



Sensitivity Analysis for the 2D Navier–Stokes Equations with Applications to Continuous Data Assimilation

Elizabeth Carlson¹ · Adam Larios¹ 

Received: 2 July 2020 / Accepted: 1 August 2021 / Published online: 27 August 2021
© The Author(s), under exclusive licence to Springer Science+Business Media, LLC, part of Springer Nature 2021

Abstract

We rigorously prove the well-posedness of the formal sensitivity equations with respect to the viscosity corresponding to the 2D incompressible Navier–Stokes equations. Moreover, we do so by showing a sequence of difference quotients converges to the unique solution of the sensitivity equations for both the 2D Navier–Stokes equations and the related data assimilation equations, which utilize the continuous data assimilation algorithm proposed by Azouani, Olson, and Titi. As a result, this method of proof provides uniform bounds on difference quotients, demonstrating parameter recovery algorithms that change parameters as the system evolves will be well behaved. Furthermore, our analysis can be extended to analyze the sensitivity of the 2D Euler equations to a viscous regularization. We also note that this appears to be the first such rigorous proof of global existence and uniqueness to strong or weak solutions to the sensitivity equations for the 2D Navier–Stokes equations (in the natural case of zero initial data), and that they can be obtained as a limit of difference quotients with respect to the viscosity.

Keywords Sensitivity analysis · Navier–Stokes equations · Continuous data assimilation · Viscosity · Reynolds number

Mathematics Subject Classification 34D06 · 35A01 · 35Q30 · 35Q35 · 37C50 · 76D03

Communicated by Alain Goriely.

✉ Adam Larios
alarios@unl.edu

¹ Department of Mathematics, University of Nebraska–Lincoln, 203 Avery Hall, Lincoln, NE 68588-0130, USA

1 Introduction

Turbulent flows are well known to be chaotic, in the sense that the solutions to mathematical models of such flows are highly sensitive to initial conditions (see, e.g., Constantin and Foias 1985; Grappin and Léorat 1991). However, sensitivity with respect to physical parameters is also an important consideration in terms of making reliably accurate predictions. Parameter sensitivity is often measured by formally considering the derivative of a solution with respect to a particular parameter; however, the only rigorous justification of this approach in the literature seems to be limited to linear equations, or nonlinear equations under assumptions on the nonlinearity which are too strong to include, e.g., the Navier–Stokes equations of fluids (see, e.g., Brewer 1982; Gibson and Clark 1997 for a semigroup theory approach). Therefore, in the present work, we provide a fully rigorous proof of the global well-posedness of the sensitivity equations for the 2D Navier–Stokes equations. Specifically, we give a rigorous proof of the existence of unique weak and strong solutions with zero¹ initial data to the associated viscosity sensitivity equations specifically for the 2D Navier–Stokes equations. Moreover, we prove that the derivative of solutions with respect to the viscosity is a limit of difference quotients corresponding to different viscosities.

We also extend our results to the case of a data assimilation algorithm. This is because the motivation for this present work arose from our recent work Carlson et al. (2020), where an algorithm was proposed to recover an unknown viscosity, or equivalently Reynolds number. This algorithm works in tandem with a data assimilation method proposed in Azouani and Titi (2014), Azouani et al. (2014). This algorithm, commonly referred to as the Azouani–Olson–Titi (AOT) or continuous data assimilation (CDA) algorithm, has seen much recent work (see, e.g., Albanez et al. 2016; Altaf et al. 2017; Bessaih et al. 2015; Biswas et al. 2019, 2018; Biswas and Martinez 2017; Carlson et al. 2020; Carlson and Larios 2020; Celik et al. 2019; Di Leoni et al. 2018; Leoni et al. 2019; Desamsetti et al. 2019; Farhat et al. 2020, 2018, 2015, 2016a, b, ?, 2017, 2019; Foias et al. 2016; Foyash et al. 2014; García-Archilla et al. 2020; Gardner et al. 2021; Gesho et al. 2016; Glatt-Holtz et al. 2014; Hudson and Jolly 2019; Ibdah et al. 2019; Jolly et al. 2019, 2017; Larios et al. 2018; Larios and Victor 2021; Lunasin and Titi 2017; Markowich et al. 2016; Mondaini and Titi 2018; Pei 2019; Rebholz and Zerfas 2021; Zerfas et al. 2019; Larios and Pei 2017 and the references therein). Specifically, Azouani et al. (2014) consider the 2D Navier–Stokes system, written abstractly in the form

$$\frac{du}{dt} = F_v(u).$$

The difficulty is that the initial data are unknown; however, it is assumed that the solution can be measured at certain points. In order to converge to the correct solution, it is proposed to instead consider the system

¹ Note that considerations of sensitivity arise in the context of perturbations; hence, the natural initial data for a sensitivity equation is identically zero data.

$$\begin{cases} \frac{dv}{dt} = F_v(v) + \mu(I_h(u) - I_h(v)) \\ v(0) = v_0, \end{cases}$$

where $\mu > 0$ is a sufficiently large positive relaxation parameter, $I_h(u)$ represents the observational measurements with sufficiently small spacing $h > 0$, v_0 is arbitrarily chosen in a certain function space, and F_v is a nonlinear, nonlocal differential operator depending on the viscosity parameter $v > 0$. I_h is an interpolation operator satisfying certain bounds (see Sect. 2); e.g., I_h could be piecewise linear interpolation at given nodes. In Azouani et al. (2014), it was proved that v converges to u exponentially fast in time in certain standard norms. Later, Carlson et al. (2020) investigated the case of an unknown viscosity, gave estimates for the resulting error in the solution, proposed an algorithm to recover the unknown viscosity, and demonstrated computationally that the algorithm converges exponentially fast in time to the correct solution. However, the algorithm in Carlson et al. (2020) introduces a discontinuous change in the viscosity during the simulation, leading to a desire to ensure that this abrupt change does not lead to undesirable behavior. Hence, we also prove that the difference quotient methods developed here can be used to prove rigorous results for the sensitivity equations of the modified system of equations via the data assimilation algorithm. For this system, we prove that the derivative of solutions with respect to the viscosity is a well-defined object which is bounded in appropriate function spaces; additionally, we prove that the corresponding sensitivity equations are globally well posed in time in an appropriate sense and that strong solutions are unique.

Furthermore, one can also perceive the importance of analyzing the sensitivity of a system of equations to a parameter-based regularization. Hence, we also prove rigorously that the sensitivity equations for a viscous regularization (i.e., a Navier–Stokes-like perturbation) of the 2D incompressible Euler equations in 2D (under periodic boundary conditions) have a unique strong solution that can be realized as a sequence of difference quotients.

Sensitivity for partial differential equations has been studied formally in many contexts; see, e.g., Anderson et al. (1999), Borggaard and Burns (1997), Breckling et al. (2018), Brewer (1982), Davis and Pahlevani (2013), Fernández and Moubachir (2002), Gibson and Clark (1997), Hamby (1994), Kim et al. (1999), Kouhi et al. (2016), Neda et al. (2016), Noacco et al. (2019), Pahlevani (2004), Pahlevani (2006), Rebholz et al. (2017), Stanley and Stewart (2002), Vemuri and Raefsky (1979), Zerfas et al. (2019). In Stanley and Stewart (2002), it was argued, though only formally, that the sensitivity equations for the steady-state 2D Navier–Stokes equations are globally well posed. Some analysis for the sensitivity equations has been carried out in the slightly more general context of a large eddy simulation (LES) model of the 2D Navier–Stokes equations in an unpublished PhD thesis (Pahlevani 2004), where a formal argument for the global existence and uniqueness of the equations was given, based on formal energy estimates.

The paper is organized as follows. In Sect. 2, we describe the mathematical framework for the problems we consider. In Sect. 3, we prove the global existence and uniqueness of solutions to the sensitivity equations. Moreover, we show that these solutions can be realized as limits of difference quotients. In Sect. 4, we extend the

results in the previous section to the context of AOT data assimilation algorithm. Finally, we summarize our results and implications of this work in Sect. 5.

2 Preliminaries

In this section, we lay out notation and recall some standard results about the incompressible Navier–Stokes equations. Proofs can be found in, e.g., Constantin and Foias (1988), Foias et al. (2001), Robinson (2001), Temam (2001), Temam (1995). Similarly, equivalent results for the modified data assimilation equations given by the AOT algorithm are stated without proof as well, since proofs were given in Azouani et al. (2014). On a general open spatial domain Ω , the incompressible Navier–Stokes equations are given by

$$u_t + (u \cdot \nabla)u = -\nabla p + \nu \Delta u + f, \quad \text{in } \Omega \times [0, T], \quad (2.1a)$$

$$\nabla \cdot u = 0, \quad \text{in } \Omega \times [0, T], \quad (2.1b)$$

$$u(x, 0) = u_0(x), \quad \text{in } \Omega. \quad (2.1c)$$

where $\nu > 0$ is the kinematic viscosity, u is the velocity, p is the (density normalized) pressure, and f is a (density normalized) body force. In this paper, we take Ω to be the torus, i.e., $\Omega = \mathbb{T}^2 = \mathbb{R}^2 / \mathbb{Z}^2$, which is an open, bounded, and connected domain with empty boundary. We denote the space

$$\mathcal{V} := \{f : \Omega \rightarrow \mathbb{R}^2 \mid f \in \dot{C}_p^\infty(\mathbb{T}^2)\},$$

where $\dot{C}_p^\infty(\mathbb{T}^2)$ is the set of all mean-free, periodic, and infinitely differentiable functions on \mathbb{T}^2 . Denote by H the closure of \mathcal{V} in $L^2(\Omega; \mathbb{R}^2)$ and by V the closure of \mathcal{V} in $H^1(\Omega; \mathbb{R}^2)$. Since H and V are closed subspaces of $L^2(\Omega; \mathbb{R}^2)$ and $H^1(\Omega; \mathbb{R}^2)$, respectively, they are Hilbert spaces which inherit inner products denoted by

$$(u, v) = \int_{\mathbb{T}^2} u \cdot v \, dx \quad ((u, v)) = \sum_{i,j=1}^2 \int_{\mathbb{T}^2} \frac{\partial u_i}{\partial x_j} \frac{\partial v_i}{\partial x_j} \, dx,$$

with the corresponding norms denoted by $|u| = \sqrt{(u, u)}$ and $\|u\| = \sqrt{((u, u))}$. We denote by $L^p(0, T; X)$ those functions which are $L^p((0, T))$ in time with values in the Banach space X in the sense of the Bochner integral; similarly for expressions such as $C^k([0, T]; X)$. We denote the dual space of V by V' .

We consider the equivalent problem applying the Leray projection to (2.1), where the Leray projection onto divergence-free vector fields is defined by $P_\sigma u := u - \nabla \Delta^{-1} \nabla \cdot u$, $P_\sigma : L^2(\Omega) \rightarrow H$. As in Azouani et al. (2014), we denote the Stokes operator $A : \mathcal{D}(A) \rightarrow H$, where $\mathcal{D}(A) := \{u \in V : Au \in H\}$ is the domain of A , and the bilinear term $B : V \times V \rightarrow V'$ as the continuous extensions of the operators A , defined on \mathcal{V} , and B , defined on $\mathcal{V} \times \mathcal{V}$, by

$$Au = -P_\sigma \Delta u \quad \text{and} \quad B(u, v) = P_\sigma(u \cdot \nabla v).$$

Due to boundedness of the domain and the mean-zero condition, the following Poincaré inequalities hold:

$$\lambda_1 \|u\|_{L^2}^2 \leq \|\nabla u\|_{L^2}^2 \quad \text{for } u \in V, \quad (2.2)$$

$$\lambda_1 \|\nabla u\|_{L^2}^2 \leq \|Au\|_{L^2}^2 \quad \text{for } u \in \mathcal{D}(A). \quad (2.3)$$

We note that, as proved in, e.g., Constantin and Foias (1988), Robinson (2001), Temam (2001), A is a linear self-adjoint positive definite operator with a compact inverse. Hence, there exists a complete orthonormal set of eigenfunctions $\{w_i\}_{i \in \mathbb{N}}$ in V such that $Aw_i = \lambda_i w_i$, where the corresponding eigenvalues $\{\lambda_i\}_{i \in \mathbb{N}}$ are positive and monotonically increasing. Hence, for $s \in \mathbb{R}$, the fractional operators A^s may be defined via $A^s w_i = \lambda_i^s w_i$ and extension to appropriate spaces by linearity and density.

