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## RESEARCH ARTICLE

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### Special Section:

Fifteen Years of THEMIS  
Mission

### Key Points:

- We report on intense sub-relativistic electron precipitation events observed by ELFIN CubeSats
- Sub-relativistic precipitation events are due to electron nonlinear Landau resonance with very oblique whistler waves
- Sub-relativistic precipitation events may indicate an enhancement of equatorial energetic electron fluxes

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## On the Nature of Intense Sub-Relativistic Electron Precipitation

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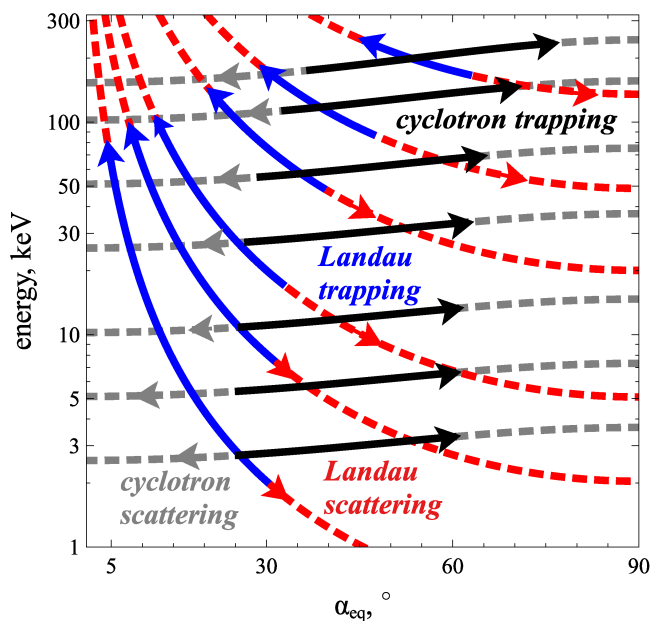
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**Abstract** Energetic electron precipitation into Earth's atmosphere is an important process for radiation belt dynamics and magnetosphere-ionosphere coupling. The most intense form of such precipitation is microbursts—short-lived bursts of precipitating fluxes detected on low-altitude spacecraft. Due to the wide energy range of microbursts (from sub-relativistic to relativistic energies) and their transient nature, they are thought to be predominantly associated with energetic electron scattering into the loss cone via cyclotron resonance with field-aligned intense whistler-mode chorus waves. In this study, we show that intense sub-relativistic microbursts may be generated via electron nonlinear Landau resonance with very oblique whistler-mode waves. We combine a theoretical model of nonlinear Landau resonance, equatorial observations of intense very oblique whistler-mode waves, and conjugate low-altitude observations of <200 keV electron precipitation. Based on model comparison with observed precipitation, we suggest that such sub-relativistic microbursts occur by plasma sheet (0.1 – 10 keV) electron trapping in nonlinear Landau resonance, resulting in acceleration to  $\lesssim 200$  keV energies and simultaneous transport into the loss cone. The proposed scenario of intense sub-relativistic ( $\lesssim 200$  keV) microbursts demonstrates the importance of very oblique whistler-mode waves for radiation belt dynamics.

## 1. Introduction

Energetic (several keV to several MeV) electron precipitation into the Earth's atmosphere is one of the most important mechanisms of electron loss from the outer radiation belt (see, e.g., reviews by Millan & Thorne, 2007; Shprits et al., 2008, and references therein). Besides the night-side curvature scattering (Artemyev, Orlova, et al., 2013; Capannolo et al., 2022; Sergeev & Tsyganenko, 1982; Yahnin et al., 1997), wave-particle resonant interaction is the main driver of such precipitation. Electrons of different energies interact resonantly with different wave modes, mainly: whistler-mode chorus waves (responsible for precipitation of sub-MeV electrons from the outer radiation belt, see, e.g., Horne and Thorne (1998); Horne and Thorne (2003); Ma, Li, Thorne, Nishimura, et al. (2016)), whistler-mode hiss waves (responsible for precipitation of sub-MeV electrons from the plasmasphere, see, e.g., Ma, Li, Thorne, Bortnik, et al. (2016); Mourenas et al. (2017)), and electromagnetic ion cyclotron (EMIC) waves (responsible for relativistic and ultra-relativistic electron precipitation from the outer radiation belts, see, e.g., Kersten et al. (2014); Blum, Halford, et al. (2015); Kubota et al. (2015); Grach and Demekhov (2020)). The most intense whistler-mode and EMIC waves interact with electrons nonlinearly (see reviews by Shklyar & Matsumoto, 2009; Albert et al., 2013; Artemyev, Neishtadt, Vainchtein, et al., 2018, and references therein) providing very intense and bursty precipitation events, called microbursts (Blum, Li, & Denton, 2015; Breneman et al., 2017; Capannolo et al., 2019; O'Brien et al., 2004; Shumko, Turner, et al., 2018; Zhang, Angelopoulos, et al., 2022). There exist many theoretical models of microbursts generated by whistler-mode waves (e.g., L. Chen et al., 2020; L. Chen et al., 2021) and EMIC waves (e.g., Kubota & Omura, 2017). Such models consider mostly field-aligned wave propagation, because the most intense waves observed are field-aligned chorus (Agapitov et al., 2013; Li, Santolík, et al., 2016; Santolík, Macúšová, et al., 2014) and EMIC (e.g., Min & Liu, 2016; Mithaiwala et al., 2013; Yu et al., 2016) waves.

Electron precipitation due to wave-particle resonant interactions corresponds to a decrease of the electron equatorial pitch-angle  $\alpha_{eq}$  and a change of electron energy. Field-aligned chorus waves interact with electrons through the first cyclotron resonance. The nonlinear regime of this resonant interaction includes phase trapping (with  $\alpha_{eq}$  and energy increase, see, e.g., Omura et al. (2015)) and phase bunching (with  $\alpha_{eq}$  and energy decrease, see,



**Figure 1.** Resonance curves (Williams & Lyons, 1974) for cyclotron resonance with field-aligned whistler-mode wave and Landau resonance with very oblique whistler-mode wave. Resonance with monochromatic waves moves electrons along the shown curves. Thus, electron transport to the loss cone (toward smaller pitch-angles) almost does not change electron energy for cyclotron resonance, but significantly increases energy for Landau resonance. Nonlinear trapping direction is shown by solid lines, whereas nonlinear scattering (phase bunching) direction is shown by dashed lines. System parameters:  $L$ -shell is 6, wave frequency to equatorial cyclotron frequency is  $f/f_{ce} = 0.35$ , wave normal angle of very oblique whistler-mode wave is  $\theta = \arccos(f/f_{ce}) - 5^\circ$ .

e.g., Albert (2000); Vainchtein et al. (2018)). The competition between phase trapping and phase bunching (Artemyev, Neishtadt, et al., 2016; Istomin et al., 1973; Shklyar, 2011; Solov'ev & Shklyar, 1986) controls the net electron transport into the loss cone, but all precipitating electrons lose energy, because precipitation is provided only by phase bunching. Thus, both linear and non-linear electron interaction with chorus via the cyclotron resonance is associated with electron deceleration. Thus, not only is net energy lost from the magnetosphere due to precipitation, but this energy is lost by equatorial electrons of equal or higher energy, resulting in a decay of the radiation belt (energetic electron) flux. Precipitation due to electron resonance with EMIC waves is due to phase trapping, which moves electrons into the loss cone (Albert & Bortnik, 2009; Grach & Demekhov, 2020; Grach et al., 2021; Kubota & Omura, 2017). Such trapping, just like electron resonance with EMIC waves in general (Ni et al., 2015; Summers & Thorne, 2003), does not change the electron energy, that is, EMIC driven precipitation events are associated with energetic electron flux decrease simply due to atmospheric loss without deceleration of equatorial electrons. As microbursts are very intense (associated with orders of magnitude increases of precipitating flux at low-altitude spacecraft measurements), this precipitation pattern is often considered as one of the main mechanisms of energetic electron loss (e.g., Thorne et al., 2005).

The energy range of microburst precipitation is quite wide (e.g., Zhang, Angelopoulos, et al., 2022): there are both relativistic microbursts (Blum, Li, & Denton, 2015; O'Brien et al., 2004), microbursts of few hundreds of keV (Capannolo et al., 2019; Shumko, Sample, et al., 2018), and even of tens of keV (Blake & O'Brien, 2016). Moreover, recent investigations show strong correlations between relativistic microbursts and tens of keV electron precipitation responsible for diffuse aurora (Y. Miyoshi et al., 2020; K. Miyoshi Y. et al., 2021; Shumko et al., 2021). Although relativistic microbursts may be generated by electron resonances with EMIC waves (Blum, Halford, et al., 2015; Grach & Demekhov, 2020; Kubota et al., 2015), sub-relativistic electrons can rarely interact resonantly with EMIC waves (e.g., L. Chen et al., 2019) and hence they are often attributed to electron scattering by chorus waves (Breneman et al., 2017; Mozer et al., 2018; Y. Miyoshi et al., 2020).

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Although field-aligned chorus waves are the most intense whistler-mode wave population in the radiation belt (Agapitov et al., 2013; Li et al., 2011), this is not the sole wave population present in this region. There is also a significant population of very oblique whistler-mode waves (almost electrostatic mode, see Artemyev, Agapitov, et al. (2016)) propagating near the resonance cone angle (Agapitov et al., 2013; Agapitov et al., 2018; Li, Santolík, et al., 2016). These oblique waves can interact resonantly with electrons through Landau resonance (e.g., Nunn & Omura, 2015; Shklyar & Matsumoto, 2009) and can accelerate electrons efficiently via phase trapping (Agapitov et al., 2014; Artemyev, Krasnoselskikh, et al., 2012; Hsieh et al., 2020; Hsieh & Omura, 2017a). The corresponding Landau resonant energies are much lower than cyclotron resonant energies with quasi-parallel waves, allowing very oblique waves to resonate with 100 eV - 10 keV electrons (Artemyev, Agapitov, et al., 2015, 2016; Mourenas et al., 2014) which can almost never be scattered by field-aligned lower-band chorus waves (Li et al., 2010). Moreover, in contrast to the cyclotron resonance associated with an energy decrease of precipitating electrons, the Landau resonant trapping is associated with energy increase (see Figure 1 for resonance curves and directions of electron motion in energy, pitch-angle space due to cyclotron and Landau resonances). In other words, precipitating electrons (those with decreasing  $\alpha_{eq}$ ) are simultaneously accelerated, corresponding to an increase of low- $\alpha_{eq}$  energetic electron flux (Agapitov et al., 2015b). Since very oblique whistler-mode waves can be quite intense (Agapitov et al., 2014; Cattell et al., 2008; Cully, Bonnell, & Ergun, 2008; Wilson et al., 2011), these waves may in principle interact nonlinearly with electrons and efficiently trap them through Landau resonance (Agapitov et al., 2015b; Mourenas et al., 2016). But can such trapping result in intense (microburst-like) electron precipitation?

