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3	Vertical Motions in Orographic Cloud Systems over the Payette River Basin.
4	Part 1: Recovery of Vertical Motions and their Uncertainty from Airborne
5	Doppler Radial Velocity Measurements
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ABSTRACT

34 Vertical motions over the complex terrain of Idaho's Payette River Basin were observed 35 by the Wyoming Cloud Radar (WCR) during 23 flights of the Wyoming King Air during the 36 SNOWIE field campaign. The WCR measured radial velocity, V_r , which includes the reflectivityweighted terminal velocity of hydrometeors (V_t) , vertical air velocity (w), horizontal wind 37 38 contributions as a result of aircraft attitude deviations, and aircraft motion. Aircraft motion was 39 removed through standard processing. To retrieve vertical radial velocity (W), V_r was corrected 40 using rawinsonde data and aircraft attitude measurements. w was then calculated by subtracting the mean W, (\overline{W}) , at a given height along a flight leg long enough for \overline{W} to equal the mean 41 reflectivity weighted terminal velocity, $\overline{V_t}$, at that height. 42

The accuracy of the w and \overline{V}_t retrievals were dependent on satisfying assumptions along a 43 44 given flight leg that the winds at a given altitude above/below the aircraft did not vary, the vertical air motions at a given altitude sum to 0 m s⁻¹, and $\overline{V_t}$ at a given altitude did not vary. The uncertainty 45 in the w retrieval associated with each assumption is evaluated. Case studies and a project wide 46 summary show that this methodology can provide estimates of w that closely match gust probe 47 48 measurements of w at the aircraft level. Flight legs with little variation in equivalent reflectivity 49 factor at a given height and large horizontal echo extent were associated with the least retrieval 50 uncertainty. The greatest uncertainty occurred in regions with isolated convective turrets or at 51 altitudes where split cloud layers were present.

52 **1. Introduction**

53 Vertical air motion is a key variable in atmospheric dynamics and cloud microphysics studies. 54 This variable is difficult to measure, and Doppler profiling radar estimates are further challenged 55 by the "contamination" of vertical air motion by the fall velocity of the radar scatters (i.e., 56 hydrometeors). Separately, there is much interest in the estimation of terminal velocity of 57 hydrometeors, as it carries information about size and riming fraction.

58 A vertically-pointing Doppler radar can provide direct measurements of vertical particle 59 motion through the depth of a cloud. For a truly vertically-pointing Doppler radar, the Doppler 60 vertical radial velocity, W, is the sum of the reflectivity-weighted terminal velocity (V_t) and the vertical air velocity (w). Approaches to retrieve w from W with ground-based radars have focused 61 62 on cold clouds and have employed different methodologies to estimate and remove V_t of ice from W (see review by Protat and Williams, 2011). One approach involves calculating the mean value 63 of W. \overline{W} , at a given altitude, over a sufficient period of time or distance, with the assumption that 64 the expected magnitude and number of updrafts and downdrafts are approximately equal so that 65 the mean value of $w, \overline{w} = 0$, and the mean value of W, \overline{W} , is equal to the mean of $V_t(\overline{V}_t)$. \overline{V}_t can 66 then be subtracted from individual values of W to retrieve estimates of w (Delanoë et al. 2007). A 67 68 second approach, applicable to hydrometeors whose fall speed is strongly dependent on size (such 69 as rain or hail), involves binning W measurements based on reflectivity. Provided that the number 70 of data points within each reflectivity bin is sufficiently large such that the number and magnitude of updrafts and downdrafts within each bin are equal, $\overline{W} = \overline{V}_t$ in each reflectivity bin and \overline{V}_t can 71 be subtracted from W based on a unique value of $\overline{V_t}$ for each reflectivity value (e.g. Orr and Kropfli 72 1999; Protat et al. 2003; Delanoë et al. 2007). A third approach, applicable in ice clouds, estimates 73 74 $\overline{V_t}$ based on a relationship between particle fall speeds and maximum dimensions of ice particles integrated over an observed or assumed ice particle size distribution (e.g., Mitchell 1996; 75 Heymsfield and Iaquinta 2000), after which \overline{V}_t is subtracted from individual measurements of W 76 77 to obtain w (Babb et al. 1999; Deng and Mace 2006). A fourth approach, applicable to liquid 78 clouds, uses Doppler spectrum and a Mie notch technique in order to retrieve w (Kollias et al. 79 2002). A fifth approach, used in airborne studies, applicable to unrimed snow, is to simply assume a constant value for V_t (e.g., Grasmick and Geerts 2020), since the reflectivity-weighted terminal 80 velocity of unrimed ice particles is often between 0.5-1.2 m s⁻¹ over much of the cloud depth in ice 81 cloud environments (e.g., Rosenow et al. 2014). In airborne radar studies, this constant value $\overline{V_t}$ 82

can be estimated as the difference between the flight-leg-mean values of *W* (obtained from radar
above and below flight level) and *w* (obtained from a gust probe) (Grasmick and Geerts 2020).

The problem of using these methods to retrieve w from an airborne platform is more complicated because of the motion of the aircraft on which the radar is mounted. When a nominally vertically-pointing beam is actually oriented in a direction other than nadir or zenith and is not orthogonal to the aircraft velocity vector, the horizontal winds and the aircraft motion will contribute to the measured radial velocity (V_r).

90 Correcting V_r for platform motion requires the radar antenna beam-pointing vector and a 91 coordinate transformation between the aircraft body fixed-reference frame and the ground-fixed 92 reference frame that considers the three-dimensional aircraft motion vector. The transformation 93 involves consideration of the pitch, roll, and yaw angles, and aircraft ground speed (Haimov and 94 Rodi 2013). However, after correction for aircraft motion, residual biases in V_r remain due to the 95 contribution of horizontal winds to V_r when the beam is not pointing at nadir or zenith. For example, Fig. 1 shows V_r , corrected for platform motion contribution, during an eastbound and 96 97 subsequent westbound flight leg over the Payette River Basin by the University of Wyoming King 98 Air (UWKA) Wyoming Cloud Radar (WCR; Pazmany et al. 1994; Wang et al. 2012) during the 99 Seeded and Natural Orographic Wintertime Clouds: the Idaho Experiment (SNOWIE; Tessendorf 100 et al. 2019). V_r in these figures is not corrected for horizontal wind contributions. These flight legs 101 (Fig. 1) illustrate an example of a deep stratiform cloud with weak echo near cloud top and possible 102 cloud top entrainment occurring, weak boundary layer turbulence, and a melting level at ~ 2.8 km 103 decreasing in altitude along the eastern end of both flight legs. The data show clear biases in V_r on 104 consecutive legs (positive eastbound and negative westbound) that resulted from horizontal winds 105 being projected into the beam.

Heymsfield (1989) was the first to develop a transformation matrix to retrieve W by removing contributions from the horizontal wind and aircraft motion components of V_r . He tested his retrieval using an idealized vertical profile of the horizontal wind field to estimate retrieval uncertainty. If the beam-pointing vector is known in ground relative coordinates, the contribution to single Doppler V_r by the horizontal wind, as measured by a rawinsonde, can be added or subtracted from individual range gates given the beam pointing direction (Geerts and Miao 2009). Both of these approaches assumed horizontal homogeneity of the horizontal winds. For airborne radars capable 113 of providing multi-Doppler measurements, W can be retrieved via two- or three-dimensional



114 Doppler velocity synthesis (Leon et al. 2006; Damiani and Haimov 2006; Hagen et al. 2021).

116 Fig. 1: V_r from IOP 23 during a consecutive east/west flight leg pair over the Payette River Basin. a) is the eastbound flight leg from 9 March 2017 22:24:00 to 22:35:42 UTC, and c) is the 117 westbound flight leg from 9 March 2017 22:39:21 to 22:57:07 UTC. b and d are contoured 118 119 frequency by altitude diagrams of V_r for the eastbound and westbound flight legs respectively, 120 binned every 0.1 m s⁻¹ every 100 m in altitude. The black vertical line on panels b and d denotes V_r of -1 m s⁻¹. In panels a and c, the dashed line is the altitude of the aircraft, and the white area 121 below is the terrain. 122

