Doublet Ground State in a Vanadium(II) Complex: Redox and Coordinative Noninnocence of Tripodal Ligand Architecture

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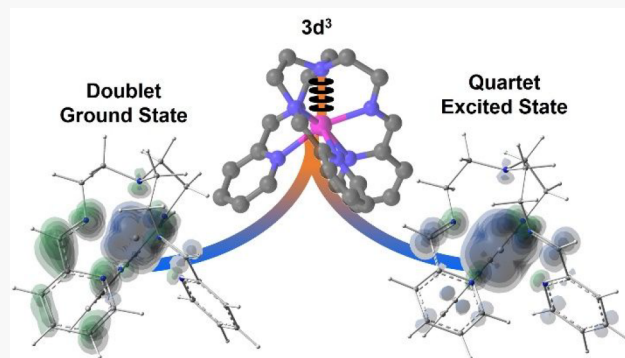Read Online

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ABSTRACT: We report on the geometric and electronic structures of a series of $V^{2+/3+}$ tren-bridged iminopyridine complexes [tren = tris(2-aminoethyl)amine] that enable the observation of an unexpected doublet ground state for a nominally $3d^3$ species. Tren undergoes condensation reactions with picolinaldehyde or methyl-6-formylnictonate to form the respective tripodal ligand sets of $(py)_3$tren and $(5-CO_2Mepy)_3$tren. The $(py)_3$tren ligand is coordinated to V^{2+} and V^{3+} metal centers to form complex salts $[\mathbf{1\text{-}H}](OTf)_2$ and $[\mathbf{1\text{-}H}](OTf)_3$, respectively ($OTf^- = CF_3SO_3^-$). For $[\mathbf{1\text{-}H}]^{2+}$, strong metal−ligand π-covalency with respect to the V^{2+} ($3d^3$) and iminopyridine ligands weakens its interelectronic repulsion. For $[\mathbf{1\text{-}H}]^{3+}$, the bridgehead nitrogen of the tren scaffold forms a seventh coordinate covalent bond with a V^{3+} ($3d^2$) metal center. The coordination of $(5-CO_2Mepy)_3$tren to a V^{2+} metal center results in the redox noninnocent and heptacoordinate compound $[\mathbf{1\text{-}CO_2Me}](OTf)_2$ with a doublet ($S = \frac{1}{2}$) ground state that we support with magnetic susceptibility and spectroscopy measurements. The complexes are uniformly characterized experimentally with single-crystal X-ray diffraction, electronic absorbance, and electrochemistry, and electronic structures are corroborated by computational techniques. We present a new computational procedure that we term the spin-optimized approximate pair (SOAP) method that enables the visualization and quantification of electron−electron interactions.
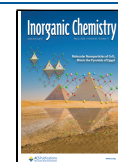

INTRODUCTION

An increasing amount of effort has been directed toward the integration of first-row transition metal ions into photocatalytic processes that are traditionally associated with precious congeners.[1−10] Earth-abundant coordination complexes strain confinement by the chemical properties of their noble metal analogues. Attributable to the weakened covalency of their metal−ligand bonds, 3d complexes possess many low-lying excited states that can have distinct electronic spin multiplicities. While these near-degenerate states are advantageous with respect to molecular magnetism,[11−14] they convolute direct mechanistic pathways for catalytic transformations.[15−24] Octahedral Cr^{3+} polypyridyl complexes have been demonstrated to facilitate photocatalytic and photoredox processes.[25−29] The low-lying ^2E excited state of d^3 complexes is classified as a spin-flip transition whose energy is approximately independent of ligand field strength but proportional to the interelectronic repulsion between its metal-centered electrons.[30,31] Weakening the interelectronic repulsion, via metal−ligand covalency, has been recently demonstrated by Wenger and Piguet[32] and Heinze[33] as a robust design principle to impact the energy of the ^2E state of a Cr^{3+} complex. The ability to modulate the energies of states of different multiplicities in first-row transition metals impacts their ability to perform catalysis.[34]

Despite potential environmental advantages, photocatalysis has not been extended to the isoelectronic (d^3) V^{2+} analogue due to greater π-bonding with its polypyridyl ligand set that results in a low-lying doublet state of both metal-centered (MC) and metal-to-ligand charge transfer (MLCT) character.[35−37] The resonance of this $^2MC/^2MLCT$ state of V^{2+} polypyridyls significantly stabilizes its doublet excited state. We report the observation of a doublet ground state in a V^{2+} tren-iminopyridine system whose ligand environment stabilizes the doublet state through several dramatic changes in the covalency of the system that are illustrated in Figure 1 and detailed herein. The characterization of the doublet state of $3d^3$ complexes has only previously been accessible through transient spectroscopic and computational techniques.[38−41]

Herein, we combine spectroscopic and theoretical methods to address the unique electronic and geometric structures of our vanadium tripodal complexes. We present a model that
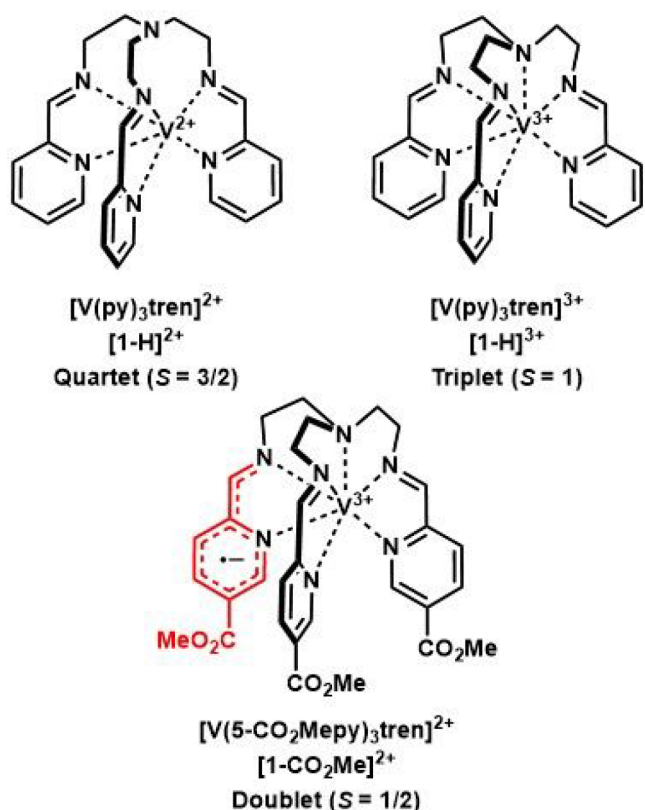
Figure 1. Illustrated geometric structures, chosen notation and abbreviations, and spin multiplicities of the three vanadium tren-iminopyridine complexes detailed in our report. Tren refers to tris(2-aminoethyl)amine, whose C_3-symmetric arrangement of primary amines undergoes condensation reactions with picolinaldehyde or methyl-6-formylnicotinate to form the respective tripodal ligand set of $(py)_3$tren and $(5-CO_2Mepy)_3$tren. The $(py)_3$tren ligand set is coordinated to V^{2+} and V^{3+} metal centers to form complexes $[1\text{-}H]^{2+}$ and $[1\text{-}H]^{3+}$, respectively. The coordination of $(5-CO_2Mepy)_3$tren to a V^{2+} metal center results in redox noninnocent complex $[1\text{-}CO_2Me]^{2+}$, whose reduced iminopyridine group is colored red.

provides visual and quantitative insight into the discrete electron−electron interactions of open-shell systems. Our spin-optimized approximate pair (SOAP) method cultivates an understanding through generation of three-dimensional (3D) graphics of the structural properties that impact the energy of spin-flip excited states. This has been a recent focus of d^2 and d^3 transition metal systems.[6,42−46] Our approach is a natural extension of the general valence bond (GVB) treatment of small open-shell molecules developed by Goddard and Bobrowicz[47−49] and finding further articulation with research from Dunning and Cooper.[50−53] Our SOAP method can be extended to systems irrespective of the total number of unpaired electrons or competing magnetic interactions, to provide new insights into the root causes of ferro- and antiferromagnetic coupling.

■ RESULTS AND DISCUSSION

V^{2+}-Iminopyridine: Metal−Ligand and π-Covalency. First, we detail the strong metal−ligand π-covalency between the V^{2+} metal center and the iminopyridine ligand set in complex $[1\text{-}H]^{2+}$, the compound notation being detailed in the legend of Figure 1. The subsequent reduction in the

interelectronic repulsion between the unpaired electrons on the paramagnetic metal center decreases the gap between the quartet ground and doublet excited states. This is manifested as a 0.09 eV reduction in exchange with respect to its bipyridyl analogue $[V(bpy)_3]^{2+}$ (bpy = 2,2′-dipyiryl), which helps facilitate the doublet ground state that we will detail with respect to $[1\text{-}CO_2Me]^{2+}$ below.

Diffraction quality single crystals of $[1\text{-}H]^{2+}$ are obtained by the slow diffusion of diethyl ether into acetonitrile, and the structure is shown in Figure 2. Structures of $[1\text{-}H]^{3+}$ and $[1\text{-}$
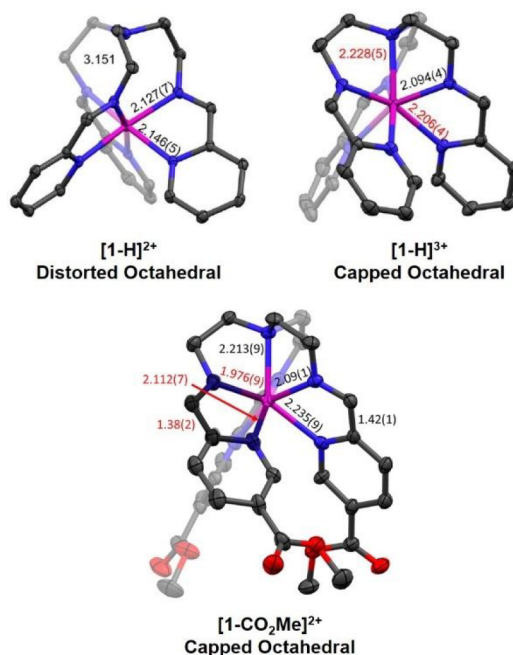

Figure 2. Experimental molecular structures of $[1\text{-}H]^{2+}$, $[1\text{-}H]^{3+}$, and $[1\text{-}CO_2Me]^{2+}$, as determined by X-ray crystallography, where hydrogen atoms, anions, and co-crystallized solvent molecules have been omitted for the sake of clarity. The notation of these complexes is detailed in the legend of Figure 1. Thermal ellipsoids are set to 40%. Assigned coordination geometries, pertinent bond distances, and interatomic separations (in angstroms) are provided. Pink, blue, black, and red ellipsoids represent vanadium, nitrogen, carbon, and oxygen, respectively.