This study aims to address this question by combining a theoretical investigation of electron phase trapping into Landau resonance, using observed wave characteristics, and low-altitude observations of electron precipitation. We show that the phase trapping into Landau resonance (Landau-trapping) can indeed create an intense electron flux peaking within the loss cone. However, in contrast to precipitation due to cyclotron resonance, precipitation due to Landau-trapping does not correspond to energetic electron flux decay. Instead, Landau-trapping precipitation is accompanied by significant electron acceleration and the formation of an additional sub-relativistic (up to  $\leq 200$  keV) electron population in the outer radiation belt.

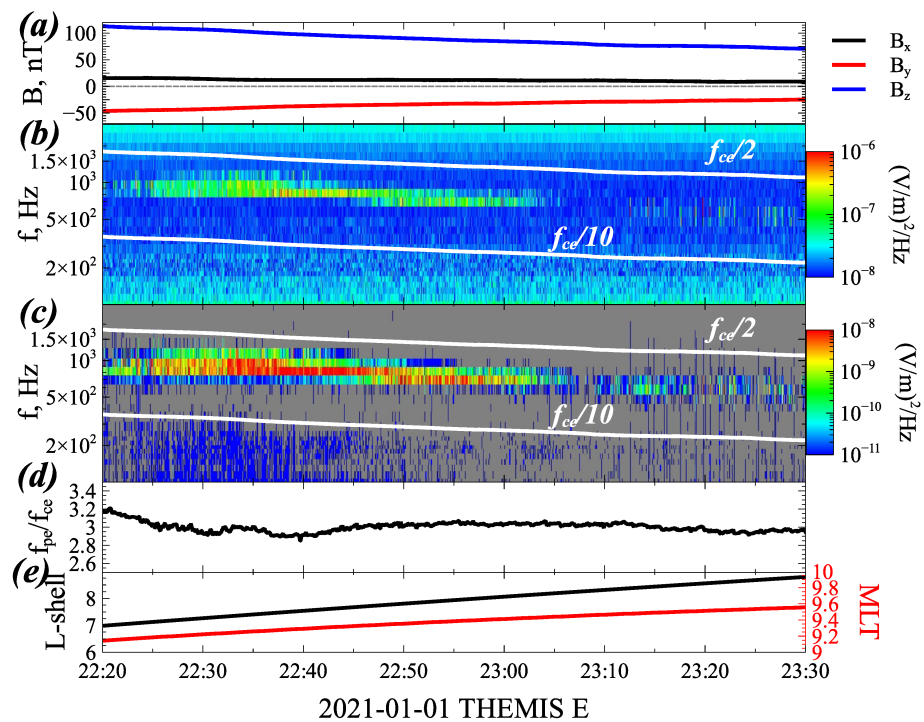
The paper consists of four sections. In Section 2, we analyze an event with Time History of Events and Macroscale Interactions during Substorms (THEMIS) spacecraft (Angelopoulos, 2008) equatorial measurements of very oblique whistler-mode waves, together with conjugate low-altitude ELFING CubeSat (Angelopoulos et al., 2020) measurements of electron precipitation. We supplement the analysis of this event by Defense Meteorological Satellite Program (DMSP) satellite measurements of  $< 10$  keV electron precipitation (Hardy et al., 1984; Rich et al., 1985), because electrons of such energies can only interact with whistler-mode lower-band chorus waves through Landau resonance. In Section 3, we describe the theoretical model of electron acceleration and precipitation by Landau-trapping. This model is based on the mapping technique for the evaluation of electron flux dynamics due to nonlinear interactions with whistler-mode waves (Artemyev et al., 2020b; Artemyev, Neishtadt, Vasiliev, Zhang, et al., 2021). We generalize the mapping technique for systems with a realistic distribution of waves (in terms of as frequency and time-series envelope amplitude), and apply it to the data set obtained from THEMIS around the equatorial plane. Our mapping simulation results of electron flux expected at ionospheric altitude are compared with ELFING observations of electron precipitation events. In Section 4, we discuss the results of this comparison and speculate about the importance of Landau resonance with very oblique whistler-mode waves for intense sub-relativistic electron precipitation.

## 2. THEMIS, ELFING, and DMSP Observations on 2021-01-01

We first analyze a precipitation event with conjugate THEMIS and ELFING observations on the dayside in the inner magnetosphere. Figure 2 is an overview of THEMIS E measurements on 2021-01-01 from 22:20 to 23:30 UT in the dayside inner magnetosphere, when the spacecraft moved outward from  $L \sim 7$  to  $L \sim 9$  ( $L$  is evaluated with the Tsyganenko (1989) magnetic field model). Fluxgate magnetometer data with 1/4s resolution (Auster et al., 2008) shows that THEMIS E was around the equator with GSM  $B_z \gg |B_x|, |B_y|$  (see panel (a)). THEMIS search coil (Le Contel et al., 2008) and electric field instrument (Bonnell et al., 2008) measurements are used to check the wave activity. Panels (b) and (c) show magnetic and electric field spectra obtained from these instruments (32 frequency channels in 10 – 4000 Hz range) in the *fff* data set (see Cully, Ergun, et al., 2008). There is strong whistler-mode wave activity within the frequency range  $ff_{ce} \in [0.2, 0.4]$ , where the electron cyclotron frequency  $f_{ce}$  is calculated from in-situ measurements from the fluxgate magnetometer. Using plasma density estimates from the spacecraft potential (Nishimura et al., 2013), we plot the plasma frequency to electron gyrofrequency ratio,  $f_{pe}/f_{ce}$ , in panel (d). This ratio is quite low ( $f_{pe}/f_{ce} \sim 3$ ) for  $L \in [6, 8]$ , suggesting atypical conditions of whistler-mode wave generation and wave resonant interaction with electrons.

Figure 3 shows several whistler-mode wave packets. We use electric field measurements in the field-aligned coordinate system, because 3D wave magnetic field measurements are not yet fully calibrated in the whistler-mode range during this interval. All wave packets show a large parallel electric field component, suggesting the presence of very oblique whistler-mode waves (Artemyev, Agapitov, et al., 2016). Using  $f_{pe}/f_{ce}$  from Figures 2d and 3d electric fields from Figure 3, we estimate the wave normal angle  $\theta$  (angle between wave-vector and the background magnetic field) for the displayed wave-packets:  $\theta \in [60, 70]$  (see details of  $\theta$  estimation method in Ni, Thorne, Meredith, Shprits, and Horne (2011); Agapitov et al. (2014)), with a Gendrin angle  $\theta_g = \arccos(2ff_{ce}) \approx 45^\circ$  and a resonance cone angle  $\theta_r = \arccos(ff_{ce}) \approx 69^\circ$ . Thus, the observed waves indeed propagate obliquely with  $\theta \in [\theta_g, \theta_r]$ , that is, in the quasi-electrostatic mode (see statistics of such whistler-mode waves in Agapitov et al. (2013); Li, Santolik, et al. (2016)), and can interact with electrons through the Landau resonance (Artemyev, Agapitov, et al., 2016; Shklyar & Matsumoto, 2009).

The generation, and propagation without significant damping, of very oblique whistler-mode waves requires a very specific type of electron distribution function (Artemyev & Mourenas, 2020; Ma et al., 2017; Mourenas et al., 2015). For the observed waves, the Landau resonant energy at the equator starts at  $m_e(\omega/k_{\parallel})^2/2 \sim 1 - 5$  keV,

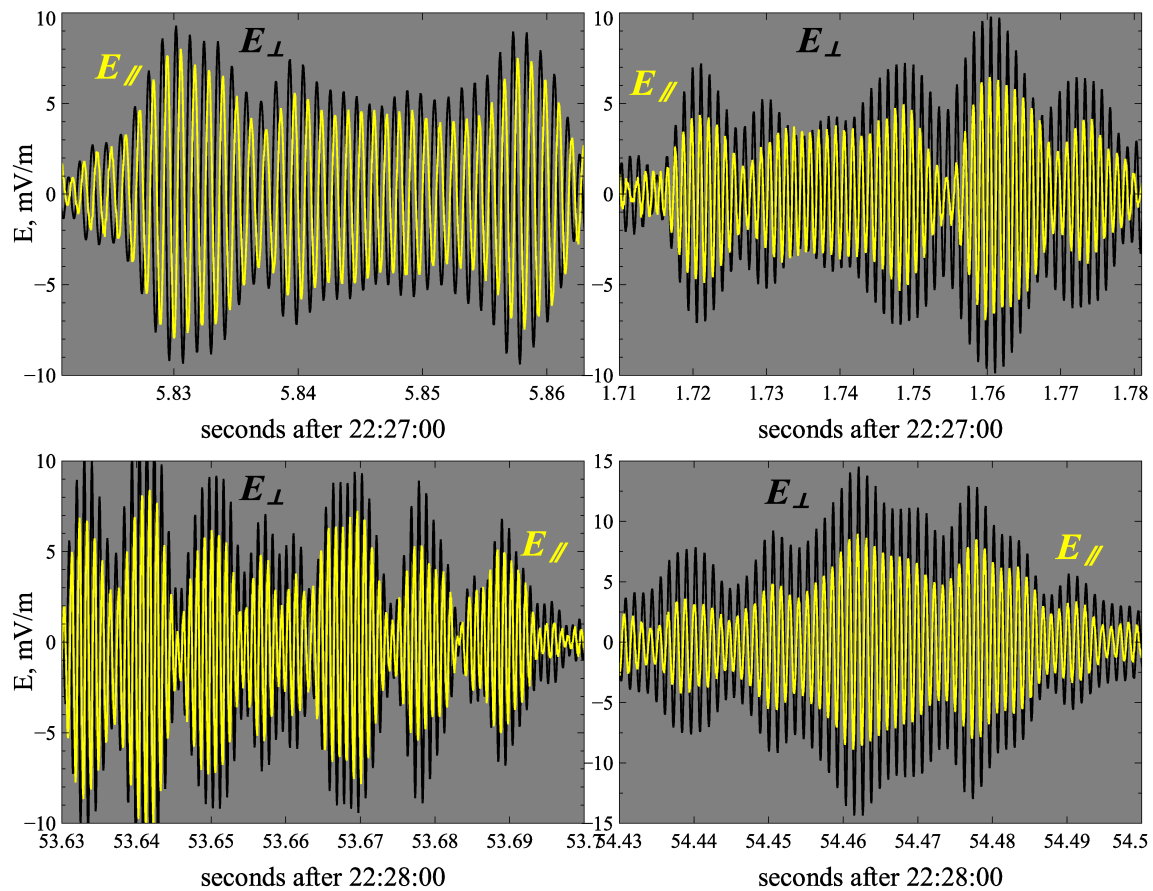


**Figure 2.** An overview of THEMIS E observations on 2021-01-01: (a) GSM magnetic fields, (b) and (c) magnetic and electric field spectra (for data format see Cully, Ergun, et al., 2008) with fractions of  $f_{ce}$  shown by horizontal curves, (d) plasma frequency to gyrofrequency ratio.

