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# Analysis of spectator chemical bonds in $S_N 2@C$ and @Si reaction mechanisms in the gas phase

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# A R T I C L E I N F O A B S T R A C T Keywords: In this work, Cl<sup>-</sup>+AR<sub>3</sub>Cl reactions (with A = C and Si, and R = H, Me, Et, Cl, and F) were investigated to evaluate the A-R spectator bond properties at the $\omega$ -B97X-D/SPK-TZP level of theory, applying the Quantum Theory of Atoms in Molecules, the Overlap Model, and Local Vibrational Mode theory. The different chemical bond analyses converge to the conclusion (in line with current literature) that the steric hindrance experienced by Cl<sup>-</sup> in S<sub>N</sub>2@C reactions can be viewed as a consequence of the greater covalent nature of C—R bonds that concentrate density in bond region more efficiently than Si—R.

# 1. Introduction

Understanding the key aspects of chemical reaction mechanisms is usually focused on the characterization of the potential energy surface (PES) [1-3] targeting in particular structural and energetic properties of the stationary points, i.e., reactants, products, intermediates, and transition states (TS). A more complete picture can be obtained by following the reaction complex (unity of reacting species) along the reaction path from entrance to exit channel, as realized in the unified reaction valley approach which correlates all chemical events including bond breaking/ formation to the curvature of the reaction path [4]. Recent studies have suggested that chemical bond analysis methods can provide details about the evolution of bonding during a chemical reaction [5-10] when applied to points along the energy profile. E.g., (i) Farfán et al. [6,7] investigated the stereoselectivity of Witting reactions applying different bond analysis methodologies. They concluded that the degree of advancement of the emerging C-C bond is intimately tied to the stereochemistry of the reaction product. (ii) Oliveira et al. [8] used a wide variety of chemical bond descriptors for a series of fluorination reactions to study the nature of X—F bonds (with X = B, C, N, O, and F) and also to understand the covalent, hypervalent, and noncovalent bonds in these compounds. [8]

An attractive target reaction for the exploration of different bond analysis methods is the bimolecular nucleophilic substitution ( $S_N$ 2) [11–16]. Classically, the mechanism of the  $S_N$ 2 reaction is described by a

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https://doi.org/10.1016/j.cplett.2021.139282 Received 26 November 2021; Accepted 6 December 2021 Available online 7 December 2021 0009-2614/© 2021 Elsevier B.V. All rights reserved. single kinetic step that comprises the attack of a nucleophile (Nu<sup>-</sup>) at the backside, opposite to the leaving group (X) of the AR<sub>3</sub>X reagent to yield the NuAR<sub>3</sub> product with configuration inversion. As shown in Fig. 1, a [Nu…AR<sub>3</sub>…X]<sup>-</sup> pentacoordinate intermediate separates the reactants and products. In the [Nu…AR<sub>3</sub>…X]<sup>-</sup> intermediate, the substituent groups (R) are forced into a planar trigonal rearrangement as a consequence of the repulsive action around the central atom (A).

Bickelhalput et al. [17,18] investigated a series of  $S_N2@A$  (A = C, Si, and P) reactions in the gas phase following the scheme depicted in Fig. 1. Fig. 2 shows four possible energy profiles for these type of reactions.

- (a) *Single-peak profile*, where the TS represents the high point of the
- (a) Single-peak profile, where the 1S represents the high point of the potential energy surface;
- (b) *Single-well profile*, where a single stable transitional complex (TC) is formed;
- (c) Double-well profile, with two stable ion-dipole complexes between the TS and the reagents/products, here labeled as reactant complex (RC) and product complex (PC);
- (d) Triple-well profile, where three minimum stationary points (RC, PC, and TC) and two maximum (pre-TS and post-TS, being before and after the TC, respectively) are obtained.

A well-marked feature of the profiles depicted in Fig. 2 consists of the nature of the pentacoordinate intermediate. For the  $S_N 2@C$  reactions in black or blue colour, the  $[Nu\cdots CR_3\cdots X]^-$  structure corresponds to a

saddle point (=TS) on the potential energy surface. For the  $S_N 2@Si$  reactions in red color the  $[Nu \cdots SiR_3 \cdots X]^-$  species often forms a stable intermediate (=TC).

