



G_2 Manifolds with Nodal Singularities Along Circles

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Received: 8 March 2019 / Published online: 23 September 2019
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Abstract

This paper finds matching building blocks for the construction of a compact manifold with G_2 holonomy and nodal singularities along circles using twisted connected sum method by solving the Calabi conjecture on certain asymptotically cylindrical manifolds with nodal singularities. However, by comparison to the untwisted connected sum case, it turns out that the obstruction space for the singular twisted connected sum construction is infinite dimensional. By analyzing the obstruction term, there are strong evidences that the obstruction may be resolved if a further gluing is performed in order to get a compact manifold with G_2 holonomy and isolated conical singularities with link $\mathbb{S}^3 \times \mathbb{S}^3$.

Keywords G_2 · Holonomy · Calabi–Yau · Edge singularity

Mathematics Subject Classification Primary 53C29 · Secondary 32Q25

1 Introduction

This paper makes progress on the construction of a compact manifold with G_2 holonomy and nodal singularities along circles using the twisted connected sum method. The twisted connected sum method was used by Kovalev [43] and Corti et al. [21] to construct smooth manifold with G_2 holonomy. In the singular case, Sect. 8.4.5 of [21] proposed three technical problems. The first problem is the construction of asymptotically cylindrical Calabi–Yau threefolds with nodal singularities. The second problem is finding matching data on the ends. The third problem is to control neck length in the gluing construction. This paper solves the first and the second problem and finds the obstruction to the third problem which may be resolved by a further gluing.

The first problem is solved by combining the Theorem 1.4 of Hein and Sun’s work [36] with Theorem D of Haskins et al. work [34]:

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Theorem 1.1 *Let $i : \mathcal{Z} \subset \mathbb{CP}^{N_1} \times \Delta$ be a flat family of projective varieties over disc Δ . Let π_1 be the projection of $\mathbb{CP}^{N_1} \times \Delta$ to $\mathbb{CP}^{N_1}$ and π_2 be the projection of $\mathbb{CP}^{N_1} \times \Delta$ to Δ . Denote $(\pi_1 \circ i)^*(\mathcal{O}(1))$ by $\mathcal{L}$. Denote $(\pi_2 \circ i)^{-1}(s)$ by Z_s . Denote $\mathcal{L}|_{Z_s}$ by L_s . Suppose that there exists an morphism $f : \mathcal{Z} \rightarrow \mathbb{CP}^1 = \mathbb{C} \cup \{\infty\}$. Denote $f^{-1}(\infty)$ by $\mathcal{S}$. Denote $\mathcal{Z} \setminus \mathcal{S}$ by $\mathcal{V}$. Denote $\mathcal{S} \cap Z_s$ by S_s . Denote $Z_s \setminus S_s$ by V_s . Suppose that V^{sing} is a finite subset of V_0 . Suppose that $\mathcal{Z}$ is smooth and the induced map $(\pi_2 \circ i)_*$ on the tangent space is surjective at each point on $\mathcal{Z} \setminus V^{\text{sing}}$. Suppose that df is not the pull back of any form on Δ at each point on $\mathcal{S}$. Suppose that the complex dimension n of Z_s is at least 3. Suppose that for each $x \in V^{\text{sing}}$, there exists a holomorphic function $\epsilon_x(s)$ with $\epsilon_x(0) = 0$ such that the germ $(\mathcal{Z}, x, \pi_2 \circ i)$ is isomorphic to the germ (C_x, o_x, π_x) , where*

$$C_{x,s} = \pi_x^{-1} C_x(s) = \left\{ z_1^2 + \cdots z_n^2 = \epsilon_x(s) \right\},$$

and o_x is the tip point of $C_x = C_{x,0}$. Assume that Ω_s is a meromorphic family of meromorphic n -forms on Z_s . Assume that Ω_s is holomorphic on $\mathcal{V} \setminus V^{\text{sing}}$. Assume that $\frac{\Omega_s}{f}$ is holomorphic near $\mathcal{S}$. Assume that the ratio of Ω_s to $\Omega_{\epsilon_x(s)}$ in Example 2.6 is holomorphic near x . For simplicity, denote $(Z_0, S_0, L_0, \Omega_0, V_0)$ by (Z, S, L, Ω, V) . Then after replacing Ω by its product with a constant, there exists an asymptotically cylindrical Calabi–Yau metric $\omega \in c_1(L)|_V$ on V such that

$$\frac{\omega^n}{n!} = \frac{i^{n^2}}{2^n} \Omega \wedge \bar{\Omega}.$$

Moreover, ω has conical singularity with rate $\lambda > 0$ and tangent cone (C_x, ω_{C_x}) at x as in Definition 2.12.

Roughly speaking, the manifold is assumed as the central fiber of a flat family of manifolds over disc such that the complex structure and the holomorphic top form Ω_s are standard near the nodal singularities as required by Hein and Sun [36]. More conditions on Ω_s and a function f are proposed to make sure the construction of the cylindrical end as required by Haskins et al. [34]. Theorem 1.1 provides the asymptotically cylindrical Calabi–Yau metric assuming the existence of such a good complex structure and the holomorphic top form Ω_s .

Then the second problem is solved for a particular example using the additional information about the quartic K3 surfaces in $\mathbb{CP}^3$ using Theorem 4.2 and Chapter 3 of [57]. Roughly speaking, the building blocks are obtained by removing K3 surfaces from the blow-ups of Fano 3-manifolds. The hyperKähler triples on the two K3 surfaces are required to be matched after a twist.

Proposition 1.2 *It is possible to find the following data with required properties:*

- (1) $X_- = \mathbb{CP}^3$. Π is a 2-plane in $\mathbb{CP}^4$. X_+ is a quartic 3-fold in $\mathbb{CP}^4$ containing Π with nine nodal singularities X_+^{sing} .
- (2) $|S_{0,\pm}, S_{\infty,\pm}| \subset |-K_{X_{\pm}}|$ are pencils with smooth base locuses $C_{\pm}$ disjoint with X_+^{sing} . $Z_{\pm}$ are the blow-up of $X_{\pm}$ at $C_{\pm}$.

- (3) $S_{\pm}$ are smooth K3 surfaces in $|S_{0,\pm}, S_{\infty,\pm}|$ disjoint with X_+^{sing} . Their proper transforms are also denoted by $S_{\pm} \subset Z_{\pm}$. $\Omega_{\pm}$ are meromorphic 3-forms on $Z_{\pm}$ with simple poles along $S_{\pm}$.
- (4) $(S_{\pm}, \omega_{S_{\pm}}, \omega_{S_{\pm}}^J + i\omega_{S_{\pm}}^K)$ are Calabi–Yau surfaces.
- (5) $\omega_{\pm}$ are asymptotically cylindrical Calabi–Yau metrics on $V_{\pm} = Z_{\pm} \setminus S_{\pm}$ with $\frac{\omega_{\pm}^3}{6} = \frac{i}{8}\Omega_{\pm} \wedge \bar{\Omega}_{\pm}$.
- (6) ω_+ has conical singularity in the sense of Definition 2.12 with the nodal singularity in Example 2.7 as the tangent cone for all $x \in V_+^{\text{sing}}$.
- (7) $K_{\pm}$ are compact subsets of $V_{\pm}$. $P_{\pm} : [1, \infty) \times \mathbb{S}^1 \times S_{\pm} \rightarrow V_{\pm} \setminus K_{\pm}$ are diffeomorphisms on the ends. $t_{\pm}$ are coordinates on $[1, \infty)$, $\vartheta_{\pm}$ are coordinates on $\mathbb{S}^1$. Up to exponentially decaying errors $\varrho_{\pm}$ and $\varsigma_{\pm}$,

$$(P_{\pm}^* \omega_{\pm}, P_{\pm}^* \Omega_{\pm}) = (\omega_{\infty,\pm}, \Omega_{\infty,\pm}) + (d\varrho_{\pm}, d\varsigma_{\pm}),$$

where

$$\omega_{\infty,\pm} = dt_{\pm} \wedge d\vartheta_{\pm} + \omega_{S_{\pm}},$$

and

$$\Omega_{\infty,\pm} = (d\vartheta_{\pm} - i dt_{\pm}) \wedge (\omega_{S_{\pm}}^J + i\omega_{S_{\pm}}^K).$$

- (8) r is a diffeomorphism from $(S_+, \omega_{S_+}, \omega_{S_+}^J, \omega_{S_+}^K)$ to $(S_-, \omega_{S_-}^J, \omega_{S_-}, -\omega_{S_-}^K)$.

Using the same notations as in Proposition 1.2, $t_{\pm}$ can be extended to non-negative smooth functions on $V_{\pm}$ such that t_+ equals to 0 near $V_+^{\text{sing}} = X_+^{\text{sing}}$. Choose $\chi = \chi(s) : \mathbb{R} \rightarrow [0, 1]$ as a smooth function satisfying $\chi(s) = 1$ for $s \leq 1$ and $\chi(s) = 0$ for $s \geq 2$. Using the data in Proposition 1.2, as in Sect. 3 of [21], for fixed large enough T , define

$$\omega_{T,\pm} = \omega_{\pm} - d((1 - \chi(t_{\pm} - T + 2))\varrho_{\pm})$$

and

$$\Omega_{T,\pm} = \Omega_{\pm} - d((1 - \chi(t_{\pm} - T + 2))\varsigma_{\pm})$$

on $V_{\pm}$. Let $M_{\pm}$ be $\mathbb{S}^1 \times V_{\pm}$. Let $\theta_{\pm}$ be the coordinates on $\mathbb{S}^1$. Define

$$\varphi_{T,\pm} = d\theta_{\pm} \wedge \omega_{T,\pm} + \text{Re } \Omega_{T,\pm}.$$

Remark that using the diffeomorphism

$$(\theta_-, t_-, \vartheta_-, x_-) = (\vartheta_+, 2T + 1 - t_+, \theta_+, r(x_+)),$$

$\varphi_{T,\pm}$ can be glued into a closed G_2 structure φ_T on M , the manifold obtained by gluing M_+ with M_- .

The construction of a G₂ manifold with nodal singularities along circles is reduced to finding a perturbation of φ_T which induces a metric with G₂ holonomy but still preserves the singularities. However, the analysis on manifolds with conical singularities along smooth submanifolds is very complicated. A slightly simpler problem is the analysis on manifolds with isolated conical singularities. Therefore, this paper starts from solving an analogous problem instead. In this case, $(Z_+, S_+, V_+, \omega_+, \Omega_+)$ is the same as in Proposition 1.2 but $(Z_-, S_-, V_-, \omega_-, \Omega_-)$ is $(Z_+, S_+, V_+, \omega_+, -\Omega_+)$ instead and r is the identity map instead. So r is a diffeomorphism from $(S_+, \omega_{S_+}, \omega_{S_+}^J, \omega_{S_+}^K)$ to $(S_-, \omega_{S_-}, -\omega_{S_-}^J, -\omega_{S_-}^K)$ instead and the gluing map is given by

$$(\theta_-, t_-, \vartheta_-, x_-) = (\theta_+, 2T + 1 - t_+, -\vartheta_+, r(x_+)).$$

Remark that in this case, the $\mathbb{S}^1$ factor with coordinate $\theta = \theta_- = \theta_+$ is global. So the theorem in this case is.

Theorem 1.3 (Doubling construction of Calabi–Yau threefolds) *For the new choice of gluing data, for sufficiently large T , there is an $\mathbb{S}^1$ -invariant perturbation φ of φ_T such that the holonomy group of the metric defined by φ is contained in $\mathrm{SU}(3) \subset \mathrm{G}_2$ and for each $x \in V_{\pm}^{\mathrm{sing}}$, there exist numbers $c_{1,x} > 0$, $c_{2,x} > 0$, $c_{3,x}$ and a homeomorphism $P_x : O_x \rightarrow U_x$ between a neighborhood $o \in O_x \subset C_x$ and $x \in U_x \subset V_{\pm}$, such that*

$$\left| \nabla_{\varphi_{\mathbb{S}^1 \times C_x}}^j ((\mathrm{Id} \times P_x)^* \varphi - \varphi_{\mathbb{S}^1 \times C_x}) \right|_{\varphi_{\mathbb{S}^1 \times C_x}} = O(r^{\lambda-j})$$

as $r \rightarrow 0$ for a positive number λ and all $j \in \mathbb{N}_0$, where

$$\varphi_{\mathbb{S}^1 \times C_x} = c_{1,x} d\theta \wedge \omega_{C_x} + c_{2,x} \mathrm{Re} \left(e^{ic_{3,x}} \Omega_{C_x} \right).$$

To emphasize that there is a global $\mathbb{S}^1$ factor, in this case, it is better to use $\mathbb{S}^1 \times M$ instead of M to denote the gluing of $\mathbb{S}^1 \times V_{\pm}$.

At the level of complex structures, this gluing result may be proved using the method developed in [45]. If this is the case, then the new content in Theorem 1.3 is an asymptotic description of the Calabi–Yau metrics as the parameter T goes to infinity.

Remark that the non-singular version of Theorem 1.3 was proved by Doi and Yotsutani [24].

A large portion of the proof of Theorem 1.3 is inspired by the work of Karigiannis and Lotay [40] and has an analogy in [40]. The main tool of the proof of Theorem 1.3 is the weighted analysis developed by Lockhart and McOwen [47] and Melrose and Mendoza [53] independently and further refined by many people. One of the key points is Theorem 2.19 when the weight changes. Another key point is the study of harmonic forms on the nodal singularity.

Back to the singular twisted connected sum case. It involves weighted analysis for manifolds with edge singularities. It was pioneered by Mazzeo [51] and followed by many people. In this paper, the analogue of Theorem 2.19 is proved. However, the obstruction space in this case is infinite dimensional.