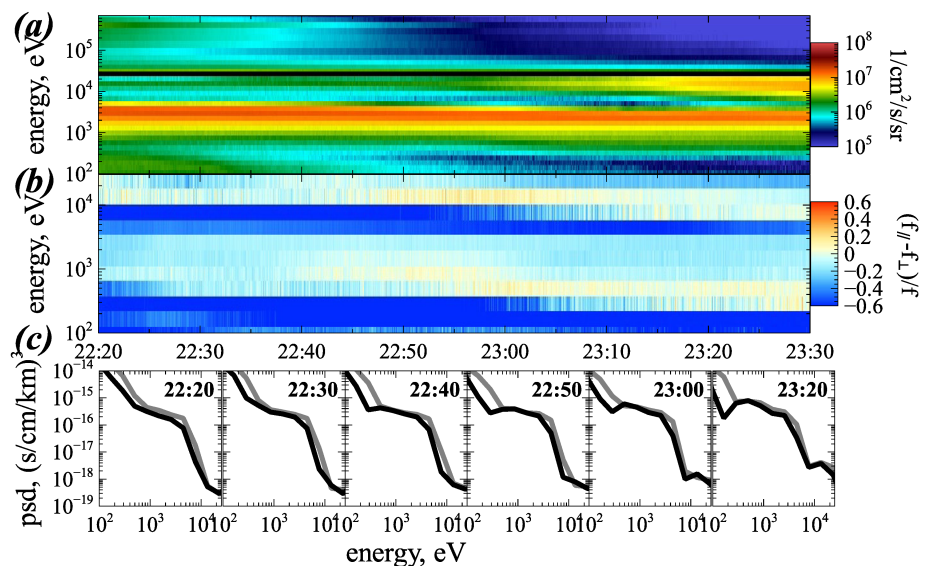
where  $\omega = 2\pi f$  is determined by measurements of wave spectrum and  $k_{\parallel}$  is determined by the wave dispersion for a measured wave normal angle. For typical electron distributions with a smoothly decreasing density toward higher parallel electron energy, Landau damping by these dense suprathermal electrons would quickly reduce the wave intensity to the noise level (Kennel, 1966; Bortnik et al., 2007; L. Chen et al., 2013). Accordingly, such very oblique waves are mainly associated with electron distributions having a plateau or weak field-aligned beams around the Landau resonant energy at the equator (Li, Mourenas, et al., 2016; R. Chen et al., 2019; Artemyev & Mourenas, 2020; Kong et al., 2021). Figure 4c shows that THEMIS E measured such a plateau both in the field-aligned and transverse phase space densities during the entire interval of whistler-mode wave observations (note also a weak field-aligned anisotropy of  $\sim 1$ –5 keV electrons in Figure 4b). The omnidirectional plateau represents the addition of a ring/shell electron distribution (see Figure 4a) with an almost null gradient in parallel velocity, that is, the absence of an intense negative distribution function gradient in parallel velocity that would provide Landau damping. Therefore, the conditions during this event (low  $f_{pe}/f_{ce}$  and absence of Landau damping) are favorable for very oblique wave generation and amplification (Artemyev, Agapitov, et al., 2016; Mourenas et al., 2015).

An interesting and important feature of the observed whistler-mode wave activity is the transient nature of wave bursts (see Figure 2b), which may be due to quasi-periodic wave generation modulated by compressional ultra-low-frequency waves (Li et al., 2011; Li, Thorne, et al., 2011; Xia et al., 2016; Zhang, Chen, et al., 2019). Such compressional waves were indeed observed concurrently by THEMIS E (not shown) over a wide frequency range. As the generation of very-oblique whistler-mode waves is associated with the Landau resonant scattering of  $<10$  keV electrons around the equator (Agapitov et al., 2016; Artemyev, Agapitov, et al., 2016; Li, Mourenas, et al., 2016), we may expect quasi-periodic enhancements of electron precipitating fluxes. To check this hypothesis, we examine measurements of precipitating electrons by the Precipitating Electron and Ion Spectrometer onboard two DMSP satellites (Hardy et al., 1984; Rich et al., 1985). Figure 5 (right column) shows DMSP electron spectra around the interval of THEMIS E observations of very oblique whistler-mode waves. There are clear spatially/temporally localized enhancements of  $\sim 5$  keV electron precipitation seen by DMSP at  $6 < L < 11$ . Such precipitation enhancements were not seen before the interval of THEMIS observations of oblique whistler-mode waves (see DMSP electron spectra in Figure 5 (left column)). Therefore, DMSP spectrometer measurements of

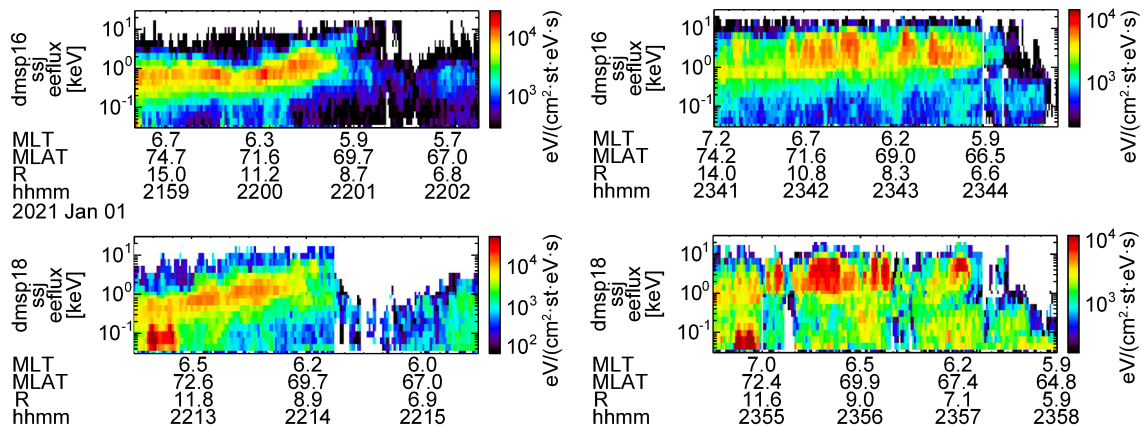




**Figure 3.** Four examples of whistler-mode wave packets during the event in Figure 2.

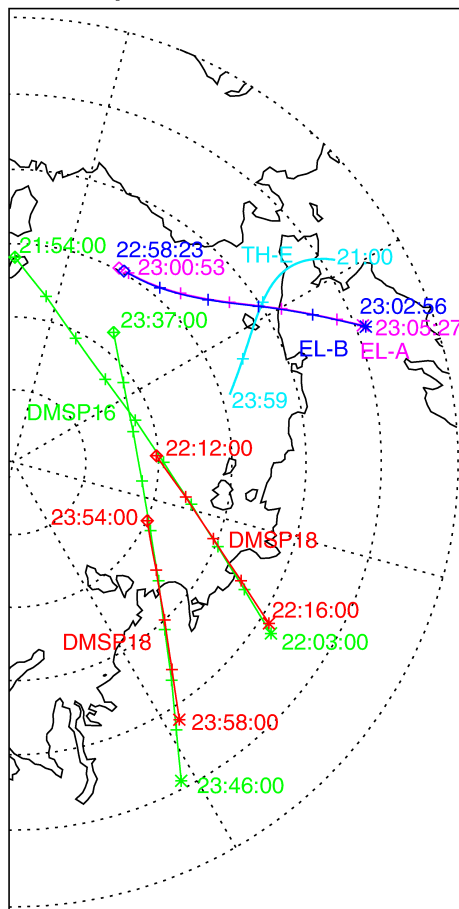


**Figure 4.** Electron flux measurements by THEMIS E during the event in Figure 2: (a) omnidirectional energy fluxes, (b) flux anisotropy in the <30 keV range, (c) parallel (black) and transverse (gray) phase space densities at several times.



**Figure 5.** Precipitating electron spectra from four Defense Meteorological Satellite Program (DMSP) orbits (two orbits of DMSP 16 and two orbits of DMSP 18) covering the interval of THEMIS and ELFIN observations: left column panels show DMSP data obtained before the detection of quasi-periodic oblique whistler-mode waves on THEMIS, whereas the right column panels show DMSP data during THEMIS observations of quasi-periodic oblique whistler-mode waves. Satellite MLT, magnetic latitude (MLT),  $L$ -shell ( $R$ ) are shown below each panel.