$CO_2Me]^{2+}$ are also included in Figure 2. While the tren-iminopyridine ligand environment is known for third-row transition metal complexes, these are the first structures reported for $V^{2+/3+}$.[54] Key bond distances and structural parameters are listed in Table 1. The imine and pyridine donor atoms of $[1\text{-}H]^{2+}$ possess nearly equivalent bond lengths. The structure displays a $V^{2+}-N_{bridge}$ distance that is intermediate between a coordinate covalent bond and a noncovalent interaction, when scaled with respect to the sum of their covalent radii $(R\sum cov^{-1})$ that is a factor of 1.41.[55]

The separation between the bridgehead nitrogen (N_{bridge}) of the tren scaffold and the transition metal center is established to be dependent on the oxidation and spin state.[57−62] We have previously reported that the N_{bridge} conformation impacts the local electronic structure of its transition metal center with respect to the photophysics of Cr^{3+} and the molecular magnetism of Co^{2+} complexes.[63,64] The value for the continuous shape measure (CShM) of a six-coordinate octahedral environment is termed $S(OC-6)$, where a greater

Table 1. Key Metric Parameters of Detailed Vanadium Complexes[a]

	$[1\text{-H}]^{2+}$	$[1\text{-H}]^{3+}$	$[1\text{-CO}_2\text{Me}]^{2+}$
V−N$_{\text{bridge}}$ (Å)	3.151	2.228(5)	2.213(9)
V−N$_{\text{imine}}$ (Å)	2.127(7)	2.094(4)	2.10(1)
V−(N$_{\text{imine}}$)$^-$ (Å)			1.976(9)
V−N$_{\text{pyridine}}$ (Å)	2.146(5)	2.206(4)	2.20(4)
V−(N$_{\text{pyridine}}$)$^-$ (Å)			2.112(7)
S(OC-6) (unitless)	1.480	2.503	2.521
S(COC-7) (unitless)	3.055	0.792	0.943

[a]Bond lengths are reported in angstroms, and S(OC-6) and S(COC-7) refer to the continuous shape measure (CShM) values for octahedral and capped octahedral coordination geometries, respectively.[56]

value indicates greater distortion from its ideal shape that is determined with SHAPE version 2.2.[56] The CShM values do not possess units for the coordination geometries considered herein because it is referenced to an ideal shape. The S(OC-6) of $[1\text{-H}]^{2+}$ is 1.480, and its primary coordination sphere is best classified as distorted octahedral. This is increased with respect to our previous analysis of the polypyridyl analogue $[\text{V(bpy)}_3]^{2+}$ whose density functional theory (DFT)-optimized structure has an S(OC-6) of 1.27.[37] This suggests greater metal−ligand π-covalency for the iminopyridine coordination environment because of the loss of degeneracy in the t_{2g} orbitals as the complex is distorted toward a trigonal, C_3-symmetric environment, where $t_{2g} \rightarrow e + a^*$.

The electronic absorbance spectra of $[1\text{-H}]^{2+}$ and its polypyridyl analogue, $[\text{V(bpy)}_3](\text{OTf})_2$ (OTf$^-$ = trifluoromethylsulfonate), are included in Figure 3. The absorbance

Figure 3. Electronic absorption spectra of $[1\text{-H}]^{2+}$, $[1\text{-H}]^{3+}$, and $[1\text{-CO}_2\text{Me}]^{2+}$, recorded in acetonitrile. The corresponding spectrum of $[\text{V(bpy)}_3](\text{OTf})_2$[35] is provided as a reference to the impact of the iminopyridine ligand environment.

data of $[1\text{-H}]^{3+}$ and $[1\text{-CO}_2\text{Me}]^{2+}$ are also included in Figure 3 and are detailed further below. $[1\text{-H}]^{2+}$ exhibits a broad electronic absorbance at approximately 750 nm that we assign as a metal-to-ligand charge transfer (MLCT). This is consistent with the transition's large molar absorptivity (ε) and the small ionization potential of a V^{2+} metal center.[36,65] The MLCT is red-shifted by 0.28 eV with respect to the

polypyridyl analogue, which we attribute to the higher electron affinity of the iminopyridine with respect to the bipyridine ligand set. Our assignment is further corroborated by time-dependent density functional theory (TD-DFT),[66] whose natural transition spin density plots (NT$\rho^{\alpha,\beta}$)[37] are included in Table S4 and Figure S13.

The computed electronic ground state of $[1\text{-H}]^{2+}$ is a quartet ($S = {}^3/_2$), the $^4A_2$ state common for d^3 systems in a nominally octahedral coordination environment.[31] The lowest-lying doublet state is ^2E that our detailed multideterminant DFT computation assigns as 0.96 eV above the ground state for $[1\text{-H}]^{2+}$.[67] We note that the A_2 and E term symbols are preserved for the C_3-symmetric environment of the tripodal ligand architecture in our detailed V^{2+} complexes. Comparatively, we have reported that $[\text{V(bpy)}_3]^{2+}$ should have a ^2E excited state of 1.10 eV.[37] We reference the spin density of ^2E of $[1\text{-H}]^{2+}$ to its $^4A_2$ ground state in our relative spin density plot ($\Delta\rho^{\alpha,\beta}$) in Figure 4.[37] While the ^2E state is traditionally
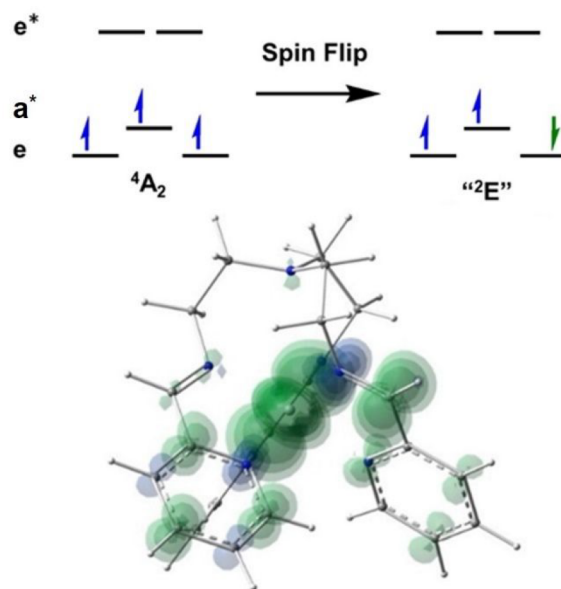

Figure 4. Electronic configuration of the lowest-lying quartet and doublet states for C_3-symmetric d^3 complexes, where the quotation marks denote a single-determinant description in the context of DFT (top). The a* orbital is destabilized by 0.22 eV with respect to the lowest-lying e orbitals, as determined from the orbital energies of the APFD-calculated quartet state. APFD spin density of the lowest-lying doublet state of the vanadium complex of $[1\text{-H}]^{2+}$ (bottom), referenced against its lowest-lying quartet state for our relative spin density ($\Delta\rho^{\alpha,\beta}$) plot. Green denotes a loss of α and an increase of β electron spin density in the lowest-lying quartet and doublet states, respectively. The contours are shown at 0.003, 0.006, and 0.012 electron bohr^{-2}.

characterized as a metal-centered spin flip, the $\Delta\rho^{\alpha,\beta}$ plot shows that it is significantly delocalized onto the iminopyridine ligand set for $[1\text{-H}]^{2+}$.

The impact of metal−ligand covalency on the ^2E excited state energy is that it weakens the interelectronic repulsion between the singly occupied metal-centered orbitals.[68−70] The unpaired electrons of the $^4A_2$ state of $[1\text{-H}]^{2+}$ occupy orbitals that are orthogonal and nearly degenerate so that Hund's first rule applies.[71−73] The basis for a quartet ground state is the three favorable exchange interactions between electrons of parallel spin (K_{ab}), detailed in section S7.1 of the Supporting

Information.[74−76] The discussion in section S7.1 similarly addresses the Jahn−Teller distortion of the ^{2}E state and the difficulty of its multideterminant character with DFT methods. The Pauli exclusion principle prohibits electrons of the same spin from occupying the same location, thus weakening a system's interelectronic repulsion via electron correlation. This can be termed Fermi correlation and is addressed in Hartree−Fock theory with respect to the antisymmetry of the Slater determinant.[77] If the unpaired electrons are restricted by symmetry from bonding, then exchange remains as the sole resource of stabilization.

Introducing the Spin-Optimized Approximate Pair (SOAP) Method. To visualize the impact of exchange on state energy differences, we present a model that provides 3D graphical representations of electron−electron interactions that we term the spin-optimized approximate pair (SOAP) method. While our current focus is the exchange interactions, our procedure similarly extends to Coulombic and more complex interelectronic interactions that are detailed in section S6. K_{ab} possesses the form of an electrostatic potential that is detailed in eq 1. While the Coulombic interaction between a pair of electrons occurs with their physically measurable charge densities, K_{ab} is the electrostatic repulsion that occurs in the unmeasurable regions where the orbitals overlap. *Exchange is a classical interaction with a quantum mechanical origin.* Dougherty has previously referred to areas of interference between orbitals as the overlap functional (f_{ab}) that is illustrated in eq 2.[78]

$$K_{ab} = \int\int \frac{\phi_a(1)\phi_b(1)\phi_a(2)\phi_b(2)}{r_{12}} \, dr_1 \, dr_2 \quad (1)$$

$$K_{ab} = \int\int \frac{f_{ab}^{\,2}}{r_{12}} \, dr_1 \, dr_2 \quad (2)$$

In Figure 5, we plot a representative exchange interaction between a pair of singly occupied metal-centered orbitals of [1-H]$^{2+}$ as an electrostatic potential overlaid on the surface density of the respective f_{ab}.[79] The exchange interaction is dominantly metal-centered because there is little overlap between the portions of the respective orbitals that are delocalized onto the ligand set. Here, blue and red represent areas in which the interference between the singly occupied orbitals displays positive and negative coherence, respectively. Because the current orbitals are orthogonal, the amount of in-phase and out-of-phase overlap will be equal.