As proved in, e.g., Constantin and Foias (1988), Robinson (2001), Temam (2001), the bilinear operator, B , has the property

$$\langle B(u, v), w \rangle_{V', V} = -\langle B(u, w), v \rangle_{V', V}, \quad (2.4)$$

for all $u, v, w \in V$, which directly implies that

$$\langle B(u, w), w \rangle_{V', V} = 0, \quad (2.5)$$

for all $u, v, w \in V$. Furthermore, we have the following inequalities:

$$|\langle B(u, v), w \rangle| \leq \|u\|_{L^\infty(\Omega)} \|v\| \|w\| \quad \text{for } u \in L^\infty(\Omega), v \in V, w \in H \quad (2.6)$$

$$|\langle B(u, v), w \rangle| \leq c|u|^{1/2} \|u\|^{1/2} \|v\| \|w\|^{1/2} \|w\|^{1/2} \quad \text{for } u, v, w \in V, \quad (2.7)$$

$$|(B(u, v), w)| \leq c|u|^{1/2} \|u\|^{1/2} \|v\|^{1/2} |Av|^{1/2} |w| \quad \text{for } u \in V, v \in \mathcal{D}(A), w \in H \quad (2.8)$$

$$|(B(u, v), w)| \leq c|u|^{1/2} |Au|^{1/2} \|v\| |w| \quad \text{for } u \in \mathcal{D}(A), v \in V, w \in H. \quad (2.9)$$

Due to the periodic boundary conditions, it also holds (in 2D) that

$$(B(w, w), Aw) = 0 \quad \text{for every } w \in \mathcal{D}(A). \quad (2.10)$$

From this and some manipulation, it follows that for $u, w \in \mathcal{D}(A)$,

$$(B(u, w), Aw) + (B(w, u), Aw) = -(B(w, w), Au). \quad (2.11)$$

Additionally, further properties of the bilinear term are stated in Lemma 2.1, which we prove using similar strategies as in Robinson (2001), Temam (2001).

Lemma 2.1 Suppose $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$ are uniformly bounded sequences in $L^2(0, T; V) \cap L^\infty(0, T; H)$. Then $\|B(a_n, b_n)\|_{L^2(0, T; V')}$ is uniformly bounded in n . Moreover, if $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$ are uniformly bounded in $L^2(0, T; \mathcal{D}(A)) \cap L^\infty(0, T; V)$, then $\|B(a_n, b_n)\|_{L^2(0, T; H)}$ is uniformly bounded in n .

Proof By the definition of the dual norm, (2.4), and (2.7)

$$\|B(a_n, b_n)\|_{V'} = \sup_{\substack{w \in V \\ \|w\|=1}} |(B(a_n, w), b_n)| \leq k|a_n|^{1/2}\|a_n\|^{1/2}|b_n|^{1/2}\|b_n\|^{1/2}.$$

Using Hölder's inequality,

$$\begin{aligned} \|B(a_n, b_n)\|_{L^2(0, T; V')}^2 &\leq \int_0^T k^2|a_n|\|a_n\|\|b_n\|\|b_n\| \, ds \\ &\leq k^2\|a_n\|_{L^\infty(0, T; H)}\|b_n\|_{L^\infty(0, T; H)}\|a_n\|_{L^2(0, T; V)}\|b_n\|_{L^2(0, T; V)}. \end{aligned}$$

Hence, since $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$ are uniformly bounded in $L^2(0, T; V) \cap L^\infty(0, T; H)$, it follows that $\|B(a_n, b_n)\|_{L^2(0, T; V')}$ is uniformly bounded in n .

Next, suppose $\{a_n\}_{n \in \mathbb{N}}$ and $\{b_n\}_{n \in \mathbb{N}}$ are bounded uniformly in $L^2(0, T; \mathcal{D}(A)) \cap L^\infty(0, T; V)$. Then by definition of the dual norm, (2.8), and (2.2),

$$\begin{aligned} \|B(a_n, b_n)\|_{L^2(0, T; H)}^2 &\equiv \int_0^T |B(a_n, b_n)|^2 \, dt \leq c \int_0^T |a_n|\|a_n\|\|b_n\|\|Ab_n\| \, dt \\ &\leq \frac{c}{\lambda_1} \|a_n\|_{L^\infty(0, T; V)}^2 \int_0^T |Ab_n|^2 \, dt \\ &= \frac{c}{\lambda_1} \|a_n\|_{L^\infty(0, T; V)}^2 \|b_n\|_{L^2(0, T; \mathcal{D}(A))}^2, \end{aligned}$$

which implies that $\|B(a_n, b_n)\|_{L^2(0, T; H)}$ is uniformly bounded in n . \square

Finally, without loss of generality, we make the assumption that $f \in L^\infty(0, T; H)$ so that $P_\sigma f = f$. This allows us to apply P_σ to (2.1) to obtain the equivalent set of equations

$$\frac{d}{dt}u + B(u, u) = -\nu Au + f, \quad \text{in } \Omega \times [0, T], \quad (2.12a)$$

$$u(x, 0) = u_0(x), \quad \text{in } \Omega. \quad (2.12b)$$

Using the following corollary of de Rham's theorem Temam (2001), Foias et al. (2001)

$$g = \nabla p \text{ with } p \text{ a distribution if and only if } \langle g, h \rangle = 0 \text{ for all } h \in \mathcal{V}, \quad (2.13)$$

one can recover the pressure term. Furthermore, we will utilize the following lemma due to Lions and Magenes as stated in Temam (2001).

Lemma 2.2 (Temam 2001) *Let V, H, V' be three Hilbert spaces such that $V \subset H \equiv H' \subset V'$, where V' is the dual of V . If a function $u \in L^2(0, T; V)$ and its time derivative $u' \in L^2(0, T; V')$, then u is almost everywhere equal to a function continuous from $[0, T]$ into H and the following equality holds in the scalar distribution sense on $(0, T)$:*

$$\frac{d}{dt}|u|^2 = 2\langle u', u \rangle.$$

For reference, we state the standard definition of weak and strong solutions to (2.12).

Definition 2.3 Let $T > 0$. Let initial data $u_0 \in H$ and forcing $f \in L^\infty(0, \infty; V')$. A *weak solution* of (2.1) is an element $u \in L^2(0, T; V) \cap C([0, T]; H)$ satisfying $\frac{du}{dt} \in L^2(0, T; V')$ and

$$\left\langle \frac{d}{dt}u, \phi \right\rangle + \langle B(u, u), \phi \rangle + \nu \langle Au, \phi \rangle = 0 \quad (2.14)$$

for a.e. $t \in [0, T]$ and for all $\phi \in V$. Furthermore, the map $u_0 \in H \mapsto u \in H$ is continuous for a.e. $t \in [0, T]$, i.e., the initial data are satisfied in the sense of $C([0, T]; H)$.

If, in addition, $f \in L^\infty(0, \infty; H)$ and $u_0 \in V$, then a *strong solution* of (2.1) is defined to be a weak solution such that $u \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ and $\frac{d\tilde{u}}{dt} \in L^2(0, T; H)$, and moreover u satisfies (2.14) for a.e. $t \in [0, T]$ and for all $\phi \in H$. Furthermore, the map $u_0 \in V \mapsto u \in V$ is continuous for a.e. $t \in [0, T]$, i.e., the initial data are satisfied in the sense of $C([0, T]; V)$.

It is well established that given a force $f \in L^2(0, \infty; H^{-1})$ and initial data $u_0 \in H$, a unique weak solution to (2.12) exists globally in time (see, e.g., Constantin and Foias 1988; Foias et al. 2001; Robinson 2001; Temam 2001). For many real-world applications, it is important to consider the sensitivity of (2.1) to the parameters since it is not necessarily the case we have an exact estimate on the said parameters; however, additional uncertainty in modeling real-world systems is introduced by the fact that we do not expect to know u_0 exactly, and so cannot compute $u(t)$ from (2.12). Hence, we will also analyze the sensitivity equations corresponding to a modified system of equations that utilizes measured data collected on the true field $u(t)$ over the time interval $[0, T]$. This modified system of equations incorporates the measured data by introducing a feedback control involving the interpolated data $I_h(u(t))$ into (2.12), resulting in the following system which was proposed and studied in Azouani et al. (2014), Azouani and Titi (2014) (see also Blömker et al. 2013).

$$v_t + (v \cdot \nabla)v = -\nabla q + \nu_{\text{appx}} \Delta v + f + \mu(I_h(u) - I_h(v)) \quad (2.15a)$$

$$v(x, 0) = v_0(x), \quad (2.15b)$$

or, with the Leray projection applied,

$$\frac{d}{dt}v + B(v, v) = -\nu_{\text{appx}}Av + f + \mu P_\sigma(I_h(u) - I_h(v)) \quad (2.16a)$$

$$v(x, 0) = v_0(x). \quad (2.16b)$$

Here, $\nu_{\text{appx}} > 0$ is a kinematic viscosity approximating $\nu > 0$, $\mu > 0$ is a positive relaxation parameter, and I_h is a linear operator satisfying

$$\|\varphi - I_h(\varphi)\|_{L^2(\Omega)}^2 \leq c_0 h^2 \|\varphi\|_{H^1(\Omega)}^2. \quad (2.17)$$

Assuming either no-slip Dirichlet or periodic boundary conditions (and allowing for even more general interpolation operators I_h), Azouani et al. (2014) proved that (2.16) has a unique solution, stated in the following theorem.

Theorem 2.4 *Suppose I_h satisfies (2.17) and $\mu c_0 h^2 \leq \nu_{\text{appx}}$, where c_0 is the constant from (2.17). Then the continuous data assimilation equations (2.16) possess a unique solution v that satisfies*

$$v \in C([0, T]; V) \cap L^2((0, T); \mathcal{D}(A)) \quad \text{and} \quad \frac{dv}{dt} \in L^2((0, T); H), \quad (2.18)$$

for any $T > 0$. Furthermore, the map $u_0 \in V \mapsto u \in V$ is continuous for a.e. $t \in [0, T]$, i.e., the initial data are satisfied in the sense of $C([0, T]; V)$.

For Eqs. (2.12) and (2.16), we denote the dimensionless Grashof numbers as

$$G_1 = \frac{1}{4\pi^2 \nu^2} \limsup_{t \rightarrow \infty} \|f(t)\|_{L^2(\Omega)} \quad (2.19)$$

$$G_2 = \frac{1}{4\pi^2 \nu_{\text{appx}}^2} \limsup_{t \rightarrow \infty} \|f(t)\|_{L^2(\Omega)}. \quad (2.20)$$

Finally, we consider the Euler equations in 2D without forcing:

$$u_t + u \cdot \nabla u + \nabla p = 0, \quad (2.21a)$$

$$\nabla \cdot u = 0 \quad (2.21b)$$

$$u(x, 0) = u_0. \quad (2.21c)$$

Applying the Leray projection yields the system

$$\frac{du}{dt} + B(u, u) = 0, \quad (2.22a)$$

$$u(x, 0) = u_0. \quad (2.22b)$$

The relevant literature on the inviscid limit of solutions to (2.12) to sufficiently smooth solutions of Euler is usually considered in the whole space \mathbb{R}^2 Kato (1972);

Swann (1971), Golovkin (1966), Ponce (1986), Constantin and Foias (1988), but we utilize the relevant theorems straightforwardly extendable to the periodic setting (see, e.g., Masmoudi 2007). We state the types of solutions to Euler that are limits of solutions to (2.1).

Theorem 2.5 (Masmoudi 2007) *Let $s > 0$ and suppose $u \in C([0, T); \mathcal{D}(A^{s/2}))$ is a solution of the Euler system (2.22) for some $T > 0$ with $u_0 \in \mathcal{D}(A^{s/2})$. Let $u_0^n \in \mathcal{D}(A^{s/2})$ such that $u_0^n \rightarrow u_0$ in $\mathcal{D}(A^{s/2})$ as $n \rightarrow \infty$. Then, for all $0 < T_0 < T$, for any sequence $v_n \rightarrow 0$, the corresponding unique sequence of solutions $u^n \in C([0, T_0]; \mathcal{D}(A^{s/2}))$ to (2.1) with $f \equiv 0$ is such that*

$$\|u^n - u\|_{L^\infty(0, T_0; \mathcal{D}(A^{s/2}))} \rightarrow 0, \quad n \rightarrow \infty.$$

For simplicity, we take $f \equiv 0$ for a direct application of this theorem from Masmoudi (2007). We also note that since we are in the periodic domain \mathbb{T}^2 instead of \mathbb{R}^2 , this implies that convergence of $u_n \rightarrow u$ also holds in $L^\infty(0, T_0; \mathcal{D}(A^{s'/2}))$ for all $s' \leq s$.