## Footprints 2021-01-01

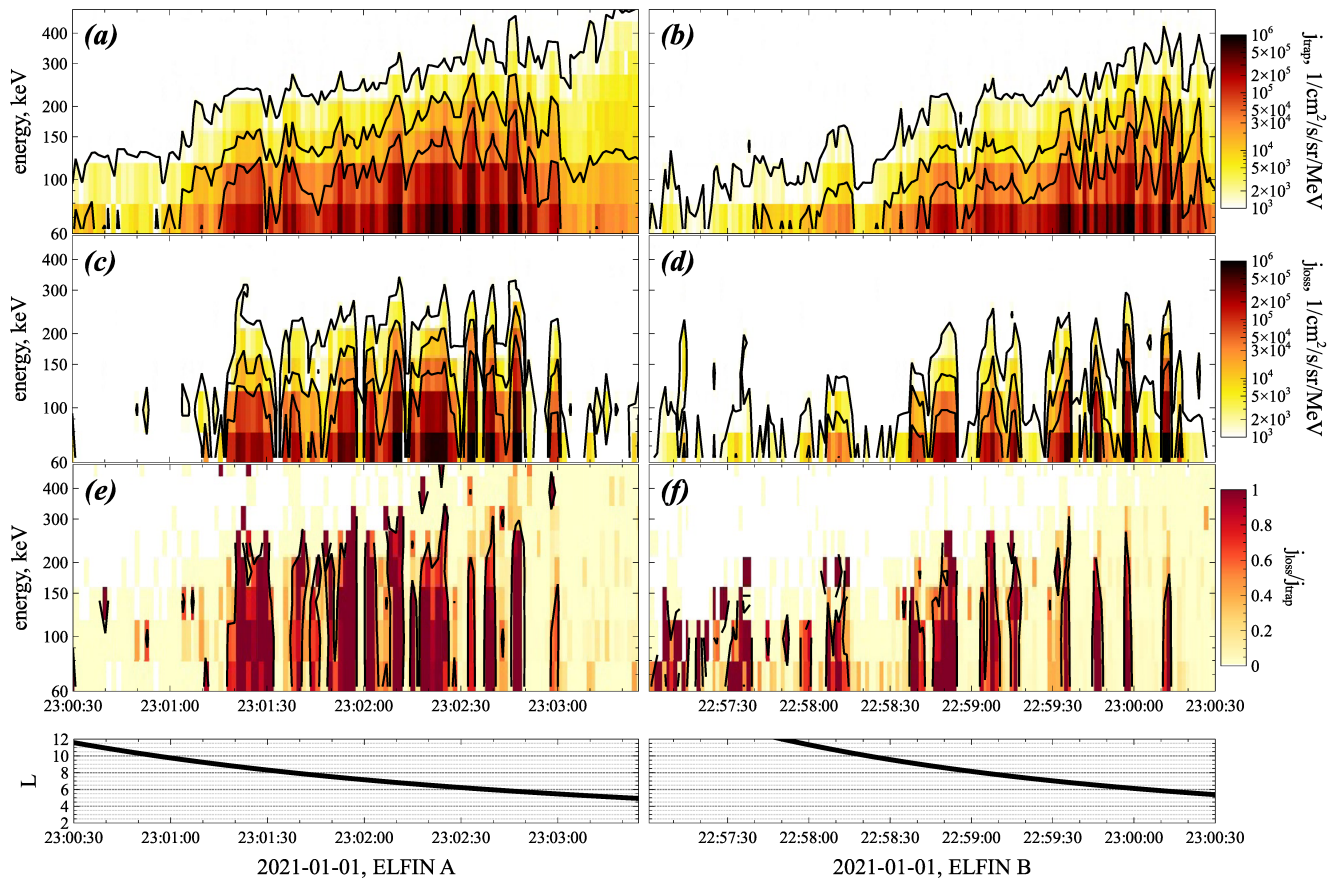


**Figure 6.** Footprints of THEMIS E, Defense Meteorological Satellite Program 16 and 18, and ELFIN A and B orbits for the event from Figures 2 and 5.

<10 keV electron precipitation are consistent with the presence of a transient (quasi-periodic) equatorial generation of very oblique whistler-mode waves.

Let us now consider conjugate low-altitude ELFIN observations and equatorial THEMIS E measurements of very oblique whistler-mode waves. There are a couple of orbits of ELFIN A and B within the same time interval and  $L$ -shell, MLT range (see Figure 6 for THEMIS, DMSP, and ELFIN orbits as a function of MLT and  $L$ -shell). Figures 7a and 7b shows that the energy spectra of trapped electron fluxes on both ELFIN CubeSats demonstrate a typical outer radiation belt structure, as indicated by the flux (and energy range) increase from  $L \sim 11$  to  $L \sim 5$  (where the plasmapause was at the time, according to the THEMIS E plasma data), and the disappearance of <500 keV fluxes due to effective scattering at these energies by plasmaspheric hiss waves (Ma, Li, Thorne, Bortnik, et al., 2016; Mourenas et al., 2017; Mourenas & Ripoll, 2012). Within the  $L$ -shell range corresponding to the outer radiation belt, ELFIN observed transient, strong precipitation of <200 keV electrons (see panels (c,d)). The timescale of these precipitation bursts is about (or even less than) ELFIN's half-spin-period, 1.5s. Both ELFIN CubeSats show similar patterns of transient precipitation, but the temporal (spatial) distributions of precipitation bursts are different on these two spacecraft moving along the same orbit with an  $\sim 2$  min time delay. Therefore, we infer that these precipitation events are probably driven by electron scattering by transient whistler-mode waves which, due to their amplitude modulation by ULF waves, do not maintain the same spatial distribution at the equator over time intervals longer than  $\sim 1$ –2 min.

Figures 7e and 7f shows the ratio of precipitating to trapped electron fluxes. Most of the precipitation bursts reach a ratio  $\sim 1$  at energies <200 keV, that is, they represent very strong precipitation with an entirely filled loss cone, roughly corresponding to the so-called strong diffusion limit (Kennel, 1969). Note that ELFIN is a spinning CubeSat (Angelopoulos et al., 2020) and that different pitch-angles are measured at slightly different times within the 1.5s half-spin-period. Thus, sub-second precipitation events may demonstrate precipitating fluxes even higher than trapped fluxes, but this can either be an effect of a specific precipitation mechanism (see Zhang, Artemyev, et al., 2022) or it can be an artifact due to aliasing, that is, the absence of



**Figure 7.** An overview of ELFIN A and B measurements for the event from Figure 2: (a, b) trapped fluxes, (c, d) precipitating fluxes, (e, f) precipitating to trapped flux ratio. Bottom panels shows ELFIN  $L$ -shells from (Tsyganenko, 1989) model.

strictly simultaneous measurements of trapped and precipitating fluxes. Nevertheless, the observed precipitation events are transient, short-scale, and very intense, as expected for microbursts.

### 3. Theory of Nonlinear Landau Resonance

In contrast to classical models of microburst precipitation events due to cyclotron resonance with field-aligned chorus waves (e.g., L. Chen et al., 2020; L. Chen et al., 2022), precipitation events in Figure 7 are observed in conjunction with equatorial measurements of very oblique whistler-mode waves (see Figure 5). Therefore, an alternative mechanism of precipitation due to the Landau resonance should be invoked here (see Figure 1 and Agapitov et al. [2015b]). To verify the efficiency of precipitation due to Landau resonance, we use the mapping technique proposed in Artemyev et al. (2020b). This technique provides an alternative to test particle simulations (widely used for investigation of nonlinear resonances, see, e.g., Tao et al. [2013]; Gan et al. [2020]; Agapitov et al. [2015b]; Allanson et al. [2020]; Zhang, Agapitov, et al. [2020]), to the Green function method (see, e.g., Omura et al. [2015]; Hsieh and Omura [2017b]), or to the inclusion of integral operators into the Fokker-Plank equation (see, e.g., Artemyev, Neishtadt, Vasiliev, and Mourenas [2018]; Vainchtein et al. [2018]). The advantages of the mapping technique are that (a) it combines the analytical theory for energy/pitch-angle change in nonlinear resonance (Neishtadt, 1975, 1999), (b) it allows to incorporate the effects of resonant interactions with a wave ensemble sharing the properties of the observed waves (Artemyev, Neishtadt, Vasiliev, Zhang, et al., 2021; Vainchtein et al., 2018; Zhang, Artemyev, et al., 2022), and (c) it is computationally fast. This technique resembles the method of stochastic differential equations for the Fokker-Plank equation (e.g., Tao et al., 2008), but it further includes nonlinear resonance effects such as phase trapping and bunching (Artemyev et al., 2020b).

The mapping technique applicability is based on a set of assumptions:

1. Although the wave ensemble (collected over a long time interval) can have a quite broad distribution over wave parameters (amplitudes, frequencies, etc.), the waves are assumed to be sufficiently spread along magnetic field lines so that there can be no resonance overlapping, that is, it is assumed that there is no destruction of nonlinear wave-particle interaction due to resonance overlap (Nunn, 1986; Tao et al., 2013). This condition is generally satisfied for narrow-band intense whistler-mode waves, but it can be violated by various secondary effects in the presence of non-resonant noise (Artemyev, Mourenas, et al., 2015; Brinca, 1972), wave phase decoherence (Zhang, Agapitov, et al., 2020), or wave-packet modulation (An et al., 2022; Gan et al., 2020; Gan et al., 2022; Hiraga & Omura, 2020; Nunn et al., 2021; Zhang, Mourenas, et al., 2020). The inclusion of such effects, which limit the efficiency of nonlinear resonant interaction, is generally possible (e.g., Artemyev, Neishtadt, Vasiliev, & Mourenas, 2021; Mourenas et al., 2018), but since it requires significant modifications to theoretical estimates, we did not include these effects into the current version of the mapping technique.
2. Each of the wave-particle resonant interactions should be considered independently, that is, we assume no correlation of electron gyrophase between two successive resonant interactions. Such correlations may significantly modify the efficiency of electron scattering by waves (see discussion in, e.g., Allanson et al., 2022; Osmane & Lejosne, 2021), but for realistic, inhomogeneous magnetic fields and realistic whistler-mode wave amplitudes, effects due to phase correlations are negligibly weak (see discussions in Artemyev et al., 2020b; Artemyev et al., 2020a).
3. The nonlinear resonant wave-particle interaction cannot transport electrons outside the (energy, pitch-angle) domain where such an interaction is possible, that is, there is no lost phase trapping/phase bunching along the electron orbit. This condition is satisfied owing to the asymptotes of (energy, pitch-angle) changes around the boundaries of the domain of nonlinear resonant interactions (see Artemyev et al. [2019] and Appendix B in Artemyev, Neishtadt, Vasiliev, Zhang, et al. [2021]).