In personal discussions with the author, Sir Simon Donaldson and Edward Witten conjectured that the nodal singularities along circles may be replaced by isolated conical singularities with the homogeneous space $(\mathrm{SU}(2) \times \mathrm{SU}(2) \times \mathrm{SU}(2))/\mathrm{SU}(2)$ as the link. As pointed out by Atiyah and Witten [3], there are three ways of resolving the cone over $(\mathrm{SU}(2) \times \mathrm{SU}(2) \times \mathrm{SU}(2))/\mathrm{SU}(2)$. In this paper, by analyzing the infinite dimensional obstruction space, there are strong evidences that this conjecture is correct. It is left for future studies.

The basic facts about G_2 structures, the nodal singularity and the weighted analysis are reviewed in Sect. 2. Theorem 1.1 is proved in Sect. 3. Proposition 1.2 is proved in Sect. 4. Harmonic forms on the nodal singularity is studied in Sect. 5. Theorem 1.3 is proved in Sect. 6. The analogue of the refined change of index formula for singular twisted connected sum, the fact that the obstruction space is infinite dimensional as well as the conjectural picture are discussed in Sect. 7.

2 Preliminaries

Definition 2.1 A G_2 structure on a 7-dimensional manifold M is defined by a 3-form φ such that at each point, after the identification of the tangent space with $\mathbb{R}^7$, there exists an element in $\mathrm{GL}(7, \mathbb{R})$ which maps φ to $e^{123} + e^{145} + e^{167} + e^{246} - e^{257} - e^{347} - e^{356}$, where $e^{ijk} = e^i \wedge e^j \wedge e^k$ and $\{e^i\}$ are the standard basis of $(\mathbb{R}^7)^*$. It induces a metric g_φ by

$$g_\varphi(u, v) \mathrm{Vol}_{g_\varphi} = \frac{1}{6} (u \lrcorner \varphi) \wedge (v \lrcorner \varphi) \wedge \varphi.$$

Definition 2.2 The G_2 structure provides a g -orthogonal decomposition of forms on M . In particular, for three forms, $\Omega^3(M) = \Omega_1^3(M) \oplus \Omega_7^3(M) \oplus \Omega_{27}^3(M)$, where $\Omega_1^3(M) = \{f\varphi : f \text{ is a function on } M\}$, $\Omega_7^3(M) = \{X \lrcorner *_{g_\varphi} \varphi : X \text{ is a vector field on } M\}$ and the orthogonal complement is $\Omega_{27}^3(M)$.

Proposition 2.3 (Lemma 3.1.1 of [38]) Denote $*_{g_\varphi} \varphi$ by $\Theta(\varphi)$, then

$$\Theta(\varphi + \gamma) = *_{g_{\varphi+\gamma}}(\varphi + \gamma) = *_{g_\varphi} \varphi + \frac{4}{3} \pi_1(\gamma) + *_{g_\varphi} \pi_7(\gamma) - *_{g_\varphi} \pi_{27}(\gamma) - Q(\gamma),$$

where π_1, π_7 and π_{27} are the orthogonal projection to Ω_1^3, Ω_7^3 and Ω_{27}^3 with respect to φ , and Q is the higher order term satisfying the estimates in Lemma 3.1.1 of [38].

Theorem 2.4 ([29]) The holonomy group of a metric g is contained in G_2 if and only if g is induced by a G_2 structure φ satisfying $d\varphi = d\Theta(\varphi) = 0$.

Therefore, it suffices to consider the moduli space of G_2 structures φ satisfying $d\varphi = d\Theta(\varphi) = 0$.

Definition 2.5 Suppose that M is a Kähler manifold with complex dimension n . Suppose that ω is a Kähler form and Ω is a holomorphic n -form. Then $(M, g, J, \omega, \Omega)$ is called a Calabi–Yau n -fold if

$$\frac{\omega^n}{n!} = \frac{i^{n^2}}{2^n} \Omega \wedge \bar{\Omega}.$$

Example 2.6 ([61]) Let $C_\epsilon = \{z_1^2 + \cdots + z_{n+1}^2 = \epsilon\} \subset \mathbb{C}^{n+1}$. When $\epsilon \neq 0$, up to scaling, the unique $\mathrm{SO}(n+1)$ -invariant asymptotically conical Calabi–Yau metric g on C_ϵ is given by

$$\omega_\epsilon = \frac{i}{2} \partial \bar{\partial} |\epsilon|^{\frac{n-1}{n}} \left(f \left(\cosh^{-1} \left(\frac{|z_1|^2 + \cdots + |z_{n+1}|^2}{|\epsilon|} \right) \right) \right),$$

where $(f'(w)^n)' = n(\sinh w)^{n-1}$, and $f(0) = f'(0) = 0$. When $\epsilon = 0$,

$$\omega_0 = \frac{i}{2} \partial \bar{\partial} \left(\frac{n}{n-1} \right)^{\frac{n+1}{n}} (|z_1|^2 + \cdots + |z_{n+1}|^2)^{\frac{n-1}{n}}.$$

Define

$$\Omega_\epsilon = \frac{dz_1 dz_2 \cdots dz_n}{z_{n+1}},$$

then by direct calculation,

$$\frac{\omega_\epsilon^n}{n!} = \frac{i^{n^2}}{2^n} \Omega_\epsilon \wedge \bar{\Omega}_\epsilon.$$

Example 2.7 Let C be the nodal singularity $\{z_1^2 + \cdots + z_4^2 = 0\} \subset \mathbb{C}^4$. Then

$$\omega_C = \frac{i}{2} \partial \bar{\partial} \left(\frac{3}{2} \right)^{\frac{4}{3}} (|z_1|^2 + \cdots + |z_4|^2)^{\frac{2}{3}}$$

and

$$\Omega_C = \frac{dz_1 dz_2 dz_3}{z_4}$$

satisfy $\frac{\omega_C^3}{6} = \frac{i}{8} \Omega_C \wedge \bar{\Omega}_C$ and therefore define a Calabi–Yau cone structure on C .

The nodal singularity $\{(z_1 + iz_2)(z_1 - iz_2) + (z_3 + iz_4)(z_3 - iz_4) = 0\}$ is birationally equivalent to its small resolution

$$\left\{ z_j \in \mathbb{C}, z \in \mathbb{C} \cup \{\infty\} : z_1^2 + z_2^2 + z_3^2 + z_4^2 = 0, z = \frac{z_1 + iz_2}{z_3 + iz_4} = -\frac{z_3 - iz_4}{z_1 - iz_2} \right\}.$$

It replaces the tip point by $\mathbb{CP}^1 = \mathbb{C} \cup \{\infty\}$. The other small resolution is given by

$$\left\{ z_j \in \mathbb{C}, z \in \mathbb{C} \cup \{\infty\} : z_1^2 + z_2^2 + z_3^2 + z_4^2 = 0, z = \frac{z_1 + iz_2}{z_3 - iz_4} = -\frac{z_3 + iz_4}{z_1 - iz_2} \right\}.$$

Both of the small resolutions are Calabi–Yau threefolds by [6].

It is easy to see that C is diffeomorphic to the cone over $\mathbb{S}^2 \times \mathbb{S}^3$. The deformation $\{z_1^2 + \cdots z_4^2 = \epsilon\}$ is diffeomorphic to $\mathbb{R}^3 \times \mathbb{S}^3$. Both of the small resolutions are diffeomorphic to $\mathbb{S}^2 \times \mathbb{R}^4$.

The following proposition is well-known:

Proposition 2.8 *For any Calabi–Yau threefold $(M, g, J, \omega, \Omega)$, define an $\mathbb{S}^1$ -invariant G_2 structure φ on $\mathbb{S}^1 \times M$ by $\varphi = d\theta \wedge \omega + \operatorname{Re}\Omega$, where θ is the standard coordinate on $\mathbb{S}^1$. It satisfies $d\varphi = d\Theta(\varphi) = 0$. On the other hand, any $\mathbb{S}^1$ -invariant G_2 structure on $\mathbb{S}^1 \times M$ satisfying $d\varphi = d\Theta(\varphi) = 0$ must come from a Calabi–Yau threefold structure on M .*

The next part is the definition of an asymptotically cylindrical Kähler manifold with conical singularities.

Definition 2.9 Let (F, g_F) be a compact Riemannian manifold. Then $C(F)$ is the set $((0, \infty) \times F) \cup \{o\}$. Let r be the coordinate on $(0, \infty)$ and define $r(o) = 0$. Then the cone metric $g_{C(F)}$ is defined by $g_{C(F)} = dr^2 + r^2 g_F$. Similarly, for any compact Riemannian manifold (F_∞, g_{F_∞}) , define the product metric on $\mathbb{R} \times F_\infty$ by $g_\infty = dt^2 + g_{F_\infty}$, where t is the coordinate on $\mathbb{R}$. The Kähler structure $J_{C(F)}$ or J_∞ is required to be compatible with $g_{C(F)}$ or g_∞ and commutes with the map $r \rightarrow cr$ or $t \rightarrow t + c$ for constants c . The Kähler form $\omega_{C(F)}$ or ω_∞ is required to be compatible with $g_{C(F)}$ and $J_{C(F)}$ or g_∞ and J_∞ . Moreover, $L_{\frac{\partial}{\partial r}} \omega_{C(F)} = 2\omega_{C(F)}$ and $L_{\frac{\partial}{\partial t}} \omega_\infty = 0$, where L means the Lie derivative. The $(n, 0)$ -forms $\Omega_{C(F)}$, Ω_∞ and the G_2 structures $\varphi_{C(F)}$, φ_∞ on $C(F)$ or $\mathbb{R} \times F_\infty$ are required to have similar properties.

Definition 2.10 A Calabi–Yau cone C with smooth cross-section and with Ricci-flat Kähler cone metric $\omega_C = \frac{i}{2} \partial \bar{\partial} r^2$ is regular if its Reeb field, i.e. the holomorphic Killing field $J(r \frac{\partial}{\partial r})$, generates a free $\mathbb{S}^1$ -action on $C \setminus \{o\}$. This exhibits C as the blow-down of the zero section of $\frac{1}{q} K_B$ for some Kähler–Einstein Fano manifold B and $q \in \mathbb{N}$. C is called strongly regular if $-\frac{1}{q} K_B$ is very ample.

Remark 2.11 Example 2.6 is strongly regular as defined in Definition 2.10.

Definition 2.12 Let V be a manifold with Kähler metric ω . $x \in V$ is called a conical singularity with rate $\nu_x > 0$ and tangent cone (C_x, ω_{C_x}) with respect to ω if there exist a Kähler metric cone $(C_x, J_{C_x}, \omega_{C_x})$ and a biholomorphism $P_x : O_x \rightarrow U_x$ between neighborhoods $o \in O_x \subset C_x$ and $x \in U_x \subset V$ such that

$$\left| \nabla_{\omega_{C_x}}^j (P^* \omega - \omega_{C_x}) \right|_{\omega_{C_x}} = O\left(r^{\nu_x - j}\right)$$

as $r \rightarrow 0$ for all $j \in \mathbb{N}_0$. Assume that the set $\{r \leq r_{0,x}\}$ is contained in O_x .

Definition 2.13 Let V be a manifold with a metric g . g is called asymptotically cylindrical with rate $\nu_\infty > 0$ and cross-section F_∞ if there exist a set U_∞ and a diffeomorphism $P_\infty : [1, \infty) \times F_\infty \rightarrow U_\infty$ such that $V \setminus U_\infty$ is bounded and

$$\left| \nabla_{g_\infty}^j (P^* g - g_\infty) \right|_{g_\infty} = O\left(e^{-\nu_\infty t}\right)$$

as $t \rightarrow \infty$ for all $j \in \mathbb{N}_0$. The asymptotically cylindrical almost complex structure J , Kähler form ω , $(n, 0)$ -form Ω and G_2 structure φ are defined similarly.

The next goal is to describe the analysis on an asymptotically cylindrical Kähler manifold V with conical singularities following Lockhart and McOwen [47]. Remark that in this paper, δ is multiplied by -1 and λ is divided by i compared to [47]. The same result was obtained by Melrose and Mendoza [53] independently.