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124 Miao et al. (2006) commented on the horizontal wind contamination of the radial velocity from 125 a non-vertically pointing airborne radar and the limited possibility of using a nearby sounding for 126 correction, but did not use it. In order to remove the horizontal wind contribution using a 127 rawinsonde, the assumption must be made that the flight legs occur over a short enough time and distance that the winds at a given altitude above/below the aircraft do not vary horizontally or 128 129 change with time. Geerts and Miao (2009) applied this approach using rawinsonde winds with the WCR to retrieve W. However, their study did not include an uncertainty analysis regarding 130 131 retrieved W and w. Geerts et al. (2011) and following studies (e.g. Bergmaier and Geerts 2016; 132 2017; 2020; Chu et al. 2017; Grasmick and Geerts 2020) have used horizontal wind profiles from 133 rawinsondes in order to estimate the aircraft cross-track wind component and reduce error when 134 performing dual-Doppler retrievals (Damiani and Haimov 2006). Bergmaier and Geerts (2017) 135 noted that the use of other soundings (with different wind profiles) taken near the same time did 136 not produce any significant differences in the recovered velocity field since aircraft attitude 137 changes were small. French et al. (2015) used dual-Doppler synthesis (Damiani and Haimov, 138 2006) and instantaneous flight-level winds to retrieve the vertical plane two-dimensional velocity 139 field. Pokharel et al. (2017) corrected nadir and zenith beams using a rawinsonde and noted that 140 when the wind profile changes dramatically along a flight leg, there is higher uncertainty in w 141 estimation especially when aircraft attitude changes were large. All of the studies quoted above use winds from a sounding profile to correct V_r "contaminated" by the horizontal wind to retrieve 142 143 W but provide limited documentation and assessment of the effect of the rawinsonde correction 144 algorithm on retrieval of w, its assumptions, and uncertainties.

An error in estimated W will also be present if there is an error in the antenna beam-pointing vector. The WCR beam-pointing vector, used herein, has been calibrated following Haimov and Rodi (2013). The maximum root mean square error in the calibrated beam-pointing angle is less than 0.1° resulting in less than 0.15 m s⁻¹ error due to residual aircraft motion after removing the aircraft motion contribution.

To explore updraft retrievals under a variety of atmospheric conditions in complex terrain, we use herein aircraft observations from UWKA flown during SNOWIE. During the campaign, 23 research flights sampled orographic cloud systems over the Payette River Basin of Idaho. During flights the WCR made measurements of V_r within orographic cloud systems at high resolution (see Sec. 2a). During each research flight the UWKA flew along fixed tracks over the Payette River Basin parallel to mid-level (~700 hPa) flow (Fig. 2). One flight leg was typically completed in 10-20 minutes with 4-hour flights typically completing a total of 10-14 flight legs.

157 Flight legs sampled a variety of fixed (tied to the orography) and transient (related to vertical 158 wind shear and conditional instability within passing weather systems) updrafts, providing a large dataset where w and \overline{V}_t retrievals were possible using a rawinsonde wind correction after applying 159 a series of assumptions regarding the horizontal wind and V_t . These assumptions include that i) 160 161 the legs occur over a short enough time and distance that the horizontal winds at a given altitude 162 above/below the aircraft did not vary horizontally or change with time, ii) the legs are also long 163 enough that the magnitudes of the updrafts and downdrafts along a flight leg at any given altitude sum to 0 m s⁻¹, and iii) that the $\overline{V_t}$ did not vary substantially along the flight legs at a given altitude. 164 165 The purpose of this study, Part 1, is to introduce and test the validity of these assumptions, and to

evaluate the retrieval of w and $\overline{V_t}$ from the WCR using data from SNOWIE. In Part 2 (Zaremba et al. 2022), we quantify the magnitude, and associated uncertainties, of fixed and transient updraft structures over the Payette River Basin during SNOWIE and relate those updraft structures to the thermodynamic environments present during the project. In Part 3 (Heimes et al. 2022), we examine the impact of fixed and transient updrafts on trajectories of ice particles created by seeding clouds in both measured and simulated updraft fields over the Payette River Basin during SNOWIE cloud seeding operations.

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Fig. 2: Domain of SNOWIE (outlined in black) in Idaho. Terrain elevation in meters above
mean sea level is contoured. Plotted in yellow are the three flight tracks flown during SNOWIE.
Rawinsondes were launched by Idaho Power Company (IPC) at Crouch, Idaho denoted by a yellow
circle.

This paper is organized as follows: Sec. 2 presents the data used in this analysis. Sec. 3 provides an overview of the retrieval methodology and assumptions required to retrieve w and $\overline{V_t}$. Sec. 4 evaluates retrieval uncertainty including a summary of uncertainty estimates using data from all research flights. Sec. 5 presents examples of retrieved w and $\overline{V_t}$ along flight legs. Sec. 6 184 examines conditions which can result in higher retrieval uncertainty and Sec. 7 quantifies retrieval
185 uncertainties for all flight legs during SNOWIE. Key conclusions are discussed in Sec. 8.

186 **2. Data**

187 a) UWKA Wyoming Cloud Radar data

188 The WCR is a 95 GHz, 3 mm wavelength, pulsed Doppler cloud radar that was flown on the 189 UWKA during SNOWIE. Data used herein are from the WCR fixed antennas nominally pointed 190 at zenith and nadir during straight, level flight. In this configuration, the WCR measured the equivalent reflectivity factor (Z_e) and V_r . The WCR reflectivity is calibrated by measuring the 191 192 return from a trihedral corner reflector with a known scattering cross-section. Error associated with 193 this calibration is estimated to be less than 2.5 dB at any distance away from the radar flight level 194 (Wendisch and Brenquier 2013, Ch. 9.5.5, pp. 509-517; Grasmick et al. 2022). The minimum 195 detectable signal was \sim -40 dBZ_e at 1 km distance away from the radar and \sim -26 dBZ_e at a distance 196 of 5 km. Data were sampled at 30 m in range along the radar beam and 4.5-7.5 m along the flight 197 track depending on the speed of the UWKA. Only straight, level flight legs were used in this 198 analysis.

199 The measured values of Z_e and V_r can be negatively impacted by attenuation, particularly in 200 water clouds. Attenuation, under these conditions, can reduce the signal strength to the point of 201 low signal to noise ratio. This would result in reduced Z_e and higher variance in V_r estimates. 202 During SNOWIE, radar echoes were almost entirely due to ice. On a few flights, a very low-level 203 melting level was present near the terrain. Protat et al. (2019) show that the two-way attenuation coefficient produced by stratiform ice particles at W-band ranges between 1 and 1.6 dB km⁻¹ for 204 205 W-band reflectivity values between 13 and 18 dBZe with an increase in attenuation with 206 reflectivity. At such high reflectivity values, W-band scattering largely falls in the Mie regime, i.e. 207 reflectivity values for weather (cm-wave) radars are much higher. Figure 3a shows a Contoured 208 Frequency by Altitude Diagram (CFAD) of WCR Z_e for the entire SNOWIE campaign (all 238 flight legs). Less than 1% of the measurements exceeded 19 dBZe and less than 5% exceeded 12 209 210 dBZe. W-band power is also attenuated by cloud droplets: the two-way path-integrated attenuation rate is about 10 dB per km per g kg⁻¹ of liquid water (Liebe et al. 1989; Vali and Haimov 2001). 211 212 A microwave radiometer was located at Horseshoe, Idaho (just southwest of Crouch, Idaho; see 213 Fig. 2). The radiometer provided measurements of vertical liquid water path every six minutes 214 during SNOWIE on station flight times (Table 1). Data were not available for the March flights

215 and were masked when liquid water was present on the radiometer dome. 90% of the values were 216 below 1.03 mm (Fig. 3b). The aircraft flew at approximately 4 km altitude during SNOWIE, with 217 about 2.5 km of cloud beneath the aircraft, and 1-6 km of cloud above the aircraft. The two-way 218 attenuation rate for different cloud depths as a function of radiometric liquid water measurements 219 appears in Fig. 3. If we assume that the cloud water was concentrated below the aircraft at warmer temperatures, the W-band two-way attenuation was less than 3.3 dB km⁻¹ for 90% of the time the 220 221 aircraft was flying. In this paper, horizontal variations in Z_e at a given level are used to quantitatively estimate uncertainties associated with retrieval of w and \overline{V}_t . Vertical variations in 222 Z_e , where attenuation would be more prevalent, are not considered in this analysis. Since 223 224 attenuation would lead to greater variation in Z_e at a given level, the effect of the attenuation would 225 be to increase the calculated uncertainty in the retrieval of w.