The mapping technique describes electron energy  $\gamma$  (electron gamma factor) and equatorial pitch-angle  $\alpha_{eq}$  changes due to resonant interactions (note that these variables remain constant between resonant interactions) as

$$\gamma_{n+1} = \gamma_n + \begin{cases} \Delta\gamma_{trap}, & \xi \in [0, \pi] \\ \Delta\gamma_{scat}, & \xi \in (\pi, 1] \end{cases}, \quad \alpha_{eq,n+1} = \alpha_{eq,n} + \begin{cases} \Delta\alpha_{eq,trap}, & \xi \in [0, \pi] \\ \Delta\alpha_{eq,scat}, & \xi \in (\pi, 1] \end{cases} \quad (1)$$

where  $n$  is the number of resonant interactions (the mapping iteration number),  $\Delta\gamma_{trap}$  and  $\Delta\gamma_{scat}$  are energy changes due to trapping and bunching (the same for  $\alpha_{eq}$ ),  $\pi \leq 1$  is the trapping probability ( $\pi$  can be interpreted as the relative number of resonant electrons which experienced trapping during a single resonant interaction), and  $\xi$  is a random value with a uniform distribution over  $[0, 1]$  (Artemyev et al., 2020b). Three main characteristics of the nonlinear resonant interactions,  $\Delta\gamma_{trap}$ ,  $\Delta\gamma_{scat}$ , and  $\pi$  (note that  $\Delta\alpha_{eq,trap}$ ,  $\Delta\alpha_{eq,scat}$  can be recalculated from  $\Delta\gamma_{trap}$ ,  $\Delta\gamma_{scat}$ ) depend on electron energy  $\gamma_n$  and equatorial pitch-angle  $\alpha_{eq,n}$ , and on wave and background plasma characteristics. The energy change due to the phase bunching (also called nonlinear scattering) is the mean energy change of all electrons that did not experience trapping within a single resonant interaction. This is a local characteristic depending on particle, wave, and background magnetic field characteristics of the resonance location (or the latitude of resonance). Examples of  $\Delta\gamma_{scat}$  evaluation for various systems with nonlinear resonance can be found in (Albert, 2002; Artemyev, Vasiliev, Mourenas, Agapitov, & Krasnoselskikh, 2014; Artemyev, Vasiliev, Mourenas, Agapitov, Krasnoselskikh, et al., 2014; Itin et al., 2000; Neishtadt et al., 2011; Shklyar, 2011), whereas the basic theory of  $\Delta\gamma_{scat}$  evaluation can be found in (Albert, 1993; Karpman et al., 1974; Karpman & Shklyar, 1977; Neishtadt, 1999). There is an important relation between  $\Delta\gamma_{scat}$  and the trapping probability:  $\pi = \nabla\Delta\gamma_{scat}$  (see Artemyev, Neishtadt, et al., 2016; Artemyev, Neishtadt, Vainchtein, et al., 2018), where the gradient is along the resonance curve  $n_r\gamma - (\omega/\Omega_{ce,eq})(\gamma^2 - 1)\sin^2\alpha_{eq}/2 = const$  (Shklyar & Matsumoto, 2009) that connects  $\gamma$  and  $\alpha_{eq}$  changes for a given wave frequency  $\omega$ , equatorial electron cyclotron frequency  $\Omega_{ce,eq}$ , and resonant number  $n_r$  ( $n_r = 1$  for the first cyclotron resonance and  $n_r = 0$  for the Landau resonance). Examples of the theoretical evaluation and numerical verification of  $\pi$  for electron resonant interaction with whistler-mode waves can be found in Artemyev, Vasiliev, Mourenas, Agapitov, and Krasnoselskikh (2014); Artemyev, Vasiliev, et al. (2015); Artemyev et al. (2020b). In contrast to  $\Delta\gamma_{scat}$  and  $\pi$ , the energy change due to trapping  $\Delta\gamma_{trap}$  is a nonlocal characteristic and depends on the wavefield distribution along magnetic field lines (see, e.g., Artemyev, Krasnoselskikh, et al., 2012; Agapitov et al., 2014; Artemyev, Neishtadt, Vasiliev, & Mourenas, 2018),  $B_w(\lambda)$  where  $\lambda$  is the magnetic latitude. Such a distribution can be derived from empirical models of whistler-mode wave



intensity (Agapitov et al., 2015a, 2018; Wang & Shprits, 2019) or from numerical tracing of whistler-mode waves (Hsieh & Omura, 2017b; Omura et al., 2015). For very oblique waves, we use a simplified model that fits the observed  $B_w(\lambda)$  distribution obtained from satellite statistics by Agapitov et al. (2018):

$$B_w = B_{w,eq} \cdot \frac{1}{2} (1 + \tanh(\lambda/\delta\lambda_1)) \cdot \exp(-\lambda^2/\delta\lambda_2^2) \quad (2)$$

where  $\delta\lambda_1$  determines the spatial extent of the wave generation region, and  $\delta\lambda_2$  determines the latitude range of wave propagation before damping at middle latitudes. This model describes waves in the  $\lambda > 0$  hemisphere, and we assume a symmetric wavefield distribution relative to the equator, that is,  $B_w(\lambda) = B_w(-\lambda)$ . Wave frequency can be assumed to be fixed along the wave propagation (or slowly varying in time, if we deal with chorus waves), whereas the wave number profile  $k(\lambda)$  can be obtained from the cold plasma dispersion (Stix, 1962) for a given wave normal angle distribution  $\theta(\lambda)$  at this particular frequency. For very oblique whistler-mode waves, we use a  $\theta = \theta_r(\lambda) - \Delta\theta$  model, with  $\theta_r = \arccos(\omega/\Omega_{ce})$  being the resonant cone angle and  $\Delta\theta = \text{const}$  the model parameter specifying the wave normal angle deviation from  $\theta_r$ . This wave normal angle model is based on previous observations of very oblique whistler-mode wave propagation around the resonance cone angle (Agapitov et al., 2013; Li, Santolik, et al., 2016). The cold plasma dispersion relation does not work for  $\Delta\theta \rightarrow 0$ , where the thermal electron contribution to the wave dispersion properties becomes crucial (Sazhin & Horne, 1990). Theoretical estimates (Artemyev, Agapitov, et al., 2016; Mourenas et al., 2014) and spacecraft observations (Ma et al., 2017) suggest that the minimum  $\Delta\theta$  can be determined from the limitation of the wave refractive index  $N = kc/\omega$  to values  $N < N_{hot} \approx 100\text{--}300$  for typical thermal electron energies in the inner magnetosphere.

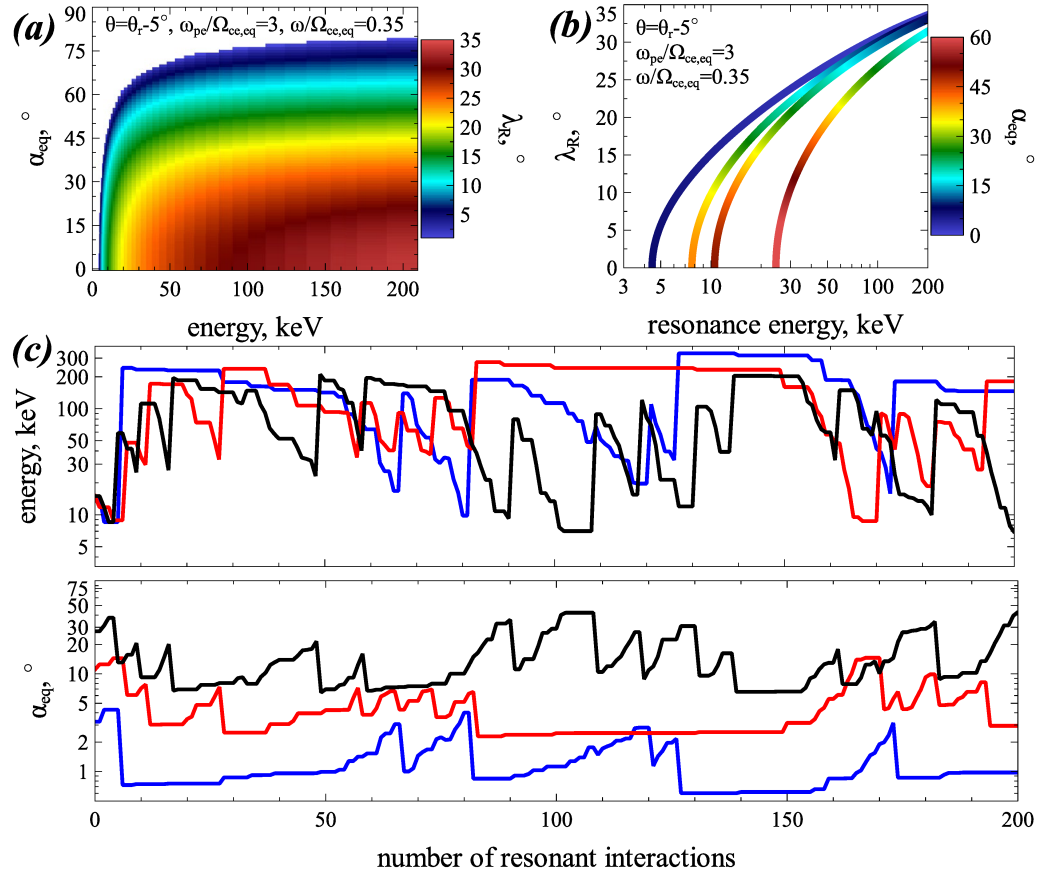
During Landau resonance ( $\omega\gamma - k_{\parallel}c\sqrt{\gamma^2 - 1}\cos\alpha = 0$ ), the electron propagates in the same direction as the wave, in contrast to cyclotron resonance characterized by opposite propagations of wave and electron. For a fixed wave frequency at Landau resonance, the resonance curve  $(\omega/\Omega_{ce,eq})(\gamma^2 - 1)\sin^2\alpha_{eq}/2 = \text{const}$  (Shklyar & Matsumoto, 2009) is given by the equation of the constant magnetic moment (we use the normalized moment  $I_x = (\omega/\Omega_{ce,eq})(\gamma^2 - 1)\sin^2\alpha_{eq}/2$ ). Combining this equation with the resonance condition, we obtain the equation for dependence of the Landau resonant energy on the magnetic latitude:

$$\gamma_R = \frac{1}{\sqrt{1 - (\omega/c k_{\parallel})^2}} \sqrt{1 + 2I_x\Omega_{ce}} \quad (3)$$

Figure 8a shows the latitude of resonance  $\lambda_R(\gamma, \alpha_{eq})$ , whereas Figure 8b shows  $\lambda_R(\gamma_R)$  for fixed  $I_x$ . Small pitch-angle ( $\alpha_{eq} < 30^\circ$ ) low-energy (<10 keV) electrons may resonate with waves around the equatorial plane ( $\lambda_R < 15^\circ$ ), whereas energetic (~150 keV) electrons reach resonance only at  $\lambda_R > 30^\circ$ . The wave intensity increases with latitudes  $dB_w/d\lambda > 0$  around the equator (Agapitov et al., 2018; Omura et al., 2008). This condition suggests that electrons may be trapped by waves. Being trapped at small  $\lambda_R$ , electrons are transported by the wave to higher latitudes (with a decrease of their pitch-angle) and escape from the resonance at a latitude where  $dB_w/d\lambda$  becomes sufficiently strong (for a precise description of the trapping and de-trapping conditions, see, e.g., Artemyev, Krasnoselskikh, et al. (2012); Artemyev, Vasiliev, et al. (2013)). Such an electron transport is associated with an energy increase, for example, a 5 keV equatorial electron can gain 100 keV before reaching  $\lambda_R \sim 30^\circ$  and finishing within the loss cone (see Figure 8b). Therefore, the trapping moves lower energy/larger pitch-angle electrons toward the larger energy/lower pitch-angle region in phase space (see schematic in Figure 1). The Landau resonance scattering moves electrons along the same resonance curve  $I_x = \text{const}$ , but in the direction opposite to the trapping motion. Therefore, Landau resonance results in electron drifts toward smaller energy/larger pitch-angle (the phase bunching effect) and more rare jumps to larger energy/smaller pitch-angle (the phase trapping effect). Figure 8c shows three examples of such electron dynamics evaluated with the mapping technique (1) for typical wave characteristics.