Interestingly, in  $S_N 2@P$  reactions (particularly for  $Nu^- + POR_2X$ ) all four types of profiles can be obtained depending on the nature of  $Nu^-$ , X, and R [16]. The energy decomposition analysis (EDA) has revealed that an interplay between the steric (Pauli) repulsion and the electronic effects can explain the distinct shapes of the PES in the gas phase  $S_N 2@A$ reactions [17–19]. Recently the quantum theory of atoms in molecules (QTAIM) [20] and natural bond orbital (NBO) [21] analysis have shown that hydrogen and  $\sigma$ -hole bonding can be responsible for the formation of triple-well profile in some  $S_N 2@Si$  reactions [9]. Alkorta and Elguero [22] studied  $S_N 2@C$ , @N, @Si, and @P reactions and pointed to the existence of valence expansion effect in the  $S_N 2@P$  reaction.

Thomas et al. [23] studied the rate of an  $S_N2@Si$  reaction with a fluoride nucleophile. The reaction rate decreased by a factor of up to 5.5 by coupling the Si—C vibrational stretching modes (Si – leaving group bonding) to the zero-point fluctuations of a resonant infrared microfluidic cavity [23]. Climent and Feist [24] reported the experimental modification of the chemical reactivity of  $S_N2@Si$  reactions under vibrational strong coupling (VSC) in microfluid cavities. They showed that certain modes of their studied  $S_N2@Si$  reactions are highly coupled among the different fragments and that it is difficult to find normal isolated stretching modes to elucidate VSC experiments. Hansen et al. [25] applied the Activation Strain Model [26] extended with EDA [27] to study the enhanced reactivity of  $\alpha$ -nucleophiles and concluded that orbital overlap between the HOMO lobe on the  $\alpha$ -nucleophile and the substrate influences Nu  $\cdots$  substrate Pauli repulsion.

Chemical bond analyses generally target  $[Nu \cdots AR_3X]^-$  or  $[NuAR_3 \cdots X]^-$  interactions missing out the analysis of spectator bonds A—R bonds. Therefore, to advance our understanding of the S<sub>N</sub>2 mechanism in the gas phase, we investigated in this work the stationary points of the Cl<sup>-</sup> + AR<sub>3</sub>Cl reactions (with A = C and Si, and R = H, Me, Et, Cl, and F) with chemical bond analyses applied to the A—R bonds. An important aspect of this study was to evaluate the nature of A—R bonding in particular at the TS and TC pentacoordinate structures. Quantum Theory of Atoms in Molecules (QTAIM) [20], Overlap Model (OP) [28], and Local Vibrational Mode (LVM) theory [29] were used to further study the A—R bond interaction along the S<sub>N</sub>2@C and S<sub>N</sub>2@Si profiles at different stages of the S<sub>N</sub>2@A reactions.

# 2. Methodology

The three bond analysis methods are briefly described in the following subsections.

# 2.1. The chemical bond overlap model

OP was recently introduced [28,30] and is implemented in our ChemBOS code. OP can be successfully applied to analyze chemical bonds in compounds of different complexity, e.g., simple diatomic molecules, coordination compounds, and solid-state systems [30–33]. It is possible to interpret 4*f*-4*f* transitions intensities in terms of covalency [32,34]. Details can be found in [28].

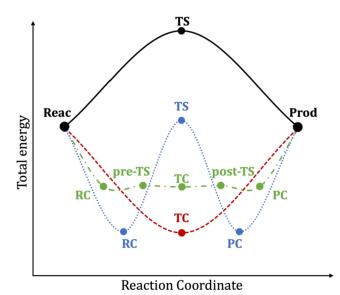


Fig. 2. Typical profiles for thermoneutral  $S_N2$  reactions in gas phase: single-peak (solid black line), single-well (dash red line), double-well (dot blue line), and triple-well (dash and dot green line). Reac = reactant, RC = reactant complex, TC = transitional complex, TS = transition state, PC = product complex, and Prod = product.