Overall, K_{ab} is a repulsive interaction whose energetic impact on the system depends on the orientation of the electron spins that is favorable for aligned spins. Metal−ligand π-covalency delocalizes the singly occupied V^{2+} orbitals onto the ligand, thus reducing the f_{ab} overlap functional and the magnitude of their exchange interactions. We note that K_{ab} can be similarly plotted as a surface density, which describes the integrand in eq 1 or 2, providing a quantitative numerical value that is provided in Figure 5. Our SOAP method provides both visual and nearly quantitative values for discrete electron−electron interactions. This is distinct from the information provided in Figure 4 in that it illustrates the magnitude of the spin-flip excited state energy and suggests that exchange interactions can be modulated via established synthetic principles for metal−ligand covalency. We also note that the spin densities of antiferromagnetically coupled systems are nonphysical and cannot be related, without projection methods, to experimental
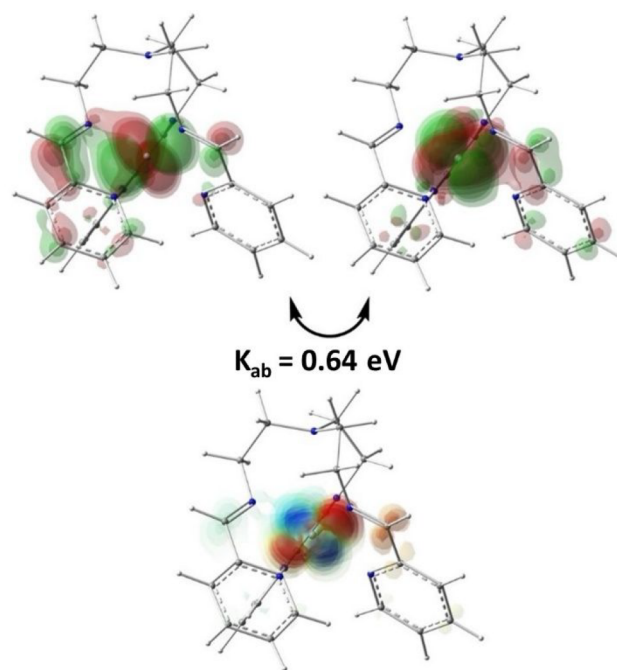

Figure 5. Two lowest-lying singly occupied molecular orbitals of the vanadium complex [1-H]$^{2+}$ that possess e-symmetry with respect to the complex's C_3 coordination environment (top). The orbitals are obtained from the corresponding CASSCF(3,8) wave function. The contours are shown from 0.025 to 0.045 in increments of 0.01 electron bohr^{-2}. Exchange potential and integrand associated with K_{ab} obtained with our SOAP method (bottom). Blue and red denote an exchange potential that results from positive and negative overlap, respectively, between the two lowest-lying singly occupied molecular orbitals shown in the top panel. The f_{ab} contours are shown for 0.001, 0.002, and 0.004 electron bohr^{-2}. The colors range from −0.235 (red) to 0.235 (blue) hartree bohr^{-2}.

results that include electron paramagnetic resonance (EPR) spectroscopy.[80,81] We will detail an antiferromagnetically coupled system with respect to the electronic structure of [1-CO$_2$Me]$^{2+}$ provided below. The orbitals for our SOAP method are generated from the natural orbitals of the complete active space self-consistent field (CASSCF) wave function in which the active space is defined as the metal-centered 3d orbitals and ligand-centered π^* orbitals.[82] The excited state energies of both quartet and doublet multiplicities are further refined with the n-electron valence state perturbation theory (NEVPT2) that accounts for dynamic correlation.[83−85] Our multireference calculations were performed with the ORCA 4.1 electronic structure software package.[86]

Heptacoordination with the Tren Scaffold. Oxidation of [1-H]$^{2+}$ with silver triflate (AgOTf) produces the V^{3+} complex [1-H]$^{3+}$. The structure (Figure 2) features a seventh coordinate covalent bond, now involving the tertiary amine (N$_{bridge}$) of the tren backbone. Relative to that of [1-H]$^{2+}$, the V−N$_{bridge}$ distance for [1-H]$^{3+}$ contracts by 0.92 Å to 2.228(5) Å, which is smaller than the sum of their covalent radii [$R\sum cov^{-1} = 0.995(2)$].[55] As shown in Table 1, whereas [1-H]$^{2+}$ shows relatively equal V−N distances, for [1-H]$^{3+}$, significant deviation is apparent in the metal−ligand bond lengths of the imine and pyridine donor atoms.

Complex [1-H]$^{3+}$ displays the shortest metal−bridgehead nitrogen contact among the reported tren systems. The structure of [1-H]$^{3+}$ is well-defined by the CShM for a capped

octahedral geometry [S(COC-7) = 0.791], an arrangement that is favored in tren-containing systems but relatively anomalous in the full catalogue of heptacoordinate complexes.[62] A greater value of S(COC-7) indicates greater distortion from the seven (hepta)-coordinate geometry of a reference capped octahedral coordination environment. While analogous Mn^{2+} complex geometries have been assigned as capped octahedral, the separation from N$_{bridge}$ is primarily an impact of crystal packing forces.[61,87,88] In contrast, as we will outline below, heptacoordination is a source of stability for the d^2 electronic configuration of V^{3+}, which will help facilitate the doublet ground state we observe for [1-CO$_2$Me]$^{2+}$.

Heptacoordination can be understood in terms of the hybridization between the metal d orbitals. Independent of metal identity, the σ-bonding abilities of the five d orbitals are inequivalent, and greatest for the d$_{z^2}$ orbital.[89,90] In octahedral symmetry, the solitary d$_{z^2}$ orbital points along a metal–ligand bond axis. Interestingly, the symmetric linear combination of the π-symmetric octahedral t$_{2g}$ orbitals (d$_{xz}$, d$_{yz}$, and d$_{xy}$) forms a hybridized orbital of pure d orbital and σ symmetry character along a 3-fold octahedral axis with a bond strength equivalent to that of the d$_{z^2}$ orbital. We refer to this orbital as d$_\sigma$, and a plot of it is shown in Figure 6.[91,92] We note that our current
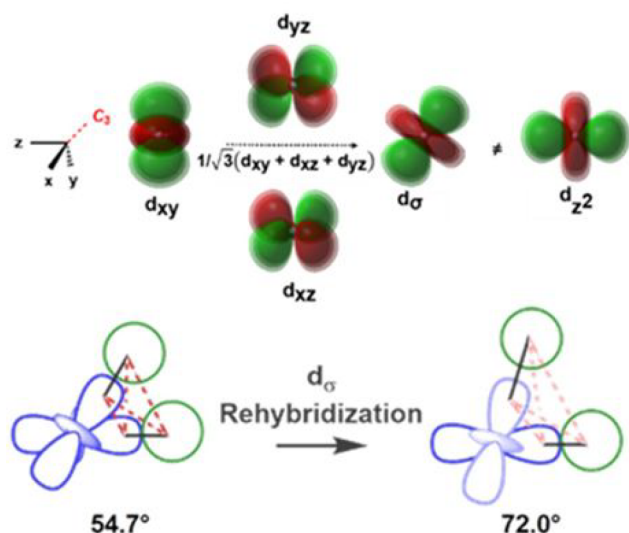

Figure 6. Symmetric hybridization among the d$_{xy}$, d$_{xz}$, and d$_{yz}$ orbitals that constitute the t$_{2g}$ set with respect to an octahedral coordination environment (top). This generates a d$_\sigma$ orbital that is equivalent to the d$_{z^2}$ orbital but rotated 54.7° along its nodal zone. The contours are shown from 0.025 to 0.045 in increments of 0.01 electron bohr^{-2}. Rehybridization of the d$_\sigma$ orbital from the ideal bond angle of 54.7° to the experimental bond angle of 72.0° to remove its angular overlap with the d$_{z^2}$ orbital (bottom). This hybridization results in the orbital character of a capped octahedral, C_3-symmetric coordination environment detailed in Figure 2.

model is distinct from a group theory approach that assigns the d orbital oriented along the C_3-symmetry axis as d$_{z^2}$ that has been previously detailed by McGarvey and Telser[93] and Atanasov and Neese.[14] The C_3-symmetric "a" orbital does not interfere with the "octahedral" orbitals of e$_g^*$ symmetry that is the basis of the following discussion.