### 3.1. Mapping Model for Observed Wave Characteristics

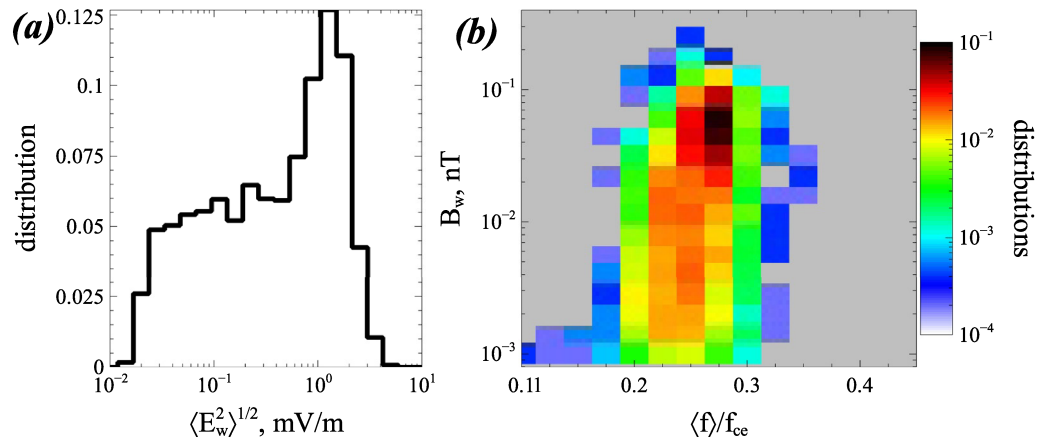
Near-equatorial spacecraft measurements of whistler-mode waves provide distributions of wave characteristics, for example, the average wave intensity for various frequency ranges,  $L$ -shells,  $MLT$ , and  $\lambda$  (Meredith et al., 2001, 2012; Agapitov et al., 2013, 2015a). Such time-averaged intensities can be directly applied to the evaluation of quasi-linear diffusion rates (e.g., Agapitov et al., 2018; Horne et al., 2013; Ma et al., 2018), but



**Figure 8.** Panel (a) shows the resonance latitudes as a function of energy and equatorial pitch-angle. Panel (b) shows variations of the trapped particle latitude with energy given by Equation 3, that is, each curve shows the resonant condition for constant  $I_r$ . Colors denote variations of the equatorial pitch-angle. Panels (c) shows energy and equatorial pitch-angle for three trajectories evaluated with the mapping technique (1).

they are not suitable for the evaluation of nonlinear wave-particle interactions. Instead of time-averaged intensities, the nonlinear resonance models require distributions (occurrence rates) of intense waves that can interact with electrons nonlinearly (Zhang et al., 2018; Zhang, Mourenas, et al., 2019). Therefore, we use THEMIS E measurements during the event in Figure 2 to obtain the distribution of wave amplitudes  $E_w$ , as inferred from the wave frequency spectrum  $\mathcal{E}_w^2(f)$ :  $E_w^2 = \int_{f_{ce}/20}^{f_{ce}/2} \mathcal{E}_w^2(f) df$ . As shown in Figure 9a, most waves have amplitudes  $<3$  mV/m. However,  $E_w$  here is inferred from the 1s-averaged spectrum (Cully, Ergun, et al., 2008) and the actual magnitude of individual wave-packets is often significantly larger (see Figure 3). To account for such an underestimation of  $E_w$  due to time averaging, we introduce a multiplicative factor  $K$ , that is, we assume that electrons interact with waves of amplitudes  $K \times E_w$ . To keep the same average intensity as without the  $K$  factor, we assume that electrons resonate with waves only a fraction  $(1/K^2)$  of the simulation time, whereas during the remaining fraction  $(1 - 1/K^2)$  of the time there is no resonant interaction.

The wave-particle resonant interaction is not determined solely by  $E_w$ , but also by the wave frequency, wave normal angle, and the  $E_w(\lambda)$  profile along the magnetic field line. The wave frequency can be determined from the wave frequency spectrum:  $\langle f \rangle = \int_{f_{ce}/20}^{f_{ce}/2} f \cdot \mathcal{E}_w^2(f) df / E_w^2$ . Figure 9b shows the  $\mathcal{P}(B_w, \langle f \rangle)$  distribution for the event in Figure 2, with  $B_w$  recalculated from  $E_w$  using the cold plasma dispersion (see examples in, e.g., Ni, Thorne, Meredith, Horne, & Shprits, 2011; Agapitov et al., 2014). Using this  $\mathcal{P}(B_w, \langle f \rangle)$  distribution, we can rewrite the mapping Equation 1 as a two-step equation. During each half of the bounce period (time interval between two successive resonances; this is further discussed in Section 4.3), we generate a random number to select a  $(B_w, \langle f \rangle)$  bin from the  $\mathcal{P}$  distribution. According to this number, we choose the wave amplitude  $B_w$  and wave frequency  $\langle f \rangle / f_{ce}$  to calculate three main characteristics of wave-particle interactions:  $\Delta \gamma_{trap}$ ,  $\Delta \gamma_{scat}$ , and  $\pi$ . Then, we make one iteration for energy and pitch-angle using Equation 1. Thus, the observed distribution of wave



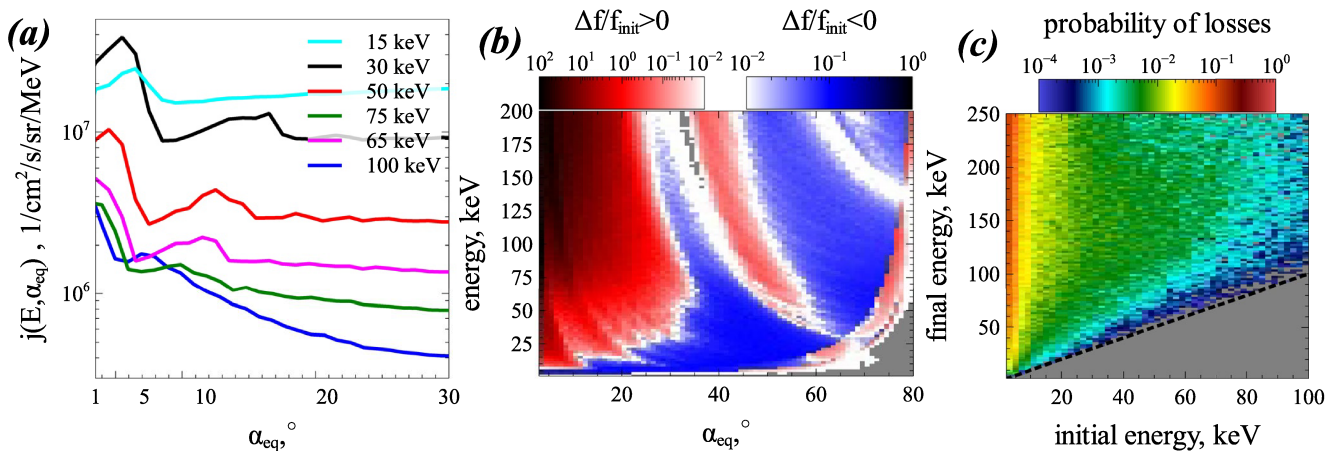
**Figure 9.** Panel (a) shows the distribution of electric field amplitudes derived from the wave frequency spectra (*fff* data) shown in Figure 2c. Panel (b) shows the distribution of wave amplitudes and frequencies for the event in Figure 2. Wave magnetic field amplitudes are converted from the wave electric field amplitudes using cold plasma dispersion relation (Tao & Bortnik, 2010).

characteristics,  $\mathcal{P}(B_w, \langle f \rangle)$ , determines  $\Delta\gamma_{trap}$ ,  $\Delta\gamma_{scat}$ , and  $\pi$  for each resonant interaction. Figure 8c shows several examples of electron trajectories calculated with the observed  $\mathcal{P}$  distribution.

### 3.2. Simulation Results

Using the mapping technique with the observed distribution  $\mathcal{P}(B_w, \langle f \rangle)$ , we integrate a large ensemble of electron trajectories having an initial (energy, pitch-angle) distribution from Figure 4. As we are interested in precipitation patterns, we restrict the integration time to  $\sim 1$  bounce period of 100 keV electrons.

All electrons with  $\alpha_{eq} < 2^\circ$  are considered to be precipitating (as appropriate at  $L < 8$ ). Figure 10a shows several pitch-angle distributions for different energies after wave-particle resonant interactions. There is a clear peak of small  $\alpha_{eq}$  electron fluxes, and this peak is most pronounced for energies higher than  $\sim 30$  keV. The formation of such a field-aligned (and partially precipitating) electron population is produced by the electron trapping acceleration in the Landau resonance. Despite the small probability of electron Landau trapping  $\pi$  (see, e.g., Artemyev, Krasnoselskikh, et al., 2012; Artemyev, Vasiliev, Mourenas, Agapitov, & Krasnoselskikh, 2014), this



**Figure 10.** Results of the mapping simulation of electron nonlinear resonant interactions for the distribution of wave amplitudes and frequencies during the event in Figure 2. Panel (a) shows pitch-angle distributions for different energies after  $\sim 10 R/c \sim 2$  s of interactions (sufficient to make at least one bounce period for all considered particles). The initial (equatorial) electron phase space density is obtained by fitting THEMIS E observations. Panel (b) shows the ratio of final (after  $\sim 10 R/c \sim 2$  s) to initial phase space densities. Panel (c) shows the probability of losses, that is, the percentage of particles in each bin that is transported into the loss cone within the simulation interval ( $\sim 10 R/c \sim 2$  s). We use a  $K = 3$  factor for this simulation, with the corresponding probability of wave-particle interactions  $1/K^2$ .

acceleration mechanism is quite effective, because waves trap particles with initially smaller energies ( $\sim$  a few keV) and accelerate them up to  $\sim 30$ – $100$  keV. Therefore, the large difference of initial phase space densities in the energy of trapping and release compensates the small probability of such trapping ( $\pi$ ) and allows for the formation of a very substantial field-aligned electron population. Figure 10b shows the ratio of final (after the wave-particle interaction) to initial phase space densities: there is a clear particle transport toward the domain of smaller pitch-angles and larger energies. An additional effect of Landau trapping acceleration is a flux enhancement at intermediate (medium to high) pitch-angles that can produce a butterfly distribution (see also Ke et al., 2022; Mourenas et al., 2016).

