + \frac{\rho_f}{\gamma} \bar{\mathbf{u}}' \cdot \mathbf{n} p \right\} dS_{\boldsymbol{\xi}} = 0.$$

It is worth mentioning that (B.1) is obtained by premultiplying the first of (5) by $\bar{\mathbf{u}}'$ and its second by $\bar{\mathbf{p}}'$; the results are then added and integrated by parts over $\mathcal{B} \setminus \bar{\Gamma}$. The duality pairing $\langle \cdot, \cdot \rangle_{\Gamma}$ on the left hand side of (B.1) is continuous with respect to \mathbf{t}^t , q^t , and p^t owing to the continuity of the trace mapping $(\mathbf{u}', \mathbf{p}') \rightarrow (\llbracket \mathbf{u}' \rrbracket, \llbracket \mathbf{p}' \rrbracket, \langle \mathbf{p}' \rangle)$ from $H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ into $\tilde{H}^{1/2}(\Gamma)^3 \times \tilde{H}^{1/2}(\Gamma) \times H^{1/2}(\Gamma)$. Now let us define $\forall (\mathbf{u}', \mathbf{p}') \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ subject to (9), the coercive operator

$$\begin{aligned} A((\mathbf{u}, p), (\mathbf{u}', \mathbf{p}')) &:= \int_{\mathcal{B} \setminus \Gamma} \{ \nabla \bar{\mathbf{u}}' : \mathbf{C} : \nabla \mathbf{u} + \frac{1}{\gamma \omega^2} \nabla \bar{\mathbf{p}}' \cdot \nabla p - \\ &(\alpha - \frac{\rho_f}{\gamma}) (\nabla \cdot \bar{\mathbf{u}}' p + \bar{\mathbf{p}}' \nabla \cdot \mathbf{u}) + \bar{\mathbf{p}}' p + \frac{\rho_f^2}{\gamma} \omega^2 \bar{\mathbf{u}}' \cdot \mathbf{u} \} dV_{\xi}. \end{aligned} \quad (\text{B.2})$$

On invoking $\mu_r = \lambda'$ i.e., $\lambda = 1$ from Section 2.1 and recalling from (4) that $|\alpha - \rho_f/\gamma| < 1$, one may observe

$$\begin{aligned} \Re(A)((\mathbf{u}, p), (\mathbf{u}, p)) &\geq \int_{\mathcal{B} \setminus \Gamma} \left\{ (\lambda - |\alpha - \frac{\rho_f}{\gamma}|) \|\nabla \cdot \mathbf{u}\|^2 + \mu \|\nabla \mathbf{u}\|^2 + \frac{\Re \gamma}{\gamma^2 \omega^2} \|\nabla p\|^2 + \right. \\ &\left. (1 - |\alpha - \frac{\rho_f}{\gamma}|) \|p\|^2 + \frac{\rho_f^2}{\gamma^2} \omega^2 \Re \gamma \|\mathbf{u}\|^2 \right\} dV_{\xi} \geq c (\|\mathbf{u}\|_{H^1(\mathcal{B} \setminus \Gamma)^3}^2 + \|p\|_{H^1(\mathcal{B} \setminus \Gamma)}^2), \end{aligned} \quad (\text{B.3})$$

for a constant $c > 0$ independent of (\mathbf{u}, p) . Next, define $\forall (\mathbf{u}', \mathbf{p}') \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ subject to (9), the compact operator

$$\begin{aligned} B((\mathbf{u}, p), (\mathbf{u}', \mathbf{p}')) &= -\langle (-\llbracket \mathbf{u} \rrbracket \cdot \mathbf{K} + \tilde{\alpha}_f \langle p \rangle \mathbf{n}, -\frac{\varkappa_f}{i\omega\Pi} \llbracket p \rrbracket, \langle p \rangle + \frac{\beta_f}{1 - \tilde{\alpha}_f \beta_f} \mathbf{n} \cdot \mathbf{K} \llbracket \mathbf{u} \rrbracket), (\llbracket \mathbf{u}' \rrbracket, \llbracket \mathbf{p}' \rrbracket, \vartheta_f \langle \mathbf{p}' \rangle) \rangle_{\Gamma} - \\ &\int_{\mathcal{B} \setminus \Gamma} \{ \rho \omega^2 \bar{\mathbf{u}}' \cdot \mathbf{u} + (M^{-1} + 1) \bar{\mathbf{p}}' p \} dV_{\xi} + \frac{\rho_f}{\gamma} \int_{\Gamma} (\langle p \rangle \llbracket \bar{\mathbf{u}}' \rrbracket + \llbracket p \rrbracket \langle \bar{\mathbf{u}}' \rangle + \langle \bar{\mathbf{p}}' \rangle \llbracket \mathbf{u} \rrbracket + \llbracket \bar{\mathbf{p}}' \rrbracket \langle \mathbf{u} \rangle) \cdot \mathbf{n} dS_{\xi}, \end{aligned} \quad (\text{B.4})$$