Given a localized molecular orbital (LMO) *l*, attributed to a chemical bond A—B in a molecular system, overlap population density maps can be generated from

$$\rho_{\rm OP}^{l}(\overrightarrow{r}) = 2\sum_{i\in \mathbf{A}}^{m} \sum_{j\in \mathbf{B}}^{m} c_{li}c_{lj}\varphi_{li}(\overrightarrow{r})\varphi_{lj}(\overrightarrow{r})$$
(1)

where *m* is the number of atomic orbitals or basis functions,  $\varphi_{li}(\vec{r})$  are the primitive or contracted functions, and  $c_{li}$  are the coefficients of the LMOs.

The overlap population  $p^l$  is then obtained from the integration of Eq. (1) for a  $N_{occ}$  LMO occupation number,

$$p^{l} = N_{occ} \cdot 2 \sum_{i \in A} \sum_{j \in B} c_{ii} c_{ij} S^{l}_{ij}$$
<sup>(2)</sup>

where  $S_{ii}^{l}$  are the overlap integrals.

The anisotropic overlap polarizability  $\overline{\alpha}_{OP}$  can be calculated from the tensor

$$\overline{\alpha}_{OP} = \frac{1}{3} \left( \alpha_{OP}^{xx} + \alpha_{OP}^{yy} + \alpha_{OP}^{zz} \right)$$
(3)

where each term is calculated by

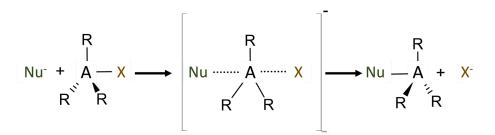


Fig. 1. Walden inversion mechanism of the  $S_N2$  reaction. R = substituent groups;  $Nu^- =$  nucleophile, X = leaving group, and A = central atom.

$$\alpha_{OP}^{\kappa\lambda} = -\frac{2}{F_{\kappa}} \left( \sum_{i \in B \atop i \in B}^{m} \sum_{j \in A \atop j \in A}^{n} c_{i}^{\cdot} c_{j}^{\cdot} \langle \varphi_{i} | \lambda | \varphi_{j} \rangle - \sum_{i \in A \atop i \in B}^{m} \sum_{j \in A \atop j \in A}^{n} c_{i}^{0} c_{j}^{0} \langle \varphi_{i} | \lambda | \varphi_{j} \rangle \right)$$
(4)

where  $c_i^0$  and  $c_i^{\prime}$  are the expansion coefficients for the unperturbed and perturbed LMO, that can be obtained by a self-consistent field (SCF) procedure followed by a localization unitary transformation. In Eq. (4),  $\kappa$ is the direction in which the electrical field  $F_r$  is applied, and can be equals to x, y or z. Similarly,  $\lambda$  (than can be x, y or z) defines the direction in which the dipole (induced by  $F_{\kappa}$ ) will be inspected.

The so-called intra-overlap Coulomb repulsion, that occurs within the overlap of the chemical bond described by the LMO l, is defined as

$$\mathbf{J}_{OP}^{l,l} = \int \rho_{OP}^{l}(\overrightarrow{r}_{12}) r_{12}^{-1} \rho_{OP}^{l}(\overrightarrow{r}_{2}) d\overrightarrow{r}_{1} d\overrightarrow{r}_{2}$$
(5)

where  $\rho_{\rm OP}^l$  is described by Eq. (1). J<sup>l</sup><sub>OP</sub>, from here denoted by J<sup>intra</sup><sub>OP</sub> is a 6D integration, that is solved numerically [35].

#### 2.2. The local vibrational modes theory

LVM, originally introduced by Konkoli and Cremer [29,36], provides through its local stretching force constant, a unique measure of the intrinsic strength of a chemical bond [37]. The local mode force constant, for a local mode *n*, is obtained via Eq. (6) [29]:

$$k_{n}^{a} = \left(d_{n}K^{-1}d_{n}^{+}\right)^{-1} \tag{6}$$

where  $d_n$  is the *n* th normal mode in internal coordinates and diagonal matrix K is the force constant in normal mode coordinates. Both quantities can be obtained via a standard normal mode analysis, available in most of the common quantum chemistry packages [29].