226 227 Fig. 3: a: CFAD of Z_e for all 238 SNOWIE research flight legs. The CFAD is binned every 100 228 m in altitude and every 1 dBZe. The frequency is normalized to 100% at each altitude bin. The 50th, 95th, and 99th quantiles are overlaid in white and labeled. b: (solid black line) distribution of 229 230 vertically integrated liquid water path from a radiometer at Horseshoe, ID during SNOWIE aircraft 231 on station sampling times in Table 1. Red lines represent the two-way path integrated attenuation for different cloud depths. The black dashed line represents the 90th percentile of the radiometer 232 233 measurements. 234

235 b) Rawinsonde data

236 Special rawinsondes were launched over the Payette River Basin to analyze the 237 thermodynamic properties and wind fields within the orographic cloud systems. The Idaho Power 238 Company (IPC) launched Lockheed Martin LMS6 rawinsondes from Crouch, Idaho (Fig. 2) at 239 regular intervals during research flights, typically every 2-3 hours. In this paper, we limit the use 240 of the rawinsondes to measurement of the winds. The manufacturer-stated accuracy of IPC rawinsondes was ± 0.2 m s⁻¹ for wind speed. Rawinsonde data collected during SNOWIE typically 241 242 had an average vertical resolution of 4 m with the sondes recording data every second, ascending 243 at \sim 5-6 m s⁻¹, and drifting an average of 12.4 km away from their launch location upon reaching 244 cloud top.

245 c) UWKA gust probe data

246 During SNOWIE, horizontal winds and vertical velocity from a gust probe mounted on a 247 nose boom on the UWKA were recorded at a rate of 1 Hz. Aircraft motion was removed from the 248 gust probe raw winds to retrieve w_{gp} using aircraft speed and acceleration obtained from the Inertial 249 Navigation System and aircraft attitude parameters. Because gust probe velocity components result 250 from the integration of accelerations, their variations are known more accurately than their long-251 track mean. Therefore, it is customary to remove the long-track mean vertical velocity. After 252 removal of leg averages, the resulting gust probe vertical velocity (w_{gp}) is accurate to at least 0.1 m s⁻¹ (Lenschow 1972; Geerts and Miao 2005). This procedure was followed in this analysis. w_{gp} 253 254 was linearly interpolated to match the sampling rate of the WCR.

255 **3.** Retrieval of vertical air motion and mean terminal fall velocity

256 The goal of this paper is to retrieve vertical profiles of vertical air motion, w, along flight 257 legs, and a profile of reflectivity-weighted mean terminal velocity averaged along a flight leg, $\overline{V_t}$, 258 from measurements of radial velocity, V_r , by an aircraft with radar beams nominally pointing at 259 nadir or zenith, and estimate associated uncertainties. The retrieval technique first involves using 260 rawinsonde-measured winds to retrieve W, the vertical hydrometeor velocity at a range gate, by 261 removing contributions of the aircraft motion and the horizontal wind to V_r due to time-dependent 262 variations in the beam-pointing vector. These result from small fluctuations in pitch, roll, and yaw 263 of the aircraft due to flight-level turbulence or pilot adjustments while flying along straight flight 264 legs.

The value of W at each range gate is retrieved using a horizontal wind profile derived from a rawinsonde, the ground-relative 3D aircraft motion vector, and the beam-pointing vector by performing a coordinate transform from the fixed aircraft reference frame to a ground reference frame. The transformation matrix (Haimov and Rodi, 2013) and method of retrieval of W from V_r are given in Appendix A.

Values of *W* were retrieved along each beam between the ground and cloud echo top except in a 250 m zone centered at flight level. *W* was then regridded to a common grid of 30 m vertical range referenced to mean sea level (all altitudes in this paper are with respect to mean sea level (MSL)). Range gates were resampled to the new grid using a nearest neighbor approach to estimate 274 $\overline{V_t}$ at a given height and preserve the original *W* values. There were *m* altitudes with valid 275 measurements between the aircraft and ground (or cloud echo top), and *n* grid points along a given 276 flight leg, so each flight leg had an *m* × *n* grid of retrieved *W*.

- 277 In order to retrieve w and $\overline{V_t}$ from W, three assumptions had to be made:
- 278 1. The horizontal wind, V_H , is invariant at a given height, *m*, along a flight leg, and 279 can be represented by the winds measured by a nearby rawinsonde.
- 280 2. The flight leg was sufficiently long such that at each m, $\overline{w_m} = \sum_{1}^{n} w_{m,n} = 0 \text{ m s}^{-1}$. 281 Simply stated, this assumption is that the sum of the magnitudes of the updrafts and 282 downdrafts along the track at any altitude average to zero.
- 283 3. The reflectivity-weighted hydrometeor terminal velocity does not vary significantly 284 along the leg at a given altitude, so that, at any point along the track, $V_{t,m} \cong \overline{V_{t,m}} =$ 285 $\sum_{1}^{n} V_{t,(m,n)}/n$.

 $\overline{V_{t\,m}} = \overline{W_m}$

286 Applying these assumptions:

- 287
- 288 and:
- 289

- $w_{(n,m)} = W_{(n,m)} \overline{V_{t,m}}$
- In this manner, *w* can be retrieved for all $m \times n$ grid points on a flight leg and a \overline{V}_t profile for that flight leg can be obtained. In the remainder of this paper, we test the validity of, and estimate the uncertainty associated with the three assumptions stated above and examine example retrievals of *w* and \overline{V}_t in orographic clouds over the Payette River Basin.
- 294 4. Retrieval uncertainty

In this section, we evaluate the uncertainty in the estimates of *w* arising from violations of the assumptions stated in the previous section.

a) Assumption 1: The horizontal wind, $\overrightarrow{V_H}$, *is invariant at a given height, m, along a flight leg, and can be represented by the winds measured by a nearby rawinsonde.*

During SNOWIE flights, project rawinsondes were launched at Crouch, Idaho (see Sec. 2c), a site close to the flight track. The rawinsonde used to retrieve *w* for each flight can be found in Table 1. The average aircraft on-station sampling time was 2.5-3.5 hours with the rawinsonde typically launched during or near the sampling period. A difficulty with quantifying uncertainty associated with this assumption was that there were no measurements of the along-track variation 304 of the horizontal wind above or below the aircraft flight level. The UWKA did however, measure 305 the zonal, u, and meridional, v, components of the wind at flight level (m_{ac}) along each flight leg. 306 These could be directly compared to sounding measurements. Flight level measurements will be 307 used to provide a best available estimate of uncertainty due to differences in measured wind along 308 a given flight leg between the aircraft and rawinsonde.