We can consider the strength of the hypervalent metal–ligand bond involving d$_\sigma$ using principles of orbital orthogonality that stem from the classic work of Pauling[89]

and Hultgren[90] that were later refined by Landis.[94−97] Irrespective of coordination number, we refer to a transition metal center as hypervalent if it possesses an orbital of σ-bonding character that is equivalent to its d$_{z^2}$ orbital. Equation 3 details the angular overlap (S_{12}) between the d$_\sigma$ and d$_{z^2}$ orbital based on the angle between the seventh ligand and the adjacent metal–ligand bond (θ). At an angle of 54.7°, the orbitals have an angular overlap of 0 so that they are orthogonal. The directional character of this d$_\sigma$ orbital is equivalent to that of the d$_{z^2}$ orbital that is related to its σ-bond strength with an interacting ligand. The orbital energy of this d$_\sigma$ orbital is +0.60Δ with respect to the ligand field splitting of its corresponding octahedral coordination environment.[98] Due to interelectronic repulsion between the metal–ligand σ-bonds, the angle expands to approximately 72° as reported by Hoffmann, Muetterties, and co-workers and illustrated in Figure 6.[99] This is consistent with the structure of [1-H]$^{3+}$ that displays an average ∠N$_{bridge}$–V–N$_{imine}$ bond angle of 71.5(3)° corresponding to an S_{12} value of 0.35 with respect to the d$_\sigma$ and d$_{z^2}$ orbitals. This posits that the strength of the hypervalent V–N$_{bridge}$ bond of [1-H]$^{3+}$ is a factor of 0.65 (1 − S_{12}) that is equal to +0.25Δ with respect to the remaining metal–ligand σ-bonds of the octahedral coordination environment.

$$S_{12} = \left| \frac{3 \cos^2 \theta - 1}{4} \right| \tag{3}$$

The ligand field transition energies are centered at 774 and 506 nm, highlighted in Figure S6, and illustrated with our NT$\rho^{\alpha,\beta}$ plots in Figure 7. These ligand field transitions are descriptive of the C_3 symmetry of the tripodal scaffold that reduces the 3-fold degeneracy of t$_{2g}$ → e + a*. This reduction in symmetry shifts the reference energy (barycenter) of the complex, stabilizing the lower- and upper-lying e orbitals. The difference in energy between the d$_\sigma$ (a*) and lowest-lying e orbitals relative to the octahedral ligand field splitting (Δ) is equal to the deviation in the angular overlap from unity that we term Δ_3 in eq 4.

$$1 - S_{12} = \left(\frac{a^* \leftarrow e}{e^* \leftarrow e} \right) = \Delta_3 \tag{4}$$

Equation 5 provides the change in ligand field stabilization energy (ΔLFSE) associated with a d^n metal center adopting a capped octahedral geometry with respect to the corresponding octahedral complex. The experimental bond angle for [1-H]$^{3+}$ of 72° suggests that heptacoordination stabilizes the complex by −0.26Δ. This converts to a hypervalent bond strength of 14.8 kcal mol^{-1} if we assume that contraction of the bridgehead nitrogen strictly impacts the energy of the a* orbital. The experimental ligand field transition of a* ← e suggests that coordination of the bridgehead nitrogen stabilizes [1-H]$^{3+}$ by an identical value of 14.8 kcal mol^{-1}, which validates our current structural model. Comparatively, a capped octahedral geometry is not anticipated for [1-H]$^{2+}$ because the d$_\sigma$ orbital is populated, heptacoordination destabilizing the complex by +0.26Δ, which is consistent with its experimental V–N$_{bridge}$ separation shown in Figure 2.

$$\Delta\text{LFSE}(\Delta) = 2\Delta_3[4N(a^*) - N(e^*) - N(e)] \tag{5}$$

Electrochemical experiments further support stabilization from oxidation-activated hypervalency. The cyclic voltammogram of [1-H]$^{2+}$ shown in Figure 8 displays a reversible 3+/2+ oxidation at −0.51 V versus Fc$^+$/Fc^0 in CH$_3$CN, along with

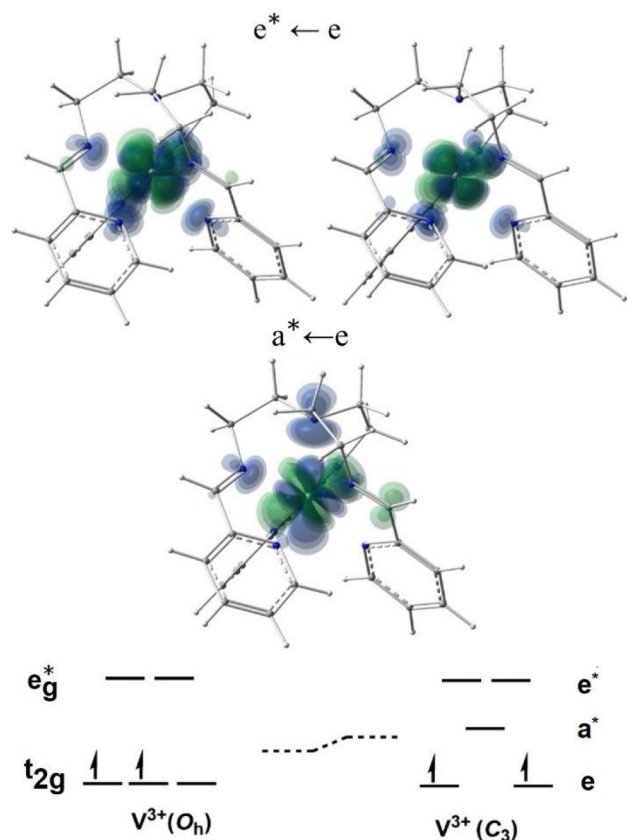**Figure 7.** Natural transition spin density $(NT\rho^{\alpha,\beta})$ plots of the ligand field transitions of the vanadium complex of $[\mathbf{1\text{-}H}]^{3+}$ that illustrate the loss of degeneracy associated with heptacoordination to the tren scaffold (top). The contours are shown for 0.003, 0.006, and 0.012 electron bohr^{-2}. Orbital energy diagrams of an ideal hexacoordinate, octahedral, and heptacoordinate, capped-octahedral, V^{3+} (d^2) complex (bottom). The shift in the barycenter (dashed lines) between the octahedral and capped octahedral coordination geometry is shown.
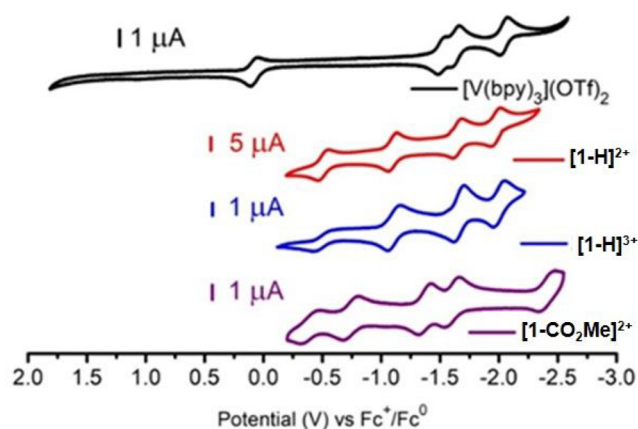

Figure 8. Cyclic voltammograms (0.1 V/s) of the cathodic reduction potentials for $[\mathbf{1\text{-}H}]^{2+}$, $[\mathbf{1\text{-}H}]^{3+}$, $[\mathbf{1\text{-}CO_2Me}]^{2+}$, and $[V(bpy)_3]$-$(OTf)_2$.[35] The experiments were performed in 0.1 M Bu$_4$NPF$_6$ and CH$_3$CN.

ligand reductions at −1.09, −1.64, and −1.98 V versus Fc$^+$/Fc^0. The corresponding V^{2+} bipyridine complex exhibits an $E_{1/2}(3+/2+)$ of 0.08 V, which suggests the tren scaffold stabilizes the V^{3+} oxidation state by 0.59 eV (13.6 kcal mol^{-1}), in agreement with the ΔLFSE values calculated above.

Compound $[\mathbf{1\text{-}H}]^{3+}$ displays identical redox properties except that the 3+/2+ couple lacks reversibility. It is interesting to note that the same trend in the electrochemical reversibility of the $E_{1/2}(3+/2+)$ value based on coordination geometry has been previously reported in characterizing spin-labile Fe^{2+} tren-iminopyridine complexes.[100,101] We report a Randles−Sevcik analysis on the 3+/2+ oxidation of $[\mathbf{1\text{-}H}]^{3+}$ in section S2.1 that determines a linear relationship between the total current with respect to the square root of the electrochemical scan rate, $\nu^{1/2}$. The linear dependence from the Randles−Sevcik analysis of the reduction with respect to the oxidation of the $E_{1/2}(3+/2+)$ couple is 3 times greater. This could suggest that the electrochemical scan rate is on the order of the kinetics of the breaking of the seventh coordinate covalent bond associated with the reduction of complex $[\mathbf{1\text{-}H}]^{3+}$ to $[\mathbf{1\text{-}H}]^{2+}$. The cyclic voltammogram of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ is included in Figure 8 and detailed below.

For the vanadium complexes reported herein, heptacoordination is a redox-mediated process. The conversion between OC-6 and COC-7 has not been previously reported. While spin-labile Fe^{2+} complexes will exhibit a contraction of the tren scaffold when undergoing a low- to high-spin transition, the Fe−N$_{bridge}$ separation does not approach the sum of their covalent radii.[60,61] Halcrow has previously reported that spin crossover is inhibited in these complexes by deviations in the Fe^{2+}−N$_{bridge}$ separation that exceed 0.5 Å upon comparison of low- and high-spin states.[102] Similar to the discussion of Mn^{2+} presented above, the equilibrium position of the N$_{bridge}$ of the tren scaffold of high-spin Fe^{2+} complexes is principally an impact of crystal packing forces.

Doublet Ground State in 1-CO$_2$Me. Lastly, we detail the roles that heptacoordination and redox noninnocence play in the formation of an unexpected doublet ground state for $[\mathbf{1\text{-}CO_2Me}]^{2+}$. For comparison with previously reported tren structures,[87] we installed an ester substituent that is located *meta* and *para* with respect to the pyridine and imine N donor atoms, respectively. The isovalent Cr^{2+} complex displayed redox noninnocence with an iminopyridine moiety that is descriptive of its triplet ground state.[103] We would anticipate this behavior to be more accessible for V^{2+} systems because their ionization potential (29.31 eV) is smaller than that of Cr^{2+} (30.96 eV). In addition, heptacoordination restores some of the loss of LFSE associated with removing an electron from a t$_{2g}$ orbital from an octahedral V^{2+} complex that is not accessible to Cr^{2+} because its d$_\sigma$ orbital would be singly occupied.