#### 2.3. The quantum theory of atoms in molecules

OTAIM was used in this work to assess the covalent character of the A-R bonds applying the Cremer-Kraka criterion [38], which is based on the local energy density  $H(r_{BCP})$  evaluated at the bond critical point BCP

$$H(r_{BCP}) = G(r_{BCP}) + V(r_{BCP})$$
<sup>(7)</sup>

where  $G(r_{BCP})$  is the kinetic energy and  $V(r_{BCP})$  is the potential energy, both calculated at the BCP position  $r_{BCP}$ . Given that  $G(r_{BCP})$  is destabilizing (and positive) and  $V(r_{BCP})$  is stabilizing (and negative), the Cremer-Kraka criterion states that  $H(r_{BCP}) < 0$  indicates a covalent interaction, while  $H(r_{BCP}) > 0$  indicates an electrostatic interaction. We also used the Laplacian  $\nabla^2 \rho(\mathbf{r}_{BCP})$ , that holds information about local charge concentration ( $\nabla^2 \rho(\mathbf{r}_{BCP}) < 0$ ) or depletion ( $\nabla^2 \rho(\mathbf{r}_{BCP}) > 0$ ) [20].

# 3. Computational procedure

Geometry optimization and frequency calculations were performed with the  $\omega$ -B97X-D [39] functional, and SPK-TZP basis set [40] using GAMESS [41] and Gaussian16 [42]. The OTAIM, OP, and LVM descriptors were calculated with the Multiwfn [43], ChemBOS [35], and LmodeA [44] software, respectively. For the overlap properties, the Pipek-Mezey molecular orbital localization [45] was applied.

#### 4. Results and discussion

The energy profiles for the reactions studied in this work leading to 44 stationary points are summarized in Fig. 3 (10 reacts, 10 prods, 5 RCs, 5 PCs, 2 pre-TSs, 2 post-TSs, 5 TSs, and 5 TCs); energies are given relative to reactants. For the S<sub>N</sub>2@C reactions, reactants and products are separated by a pentacoordinate transition state (TS), corresponding to a saddle point on the PES, with a single-peak profile for R=F and Cl, and a double-well profile for R=H, Me, and Et. In contrast, for  $S_N2@Si$ reactions, a stable pentacoordinate transition complex (TC) is observed for all structures with R=H, Me, Et, F, and Cl, with a single-well profile for R=H, F, and Cl, and a triple-well for R=Me and Et with a pre-TS and a post-TS in each case.

These findings are in line with earlier work [12], as well as the relative energies shown in Fig. 3 [22]. The stabilization in the doublewell and triple-well profiles seems to be a consequence of an ion-dipole interaction. The TS/TC formation found in this work can be associated with steric effects matching the findings of Bento and Mathias [12].

All stationary points were characterized as local energy minima (Reac, Prod, RC, PC, and TC) or transition states of first order (pre-TS, TS, and post-TS) with the normal mode of the imaginary harmonic frequency being associated with the nucleophilic attack path. The A-R bonds in Reacs, TCs, and TSs were analyzed using the three bond analyses described above. The results are summarized in Table 1, depicting overlap properties, QTAIM descriptors and local modes force constants.

In the following main bond property trends of the A-R spectator bonds are discussed focusing on two major aspects: i) how does each analysis methodology characterize bonding in C-R and Si-R bonds, and ii) what is the relationship between the three analysis methods applied in this work.