3 Sensitivity for 2D Navier–Stokes & The Inviscid Limit

In this section, we analyze the sensitivity of w to the viscosity by considering individually the sensitivity of u and v to the viscosity. We wish to consider taking a derivative of equations (2.12a) and (2.16a) with respect to the viscosity. This has been done formally in many works on sensitivity (see, e.g., Anderson et al. 1999; Borggaard and Burns 1997; Breckling et al. 2018; Davis and Pahlevani 2013; Fernández and Moubachir 2002; Hamby 1994; Kim et al. 1999; Kouhi et al. 2016; Pahlevani 2004, 2006; Vemuri and Raefsky 1979), yielding what are known as the *sensitivity equations*. However, to the best of our knowledge, a rigorous treatment has yet to appear in the literature. Therefore, we provide a rigorous justification here of the existence and uniqueness of weak and strong solutions to the sensitivity equations in the case of zero initial data, which is the natural data for the sensitivity equation, as discussed below. Moreover, we prove that these solutions can be realized as limits of difference quotients of Navier–Stokes solutions with respect to different viscosities. Indeed, this is the method of our existence proofs, rather than using, e.g., Galerkin methods, fixed-point methods, etc. Proofs using limits of difference quotients have appeared in the literature before, such as in standard proofs of elliptic regularity and the corresponding result for the Stokes equations. However, in the present context (i.e., the time-dependent sensitivity equations for 2D Navier–Stokes), we believe such a proof strategy is novel.

Working formally for a moment, we take the derivative of (2.1) with respect to ν , and denote (again, formally) $\tilde{u} := \frac{du}{d\nu}$ and $\tilde{p} := \frac{dp}{d\nu}$, to obtain

$$\tilde{u}_t + \tilde{u} \cdot \nabla u + u \cdot \nabla \tilde{u} - \nu \Delta \tilde{u} - \Delta u + \nabla \tilde{p} = 0, \quad (3.1a)$$

$$\nabla \cdot \tilde{u} = 0. \quad (3.1b)$$

These are known as the sensitivity equations for the Navier–Stokes equations. Similarly we formally take the derivative of (2.15) with respect to v_{appx} , denoting $\tilde{v} := \frac{dv}{dv_{\text{appx}}}$ and $\tilde{q} := \frac{dq}{dv_{\text{appx}}}$,

$$\tilde{v}_t + \tilde{v} \cdot \nabla v + v \cdot \nabla \tilde{v} - v_{\text{appx}} \Delta \tilde{v} - \Delta v + \nabla \tilde{q} = \mu I_h(\tilde{u} - \tilde{v}), \quad (3.2a)$$

$$\nabla \cdot \tilde{v} = 0. \quad (3.2b)$$

Below, we prove some well-posedness results for these systems in the case of zero initial data. We begin by defining what we mean by solutions.

Remark 3.1 The following proofs follow mostly standard techniques; however, they establish that the solutions of the sensitivity equations can indeed be realized as limits of difference quotients, which is the first time this has been done rigorously. Moreover, we note that the analysis for the sensitivity of the Euler equations with respect to viscous perturbations is done via methods typically used in Navier–Stokes-type analysis rather than methods used in the analysis of the Euler equations. Furthermore, this method of proof highlights the independence of the initial data on the parameter, and specifies what one should mean when referring to the sensitivity equations.

Definition 3.2 Let $T > 0$. Let $u \in L^2(0, T; V) \cap C([0, T]; H)$ be a weak solution to (2.1). A *weak solution* of (3.1) is a weak solution in the sense of Definition 2.3, but with the equation instead given by

$$\left\langle \frac{d}{dt} \tilde{u}, \phi \right\rangle + \langle B(\tilde{u}, u), \phi \rangle + \langle B(u, \tilde{u}), \phi \rangle + v \langle A\tilde{u}, \phi \rangle + \langle Au, \phi \rangle = 0. \quad (3.3)$$

If, in addition, $\tilde{u}_0 \in V$, and u is a strong solution to (2.1), then we similarly define a *strong solution* of (3.1) in the same sense as Definition 2.3 with respect to the above equation.

As discussed in Remark 3.6, we only give a definition of strong solutions for the assimilation equations to avoid redundancy.

Definition 3.3 Let $T > 0$. Let v be a strong solution to (2.16) with initial data $v_0 \in V$ and forcing $f \in L^\infty(0, \infty; H)$. A strong solution of (3.2) is a strong solution in the sense of Definition 2.3 with the equation given by

$$\left\langle \frac{d}{dt} \tilde{v}, \phi \right\rangle + \langle B(\tilde{v}, v), \phi \rangle + \langle B(v, \tilde{v}), \phi \rangle + v_{\text{appx}} \langle A\tilde{v}, \phi \rangle + \langle Av, \phi \rangle = \mu \langle I_h(\tilde{u} - \tilde{v}), \phi \rangle.$$

Before we prove the existence and uniqueness of solutions with zero initial data to these equations, we first consider equations for the difference quotients. Note that, since these are simple arithmetic operations on the Navier–Stokes equations, the manipulations can be performed rigorously, not just formally. To this end, let (u_1, p_1) be a strong solution to (2.1) with viscosity ν_1 and (u_2, p_2) be a strong solution to (2.1) with viscosity ν_2 with the same initial data. We take the difference of the two versions of

(2.1), each with viscosities ν_1 and ν_2 . We then divide by the difference in viscosities, yielding the system

$$w_t + u_2 \cdot \nabla w + w \cdot \nabla u_1 - \nu_2 \Delta w - \Delta u_1 + \nabla P = 0, \quad (3.4a)$$

$$\nabla \cdot w = 0, \quad (3.4b)$$

$$w(x, 0) = 0, \quad (3.4c)$$

where $w = \frac{u_1 - u_2}{\nu_1 - \nu_2}$ and $P := \frac{p_1 - p_2}{\nu_1 - \nu_2}$. As defined, w is a strong solution to (3.4), and note that $u_1 = (\nu_1 - \nu_2)w + u_2$. Additionally, $w \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ and $\frac{dw}{dt} \in L^2(0, T; H)$. However, we need to establish that w is the unique solution to (3.4), which is the content of Lemma 3.4.

Lemma 3.4 *Let $T > 0$ be given, and let $u_1, u_2 \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ be strong solutions to (2.16), with viscosities ν_1 and ν_2 , respectively. There exists one and only one solution w to (3.4) that lies in $L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$, i.e., for all $\phi \in H$,*

$$\left(\frac{d}{dt} w, \phi \right) + (B(w, u_1), \phi) + (B(u_2, w), \phi) + \nu_2 (Aw, \phi) + (Au_1, \phi) = 0,$$

where $\frac{d}{dt} w \in L^2(0, T; H)$. Moreover, the map $u_0 \in V \mapsto u \in V$ is continuous for a.e. $t \in [0, T]$, i.e., the solution depends continuously on the initial data.

Next, we consider difference quotients for the assimilation system (2.16). Let (v_1, q_1) be the strong solution to (2.16) with viscosity $\nu_{\text{appx},1}$ and (v_2, q_2) be the strong solution to (2.16) with viscosity $\nu_{\text{appx},2}$. Subtracting the two equations and dividing by the difference in the viscosities yields

$$\tilde{w}_t + \tilde{w} \cdot \nabla v_1 + v_2 \cdot \nabla \tilde{w} - \nu_{\text{appx},2} \Delta \tilde{w} - \Delta v_1 + \nabla Q = \mu I_h(w - \tilde{w}) \quad (3.5a)$$

$$\nabla \cdot \tilde{w} = 0 \quad (3.5b)$$

$$\tilde{w}(x, 0) = 0, \quad (3.5c)$$

where $\tilde{w} := \frac{v_1 - v_2}{\nu_{\text{appx},1} - \nu_{\text{appx},2}}$ and $Q := \frac{q_1 - q_2}{\nu_{\text{appx},1} - \nu_{\text{appx},2}}$. It follows from the definition that \tilde{w} is a strong solution to (3.5). Moreover, note that $v_1 = (\nu_{\text{appx},1} - \nu_{\text{appx},2})\tilde{w} + v_2$. Additionally, $\tilde{w} \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ and $\frac{d}{dt} \tilde{w} \in L^2(0, T; H)$.

Lemma 3.5 *Let $T > 0$ be given, and let $v_1, v_2 \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ be strong solutions to (2.16), with viscosities $\nu_{\text{appx},1}$ and $\nu_{\text{appx},2}$, respectively. There exists a unique strong solution w to (3.5) that lies in $L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$, in the sense that for all $\phi \in H$,*

$$\begin{aligned} & \left(\frac{d}{dt} \tilde{w}, \phi \right) + (B(v_2, \tilde{w}), \phi) + (B(\tilde{w}, \nabla v_1), \phi) + \nu_{\text{appx},2} (A\tilde{w}, \phi) + (Av_1, \phi) \\ &= \mu (P_\sigma I_h(w - \tilde{w}), \phi), \end{aligned}$$

where $\frac{d}{dt}\tilde{w} \in L^2(0, T; H)$. Moreover, the map $u_0 \in V \mapsto u \in V$ is continuous for a.e. $t \in [0, T]$, i.e., the solution depends continuously on the initial data.

Remark 3.6 The proofs of the above two lemmata are very similar; hence, we only present the proof of Lemma 3.5. Moreover, we also note that in the case $\mu = 0$, the proof of Lemma 3.4 holds *mutatis mutandis* in the case where $u_1, u_2 \in C([0, T]; H) \cap L^2(0, T; V)$ are only assumed to be weak solutions to the 2D Navier–Stokes equations, and then one obtains uniqueness of weak solutions to (3.4) in the class $C([0, T]; H) \cap L^2(0, T; V)$. However, in the case $\mu > 0$, the notion of weak solutions for the assimilation equations (2.16) has not been established in the literature for general interpolants I_h , and therefore, we assume that the solutions v_1 and v_2 are strong solutions to (2.16), and prove the uniqueness of strong solutions to (3.4). The existence of weak solutions is attainable if further assumptions are made on I_h (e.g., if I_h is a projection onto low Fourier modes) but we do not pursue this here as we want to focus on general interpolants.

Proof Suppose there exist two solutions \tilde{w}_1 and \tilde{w}_2 . We consider the difference of the equations

$$\frac{d}{dt}\tilde{w}_1 + B(\tilde{w}_1, v_1) + B(v_2, \tilde{w}_1) + v_{\text{appx},2}A\tilde{w}_1 + Av_1 = \mu P_\sigma I_h(w - \tilde{w}_1) \quad (3.6)$$

and

$$\frac{d}{dt}\tilde{w}_2 + B(\tilde{w}_2, v_1) + B(v_2, \tilde{w}_2) + v_{\text{appx},2}A\tilde{w}_2 + Av_1 = \mu P_\sigma I_h(w - \tilde{w}_2). \quad (3.7)$$

Namely, denoting $W := \tilde{w}_1 - \tilde{w}_2$ yields

$$\frac{d}{dt}W + B(W, v_1) + B(v_2, W) + v_{\text{appx},2}AW = -\mu P_\sigma I_h(W) \quad (3.8)$$

with $W(0) = 0$. So, W must be a solution to the above equation. Taking the inner product with AW , one can follow the argument in Azouani et al. (2014) line by line to obtain to obtain that

$$\|W(t)\|^2 \leq \|W(0)\|^2 e^{ct} \quad (3.9)$$

for some constant $c > 0$. Since $W(0) = 0$, this implies that $W \equiv 0$. Hence, solutions to (3.5) are unique. Moreover, (3.9) gives the continuous dependence on initial data. \square

Since systems (3.4) and (3.5) have unique strong solutions for every $\nu > 0$, we want to show that, as $\nu \rightarrow \nu_0$ for some fixed $\nu_0 > 0$, the solutions to these equations converge to the unique strong solutions of the respective equations (in the sense of Definitions 3.2 and 3.3) of the formal sensitivity equations (3.1) and (3.2) with initial data $u_0 \equiv 0$ in Theorems 3.9 and 4.1. We additionally prove that weak solutions exist for the sensitivity equations (3.1) with initial data $u_0 \equiv 0$ in Theorem 3.7. Via this method of

proof, we show that the solution \tilde{u} is a Fréchet derivative in the sense of $L^2(0, T; H)$. Specifically, for weak solutions, $v \in \mathbb{R}^+ \mapsto u \in L^2(0, T; H)$, and hence the *Fréchet derivative in $L^2(0, T; H)$* will be defined as $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; H)$ such that $\mathcal{A}(\delta) \mapsto \delta\tilde{u}$. Analogously, for strong solutions, $v \in \mathbb{R}^+ \mapsto u \in L^2(0, T; V)$, and hence, the *Fréchet derivative in $L^2(0, T; V)$* will be defined as $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$ such that $\mathcal{A}(\delta) \mapsto \delta\tilde{u}$.

Theorem 3.7 *Let $T > 0$ and $\{v_n\}_{n \in \mathbb{N}}$ be a sequence such that $v_n \rightarrow v$ as $n \rightarrow \infty$. Let*

- *u be a solution to (2.1) with viscosity v, forcing $f \in L^\infty(0, \infty; H)$, and initial data $u_0 \in V$;*
- *u_n solve (2.1) with viscosity v_n , forcing $f \in L^\infty(0, \infty; H)$, and initial data $u_0 \in V$;*
- *$\{w_n\}_{n \in \mathbb{N}}$ be a sequence of strong solutions to (3.4) with $w_n(0) = 0$ with corresponding viscosity v_n and solutions to (2.1) u_n and u.*

Then there exists a subsequence of $\{w_n\}_{n \in \mathbb{N}}$ that converges in $L^2(0, T; H)$ to a unique weak solution \tilde{u} of (3.1) with initial data $u_0 \equiv 0$. Furthermore, the operator $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; H)$ given by $\mathcal{A}(\delta) = \tilde{u}\delta$ is the Fréchet derivative in the sense of $L^2(0, T; H)$ of the solution u with respect to v.