Definition 2.14 Assume that V is asymptotically cylindrical with conical singularities at V^{sing} . Assume that U_x and U_∞ are disjoint. Let N_2 be the number of points in V^{sing} . Assume that $\nu > 0$ is smaller than the minimum of $\{\nu_1, \dots, \nu_{N_2}, \nu_\infty\}$. For any $x \in V^{\text{sing}}$, choose r_x as a smooth function with range $[0, 2r_{0,x}]$ such that $r_x = 2r_{0,x}$ outside U_x and $r_x = r$ when $r \leq r_{0,x}$. Extend t to a non-negative smooth function on V such that t equals to 0 on U_x for all $x \in V^{\text{sing}}$. For any $\delta = (\delta_1, \dots, \delta_{N_2}, \delta_\infty) \in \mathbb{R}^{N_2+1}$, using the metric ω , define the weighted L^2 space L_δ^2 by

$$\|\gamma\|_{L_\delta^2} = \left(\int_V \left| \prod_{i=1}^{N_2} r_{x_i}^{-\delta_i} e^{\delta_\infty t} \gamma \right|^2 \prod_{i=1}^{N_2} r_{x_i}^{-2n} \right)^{\frac{1}{2}},$$

where n is the complex dimension of V . Assume that k is a large enough integer. Define the weighted Hilbert space by

$$\|\gamma\|_{W_\delta^{k,2}} = \left(\sum_{j=0}^k \|\nabla^j \gamma\|_{L_{(\delta_1-j, \dots, \delta_{N_2}-j, \delta_\infty)}^2}^2 \right)^{\frac{1}{2}}.$$

Roughly speaking, it means γ has rate $r_{x_i}^{\delta_i}$ near each $x_i \in V^{\text{sing}}$ and rate $e^{-\delta_\infty t}$ on the end. In general, one can define $W_\delta^{k,p}$ and $C_\delta^{k,\alpha}$ spaces. However, the $W_\delta^{k,2}$ space is enough for this paper. Define the space C_δ^∞ as the intersection of $W_\delta^{k,2}$ for all k . Choose $\chi = \chi(s) : \mathbb{R} \rightarrow [0, 1]$ as a smooth function satisfying $\chi(s) = 1$ for $s \leq 1$ and $\chi(s) = 0$ for $s \geq 2$. Denote $\chi(\frac{2r_{x_i}}{r_{0,x_i}})$ by χ_i . Denote $1 - \chi(t)$ by χ_∞ . It is also useful to consider spaces like

$$\left(\bigoplus_{i=1}^{N_2} \mathbb{R} \chi_i \right) \oplus W_\delta^{k,2}$$

for positive δ_i . For constants $c_1, \dots, c_{N_2}$, define

$$\left\| \sum_{i=1}^{N_2} c_i \chi_i + \gamma \right\|_{\left(\bigoplus_{i=1}^{N_2} \mathbb{R} \chi_i \right) \oplus W_\delta^{k,2}} = \sum_{i=1}^{N_2} |c_i| + \|\gamma\|_{W_\delta^{k,2}}.$$

Now let D be the Laplacian operator $\Delta = dd^* + d^*d$ or the operator $d + d^*$ acting on the direct sum of all odd degree forms. Then the order m of D is 2 or 1 respectively.

Definition 2.15 $\lambda_i \in \mathbb{C}$ is called a critical rate for D near x_i if there exists

$$\gamma = \sum_{p=0}^{p_{i,\lambda_i}} e^{\lambda_i \log r} (-\log r)^p \gamma_{i,\lambda_i,p}$$

in the kernel of $D_{C_{x_i}}$ for non-zero $\gamma_{i,\lambda_i,p_{i,\lambda_i}}$. $\lambda_\infty \in \mathbb{C}$ is called a critical rate for D near infinity if there exists

$$\gamma = \sum_{p=0}^{p_{\infty,\lambda_\infty}} e^{-\lambda_\infty t} t^p \gamma_{\infty,\lambda_\infty,p}$$

in the kernel of $D_{[0,\infty) \times F_\infty}$ for non-zero $\gamma_{\infty,\lambda_\infty,p_{\infty,\lambda_\infty}}$. Define $\mathcal{K}_i(\lambda_i)$ or $\mathcal{K}_\infty(\lambda_\infty)$ as the space of such γ . Define the multiplicity $d_i(\lambda_i)$ or $d_\infty(\lambda_\infty)$ as the dimension of $\mathcal{K}_i(\lambda_i)$ or $\mathcal{K}_\infty(\lambda_\infty)$. $\delta_i \in \mathbb{R}$ or $\delta_\infty \in \mathbb{R}$ is called critical near x_i or infinity if it is the real part of a critical λ_i or λ_∞ . δ is called critical if either at least one of the δ_i is critical near x_i or δ_∞ is critical near infinity. For non-critical weights $\delta_i > \delta'_i$, define $N(\delta', \delta)$ by

$$N(\delta', \delta) = \sum_{i=1}^{N_2} \sum_{\delta'_i < \operatorname{Re}(\lambda''_i) < \delta_i} d_i(\lambda''_i) + \sum_{\delta'_\infty < \operatorname{Re}(\lambda''_\infty) < \delta_\infty} d_\infty(\lambda''_\infty).$$

The main theorem of Lockhart and McOwen [47] is the following:

Theorem 2.16 $D : W_{\delta}^{k,2} \rightarrow W_{(\delta_1-m, \dots, \delta_{N_2}-m, \delta_\infty)}^{k-m,2}$ is Fredholm if and only if δ is non-critical. Moreover, the Fredholm index $i_\delta(\Delta)$ is independent of k . For non-critical weights $\delta_i > \delta'_i$ and $\delta_\infty > \delta'_\infty$, $i_{\delta'}(\Delta) - i_\delta(\Delta) = N(\delta', \delta)$.

In the first paragraph of p. 420 of [47], Lockhart and McOwen used the following theorem of [42,50]:

Theorem 2.17 The operator $D_{C_{x_i}} : W_{\delta_i}^{k,2}(C_{x_i}) \rightarrow W_{\delta_i-m}^{k-m,2}(C_{x_i})$ or the operator $D_\infty : W_{\delta_\infty}^{k,2}(\mathbb{R} \times F_\infty) \rightarrow W_{\delta_\infty-m}^{k-m,2}(\mathbb{R} \times F_\infty)$ is an isomorphism for non-critical δ_i or non-critical δ_∞ .

In Sect. 5 of [47], Lockhart and McOwen used the following theorem of [1,42,50]:

Theorem 2.18 Suppose that $\delta_i > \delta'_i$ or $\delta_\infty > \delta'_\infty$ are non-critical, then for any $\gamma \in W_{\delta'_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$ or $\gamma \in W_{\delta'_\infty}^{k,2}([1, \infty) \times F_\infty)$ satisfying $D_{C_{x_i}}\gamma = 0$ or $D_\infty\gamma = 0$, there exists γ' in the direct sum of $\mathcal{K}_i(\lambda''_i)$ for all $\delta'_i < \operatorname{Re}(\lambda''_i) < \delta_i$ or $\mathcal{K}_\infty(\lambda''_\infty)$ for all $\delta'_\infty < \operatorname{Re}(\lambda''_\infty) < \delta_\infty$ such that $\gamma - \gamma' \in W_{\delta_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$ or $W_{\delta_\infty}^{k,2}([1, \infty) \times F_\infty)$. Moreover,

$$\|\gamma\| \left(\oplus_{\delta'_i < \operatorname{Re}(\lambda''_i) < \delta_i} \chi_i \mathcal{K}_i(\lambda''_i) \right) \oplus W_{\delta'_i}^{k,2} \leq C \|\gamma\|_{W_{\delta'_i}^{k,2}}$$

or

$$\|\gamma\|_{(\oplus_{\delta'_\infty < \operatorname{Re}(\lambda''_\infty) < \delta_\infty} \chi_\infty \mathcal{K}_\infty(\lambda''_\infty)) \oplus W_{\delta_\infty}^{k,2}} \leq C \|\gamma\|_{W_{\delta'_\infty}^{k,2}}.$$

Let ν be ν_x in Definition 2.12 or ν_∞ in Definition 2.13. For all critical λ_i or λ_∞ , choose a non-critical weight $\delta_i \in (\operatorname{Re}(\lambda_i), \operatorname{Re}(\lambda_i) + \nu)$ or a non-critical weight $\delta_\infty \in (\operatorname{Re}(\lambda_\infty), \operatorname{Re}(\lambda_\infty) + \nu)$. By Theorem 2.17, the maps

$$D_{C_{x_i}} : W_{\delta_i}^{k,2}(C_{x_i}) \rightarrow W_{\delta_i-m}^{k-m,2}(C_{x_i})$$

and

$$D_\infty : W_{\delta_\infty}^{k,2}(\mathbb{R} \times F_\infty) \rightarrow W_{\delta_\infty-m}^{k-m,2}(\mathbb{R} \times F_\infty)$$

are isomorphisms. Therefore, after shrinking r_{0,x_i} or replacing t by $t + T$ if necessary, the maps

$$\chi_i D + (1 - \chi_i) D_{C_{x_i}} : W_{\delta_i}^{k,2}(C_{x_i}) \rightarrow W_{\delta_i-m}^{k-m,2}(C_{x_i})$$

and

$$\chi_\infty D + (1 - \chi_\infty) D_\infty : W_{\delta_\infty}^{k,2}(\mathbb{R} \times F_\infty) \rightarrow W_{\delta_\infty-m}^{k-m,2}(\mathbb{R} \times F_\infty)$$

are also isomorphisms. Remark that for each $\gamma \in \mathcal{K}_i(\lambda_i)$ or $\gamma \in \mathcal{K}_\infty(\lambda_\infty)$, $-D(\gamma \chi_i) \in W_{\delta_i-m}^{k-m,2}$ or $-D(\gamma \chi_\infty) \in W_{\delta_\infty-m}^{k-m,2}$. Let

$$\mathcal{P}_i(\lambda_i) = \{\gamma + (\chi_i D + (1 - \chi_i) D_{C_{x_i}})^{-1}(-D(\gamma \chi_i)), \gamma \in \mathcal{K}_i(\lambda_i)\}$$

and

$$\mathcal{P}_\infty(\lambda_\infty) = \{\gamma + (\chi_\infty D + (1 - \chi_\infty) D_\infty)^{-1}(-D(\gamma \chi_\infty)), \gamma \in \mathcal{K}_\infty(\lambda_\infty)\}.$$

Then $D\gamma = 0$ near x_i or infinity for any γ in $\mathcal{P}_i(\lambda_i)$ or $\mathcal{P}_\infty(\lambda_\infty)$. Moreover, for any γ in $\mathcal{K}_i(\lambda_i)$ or $\mathcal{K}_\infty(\lambda_\infty)$, there exists a unique element γ' in $\mathcal{P}_i(\lambda_i)$ or $\mathcal{P}_\infty(\lambda_\infty)$ such that $\gamma' - \gamma \in W_{\delta_i}^{k,2}(C_{x_i})$ or $W_{\delta_\infty}^{k,2}(\mathbb{R} \times F_\infty)$. Conversely, for any γ' in $\mathcal{P}_i(\lambda_i)$ or $\mathcal{P}_\infty(\lambda_\infty)$, there exists a unique element γ in $\mathcal{K}_i(\lambda_i)$ or $\mathcal{K}_\infty(\lambda_\infty)$ such that $\gamma' - \gamma \in W_{\delta_i}^{k,2}(C_{x_i})$ or $W_{\delta_\infty}^{k,2}(\mathbb{R} \times F_\infty)$. Remark that the definition of $\mathcal{P}_i(\lambda_i)$ or $\mathcal{P}_\infty(\lambda_\infty)$ can be changed to any set satisfying the properties above.

The following theorem can be proved using Theorems 2.17 and 2.18:

Theorem 2.19 *Suppose that $\delta_i > \delta'_i$ or $\delta_\infty > \delta'_\infty$ are non-critical, then for any $\gamma \in W_{\delta'_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$ satisfying $D\gamma \in W_{\delta_i-m}^{k-m,2}$ or $\gamma \in W_{\delta'_\infty}^{k,2}([1, \infty) \times F_\infty)$ satisfying $D\gamma \in W_{\delta_\infty-m}^{k-m,2}$, there exists γ' in the direct sum of $\chi_i \mathcal{P}_i(\lambda''_i)$ for all critical*

$\delta'_i < \operatorname{Re}(\lambda''_i) < \delta_i$ or $\chi_\infty \mathcal{P}_\infty(\lambda''_\infty)$ for all critical $\delta'_\infty < \operatorname{Re}(\lambda''_\infty) < \delta_\infty$ such that $\gamma - \gamma' \in W_{\delta_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$ or $W_{\delta_\infty}^{k,2}([1, \infty) \times F_\infty)$. Moreover,

$$\|\gamma\|_{(\oplus_{\delta'_i < \operatorname{Re}(\lambda''_i) < \delta_i} \chi_i \mathcal{P}_i(\lambda''_i)) \oplus W_{\delta_i}^{k,2}} \leq C(\|\gamma\|_{W_{\delta'_i}^{k,2}} + \|D\gamma\|_{W_{\delta_i-m}^{k-m,2}})$$

or

$$\|\gamma\|_{(\oplus_{\delta'_\infty < \operatorname{Re}(\lambda''_\infty) < \delta_\infty} \chi_\infty \mathcal{P}_\infty(\lambda''_\infty)) \oplus W_{\delta_\infty}^{k,2}} \leq C(\|\gamma\|_{W_{\delta'_\infty}^{k,2}} + \|D\gamma\|_{W_{\delta_\infty-m}^{k-m,2}}).$$

Proof The proof of this theorem is essentially due to Lockhart and McOwen [47]. There were lots of closely related theorems due to many authors, for example, Proposition 4.21 of [40] and Proposition 2.9 of [36]. For the reader's convenience, a proof is included here without claiming any originality.