The time-scale of such electron acceleration and precipitation is about a bounce period, that is, the time-scale of electron transport in the Landau trapping. This represents an insufficiently long time for nonlinear wave-particle interactions to establish a fine balance between trapping and phase bunching (nonlinear scattering) and, thus, electron transport toward the loss cone (due to trapping) may not be fully compensated by electron transport to higher pitch-angles (due to bunching). After a sufficiently long wave-particle interaction, such a fine balance will establish a new distribution function without strong gradients along the resonance curve  $I_x = \text{const}$  (Artemyev et al., 2019). Thus, electron precipitation associated with the Landau trapping acceleration should be provided by bursts of very oblique whistler-mode waves, and they can share properties of microbursts. To demonstrate that such precipitation events do not decrease the  $\sim 100$  keV flux population, we plot the probability of electron transport into the loss cone and the final electron energy within the loss cone as a function of initial electron energy. Figure 10c shows 2D initial, final energy space, where each bin shows the percentage of particles within the loss cone with fixed initial and final energies, showing for instance that if electrons with initial energy  $\sim 10$  keV are accelerated to  $\sim 100$  keV, about 5% (yellow color of the color bar) of them appear in the loss cone. Figure 10c shows that precipitating electrons (within the loss cone) with a final energy of  $\sim 100$  keV had very rarely an initial energy of  $> 80$  keV. Figure 10c further shows that such precipitating electrons with a final energy of  $\sim 100$  keV mainly result from trapping, acceleration, and transport via Landau resonance of initially  $\sim 5$  keV electrons (more than 20% of such electrons are transported into the loss cone),  $\sim 10$  keV electrons (more than 5% of such electrons are transported into the loss cone), and  $\sim 20$  keV electrons (more than 1% of such electrons are transported into the loss cone). However, this also demonstrates that most ( $> 80\%$ ) of these trapped electrons actually do not precipitate, but are released from trapping at small pitch-angles outside the loss cone (see Figure 10b). Therefore, intense (microburst) precipitation events can indicate an effective electron acceleration and an increase of  $\sim 50$ – $200$  keV electron fluxes at small pitch-angles, rather than being a signal of electron flux depletion. The accelerated electrons ending up outside the loss cone, however, have already small pitch-angles, and can be later scattered into the loss cone by much weaker field-aligned whistler-mode waves (see discussions of such double-resonance mechanism of electron losses due to Landau trapping and cyclotron scattering in Mourenas et al., 2016; Ma, Mourenas, et al., 2016; Hsieh et al., 2021).

## 4. Discussion

In this section, we discuss the implications of sub-relativistic electron precipitation driven by nonlinear Landau trapping due to very oblique whistler-mode waves. We compare this proposed microburst precipitation mechanism with the more commonly discussed precipitation mechanism, nonlinear cyclotron resonance scattering of electrons by field-aligned whistler-mode waves (e.g., Breneman et al., 2017; L. Chen et al., 2020). We also discuss possible limitations of the proposed mechanism of nonlinear Landau trapping.

### 4.1. Landau Trapping Versus Cyclotron Phase Bunching

In this study, we examined the Landau resonance of  $\sim 1$ – $100$  keV electrons and very oblique whistler-mode waves. Trapping into Landau resonance results in electron acceleration and pitch-angle decrease, that is, electrons are transported by the waves to the loss cone. The alternative, and more widely investigated, mechanism of electron precipitation by intense whistler-mode waves is phase bunching due to cyclotron resonance (Bell, 1984; Karpman et al., 1974; Karpman & Shklyar, 1977). In contrast to Landau trapping, cyclotron phase bunching is associated with an electron energy decrease (e.g., Albert, 2000), so that precipitating electrons provide whistler-wave amplification (e.g., Demekhov, Taubenschuss, & Santolík, 2017; Omura, 2021; Shklyar, 2017; Tao et al., 2021). There are two questions that merit further discussion at this point: (a) can we distinguish precipitation due to Landau

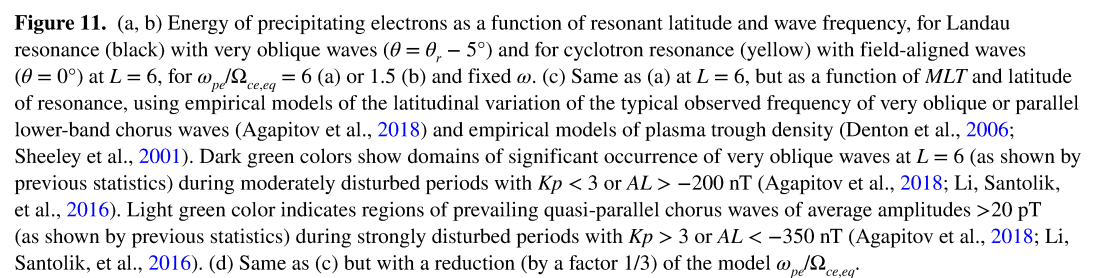


trapping from that which is due to cyclotron scattering? (b) would cyclotron scattering result in a depletion of radiation belt energetic electron fluxes, contrary to Landau trapping which do not deplete such fluxes?

To address the first question, we analyze resonant energies for both Landau trapping and cyclotron phase bunching. For monochromatic waves, the resonant wave-particle interaction conserves the invariant  $m_e c^2 n_r \gamma - I_x \omega$ , with  $n_r = 0$  for the Landau resonance and  $n_r = 1$  for the cyclotron resonance (e.g., Shklyar & Matsumoto, 2009). These invariants define resonance curves in the energy and pitch-angle space, and resonant interaction moves particles along these curves (Summers et al., 1998). Figure 1 shows the resonance curves for cyclotron and Landau resonance: Landau trapping results in a significant energy increase and transport into the loss cone, whereas cyclotron bunching near the loss cone does not change the electron energy significantly. Therefore, we examine resonant energies in the case of cyclotron bunching and energies at the time of release from resonance in the case of Landau trapping. We first consider the 2D parametric space of wave frequency,  $\omega/\Omega_{ce,eq}$ , and latitude of precipitation,  $\lambda_R$  (latitude where electrons are moved into the loss cone). This is the latitude of cyclotron resonance for phase bunching, or the latitude of electron release from Landau trapping. Figures 11a and 11b shows precipitating resonant electron energies in  $(\omega/\Omega_{ce,eq}, \lambda_R)$  space for cyclotron bunching by field-aligned waves (yellow curves) and for Landau trapping by very oblique waves (black curves). We show the results for two values of  $\omega_{pe}/\Omega_{ce,eq}$ , corresponding to the typical plasma density at  $L \sim 6$  with  $\omega_{pe}/\Omega_{ce,eq} = 6$  (Sheeley et al., 2001), and to rarefied plasma conditions  $\omega_{pe}/\Omega_{ce,eq} = 1.5$ . Figures 11a and 11b shows that for realistic latitudes  $\lambda_R < 40^\circ$  (see the wave intensity distribution with latitude in Agapitov et al., 2018), the Landau trapping would be effective only for low plasma density and can provide rapid electron precipitation at  $< 200$  keV, whereas relativistic electron precipitation should be attributed to the cyclotron scattering at  $\lambda_R \gtrsim 30^\circ$ . Therefore, if waves are confined around the equator (i.e.,  $\lambda_R < 20^\circ$ , as is typical for intense nightside whistler-mode waves; see Meredith et al. (2012); Agapitov et al. (2013)), only cyclotron bunching can provide  $< 200$  keV microbursts. Moreover the energy ranges and, possibly, the *MLT* ranges of microbursts, could be used to separate the two precipitation mechanisms: Landau trapping and cyclotron bunching.

To further investigate the *MLT* dependence, Figures 11c and 11d displays electron precipitating (resonant) energies in the  $(MLT, \lambda_R)$  space for cyclotron bunching by field-aligned waves (yellow curves) and for Landau trapping by very oblique waves (black curves) at  $L = 6$ . Precipitating electron energies are now calculated using empirical models (Agapitov et al., 2018) for the latitudinal variation of the power-weighted mean frequency  $\omega/\Omega_{ce,eq} = \max(0.17, 0.41 - 0.012 \cdot \lambda)$  of very oblique lower-band chorus waves and of the power-weighted mean frequency  $\omega/\Omega_{ce,eq} = \max(0.1, 0.35 - 0.0125 \cdot \lambda)$  of quasi-parallel chorus waves (see also Bunch et al., 2013), along with an empirical model of the equatorial plasma trough density (Sheeley et al., 2001), and an empirical model of the plasma density latitudinal variation (Denton et al., 2006). In Figures 11c and 11d, the dark green region shows the  $(MLT, \lambda_R)$  domain where very oblique waves have significant occurrence rates (sometimes even higher than for parallel waves) in statistics near  $L = 6$ , which mostly occurs during moderately disturbed periods with  $Kp < 3$  or  $AL > -200$  nT (Agapitov et al., 2018; Li, Santolik, et al., 2016). Light green regions correspond to the  $(MLT, \lambda_R)$  domains of prevailing quasi-parallel chorus waves of average amplitudes  $> 20$  pT at  $L = 6$ , which occurs during strongly disturbed periods with  $Kp > 3$  or  $AL < -350$  nT (Agapitov et al., 2018). Figure 11c shows that low energy ( $< 25$  keV) microbursts due to Landau resonant interactions with very oblique waves should mainly occur at 21–13 *MLT* during moderately disturbed periods (with  $Kp < 3$  or  $AL > -200$  nT) with typical plasma trough density levels (with  $\omega_{pe}/\Omega_{ce,eq} \sim 6$ ). In contrast, cyclotron resonance with parallel waves should lead to microbursts of much higher energies, up to  $\sim 150$ – $250$  keV at 4–12 *MLT*, more frequent during strongly disturbed periods (with  $Kp > 3$  or  $AL < -350$  nT).