In agreement with the crystal structure of $[\mathbf{1\text{-}H}]^{3+}$, $[\mathbf{1\text{-}CO_2Me}]^{2+}$ displays structural properties that are consistent with a V^{3+} species, rather than a V^{2+}, oxidation state (Table 1). The V−N$_{bridge}$ separation of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ is 2.213(9) Å, equal to 0.988(4) $R\sum$cov^{-1}. The CShM classifies $[\mathbf{1\text{-}CO_2Me}]^{2+}$ as a capped octahedral geometry in which S(COC-7) is equal to 0.908 that is more distorted than $[\mathbf{1\text{-}H}]^{3+}$. The anticipated C_3 symmetry is not observed for $[\mathbf{1\text{-}CO_2Me}]^{2+}$, as the V−N$_{imine}$ and V−N$_{pyridine}$ bond lengths of the trigonal ligand set are inequivalent. Two of the coordinated groups have average V−N$_{imine}$ and V−N$_{pyridine}$ bond lengths of 2.10(1) and 2.20(4) Å, respectively, that are consistent with a V^{3+} metal center. In those two iminopyridine ligand arms, the C$_{imine}$−C$_{pyridine}$ bond distance is 1.42(1) Å. The C$_{imine}$−C$_{pyridine}$ bond length of the remaining iminopyridine group contracts to 1.38(2) Å, suggestive of the partial double-bond character that accompanies a one-electron reduction.[104] The N donor atoms display

contracted V−N$_{imine}$ and V−N$_{pyridine}$ bond lengths of 1.976(9) and 2.112(7) Å, respectively, which can be attributed to greater Coulombic attraction between the cationic transition metal center and anionic ligand radical. The ester substituents are in the plane of the iminopyridine groups, and their dihedral angle (φ) is 14(4)°, whose π^* orbitals of the ester functional group are conjugated with that of the aromatic ligand set.

Collectively, the structural features of [1-CO$_2$Me]$^{2+}$ are consistent with a V^{3+} metal center resulting from redox noninnocence of an iminopyridine moiety. Literature precedent suggests that the spin centers should be antiferromagnetically coupled, resulting in a doublet ground state.[105,106] To support this spin assignment, we analyzed the product of the magnetic susceptibility (χ_M) and temperature (T) of [1-CO$_2$Me]$^{2+}$ as a function of temperature in Figure 9. At low
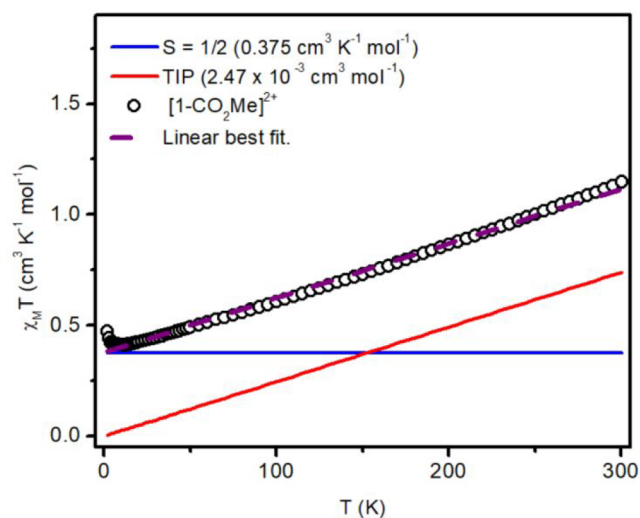

Figure 9. Temperature dependence of the magnetic susceptibility for a powdered sample of [1-CO$_2$Me]$^{2+}$, collected with a 1 kOe measurement field. The linear best fit of the experimental data is $\chi_M T = 0.377 + (2.47 \times 10^{-3})T$ ($R^2 = 0.994$).

temperatures, the minimum value of $\chi_M T$ is 0.41 cm^3 K^{-1} mol^{-1}, suggesting an isotropic $S = 1/2$ state. This conclusively illustrates that this nominally d^3 system does not possess a quartet ground state. This is further supported by electron paramagnetic resonance (EPR) spectroscopy data that are collected in section S.14. We are surprised to find that the $\chi_M T$ value increases linearly with temperature above 50 K, inconsistent with the curvature anticipated for (weaker) antiferromagnetic exchange coupling between metal and ligand spin centers. We attribute the linear response to temperature-independent magnetism (TIP) that results from second-order Zeeman coupling between the ground and excited states. The magnetic susceptibility of [V(bpy)$_3$](ClO$_4$)$_2$ has previously been reported, and its $\chi_M T$ is 1.79 cm^3 mol^{-1} K^{-1}, which is comparable with an $S = 3/2$ spin-only value of 1.88 cm^3 mol^{-1} K^{-1}.[107] The temperature dependence of the magnetic susceptibility provided therein does not support the presence of TIP. Analogous magnetic properties of [V(bpy)$_3$](ClO$_4$)$_2$ are observed for Cr^{3+} and Mn^{4+} complexes, which suggests the electronic structure of [1-CO$_2$Me]$^{2+}$ is unique with respect to other 3d^3 transition metal complexes because of its $S = 1/2$ ground state and large TIP.

While the TIP value is 15−20 times larger than is traditionally assigned to representative octahedral Co^{3+}

complexes,[108] our assignment is consistent with previous reports of redox noninnocent octahedral Ti^{3+} and V^{4+} diimine systems whose anomalous magnetic properties were attributed to trigonal distortion.[109,110] The magnitude of TIP is inversely proportional to the excited state energies. Because of the redox-active transition metal center and ligand sets, [1-CO$_2$Me]$^{2+}$ has a densely populated excited state manifold. This is detailed with respect to the NEVPT2(3,8) multireference calculations of [1-CO$_2$Me]$^{2+}$ that are collected in Table S9 and Figure S17. Unless otherwise stated, the V−N$_{bridge}$ bond length of the lowest-lying doublet state of the V^{2+} complexes was constrained during the DFT geometry optimization to the crystallographic distance of 2.25 Å.

The electronic absorption spectra of [1-H]$^{2+}$, [1-H]$^{3+}$, [1-CO$_2$Me]$^{2+}$, and our previously detailed polypyridyl analogue, [V(bpy)$_3$](OTf)$_2$, are collected in Figure 3.[36,37] Along with the ligand field transitions detailed above, [1-H]$^{3+}$ contains a high-energy MLCT at 378 nm, consistent with the large ionization energy of the V^{3+} ion. The broad peak observed in the near-IR region is indicative of a $\pi^* \leftarrow \pi^*$ transition of a ligand radical that is consistent with the crystallographic structure of [1-CO$_2$Me]$^{2+}$. Similarly, we assign the absorbance with a λ_{max} at 397 nm to the V^{3+} MLCT that is red-shifted with respect to the unsubstituted complex, [1-H]$^{3+}$. The TD-DFT results and NT$\rho^{\alpha,\beta}$ plots are provided in section S.7.2 and Figure S15.

Interestingly, we find that [1-CO$_2$Me]$^{2+}$ displays spectroscopic signatures that are diagnostic of both V^{2+} and V^{3+} formal oxidation states. Along with the ligand field transitions detailed above, [1-H]$^{3+}$ contains a high-energy MLCT at 378 nm, consistent with the large ionization energy of the V^{3+} ion. We assign the absorbance of [1-CO$_2$Me]$^{2+}$ with a λ_{max} at 397 nm that we attribute to the V^{3+} MLCT. The broad peak observed in the near-IR region is indicative of a $\pi^* \leftarrow \pi^*$ transition of a ligand radical between the reduced and an unreduced iminopyridine group. Both spectroscopic features are consistent with the crystallographic structure of [1-CO$_2$Me]$^{2+}$. The intense absorbance at 785 nm resembles the MLCT of [1-H]$^{2+}$ whose energy is red-shifted by 0.12 eV. This suggests that both ^{4}A$_2$ and the crystallographically and magnetically observed doublet state are present in solution. We corroborate this proposal with NEVPT2(3,8) calculations of [1-CO$_2$Me]$^{2+}$ in its doublet and quartet optimized geometries that are collected in Figure S18. We do not observe a thermal dependence of the electronic absorbance in a 1:4 (v:v) methanol/ethanol mixture at 298 and 77 K that could be attributed to the solvatochromism that is detailed in sections S3.2 and S3.3. We calculate that the enthalpy of the unconstrained optimized geometry of the [1-CO$_2$Me]$^{2+}$ doublet state of is 1.83 kcal mol^{-1} higher in energy than its quartet state. The unconstrained optimization lengthens the V−N$_{bridge}$ bond from 2.21 to 2.42 Å and lowers the enthalpy of the doublet state by −2.85 kcal mol^{-1}. We note that hybrid DFT functionals routinely overstabilize spin states of greater multiplicity, which supports our proposal that the quartet and doublet states of [1-CO$_2$Me]$^{2+}$ are nearly degenerate.[111,112]

The experimentally observed absorption at 467 nm is distinct to [1-CO$_2$Me]$^{2+}$. The low energy of the transition suggests that the excitation is of triplet, rather than singlet, character that is magnetically coupled to the three-electron paramagnetic system. The absorbance values of nominally spin-forbidden transitions in aromatic systems have been previously reported in the presence of O$_2$ and attributed to their triplet ground states.[113] We follow up our TD-DFT study
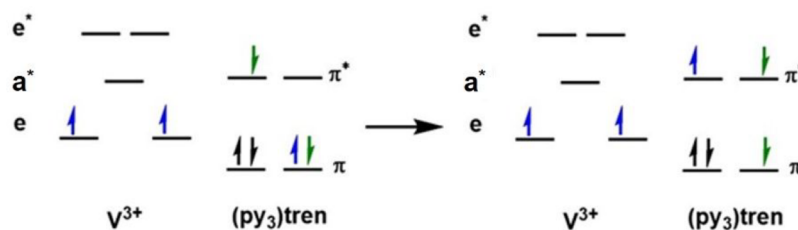
Figure 10. Orbital depiction of the ground and excited states of the intraligand $^2(^3IL)$ transition of $[\mathbf{1\text{-}CO_2Me}]^{2+}$.

with an NEVPT2(5,5) study to address the multireference character of this transition. The resulting spin eigenfunctions for this excited state suggest that the reduced iminopyridine undergoes a spin flip during the transition that changes its magnetic coupling with the V^{3+} metal center from an antiferromagnetic (^{2}MLCT) to a ferromagnetic (^{4}MLCT) interaction that is detailed in Table S14. The ^{4}MLCT state antiferromagnetically couples with the triplet-intraligand transition (^{3}IL), resulting in an overall doublet state whose absorbance is spin allowed; we refer to it as $^2(^3IL)$, and its orbital depiction is provided in Figure 10.