Suppose that this theorem is not true. Choose δ_i as a non-critical value for the failure of this theorem such that it is smaller than the infimum of all such δ_i plus $\frac{\nu}{2}$. Then choose a non-critical $\delta'_i \in (\delta_i - \nu, \delta_i - \frac{\nu}{2})$. Assume that $\gamma \in W_{\delta'_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$ and $D\gamma \in W_{\delta_i-m}^{k-m,2}$. Then $D\gamma \in W_{\delta'_i-m}^{k-m,2}$. So there exists γ'' in the direct sum of $\chi_i \mathcal{P}_i(\lambda_i)$ for all $\delta'_i < \operatorname{Re}(\lambda_i) < \delta'_i$ such that $\gamma - \gamma'' \in W_{\delta'_i}^{k,2}(C_{x_i} \cap \{r \leq r_{0,x_i}\})$. Since $D\gamma''$ vanishes on a neighborhood of x_i , it is easy to see that $D(\gamma - \gamma'') \in W_{\delta_i-m}^{k-m,2}$. So it is also true that $D_{C_{x_i}}(\gamma - \gamma'') \in W_{\delta_i-m}^{k-m,2}$. Remark that $\gamma - \gamma'' - D_{C_{x_i}}^{-1}(\chi_i(\gamma - \gamma''))$ satisfies the hypothesis of Theorem 2.18 in a smaller neighborhood of x_i . So it can be written as the sum of elements in $\chi_i \mathcal{K}_i(\lambda_i)$ for all $\operatorname{Re}(\lambda_i) \in (\delta'_i, \delta_i)$ plus an element in $W_{\delta_i}^{k,2}$. The construction of $\mathcal{P}_i(\lambda_i)$ provides a contradiction to the definition of δ_i . The argument near infinity is similar. $\square$

The following corollary is a refinement of Theorem 2.16:

Corollary 2.20 Suppose that $\delta_i > \delta'_i$ or $\delta_\infty > \delta'_\infty$ are non-critical, then the induced maps

$$D : \left(\bigoplus_{\substack{i=1, \dots, N_2, \infty \\ \delta_i > \operatorname{Re}(\lambda_i) > \delta'_i}} \chi_i \mathcal{P}_i(\lambda_i) \right) \oplus W_{\delta}^{k,2} \rightarrow W_{(\delta_1-m, \dots, \delta_{N_2}-m, \delta_\infty)}^{k-m,2}$$

and

$$D : W_{\delta'}^{k,2} \rightarrow W_{(\delta'_1-m, \dots, \delta'_{N_2}-m, \delta'_\infty)}^{k-m,2}$$

commute with the inclusion maps from $(\bigoplus_{\substack{i=1, \dots, N_2, \infty \\ \delta_i > \operatorname{Re}(\lambda_i) > \delta'_i}} \chi_i \mathcal{P}_i(\lambda_i)) \oplus W_{\delta}^{k,2}$ to $W_{\delta'}^{k,2}$ and from $W_{(\delta_1-m, \dots, \delta_{N_2}-m, \delta_\infty)}^{k-m,2}$ to $W_{(\delta'_1-m, \dots, \delta'_{N_2}-m, \delta'_\infty)}^{k-m,2}$. Moreover, the inclusion maps induce isomorphisms on the kernels and cokernels of D .

Proof It is trivial to see that the inclusion map commutes with D and induces maps between kernels of D or between cokernels of D . It is also trivial to see that the induced map is injective between the kernels of D . Suppose that $\gamma \in W_{\delta'}^{k+m,2}$ and $D\gamma = 0$. By Theorem 2.19, it is in the image of the map between the kernels induced by the inclusion map. On the other hand, for any element $\gamma \in W_{(\delta'_1-m, \dots, \delta'_{N_2}-m, \delta'_\infty)}^{k-m,2}$, define γ' by

$$\gamma' = \sum_{i=1}^{N_2} \chi_i (\chi_i D + (1 - \chi_i) D_{C_{x_i}})^{-1} (\chi_i \gamma) + (\chi_\infty D + (1 - \chi_\infty) D_\infty)^{-1} (\chi_\infty \gamma).$$

Then $\gamma' \in W_{\delta'}^{k,2}$. Moreover, $\gamma - D\gamma'$ vanishes near x_i and infinity. This proves the surjectivity of the map between cokernels induced by the inclusion map. The injectivity of this map is an immediate corollary of Theorem 2.19. $\square$

For the Laplacian operator acting on functions, as in Proposition 2.9 of [36], any element γ in $\mathcal{K}_i(\lambda_i)$ or $\mathcal{K}_\infty(\lambda_\infty)$ can be written as a generalized Fourier series. Then the following proposition follows easily from the explicit solution of the ordinary differential equation as well as Theorem 2.14 of [36]:

Proposition 2.21 *For the Laplacian operator acting on functions, the following statements are true:*

- (1) Any critical $\lambda_i \in \mathbb{C}$ or $\lambda_\infty \in \mathbb{C}$ is in fact real.
- (2) $\mathcal{K}_i(\lambda_i)$ has no $(-\log r)^p$ terms.
- (3) $\mathcal{K}_\infty(0) = \text{Span}\{1, t\}$.
- (4) There is no critical rate in $(-2n + 2, 0)$ near x_i .
- (5) $\mathcal{K}_i(0) = \text{Span}\{1\}$.
- (6) There is no critical rate in $(0, 1]$ near x_i .
- (7) If $\lambda_i \in (1, 2)$, then any element in $\mathcal{K}_i(\lambda_i)$ is pluriharmonic.
- (8) Any element in $\mathcal{K}_i(2)$ can be written as a pluriharmonic function in $\mathcal{K}_i(2)$ plus a $J(r \frac{\partial}{\partial r})$ -invariant function in $\mathcal{K}_i(2)$.
- (9) Denote the direct sum of pluriharmonic functions in $\mathcal{K}_i(\lambda_i)$ with rates $\lambda_i \in (1, 2]$ by $\mathcal{P}_i$. In Corollary 2.20, $\mathcal{P}_i(\lambda_i)$ can be replaced by the corresponding $\mathcal{P}_i$ because any pluriharmonic function is harmonic with respect to both $\omega_{C_{x_i}}$ and ω . Denote the space of $J(r \frac{\partial}{\partial r})$ -invariant functions in $\mathcal{K}_i(2)$ by $\mathcal{H}_i$. Remark that the difference between $\mathcal{H}_i$ and corresponding element in $\mathcal{P}_i(2)$ lies in $W_{2+\frac{\nu}{2}}^{k,2}$ near x_i .

The next goal is the analysis on manifolds with edge singularities. It is required that the smooth part of the manifold can be viewed as a manifold with boundary and the boundary is a fibration. As a special case, assume that the boundary is the trivial fibration over $\mathbb{S}^1$ with fiber F , in other words, is $F \times \mathbb{S}^1$. So in this special case, the singular manifold looks like $C(F) \times \mathbb{S}^1$ locally.

In the pioneering work of Mazzeo [51], the weighted Sobolev space is defined using the same formula as in Definition 2.14. Elliptic differential operators like $\Delta : W_\delta^{k,2} \rightarrow W_{\delta-2}^{k-2,2}$ are studied in [51]. The main result of [51] discusses whether the

elliptic differential operator is Fredholm or not. On the other hand, Theorem 7.14 of [51] is closely related to Theorem 2.19.

Remark that there are different versions of weighted analysis by changing the definition slightly. One way is to change the domain and range of Δ as Cheeger [8] and Hunsicker and Mazzeo [37] did when studying Hodge theory on manifolds with edge singularities. The other way is to change the definition of $C_\delta^{k,\alpha}$ as Chen et al. [15–17] did when using manifolds with edge singularities to study the Kähler–Einstein problem.

3 Asymptotically Cylindrical Calabi–Yau Manifolds with Isolated Conical Singularities

In this section, Theorem 1.1 is proved as a combination of [36] and [34]. There are several major technical problems in this process. Firstly, it is necessary to find a good substitute for the finiteness of diameter property in [36]. Secondly, it is not clear how to get the generalization of [28] on the existence of weak solutions because the weak solution in the sense of current is too weak to apply the standard analysis for the asymptotically cylindrical manifolds. Thirdly, the openness part of [36] uses a non-standard weighted analysis and therefore does not have a simple generalization to the non-compact case. Finally, as personally communicated to the author by Hein–Sun, the proof of Proposition 3.2 of [36] requires a little more explanation. In fact, one needs to show that on the central fiber we can assume a priori the existence of $K_1 > 0$ such that the K -inequalities hold for $s = 0$, $t \in [0, 1]$ with this K_1 , and one needs to take K to be bigger than this K_1 . In the setting studied in the paper of Hein–Sun, this follows directly from the results of [28] since one can work on the fixed variety X_0 . In our setting, it is necessary to find out solutions to the above problems.

In this section, the notations in Theorem 1.1 are still used. As in [34], near S_s , let (t, ϑ) be the coordinates on $[T, \infty) \times \mathbb{S}^1$ such that f can be written as $e^{t+i\vartheta}$ on V_s near S_s . Using the diffeomorphisms between fibers of f near S_s , there exists a smooth family of local diffeomorphisms between $f^{-1}(\{|\frac{1}{f}| \leq e^{-T}\}) \cap Z_s$ and $\{|\frac{1}{f}| < e^{-T}\} \times S_s$. Its restriction yields a smooth family of diffeomorphisms

$$\Phi_{\infty,s} : [T, \infty) \times \mathbb{S}^1 \times S_s \rightarrow U_{\infty,s} = f^{-1} \left(\left\{ 0 < \left| \frac{1}{f} \right| \leq e^{-T} \right\} \right) \cap Z_s \subset V_s.$$

Using Yau’s solution to the Calabi conjecture, it is easy to choose a holomorphic family of nowhere vanishing holomorphic $(n-1)$ -forms Ω_{S_s} and a family of Ricci-flat metrics $\omega_{S_s} \in [c_1(L)|_{S_s}]$ on S_s such that after replacing Ω_s , $\Omega_{\infty,s}$ by their products with constants,

$$(\omega_{\infty,s}, \Omega_{\infty,s}) = (dt \wedge d\vartheta + \omega_{S_s}, (d\vartheta - idt) \wedge \Omega_{S_s})$$

satisfy $\omega_{\infty,s}^n = i^{n^2} \Omega_{\infty,s} \wedge \bar{\Omega}_{\infty,s}$ and $\Phi_{\infty,s}^* \Omega_s = \Omega_{\infty,s} + d\zeta_s$ with

$$|\nabla_{\omega_{\infty,s}}^j \zeta_s|_{\omega_{\infty,s}} \leq C_j e^{-\nu_{\infty} t}$$

for all $t > T$ and $j \in \mathbb{N}_0$, where ν_{∞} , C_j , T are positive constants independent of s . Remark that on the cylinder $[T, \infty) \times \mathbb{S}^1 \times S_s$, any exponentially decaying closed form can be written as d of an exponentially decaying form.

The next goal is to construct a family of background metrics $\hat{\omega}_s$. Using Part 1 of Sect. 4.2 of [34], the pull back of Fubini–Study metric can be modified to a smooth family of metrics $\hat{\omega}_s$ by adding a smooth family of $i\partial\bar{\partial}$ exact forms supported in $U_{\infty,s}$ such that $\Phi_{\infty,s}^* \hat{\omega}_s = \omega_{\infty,s} + d\varrho_s$ with $|\nabla_{\omega_{\infty,s}}^j \varrho_s|_{\omega_{\infty,s}} \leq C_j (e^{-\nu_{\infty} t})$ for all $t > T$ and $j \in \mathbb{N}_0$ after modifying the positive constants ν_{∞} , C_j and T if necessary. Moreover, $\hat{\omega}_s$ can be chosen to satisfy

$$\int_V \frac{\hat{\omega}_s^n}{n!} - \frac{i^{n^2}}{2^n} \Omega_s \wedge \bar{\Omega}_s = 0.$$

Remark that the assumption that Ω_s are comparable to explicit n -forms $\Omega_{\epsilon_x(s)}$ near $x \in V^{\text{sing}}$ means that the singularities are harmless.

By assumption, for each $x_i \in V^{\text{sing}}$, there exists a local biholomorphism $\Phi_{x_i} : \mathcal{U}_{x_i} \rightarrow \mathcal{O}_{x_i}$ between a neighborhood $\mathcal{U}_{x_i}$ of x_i in $\mathcal{Z}$ and a neighborhood $\mathcal{O}_{x_i}$ of the vertex o_{x_i} in $\mathcal{C}_{x_i}$. Assume that $\mathcal{U}_{x_i}$ and $U_{\infty,s}$ are disjoint. As in Proposition 2.4 of [2], there exists a bounded family of functions $\psi_{1,s}$ such that $\psi_{1,s}$ is smooth outside x_i and $\hat{\omega}_s + i\partial\bar{\partial}\psi_{1,s}$ equals to $\Phi_{x_i}^* \omega_{\epsilon_{x_i}(s)}$ after shrinking the neighborhood $\mathcal{U}_{x_i}$ of x_i . Assume that $\psi_{1,s}$ is supported in $\mathcal{U}_{x_i}$ before the shrinking of $\mathcal{U}_{x_i}$. In particular, the condition

$$\int_V \frac{(\hat{\omega}_s + i\partial\bar{\partial}\psi_{1,s})^n}{n!} - \frac{i^{n^2}}{2^n} \Omega_s \wedge \bar{\Omega}_s = 0$$

is still true.

Assume that ν is small enough. Let F_s be the function in $(\oplus_{i=1}^{N_2} \mathbb{R}\chi_i) \oplus C_{\nu, \dots, \nu, \nu}^{\infty}$ satisfying

$$\frac{(\hat{\omega}_s + i\partial\bar{\partial}\psi_{1,s})^n}{n!} = e^{F_s} \frac{i^{n^2}}{2^n} \Omega \wedge \bar{\Omega}.$$

F_s is the real part of a holomorphic function on $\mathcal{U}_{x_i}$. Denote $\mathcal{U}_{x_i} \cap Z_s$ by $U_{x_i,s}$. Then it is easy to find a family of functions $F_{\tau,s}$ on $V_s = Z_s \setminus S_s$ for all $\tau \in [0, 1]$ continuous in $(\oplus_{i=1}^{N_2} \mathbb{R}\chi_i) \oplus C_{\nu, \dots, \nu, \nu}^{\infty}$ topology such that $F_{\tau,s} = \tau F_s$ on $U_{x_i,s}$, $F_{1,s} = F_s$ on $x \in \mathcal{Z}$, $F_{0,s} = 0$ and

$$\int_{V_s} (e^{F_{\tau,s}} - 1) \frac{i^{n^2}}{2^n} \Omega_s \wedge \bar{\Omega}_s = 0.$$

As in [34,36], define $\mathcal{T} \subset [0, 1]$ as the set of τ such that

$$\frac{(\hat{\omega}_0 + i\partial\bar{\partial}\psi_{\tau,0})^n}{n!} = e^{F_{\tau,0}} \frac{i^{n^2}}{2^n} \Omega \wedge \bar{\Omega}$$

has a bounded and smooth solution $\psi_{\tau,0}$ on $V \setminus V^{\text{sing}}$ such that $\psi_{\tau,0} \in C_v^\infty$ near infinity and $\hat{\omega}_0 + i\partial\bar{\partial}\psi_{\tau,0}$ has conical singularity at each $x_i \in V^{\text{sing}}$ with tangent cone the same as Example 2.7. It is clear that $1 \in \mathcal{T}$.