Fig. 4 visualizes the main results shown in Table 1. The general trend

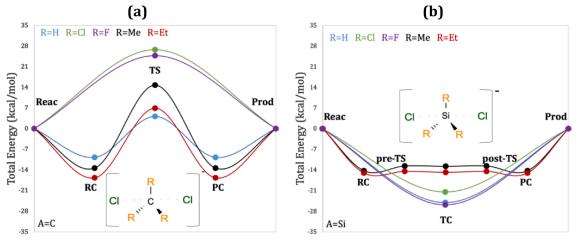


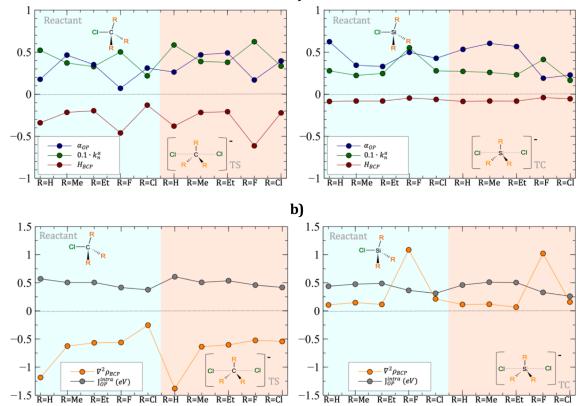
Fig. 3. Energy profiles obtained for the  $Cl^- + AR_3Cl S_N 2$  reactions, where A = C, Si, and R = H, F, Cl, Me, and Et. Reac = reactants, RC = reactant complex, pre-TS = pre transition state, TS = transition state, post-TS = post transition state; TC = transition complex, PC = product complex, and Prod = products.

#### Table 1

Spectator bonds (A—R) properties obtained for the Reac, TC, and TS structures: bond distance D (in Å); overlap density  $\rho_{OP}$  (in e), polarizability  $\alpha_{OP}$  (in Å<sup>3</sup>), and repulsion  $J_{OP}^{intra}$  (in  $E_h$ ); density  $\rho_{BCP}$  (in  $e/a_0^3$ ), Laplacian of the density  $\nabla^2 \rho_{BCP}$  (in  $e/a_0^5$ ), and local energy density  $H_{BCP}$  (in  $E_h/a_0^3$ ) at the critical point; local mode force constant  $k_n^a$  (in mDyn/Å).

					Overlap properties			QTAIM			LVM
Entry		System	Bond	D	$\rho_{OP}$	$\alpha_{OP}$	J <sup>intra</sup>	$\rho_{BCP}$	$\nabla^2 \rho_{BCP}$	H <sub>BCP</sub>	$k_n^a$
1	Reactant	CH <sub>3</sub> Cl	C–H	1.091	0.890	0.177	0.570	0.292	-1.184	-0.340	5.218
2		CMe <sub>3</sub> Cl	C–C	1.530	0.815	0.465	0.505	0.249	-0.625	-0.216	3.708
3		CEt <sub>3</sub> Cl	C–C	1.541	0.817	0.352	0.505	0.239	-0.568	-0.197	3.284
4		CF <sub>3</sub> Cl	C–F	1.338	0.647	0.070	0.415	0.296	-0.562	-0.461	5.022
5		CCl <sub>3</sub> Cl	C–Cl	1.782	0.713	0.311	0.376	0.196	-0.255	-0.130	2.179
6	Transition state	$Cl \cdots CH_3 \cdots Cl$	C–H	1.075	0.924	0.263	0.607	0.307	-1.381	-0.380	5.845
7		$Cl \cdots CMe_3 \cdots Cl$	C–C	1.490	0.795	0.468	0.507	0.250	-0.637	-0.217	3.885
8		$Cl \cdots CEt_3 \cdots Cl$	C–C	1.492	0.826	0.491	0.536	0.245	-0.604	-0.208	3.790
9		$Cl \cdots CF_3 \cdots Cl$	C–F	1.290	0.664	0.169	0.459	0.351	-0.524	-0.614	6.240
10		$\mathrm{Cl}\cdots\mathrm{CCl}_3\cdots\mathrm{Cl}$	C–Cl	1.688	0.741	0.395	0.417	0.256	-0.543	-0.222	3.341
11	Reactant	SiH <sub>3</sub> Cl	Si–H	1.487	0.851	0.624	0.438	0.125	0.106	-0.086	2.779
12		SiMe <sub>3</sub> Cl	Si–C	1.881	0.854	0.344	0.475	0.124	0.147	-0.081	2.232
13		SiEt <sub>3</sub> Cl	Si–C	1.898	0.879	0.330	0.488	0.122	0.115	-0.081	2.459
14		SiF <sub>3</sub> Cl	Si–F	1.583	0.592	0.498	0.365	0.147	1.085	-0.046	5.507
15		SiCl <sub>3</sub> Cl	Si–Cl	2.040	0.682	0.427	0.312	0.107	0.213	-0.062	2.782
16	Transition complex	$Cl \cdots SiH_3 \cdots Cl$	Si–H	1.484	0.882	0.533	0.460	0.125	0.114	-0.086	2.713
17		$Cl \cdots SiMe_3 \cdots Cl$	Si–C	1.903	0.908	0.604	0.511	0.124	0.119	-0.081	2.585
18		$Cl \cdots SiEt_3 \cdots Cl$	Si–C	1.921	0.904	0.567	0.504	0.119	0.066	-0.082	2.309
19		$Cl \cdots SiF_3 \cdots Cl$	Si–F	1.621	0.584	0.190	0.329	0.142	1.021	-0.040	4.129
20		$Cl \cdots SiCl_3 \cdots Cl$	Si–Cl	2.128	0.637	0.228	0.260	0.098	0.158	-0.055	1.649