Remark 3.8 Theorem 3.7 holds for more general initial data via a similar proof under certain assumptions on the initial data; indeed, assuming that the initial data of (2.1) has a derivative with respect to the viscosity and that it lies in an appropriate space and is a limit of difference quotients of the initial data corresponding to the sequence of solutions of (2.1), one can show a similar result. However, note that the sensitivity equations are a model for the evolution of the instantaneous change in a solution with respect to changes in the viscosity, hence the natural initial condition to consider is the case of identically zero initial data. Thus, to avoid obfuscation, we work in the natural setting of identically zero initial data; however, we provide a definition of weak and strong solutions for above general initial data. Similar remarks holds for all theorems below.

Proof Let $T > 0$ be given. Let N sufficiently large such that for all $n > N$, $\{v_n\}_{n \in \mathbb{N}} \subset (\frac{v}{2}, \frac{3v}{2})$. Then, one can follow the proof of strong solutions for (2.1) as in, e.g., Constantin and Foias (1988), Foias et al. (2001), Robinson (2001), Temam (2001), to obtain bounds on $\{u_n\}$ for $n > N$ in the appropriate spaces that are independent of v_n :

$$\|u_n\|_{L^\infty(0, T; V)}^2 \leq \|u_n(0)\|^2 + \frac{\|f\|_{L^2(0, T; H)}^2}{v_n} \leq \|u_0\|^2 + \frac{2\|f\|_{L^2(0, T; H)}^2}{v}$$

and

$$\|u_n\|_{L^2(0, T; \mathcal{D}(A))}^2 \leq \frac{1}{v_n} \|u_n(0)\|^2 + \frac{\|f\|_{L^2(0, T; H)}^2}{v_n^2} \leq \frac{2}{v} \|u_0\|^2 + \frac{4\|f\|_{L^2(0, T; H)}^2}{v^2}.$$

Note that $\|f\|_{L^2(0, T; H)}^2 < \infty$ since all bounded functions are locally integrable. Hence there is a subsequence that is relabeled $u_n \rightarrow u_1$ in $L^2(0, T; V)$ for some function u_1 . Continuing to follow the proof of strong solutions for (2.1) as in, e.g., Constantin and Foias (1988), Foias et al. (2001), Robinson (2001), Temam (2001), we note that $\frac{du_n}{dt}$ is uniformly bounded in n in $L^2(0, T; H)$. Hence, we can find a subsequence which we relabel $\{u_n\}$ such that

$$\begin{aligned} \frac{du_n}{dt} &\rightarrow \frac{du_1}{dt} && \text{in } L^2(0, T; H) \\ v_n Au_n &\rightarrow v Au_1 && \text{in } L^2(0, T; H) \\ B(u_n, u_n) &\rightarrow B(u_1, u_1) && \text{in } L^2(0, T; H). \end{aligned}$$

Indeed, u_1 satisfies (2.1) with corresponding viscosity v and thus, by uniqueness and the fact that $u_n \rightarrow u_1$ in $L^2(0, T; V)$, it follows that $u_1 = u$.

Let w_n be the strong solution to (3.4) with $v = v_n$. Taking the action of (3.4) on w_n and using (2.7), (2.5), Lemma 2.2, and Young's inequality, we obtain

$$\frac{1}{2} \frac{d}{dt} |w_n|^2 + v_n \|w_n\|^2 \leq \frac{c^2}{v_n} \|u_1\|^2 |w_n|^2 + \frac{v_n}{4} \|w_n\|^2 + \frac{1}{2v_n} \|u_1\|^2 + \frac{v_n}{2} \|w_n\|^2.$$

Hence,

$$\frac{1}{2} \frac{d}{dt} |w_n|^2 + \frac{v_n}{4} \|w_n\|^2 \leq \frac{c^2}{v_n} \|u\|^2 |w_n|^2 + \frac{1}{2v_n} \|u\|^2. \quad (3.10)$$

Dropping the second term on the left-hand side, we obtain

$$\frac{1}{2} \frac{d}{dt} |w_n|^2 \leq \frac{c^2}{v_n} \|u\|^2 |w_n|^2 + \frac{1}{2v_n} \|u\|^2.$$

Taking the integral with respect to time on $[0, T]$ and applying Grönwall's inequality, then for a.e. $t \in [0, T]$,

$$\begin{aligned} |w_n(t)|^2 &\leq \left[\frac{1}{v_n} \int_0^T \|u\|^2 dt \right] \exp \left(\int_0^T \frac{2c^2}{v_n} \|u\|^2 dt \right) \\ &\leq \left[\frac{2}{v} \int_0^T \|u\|^2 dt \right] \exp \left(\int_0^T \frac{4c^2}{v} \|u\|^2 dt \right) =: K_1. \end{aligned}$$

Since $u \in L^2(0, T; V)$, then w_n is bounded uniformly in $L^\infty(0, T; H)$.

Next, refraining from dropping the second term on the left-hand side of (3.10), we estimate

$$\begin{aligned} \frac{\nu_n}{4} \int_0^T \|w_n\|^2 dt &\leq \frac{c^2}{\nu_n} \int_0^T \|u\|^2 |w_n|^2 dt + \frac{1}{2\nu_n} \int_0^T \|u\|^2 dt \\ &\leq K_1 \frac{c^2}{\nu_n} \int_0^T \|u\|^2 dt + \frac{1}{2\nu_n} \int_0^T \|u\|^2 dt \end{aligned}$$

Rewriting, we obtain

$$\begin{aligned} \int_0^T \|w_n\|^2 dt &\leq K_1 \frac{4c^2}{\nu_n^2} \int_0^T \|u\|^2 dt + \frac{2}{\nu_n^2} \int_0^T \|u\|^2 dt \\ &\leq K_1 \frac{16c^2}{\nu^2} \int_0^T \|u\|^2 dt + \frac{8}{\nu^2} \int_0^T \|u\|^2 dt \end{aligned}$$

Thus, w_n is bounded above uniformly in $L^2(0, T; V)$ with respect to n . Hence, by the Banach–Alaoglu theorem, there exists a subsequence, relabeled as (w_n) , such that

$$w_n \xrightarrow{*} w \text{ in } L^\infty(0, T; H) \quad \text{and} \quad w_n \rightharpoonup w \text{ in } L^2(0, T; V). \quad (3.11)$$

Using (3.11), note that all uniform bounds in n on the terms in (3.4) in $L^2(0, T; V')$ are obtained in a similar manner to the proof of weak solutions for (2.1) except for the term $B(u_n, w_n)$. However, by Lemma 2.1,

$$\|B(u_n, w_n)\|_{L^2(0, T; V')} \leq k \|u_n\|_{L^\infty(0, T; H)} \|w_n\|_{L^\infty(0, T; H)} \|u_n\|_{L^2(0, T; V)} \|w_n\|_{L^2(0, T; V)},$$

and due to the following standard bounds on u_n (which can be found in Constantin and Foias 1988; Foias et al. 2001; Robinson 2001; Temam 2001, etc.) and the fact that $\nu_n \in (\frac{\nu}{2}, \frac{3\nu}{2})$,

$$\|u_n\|_{L^\infty(0, T; H)}^2 \leq |u_0^n|^2 + \frac{\|f\|_{L^\infty(0, T; H)}}{\lambda_1^2 \nu_n^2} \leq |u_0|^2 + \frac{4\|f\|_{L^\infty(0, T; H)}}{\lambda_1^2 \nu^2}$$

and

$$\|u_n\|_{L^2(0, T; V)}^2 \leq \frac{1}{\nu_n} |u^n(0)|^2 + \frac{\|f\|_{L^\infty(0, T; H)}^2}{\lambda_1 \nu_n^2} T \leq \frac{2}{\nu} |u_0|^2 + \frac{4\|f\|_{L^\infty(0, T; H)}^2}{\lambda_1 \nu^2} T,$$

and thus $\|B(u_n, w_n)\|_{L^2(0, T; V')}$ is bounded above uniformly in n independent of ν_n . Hence, independent of ν_n , $d w_n/dt$ is bounded uniformly in n . Thus, by the Aubin compactness theorem, $w_n \rightarrow w$ strongly in $L^2(0, T; H)$. Utilizing these bounds and the convergence properties, we can follow the usual arguments to state there exists a weak solution in the sense of Definition 3.2, where weak continuity can be proved directly in the same manner as for weak solutions to (2.12) and strong continuity follows from Lemma 2.2. Furthermore, weak continuity in H follows due to the bounds on each of the above terms. The initial condition is satisfied by construction. To prove

uniqueness, suppose that there exist two weak solutions \tilde{u}_1 and \tilde{u}_2 . We consider the difference of the equations

$$\frac{d}{dt}\tilde{u}_1 + B(\tilde{u}_1, u) + B(u, \tilde{u}_1) + \nu A\tilde{u}_1 + Au_1 = 0$$

and

$$\frac{d}{dt}\tilde{u}_2 + B(\tilde{u}_2, u) + B(u, \tilde{u}_2) + \nu A\tilde{u}_2 + Au_1 = 0,$$

which, defining $U := \tilde{u}_1 - \tilde{u}_2$, yields

$$\frac{d}{dt}U + B(U, u) + B(u, U) + \nu AU = 0$$

with $U(0) = 0$. Taking the action on U , using (2.5), and applying the Lions–Magenes Lemma 2.2,

$$\frac{1}{2} \frac{d}{dt}|U|^2 + \langle B(U, u), U \rangle + \nu \|U\|^2 = 0.$$

Thus, by (2.7) and Young's inequality,

$$\frac{1}{2} \frac{d}{dt}|U|^2 + \nu \|U\|^2 \leq c\|U\||U||u| \leq \frac{c^2}{\nu} \|u\|^2 |U|^2 + \frac{\nu}{2} \|U\|^2.$$

Dropping the second term and applying Grönwall's inequality, for a.e. $0 \leq t \leq T$,

$$|U(t)|^2 \leq |U(0)|^2 \exp\left(\int_0^T \frac{c^2}{2\nu} \|u\|^2 dt\right).$$

Thus, $\exp\left(\int_0^T \frac{c^2}{\nu} \|u\|^2 dt\right) < \infty$ for all $T > 0$ and $U(0) = 0$. Hence, $\|U\|_{L^\infty(0, T; H)} = 0$, which implies that $U \equiv 0$. Hence, weak solutions to (3.4) are unique.

Finally, we want to show that the sequence of difference quotients defines a Fréchet derivative. Let $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; H)$, $\mathcal{A}(\delta) \equiv \delta\tilde{u}$. Note that \mathcal{A} is a bounded linear operator since \tilde{u} is bounded in $L^2(0, T; H)$. Let $\delta_n = \nu_n - \nu$, so that $\delta_n \rightarrow 0$ as $n \rightarrow \infty$. Note that for any sequence $\nu_n \rightarrow \nu$, there is a subsequence such that $w_{n_k} \rightarrow \tilde{u}$ in $L^2(0, T; H)$. However, by the fact that for every subsequence of $\{w_n\}_{n \in \mathbb{N}}$ we can find a convergent subsequence, we can conclude that $w_n \rightarrow \tilde{u}$ in $L^2(0, T; H)$. We rewrite $w_n \rightarrow \tilde{u}$ in $L^2(0, T; H)$ as

$$\left\| \frac{u(\nu + \delta_n) - u(\nu)}{\delta_n} - \tilde{u} \right\|_{L^2(0, T; H)} = \frac{1}{|\delta_n|} \|u(\nu_1 + \delta_n) - u(\nu_1) - \delta_n \tilde{u}\|_{L^2(0, T; H)} \rightarrow 0.$$

Therefore, \mathcal{A} defines a Fréchet derivative from \mathbb{R} to $L^2(0, T; H)$. \square

Theorem 3.9 Let $T > 0$ and $\{v_n\}_{n \in \mathbb{N}}$ be a sequence such that $v_n \rightarrow v$ as $n \rightarrow \infty$. Let

- u be the solution to (2.1) with viscosity v , forcing $f \in L^\infty(0, \infty; H)$, and initial data u_0 ;
- u_n solve (2.1) with viscosity v_n , forcing $f \in L^\infty(0, \infty; H)$, and initial data $u_0 \in V$;
- $\{w_n\}_{n \in \mathbb{N}}$ be a sequence of strong solutions to (3.4) with $w_n(0) = 0$ with corresponding viscosity v_n and solutions to (2.1) u_n and u .

