

# Tunable optical switches based on spin valley quantum coherence in hybrid WS<sub>2</sub>-metallic nanoantenna systems

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**Abstract:** A tunable optical switching process based on spin valley quantum coherence in a hybrid system consisting of a WS<sub>2</sub> monolayer and a metallic nanoantenna is investigated. This process is induced by plasmonically-mediated intervalley exciton-plasmon coupling which is shown to be quite resilient against ultrafast valley decoherence. © 2022 The Author(s)

## 1. Introduction

Optical and material properties of monolayers (MLs) of transition metal dichalcogenides (TMD) hold the promise advanced device applications, ranging from detectors, to light emitters, memory devices, excitonic lasers, etc [1]. Such materials have direct band gaps, high absorption coefficients, and very large spin-orbit coupling [2]. These features have made such materials unique hosts for investigation of coupling of valley excitons with localized surface plasmons resonances (LSPRs) of metallic nanoantennas (mANTs), supporting Rabi splitting, dressed states, and polaritons [3]. TMD-mANT systems can also lead to gain without inversion [4], and coherent transport of energy. [5] In this contribution we study application of such systems for coherent optical switches that can be tuned by varying the intensity of a laser field. Such switches, which are based on intervalley quantum coherence mediated by spin valley exciton-plasmon coupling, offer significant resiliency against valley decoherence.

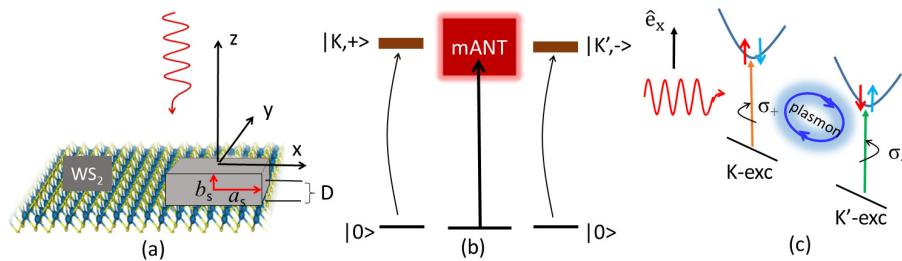


Fig. 1. (a) Schematic of the WS<sub>2</sub>-mANT system interacting with a laser field polarized along the x-axis. (b) Electronic structures of spin polarized valley K-exc and K'-exc in the presence of mANT. (c) Schematic of the plasmonically-mediated intervalley coupling.

## 2. Review of theory

We consider an Ag mANT with semimajor and semiminor axes of  $a_s$  and  $b_s$  placed on the surface of a WS<sub>2</sub> ML with the center-to-center distance of  $D$  (Fig.1a). The materials surrounding these are assumed to have dielectric constant  $\epsilon_0$ . The WS<sub>2</sub>-mANT system interacts with a laser field linearly polarized along the x-axis  $\mathbf{E} = E_0(t) \cos(\omega t) \hat{\mathbf{e}}_x$ . Here  $\omega$  is the frequency of the laser,  $E_0(t)$  is its time-dependent amplitude, and  $\hat{\mathbf{e}}_x$  is the polarization unit vector along the x-axis. Under these conditions the interaction Hamiltonian of the excitons associated with K valley (K-exc) and K' valley (K'-exc) can have the following form:

$$\hat{H}_{int}^{K,(K')} = -\frac{1}{\sqrt{2}}(\mu_+ E_K a_+^\dagger a_0 + \mu_- E_{K'} b_-^\dagger b_0) + H.C. \quad (1)$$

Here  $a_+^\dagger$  ( $b_-^\dagger$ ) and  $a_0$  ( $b_0$ ) are creator and annihilation operators for K-exc (K'-exc) and ground states.  $\mu_+$  and  $\mu_-$  are dipole moments associated the K-exc and K'-exc transitions (Fig.1b).  $E_K$  and  $E_{K'}$  refer to the field experienced

by K-exc and K'-exc, respectively [5]. The density matrix equations for K and K' valleys are given by:

$$\frac{d\rho^{K(K')}}{dt} = -i[H_0 + H_{int}^{K,(K')}, \rho^{K(K')}(t)] + [\mathcal{L}\rho^{K(K')}]_{damp}. \quad (2)$$

Here  $[\mathcal{L}\rho^{K(K')}]_{damp}$  refers to the damping terms which contain pure dephasing rate,  $\gamma_p^{K(K')}$ .  $H_0$  is the Hamiltonian of the WS<sub>2</sub> ML in the absence of the laser field and the mANT. Solving these equation under quasi-steady state condition allows one to find the density matrix elements and coherently normalized plasmonic field enhancement factor given by  $P_{coh}^{K(K')} = |E_{K,(K')}|^2/|E_0|^2$ .