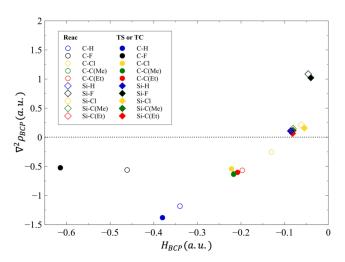
**Fig. 4.** Graphical representation of the spectator bonds (A–R) properties for the Reac (left side, light blue background), TC/TS (right side, light red background) structures: a) overlap polarizability  $\alpha_{OP}$  (in Å<sup>3</sup>), local energy density  $H_{BCP}$  (in  $E_h/a_0^3$ ) at the critical point, and repulsion  $J_{DP}^{inpra}$  (in  $E_h$ ); density  $\rho_{BCP}$  (in  $e/a_0^3$ ), Laplacian of the density  $\nabla^2 \rho_{BCP}$  (in  $e/a_0^5$ ), and; local mode force constant  $k_n^a$  (in mDyn/Å) – scaled by 0.1 for visualization purposes.

for the C—R bonds is an increase of  $H_{BCP}$ , i.e., reduced covalent character is in line with a decrease in the local force constant  $k_n^a$ , which indicates a decrease in the bond strength. An increase of  $\alpha_{OP}$  indicates that the overlap density is being more polarizable. For Si—R bonds, the  $H_{BCP}$  values are smaller than those of the C—R and closer to 0.0  $E_h/a_0^3$ , reflecting the well-known fact that Si—R bonds are generally less covalent than the C—R bonds [46]. Variations in  $\alpha_{OP}$  and  $k_n^a$  show a parallel trend for Si—R bonds of the reactants, and an opposite trend for the TC structures.

Fig. 4b depicts the  $\nabla^2 \rho_{BCP}$  and  $J_{OP}^{intra}$  properties for the studied bonds. C—R bonds exhibit an inverse relationship between  $\nabla^2 \rho_{BCP}$  and  $J_{OP}^{intra}$ . This is expected, given that a less negative  $\nabla^2 \rho_{BCP}$  indicates a charge depletion at the critical point, and that chemical bonds with a charge depletion tend to have smaller  $J_{OP}^{intra}$  (less overlap density intra-repulsion) values [28].

Fig. 5 visualizes the relationship between  $H_{BCP}$  and  $\nabla^2 \rho_{BCP}$ . All C—R chemical bonds (empty and filled circles in Fig. 5) have negative  $\nabla^2 \rho_{BCP}$  values while the Si—R values are positive. The  $H_{BCP}$  for Si—R bonds are in range  $-0.1 E_h/a_0^3 < H_{BCP} < 0.0 E_h/a_0^3$ , while for C—R bonds are in range  $-0.7 E_h/a_0^3 < H_{BCP} < -0.1 E_h/a_0^3$  again reflecting the fact that C—R bonds are generally more covalent than their higher homologues, concentrating more electron density in the interatomic region. There is no significant difference between Si—R bonds are generally more covalent in the TS-TC, whereas QTAIM indicates that C—R bonds are generally more covalent in the TS/TC than in the Reac structure (filled circles with more negative  $H_{BCP}$  and  $\nabla^2 \rho_{BCP}$  values than empty circles). C—F bonds in reactant and TS are the exceptions in which the more covalent bond (smaller  $H_{BCP}$ ) does not exhibit more density concentration in the bond region (smaller  $\nabla^2 \rho_{BCP}$ ), see Fig. 5, when compared with C—C bonds in reactants.