The openness of $\mathcal{T}$ can be proved using the following proposition as in [36]:

Proposition 3.1 *The Laplacian operator is a bijective map between the set*

$$W_{(2+v,\dots,2+v,v)}^{k,2} \oplus \left(\oplus_{i=1}^{N_2} \chi_i \left(\mathcal{P}_i \oplus \mathbb{R}r_{x_i}^2 \oplus \mathcal{H}_i \right) \right)$$

and the set

$$\left\{ \gamma \in W_{v,\dots,v,v}^{k-2,2} \oplus \left(\oplus_{i=1}^{N_2} \mathbb{R}\chi_i \right) : \int_V \gamma \omega^n = 0 \right\}.$$

Proof Remark that $\mathcal{H}_i$ and $\mathcal{P}_i$ were defined in Proposition 2.21.

In the compact case, this proposition was proved in Hein–Sun’s paper [36] using a non-standard weighted norm. It is not clear how to do the analogue here. However, this proposition can be proved using Corollary 2.20 and Proposition 2.21 as an analogue of Proposition 2.7 of [34]:

Consider the Laplacian operator acting on functions from $W_{(-n+1,\dots,-n+1,v)}^{k,2}$ to $W_{(-n+1,\dots,-n+1,v)}^{k-2,2}$. Remark that the weight is non-critical by Proposition 2.21 if v is small enough. The dual operator is the Laplacian operator acting on functions from $W_{(-n+1,\dots,-n+1,-v)}^{2-k,2}$ to $W_{(-n+1,\dots,-n+1,-v)}^{-k,2}$. Thus

$$i_{(-n+1,\dots,-n+1,v)}(\Delta) = -i_{(-n+1,\dots,-n+1,-v)}(\Delta).$$

Therefore, by Theorem 2.16 and Proposition 2.21,

$$i_{(-n+1,\dots,-n+1,v)}(\Delta) = -\frac{1}{2} \sum_{-v < \text{Re}(\lambda_\infty) < v} d(\lambda_\infty) = -1.$$

Using Corollary 2.20 and Proposition 2.21, the index of

$$\Delta : W_{(2+v,\dots,2+v,v)}^{k,2} \oplus \left(\oplus_{i=1}^{N_2} \chi_i (\mathcal{P}_i \oplus \mathcal{H}_i) \right) \rightarrow W_{v,\dots,v,v}^{k-2,2}$$

is also -1. Using the fact that $-\Delta(\chi_i r_{x_i}^2) - 4n\chi_i \in W_{v,\dots,v,v}^{k-2,2}$, it is easy to see that the index of

$$\Delta : W_{(2+v,\dots,2+v,v)}^{k,2} \oplus \left(\oplus_{i=1}^{N_2} \chi_i (\mathcal{P}_i \oplus \mathbb{R}r_{x_i}^2 \oplus \mathcal{H}_i) \right) \rightarrow W_{v,\dots,v,v}^{k-2,2} \oplus \left(\oplus_{i=1}^{N_2} \mathbb{R}\chi_i \right)$$

is also -1 . Remark that for any function ψ in the kernel of this operator, the decay condition of ψ near x_i and infinity insures that the boundary term in the integral $\int |\nabla \psi|^2 - \psi \Delta \psi$ vanishes. This implies that $\nabla \psi = 0$, so $\psi = 0$ by the decay condition near infinity. Another integration by parts shows that the integral of any function in the image is 0. Now, this proposition is an immediate corollary of the fact that the index is -1 . $\square$

Now assume that $\{\tau_i\} \subset \mathcal{T} \rightarrow \tau_\infty$. It suffices to show that $\tau_\infty \in \mathcal{T}$. Using Theorem 4.1 of [34], for all $s \neq 0$, there exists $\psi_{\tau,s} \in C^\infty_v$ such that

$$\frac{(\hat{\omega}_s + i\partial\bar{\partial}\psi_{\tau,s})^n}{n!} = e^{F_{\tau,s}} \frac{i^{n^2}}{2^n} \Omega_s \wedge \bar{\Omega}_s.$$

The goal is to obtain uniform estimates on ψ_{τ_j,s_j} for a sequence $s_j \rightarrow 0$ so that the limit is expected to be $\psi_{\tau_\infty,0}$.

The starting point is the C^0 -estimate as in Step 1 of [34]. It requires an estimate using $\omega_{1,s} = \hat{\omega}_s + i\partial\bar{\partial}\psi_{1,s}$ as the background metric. Recall the definition of t in Definition 2.14. Using this notation, the required estimate can be stated as the following:

Proposition 3.2 *For all small enough $s \neq 0$ and $\mu > 0$, there exist constants $C_\mu, C_{\mu,\sigma}$ and a family of piecewise constant positive functions $\xi_{\mu,s}$ on V_s with $C_\mu^{-1}e^{-2\mu t} \leq \xi_{\mu,s} \leq C_\mu e^{-2\mu t}$ and $\int_{V_s} \xi_{\mu,s} \omega_{1,s}^n = 1$ such that*

$$\|e^{-\mu t}(u - \bar{u}_\mu)\|_{L^{2\sigma}(\omega_{1,s})} \leq C_{\mu,\sigma} \|\nabla^{\omega_{1,s}} u\|_{L^2(\omega_{1,s})}$$

for all $\sigma \in [1, \frac{n}{n-2}]$ and all $u \in C_0^\infty(V_s)$ where $\bar{u}_\mu = \int_{V_s} u \xi_{\mu,s} \omega_{1,s}^n$.

Proof Recall that $\hat{\omega}_s + i\partial\bar{\partial}\psi_{1,s}$ is Ricci flat in $\mathcal{U}_{x_i}$. It is clear that its diameter and volume on $U_{x_i,s}$ have two-sided bounds for small enough s . In particular, it has a Sobolev bound on $U_{x_i,s}$ uniform in s . On the region $(Z_s \setminus U_{\infty,s}) \cup \bigcup_{i=1}^{N_2} U_{x_i,s}$, the metrics are smooth, so its diameter and volume also have two-sided bounds and moreover the Sobolev bound is also uniform in s for small enough s . Now the proposition follows as the proof of Proposition 4.21 of [34]. $\square$

Using $\omega_{1,s}$ as the background metric, Step 1 of the proof of Theorem 4.1 of [34] can be applied without change. In particular, the potential $\psi_{\tau,s} - \psi_{1,s}$ has a C^0 -bound uniform in τ and s for $\tau \in [0, 1]$ and small non-zero s . However, recall that $\psi_{1,s}$ already has a uniform C^0 -bound. Therefore, the potential $\psi_{\tau,s}$ also has a uniform C^0 -bound. The next goal is the proof of the following C^2 -bound for the metric $\omega_{\tau,s} = \hat{\omega}_s + i\partial\bar{\partial}\psi_{\tau,s}$:

Proposition 3.3 *For all $\tau \in [0, 1]$, all small enough non-zero s , there exists a constant C independent of τ and s such that $C^{-1}\hat{\omega}_s \leq \omega_{\tau,s}$ on V_s . Moreover, for any closed subset $\mathcal{K}$ of $\mathcal{V} \setminus V^{\text{sing}}$, there exists a constant $C_{\mathcal{K}}$ only depending on $\mathcal{K}$ such that in addition $\omega_{\tau,s} \leq C_{\mathcal{K}}\hat{\omega}_s$ on $K_s = \mathcal{K} \cap V_s$.*

Proof Viewing the identity map as a harmonic map from $(V_s, \omega_{\tau,s})$ to $(V_s, \hat{\omega}_s)$, the Eells–Sampson’s Bochner type formula (Eq. (16) of [27], see also Eq. (3) of [18] and Theorem 4.1 of [49]) implies that

$$-\Delta_{\omega_{\tau,s}} \log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s \geq C(-|\operatorname{Ric}_{\omega_{\tau,s}}|_{\omega_{\tau,s}} - |\operatorname{Rm}_{\hat{\omega}_s}|_{\hat{\omega}_s} \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s).$$

Remark that $-\Delta|z|^2 = 4$ on $\mathbb{C}$ using the Laplacian operator $\Delta = d^*d + dd^*$ acting on 0-forms. The Riemannian curvature of $\hat{\omega}_s$ is uniformly bounded on $U_{\infty,s}$. The Ricci curvature of $\omega_{\tau,s}$ is bounded by $C\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$ because the Ricci form $i\partial\bar{\partial} \log F_{\tau,s}$ is bounded using $\hat{\omega}_s$ -norm. Therefore,

$$-\Delta_{\omega_{\tau,s}} \log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s \geq -C\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$$

on $U_{\infty,s}$.

On the compact part $V_s \setminus U_{\infty,s}$, the curvature of $\hat{\omega}_s$ is no longer bounded. However, recall that $\hat{\omega}_s$ equals to the pull back of the Fubini–Study metric in this region. Thus, as in the proof of Lemma 3.2 of [58], embed V_s into $\mathbb{CP}^{N_1}$ and view the composition of the embedding with the identity map as a harmonic map from $(V_s, \omega_{\tau,s})$ to $(\mathbb{CP}^{N_1}, \omega_{FS})$. In this case, the Eells–Sampson’s Bochner type formula implies that

$$-\Delta_{\omega_{\tau,s}} \log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s \geq C(-|\operatorname{Ric}_{\omega_{\tau,s}}|_{\omega_{\tau,s}} - |\operatorname{Rm}_{\omega_{FS}}|_{\omega_{FS}} \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s).$$

The Riemannian curvature of the Fubini–Study metric is indeed bounded. The Ricci curvature of $\omega_{\tau,s}$ is also bounded by $C\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$. Thus

$$-\Delta_{\omega_{\tau,s}} \log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s \geq -C\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$$

is true on V_s .

As in the proof of Lemma 3.2 of [58], using the formula

$$-\Delta_{\omega_{\tau,s}} \psi_{\tau,s} = 2n - 2\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s,$$

it is easy to see that

$$-\Delta_{\omega_{\tau,s}} (\log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s - C\psi_{\tau,s}) \geq C\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s - 2Cn.$$

Since $\omega_{\tau,s}$ and $\hat{\omega}_s$ are asymptotically cylindrical, for a large enough number $T_{\tau,s}$ depending on τ and s , for all $t \geq T_{\tau,s}$, $\log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s - C\psi_{\tau,s} \leq \log(2n)$. So either $\log \operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s - C\psi_{\tau,s}$ is bounded above by $\log(3n)$ or it attains its maximum at a point on V_s . In either cases, there is a uniform upper bound of $\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$ independent of τ and s using the C^0 -bound of $\psi_{\tau,s}$. The lower bound comes from the upper bound of $\operatorname{tr}_{\omega_{\tau,s}} \hat{\omega}_s$ and the bound of $F_{\tau,s}$ on K_S . $\square$

Let $\mathcal{K}$ be the set $\{t \geq 1\}$, then there is a C^2 -bound on $\mathcal{K}$. The uniform $C^{2,\alpha}$ -estimate independent of τ , s and $T \geq 1$ on $\{T \leq t \leq T+1\}$ for real Monge–Ampère equation was done by Evans–Krylov–Trudinger. See Sect. 17.4 of [32] for details. In complex

case, the arguments in Sect. 17.4 of [32] still work. An alternative way to achieve the $C^{2,\alpha}$ -estimate on $\{T \leq t \leq T+1\}$ was done by Theorem 1.5 of [14] using the rescaling argument. Now it is standard to get a C^∞ -bound of $\psi_{\tau,s}$ on $\{T \leq t \leq T+1\}$ independent of τ, s and $T \geq 1$ through Schauder estimates. Using Steps 3 and 4 of proof of Theorem 4.1 of [34], there is a C_v^∞ -bound of $\psi_{\tau,s}$ on $\{t \geq 1\}$. The same argument implies the C^∞ -estimate on $\mathcal{K}$ with bound depending on $\mathcal{K}$ but independent of τ and s for all compact subset $\mathcal{K}$ of $\mathcal{V} \setminus V^{\text{sing}}$.

Recall that there exists a smooth family of diffeomorphisms between S_0 and S_s . Its product with $[T, \infty) \times \mathbb{S}^1$ is a smooth family of diffeomorphisms from $\{t \geq T\} \cap V_0$ to $\{t \geq T\} \cap V_s$ such that it is the identity map when $s = 0$. For any closed subset K of $V_0 \setminus V^{\text{sing}}$, there exists a smooth family of embeddings $\hat{\Phi}_{K,s} : K \rightarrow V_s$ such that it is the identity map when $s = 0$ and is the given diffeomorphism when restricted to $t \geq T$. Using $\hat{\Phi}_{K,s}$, it is possible to talk about the C_{loc}^∞ convergence as the following:

Proposition 3.4 *For all $\tau_i \rightarrow \tau_\infty$ and $s_i \rightarrow 0$, there exists a subsequence τ_{i_k} and s_{i_k} and a metric $\tilde{\omega} = \hat{\omega}_0 + i\partial\bar{\partial}\tilde{\psi}$ on $V_0 \setminus V^{\text{sing}}$ such that for any closed subset K of $V_0 \setminus V^{\text{sing}}$, $\hat{\Phi}_{K,s}^* \psi_{\tau_{i_k}, s_{i_k}} \rightarrow \tilde{\psi}$ in C_v^∞ -sense on K . Moreover,*

$$\frac{\tilde{\omega}^n}{n!} = e^{F_{\tau_\infty, 0}} \frac{i^{n^2}}{2^n} \Omega_0 \wedge \bar{\Omega}_0$$

in weak sense.