The cyclic voltammogram of $[\mathbf{1\text{-}CO_2Me}]^{2+}$, shown in Figure 9, resembles those of $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}H}]^{3+}$. $[\mathbf{1\text{-}CO_2Me}]^{2+}$ is oxidized at -0.39 V and reduced at -0.73 V versus Fc^+/Fc^0. The $E_{1/2}(3+/2+)$ and $E_{1/2}(2+/1+)$ redox events of compound $[\mathbf{1\text{-}CO_2Me}]^{2+}$ are 0.12 and 0.36 V anodic relative to the parent compounds, respectively, which we attribute to the electron-withdrawing properties of the ester substituents. Importantly, the open circuit potential does not sit at the $E_{1/2}$ value of the observed peaks, which suggests that there is one dominant species in solution and no impurity, such as a di- and tricationic reduced−oxidized pair. The $E_{1/2}(3+/2+)$ redox event displays quasi-reversibility, consistent with the electrochemical behavior of $[\mathbf{1\text{-}H}]^{3+}$. This suggests that $[\mathbf{1\text{-}CO_2Me}]^{2+}$ maintains the crystallographically observed capped octahedral geometry in solution. The similar electrochemical properties suggest that the products of both reduction and oxidation of $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$ are analogous, which is supported by spectroelectrochemistry data that are provided in section S2.2. The data suggest that the structural features of $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$ are attributable to different electronic states of the +2 system, which correspond to the lowest-lying quartet and doublet states of a 3d^3 system, respectively.

Treatment of Multielectron Antiferromagnetism with the SOAP Method. To elucidate this complex spin system, we expand upon Hund's first rule and consider the electronic structure of the doublet ground state of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ that can be described in local terms as a V^{3+} (d^2) metal center and anionic ligand radical (π^{*-}).

The metal center's two unpaired electrons and anionic ligand radical occupy nondegenerate orbitals that can be antiferromagnetically coupled, the doublet being lower in energy than the quartet state. Both solid state and molecular models of antiferromagnetic coupling require overlap between magnetic orbitals.[114−117] Orbital orthogonality is a requirement of most electronic structure methods, spin unrestricted broken symmetry models being the exception.[118−120] The overlap between unpaired electrons that engender antiferromagnetic coupling can be found by using multideterminant methods. Multireference NEVPT2(3,8) calculations suggest that the doublet ground state is stabilized by 0.55 eV with respect to the ^{4}MLCT state, that is equal to a magnetic

coupling of $J = -1480$ cm^{-1} with respect to a $-2J$ formalism for the Heisenberg−Dirac−van Vleck Hamiltonian.[84,85]

We describe an antiferromagnetic interaction as a configuration interaction involving metal-centered (φ_a) and an anionic ligand radical (φ_b) orbital that are partially occupied. φ_a and φ_b are orthogonal one-electron wave functions.[48,49,51,121] This is detailed in eq 6, where the pair coefficients (c_1 and c_2) are solved variationally with respect to the lowest-energy doublet state. The first configuration by itself is termed a perfect pairing model in which electrons that are singlet-coupled are orthogonal to the spatial and spin wave functions of the system's remaining electrons.[50] Our SOAP method expands on this conventional description of two electrons[117] by addressing the spin recoupling involving the remaining metal-centered unpaired electron (φ_c).[122] This recoupling is accomplished with inclusion of a third configuration in eq 6, where φ_a and φ_b are triplet coupled via their magnetic interaction with φ_c. The spin coefficients (ε_1 and ε_2) are normalized with respect to the two spin eigenfunctions that are shown in eq 6, where α and β are electron spins.

$$\psi(SOAP) = \varepsilon_1[(c_1\phi_a\phi_a - c_2\phi_b\phi_b)\phi_c](\alpha\beta - \beta\alpha)\alpha$$
$$+ \varepsilon_2[(\phi_a\phi_b - \phi_b\phi_a)\phi_c](2\alpha\alpha\beta - \alpha\beta\alpha - \beta\alpha\alpha)$$
(6)

The coupling between this new configuration and a perfect pairing model is provided by the meta-exchange interaction (M_c) in eqs 7 and 8. This electron−electron interaction reflects the overlap between the overlap functionals that involve the third unpaired electron (f_{ac} and f_{bc}). If the three orbitals are mutually orthogonal, then the energy expression of eq 6 simplifies to the ^2E excited state energy detailed with respect to $[\mathbf{1\text{-}H}]^{2+}$.

$$M_c = \int\int \frac{\phi_a(1)\phi_c(1)\phi_b(2)\phi_c(2)}{r_{12}}\, dr_1\, dr_2$$
(7)

$$M_c = \int\int \frac{f_{ac}f_{bc}}{r_{12}}\, dr_1\, dr_2$$
(8)

For visualization and insight, the two paired orthogonal orbitals φ_a and φ_b can be transformed into equivalent nonorthogonal one-electron wave functions, φ_x and φ_y, respectively, as detailed in eqs 9−11 with respect to the pair coefficients c_1 and c_2. The resulting overlap of φ_x and φ_y in the SOAP wave function is 0.23, which classifies $[\mathbf{1\text{-}CO_2Me}]^{2+}$ as showing an antiferromagnetic coupling, rather than a covalent bond, between its V^{3+} metal center and anionic ligand radical. Our calculated nonorthogonal orbitals of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ are presented in Figure 11, along with their originating orthogonal orbitals from the SCF framework. Our SOAP method assigns

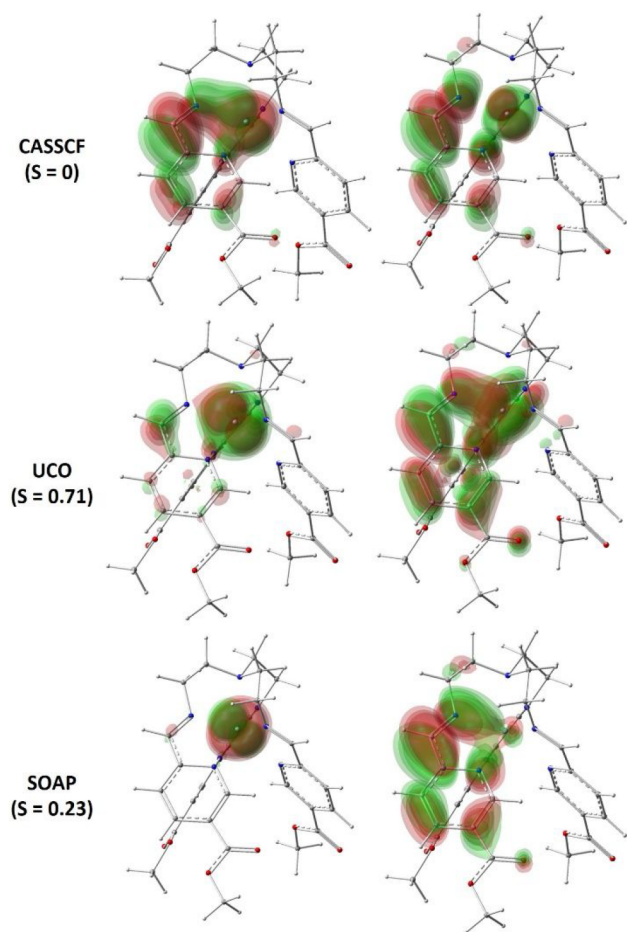$$\phi_x = \frac{\sqrt{c_1}\,\phi_a + \sqrt{c_2}\,\phi_b}{\sqrt{c_1 + c_2}} \tag{9}$$

$$\phi_y = \frac{\sqrt{c_1}\,\phi_a - \sqrt{c_2}\,\phi_b}{\sqrt{c_1 + c_2}} \tag{10}$$

$$S = \frac{c_1 - c_2}{c_1 + c_2} \tag{11}$$

The observation of a doublet electronic configuration as the ground state is unique to $[\mathbf{1\text{-}CO_2Me}]^{2+}$ with respect to $[\mathbf{1\text{-}H}]^{2+}$. To improve our understanding of this, we consider the NEVPT2(3,8) transitions of the quartet-optimized geometry of $[\mathbf{1\text{-}CO_2Me}]^{2+}$, which is consistent with a local V^{2+} metal center, analogous to the crystallographic structure of $[\mathbf{1\text{-}H}]^{2+}$. On the basis of analogous Cr^{3+} complexes, we would expect to observe five near-degenerate low-lying doublets that are attributed to the metal-centered ^2E and $(^2E + ^2A_2)$ states of a C_3-symmetric d^3 system. However, as previously reported for $[V^{2+}(bpy)_3]$,[36] we calculate six doublet states that are below approximately 1.5 eV for both $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$. The states exhibit significant multireference mixing between the metal-centered and π^* orbitals, which describe resonance stabilization and static correlation. Rather than dominantly metal-based states, the doublet excited states exhibit coupling between their metal-centered and charge transfer states that we term $^2MC/^2MLCT$.[36] The CASSCF(3,8) wave function determines that the ^2MC and ^2MLCT state functions contribute approximately equally to the lowest-lying doublet excited state for both $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$, whose excited state energies are 1.05 and 0.88 eV, respectively.

We also observe a low-lying ^4MLCT state, strongly stabilized by dynamic correlation, that is 1.75 and 0.74 eV above the ground states for $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$, respectively. Heptacoordination with the tren scaffold and antiferromagnetic coupling with the iminopyridine anionic ligand radical in the doublet ground state geometry should collectively stabilize the doublet configuration by approximately 1.2 eV. These features place the energy of the ^2MLCT excited state below that of the $^4A_2$ ground state for $[\mathbf{1\text{-}CO_2Me}]^{2+}$, which would not be observed in the case of $[\mathbf{1\text{-}H}]^{2+}$. Relative to the parent ligand set, we calculate that the 5-CO_2Me substituent stabilizes the electron affinity of the uncoordinated ligand by 0.58 eV. Rotating the ester substituent so that its π-system is orthogonal to the iminopyridine ligand set diminishes the electron affinity stabilization to -0.18 eV. This suggests that the ester substituent serves as a π-acid that stabilizes the charge transfer state of $[\mathbf{1\text{-}CO_2Me}]^{2+}$, resulting in the relative energies of the lowest-lying quartet and doublet states.