However, Figure 11d also indicates that Landau resonant interactions with very oblique waves can produce microbursts of higher energies at  $L = 6$ , extending up to  $\sim 100$ – $150$  keV at 21–13 *MLT* during moderately disturbed periods (with  $Kp < 3$  or  $AL > -200$  nT) of strongly reduced plasma density, such that  $\omega_{pe}/\Omega_{ce,eq} \sim 2$ . These unusual conditions (Sheeley et al., 2001) can be encountered in the aftermath of a strong solar wind dynamic pressure pulse, as for the 2021-01-01 event in Figures 2–7 which took place 11 days after a peak of the solar wind dynamic pressure at 13.9 nPa (similar peaks preceding the events have been shown in Zhang, Artemyev, et al., 2022). For such a low plasma density, cyclotron resonance with parallel waves could easily produce 1.0–1.5-MeV microbursts, more frequent during  $Kp > 3$  periods. The results in Figure 11 therefore help to distinguish microbursts driven by very oblique or parallel whistler-mode waves. In addition, Figures 11c and 11d shows that a lower plasma trough density in the 21–08 *MLT* sector (Sheeley et al., 2001), combined with the higher occurrences of



To address the second question, let us examine in detail the electron nonlinear resonant interaction with field-aligned whistler-mode waves. This interaction combines phase bunching (energy and pitch-angle decrease) and phase trapping (energy and pitch-angle increase) (e.g., Shklyar & Matsumoto, 2009). Trapping and bunching are in a fine balance that ultimately flattens the initial phase space density gradients along the resonance curves,  $m_e c^2 \gamma - I_x \omega = \text{const}$  (e.g., Artemyev et al., 2019; Artemyev, Neishtadt, et al., 2016). Electron precipitation due to phase bunching is accompanied by acceleration of other electrons that experience phase trapping. The cyclotron phase trapping can provide an electron energy increase up to hundreds of keV, even up to a few MeVs at  $\alpha_{ra} > 60^\circ$  (e.g., Hsieh et al., 2020; Hsieh & Omura, 2017a; Omura et al., 2007). Therefore, microbursts driven

by Landau trapping should be associated with the formation of an accelerated field-aligned (<200 keV) electron population, whereas microbursts driven by cyclotron bunching should be associated with the formation of a transversely accelerated (up to MeVs) electron population. Accordingly, the actual contribution of microbursts to the depletion of the outer radiation belt ought to be investigated through detailed comparisons of simulations with observed electron flux dynamics (in different energy and pitch-angle domains) in response to nonlinear resonant interactions.

#### 4.2. On the Importance of Wave Coherence

All existing theoretical models of electron nonlinear resonant interactions with whistler-mode waves operate with highly coherent monochromatic ( $\partial\omega/\partial t = 0$ ) or weakly non-monochromatic (i.e., with  $\partial\omega/\partial t \neq 0$ , but  $\partial\omega/\partial t \ll \omega^2$ ) waves (see, e.g., Demekhov et al., 2006; Shklyar & Matsumoto, 2009; Albert et al., 2013; Omura et al., 2015; Artemyev, Neishtadt, Vainchtein, et al., 2018, and references therein). However, numerical simulations (see Zhang et al., 2021; Mourenas et al., 2022, and references therein) and spacecraft observations (e.g., Zhang, Mourenas, et al., 2020; Zhang, Agapitov, et al., 2020, and references therein) show that intense whistler-mode waves often propagate in weakly coherent, short wave-packets, with only ~5–10 wave periods within a packet. Such a strong wave amplitude modulation is typical for both quasi-parallel (e.g., Santolík, Kletzing, et al., 2014; Tsurutani et al., 2011; Tsurutani et al., 2020) and very oblique (e.g., Agapitov et al., 2014) waves. Loss of coherence and strong wavefield modulation are critically important for the nonlinear resonant interaction, because these effects destroy long-term resonances and limit trapping acceleration (Allanson et al., 2021; An et al., 2022; Artemyev, Mourenas, et al., 2015; Brinca, 1972; Gan et al., 2022; Tao et al., 2013). Therefore, an accurate quantification of nonlinear resonant interactions for electron precipitation via Landau trapping requires a further generalization of theoretical models for weakly coherent waves.

#### 4.3. On the Resonance Overlapping

In our mapping simulation, we set the time interval between two successive resonant interactions as one half of the bounce period, that is, we assume that there is only one resonance during this interval. This assumption is supported by observations of intense whistler-mode wave packets that often propagate at some distance from each other, each packet occupying a sufficient fraction of the magnetic field line to exclude the possibility of electrons reaching resonance with two similar wave packets within a half bounce period (see discussions in, e.g., Artemyev, Krasnoselskikh, et al., 2012; Zhang et al., 2018). However, this does not exclude the possibility of multiple resonances (attaining  $|n_r| \gg 1$ ) with the same wave-packet (e.g., Artemyev, Agapitov, et al., 2012; Shklyar, 1981; Shklyar, 2009). The effects from each of these cyclotron resonances can sometimes be comparable (e.g., Artemyev, Mourenas, et al., 2013; Mourenas et al., 2012). For weakly oblique whistler-mode waves, the contribution of high-order cyclotron resonances to electron diffusion is quite low (e.g., Shprits & Ni, 2009), but diffusion by very oblique whistler-mode waves can become very effective if these waves are not purely electrostatic (see discussions in Artemyev, Agapitov, et al., 2016; Li et al., 2014; Mourenas et al., 2014; Albert, 2017). Therefore, the inclusion of higher-order resonance effects into the mapping model of electron resonances with whistler-mode waves can potentially alter the rates of electron precipitation and acceleration via nonlinear Landau trapping (see discussions in Nunn & Omura, 2015; Hsieh & Omura, 2017a). For typical plasma and magnetic field conditions, these higher-order resonances do not overlap (Solovév & Shklyar, 1986) and can be considered separately, without destruction of the nonlinear Landau trapping acceleration. The resonance overlapping, the most destructive effect for nonlinear wave-particle interaction, may appear only for a sufficiently broad wave spectrum, where simultaneous waves with different frequencies would destroy long-term resonances with electrons (e.g., Karpman, 1974; Shapiro & Sagdeev, 1997). For electron interaction with whistler-mode waves, such resonance overlapping significantly reduces the efficiency of the trapped electron acceleration (Allanson et al., 2020; An et al., 2022; Gan et al., 2022; Tao et al., 2013; Zhang, Agapitov, et al., 2020). Therefore, efficient and short-lived microburst electron precipitation events should probably be attributed to electron resonant interactions with highly coherent, narrow band whistler-mode waves.

#### 4.4. On Importance of Wave Propagation Properties

The proposed model of Landau trapping acceleration is based on the assumption that very oblique whistler-mode waves can reach middle latitudes despite their obliqueness. This assumption is based on spacecraft statistics of very oblique whistler-mode wave, which indeed demonstrate the presence of these very oblique waves at middle to high latitudes (see Agapitov et al., 2012; Agapitov et al., 2013). Nevertheless, this point requires an additional discussion. The general understanding of whistler-mode wave generation at the equator and propagating along magnetic field lines suggests that only ducted field-aligned propagating waves may reach middle latitudes (e.g., Pasmanik & Trakhtengerts, 2005; Streltsov et al., 2012; Demekhov, Manninen, et al., 2017; Pasmanik & Demekhov, 2020; Ke et al., 2021; R. Chen, Gao, Lu, Chen, et al., 2021), whereas the wave normal angle of non-ducted waves rapidly increases during their propagation to higher latitudes, leading to significant deviation from the source region's magnetic field lines and strong damping of the waves before they reach middle/high latitudes (e.g., Breuillard et al., 2013; L. Chen et al., 2013; Watt et al., 2013; L. Chen et al., 2021). These simulation results, however, were obtained for waves generated with small wave normal angles at the equator. In contrast, our wave model is based on a different population of waves, observed with very oblique wave normal angles already near their equatorial source region in the presence of a reduced Landau damping due to a small plateau/beam electron population (e.g., Agapitov et al., 2016; Li, Mourenas, et al., 2016; Mourenas et al., 2015). Such very oblique waves can propagate with little attenuation up to middle/high latitudes due to their reduced Landau damping (Agapitov et al., 2013, 2018; Ma et al., 2017), and they may also propagate guided inside field-aligned ducts or micro-ducts with appropriate parameters (Streltsov et al., 2006; Streltsov & Goyal, 2021; R. Chen, Gao, Lu, Tsurutani, & Wang, 2021). Nevertheless, further investigations of the spatial scales occupied by these very oblique waves, their cross-field coherence scale, and its evolution with wave propagation (see discussion in Tsurutani et al., 2011) are required to estimate the related possible constraints on electron acceleration through Landau trapping due to cross-field motion and decoherence of very oblique waves at middle latitudes.

### 5. Conclusions

In this study, we investigated intense sub-relativistic electron precipitation (microbursts) observed by low-altitude ELFEN spacecraft in conjunction with equatorial THEMIS E spacecraft observations of very oblique whistler-mode (lower-band) chorus waves. Combining ELFEN and THEMIS E measurements with DMSP observations, and the theoretical model of electron resonant interaction with very-oblique whistler-mode waves, we have drawn the following conclusions:

1. Intense, short-lived precipitation of  $\lesssim 200$  keV electrons can be produced by electron nonlinear Landau resonance with very oblique whistler-mode waves. The associated Landau resonant trapping results in electron acceleration and transport into the loss cone: electrons with initial energies of a few keV are trapped into the Landau resonance and accelerated up to  $\lesssim 200$  keV.
2. Such precipitation of  $\lesssim 200$  keV particles cannot deplete trapped electron fluxes at the same energy, but rather indicate an enhancement of the trapped  $\lesssim 200$  keV electron population at small pitch-angles.
3. Sub-relativistic ( $\lesssim 200$  keV) precipitation events observed on low-altitude spacecraft are a good indicator of such near-equatorial electron acceleration through Landau trapping by very oblique waves.
4. Sub-relativistic ( $\lesssim 200$  keV) electron precipitation through Landau trapping by very oblique waves requires low plasma trough densities during moderately disturbed periods (e.g., in the aftermath of a solar wind dynamic pressure pulse), and could lead to higher precipitating electron energies at 04–08 MLT. In contrast, relativistic ( $\sim 500$  keV) electron precipitation through cyclotron bunching by parallel waves can occur for both low and average plasma trough densities.
5. The maximum energy of precipitating electrons is roughly 10 times smaller for Landau trapping by very oblique waves than for cyclotron bunching by parallel waves at  $L \sim 6$ . Microbursts driven by Landau trapping of very oblique waves should be more frequent during weakly disturbed periods ( $Kp < 3$ ), whereas microbursts driven by cyclotron bunching by parallel waves should be more frequent/efficient during disturbed periods ( $Kp > 3$ ).
6. The efficiency of nonlinear Landau resonant interaction, resulting in intense (microburst) precipitation, shows that the population of very oblique whistler-mode waves may significantly alter sub-relativistic electron flux dynamics and should be included into radiation belt models.



## Data Availability Statement

ELFIN data is available at <https://data.elfin.ucla.edu/>, THEMIS data is available at <http://themis.ssl.berkeley.edu>. Data analysis was done using SPEDAS V4.1 Angelopoulos et al. (2019) available at <https://spedas.org/>.

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