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IOP	Rawinsonde Launch Time	On-Station Flight Time
1	8 January 2017 04 UTC	0246 UTC – 0540 UTC
2	9 January 2017 05 UTC	0435 UTC – 0727 UTC
3	11 January 2017 03 UTC	0247 UTC – 0528 UTC
4	18 January 2017 22 UTC	2017 UTC – 2351 UTC
5	19 January 2017 16 UTC	1541 UTC – 1822 UTC
6	20 January 2017 00 UTC	2245 UTC – 0152 UTC
7	21 January 2017 21 UTC	2249 UTC – 0128 UTC
8	22 January 2017 18 UTC	2118 UTC – 0020 UTC
9	31 January 2017 16 UTC	2021 UTC – 2201 UTC
10	3 February 2017 21 UTC	2006 UTC – 2125 UTC
11	4 February 2017 23 UTC	2209 UTC – 0107 UTC
12	7 February 2017 22 UTC	2013 UTC – 2303 UTC
13	16 February 2017 20 UTC	2343 UTC – 0041 UTC
14	18 February 2017 20 UTC	2139 UTC – 0041 UTC
15	19 February 2017 15 UTC	1744 UTC – 2046 UTC
16	20 February 2017 15 UTC	1447 UTC – 1721 UTC
17	21 February 2017 16 UTC	1443 UTC – 1807 UTC
19	4 March 2017 14 UTC	1330 UTC – 1642 UTC
20	5 March 2017 13 UTC	1213 UTC – 1444 UTC
21	7 March 2017 14 UTC	1423 UTC – 1736 UTC
22	9 March 2017 16 UTC	1422 UTC – 1646 UTC
23	9 March 2017 23 UTC	2019 UTC – 2332 UTC
24	16 March 2017 03 UTC	0108 UTC - 0420 UTC

310

311 Table 1: Rawinsonde launched at Crouch used to retrieve *w* during each SNOWIE Intensive312 Operation Period (IOP) research flight.

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The only measure of the variability of the horizontal winds along the cross section were made by the aircraft at flight level. Flight level winds were compared to those measured by the rawinsonde at m_{ac} . The average difference and standard deviation between wind speed measured at flight level and wind speed measured by the rawinsonde for the entire SNOWIE field campaign was $u = -0.7 \pm 4.2$ m s⁻¹ and $v = 0.0 \pm 3.7$ m s⁻¹ (Fig. 4a,b). An uncertainty estimate (Δw) using 319 Appendix A Eq. 4 was calculated for each beam along a given flight leg using the difference in 320 measured wind speeds between the sounding and gust probe to calculate the difference in w. In 321 order to match rawinsonde winds with those measured at the aircraft flight level, sounding data 322 were linearly interpolated to the nearest 0.1 m. Δw was calculated for both the nadir and zenith 323 beams.

324 To estimate uncertainty along a given flight leg associated with assumption 1 the standard deviation of Δw ($\sigma_{w,1}$) for all nadir and zenith beams along a given leg was calculated. Figure 4c,d 325 shows the distribution of Δw for all beams during the SNOWIE field campaign associated with 326 violations of assumption 1. Δw had a standard deviation of 0.15 m s⁻¹ for all zenith beams and 327 0.14 m s⁻¹ for the nadir beams. $\sigma_{w,1}$ had a median of 0.06 m s⁻¹ for all flight legs for both nadir 328 and zenith beams (Fig. 4e,f). For most flight legs the uncertainty associated with $\sigma_{w,1}$ was less 329 than 0.1 m s⁻¹ (Fig. 4g). Intensive operation period (IOP) 13 may have had a greater uncertainty 330 because the rawinsonde used to retrieve w was launched ~ 4 hours before the research flight, the 331 largest difference for any IOP. In Sec. 5 we quantify $\sigma_{w,1}$ for specific flight legs during SNOWIE. 332



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Fig. 4: (a) Δu : Difference in the *u* component of the wind speed measured by the aircraft at m_{ac} and the *u* component of the wind speed measured by the rawinsonde used to retrieve *w*. (b) Δv : Same as (a) except for the *v* component of the wind. (c) Δw for all zenith beams calculated using the difference in aircraft measured and sounding measured wind speed. (d) same as (c) except for the nadir beam. (e) $\sigma_{w,1}$ for zenith beams for all flight legs during SNOWIE. (f) Same as e except for nadir beams. (g) boxplot of $\sigma_{w,1}$ for nadir and zenith beams for all research flights during SNOWIE.

341 b) Assumption 2:
$$\overline{w_m} = \sum_{1}^{n} w_{m,n} = 0 \text{ m s}^{-1} \text{ along a flight leg}$$

To estimate horizontal cloud extent needed at a given grid level *m* for $\overline{w_m} = 0$ m s⁻¹, the great circle distance between each grid column *n* on a given flight leg was calculated based on the latitude and longitude of the UWKA. Each flight leg was then broken up into units 2-120 km in length (in 2 km increments). For example, a flight leg 100 km in length was broken up into fifty 2 km units, twenty-five 4 km units, sixteen 6 km units, etc. Units could only be as long as the 347 flight leg's maximum length. Only grid levels where radar echo was present across the entire flight 348 leg were used in the uncertainty analysis. \overline{w} was calculated at each *m*, along each unit. Figure 5 349 shows \overline{w} for all segment lengths during SNOWIE normalized as a percentage of a given segment length (binned every 0.05 m s⁻¹ and 2 km). The standard deviation of each segment length was 350 then calculated and used to estimate the standard deviation of w associated with assumption 2 351 352 $(\sigma_{w,2})$ (see Fig. 5). The analysis shows that if horizontal echo extent was 2 km at a given m, $\sigma_{w,2}$ would be 0.46 m s⁻¹, while if the echo extent was 80 km at a given *m*, $\sigma_{w,2}$ would be 0.03 m 353 s⁻¹. $\sigma_{w,2}$ approaches 0 m s⁻¹ at leg lengths >80 km. The source of the increase in $\sigma_{w,2}$ beyond 80 354 355 km is uncertain, but may be related to the broader effect of ascent across the entire mountain massif 356 of Idaho.

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Fig. 5: \overline{w} found using methodology presented in Sec. 4 for different segment lengths (2 km intervals). a: $\sigma_{w,2}$ for all segment lengths. b: distribution of \overline{w} for each segment normalized as a percentage of all segments with a given segment length (binned every 0.05 m s⁻¹ and 2 km).

362 c) Assumption 3: $V_{t,m} = \overline{V_{t,m}} = \sum_{1}^{n} V_{t,(m,n)}/n$ along a flight leg

363 Periods with large Z_e variation at a given level (m) would be expected to have large variation in $\overline{V_{t,m}}$. In order to estimate the standard deviation of w associated with assumption 3 ($\sigma_{w,3}$), the 364 365 retrieved w at the first valid range gate above and below the UWKA were averaged and compared to vertical velocity measured by the gust probe at flight level (w_{gp}) . w_{gp} was linearly interpolated 366 to match the sampling rate of the WCR. The standard deviation of $|w - w_{gp}|$, $(\sigma_{|w-w_{gp}|})$, along 367 a given flight leg was compared to the standard deviation of Z_{e} , (using logarithmic units), $\sigma_{Z_{e}}$, 368 369 directly above and below the aircraft, also along a given flight leg, using the same range gates as 370 radar-retrieved w (Fig. 6). $\sigma_{|w-w_{gp}|}$ was then estimated using a least squares line of best fit where

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$$\sigma_{w,3} = \sigma_{|w-w_{qp}|} = 0.016 * \sigma_{Z_e} + 0.126$$

372 At a given level *m*, σ_{Z_e} was calculated and the relationship above was used to estimate $\sigma_{w,3}$. For

example, σ_{Z_e} of 15 dBZ_e would have a $\sigma_{w,3}$ of 0.37 m s⁻¹, and a σ_{Z_e} of 1 dBZ_e would have a $\sigma_{w,3}$

$$3'/4$$
 of 0.14 m s⁻¹



Fig. 6: A comparison of $\sigma_{|w-w_{gp}|}$ and σ_{Z_e} for all flight legs during SNOWIE. The black line represents the best fit line which is $\sigma_{w,3}$ or $\sigma_{|w-w_{gp}|} = 0.016 * \sigma_{Z_e} + 0.126$.