Then there exists a subsequence of $\{w_n\}_{n \in \mathbb{N}}$ that converges in $L^2(0, T; V)$ to a unique strong solution \tilde{u} of (3.1) with initial data $u_0 \equiv 0$. Furthermore, the operator $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$ given by $\mathcal{A}(\delta) = \tilde{u}\delta$ is the Fréchet derivative in $L^2(0, T; V)$ of the solution u with respect to v .

Proof Let $T > 0$ be given, and let $N > 0$ be large enough that $n > N$ implies $\{v_n\} \subset (\frac{v_1}{2}, \frac{3v_1}{2})$. Then by the argument in Theorem 3.7, we can obtain a subsequence which we relabel $\{u_n\}$ such that $u_n \rightarrow u$ in $L^2(0, T; V)$.

Consider w_n to be the strong solution to (3.4) with viscosity v_n . Taking a justified inner product of (3.4) with Aw_n ,

$$\frac{1}{2} \frac{d}{dt} \|w_n\|^2 + v_n |Aw_n|^2 = -(B(w_n, u), Aw_n) - (B(u_n, w_n), Aw_n) - (Au, Aw_n).$$

Applying Young's inequality, we obtain

$$\frac{1}{2} \frac{d}{dt} \|w_n\|^2 + \frac{v_n}{2} |Aw_n|^2 \leq -(B(w_n, u), Aw_n) - (B(u_n, w_n), Aw_n) + \frac{1}{2v_n} |Au|^2.$$

Applying (2.6) to the second bilinear term,

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|w_n\|^2 + \frac{v_n}{2} |Aw_n|^2 &\leq -(B(w_n, u), Aw_n) + \|u_n\|_{L^\infty(\Omega)} \|w_n\| |Aw_n| + \frac{1}{2v_n} |Au|^2 \\ &\leq \frac{2k^2}{v_n} |u_n| |Au_n| \|w_n\|^2 + \frac{v_n}{8} |Aw_n|^2 \\ &\quad - (B(w_n, u), Aw_n) + \frac{1}{2v_n} |Au|^2 \end{aligned}$$

and applying (2.8) to the first bilinear term,

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|w_n\|^2 + \frac{3v_n}{8} |Aw_n|^2 &\leq \frac{2k^2}{v_n} |u_n| |Au_n| \|w_n\|^2 \\ &\quad + c |w_n|^{1/2} \|w_n\|^{1/2} \|u\|^{1/2} |Au|^{1/2} |Aw_n| + \frac{1}{2v_n} |Au|^2 \\ &\leq \frac{2k^2}{v_n} |u_n| |Au_n| \|w_n\|^2 + \frac{2c^2}{\lambda_1 v_n} \|w_n\|^2 \|u\| |Au| \\ &\quad + \frac{v_n}{8} |Aw_n|^2 + \frac{1}{2v_n} |Au|^2 \end{aligned}$$

which can be rewritten as

$$\frac{d}{dt} \|w_n\|^2 + \frac{\nu_n}{2} |Aw_n|^2 \leq \left(\frac{4k^2}{\nu_n} |u_n| |Au_n| + \frac{4c^2}{\lambda_1 \nu_n} \|u\| |Au| \right) \|w_n\|^2 + \frac{1}{\nu_n} |Au|^2.$$

Integrating on both sides in time, with $0 \leq t \leq T$,

$$\begin{aligned} \|w_n(t)\|^2 + \frac{\nu_n}{2} \int_0^t |Aw_n|^2 ds &\leq \frac{1}{\nu_n} \int_0^t |Au(s)|^2 ds \\ &+ \int_0^t \left(\frac{4k^2}{\nu_n} |u_n(s)| |Au_n(s)| + \frac{4c^2}{\lambda_1 \nu_n} \|u(s)\| |Au(s)| \right) \|w_n(s)\|^2 ds \end{aligned}$$

Dropping the second term on the left-hand side, we apply Grönwall's inequality to obtain

$$\begin{aligned} \|w_n(t)\|^2 &\leq \alpha_n(t) \exp \left(\int_0^t \frac{4k^2}{\nu_n} |u_n(s)| |Au_n(s)| + \frac{4c^2}{\lambda_1 \nu_n} \|u(s)\| |Au(s)| ds \right) \\ &\leq \alpha(t) \exp \left(\int_0^t \frac{8k^2}{\nu} |u_n(s)| |Au_n(s)| + \frac{8c^2}{\lambda_1 \nu} \|u(s)\| |Au(s)| ds \right). \end{aligned}$$

where $\alpha_n(t) := \frac{2}{\nu_n} \int_0^t |Au(s)|^2 ds \leq \alpha(t) := \frac{4}{\nu} \int_0^t |Au(s)|^2 ds$. Since

$$\int_0^T |Au_n|^2 ds \leq \|u_0\|^2 + \frac{\|f\|_{L^2(0,T;H)}^2}{\nu_n^2}$$

as proved in, e.g., Constantin and Foias (1988), Robinson (2001), Foias et al. (2001), Temam (2001), then

$$\begin{aligned} \sup_{t \in [0,T]} \|w_n(t)\|^2 &\leq \alpha(T) \exp \left(\frac{8k^2}{\lambda_1^2 \nu} \int_0^T |Au_n|^2 ds + \frac{8c^2}{\lambda_1 \nu} \|u(s)\| |Au(s)| ds \right) \\ &\leq \alpha(T) \exp \left(\frac{8k^2}{\lambda_1^2 \nu} \left[\|u_0\|^2 + \frac{\|f\|_{L^2(0,T;H)}^2}{\nu_n^2} \right] \right) \\ &\quad + \alpha(T) \exp \left(\int_0^T \frac{8c^2}{\lambda_1 \nu} \|u(s)\| |Au(s)| ds \right) \\ &\leq \alpha(T) \exp \left(\frac{8k^2}{\lambda_1^2 \nu} \left[\|u_0\|^2 + \frac{4\|f\|_{L^2(0,T;H)}^2}{\nu^2} \right] \right) \\ &\quad + \alpha(T) \exp \left(\int_0^T \frac{8c^2}{\lambda_1 \nu} \|u(s)\| |Au(s)| ds \right) \end{aligned}$$

This implies that $w_n \in L^\infty(0, T; V)$ and $\{w_n\}$ is uniformly bounded in this space.

Additionally, considering again the inequality

$$\begin{aligned} \|w_n(t)\|^2 + \frac{\nu_n}{2} \int_0^t |Aw_n|^2 ds &\leq \frac{1}{\nu_n} \int_0^t |Au(s)|^2 ds \\ &+ \int_0^t \left(\frac{4k^2}{\nu_n} |u_n(s)| |Au_n(s)| + \frac{4c^2}{\lambda_1 \nu_n} \|u(s)\| |Au(s)| \right) \|w_n(s)\|^2 ds, \end{aligned}$$

we set $t = T$, drop the first term on the left-hand side, and bound the above viscosity to obtain

$$\begin{aligned} \int_0^T |Aw_n|^2 ds &\leq \frac{8}{\nu^2} \left(\int_0^T |Au(s)|^2 ds \right) \\ &+ \int_0^T \left(\frac{32k^2}{\lambda_1 \nu^2} |Au_n(s)|^2 + \frac{32c^2}{\lambda_1 \nu^2} \|u(s)\| |Au(s)| \right) \|w_n(s)\|^2 ds. \end{aligned}$$

By the fact that $\{\|u_n\|_{L^2(0, T; \mathcal{D}(A))}\}$ is bounded above in n , as demonstrated in Theorem 3.7, and the result that $\{\|w_n\|_{L^\infty(0, T; V)}\}$ is bounded above uniformly in n , it follows that $\{\|w_n\|_{L^2(0, T; \mathcal{D}(A))}\}$ is bounded above uniformly in n . Since $\{w_n\}$ is bounded above uniformly in n in both $L^\infty(0, T; V)$ and $L^2(0, T; \mathcal{D}(A))$, then we can conclude that there exists a subsequence, which we relabel as $\{w_n\}$, such that

$$w_n \xrightarrow{*} w \text{ in } L^\infty(0, T; V) \text{ and } w_n \rightharpoonup w \text{ in } L^2(0, T; \mathcal{D}(A)). \quad (3.12)$$

Using (3.12), note that all uniform bounds in n on the terms in (3.4) in $L^2(0, T; H)$ are obtained in a similar manner to the proof of strong solutions for the (2.1) and are independent of ν_n except for the bilinear terms. The bilinear terms are bounded uniformly in $L^2(0, T; H)$ with respect to n , due to Lemma 2.1. Hence, $\frac{dw_n}{dt}$ is bounded above uniformly in n in $L^2(0, T; H)$. Thus, by the Aubin Compactness Theorem, $w_n \rightarrow w$ strongly in $L^2(0, T; V)$. Utilizing these bounds and convergence rates, it is classical to show that w is a strong solution in the sense of Definition 3.2. Specifically, as in, e.g., Robinson (2001), Temam (2001), Foias et al. (2001), Constantin and Foias (1988), $w \in C([0, T]; V)$. The initial condition is also satisfied by construction. Uniqueness holds due to the results in Theorem 3.7.

Finally, following a similar argument as in the proof of Theorem 3.7, one can show that the sequence of difference quotients gives rise to a Fréchet derivative $\mathcal{A}(\delta) \equiv \delta \tilde{u}$, where $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$. Let $\mathcal{A}(\delta) = \delta \tilde{u}$. Let $\delta_n = \nu_n - \nu$, so that $\delta_n \rightarrow 0$ as $n \rightarrow \infty$. Note that for any sequence $\nu_n \rightarrow \nu$, there is a subsequence such that $w_{n_k} \rightarrow \tilde{u}$ in $L^2(0, T; V)$. However, by the fact that for every subsequence of $\{w_n\}_{n \in \mathbb{N}}$ we can find a convergence subsequence, one can in fact conclude that the entire sequence $w_n \rightarrow \tilde{u}$ in $L^2(0, T; V)$. We rewrite $w_n \rightarrow \tilde{u}$ in $L^2(0, T; V)$ as

$$\left\| \frac{u(\nu_1 + \delta_n) - u(\nu_1)}{\delta_n} - \tilde{u} \right\|_{L^2(0, T; V)} = \frac{1}{|\delta_n|} \|u(\nu_1 + \delta_n) - u(\nu_1) - \delta_n \tilde{u}\|_{L^2(0, T; V)} \rightarrow 0.$$

□

Remark 3.10 Ideally, one would like to show that the difference quotients of weak solutions give rise to a Fréchet derivative in the sense of the Leray space $L^2(0, T; V) \cap L^\infty(0, T; H)$, rather than just the space $L^2(0, T; H)$. Unfortunately, with the given hypotheses the method of proof employed here does not seem to allow for a proof in this context, since we only have weak or weak-* convergence in the relevant spaces. It may be possible to prove such a result, using, e.g., the methods described in Singler (2008), but such a proof would be lengthy and distract from the main focus of the present work. Hence, we plan to explore these details in a future work.

Finally we consider the sensitivity of the Euler equations to a viscous regularization. This leads to the following statement of the formal sensitivity equations:

$$\tilde{u}_t + \tilde{u} \cdot \nabla u + u \cdot \nabla \tilde{u} - \Delta u + \nabla \tilde{p} = 0, \quad (3.13a)$$

$$\nabla \cdot \tilde{u} = 0, \quad (3.13b)$$

where (u, p) is the solution to the Euler equations. With (u_ν, p_ν) is the solution to (2.1) with viscosity ν , the corresponding system of difference quotient equations is

$$w_t + u_\nu \cdot \nabla w + w \cdot \nabla u - \Delta u + \nabla P = 0, \quad (3.14a)$$

$$\nabla \cdot w = 0, \quad (3.14b)$$

$$w(x, 0) = 0, \quad (3.14c)$$

where $w = \frac{u_\nu - u}{\nu}$ and $P = \frac{p_\nu - p}{\nu}$.

We only utilize the least amount of regularity required on the initial data in Theorem 2.5 to preserve the convergence of strong solutions to the corresponding systems of difference quotient equations in the same spaces as in Theorem 3.9. We state the definition of strong solutions to (3.13) in the same sense as for (3.5) to illustrate the minimal needed additional regularity for the convergence of solutions to (3.14) to those of (3.13).

Definition 3.11 Let $T > 0$, and $s > 2$. Let $u \in C([0, T]; \mathcal{D}(A^{s/2}))$ be a strong solution to (2.21) with initial data $u_0 \in \mathcal{D}(A^{s/2})$ and $f \equiv 0$. A *strong solution* of (3.13) is an element $\tilde{u} \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ satisfying $\frac{d\tilde{u}}{dt} \in L^2(0, T; H)$ and

$$\left\langle \frac{d}{dt} \tilde{u}, \phi \right\rangle + \langle B(\tilde{u}, u), \phi \rangle + \langle B(u, \tilde{u}), \phi \rangle + \langle Au, \phi \rangle = 0 \quad (3.15)$$

for a.e. $t \in [0, T]$, for all $\phi \in H$, and initial data $\tilde{u}_0 \in \mathcal{D}(A^{s/2})$. Moreover, the map $\tilde{u}_0 \in H \mapsto \tilde{u} \in H$ is continuous for a.e. $t \in [0, T]$, i.e., the solution depends continuously on the initial data.