Remark 3.5 For bounded plurisubharmonic functions on a smooth manifold, the Monge–Ampère equation makes sense in the weak sense. A weak solution on a singular manifold is defined as the solution whose pull back to the resolution of the singular manifold satisfies the Monge–Ampère equation in weak sense.

Proof For any $T_0 \geq T$, using the diagonal argument and the uniform C^∞ -bound on $\{T_0 \leq t \leq T_0+1\}$, it is easy to find a subsequence C^∞ -converging on $\{T_0 \leq t \leq T_0+1\}$. The limit belongs to C_v^∞ and satisfies the equation. The convergence in C_v^∞ -norm follows from the weighted analysis on the cylinder. The pre-compactness on $K \cap \{t \leq T\}$ follows from the C_{loc}^∞ -estimate of $\psi_{\tau,s}$ as in Theorem 1.4 of [58].

The limit $\tilde{\omega}$ satisfies the equation locally on $V \setminus V^{\text{sing}}$. In [36, 58], they claim that the compact analogue is a weak solution without proof. In an email from Hein and Sun to the author, they provided the following explanation:

$\tilde{\omega}$ can be pulled back to the resolution of V so that the equation is satisfied locally on the resolution except on the exceptional divisor. However, the pull back of $\tilde{\omega}$ can be written as $i\partial\bar{\partial}$ of a bounded plurisubharmonic function locally except on the exceptional divisor. By Sect. 5 of [23], the bounded plurisubharmonic function can be extended to the resolution. By Proposition 4.6.4 of [41], the extension satisfies the Monge–Ampère equation in the weak sense because the Monge–Ampère mass on the exceptional divisor vanishes. $\square$

On the other hand, it is also interesting to consider the Gromov–Hausdorff limit of a subsequence of $\omega_{\tau_{i_k}, s_{i_k}}$. To get started, using the C_{loc}^∞ -convergence of the metric outside the singularity, there exists $d > 0$ such that for any $x \in V^{\text{sing}}$ and any point

$q \in U_{x,s_{i_k}}$ with $\text{dist}_{\omega_{\tau_{i_k},s_{i_k}}}(q, \partial U_{x,s_{i_k}}) = d$, the volume of the ball $B_{\omega_{\tau_i,s_i}}(q, d)$ has a positive lower bound. Therefore, for any $x \in V^{\text{sing}}$ and any point $p \in U_{x,s_{i_k}}$, let D be the distance of p to $\partial U_{x,s_{i_k}}$ using $\omega_{\tau_{i_k},s_{i_k}}$. Then if $D > 3d$, choose q as the point on the minimal geodesic of length D joining p and $\partial U_{x,s_{i_k}}$ such that the distance from p to q is $D - d$. Therefore, using the fact that $\omega_{\tau_{i_k},s_{i_k}}$ is Ricci flat on $U_{x,s_{i_k}}$, the volume comparison implies that using $\omega_{\tau_{i_k},s_{i_k}}$,

$$\text{Vol}(B(p, D - 2d)) \geq \frac{\text{Vol}(B(p, D) \setminus B(p, D - 2d))}{\left(\frac{D}{D-2d}\right)^{2n} - 1} \geq \frac{CD \text{Vol}(B(q, d))}{d}.$$

Since the volume of $U_{x,s_{i_k}}$ is bounded, D is also bounded. Using the standard ϵ -net argument, it is easy to see that $(V_{s_{i_k}}, \omega_{\tau_{i_k},s_{i_k}})$ has a Gromov–Hausdorff limit (X, d_X) with Gromov–Hausdorff approximation equalling to $\hat{\Phi}_{K,s_{i_k}}$ when restricted to the set $K = V_0 \setminus \bigcup_{x \in V^{\text{sing}}} U_{x,0}$. Remark that the space (X, d_X) is non-compact and therefore the result of Donaldson–Sun cannot be applied directly. In order to solve this problem, the Gromov–Hausdorff limit of the metrics defined by the $(1,1)$ -forms $(\pi_1 \circ i)^* \omega_{FS} + i\partial\bar{\partial}\psi_{\tau_{i_k},s_{i_k}}$ is considered instead. This $(1,1)$ -form is positive on $\{t > T_0\}$ for large enough T_0 , but may not be positive on $\{t \leq T_0\}$. However, by checking the difference between $(\pi_1 \circ i)^* \omega_{FS}$ and $\hat{\omega}_{s_{i_k}}$ on the set $\{t \leq T_0\}$ carefully, there exists a bump $(1,1)$ -form β on $\mathbb{CP}^1$ such that for large enough number N_3 , $(\pi_1 \circ i)^* \omega_{FS} + N_3 f^* \beta \geq \hat{\omega}_{s_{i_k}}$ on $\{t \leq T_0\}$. Thus $(\pi_1 \circ i)^* \omega_{FS} + i\partial\bar{\partial}\psi_{\tau_{i_k},s_{i_k}} + N_3 f^* \beta$ is positive. Without loss of generality assume that $N_3 \beta = N_4 [c_1(\mathcal{O}(1))]$ for an integer N_4 . Thus, after modifying the metric $\omega_{\tau_{i_k},s_{i_k}}$ near infinity, there exists a compact metric on Z_s in the cohomology class $c_1((\pi_1 \circ i)^*(\mathcal{O}(1)) \otimes (f^* \mathcal{O}(1))^{\otimes N_4})$ such that the diameter, the volume and the Ricci curvature have two-sided bounds. The Gromov–Hausdorff limit is isometric to (X, d_X) except on the end $U_{\infty,s_{i_k}}$. Therefore, even though (X, d_X) is non-compact, Cheeger–Colding theory [10–12] and Donaldson–Sun theory [25,26] can still be used.

As in [58], (X, d_X) is isometric to the completion of the metric $\tilde{\omega}$. In particular, when fixing $\tau_i = \tau_\infty \in \mathcal{T}$, using the uniqueness of weak solutions, d_X is isometric to $\omega_{\tau_i,0}$. Now assume that instead $\tau_i \rightarrow \tau_\infty \in \mathcal{T}$, then by choosing small enough s_i , (X, d_X) is also the Gromov–Hausdorff limit of $\omega_{\tau_i,0}$. The rest parts of [36] can be applied without change.

4 The Matching Problem

This section solves the matching problem. The starting point is the review of the matching problem of smooth manifolds in [21] using a particular example:

Example 4.1 (Example 7.3 of [20]) Fix a 2-plane $\Pi \subset \mathbb{CP}^4$. Let $X_+ \subset \mathbb{CP}^4$ be a general quartic threefold containing Π . It has nine nodal singularities. The blowing up of X_+ over Π yields a non-singular threefold $Y_+ \rightarrow X_+$ with nine $(-1, -1)$ -curves resolving the nine ordinary double points of X_+ on Π . In Example 7.3 of [20], Corti–Haskins–Nordström–Pacini prove that $N_+ := H^2(Y_+, \mathbb{Z}) = \mathbb{Z}^2$ with basis $\tilde{\Pi}$

(the proper transform of Π) and $-K_{Y_+}$. The quadratic form $[.] \cup [.] \cup -K_{Y_+}$ in this basis equals to

$$\begin{bmatrix} -2 & 1 \\ 1 & 4 \end{bmatrix}.$$

It is easy to see that $-K_{Y_+}$ also equals to the pull back of the $\mathcal{O}(1)$ -bundle of $\mathbb{CP}^4$.

On the other hand, choose $Y_- = X_- = \mathbb{CP}^3$. Then $N_- := H^2(Y_-, \mathbb{Z}) = \mathbb{Z}$ with base $\mathcal{O}(1)$. The anti-canonical divisor $-K_{Y_-} = \mathcal{O}(4)$. The quadratic form $c_1(\mathcal{O}(1)) \cup c_1(\mathcal{O}(1)) \cup -K_{Y_-}$ in this basis equals to 4 times the identity.

Choose smooth anti-canonical divisors $S_{\pm}$ of $Y_{\pm}$. It is easy to find other anti-canonical divisors $S_{0,\pm}$ intersecting $S_{\pm}$ transversally. Let $C_{\pm}$ be their intersections. Then the ratios of the corresponding sections provide holomorphic functions $f_{\pm}$ from $\tilde{Z}_{\pm}$ to $\mathbb{CP}^1$ with $\{f_{\pm} = \infty\}$ equal to the proper transforms $S_{\infty,\pm}$ of $S_{\pm}$, where $\tilde{Z}_{\pm}$ is the blowing up of $Y_{\pm}$ at $C_{\pm}$. Define Z_+ as the blowing up of X_+ at C_+ . Define Z_- as $\tilde{Z}_-$. Let $\Omega_{\pm}$ be the meromorphic 3-forms on $\tilde{Z}_{\pm}$ with simple poles along $S_{\infty,\pm}$. Their residues $\omega_{S_{\pm}}^J + i\omega_{S_{\pm}}^K$ are nowhere vanishing $(2,0)$ -forms on $S_{\infty,\pm} = S_{\pm}$. The main goal in the smooth case is to find Kähler classes on $\tilde{Z}_{\pm}$ and a diffeomorphism r from S_+ to S_- such that the unique Ricci-flat metrics $\omega_{S_{\pm}}$ on the restrictions of the Kähler classes on $S_{\pm}$ satisfy

$$(\omega_{S_+}, \omega_{S_+}^J, \omega_{S_+}^K) = (r^* \omega_{S_-}^J, r^* \omega_{S_-}, -r^* \omega_{S_-}^K).$$

Remark that $Y_{\pm}$ are simply-connected. By Lefschetz hyperplane theorem, $S_{\pm}$ are also simply-connected. Since they have trivial canonical line bundles using the adjoint formula, they are all K3 surfaces. Therefore, the following theorem about the moduli space of marked hyperkähler structures on the K3 surface proved by [5,48,56,60,63] can be applied:

Theorem 4.2 ([4]) *Let S be the smooth 4-manifold which underlies the minimal resolution of $\mathbb{T}^4/\mathbb{Z}_2$. Let W be the space of three cohomology classes $[\alpha_1], [\alpha_2], [\alpha_3]$ in $H^2(S, \mathbb{R})$ which satisfy the following conditions:*

(1) *(Integrability)*

$$\int_M \alpha_i \wedge \alpha_j = 0$$

for $i \neq j$ and

$$\int_M \alpha_1^2 = \int_M \alpha_2^2 = \int_M \alpha_3^2 > 0.$$

(2) *(Nondegeneracy) For any $[\Sigma] \in H_2(S, \mathbb{Z})$ with $[\Sigma]^2 = -2$, there exists $i \in \{1, 2, 3\}$ with $[\alpha_i][\Sigma] \neq 0$.*

W has two components W^+ and W^- . For any $([\alpha_1], [\alpha_2], [\alpha_3]) \in W^+$, there exists on S a hyperkähler structure for which the cohomology classes of the Kähler forms $[\omega_i]$ are the given $[\alpha_i]$. It is unique up to tri-holomorphic isometries which induce identity on $H_2(S, \mathbb{Z})$. Moreover, any hyperkähler structure on K3 surface must be constructed by this way.

Thus, it suffices to find out the matching cohomology classes. In general, the matching data can only be found in the deformation classes of $\tilde{Z}_\pm$.

Remark that all K3 surfaces are diffeomorphic to S . Denote $H^2(S, \mathbb{Z})$ by L . L is a lattice

$$L = -E_8 \oplus -E_8 \oplus \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}^{\oplus 3}$$

using the intersection form as the quadratic form. For example, see [59] for a concrete description of the K3 lattice. It is clear that the quadratic form in fact acts on $L \otimes \mathbb{R}$. The set of elements in $L \otimes \mathbb{R}$ for which the square using this quadratic form is positive is called the positive cone. It is easy to get the following proposition:

Proposition 4.3 *The lattice $N_+ \oplus N_-$ with quadratic form $\begin{bmatrix} -2 & 1 & 0 \\ 1 & 4 & 0 \\ 0 & 0 & 4 \end{bmatrix}$ can be embedded into L .*

Proof Let B_1, B_2, B_3 and C_1, C_2, C_3 be the basis in the last three components. Embed $\tilde{\Pi}$ into a simple root of the first $-E_8$ component. The adjacent simple root of $-E_8$ is an element whose square equals to -2 and its product with $\tilde{\Pi}$ is 1. Let $-K_{Y_+}$ be the sum of this element with $B_1 + C_1 + B_2 + C_2 + B_3 + C_3$, then it is clear that $(-K_{Y_+})^2 = 4$. Now let $-\frac{1}{4}K_{Y_-} = B_1 + C_1 - B_2 - C_2$. $\square$

The following key proposition was used by [20]:

Proposition 4.4 (Proposition 6.9 of [21]) *Fix the embeddings $N_\pm \subset L$ as in Proposition 4.3. Let $D_{N_\pm}$ be the Griffiths domains $\{\Pi \in \mathbb{P}(N_\pm^\perp \otimes \mathbb{C}) : \Pi \wedge \bar{\Pi} > 0\}$. Let $\mathcal{Y}_\pm$ be the deformation types of $Y_\pm$ such that there exist anti-canonical K3 divisors $S_\pm$ on $Y_\pm$ with $N_\pm$ -polarised markings $h_\pm : L \cong H^2(S_\pm, \mathbb{Z})$, which means, by definition, the restriction maps $H^2(Y_\pm, \mathbb{Z}) \rightarrow H^2(S_\pm, \mathbb{Z})$ are equivalent to the inclusions $N_\pm \hookrightarrow L$ for the chosen isomorphisms $N_\pm \cong H^2(Y_\pm, \mathbb{Z})$ and $h_\pm$. Then there exist*

- (1) *sets $U_{\mathcal{Y}_\pm} \subset D_{N_\pm}$ with complement locally finite unions of complex analytic submanifolds of positive codimensions;*
- (2) *open subcones $\text{Amp}_{\mathcal{Y}_\pm}$ of the positive cones of $N_\pm \otimes \mathbb{R}$ with the following property: for any $\Pi_\pm \in U_{\mathcal{Y}_\pm}$ and $k_\pm \in \text{Amp}_{\mathcal{Y}_\pm}$, there exist $Y_\pm \in \mathcal{Y}_\pm$, smooth anti-canonical divisors $S_\pm$, and $N_\pm$ -polarized markings $h_\pm : L \rightarrow H^2(S_\pm, \mathbb{Z})$ such that $h_\pm(\Pi_\pm) = H^{2,0}(S_\pm)$ and $h_\pm(k_\pm)$ are the restrictions to $S_\pm$ of Kähler classes on $Y_\pm$.*

For any $\Pi_\pm \in U_{\mathcal{Y}_\pm}$ and $k_\pm \in \text{Amp}_{\mathcal{Y}_\pm}$, choose $(2,0)$ -forms $\omega_{S_\pm}^J + i\omega_{S_\pm}^K$ in $H^{2,0}(S_\pm)$ and denote $h_\pm(k_\pm)$ by $\omega_{S_\pm}$. Then there exist Kähler classes on $\tilde{Z}_\pm$ such that their restrictions to $S_\pm$ are also $\omega_{S_\pm}$.