- 378 *d)* Total uncertainty of assumptions 1-3
- 379 The total retrieval uncertainty (σ_T) at a given *m* was estimated as:

380
$$\sigma_T = \sqrt{\sigma_{w,1}^2 + \sigma_{w,2}^2 + \sigma_{w,3}^2}$$

381 This assumes that the three errors are independent. Sec. 5 and 6 show examples of *w* retrieval

and σ_T from specific flight legs during SNOWIE research flights. A summary of uncertainty for

all SNOWIE research flights is presented in Sec. 7.

5. Retrieval examples and uncertainty estimates

385 *a)* IOP 23 9 March 2017

386 IOP 23 sampled a deep stratiform cloud within southwest flow over the Payette River Basin. Figure 7a shows an eastbound flight leg (22:24:40 to 22:35:35 UTC) over the terrain and reveals 387 that V_r was ~1 m s⁻¹ near cloud echo top. A Contour Frequency by Altitude diagram (CFAD) in 388 Fig. 7b show V_r decreasing with depth beneath cloud echo top. Applying the retrieval methodology 389 390 in Sec. 3, the retrieved w in Fig. 7c reveals weak orographic ascent extending through the cloud 391 echo top on the windward (west) sides of mountains on the flight leg, and downslope flow on the 392 leeward sides, with updrafts and downdrafts on the order of ± 0.5 m s⁻¹ (Figs. 7d). Turbulence, 393 and Z_e approaching the minimal detectable signal, is influencing V_r near cloud echo top. Turbulence in the boundary layer is also evident, with the CFAD in Fig. 7d showing w ranging 394 from $\pm 3 \text{ m s}^{-1}$ in both these regions. The *w* field in Fig. 7c suggests that it is meaningful to separate 395 396 between *fixed updrafts*, i.e. part of the stationary wave pattern tied to the orography, and *transient* 397 updrafts, i.e. more short-lived, advecting features associated with turbulence (in this case), or with 398 small-scale convective instabilities within passing weather systems updrafts. This distinction is 399 explored in detail in Part 2.

Note that at and below the melting level, errors in the retrieval are evident due to nonuniformity of the melting level. The cause of these errors will be addressed in the Sec. 6. The retrieved \overline{V}_t was -0.2 m s⁻¹ near cloud echo top, decreasing with depth beneath cloud echo top to ~ -1 m s⁻¹, and further decreasing to ~-4 m s⁻¹ beneath the melting level.

404 Examining uncertainty associated with each of the three assumptions:

405 Assumption 1: $\sigma_{w,1} =\pm 0.03 \text{ m s}^{-1}$ as a consequence of differences between wind speed 406 measured at flight level and the rawinsonde (Fig. 7i).

407 Assumption 2: $\sigma_{w,2} = \pm 0.03 \text{ m s}^{-1}$ between 2.3 and 7.9 km (Fig 7j) where echo extent 408 encompassed the entirety of the 82.8 km flight leg. $\sigma_{w,2} = \pm 0.47 \text{ m s}^{-1}$ near the surface and 409 within 200 m of cloud echo top, where echo extent was less than the length of the flight leg 410 due to the terrain and cloud top variability. The implication is that assumption 2 loses 411 validity (and w cannot be reliably retrieved) near the surface in the presence of complex 412 terrain and near cloud top when the cloud top is uneven and Z_e approaches the minimum 413 detectable signal, leading to increased variance in measured V_r .





Fig. 7: Eastbound flight leg during IOP 23 from 9 March 2017 22:24:40 to 22:35:35 UTC. First 415 row: (a) uncorrected V_r and (b) uncorrected V_r CFAD binned every 0.1 m s⁻¹ and 100 m in altitude. 416 Second row: (c) retrieved w and (d) accompanying CFAD binned every 0.1 m s⁻¹ and 100 m in 417 altitude. Third row: (e) Z_e cross section and (f) Z_e CFAD binned every 1 dBZ_e and 100 m in altitude. 418 419 Fourth row: (g) comparison of w_{gp} (red) at flight level with radar retrieved w (black) and (h) is the 420 retrieved \overline{V}_t profile. Fifth row: (i) ΔW for all nadir and zenith beams, (j) $\sigma_{w,2,m}$, (k) $\sigma_{Z_{\rho},m}$ (black) and $\sigma_{w,3,m}$ (red), and (l) $\sigma_{T,m}$. West is on the left and east on the right of the radar 421 422 cross sections.

423

424 Assumption 3: $\sigma_{w,3} = \pm 0.13 - 0.19 \text{ m s}^{-1}$ between 1.5-7.6 km. In this layer, the Z_e CFAD 425 (Fig. 7f) showed minimal variation in Z_e at all m with $\sigma_{Z_e} < 5$ dBZ_e (Fig. 7k). This suggests 426 that the ice particle ensemble observed were likely undergoing similar growth mechanisms 427 at a given altitude. At 8 km and within 240 m near the surface, $\sigma_{w,3} = \pm 0.20$ -0.25 m s⁻¹. 428 As a result, ± 0.14 m s⁻¹ $< \sigma_T < \pm 0.21$ m s⁻¹ between 1.9 and 7.6 km increasing within 200 m 429 of cloud echo top to ± 0.48 m s⁻¹ and to ± 0.50 within 120 m of the surface (Fig. 71).

430 To determine if σ_T was a reasonable estimate of uncertainty on a flight leg, w_{gp} was taken as 431 truth and the mean absolute error (*MAE*) was calculated along a given flight leg as follows:

$$MAE = \frac{\sum_{i=1}^{k} |w_{gp} - w|}{k}$$

where *k* was the number of collocated gust probe/radar retrieval measurements along a given flight leg and the average radar retrieved *w* is from the first range gates above and below the aircraft averaged together. The comparison shows a close correspondence in time between the retrieved *w* and measured w_{gp} (Fig. 7g). *MAE* was 0.05 m s⁻¹ along the flight leg, less than σ_T at the same altitude (±0.21 m s⁻¹).

438 b) IOP 21 7 March 2017

439 IOP 21 sampled a cloud system over the Payette River Basin that was on the north side of an 440 extratropical cyclone. A V_r cross section from a northeast to southwest (16:13:00 to 16:27:50 UTC) 441 flight leg revealed a split cloud layer (Fig. 8a). The lower cloud layer was predominantly stratiform 442 with boundary layer turbulence present in the lowest kilometer near the terrain.

443 Retrieved w in Figs. 8c,d reveals broadscale orographic lift with updrafts on the windward, 444 and downdrafts on the leeward sides of the mountains. Three regions of stronger updrafts ranging 445 from -3 to 3 m s⁻¹ are evident, the first near the terrain (< 2.5 km), the second near the top and base of the split in the cloud layers (between 5-6 km) and the third near cloud echo top (near 9 km) (Fig. 446 447 8d). At low levels, these are associated with terrain-induced turbulence and terrain-driven eddy 448 dipoles (the wind is left to right). At higher levels, the main mechanism is mixing, likely due to evaporation and radiative cooling along cloud/clear boundaries, as well as Z_e approaching the 449 minimum detectable signal. \overline{V}_t was more variable in this case due to the split cloud layer, 450 decreasing from near -0.1 m s⁻¹ to \sim -1 m s⁻¹ between cloud echo top and the base of the lower 451 452 cloud layer (Fig. 8h).



453

Fig. 8: Same as Fig. 7 except a southwest bound flight leg from IOP 21 from 7 March 2017 16:13:00 to 16:27:50 UTC. Southwest is on the left on panels (a), (c), (e), and (g).

456

457 Examining uncertainty associated with each of the three assumptions:

458 Assumption 1: $\sigma_{w,1} = \pm 0.08 \text{ m s}^{-1}$, a consequence of differences between wind speed 459 measured at flight level and the rawinsonde.