Figure 11. Singly occupied orbitals that describe the antiferromagnetic coupling of the vanadium complex in $[\mathbf{1\text{-}CO_2Me}]^{2+}$. This is presented as orthogonal bonding (φ_a) and antibonding (φ_b) orbitals (top) with respect to the state specific CASSCF(3,3) wave function of the lowest-lying doublet. Unrestricted corresponding orbitals (UCO's) of the TPSSh BS-DFT wave function (middle). Equations 9 and 10 transform those orbitals into nonorthogonal and localized metal-centered (φ_x) and anionic ligand radical (φ_y) orbitals (bottom) with respect to their calculated pair coefficients, c_1 and c_2. The overlap between the two orbitals (S) is provided in parentheses with respect to their pair coefficients illustrated in eq 11. The contours are shown from 0.025 to 0.045 in increments of 0.01 electron bohr^{-2}.

an antiferromagnetic coupling of $J = -1490$ cm^{-1} that accurately reproduces the sign and magnitude of the magnetic coupling of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ (the energy expressions are further detailed in section S7.2). We also use the broken symmetry method (BS-DFT) with the TPSSh hybrid-DFT functional to calculate the magnetic coupling of $[\mathbf{1\text{-}CO_2Me}]^{2+}$. The TPSSh functional has been shown to accurately treat the spin state energetic of diverse multielectron systems.[123,124] The DFT-BS procedure calculates a J of -1600 cm^{-1} that is consistent with the methods detailed above. The unrestricted corresponding orbital (UCO) transformation of the DFT-BS wave function assigns an overlap of 0.70 for the nonorthogonal electron pair that is shown in Figure 11. We attribute the larger calculated overlap for the UCO transformation (with respect to our SOAP method) to its lack of accounting for the third spin coupling configuration detailed in eq 6.

CONCLUSIONS AND OUTLOOK

We have established that the doublet ground state of $[\mathbf{1\text{-}CO_2Me}]^{2+}$ is stabilized by a large electron affinity for the *m*-methyl ester-substituted iminopyridine, strong antiferromagnetic coupling with the anionic ligand radical, and heptacoordination with the tren scaffold. Concurrently, the stability of the ferromagnetic quartet state is decreased due to the delocalization of its metal-centered orbitals, which weakenes exchange interactions. Stabilization of the doublet ground state is further attributed to the electronic coupling, resonance stabilization, between the doublet metal-centered (^2MC) and metal-to-ligand charge transfer (^2MLCT), which

we term ^{2}MC/^{2}MLCT. Collectively, these features combine to invert the ground and excited states for $[\mathbf{1\text{-}CO_2Me}]^{2+}$, as observed crystallographically, spectroscopically, and magnetically.

States of distinct multiplicities are differentially impacted by resonance. This is foundational to chemistry, the Heitler–London treatment of H_2 assigning resonance and configuration interactions as the basis of the attractive and dominantly repulsive character of its singlet ($\alpha\beta$) and triplet ($\alpha\alpha$) states, respectively.[125] While intuitive for a covalent bond, resonance in the ground states of paramagnetic coordination complexes is relatively obscured. The stability of the high-spin state for $[\mathbf{1\text{-}CO_2Me}]^{2+}$ predicted from Hund's first rule is complicated by nondegenerate and nonorthogonal orbitals that possess competing magnetic interactions.[126] Herein, we uncover a unique phenomenon in which resonance stabilization causes an unexpected ground ($S = \frac{1}{2}$) state of a nominally V^{2+} ($3d^3$) transition metal complex.

The conditions for this doublet ground state to manifest are perfectly satisfied with the ligand environment of $[\mathbf{1\text{-}CO_2Me}]^{2+}$. Alvarez has previously performed an extensive characterization of the crystallographic structures of heptacoordinate transition metal complexes.[62] Heptacoordinate structures are largely observed in group 3–6 elements with a d^0–d^4 electronic population whose conditions are met for $V^{2+/3+}$. While the geometry of heptacoordinate structures is traditionally dynamic, possessing low barriers of conversion, the C_3 symmetry and multidentate chelation of the tren scaffold enables achievement of a capped octahedral coordination environment.[127] Coordination between the V^{3+} metal center and bridgehead nitrogen stabilizes the doublet state by approximately 15 kcal mol^{-1}.

Lastly, the redox noninnocence of the iminopyridine groups provides $[\mathbf{1\text{-}CO_2Me}]^{2+}$ a mechanism for accessing its doublet ground state. The strong electronic coupling between the metal-centered orbitals and anionic ligand radical results in resonance stabilization that assigns the ^{2}MC/^{2}MLCT as the electronic ground state. The ester substituent magnifies this effect, consistent with our previous studies of the emission properties of analogous Cr^{3+} complexes, which suggest greater metal–ligand covalency.[63,128]

We briefly summarize the advancements that our SOAP method affords for the theoretical treatment of open-shell systems. We note that, with select exceptions,[129−132] models for antiferromagnetism have focused on two-electron systems.[117] Because our method decomposes the magnetism of a system into its pairwise components, it can be employed to assign competing magnetic interactions in a multielectron system.[133,134] We establish that metal–ligand covalency affects the magnitude of electron–electron exchange interactions, thus modulating the strength of the high-spin (ferromagnetic) state. Greater metal–ligand covalency concurrently stabilizes charge transfer states in the molecule. The observation of the $^4A_2$ and ^{2}MC/^{2}MLCT ground states for $[\mathbf{1\text{-}H}]^{2+}$ and $[\mathbf{1\text{-}CO_2Me}]^{2+}$, respectively, alludes to a promising design principle for stabilizing the doublet ground state based on established synthetic principles of ligand substituent electronics. An analogous proposal for stabilizing the quartet state has recently been made by Wenger.[9]

The variational approach of our SOAP method can address the avenues for resonance stabilization and magnetic coupling that are accessible to open-shell systems. This feature also provides the opportunity to transform computationally convenient orthogonal orbitals to nonorthogonal orbitals, orbital overlap being the requirement for antiferromagnetic coupling. From a practical standpoint, the calculation time of the interelectronic terms scales with the volume of an orbital's grid rather than the total number of electrons. Because the only required input is the Gaussian cube of a system's singly occupied orbitals, the procedure is agnostic to the originating electronic structure software.
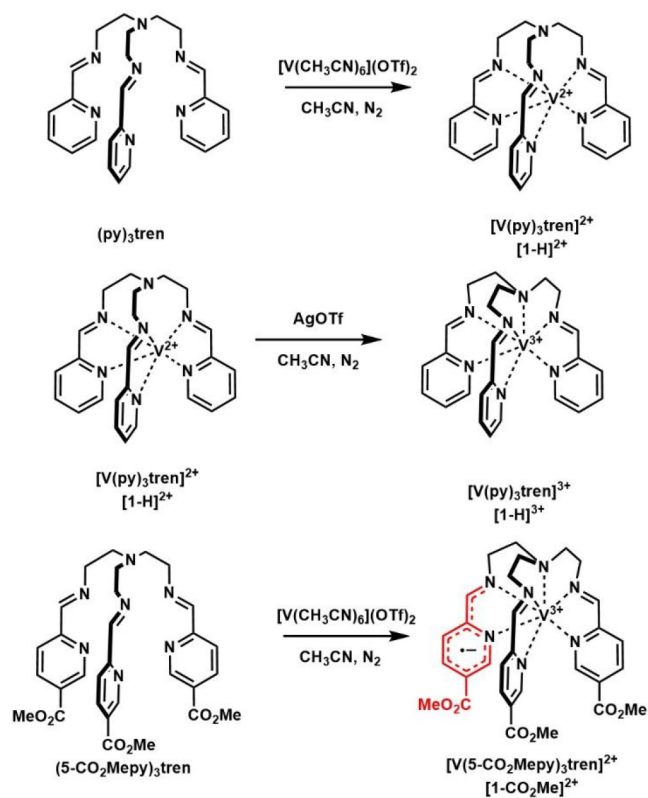
■ EXPERIMENTAL SECTION

Physical Methods. Absorption spectra were recorded with a Hewlett-Packard 8453 spectrometer in quartz cuvettes with a 1 cm path length. Infrared spectra were recorded with either a Nicolet 380 FT-IR or a Bruker TENSOR II spectrometer. Mass spectrometry measurements were performed in positive ion or negative ion mode on a Thermo-Finnigan LTQ mass spectrometer equipped with an analytical electrospray ion source using a spray voltage of 2.5 V and a capillary temperature of 175 °C.

Preparation of Compounds. Manipulations and syntheses of all metal complexes were performed inside a dinitrogen-filled glovebox (MBRAUN Labmaster 130). All solvents were sparged with dinitrogen, passed over molecular sieves, and degassed prior to use. The compounds $[V(CH_3CN)_6](OTf)_2$,[135] (py$_3$)tren,[136] and (5-CO_2Mepy)$_3$tren[87] were all prepared according to literature procedures.

Compound $[\mathbf{1\text{-}H}](OTf)_2$ can be generated under multiple conditions, although we have found that the choice of starting material influences the purity of the products. Use of the V^{3+} starting material $[V(CH_3CN)_6Cl_3]$ affords a green solution with electronic absorption properties similar to those of compound $[\mathbf{1\text{-}H}](OTf)_2$ as measured in CH_3CN. When the reaction solution is changed from acetonitrile to tetrahydrofuran, a brown product resembling compound $[\mathbf{1\text{-}H}](OTf)_3$ is produced. The yield of $[\mathbf{1\text{-}H}](OTf)_3$ increases with a decrease in temperature as reactions carried out at −77 °C afford ∼88% of the crude brown precipitate. Due to these complications, we investigated $[V(CH_3CN)_6](BPh_4)_2$ as a V^{2+} starting material. We found $[V(py)_3tren]^{2+}$ is present on the basis of the electronic absorption spectrum, but the molar absorptivity is much lower than that of compound $[\mathbf{1\text{-}H}](OTf)_2$ in the same solvent, indicating a mixture of products. Eventually, we found compound $[\mathbf{1\text{-}H}](OTf)_2$ to be readily prepared as a pure compound by using the starting material $[V(CH_3CN)_6](OTf)_2$, according to the procedure described below (Scheme 1).