The following proposition was proved in [21] using Proposition 4.4:

Proposition 4.5 (Proposition 6.18 of [21]) *There exist $\Pi_{\pm}$ and $k_{\pm}$ such that the corresponding $Y_{\pm}$, $S_{\pm}$, $h_{\pm}^{-1} : H^2(S_{\pm}, \mathbb{Z}) \rightarrow L$, $\omega_{S_{\pm}}$, $\omega_{S_{\pm}}^J$, $\omega_{S_{\pm}}^K$ satisfy*

$$h_+^{-1}([\omega_{S_+}], [\omega_{S_+}^J], [\omega_{S_+}^K]) = h_-^{-1}([\omega_{S_-}^J], [\omega_{S_-}], -[\omega_{S_-}^K]).$$

Proof The proof due to [21] is sketched here in order to see how to adjust it to the singular situation.

Let $T(\mathbb{R})$ be the subspace $N_+^{\perp} \cap N_-^{\perp}$ of $L \otimes \mathbb{R}$. Consider the real manifold

$$\mathcal{A} = \mathbb{S}(\text{Amp}_{Y_+}) \times \mathbb{S}(\text{Amp}_{Y_-}) \times \mathbb{S}(T(\mathbb{R})),$$

where $\mathbb{S}(\text{Amp}_{Y_+})$, $\mathbb{S}(\text{Amp}_{Y_-})$ or $\mathbb{S}(T(\mathbb{R}))$ means the set of elements in Amp_{Y_+} , Amp_{Y_-} or $T(\mathbb{R})$ whose square equals to 1 with respect to the quadratic form. There are two projections $\text{pr}_{\pm} : \mathcal{A} \rightarrow D_{N_{\pm}}$, $(k_+, k_-, k) \rightarrow \langle k_{\mp} \pm ik \rangle$, where $D_{N_{\pm}}$ are the Griffiths period domains. The key point of the proof due to [21] is the fact that the real analytic embedded submanifolds $\mathbb{S}(\text{Amp}_{Y_{\mp}}) \times \mathbb{S}(T(\mathbb{R}))$ of $D_{N_{\pm}}$ are totally real with maximal dimensions. Therefore, it is easy to see that the set $\text{pr}_+^{-1}(U_{Y_+}) \cap \text{pr}_-^{-1}(U_{Y_-})$ is non-empty. $\square$

Up to here, it has been shown how to find the matching data for smooth asymptotically cylindrical Calabi–Yau manifolds. Remark that $\mathbb{S}(\text{Amp}_{Y_-})$ has a single point. Denote it by k_- . In order to get the matching data for the manifolds with nodal singularities, the cohomology class k_+ defined as $h_+^{-1}[\omega_{S_+}]$ must come from the restriction of $-\frac{1}{2}c_1(K_{Y_+})$. This means that even though for any k in the complement of real submanifolds with smaller dimensions in $\mathbb{S}(T(\mathbb{R}))$, it is still true that $\langle k_- + ik \rangle \in U_{Y_+}$, in general, $\langle k_+ - ik \rangle$ may not be in U_{Y_-} due to the restriction on the value of k_+ . This problem is similar to the “handcrafted gluing” problem in [21]. The following well-known lemma can be used solve this problem:

Lemma 4.6 (Chapter 3 of [57], cited as Lemma 7.15 of [21]) *Let S be a K3 surface, and let A be a nef line bundle on S with $A^2 > 0$ (i.e., A is nef and big). Then either*

- (1) $|A|$ is monogonal, that is, $A = aE + \Gamma$, where E and Γ are holomorphic curves with $E^2 = 0$, $E \cdot \Gamma = 1$, $\Gamma^2 = -2$, and $a = 1, 2, 3, \dots$, or
- (2) $|A|$ has no fixed point, is base point free and either:
 - (2.1) the morphism given by $|A|$ is birational onto its image and an isomorphism away from a finite union of -2 curves, or
 - (2.2) A is hyperelliptic, that is, one of the following cases holds: (2.2.1) $A^2 = 2$ and S is a double cover of $\mathbb{CP}^2$; (2.2.2) $A = 2B$ with $B^2 = 2$ and S is a double cover of the Veronese surface; or (2.2.3) S has an elliptic pencil E with $A \cdot E = 2$.

Motivated by Lemma 4.6, the first thing to check is linear combinations $a\tilde{\Pi} - bK_{Y_+} - \frac{c}{4}K_{Y_-} \in L$,

$$\left(a\tilde{\Pi} - bK_{Y_+} - \frac{c}{4}K_{Y_-}\right)^2 = -2$$

and

$$\left(a\tilde{\Pi} - bK_{Y_+} - \frac{c}{4}K_{Y_-}\right) \cdot (-K_{Y_+}) = 0.$$

The first condition implies that $a, b, c \in \mathbb{Z}$. The second condition implies that $-2a^2 + 2ab + 4b^2 + 4c^2 = -2$. The third condition implies that $a + 4b = 0$. Therefore $-36b^2 + 4c^2 = -2$. This is impossible. Thus, for all $C \in L$ satisfying $C^2 = -2$ and $C \cdot (-K_{Y_+}) = 0$, C cannot be in $(N_+ \otimes \mathbb{R}) \oplus (N_- \otimes \mathbb{R})$. So $C^\perp$ intersects $T(\mathbb{R})$ transversally. There is no difficulty to find $k \in \mathbb{S}(T(\mathbb{R}))$ such that k does not lie in $C^\perp$ for all such C and $\langle k_- + ik \rangle \in U_{\mathcal{Y}_+}$ is still true. Similarly, it is possible to assume that for all $E \in L$ satisfying $E^2 = 0$, $E \cdot (-K_{Y_+}) = 0$ and $E \cdot (-\frac{1}{4}K_{Y_-}) = 2$, k does not lie in $C^\perp$.

By Theorem 4.2, there exists a K3 surface S_- with a marking h_- such that

$$h_- \left(\left[\omega_{S_-}^J \right], \left[\omega_{S_-} \right], - \left[\omega_{S_-}^K \right] \right) = (k_+, k_-, k).$$

It is smooth and does not contain any -2 curve because k_+ and k cannot vanish simultaneously on it. $A = -\frac{1}{4}c_1(K_{Y_-}) = 2k_-$ lies in $H^2(S_-, \mathbb{Z})$ and is a Kähler $(1,1)$ -class on S_- , so A is a nef line bundle on S_- with $A^2 = 4 > 0$. By Lemma 4.6, the morphism given by $|A|$ is an isomorphism onto its image. By Kodaira vanishing theorem and Riemann–Roch theorem, the morphism given by $|A|$ is in fact an isomorphism onto a smooth quartic surface in $\mathbb{CP}^3$. This solves the matching problem.

In an email from Nordström to the author, he said that Lemma 2.4 of [31] and Lemma 5.18 of [22] may provide more examples of matching data for the singular twisted connected sum problem.

5 Harmonic Forms on the Nodal Cone

This section deals with the homogeneous harmonic forms on strongly regular Calabi–Yau cones $C = C(F)$ with complex dimension 3 defined in Definition 2.10. Assume that F is not the sphere $\mathbb{S}^5$ with the standard metric. Some results in this section have been proved in [7, 9, 30, 40].

The starting point is the definition of homogeneous forms on C .

Definition 5.1 A p -form $\gamma = r^\lambda(r^{p-1}dr \wedge \alpha + r^p\beta)$ is called homogeneous of rate λ if $\frac{\partial}{\partial r}\alpha = \frac{\partial}{\partial r}\beta = 0$.

A direct calculation shows the following:

Proposition 5.2 Let $\gamma = r^\lambda(r^{p-1}dr \wedge \alpha + r^p\beta)$ be a p -form on C . Let d_C , d_C^* and Δ_C be the exterior derivative, its adjoint and the Hodge Laplacian operator on C using the cone metric. Let d_F , d_F^* , Δ_F be the exterior derivative restricted to $\{r = r_0\} \simeq F$ for each $r_0 > 0$, its adjoint using the metric on $F = \{r = 1\}$ and the Hodge Laplacian

operator $\Delta_F = d_F d_F^* + d_F^* d_F$, then

$$\begin{aligned} d_C \gamma &= r^{\lambda+p-1} dr \wedge \left((\lambda + p) \beta + r \frac{\partial}{\partial r} \beta - d_F \alpha \right) + r^{\lambda+p} d_F \beta, \\ d_C^* \gamma &= r^{\lambda+p-3} dr \wedge (-d_F^* \alpha) + r^{\lambda+p-2} \left(-(\lambda - p + 6) \alpha - r \frac{\partial}{\partial r} \alpha + d_F^* \beta \right), \\ \Delta_C \gamma &= r^{\lambda+p-3} dr \wedge \left(\Delta_F \alpha - (\lambda + p - 2)(\lambda - p + 6) \alpha - \left(r \frac{\partial}{\partial r} \right)^2 \alpha \right. \\ &\quad \left. - (2\lambda + 4) r \frac{\partial}{\partial r} \alpha - 2 d_F^* \beta \right) + r^{\lambda+p-2} (\Delta_F \beta - (\lambda + p)(\lambda - p + 4) \beta \\ &\quad - \left(r \frac{\partial}{\partial r} \right)^2 \beta - (2\lambda + 4) r \frac{\partial}{\partial r} \beta - 2 d_F \alpha). \end{aligned}$$

So the homogeneous harmonic forms are closely related to the eigenforms on the link F .

Definition 5.3 Using Hilbert–Schmidt theorem, define $\phi_{0,j}$ as the orthogonal basis of $L^2(\Lambda^0(F))$ with $\Delta_F \phi_{0,j} = \mu_{0,j} \phi_{0,j}$. Then $d_F \phi_{0,j}$ are orthogonal to each other because

$$(d_F \phi_{0,j}, d_F \phi_{0,j'})_{L^2} = (d_F^* d_F \phi_{0,j}, \phi_{0,j'})_{L^2} = \mu_{0,j} (\phi_{0,j}, \phi_{0,j'})_{L^2} = 0$$

if $j \neq j'$. By Hilbert–Schmidt theorem applied to the Laplacian operator acting on 1-forms, it is possible to define $\phi_{1,j}$ such that $d_F \phi_{0,j}$ for $j = 2, 3, \dots$ and $\phi_{1,j}$ for $j = 1, 2, \dots$ are orthogonal basis of $L^2(\Lambda^1(F))$ with

$$\Delta_F d_F \phi_{0,j} = \mu_{0,j} d_F \phi_{0,j}$$

and

$$\Delta_F \phi_{1,j} = \mu_{1,j} \phi_{1,j}.$$

It is clear that $d_F^* d_F \phi_{0,j} = \mu_{0,j} \phi_{0,j}$ while $d_F^* \phi_{1,j} = 0$.

Inductively, for $p = 0, 1, 2, \dots, 5$, $d_F \phi_{p,j}$ are orthogonal to each other, so it is possible to define $\phi_{p+1,j}$ such that $d_F \phi_{p,j}$ for $j = h_p + 1, h_p + 2, \dots$ and $\phi_{p+1,j}$ for $j = 1, 2, 3, \dots$ are orthogonal basis of $L^{p+1}(\Lambda^2(F))$ with

$$\Delta_F d_F \phi_{p,j} = \mu_{p,j} d_F \phi_{p,j}$$

and

$$\Delta_F \phi_{p+1,j} = \mu_{p+1,j} \phi_{p+1,j},$$

where h_p is the dimension of the cohomology group $H^p(F, \mathbb{R})$.