where

$$\begin{aligned} |B((\mathbf{u}, p), (\mathbf{u}', \mathbf{p}'))| &\leq c \left\{ \|\mathbf{u}\|_{L^2(\mathcal{B} \setminus \Gamma)^3} \|\mathbf{u}'\|_{L^2(\mathcal{B} \setminus \Gamma)^3} + \|p\|_{L^2(\mathcal{B} \setminus \Gamma)} \|\mathbf{p}'\|_{L^2(\mathcal{B} \setminus \Gamma)} + \right. \\ &\|\llbracket \mathbf{u} \rrbracket\|_{L^2(\Gamma)^3} \|\llbracket \mathbf{u}' \rrbracket\|_{L^2(\Gamma)^3} + \|\llbracket p \rrbracket\|_{L^2(\Gamma)} \|\llbracket \mathbf{p}' \rrbracket\|_{L^2(\Gamma)} + \|\langle p \rangle\|_{L^2(\Gamma)} \|\langle \mathbf{p}' \rangle\|_{L^2(\Gamma)} + \\ &\left. \|\langle p \rangle\|_{L^2(\Gamma)} \|\llbracket \mathbf{u}' \rrbracket\|_{L^2(\Gamma)^3} + \|\llbracket p \rrbracket\|_{L^2(\Gamma)} \|\langle \bar{\mathbf{u}}' \rangle\|_{L^2(\Gamma)^3} + \|\llbracket \mathbf{u} \rrbracket\|_{L^2(\Gamma)^3} \|\langle \mathbf{p}' \rangle\|_{L^2(\Gamma)} + \|\langle \mathbf{u} \rangle\|_{L^2(\Gamma)^3} \|\llbracket \mathbf{p}' \rrbracket\|_{L^2(\Gamma)} \right\}, \end{aligned}$$

for a constant c independent of (\mathbf{u}, p) and $(\mathbf{u}', \mathbf{p}')$. Then, the compact embedding of $H^1(\mathcal{B} \setminus \bar{\Gamma})$ into $L^2(\mathcal{B} \setminus \bar{\Gamma})$ and the compactness of the trace operator indicates that B defines a compact perturbation of $A((\mathbf{u}, p), (\mathbf{u}', \mathbf{p}'))$. Therefore, on letting $R \rightarrow \infty$, one may observe that the problem (B.1) is of Fredholm type, and thus, is well-posed as soon as the uniqueness of a solution is guaranteed. Assume that $(\mathbf{t}^t, q^t, p^t) = \mathbf{0}$. Then, on invoking (4) and (12), one may show that when $R \rightarrow \infty$,

$$-\frac{\omega}{\kappa} \|\mathbf{q}\|_{L^2(\mathbb{R}^3 \setminus \Gamma)^3}^2 + \Im \langle \mathfrak{P}(\llbracket \mathbf{u} \rrbracket, \llbracket p \rrbracket, -\llbracket q \rrbracket), (\llbracket \mathbf{u} \rrbracket, \llbracket p \rrbracket, -\llbracket q \rrbracket) \rangle_{\Gamma} = 0, \quad (\text{B.5})$$

where

$$\mathbf{q} := \frac{1}{\gamma \omega^2} (\nabla p - \rho_f \omega^2 \mathbf{u}) \quad \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}. \quad (\text{B.6})$$

It should be mentioned that (B.5) is obtained by (a) premultiplying the first of (5) by $\bar{\mathbf{u}} \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3$, (b) conjugating the second of (5) and multiplying the results by $-p \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$, and (c) integrating by parts

over $\mathcal{B} \setminus \bar{\Gamma}$ and adding the results. On letting $R \rightarrow \infty$, one finds (B.5) indicating that $\mathbf{q} = \mathbf{0}$ in $\mathbb{R}^3 \setminus \bar{\Gamma}$ by the premise of Remark 2.3. As a result, the first equation in (5) will be reduced to the Navier equation that is subject to the (exponentially decaying) poroelastodynamics radiation conditions (9) which implies that $\mathbf{u} = \mathbf{0}$ in $\mathbb{R}^3 \setminus \bar{\Gamma}$. The second equation in (5) then reads $p = 0$ in $\mathbb{R}^3 \setminus \bar{\Gamma}$ which completes the proof for the uniqueness of a scattering solution, and thus, substantiates the wellposedness of the forward problem.