460 Assumption 2: $\sigma_{w,2} = \pm 0.04 \text{ m s}^{-1}$ between 2.2-5 km and 7.4-8.5 km MSL where echo extent 461 encompassed the entirety of the 76.1 km southwest flight leg. $\sigma_{w,2} = \pm 0.46 \text{ m s}^{-1}$ near cloud 462 echo top (9.4 km) and at the surface (Figs. 8j). The split cloud layer present between 6-7 463 km had $\sigma_{w,2} = \pm 0.09$ -0.22 m s⁻¹. Echo extent was limited in the split cloud layer to regions 464 where the upper cloud layer was precipitating into the lower cloud layer. 465 Assumption 3: The Z_e CFAD (Fig. 8f) showed more variation at all m as a result of the layering structure. σ_{Z_e} was largest between surface and 2.5 km reaching 13.3 dBZ_e with $\sigma_{w,3}$ 466 = ±0.34 m s⁻¹. σ_{Z_e} = 3.4-8 dBZe between 3.4-6.5 km with $\sigma_{w,3}$ = ± 0.18-0.24 m s⁻¹ (Fig. 8k). 467 As a result, $\sigma_T = \pm 0.46$ -0.51 m s⁻¹ within 200 m of the surface, decreasing to ± 0.19 m s⁻¹ 468 469 at 2.75 km (the minimum in the lower cloud layer), and then increasing toward cloud echo top of the lower cloud layer to ± 0.24 m s⁻¹ (6.5 km) (Fig. 81). $\sigma_T = \pm 0.29$ to 0.49 m s⁻¹ within 200 m of 470 cloud echo top of the upper cloud layer and had a minimum $\sigma_T = 0.21 \text{ m s}^{-1}$ at 8.5 km. Comparison 471 of w and w_{gp} show close correspondence near flight level (Fig. 8g), with MAE = 0.15 m s⁻¹ along 472 the flight leg. The *MAE* was less than σ_T (± 0.24 m s⁻¹) at flight level. 473

474 *c) IOP 20 5 March 2017*

475 IOP 20 sampled a complex cloud system over the Payette River Basin with high-amplitude 476 gravity wave signatures in the V_r field between 4.5 km and 6.0 km and elevated convection 477 apparent between 6 km and cloud echo top. Boundary layer turbulence was also present in the 478 lowest 1 km above the terrain. The retrieved w showed that the gravity waves had maximum updrafts and downdrafts ranging from -7 m s⁻¹ to 7 m s⁻¹ (Fig. 9c,d). Vertical drafts were regularly 479 spaced, and not related to the terrain as will be discussed in Part 2. Retrieved \overline{V}_t was ~ -0.2 m s⁻¹ 480 near cloud echo top (8 km), decreasing with depth beneath cloud echo top to \sim -1 m s⁻¹ at 2.5 km 481 482 (Fig. 9h).

483 Examining uncertainty associated with each of the three assumptions:

484 Assumption 1: $\sigma_{w,1} = \pm 0.23$ m s⁻¹, a consequence of differences between wind speed 485 measured at flight level and the rawinsonde (Fig 9i).

486 Assumption 2: $\sigma_{w,2} \le \pm 0.1 \text{ m s}^{-1}$ between 1.9 and 8.2 km where echo extent encompassed 487 the entirety of the 72.1 km flight leg (Fig. 9j). $\sigma_{w,2} = \pm 0.46 \text{ m s}^{-1}$ near the surface and cloud 488 echo top.

489 Assumption 3: Z_e at all levels *m* had larger σ_{Z_e} than the previous two cases (Fig. 9k), 490 increasing aloft where elevated convection was located (Fig. 9f). σ_{Z_e} exceeded 5 dBZ_e 491 throughout cloud depth except near cloud echo top. σ_{Z_e} increased near the surface to 15.2 492 dBZ_e at 2.5 km. $\sigma_{w,3}$ increased with depth from \pm 0.21 m s⁻¹ near cloud echo top (8 km) to 493 \pm 0.34 m s⁻¹ at ~2.5 km (Fig. 9i). 494 The result was that $\sigma_T = \pm 0.20$ -0.40 m s⁻¹ between 1.8 and 7.8 km increasing to ± 0.49 m 495 s⁻¹ at cloud echo top and ± 0.58 m s⁻¹ at the surface (Fig. 91). MAE = 0.29 m s⁻¹ along the flight leg 496 and MAE was less than $\sigma_T (\pm 0.43 \text{ m s}^{-1})$ at the same level.



497

Fig. 9: Same as Fig. 7 but for a southwest flight leg from IOP 20 from 5 March 2017 13:32:50 to 13:46:50 UTC.

500

501 6. Examining scenarios with large uncertainty

The retrieval of *w* using the sounding correction for horizontal wind contribution presented above is best applied in deep stratiform cloud systems with uniform cloud coverage at all levels where large variations in particle V_t at a given level are not present. In the SNOWIE data, one situation was found to violate the assumption of constant \bar{V}_t at a given a level, specifically a sloped melting level. Two additional situations were found to violate the assumption of uniform horizontal cloud extent: convective turrets and split layers. We illustrate these three problemsbelow.

509 a) Sloped Melting Level

Figure 10 shows an example from IOP 23 where a sloped melting level was present decreasing in altitude from 2.8 km to 2 km as the aircraft traveled west to east over the Payette River Basin. Horizontal variation in the melting layer level resulted in an underestimation along the melting level of w on the western end of the flight leg and an overestimate of w along the eastern end of the flight leg, a result of over or under correction based on the subtraction of $\overline{V_t}$ along a non-homogenous feature at a constant grid level m.



517 Fig. 10: Westbound flight leg during IOP 23 on 9 March 2017 from 23:15:00 to 23:32:30 UTC.

- 518 (a) Z_e . The dashed black line is the aircraft flight level, (b) w, (c) comparison of w_{gp} (red) measured 519 at flight level and radar retrieved w (black) from nearest range gates above and below the aircraft,
- 520 (d) ΔW for nadir and zenith pointing beams, (e) $\sigma_{w,2,m}$, (f) $\sigma_{Z_{e,m}}$ (black) and $\sigma_{w,3,m}$ (red), (g) $\sigma_{T,m}$.

521 b) Convection

522 Figure 11a shows an example of elevated convection during IOP 12. The horizontal extent 523 of cloud top echo was limited, consisting only of an elevated convective turret ~25 km wide. 524 Examining uncertainty associated with assumptions 2 and 3:

- Assumption 2: $\sigma_{w,2} = \pm 0.15$ -0.46 m s⁻¹ between 5 and 6.6 km due to echo extent (Fig. 11d). 525
- Assumption 3: The elevated convective turnet had $\sigma_{Z_e} = 11-14 \text{ dBZ}_e$ between 4-6 km with 526 $\sigma_{w,3} = \pm 0.3 - 0.35 \text{ m s}^{-1}$ (Fig. 11e).
- 527
- As a result, σ_T was estimated to be ± 0.38 m s⁻¹ near flight level. There were also large differences 528
- between retrieved w and w_{gp} measured at flight level with MAE = 0.62 m s⁻¹ (Fig. 11c). These 529 530 results show the impact of small echo extent along a given *m* on retrieval uncertainty.



531

532 Fig. 11: Same as Fig. 10 except for an eastbound flight leg during IOP 12 on 7 February 2017

533 from 20:50:00 to 20:59:30 UTC.