Remark 3.12 We do not extend this analysis to the equations involving the data assimilation algorithm, as the Euler equations are not dissipative.

Definition 3.13 Let $T > 0$, and $s > 2$, and suppose $u, u_v \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$. A strong solution to (3.14) is an element $w \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$, such that $\frac{d}{dt}w \in L^2(0, T; H)$, and for all $\phi \in H$ and for a.e., $t \in [0, T]$,

$$\left(\frac{d}{dt}w, \phi \right) + (B(w, u), \phi) + (B(u_v, w), \phi) + (Au, \phi) = 0.$$

Moreover, the map $u_0 \in H \mapsto u \in H$ is continuous for a.e. $t \in [0, T]$, i.e., the solution depends continuously on the initial data.

Lemma 3.14 Let $T > 0$ be given, and let $u, u_1 \in L^2(0, T; \mathcal{D}(A)) \cap C([0, T]; V)$ be strong solutions to (2.1) and (2.21), respectively. Then there exists a unique strong solution to (3.14).

Proof Since $\Delta u \in L^2(0, T; H)$, and u_v is sufficiently smooth, existence follows exactly as in the standard theory for the 2D Navier–Stokes equations. Uniqueness is also similar, but we carry out the short proof to show the dependencies on various quantities. Hence, suppose w_1 and w_2 are strong solutions to (3.14). Let $U := w_1 - w_2$, so that

$$\frac{d}{dt}U + B(U, u) + B(u_1, U) = 0$$

with $U(0) = 0$ in V . Taking the H inner product with U ,

$$\frac{1}{2} \frac{d}{dt} |U|^2 = -(B(U, u), U) \leq \|\nabla u\|_{L^\infty} |U|^2,$$

and applying Grönwall's inequality yields

$$|U(t)|^2 \leq |U(0)|^2 e^{2 \int_0^t \|\nabla u(s)\|_{L^\infty} ds}. \quad (3.16)$$

Since $u \in C([0, T], \mathcal{D}(A^{s/2}))$, $s > 2$, and $U(0) = 0$, we necessarily have uniqueness. Furthermore, note that (3.16) implies that the map $u_0 \in H \mapsto u \in H$ is continuous. \square

Remark 3.15 In the following theorem, we note that it is worth noting that the analysis can be carried out in a similar manner as when the viscosity is nonzero for the formal sensitivity equations, as one does not encounter the difficulties expected in energy estimates for Euler. The reason for this is that the formal derivative of the viscous regularization results in a diffusive term that is independent of the viscosity.

Theorem 3.16 Let $T > 0$ and $\{v_n\}_{n \in \mathbb{N}}$ be a sequence such that $v_n \rightarrow 0$ as $n \rightarrow \infty$. Let

- u be the solution to (2.21) with forcing $f \equiv 0$ and initial data $u_0 \in \mathcal{D}(A^{s/2})$, $s > 2$,

- u_n solves (2.1) with viscosity ν_n , forcing $f \equiv 0$, and initial data $u_0 \in \mathcal{D}(A^{s/2})$, $s > 2$,
- $\{w_n\}$ be a sequence of strong solutions to (3.14) with corresponding viscosity ν_n , solution to (2.1) u_n , and solution to (2.22) u .

Then $\{w_n\}_{n \in \mathbb{N}}$ that converges in $L^2(0, T; V)$ to a unique strong solution \tilde{u} of (3.13) with initial data $u_0 \equiv 0$. Furthermore, the operator $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$ given by $\mathcal{A}(\delta) = \tilde{u}\delta$ is the Fréchet derivative in $L^2(0, T; V)$ of the weak solution u with respect to ν .

Proof Taking the difference between (2.1) and (2.21), and dividing by the viscosity ν_n , we obtain the system of difference quotient equations, we obtain

$$\frac{d}{dt}w_n + B(w_n, u) + B(u_n, w_n) + Aw_n + Au = 0.$$

Taking the inner product with Aw_n then applying (2.7) and Young's inequality,

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|w_n\|^2 + |Aw_n|^2 &= -(B(w_n, u), Aw_n) - (B(u_n, w_n), Aw_n) - (Au, Aw_n) \\ &\leq c|w_n|^{1/2} \|w_n\|^{1/2} \|u\|^{1/2} |Au|^{1/2} |Aw_n| \\ &\quad + c|u_n|^{1/2} \|u_n\|^{1/2} \|w_n\|^{1/2} |Aw_n|^{1/2} |Aw_n| + |Au| |Aw_n| \\ &\leq \left(\frac{c^2}{\sqrt{\lambda_1}} \|u\| |Au| \right) \|w_n\|^2 + \frac{1}{4} |Aw_n|^2 \\ &\quad + \frac{(32c/3)^4}{4} |u_n|^2 \|u_n\|^2 \|w_n\|^2 + \frac{1}{8} |Aw_n|^2 + 2|Au|^2 \\ &\quad + \frac{1}{8} |Aw_n|^2, \end{aligned}$$

which implies

$$\begin{aligned} \frac{d}{dt} \|w_n\|^2 + |Aw_n|^2 &\leq \left(\frac{2c^2}{\sqrt{\lambda_1}} \|u\| |Au| \right) \|w_n\|^2 \\ &\quad + \frac{(32c/3)^4}{2} |u_n|^2 \|u_n\|^2 \|w_n\|^2 + 4|Au|^2. \end{aligned}$$

Dropping the second term on the left-hand side and applying Grönwall's inequality, for any $0 \leq t \leq T$,

$$\|w_n(t)\|^2 \leq 4 \int_0^t |Au(s)|^2 ds \exp \left(\int_0^t \frac{2c^2}{\sqrt{\lambda_1}} \|u(\tau)\| |Au(\tau)| + \frac{(32c/3)^4}{2} |u_n(\tau)|^2 \|u_n(\tau)\|^2 d\tau \right), \quad (3.17)$$

which is bounded due to the fact that $u \in C([0, T]; \mathcal{D}(A^{s/2}))$ and u_n is bounded uniformly as well in $C([0, T]; \mathcal{D}(A^{s/2}))$. Furthermore, retaining the above second

term and integrating, we obtain

$$\begin{aligned}
& \int_0^t |Aw_n(s)|^2 ds \\
& \leq \int_0^t \left(\frac{2c^2}{\sqrt{\lambda_1}} \|u(s)\| |Au(s)| + \frac{(32c/3)^4}{2} |u_n(s)|^2 \|u_n(s)\|^2 \|w_n(s)\|^2 \right) ds \\
& \quad + 4 \int_0^t |Au(s)|^2 ds \\
& \leq \|w_n(s)\|_{L^\infty(0,T;V)}^2 \int_0^t \left(\frac{2c^2}{\sqrt{\lambda_1}} \|u(s)\| |Au(s)| + \frac{(32c/3)^4}{2} |u_n(s)|^2 \|u_n(s)\|^2 \right) ds \\
& \quad + 4 \int_0^t |Au(s)|^2 ds,
\end{aligned}$$

which is bounded by (3.17). By Theorem 2.5 we have that $u \in L^2(0, T; \mathcal{D}(A))$ and that u_n is bounded uniformly in $C([0, T]; V) \cap L^2(0, T; \mathcal{D}(A))$ with respect to n . The convergence results in Theorem 2.5 can also be applied such that following the exact same arguments as in Theorem 3.9, we obtain the existence of a strong solution \tilde{u} as per Definition 3.13. Uniqueness and continuity with respect to the initial data also follows as in Lemma 3.14. Finally, defining $\mathcal{A}(\delta) \equiv \delta \tilde{u}$, $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$, we also obtain a formal Fréchet derivative in the same manner as in Theorem 3.9. \square

4 Extension to a Data Assimilation Algorithm

In this section, we extend our analysis to the context of a data assimilation algorithm, as discussed in the introduction.

Theorem 4.1 *Let $T > 0$ and $\{v_n\}_{n \in \mathbb{N}}$ be a sequence such that $v_n \rightarrow v_{appx}$, $v_{appx} > 0$, as $n \rightarrow \infty$. Choose μ and h such that $4\mu c_0 h^2 \leq v_n \leq \frac{3v_{appx}}{2}$. Let*

- v be the solution to (2.16) with viscosity v_{appx} , forcing $f \in L^\infty(0, \infty; H)$, and initial data v_0 ;
- \tilde{u} be a strong solution to (3.1);
- v_n solve (2.16) with viscosity v_n , forcing $f \in L^\infty(0, \infty; H)$, and initial data $v_0 \in V$;
- $\{\tilde{w}_n\}_{n \in \mathbb{N}}$ be a sequence of strong solutions to (3.5) with $\tilde{w}_n(0) = 0$ with corresponding viscosity v_n and solutions to (2.16) u_n and u .

Then there exists a subsequence of $\{\tilde{w}_n\}_{n \in \mathbb{N}}$ that converges in $L^2(0, T; V)$ to a unique solution \tilde{v} of (3.2) with initial data $u_0 \equiv 0$. Furthermore, the operator $\mathcal{A} : \mathbb{R} \rightarrow L^2(0, T; V)$ given by $\mathcal{A}(\delta) = \tilde{u}\delta$ is the Fréchet derivative in $L^2(0, T; V)$ of the solution u with respect to v_{appx} .

Proof Let $T > 0$. Note that since $\{v_n\} \subset \left(\frac{v_{appx}}{2}, \frac{3v_{appx}}{2}\right)$ for $n > N$ for some sufficiently large N , we can follow the proof of strong solutions for (2.16) in Azouani et al. (2014) to obtain bounds on $\{v_n\}_{n > N}$ in the appropriate spaces that are independent of

v_n . First, we note that Azouani et al. (2014) quickly proves $|f + \mu P_\sigma I_h(u_n)|^2 \leq M_n$ since $|P_\sigma I_h(u_n)|^2 \leq (c_0\sqrt{h} + \lambda_1^{-1/2})\|u_n\|^2$. However, since u_n is bounded above uniformly in n (see the proof of Theorem 3.7), we have that $|f + \mu P_\sigma I_h(u_n)|^2 \leq m$ for some m independent of n . Thus, we have the following bounds from Azouani et al. (2014) bounded above uniformly in n :

$$\|v_n\|_{L^\infty(0,T;H)}^2 \leq |v_n(0)|^2 + \frac{m}{\mu v_n \lambda_1} \leq |v_0|^2 + \frac{2m}{\mu v_{\text{appx}} \lambda_1}, \quad (4.1)$$

$$\|v_n\|_{L^2(0,T;V)}^2 \leq \frac{1}{v_n} |v_n(0)|^2 + \frac{T}{\mu v_n} m \leq \frac{2}{v_{\text{appx}}} |v_0|^2 + \frac{2T}{\mu v_{\text{appx}}} m, \quad (4.2)$$

and, defining

$$\exp\left\{\frac{c}{(\nu_2)_n^3} \int_0^T |v_n|^2 \|v_n\|^2 ds\right\} \equiv \frac{1}{\psi_{v_n}^n(T)} \leq \frac{1}{\psi_{v_{\text{appx}}}^n(T)} \equiv \exp\left\{\frac{8c}{\nu_{\text{appx}}^3} \int_0^T |v_n|^2 \|v_n\|^2 ds\right\},$$

which is bounded above uniformly in n due to (4.1) and (4.2),

$$\|v_n\|_{L^\infty(0,T;V)}^2 \leq \frac{1}{\psi_{v_n}^n(T)} \left[|v_n(0)|^2 + \frac{4T}{\nu_n} m \right] \leq \frac{1}{\psi_{v_{\text{appx}}}^n(T)} \left[|v_0|^2 + \frac{8T}{\nu_{\text{appx}}} m \right]. \quad (4.3)$$

Furthermore,

$$\begin{aligned} \|v_n\|_{L^2(0,T;\mathcal{D}(A))}^2 &\leq \frac{1}{\nu_n} \|v_n(0)\|^2 + \frac{c}{(\nu_2)_n^3} \int_0^T (|v_n|^2 \|v_n\|^4 + \frac{4}{\nu_n} |f + P_\sigma I_h(u_n)|^2) ds \\ &\leq \frac{2}{\nu_{\text{appx}}} \|v_0\|^2 + \frac{8c}{\nu_{\text{appx}}^3} \int_0^T |v_n|^2 \|v_n\|^4 ds + \frac{8T}{\nu_{\text{appx}}} m, \end{aligned}$$

which is bounded above uniformly in n due to (4.1), (4.2), (4.3). Hence, we will obtain a subsequence that is relabeled $v_n \rightarrow v_1$ in $L^2(0, T; V)$ for some function v_1 . Indeed, we see that by identical arguments presented in Theorem 3.7, $v_1 = v$. Also due to Poincaré's inequality, we obtain that $v_n \rightarrow v$ in $L^2(0, T; H)$.