Appendix C. Proof of Lemma 2.1

Let us define

$$\begin{aligned}\boldsymbol{\sigma}^\iota &:= \mathbf{C} : \nabla \mathbf{u}^\iota - \alpha p^\iota \mathbf{I}_{3 \times 3}, & \zeta^\iota &:= M^{-1} p^\iota + \alpha \nabla \cdot \mathbf{u}^\iota & \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \\ \boldsymbol{\sigma}_a &:= \mathbf{C} : \nabla \mathbf{u}_a - \alpha p_a \mathbf{I}_{3 \times 3}, & \zeta_a &:= M^{-1} p_a + \alpha \nabla \cdot \mathbf{u}_a & \text{in } \mathbb{R}^3 \setminus \bar{\Gamma},\end{aligned}$$

where $(\mathbf{u}^\iota, p^\iota) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\mathcal{G}})$ and $(\mathbf{u}_a, p_a) \in H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})^3 \times H_{\text{loc}}^1(\mathbb{R}^3 \setminus \bar{\Gamma})$ solve (3) and (22) respectively. Then, the poroelastic reciprocal theorem [41] reads

$$\nabla \bar{\mathbf{u}}_a : \boldsymbol{\sigma}^\iota + \bar{\zeta}_a p^\iota = \bar{\boldsymbol{\sigma}}_a : \nabla \mathbf{u}^\iota + \bar{p}_a \zeta^\iota. \quad (\text{C.1})$$

Now, on setting

$$\mathbf{q}^\iota := \frac{1}{\gamma \omega^2} (\nabla p^\iota - \rho_f \omega^2 \mathbf{u}^\iota), \quad \mathbf{q}_a := \frac{1}{\bar{\gamma} \omega^2} (\nabla p_a - \rho_f \omega^2 \mathbf{u}_a) \quad \text{in } \mathbb{R}^3 \setminus \bar{\Gamma}, \quad (\text{C.2})$$

observe that

$$\begin{aligned}\int_{\mathbb{R}^3 \setminus \Gamma} \{ \nabla \bar{\mathbf{u}}_a : \boldsymbol{\sigma}^\iota + \bar{\zeta}_a p^\iota \} dV_\xi &:= - \int_\Gamma \{ \llbracket \bar{\mathbf{u}}_a \rrbracket \cdot \mathbf{t}^\iota - \llbracket \bar{\mathbf{q}}_a \rrbracket p^\iota \} dS_\xi + \int_{\mathcal{G}} \bar{\mathbf{u}}_a \cdot \mathbf{g}_s dS_\xi + \\ &\quad \rho \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \bar{\mathbf{u}}_a \cdot \mathbf{u}^\iota dV_\xi + \rho_f \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \{ \bar{\mathbf{u}}_a \cdot \mathbf{q}^\iota + \bar{\mathbf{q}}_a \cdot \mathbf{u}^\iota \} dV_\xi + \gamma \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \bar{\mathbf{q}}_a \cdot \mathbf{q}^\iota dV_\xi, \\ \int_{\mathbb{R}^3 \setminus \Gamma} \{ \bar{\boldsymbol{\sigma}}_a : \nabla \mathbf{u}^\iota + \bar{p}_a \zeta^\iota \} dV_\xi &:= \int_\Gamma \llbracket \bar{p}_a \rrbracket q^\iota dS_\xi - \int_{\mathcal{G}} \bar{p}_a \cdot \mathbf{g}_f dS_\xi + \\ &\quad \rho \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \bar{\mathbf{u}}_a \cdot \mathbf{u}^\iota dV_\xi + \rho_f \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \{ \bar{\mathbf{u}}_a \cdot \mathbf{q}^\iota + \bar{\mathbf{q}}_a \cdot \mathbf{u}^\iota \} dV_\xi + \gamma \omega^2 \int_{\mathbb{R}^3 \setminus \Gamma} \bar{\mathbf{q}}_a \cdot \mathbf{q}^\iota dV_\xi.\end{aligned}$$

Then, by invoking (19), (21), and (22), it is evident that

$$\langle \mathcal{S}(\mathbf{g}), \mathbf{a} \rangle_\Gamma = (\mathbf{g}, \mathcal{S}^*(\mathbf{a}))_{L^2(\mathcal{G})},$$

which completes the proof.

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