534 c) Split Layers

535 IOP 3 sampled a split cloud layer along the western end of the flight leg and deep stratiform 536 cloud along the eastern end (Fig. 12a, b). Along the western end of the flight leg, the lower cloud 537 layer had Kelvin-Helmholtz waves near cloud echo top. The retrieved w along the western end of 538 the flight leg had greater uncertainty due to inhomogeneities in Z_e along the leg. The value of $\sigma_{Z_{a}}$ was < 5 dB above 6.1 km but beneath 6.1 km, the western half of the flight leg had low average 539 Z_e (< 0 dBZ_e) and high Z_e on the eastern half of the flight leg (> 0 dBZ_e) associated with the deep 540 541 precipitating orographic cloud. Throughout the split layer $\sigma_{Z_e} > 10$ dB reaching 25.8 dB at 4.3 km. As a result, $\sigma_{w,3} = \pm 0.5$ m s⁻¹ at flight level (Fig. 12e). In this case there were also larger 542 differences between retrieved w and w_{gp} measured at flight level with MAE = 0.34 m s⁻¹ (Fig. 12c). 543 544 The examples illustrated show that different events may have different sources of uncertainty when 545 retrieving w.

546 7. Summary of *w* retrieval uncertainty on all flight legs

The retrieval method presented in Sec. 3 was applied to obtain w and \overline{V}_t for the entire 547 SNOWIE dataset. Figure 13 shows a summary of σ_T as a function of all retrieved w values during 548 549 the 23 SNOWIE research flights. Figure 13b shows that 67% of retrieved w values had updrafts/downdrafts between ± 0.71 m s⁻¹ over the Payette River Basin and 95% of retrieved w 550 551 values had w between ± 1.42 m s⁻¹. Most updrafts were relatively weak associated with stratiform 552 ascent/descent within fixed orographically-induced waves. The range of σ_T typically increased as the magnitude of updrafts/downdrafts increased. Median σ_T increased from 0.22 m s⁻¹ to 0.39 m 553 s⁻¹ as w increased from 0 m s⁻¹ to ± 10 m s⁻¹. The 95th percentile of σ_T increased from 0.43 m s⁻¹ to 554 0.62 m s⁻¹ as w magnitude increased from 0 m s⁻¹ to ± 10 m s⁻¹. The 5th percentile increased slightly 555 from 0.18 m s⁻¹ to 0.23 m s⁻¹ as w increased from 0 m s⁻¹ to ± 10 m s⁻¹. Stronger updrafts and 556 557 downdrafts were typically associated with a wider range of uncertainties. For example, an updraft between $4.5 - 5 \text{ m s}^{-1}$ had a median σ_T of ~0.4 m s⁻¹, while an updraft of $0 - 0.5 \text{ m s}^{-1}$ had a median 558 559 $\sigma_{\rm T}$ of ~0.2 m s⁻¹.

560 To further evaluate the performance of the *w* retrieval, each flight leg during SNOWIE was 561 interrogated. For these legs, retrieved *w* above the aircraft and below the aircraft was averaged and 562 compared to w_{gp} , resulting in 59,701 collocated *w* and w_{gp} measurements (Fig. 14). The absolute 563 difference $|w-w_{gp}|$ was calculated for all samples and the distribution of these values is shown as 564 in Fig. 14. The median $|w-w_{gp}|$ was 0.18 m s⁻¹ and the mean was 0.27 m s⁻¹. 565



Fig. 12: Same as Fig. 10 except for a westbound flight leg during IOP 3 on 11 January 2017 from03:07:00 to 03:24:00 UTC.



570 Fig. 13: a) Boxplots of σ_{T} for all *w* values during SNOWIE binned every 0.5 m s⁻¹. Orange lines 571 denote median σ_{T} for a given *w* bin. The upper bound of the box represents the 75th percentile of 572 σ_{T} while the lower bound of the box represents the 25th percentile of σ_{T} . The whiskers represent 573 the 5th and 95th percentile of σ_{T} . b) the percentage of *w* values sampled during SNOWIE.

574

569



575

576 Fig. 14: $|w - w_{gp}|$ for all collocated UWKA *w* retrievals and w_{gp} data points as a percentage on the 577 left axis. The black curve and right axis show the cumulative percentage. The vertical black line 578 represents the median $|w - w_{gp}|$.

580 8. Conclusions and Discussion

This paper presented an analysis of uncertainties associated with assumptions made when retrieving vertical air motion (*w*) and mean profiles of reflectivity-weighted terminal velocity (\overline{V}_t) from airborne measurements of Doppler radar radial velocity (V_r) from the nadir and zenith WCR antennas. This retrieval methodology and its assumptions are directly applicable to any airborne vertically-pointing Doppler radars. Assumptions were tested in orographic clouds over the Payette River Basin of Idaho sampled during the Seeded and Natural Orographic Wintertime Clouds: the Idaho Experiment (SNOWIE).

The retrieval technique for extracting w and $\overline{V_t}$ from V_r involves correcting V_r for known pitch, 588 589 roll, and yaw angle deviations due to aircraft motion using the magnitude of the horizontal wind 590 components (u, v) at a given height measured independently by a rawinsonde. This allows for the 591 retrieval of vertical radial velocity, W, effectively the hydrometeor vertical velocity, from which w and $\overline{V_t}$ can be retrieved. The accuracy of the retrieval of w and $\overline{V_t}$ was assessed and shown to be 592 593 dependent on satisfying assumptions that (a) the flight legs occur over a short enough time and 594 distance that the along and across track winds at a given altitude above/below the aircraft do not 595 vary horizontally or change with time, (b) the legs are long enough for the magnitudes of the updrafts and downdrafts at any given altitude to sum to 0 m s⁻¹, and (c) that the reflectivity-596 weighted hydrometeor \overline{V}_t does not vary substantially at a given altitude, such that $V_{t,m}$ at any point 597 along the flight leg can be approximated by $\overline{V_{t.m}}$. A method to estimate the uncertainty in the 598 599 retrieval of w as a function of altitude was presented based on an evaluation of these assumptions. 600 Each of these assumptions were evaluated quantitatively for example case studies and for the entire 601 project dataset.

Case studies from SNOWIE research flights show that this methodology can provide estimates of w that closely matched measurements at the aircraft level. Deep stratiform precipitation with a rather flat cloud top and little Z_e variation at a given height is associated with the least retrieval uncertainty. The greatest uncertainty occurred in regions with isolated convective turrets, and at altitudes where split cloud layers were evident. Greater uncertainty also occurred in the presence a sloped melting level. Assumption (b) loses validity, and w cannot be reliably retrieved, near cloud top, and, in the presence of complex terrain, near the surface. In Part 2 (Zaremba et al. 2022), we apply this retrieval technique to examine representative
fixed and transient updraft structures present over the Payette River Basin of Idaho during
SNOWIE and their relationship to thermodynamic forcing.

612

APPENDIX

613 Appendix A: Retrieval of W using a rawinsonde correction

614 The goal is to retrieve the vertical hydrometeor motion, $W = w - V_t$, along a radar beam 615 from measured radial velocity V_r by removing contributions to V_r by aircraft motion and the 616 horizontal wind.

617 The correction involves application of the transformation matrix, T_{A2G} , from aircraft to 618 ground relative coordinates (x east-west, y north-south, and z up-down). T_{A2G} is the inverse of 619 Haimov and Rodi (2013) where:

$$620 \quad T_{A2G} = \begin{pmatrix} t_{11} & t_{12} & t_{13} \\ t_{21} & t_{22} & t_{23} \\ t_{31} & t_{32} & t_{33} \end{pmatrix}$$

$$621 \quad = \begin{pmatrix} \sin(h)\cos(p) & \cos(h)\cos(p) & \sin(p) \\ \cos(h)\cos(r) + \sin(h)\sin(p)\sin(r) & -\sin(h)\cos(r) + \cos(h)\sin(p)\sin(r) & -\cos(p)\sin(r) \\ -\cos(h)\sin(r) + \sin(h)\sin(p)\cos(r) & \sin(h)\sin(r) + \cos(h)\sin(p)\cos(r) & -\cos(p)\cos(r) \end{pmatrix}$$

622 and p, h, and r, are the pitch, heading, and roll of the aircraft measured by the navigation system.