Let $\{\tilde{w}_n\}_{n \in \mathbb{N}}$ be a sequence of solutions to (3.5). We consider the Leray projection of (3.5):

$$\frac{d}{dt} \tilde{w}_n + B(\tilde{w}_n, v) + B(v_n, \tilde{w}_n) + \nu_n A \tilde{w}_n + Av = \mu P_\sigma I_h(w_n - \tilde{w}_n).$$

The existence proof for (3.2) closely follows the proof of Theorem 3.9, with some modifications on the bounds of \tilde{w}_n which we show below. Taking the inner product with $A \tilde{w}_n$ and proceeding as in the proof of Theorem 3.9, we obtain

$$\frac{1}{2} \frac{d}{dt} \|\tilde{w}_n\|^2 + \frac{\nu_n}{4} |A \tilde{w}_n|^2 \leq \left(\frac{2k^2}{\nu_n} |v_n| |Av_n| + \frac{2c^2}{\lambda_1 \nu_n} \|v\| |Av| \right) \|\tilde{w}_n\|^2$$

$$+ \frac{1}{2\nu_n} |Av|^2 + \mu(I_h(w_n - \tilde{w}_n), A\tilde{w}_n). \quad (4.4)$$

We slightly modify the inequalities obtained in Azouani et al. (2014) for the interpolant term,

$$\begin{aligned} -\mu(I_h(\tilde{w}_n), A\tilde{w}_n) &\leq \frac{4\mu^2}{\nu_n} |\tilde{w}_n - I_h(\tilde{w}_n)|^2 + \frac{\nu_n}{16} |A\tilde{w}_n|^2 - \mu \|\tilde{w}_n\|^2 \\ &\leq \frac{4\mu^2 c_0 h^2}{\nu_n} \|\tilde{w}_n\|^2 + \frac{\nu_n}{16} |A\tilde{w}_n|^2 - \mu \|\tilde{w}_n\|^2 \\ &\leq \frac{\nu_n}{16} |A\tilde{w}_n|^2. \end{aligned}$$

Also,

$$\mu|(I_h(w_n), A\tilde{w}_n)| \leq \frac{4\mu^2}{\nu_n} |w_n|^2 + \frac{\nu_n}{16} |A\tilde{w}_n|^2.$$

Using these inequalities in (4.4):

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\tilde{w}_n\|^2 + \frac{\nu_n}{8} |A\tilde{w}_n|^2 &\leq \left(\frac{2k^2}{\nu_n} |v_n| |Av_n| + \frac{2c^2}{\lambda_1 \nu_n} \|v\| |Av| \right) \|\tilde{w}_n\|^2 \\ &\quad + \frac{1}{2\nu_n} |Av|^2 + \frac{4}{\nu_n} |w_n|^2 \\ &\leq \left(\frac{2k^2}{\nu_n} |v_n| |Av_n| + \frac{2c^2}{\lambda_1 \nu_n} \|v\| |Av| \right) \|\tilde{w}_n\|^2 \\ &\quad + \frac{1}{2\nu_n} |Av|^2 + \frac{4}{\lambda_1^2 \nu_n} |Aw_n|^2. \end{aligned}$$

Following identical arguments as in Theorem 3.9 with

$$\alpha_n(t) := \frac{1}{2\nu_n} |Av_1|^2 + \frac{4}{\lambda_1^2 \nu_n} |Aw_n|^2 \leq \alpha(t) := \frac{1}{\nu_{\text{appx}}} |Av|^2 + \frac{8}{\lambda_1^2 \nu_{\text{appx}}} |Aw_n|^2,$$

along with the fact that $P_\sigma I_h(w_n - \tilde{w}_n)$ is bounded uniformly in n in $L^2(0, T; H)$, we obtain a subsequence relabeled $\tilde{w}_n \rightarrow \tilde{w}$ in $L^2(0, T; V)$. Indeed, let $\phi \in L^2(0, T; H)$; then

$$\begin{aligned} \int_0^T (P_\sigma I_h(w_n - \tilde{w}_n) - P_\sigma I_h(w - \tilde{w}), \phi) ds &\leq \int_0^T |I_h(w_n - \tilde{w}_n) - I_h(w - \tilde{w})| |\phi| ds \\ &\leq \int_0^T |I_h(w_n - w) - I_h(\tilde{w}_n - \tilde{w})| |\phi| ds \\ &\leq \int_0^T |[(w_n - w) - (\tilde{w}_n - \tilde{w})] - I_h((w_n - w) - (\tilde{w}_n - \tilde{w}))| |\phi| ds \end{aligned}$$

$$\begin{aligned}
& + \int_0^T |(w_n - w) - (\tilde{w}_n - \tilde{w})| |\phi| ds \\
& \leq \sqrt{c_0} h \int_0^T \|(w_n - w) - (\tilde{w}_n - \tilde{w})\| |\phi| ds + \frac{1}{\lambda_1^{1/2}} \int_0^T \|(w_n - w) - (\tilde{w}_n - \tilde{w})\| |\phi| ds \\
& \leq \sqrt{c_0} h \|w_n - w\|_{L^2(0, T; V)} \|\phi\|_{L^2(0, T; H)} + \frac{1}{\lambda_1^{1/2}} \|\tilde{w}_n - \tilde{w}\|_{L^2(0, T; V)} \|\phi\|_{L^2(0, T; H)}.
\end{aligned}$$

Additionally, since we now have that $\tilde{w}_n \rightarrow w$ in $L^2(0, T; V)$, then $P_\sigma I_h(w_n - \tilde{w}_n) \rightarrow P_\sigma I_h(w - \tilde{w})$ in $L^2(0, T; H)$ and we conclude \tilde{w} is a strong solution in the sense of Definition 3.3.

To show that the solutions are unique, we consider the difference of the equations

$$\frac{d}{dt} \tilde{v}_1 + B(\tilde{v}_1, v) + B(v, \tilde{v}_1) + \nu_{\text{appx}} A \tilde{v}_1 + A v = \mu P_\sigma I_h(\tilde{u} - \tilde{v}_1)$$

and

$$\frac{d}{dt} \tilde{v}_2 + B(\tilde{v}_2, v) + B(v, \tilde{v}_2) + \nu_{\text{appx}} A \tilde{v}_2 + A v = \mu P_\sigma I_h(\tilde{u} - \tilde{v}_2)$$

which, defining $W := \tilde{v}_1 - \tilde{v}_2$, yields

$$\frac{d}{dt} W + B(W, v) + B(v, W) + \nu_{\text{appx}} A W = -\mu P_\sigma I_h(W)$$

with $W(0) = 0$. So, W must be a solution to the above equation. Taking the action on W and applying the Lions–Magenes Lemma 2.2,

$$\frac{1}{2} \frac{d}{dt} |W|^2 + \langle B(W, v), W \rangle + \nu_{\text{appx}} \|W\|^2 = \langle -\mu P_\sigma I_h(W), W \rangle$$

which implies that

$$\begin{aligned}
\frac{1}{2} \frac{d}{dt} |W|^2 + \nu_{\text{appx}} \|W\|^2 & \leq c \|v\| |W| \|W\| + \mu (\sqrt{c_0} h + \lambda_1^{-1}) \|W\| |W| \\
& \leq \frac{\mu^2 (\sqrt{c_0} h + \lambda_1^{-1})^2}{\nu_{\text{appx}}} |W|^2 + \frac{\nu}{4} \|W\|^2 \\
& \quad + \frac{c^2}{2 \nu_{\text{appx}}} \|v\|^2 |W|^2 + \frac{\nu_{\text{appx}}}{2} \|W\|^2.
\end{aligned}$$

Thus,

$$\frac{d}{dt} |W|^2 \leq \left(\frac{\mu^2 (\sqrt{c_0} h + \lambda_1^{-1})^2}{\nu_{\text{appx}}} + \frac{c^2}{2 \nu_{\text{appx}}} \|v\|^2 \right) |W|^2$$

and Grönwall's inequality implies, for a.e. $0 \leq t \leq T$,

$$|W(t)|^2 \leq |W(0)|^2 \exp\left(\int_0^T \frac{\mu^2(\sqrt{c_0}h + \lambda_1^{-1})^2}{\nu_{\text{appx}}} + \frac{c^2}{2\nu_{\text{appx}}} \|v\|^2 dt\right).$$

But $W(0) = 0$, and thus $\|W\|_{L^\infty(0,T;H)} = 0$ implies that $W \equiv 0$. Hence, solutions to (3.5) are unique.

Finally, we want to show that the sequence of difference quotients defines a Fréchet derivative. Following a similar argument as in Theorem 3.7, the Fréchet derivative $\mathcal{A}(\delta) \equiv \delta \tilde{u}$ maps $\mathbb{R} \rightarrow L^2(0, T; V)$. \square

5 Conclusion

In this article, we proved well-posedness of the sensitivity equations for the 2D incompressible Navier–Stokes equations and the associated AOT data assimilation system, as well as a viscous regularization of Euler. Specifically, we proved the existence and uniqueness of global solutions to these equations. A byproduct of the proof is that the sensitivity of solutions to the equations involved in the algorithm are bounded in appropriate spaces. Hence, changing the viscosity, or equivalently the Reynolds number, mid-simulation as in Carlson et al. (2020) does not result in major aberrations in the solution. We note that in the present context, our proof is somewhat nonstandard, in that we proved the existence by showing that the difference quotients converge to a solution of the equations. We believe this is the first such rigorous proof that the sensitivity equations for the 2D Navier–Stokes equations are globally well posed (even in the inviscid limit), although formal proofs have been given in other works, cited above.

Acknowledgements The authors would like to thank the anonymous referees for helpful comments and suggestions that significantly improved the quality of this manuscript. E.C. would like to give thanks for the kind hospitality of the COSIM group at Los Alamos National Laboratory where some of this work was completed. The research of E.C. was supported in part by the NSF GRFP grant no. 1610400. The research of A.L. was supported in part by the NSF grants no. DMS-1716801 and CMMI-1953346.

References

- Albanez, D.A., Nussenzveig Lopes, H.J., Titi, E.S.: Continuous data assimilation for the three-dimensional Navier–Stokes- α model. *Asymptotic Anal.* **97**(1–2), 139–164 (2016)
- Altaf, M.U., Titi, E.S., Knio, O.M., Zhao, L., McCabe, M.F., Hoteit, I.: Downscaling the 2D Benard convection equations using continuous data assimilation. *Comput. Geosci.* **21**(3), 393–410 (2017)
- Anderson, K., Newman, J.C., Whitfield, D.L., Nielsen, E.J.: Sensitivity analysis for Navier–Stokes equations on unstructured meshes using complex variables. *AIAA J.* **39**, 11 (1999)
- Azouani, A., Titi, E.S.: Feedback control of nonlinear dissipative systems by finite determining parameters—a reaction-diffusion paradigm. *Evol. Equ. Control Theory* **3**(4), 579–594 (2014)
- Azouani, A., Olson, E., Titi, E.S.: Continuous data assimilation using general interpolant observables. *J. Nonlinear Sci.* **24**(2), 277–304 (2014)
- Bessaih, H., Olson, E., Titi, E.S.: Continuous data assimilation with stochastically noisy data. *Nonlinearity* **28**(3), 729–753 (2015)

Biswas, A., Martinez, V.R.: Higher-order synchronization for a data assimilation algorithm for the 2D Navier–Stokes equations. *Nonlinear Anal. Real World Appl.* **35**, 132–157 (2017)

Biswas, A., Hudson, J., Larios, A., Pei, Y.: Continuous data assimilation for the 2D magnetohydrodynamic equations using one component of the velocity and magnetic fields. *Asymptot. Anal.* **108**(1–2), 1–43 (2018)

Biswas, A., Foias, C., Mondaini, C.F., Titi, E.S.: Downscaling data assimilation algorithm with applications to statistical solutions of the Navier–Stokes equations. In: *Annales de l’Institut Henri Poincaré C, Analyse non linéaire*, pp. 295–326. Elsevier (2019)

Blömker, D., Law, K., Stuart, A.M., Zygalakis, K.C.: Accuracy and stability of the continuous-time 3DVAR filter for the Navier–Stokes equation. *Nonlinearity* **26**(8), 2193–2219 (2013)

Borggaard, J., Burns, J.: A PDE sensitivity equation method for optimal aerodynamic design. *J. Comput. Phys.* **136**(2), 366–384 (1997)

Breckling, S., Neda, M., Pahlevani, F.: A sensitivity study of the Navier–Stokes- α model. *Comput. Math. Appl.* **75**(2), 666–689 (2018)

Brewer, D.: The differentiability with respect to a parameter of the solution of a linear abstract Cauchy problem. *SIAM J. Math. Anal.* **13**(4), 607–620 (1982)