Recently, Blokke and coworkers [47] applied the Activation Strain Model [26] extended with EDA [27] to study archetypal H<sub>3</sub>C—CH<sub>3</sub>, H<sub>3</sub>C—F, and H<sub>3</sub>C—Cl molecules, among others, revealing that from C—F to C—Cl the expected bond weakening effect is caused primary by Pauli repulsion, not because of decreasing electronegativity difference. These authors pointed to the increase in effective atomic size from F to Cl and its more extended atomic orbitals which lead to an increase in occupied molecular orbitals overlap (S) [47]. Similarly, the results reported in the present work corroborate these trends. In C—F bonds (Entries 4, and 9 in Table 1), despite having smaller equilibrium bond distances than C—Cl bonds (Entries 5, and 10 in Table 1), the  $\rho_{OP}$  values are smaller. A more diffuse valence shell interaction causes the overlap polarizability  $a_{OP}$ values for the C—Cl bonds to be greater than that of the C—F



**Fig. 5.** QTAIM descriptors  $\nabla^2 \rho_{BCP}$  (in  $e/a_0^5$ ), and  $H_{BCP}$  (in  $E_h/a_0^3$ ) of the spectator chemical bonds A–R in reactants AR<sub>3</sub>Cl and TS/TC Cl  $\cdots$  AR<sub>3</sub>  $\cdots$  Cl, (X = C, Si).

counterparts. As a consequence, despite being more covalent (smaller  $H_{BCP}$ ), C—F bonds exhibit a bond charge concentration ( $\nabla^2 \rho_{BCP}$ ) similar to C—C bonds, and smaller overlap density  $\rho_{OP}$  and polarizability  $\alpha_{OP}$ .

The overlap polarizability  $\alpha_{OP}$  concept introduced by Malta and coworkers in 2002 [30] revealed an approximated analytical expression for  $\alpha_{OP}$ , in a relationship with overlap integral  $\rho$ , bond distance R, and HOMO-LUMO gap  $\Delta \varepsilon$  quantities of type:  $\alpha_{OP} R^2 \rho^2 / \Delta \varepsilon$ . Interestingly, Blokker and coworkers [47] pointed to a relationship between the orbital interaction  $\Delta E_{oi}$  (that accounts for electron-pair interaction, charge transfer, and polarization), and the orbital overlap S (direct), and the orbital–orbital energy gap (inverse).

EDA [17–19] has revealed that the steric (Pauli) repulsion and electronic effects can explain the distinct shapes of the PES in the gas phase  $S_N 2@A$  reactions. Our results support these findings, given that the steric hindrance experienced by Cl<sup>-</sup> in  $S_N 2@C$  reactions can be viewed as a consequence of the greater covalent nature of C—R bonds, that concentrates density along the chemical bond more efficiently than in Si—R. Also, C—R bonds are generally shorter than Si—R bonds, therefore the CR<sub>3</sub> core is more compact than the SiR<sub>3</sub> core.

Distance D and  $k_n^a$  values for the C—R and Si—R bonds in reactants and TS (see Table 1, entries 1–10, and Fig. 6a) are qualitatively connected via a generalized Badger type [48] relationship. It is noteworthy that the shorter bond is not always the stronger bond, as documented in literature [48–51]. As expected, C—R bonds investigated in this work are stronger (larger  $k_n^a$  values) than their Si—R counterparts, see Fig. 6b. Also, the  $k_n^a$  values for C—R bonds are larger in reactants than in TS/TC stationary points. The results found in this work corroborate the literature [47] values of bond dissociation energies (BDE) increasing from H<sub>3</sub>C—CH<sub>3</sub> to H<sub>3</sub>C—F and decreasing from H<sub>3</sub>C—F to H<sub>3</sub>C—CL.Fig. 6.