### 3. Tunable coherent optical switching

For simulation we consider  $\mu_- = \mu_+ = 1.12 \text{ eV nm}$ , the linewidths of K-exc and K'-exc equal to 20 meV,  $\omega_K = 2.006 \text{ eV}$ , and the radiative lifetime of excitons equal to 1.5 ns [4]. Additionally, for the Ag mANT we assume  $a_s = 9$  and  $b_s = 4 \text{ nm}$ , and  $\epsilon_0 = 3.61$  (SiN). We also set  $D = 6 \text{ nm}$  and the dielectric constant of the WS<sub>2</sub> ML,  $\epsilon_{ML}$  is  $\sim 13.7$ . Fig.2a shows the results for  $P_{coh}^K$  as a function of detuning of the laser field from the K-exc transition ( $\Delta_{+0}$ ) for different values of the incident laser intensity (legends in  $\text{W/cm}^2$ ). The results suggest sudden changes (arrow) associated with switching between two collective states of the system. Figs.2b-2d show variation of the K-exc population, and real and imaginary parts of valley polarization ( $\rho_{+,0}^K$ ), respectively. Note that as the laser intensity changes the position of transition associated with the switching is shifted, suggesting its tunability.

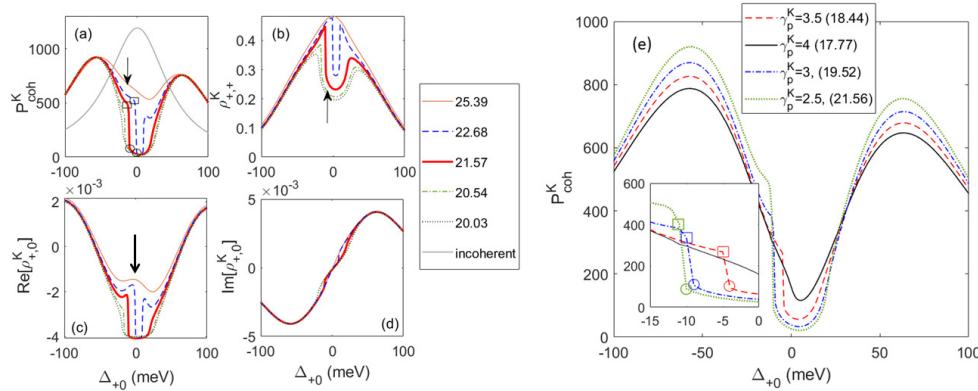


Fig. 2. Variation  $P_{coh}^K$  (a), K-exc population (b),  $\text{Re}[\rho_{+,0}^K]$  (c) and  $\text{Im}[\rho_{+,0}^K]$  (d) as a function  $\Delta_{+0}$  for various laser intensity (legends in  $\text{W/cm}^2$ ). (e) shows  $P_{coh}^K$  as a function of  $\Delta_{+0}$  for  $\gamma_p^K = 4$  (solid line), 3.5 (dashed line), 3 (dashed-dotted line), 2.5  $\text{ps}^{-1}$  (dotted line). Numbers in parentheses in the legend of (e) refer to the values of laser intensity in  $\text{W/cm}^2$ .

Fig.2e shows variations of  $P_{coh}^K$  for various values of valley pure dephasing rate ( $\gamma_p^K$ ). This parameter includes contributions of various processes that can lead to polarization decoherence. The results show that the switching process can occur even when  $\gamma_p^K = 4 \text{ ps}^{-1}$ , indicating a good resiliency against valley decoherence.

### 4. Conclusions

The potential of TMD-mANT systems for coherent devices based on spin valley quantum coherence was studied. Such systems can sustain coherent exciton-plasmon coupling under ultrafast polarization dephasing (300 fs).

### References

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