Carlson, E., Larios, A.: Super-exponential convergence of certain nonlinear algorithms for continuous data assimilation. (2020). (preprint)

Carlson, E., Hudson, J., Larios, A.: Parameter recovery for the 2 dimensional Navier–Stokes equations via continuous data assimilation. *SIAM J. Sci. Comput.* **42**(1), A250–A270 (2020)

Celik, E., Olson, E., Titi, E.S.: Spectral filtering of interpolant observables for a discrete-in-time downscaling data assimilation algorithm. *SIAM J. Appl. Dyn. Syst.* **18**(2), 1118–1142 (2019)

Constantin, P., Foias, C.: Global Lyapunov exponents, Kaplan–Yorke formulas and the dimension of the attractors for 2D Navier–Stokes equations. *Commun. Pure Appl. Math.* **38**(1), 1–27 (1985)

Constantin, P., Foias, C.: Navier–Stokes Equations. Chicago Lectures in Mathematics. University of Chicago Press, Chicago, IL (1988)

Davis, L., Pahlevani, F.: Parameter sensitivity of an eddy viscosity model: analysis, computation and its application to quantifying model reliability. *Int. J. Uncertain. Quantif.* **3**(5), 397–419 (2013)

Desamsetti, S., Dasari, H.P., Langodan, S., Knio, O., Hoteit, I., Titi, E.S.: Efficient dynamical downscaling of general circulation models using continuous data assimilation. *Q. J. R. Meteorol. Soc.* **145**, 3175–3194 (2019). <https://doi.org/10.1002/qj.3612>. <http://hdl.handle.net/10754/656325>

Di Leoni, P.C., Mazzino, A., Biferale, L.: Inferring flow parameters and turbulent configuration with physics-informed data assimilation and spectral nudging. *Phys. Rev. Fluids* **3**(10), 104604 (2018)

Farhat, A., Jolly, M.S., Titi, E.S.: Continuous data assimilation for the 2D Bénard convection through velocity measurements alone. *Phys. D* **303**, 59–66 (2015)

Farhat, A., Lunasin, E., Titi, E.S.: Abridged continuous data assimilation for the 2D Navier–Stokes equations utilizing measurements of only one component of the velocity field. *J. Math. Fluid Mech.* **18**(1), 1–23 (2016)

Farhat, A., Lunasin, E., Titi, E.S.: Data assimilation algorithm for 3D Bénard convection in porous media employing only temperature measurements. *J. Math. Anal. Appl.* **438**(1), 492–506 (2016)

Farhat, A., Lunasin, E., Titi, E.S.: On the Charney conjecture of data assimilation employing temperature measurements alone: the paradigm of 3D planetary geostrophic model. *Math. Clim. Weather Forecast.* **2**(1), 61–74 (2016). <https://doi.org/10.1515/mcwf-2016-0004>

Farhat, A., Lunasin, E., Titi, E.S.: Continuous data assimilation for a 2D Bénard convection system through horizontal velocity measurements alone. *J. Nonlinear Sci.*, 1–23 (2017)

Farhat, A., Johnston, H., Jolly, M., Titi, E.S.: Assimilation of nearly turbulent Rayleigh–Bénard flow through vorticity or local circulation measurements: a computational study. *J. Sci. Comput.* **77**(3), 1519–1533 (2018)

Farhat, A., Lunasin, E., Titi, E.S.: A data assimilation algorithm: the paradigm of the 3D Leray- α model of turbulence. In: *Partial Differential Equations Arising from Physics and Geometry*, vol. 450, pp. 253–273 (2019)

Farhat, A., Glatt-Holtz, N.E., Martinez, V.R., McQuarrie, S.A., Whitehead, J.P.: Data assimilation in large Prandtl Rayleigh–Bénard convection from thermal measurements. *SIAM J. Appl. Dyn. Syst.* **19**(1), 510–540 (2020)

Fernández, M.A., Moubachir, M.: Sensitivity analysis for an incompressible aeroelastic system. *Math. Models Methods Appl. Sci.* **12**, 1109–1130 (2002)

Foias, C., Manley, O., Rosa, R., Temam, R.: Navier–Stokes Equations and Turbulence. Encyclopedia of Mathematics and its Applications, vol. 83. Cambridge University Press, Cambridge (2001)

Foias, C., Mondaini, C.F., Titi, E.S.: A discrete data assimilation scheme for the solutions of the two-dimensional Navier–Stokes equations and their statistics. *SIAM J. Appl. Dyn. Syst.* **15**(4), 2109–2142 (2016)

Foyash, K., Dzholl, M.S., Kravchenko, R., Titi, È.S.: A unified approach to the construction of defining forms for a two-dimensional system of Navier–Stokes equations: the case of general interpolating operators. *Uspekhi Mat. Nauk* **69**(2(416)), 177–200 (2014)

García-Archilla, B., Novo, J., Titi, E.S.: Uniform in time error estimates for a finite element method applied to a downscaling data assimilation algorithm for the Navier–Stokes equations. *SIAM J. Numer. Anal.* **58**(1), 410–429 (2020)

Gardner, M., Larios, A., Rebholz, L.G., Vargun, D., Zerfas, C.: Continuous data assimilation applied to a velocity–vorticity formulation of the 2d Navier–Stokes equations. *Electron. Res. Arch.* **29**(3), 2223–2247 (2021)

Gesho, M., Olson, E., Titi, E.S.: A computational study of a data assimilation algorithm for the two-dimensional Navier–Stokes equations. *Commun. Comput. Phys.* **19**(4), 1094–1110 (2016)

Gibson, J., Clark, L.: Sensitivity analysis for a class of evolution equations. *J. Comput. Phys.* **136**(2), 366–384 (1997)

Glatt-Holtz, N., Kukavica, I., Vicol, V., Ziane, M.: Existence and regularity of invariant measures for the three dimensional stochastic primitive equations. *J. Math. Phys.* **55**(5), 051504, 34 (2014)

Golovkin, K.K.: Vanishing viscosity in Cauchy's problem for hydrodynamics equations. *Proc. Steklov Inst. Math.* **92**, 33–53 (1966)

Grappin, R., Léorat, J.: Lyapunov exponents and the dimension of periodic incompressible Navier–Stokes flows: numerical measurements. *J. Fluid Mech.* **222**, 61–94 (1991)

Hamby, D.: A review of techniques for parameter sensitivity analysis of environmental models. *Environ. Monit. Assess.* **32**(2), 135–154 (1994)

Hudson, J., Jolly, M.: Numerical efficacy study of data assimilation for the 2D magnetohydrodynamic equations. *J. Comput. Dyn.* **6**(1), 131–145 (2019)

Ibdah, H.A., Mondaini, C.F., Titi, E.S.: Fully discrete numerical schemes of a data assimilation algorithm: uniform-in-time error estimates. *IMA J. Numer. Anal.* **11** 2019. drz043

Jolly, M.S., Martinez, V.R., Titi, E.S.: A data assimilation algorithm for the subcritical surface quasi-geostrophic equation. *Adv. Nonlinear Stud.* **17**(1), 167–192 (2017)

Jolly, M.S., Martinez, V.R., Olson, E.J., Titi, E.S.: Continuous data assimilation with blurred-in-time measurements of the surface quasi-geostrophic equation. *Chin. Ann. Math. Ser. B* **40**(5), 721–764 (2019)

Kato, T.: Nonstationary flows of viscous and ideal fluids in R^3 . *J. Funct. Anal.* **9**, 296–305 (1972)

Kim, H., Kim, C., Rho, O.-H., Dong Lee, K.: Aerodynamic sensitivity analysis for Navier–Stokes equations. *J. KSIAM* **3**, 161–171 (1999)

Kouhi, M., Houzeaux, G., Cucchietti, F., Vázquez, M.: Implementation of discrete adjoint method for parameter sensitivity analysis in chemically reacting flows. In: 57th AIAA/ASCE/AHS/ASC Structures, Structural Dynamics, and Materials Conference, AIAA SciTech Forum (2016)

Larios, A., Pei, Y.: Nonlinear continuous data assimilation. (2017). [arXiv:1703.03546](https://arxiv.org/abs/1703.03546)

Larios, A., Victor, C.: Continuous data assimilation with a moving cluster of data points for a reaction diffusion equation: a computational study. *Commun. Comput. Phys.* **29**, 1273–1298 (2021)

Larios, A., Rebholz, L.G., Zerfas, C.: Global in time stability and accuracy of IMEX-FEM data assimilation schemes for Navier–Stokes equations. *Comput. Methods Appl. Mech. Eng.* **345**, 1077–1093 (2018). <https://doi.org/10.1016/j.cma.2018.09.004>. <http://www.sciencedirect.com/science/article/pii/S004578251830447X>

Leoni, D., Clark, P., Mazzino, A., Biferale, L.: Synchronization to big-data: nudging the Navier–Stokes equations for data assimilation of turbulent flows. *arXiv preprint arXiv:1905.05860*, (2019)

Lunasin, E., Titi, E.S.: Finite determining parameters feedback control for distributed nonlinear dissipative systems—a computational study. *Evol. Equ. Control Theory* **6**(4), 535–557 (2017)

Markowich, P.A., Titi, E.S., Trabelsi, S.: Continuous data assimilation for the three-dimensional Brinkman–Forchheimer-extended Darcy model. *Nonlinearity* **29**(4), 1292–1328 (2016)

Masmoudi, N.: Remarks about the inviscid limit of the Navier–Stokes system. *Commun. Math. Phys.* **270**, 777–788 (2007)

Mondaini, C.F., Titi, E.S.: Uniform-in-time error estimates for the postprocessing Galerkin method applied to a data assimilation algorithm. *SIAM J. Numer. Anal.* **56**(1), 78–110 (2018)

Neda, M., Pahlevani, F., Rebholz, L.G., Waters, J.: Sensitivity analysis of the grad-div stabilization parameter in finite element simulations of incompressible flow. *J. Numer. Math.* **24**(3), 189–206 (2016)

Noacco, V., Sarrazin, F., Pianosi, F., Wagener, T.: Matlab/r workflows to assess critical choices in global sensitivity analysis using the safe toolbox. *MethodsX* **6**, 2258–2280 (2019)

Pahlevani, F.: Sensitivity analysis of eddy viscosity models. ProQuest LLC, Ann Arbor, MI, 2004. Thesis (Ph.D.)—University of Pittsburgh

Pahlevani, F.: Sensitivity computations of eddy viscosity models with an application in drag computation. *Int. J. Numer. Methods Fluids* **52**(4), 381–392 (2006)

Pei, Y.: Continuous data assimilation for the 3D primitive equations of the ocean. *Commun. Pure Appl. Math.* **18**(2), 643 (2019)

Ponce, G.: On two-dimensional incompressible fluids. *Comm. Part. Diff. Eqs.* **11**, 483–511 (1986)

Rebholz, L.G., Zerfas, C.: Simple and efficient continuous data assimilation of evolution equations via algebraic nudging. *Numer. Methods Partial Differ. Equ.* **37**, 1–25 (2021)

Rebholz, L., Zerfas, C., Zhao, K.: Global in time analysis and sensitivity analysis for the reduced NS- α model of incompressible flow. *J. Math. Fluid Mech.* **19**(3), 445–467 (2017)

Robinson, J.C.: Infinite-Dimensional Dynamical Systems. Cambridge Texts in Applied Mathematics. Cambridge University Press, Cambridge, 2001. An Introduction to Dissipative Parabolic PDEs and the Theory of Global Attractors

Singler, J.: Differentiability with respect to parameters of weak solutions of linear parabolic equations. *Math. Comput. Model.* **47**(3), 422–430 (2008)

Stanley, L.G., Stewart, D.L.: Design Sensitivity Analysis: Computational Issues of Sensitivity Equation Methods, volume 25 of Frontiers in Applied Mathematics. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 2002. Computational issues of sensitivity equation methods

Swann, H.G.: The convergence with vanishing viscosity of nonstationary Navier–Stokes flow to ideal flow in \mathbb{R}^3 . *Trans. Am. Math. Soc.* **157**, 373–397 (1971)

Temam, R.: Navier–Stokes Equations and Nonlinear Functional Analysis, volume 66 of CBMS-NSF Regional Conference Series in Applied Mathematics. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, second edition (1995)

Temam, R.: Navier–Stokes Equations: Theory and Numerical Analysis. AMS Chelsea Publishing, Providence, RI, 2001. Theory and numerical analysis, Reprint of the 1984 edition

Vemuri, V., Raefsky, A.: On a new approach to parameter estimation by the method of sensitivity functions. *Int. J. Syst. Sci.* **10**(4), 395–407 (1979)

Zerfas, C., Rebholz, L.G., Schneier, M., Iliescu, T.: Continuous data assimilation reduced order models of fluid flow. *Comput. Methods Appl. Mech. Eng.* **357**, 112596, 18 (2019)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.