Following this trend, the overlap properties point to C—R bonds with greater density concentration than for Si—R bonds, as represented in overlap density maps in Fig. 7. Also, the overlap Coulomb repulsion energy  $(J_{OP}^{imra})$  is generally greater in C—R than in Si—R in AR<sub>3</sub>Cl reactant systems, as can be seen in Table 1 and Fig. 4. As can be seen in Fig. 7, the overlap density maps of C—R bonds in both Reac and TS show a density concentration when compared with Si—R bonds.

It has to be pointed out that the overlap density maps must be analysed together with the overlap properties shown in Table 1; e.g., A—F bonds appear to have a large overlap density concentration, but if the corresponding  $\rho_{OP}$  and  $J_{OP}^{intra}$  values (entries 4, 9, 14, and 19 in Table 1) are taken into consideration, one sees that the overlap density and intrarepulsion are smaller than for other bonds.

# 5. Conclusions

Cl<sup>-</sup> + AR<sub>3</sub>Cl reactions (where A=C and Si, and R=H, Me, Et, Cl, and F) were investigated and A-R chemical bond properties were calculated using QTAIM, OP, and LVM. The different chemical bond analysis tools used in this work converge to the conclusion (in line with current literature) that the steric hindrance experienced by Cl<sup>-</sup> in S<sub>N</sub>2@C reactions can be viewed as a consequence of the greater covalent nature of C-R bonds, that concentrate density along the chemical bond more efficiently than in Si-R. In summary,  $H_{BCP}$  values for C-R bonds are more negative than for Si—R bonds.  $\nabla^2 \rho_{BCP}$  are negative for C—R bonds and positive for Si-R bonds. The C-F bonds in reactant and TS were found to be exception cases in  $H_{BCP}$  vs.  $\nabla^2 \rho_{BCP}$  observed trend, what is connected with the smaller  $\rho_{OP}$  and  $\alpha_{OP}$  (compared to C—Cl counterpart) values. The local vibrational analysis reveals that distance D and  $k_n^a$ values for the studied bonds in reactants and TS/TC stationary points are qualitatively connected via a generalized Badger type relationship. The OP model points to a more efficient overlap density concentration in C-R bonds, when compared with Si-R ones, also supporting the QTAIM and LVM indicatives.

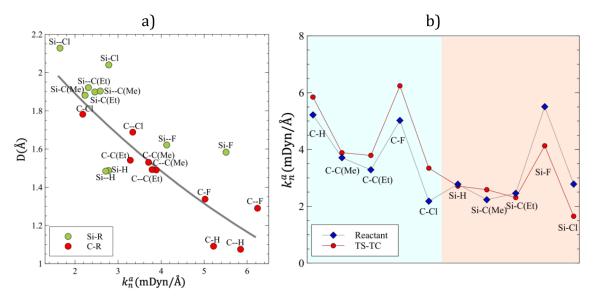


Fig. 6. (a) A–R distance D vs. Local stretching force constant  $k_n^a$  (in mDyn/Å), and (b)  $k_n^a$ ()A–R) for reactants and TS/TC structures.

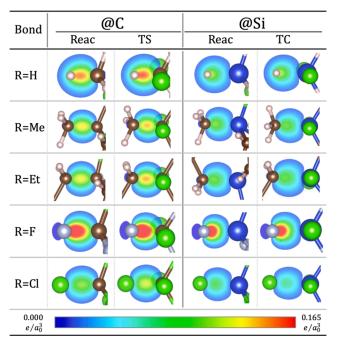


Fig. 7. Overlap density maps for C—H and Si—H bonds in  $\rm CH_3Cl$  and  $\rm SiH_3Cl$  structures.

# CRediT authorship contribution statement

**Carlos V. Santos-Jr:** Software, Methodology, Formal analysis, Investigation, Writing – original draft. **Miguel A. F. de Souza:** Methodology, Writing – original draft, Supervision. **Elfi Kraka:** Methodology, Writing – original draft, Writing – review & editing, Funding acquisition. **Renaldo T. Moura Jr:** Conceptualization, Software, Methodology, Formal analysis, Investigation, Writing – original draft, Writing – review & editing.

#### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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