$[V(py_3)tren](OTf)_2$ $[[\mathbf{1\text{-}H}](OTf)_2]$. A green solution of $[V(CH_3CN)_6](OTf)_2$ (629 mg, 1.06 mmol) and (py$_3$)tren (444 mg, 1.08 mmol) in 8 mL of CH_3CN was stirred for 4 h. The solvent was removed under reduced pressure and washed with diethyl ether (3 × 3 mL) to afford a green powder. Dark green plate crystals were grown via diffusion of a diethyl ether into a concentrated solution of CH_3CN to yield 665 mg (0.872 mmol, 83% yield) of dark green plate crystals: UV−vis (CH_3CN) λ_{max} (nm) [ε_M (M^{-1} cm^{-1})] 213 (19967), 276 (17362), 402 (1978), 655 (3836), 729 (5219); IR (KBr pellet) $\nu_{C=N}$ 1599 cm^{-1}; ESI-MS(+) (CH_3CN) m/z 232.09 ($[\mathbf{1\text{-}H}](OTf)_2$ − 2OTf$^-$)$^{2+}$. Anal. Calcd for $C_{26}H_{27}F_6N_7O_6S_2V$ $[[\mathbf{1\text{-}H}](OTf)_2]$: C, 40.95; H, 3.57; N, 12.86. Found: C, 40.87; H, 3.56; N, 12.84.

$[V(py_3)tren](OTf)_3$ $[[\mathbf{1\text{-}H}](OTf)_3]$. A clear solution of AgOTf (129 mg, 0.503 mmol) in 3 mL of CH_3CN was added to a solution of $[V(py_3)tren](OTf)_2$ (383 mg, 0.502 mmol) in 8 mL of CH_3CN. Immediately, a Ag precipitate was observed as the color of the solution changed from green to brown. The mixture was stirred for 1 h. The reaction mixture was filtered through Celite to remove Ag metal, and the filtrate was concentrated under reduced pressure. Crystals were grown from diffusion of diisopropyl ether into a concentrated CH_3CN solution to yield 303 mg (0.332 mmol, 66% yield) of brown needle crystals: UV−vis (CH_3CN) λ_{max} (nm) [ε_M (M^{-1} cm^{-1})] 288 (15200), 369 (3760), 449 (788), 506 (351), 774 (26); IR (KBr pellet) $\nu_{C=N}$ 1605 cm^{-1}; ESI-MS(+) (CH_3CN) m/z 154.75 ($[\mathbf{1\text{-}H}]$ − 3OTf$^-$)$^{3+}$, 762.09 ($[\mathbf{1\text{-}H}](OTf)_3$ − 2OTf$^-$)$^{2+}$. Anal.

Scheme 1. Preparation of [1-H](OTf)₂, [1-H](OTf)₃, and [1-CO₂Me](OTf)₂

Calcd for C₂₇H₂₇F₉N₇O₉S₃V [[1-H](OTf)₃]: C, 35.57; H, 2.99; N, 10.75. Found: C, 35.44; H, 3.01; N, 10.78.

[V(5-CO₂Mepy)₃)tren](OTf)₂ [[1-CO₂Me](OTf)₂]. A solution of [V(CH₃CN)₆](OTf)₂ (242 mg, 0.406 mmol) and the ligand (5-CO₂Mepy)₃)tren (240 mg, 0.408 mmol) in 8 mL of CH₃CN was stirred for 4 h. The solvent was removed under reduced pressure and washed with diethyl ether (3 × 3 mL). Crystals were grown from diffusion of diethyl ether into a concentrated solution of methanol to yield 195 mg (0.208 mmol, 51% yield) of black plate crystals: UV−vis (CH₃CN) λ_{max} (nm) [ε_M (M⁻¹ cm⁻¹)] 205 (55300), 236 (37500), 287 (29900), 396 (4620), 667 (3900), 785 (5440), 1024 (1210); IR (KBr pellet) $\nu_{C=O}$ 1725 cm⁻¹, $\nu_{C=N}$ 1600 cm⁻¹; ESI-MS(+) (CH₃CN) m/z 319.25 ([1-CO₂Me](OTf)₂ − 2OTf⁻)²⁺. Anal. Calcd for C₃₂H₃₃F₆N₇O₁₂S₂V [[1-CO₂Me](OTf)₂]: C, 41.03; H, 3.55; N, 10.47. Found: C, 40.68; H, 2.72; N, 10.42.

Crystallography. Key structural data for compounds [1-H]-(OTf)₂, [1-H](OTf)₃, and [1-CO₂Me](OTf)₂ are listed in Table S1. Crystals were coated in Paratone oil, supported on Cryoloops, and mounted on either a Bruker Kppa Apex 2 instrument with a CCD diffractometer [[1-H](OTf)₂] or a Bruker D8 Quest ECO instrument with a Photon 50 CMOS diffractometer [[1-H](OTf)₃ and [1-CO₂Me](OTf)₂] under a stream of cold nitrogen. All data collections were performed with Mo Kα radiation and a graphite monochromator. Initial lattice parameters were determined from a minimum of 310 reflections harvested from 24 frames; these parameters were later refined against all data. Data sets were collected targeting full coverage and 4-fold redundancy. Data were integrated and corrected for absorption effects with the APEX 3 software packages. Structures were determined by direct methods and refined with the SHELXTL software package. Displacement parameters for all non-hydrogen atoms were refined anisotropically. All hydrogens were assigned to ideal positions and refined using a riding model with an isotropic thermal parameter 1.2 times greater than that of the attached carbon atom.

Electrochemistry. Voltammograms (Figure 8) were recorded with either a CH Instruments 1230A or a Gamry Reference 600

potentiostat under a dinitrogen/argon atmosphere. All experiments used a 0.1 M tetrabutylammonium hexafluorophosphate (Bu₄NPF₆) solution in an acetonitrile with a 0.25 mm Pt disk working electrode, a Ag wire quasi-reference electrode, and a Pt wire auxiliary electrode. Reported potentials (Table S2) are referenced to the ferrocenium/ferrocene {[(C₅H₅)₂Fe]⁺/[(C₅H₅)₂Fe], Fc⁺/⁰} redox couple and were determined by adding ferrocene as an internal standard at the conclusion of each electrochemical experiment.

Magnetic Measurements. Solid state magnetic property data for compound [1-CO₂Me](OTf)₂ were collected using a Quantum Design MPMS XL SQUID magnetometer. Powdered microcrystalline samples were loaded into polyethylene bags and inserted into a straw before being transported to the magnetometer. The presence of ferromagnetic impurities was probed by a variable field analysis (0 to 10 kOe) of the magnetization at 100 K (Figure S10). The lack of curvature in the plots of M versus H for [1-CO₂Me](OTf)₂ indicate the absence of significant ferromagnetic impurities. Magnetic susceptibility data were collected at temperatures ranging from 2 to 300 K. Data were corrected for the diamagnetic contributions of the sample holder and bag by subtracting empty containers. Corrections for the sample were calculated from Pascal's constants.[137]

Computational Procedure. The following calculations were performed with the Gaussian16 electronic structure software package.[138] We employed the APFD hybrid DFT functional[139] and a 6-311+g(d) basis set[140] with its parametrized empirical dispersion correction in a PCM acetonitrile continuum solvent.[141] The quartet and doublet states of [1-H](OTf)₂ and [1-CO₂Me](OTf)₂ were optimized where the V−N_bridge separation was constrained to their crystallographic values of 3.15 and 2.21 Å, respectively. The structure of [1-H](OTf)₃ was optimized as a triplet with no geometric constraints.[66] The uncoordinated iminopyridine ligand sets were optimized where the torsion of the bidentate binding group was constrained at 0° for the neutral and singly reduced (monoanionic) states.

A frequency analysis was performed on the quartet and doublet states of [1-CO₂Me]²⁺. We report a single imaginary frequency along the constrained V−N_bridge bond of the doublet state. The unconstrained optimization of the [1-CO₂Me]²⁺ state removes the imaginary frequency and decreases the system's enthalpy by −2.9 kcal mol⁻¹. This suggests that the vibrational mode has a soft potential energy. We also report that the bond length increases from the crystallographic value of 2.21 to 2.42 Å. Our calculations refer to the [1-CO₂Me]²⁺ crystallographic V−N_bridge bond length unless otherwise stated.

The energy of the ²E excited state with respect to the quartet geometry of [1-H](OTf)₂ and [1-CO₂Me](OTf)₂ was multiplied by a factor of ³/₂ to account for the single-determinant character of the DFT-based wave function. The foundation of our multideterminant DFT correction has been detailed elsewhere.[37] The doublet wave function was verified as the ground state through the use of a wave function stability check: single excitations with respect to the trial wave function were included to test for a lower-energy solution, the 'stable=opt' command in g16.

The following calculations were performed with the ORCA 4.1 electronic structure software program.[86] We used multireference techniques with a CASSCF(3,8) model in which the active space was selected to be the five 3d orbitals and the three lowest-lying π^* orbitals. The def2-TZVP and def2-TZVPP basis sets[142] were used for the non-metal and metal atoms, respectively, with a CPCM acetonitrile continuum solvent. The final energies for the excited states were obtained with the NEVPT2 technique with the same active space and conditions.[83−85]

The electron−electron integrals detailed for our SOAP method were calculated, and their cubes generated from their CASSCF(3,8) orbitals that were state specific with respect to their electronic ground state. We selected a grid size of 90 × 90 × 90 for fine resolution. The orbital energies of the magnetic orbitals were approximated from the converged Restricted Open Hartree−Fock (ROHF) wave function of the quartet state. The coefficients used for the SO method ($A_1−A_3$)

were obtained from the lowest-lying doublet of the corresponding CASSCF(3,3) wave function.

The continuous shape measures (CShM) of the experimental crystallographic coordinates and DFT-optimized geometries were evaluated with SHAPE version 2.2.[56]

■ ASSOCIATED CONTENT

ⓈⒾ Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.inorgchem.1c03418.

> Crystal structure details, Randles−Sevcik analysis, spectroelectrochemistry, temperature-dependent UV−vis, electronic absorbance solvatochromism, magnetic measurements, geometry-optimized total energies, TD-DFT, NT$\rho^{\alpha,\beta}$ plots, NEVPT2(3,8) transition energies, CASSCF(3,8) active space orbitals, procedure for the SOAP method and calculating doublet energies of three-electron systems (PDF)
>
> Cartesian coordinates of the DFT geometry-optimized complexes (ZIP)

Accession Codes

CCDC 2114725−2114727 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via www.ccdc.cam.ac.uk/data_request/cif, or by emailing data_request@ccdc.cam.ac.uk, or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

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Notes

The authors declare no competing financial interest.

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