Let \vec{b} be the calibrated beam-pointing vector in aircraft coordinates, $\vec{b_g}$, the beam-pointing vector in ground relative coordinates (where $\vec{b_g} = \vec{b} T_{A2G}$), $\vec{V_{ac}}$, the aircraft velocity vector in ground coordinates, and $\vec{V_s}$, the mean scatterer velocity vector in ground-relative coordinates where $\vec{V_s} = \vec{V} + \vec{V_t}$, where \vec{V} is the 3D wind vector, and $\vec{V_t}$ is the pulse-volume average terminal velocity vector. V_r is equivalent to:

628

$$V_r = \overrightarrow{b_g} \cdot \left(\overrightarrow{V_{ac}} + \overrightarrow{V_s} \right) = \left(\overrightarrow{b} T_{A2G} \right) \cdot \left(\overrightarrow{V} + \overrightarrow{V_t} + \overrightarrow{V_{ac}} \right)$$
(1)

629 The vectors in the x, y, and z directions in equation 1 are:

$$\vec{b} = (b_x, b_y, b_z)$$

$$\vec{V} = (u, v, w)$$

$$\overrightarrow{V_t} = (0, 0, -V_t)$$

 $\overline{V_{ac}} = (V_{ax}, V_{ay}, V_{az})$

634 Multiplying the beam pointing vector by the transformation matrix results in:

635
$$(\vec{b} T_{A2G}) = b_x t_{11} + b_y t_{21} + b_z t_{31} + b_x t_{12} + b_y t_{22} + b_z t_{23} + b_x t_{31} + b_y t_{32} + b_z t_{33}$$

636 The dot product of the beam transformation vector and the wind vector (\vec{V}) is:

637
$$(\vec{b} T_{A2G}) \cdot (\vec{V}) = (\vec{b} T_{A2G}) \cdot (u, v, w)$$

638
$$= b_x(t_{11}u + t_{12}v + t_{13}w) + b_y(t_{21}u + t_{22}v + t_{23}w) + b_z(t_{31}u + t_{32}v)$$

639 $+ t_{33}w$)

640 This can be simplified as:

641
$$(\bar{b} T_{A2G}) \cdot (u, v, w) = b_{t1}u + b_{t2}v + b_{t3}w$$

642 where:

643 643 $b_{t1} = b_x t_{11} + b_y t_{21} + b_z t_{31}$ 644 $b_{t2} = b_x t_{12} + b_y t_{22} + b_z t_{32}$ 645 $b_{t3} = b_x t_{13} + b_y t_{23} + b_z t_{33}$

646 The dot product of the beam transformation vector and terminal velocity vector $(\vec{V_t})$ is:

 $(\vec{b} \ T_{A2G}) \ \cdot (\vec{V_t}) = -b_{t3}V_t$

648 The dot product of the beam transformation vector and the aircraft motion vector $(\overrightarrow{V_{ac}})$ is:

$$(\vec{b} \ T_{A2G}) \cdot (\vec{V_{ac}}) = b_{t1}V_{ax} + b_{t2}V_{ay} + b_{t3}V_{az}$$

650 so that (1) becomes:

651
$$V_r = b_{t1}u + b_{t2}v + b_{t3}w - b_{t3}V_t + b_{t1}V_{ax} + b_{t2}V_{ay} + b_{t3}V_{az}$$
(2)

652 Solving for $w - V_t$ or W (vertical radial velocity) at each range gate:

653
$$W = w - V_t = \frac{V_r - (b_{t1}u + b_{t2}v + b_{t1}V_{ax} + b_{t2}V_{ay} + b_{t3}V_{az})}{b_{t3}}$$

654

The University of Wyoming King Air facility provides Level 1 and Level 2 data that is corrected for aircraft motion but not the horizontal wind contribution. The radial velocity provided by the facility, V'_r , is:

658
$$V'_{r} = V_{r} - b_{t1}V_{ax} + b_{t2}V_{ay} + b_{t3}V_{az} = b_{t1}u + b_{t2}v + b_{t3}w - b_{t3}V_{t}$$
(3)

For the provided data, corrected for aircraft motion, the retrieval of W for a single range gate becomes:

661 $W = \frac{V_r' - (b_{t1}u + b_{t2}v)}{b_{t3}}$ (4)

663	Appendix B: List of variables and their descriptions				
664	$ec{b}$	calibrated beam pointing vector in aircraft coordinates			
665	b_x	beam vector in x direction			
666	b_{v}	beam vector in y direction			
667	b_z	beam vector in z direction			
668	$\overrightarrow{b_g}$	calibrated beam pointing vector in ground relative coordinates			
669	h	heading			
670	т	a given height (altitude) index			
671	m_{ac}	aircraft altitude index			
672	MAE	mean absolute error			
673	n	beam index			
674	p	pitch			
675	r	roll			
676	$\sigma_{\rm T}$	total uncertainty			
677	$\sigma_{w,1}$	uncertainty due to assumption 1			
678	$\sigma_{w,2}$	uncertainty due to assumption 2			
679	$\sigma_{w,3}$	uncertainty due to assumption 3			
680	$\sigma_{ w-w_{gp} }$	standard deviation of the absolute value of vertical air velocity minus			
681		vertical air velocity measured by the gust probe			
682	$\sigma_{\Delta w}$	standard deviation of the difference in vertical radial velocity as a result of			
683		differences in rawinsonde measured winds and aircraft measured winds for			
684		all beams along a given flight leg			
685	σ_{Z_e}	standard deviation of equivalent reflectivity factor			
686	T_{A2G}	transformation matrix from aircraft to ground coordinates			
687	u	zonal wind component			
688	Δu	difference in the zonal wind component between the aircraft and sounding			
689		at the altitude of the aircraft for a given beam			
690	v	meridional wind component			
691	Δv	difference in the meridional wind component between the aircraft and			
692		sounding at the altitude of the aircraft for a given beam			
693	\overline{V}	wind vector in aircraft relative coordinates			
694	$\overline{V_{ac}}$	aircraft velocity vector in ground coordinates			
695	V_{ax}	aircraft velocity in x direction			
696	V_{ay}	aircraft velocity in y direction			
697	V_{az}	aircraft velocity in z direction			
698	$\overrightarrow{V_H}$	horizontal wind vector			
699	V_r	measured Doppler radial velocity			
700	V_r'	radial velocity corrected for aircraft motion but not horizontal wind			
701	,	contribution provided by the UWKA facility.			
702	$\overline{V_r}$	mean measured Doppler radial velocity			
703	$\overrightarrow{V_s}$	mean scatter velocity vector in ground relative coordinates			
704	V_t	reflectivity weighted terminal velocity of hydrometeors			
705	V_{tm}	reflectivity weighted terminal velocity at a given height			
706	$V_{t m n}$	reflectivity weighted terminal velocity at a given height along a given beam			
	0,110,10				

707	$\overrightarrow{V_t}$	pulse-volume average terminal velocity vector
708	$\overline{V_{t.m}}$	mean reflectivity weighted terminal velocity at a given height
709	$\overline{V_t}$	mean reflectivity weighted terminal velocity of hydrometeors
710	W	vertical air velocity
711	W_{gp}	vertical air velocity measured by gust probe
712	\overline{W}	mean vertical air velocity
713	$W_{m,n}$	vertical air velocity at a given height along a given beam
714	$\overline{W_m}$	mean vertical air velocity at a given height
715	W	vertical radial velocity
716	Δw	difference in vertical air velocity as a result of differences in rawinsonde
717		measured winds and aircraft measured winds
718	$W_{m,n}$	vertical component of radial velocity at a given height along a given beam
719	\overline{W}	mean vertical component of radial velocity
720	$\overline{W_m}$	mean vertical component of radial velocity at a given height
721	x	east-west direction
722	У	north-south direction
723	Z	up-down direction
724	Z_e	equivalent reflectivity factor
725		

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737 (https://data.eol.ucar.edu/master_lists/generated/snowie/) maintained by the Earth Observing

T38 Laboratory (EOL) at the National Center for Atmospheric Research (NCAR).

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