The relationship between homogeneous harmonic forms on C and eigenforms on F are given by the following:

Proposition 5.4 Choose $\lambda = -2$ in Proposition 5.2, then

$$\Delta_C \gamma = r^{p-5} dr \wedge \left(\Delta_F \alpha + (p-4)^2 \alpha - \left(r \frac{\partial}{\partial r} \right)^2 \alpha - 2d_F^* \beta \right) + r^{p-4} \left(\Delta_F \beta + (p-2)^2 \beta - \left(r \frac{\partial}{\partial r} \right)^2 \beta - 2d_F \alpha \right).$$

So it is important to study the eigenvalues of the self-adjoint operator

$$(\alpha, \beta) \rightarrow (\Delta_F \alpha + (p-4)^2 \alpha - 2d_F^* \beta, \Delta_F \beta + (p-2)^2 \beta - 2d_F \alpha)$$

from a subspace of $L^2(\Lambda^{p-1} F \oplus \Lambda^p F)$ to another subspace of $L^2(\Lambda^{p-1} F \oplus \Lambda^p F)$. It is easy to see that the eigenforms are, up to linear combinations,

- (1) $(d_F \phi_{p-2,j}, 0)$ with eigenvalue $\mu_{p-2,j} + (p-4)^2$,
- (2) $(0, \phi_{p,j})$ with eigenvalue $\mu_{p,j} + (p-2)^2$,
- (3) $(\phi_{p-1,j}, 0)$ with eigenvalue $(p-4)^2$ if $\mu_{p-1,j} = 0$,
- (4) $(\phi_{p-1,j}, \frac{3-p \pm \sqrt{(p-3)^2 + \mu_{p-1,j}}}{\mu_{p-1,j}} d_F \phi_{p-1,j})$ with eigenvalue

$$\left(\sqrt{(p-3)^2 + \mu_{p-1,j} \mp 1} \right)^2$$

if $\mu_{p-1,j} \neq 0$.

Using the identification $\gamma = r^{-2}(r^{p-1} dr \wedge \alpha + r^p \beta)$ between $\gamma \in \Lambda^p(C(F))$ with $(\alpha, \beta) \in \Lambda^{p-1} F \oplus \Lambda^p F$, the eigenforms are denoted by $\hat{\phi}_{p,j} \in \Lambda^p(C(F))$ with eigenvalues $\hat{\mu}_{p,j}$. By Hilbert–Schmidt theorem, they form an L^2 basis. By Proposition 5.5, any homogeneous form is harmonic if and only if it is the linear combination of $r^{\pm \sqrt{\hat{\mu}_{p,j}}} \hat{\phi}_{p,j}$.

A more precise decomposition is the following:

Proposition 5.5 A harmonic homogeneous p -form $\gamma = r^\lambda(r^{p-1} dr \wedge \alpha + r^p \beta)$ can be written as a linear combination of

- (1) closed but not coclosed harmonic homogeneous forms $r^\lambda r^{p-1} dr \wedge d_F \phi_{p-2,j}$ with $\mu_{p-2,j} = (\lambda + p - 2)(\lambda - p + 6) \neq 0$,
- (2) closed but not coclosed harmonic homogeneous forms $r^\lambda r^{p-1} dr \wedge \phi_{p-1,j}$ with $\mu_{p-1,j} = \lambda + p - 2 = 0$, and $\lambda \neq -2$.
- (3) closed and coclosed homogeneous forms $r^\lambda r^{p-1} dr \wedge \phi_{p-1,j}$ with

$$\mu_{p-1,j} = \lambda - p + 6 = 0,$$

(4) *closed and coclosed homogeneous forms*

$$r^\lambda((\lambda + p)r^{p-1}dr \wedge \phi_{p-1,j} + r^p d_F \phi_{p-1,j})$$

with $\mu_{p-1,j} = (\lambda + p)(\lambda - p + 6) \neq 0$,

(5) *neither closed nor coclosed homogeneous harmonic forms*

$$r^\lambda(-(\lambda - p + 4)r^{p-1}dr \wedge \phi_{p-1,j} + r^p d_F \phi_{p-1,j})$$

with $\mu_{p-1,j} = (\lambda + p - 2)(\lambda - p + 4) \neq 0$ and $\lambda \neq -2$,

(6) *closed and coclosed homogeneous forms $r^\lambda(r^p \phi_{p,j})$ with $\mu_{p,j} = \lambda + p = 0$, and*

(7) *coclosed but not closed harmonic homogeneous forms $r^\lambda(r^p \phi_{p,j})$ with*

$$\mu_{p,j} = (\lambda + p)(\lambda - p + 4)$$

and $\lambda + p \neq 0$.

Proof Use Hilbert–Schmidt theorem to write α and β as the generalized Fourier series

$$\alpha = \sum_{j=h_{p-2}+1}^{\infty} \alpha_{p-2,j} d_F \phi_{p-2,j} + \sum_{j=1}^{\infty} \alpha_{p-1,j} \phi_{p-1,j},$$

and

$$\beta = \sum_{j=h_{p-1}+1}^{\infty} \beta_{p-1,j} d_F \phi_{p-1,j} + \sum_{j=1}^{\infty} \beta_{p,j} \phi_{p,j}.$$

Since $d_F \phi_{p-2,j}$ is perpendicular to $d_F^* \beta$, the harmonic assumption implies that $\alpha_{p-2,j} = 0$ unless $\mu_{p-2,j} = (\lambda + p - 2)(\lambda - p + 6)$. When $\mu_{p-1,j} = 0$, using the fact that $\phi_{p-1,j}$ is also perpendicular to $d_F^* \beta$, the harmonic assumption implies that $\alpha_{p-1,j} = 0$ unless $(\lambda + p - 2)(\lambda - p + 6) = 0$. Similarly, $\beta_{p,j} = 0$ unless $\mu_{p,j} = (\lambda + p)(\lambda - p + 4)$.

For $j = h_{p-1} + 1, h_{p-1} + 2, \dots$, the equation

$$\Delta_F \beta - (\lambda + p)(\lambda - p + 4)\beta - 2d_F \alpha = 0$$

implies that $\alpha_{p-1,j} = \frac{1}{2}(\mu_{p-1,j} - (\lambda + p)(\lambda - p + 4))\beta_{p-1,j}$. So the equation

$$\Delta_F \alpha - (\lambda + p - 2)(\lambda - p + 6)\alpha - 2d_F^* \beta = 0$$

implies that

$$\frac{1}{4}(\mu_{p-1,j} - (\lambda + p - 2)(\lambda - p + 6))(\mu_{p-1,j} - (\lambda + p)(\lambda - p + 4)) = \mu_{p-1,j}$$

unless $\alpha_{p-1,j} = \beta_{p-1,j} = 0$. Note that

$$\frac{1}{4}(\mu_{p-1,j} - (\lambda + p - 2)(\lambda - p + 6))(\mu_{p-1,j} - (\lambda + p)(\lambda - p + 4)) = \mu_{p-1,j}$$

if and only if $\mu_{p-1,j} = (\lambda + p)(\lambda - p + 6)$ or $\mu_{p-1,j} = (\lambda + p - 2)(\lambda - p + 4)$. $\square$

Recall the definition of $\mathcal{K}_i(\lambda_i)$ in Definition 2.15. It is easy to prove the following proposition:

Proposition 5.6 *Suppose that*

$$\gamma \in \mathcal{K}_i(\lambda_i)$$

for the Hodge Laplacian acting on p -forms. Then up to linear combinations, either

- (1) $\gamma = r^{\pm\sqrt{\hat{\mu}_{p,j}}}\hat{\phi}_{p,j}$ *is homogeneous with $\hat{\mu}_{p,j} \in \mathbb{R}$, or*
- (2) $\gamma = \log r \hat{\phi}_{p,j}$ *with $\hat{\mu}_{p,j} = 0$.*

Proof Write γ as

$$\gamma = \sum_{j=1}^{\infty} \gamma_j(r) \hat{\phi}_{p,j},$$

then

$$\left(\hat{\mu}_{p,j} - \left(r \frac{d}{dr} \right)^2 \right) \gamma_j = 0.$$

If $\hat{\mu}_{p,j} \neq 0$, γ_j is the linear combination of $r^{\pm\sqrt{\hat{\mu}_{p,j}}}$. When $\hat{\mu}_{p,j} = 0$, γ_j is the linear combination of 1 and $\log r$. $\square$

The next goal is the estimate of eigenvalues:

Proposition 5.7 (Obata [55]) $\mu_{0,1} = 0$ and $\phi_{0,1} = 1$. For all $j = 2, 3, 4, \dots$, $\mu_{0,j} > 5$.

Proof This follows from [55] because the metric on F is Einstein with scalar curvature 20 and F is not isometric to the sphere. $\square$

Proposition 5.8 $\mu_{1,j} \geq 8$ for all $j = 1, 2, 3, \dots$, moreover, when $\mu_{1,j} = 8$, then $r^2 \phi_{1,j}^{\#}$ is a Killing vector field on C , where $\phi_{1,j}^{\#}$ means the metric dual using g_C .

Proof This is similar to Lemma 3.11 of [40]. $\square$

Proposition 5.9 If $\phi_{2,j}$ is a primitive $(1,1)$ -form on C , then either $\mu_{2,j} = 0$ or $\mu_{2,j} \geq 9$.

Proof It is proved in the proof of Proposition 4.9. (iii) of [30]. $\square$

The homogeneous harmonic forms on C can be studied using the estimates of $\mu_{p,j}$.

Proposition 5.10 *Let γ be a homogeneous harmonic 1-form on C with rate in $[-3, 0]$, then $\gamma = 0$.*

Proof This is an immediate corollary of Propositions 5.5, 5.7, and 5.8. $\square$

Recall the following theorem essentially due to Cheeger–Tian:

Proposition 5.11 (Theorem 7.27 of [13], see also Lemma 2.17 of [36]). *Let γ be a homogeneous 1-form on C with rate in $(0, 1]$. Then γ is harmonic if and only if, up to linear combinations, either $\gamma = d(r\sqrt{\mu_{0,j}+4}-2)\phi_{0,j}$ with $\mu_{0,j} \in (5, 12]$ or $\gamma = r^2\phi_{1,j}$, where $\mu_{1,j} = 8$ or $\gamma = r dr$.*

Remark that Proposition 5.11 has been adjusted to the strongly regular Calabi–Yau cone case.

Lemma 5.12 *Choose $r\phi_{1,1} = e^1 = -Je^0 = -Jdr$. Then $\mu_{1,1} = 8$. If $\mu_{1,j} = 8$, then there exists a constant k_j such that the Lie derivatives*

$$L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}\omega = L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}Re\Omega = L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}Im\Omega = 0,$$

where $^\#$ means the metric duals using g_C .

Proof By Proposition 5.8, $r^2\phi_{1,j}^\#$ preserves the metric g_C . Let $e^{sr^2\phi_{1,j}^\#}$ be the one-parameter subgroup generated by $r^2\phi_{1,j}^\#$, then $e^{sr^2\phi_{1,j}^\#}$ preserves the metric g_C . Since ω, Ω are parallel unit-length forms on C , $(e^{sr^2\phi_{1,j}^\#})^*\omega$ and $(e^{sr^2\phi_{1,j}^\#})^*\Omega$ are also parallel. The holonomy group of g_C equals to $SU(3)$ instead of a proper subgroup of $SU(3)$, so the only unit-length parallel 2-forms are $\pm\omega$. By continuity, $(e^{sr^2\phi_{1,j}^\#})^*(\omega) = \omega$ for all s . So $L_{r^2\phi_{1,j}^\#}\omega = 0$. Any unit-length parallel 3-form must be $e^{i\theta_j(s)}\Omega$. So after differentiating, there exists a constant k_j such that $L_{r^2\phi_{1,j}^\#}Re\Omega = 3k_jIm\Omega$ and $L_{r^2\phi_{1,j}^\#}Im\Omega = -3k_jRe\Omega$. When $j = 1$, $r^2\phi_{1,1}^\# = rJ\frac{\partial}{\partial r}$. So

$$L_{r^2\phi_{1,1}^\#}Re\Omega = d(r^2\phi_{1,1}^\# \lrcorner Re\Omega) = d\left(r\frac{\partial}{\partial r} \lrcorner Im\Omega\right) = 3Im\Omega.$$

Similarly $L_{r^2\phi_{1,1}^\#}Im\Omega = -3Re\Omega$. So

$$L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}\omega = L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}Re\Omega = L_{r^2(\phi_{1,j}-k_j\phi_{1,1})^\#}Im\Omega = 0.$$

$\square$

Proposition 5.13 *A homogeneous 1-form γ with rate $\lambda \in [-3, 1]$ is harmonic if and only if up to linear combinations, either*

- (1) $\gamma = d_C(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j})$ with $\mu_{0,j} \in (5, 12]$,
- (2) $\gamma = d_C r^2$,
- (3) $\gamma = (d_C r^2)^\# \lrcorner \omega$, $d_C \gamma = 4\omega$ and $d_C((d_C r^2)^\# \lrcorner \text{Re } \Omega) = 6\text{Re } \Omega$ or
- (4) $\gamma = (d_C(r^2\phi_{0,j}))^\# \lrcorner \omega$ with $\mu_{0,j} = 12$, $\lambda = 1$ and $d_C((d_C r^2\phi_{0,j})^\# \lrcorner \text{Re } \Omega)$ equals to a linear combination of $\text{Re } \Omega$ and $\text{Im } \Omega$.

Proof By Propositions 5.10 and 5.11, up to linear combinations, either (1) or (2) holds or

$$\gamma = r^2\phi_{1,j} = (Jr^2\phi_{1,j})^\# \lrcorner \omega$$

with $\mu_{1,j} = 8$. Remark that

$$d_C((Jr^2\phi_{1,j})^\# \lrcorner \text{Re } \Omega) = d_C((r^2\phi_{1,j})^\# \lrcorner \text{Im } \Omega)$$

is a multiple of $\text{Re } \Omega$ by Proposition 5.12.

The 1-form $r^2\phi_{1,j}$ is harmonic. By the $\text{SU}(3)$ structure, $Jr^2\phi_{1,j}$ is also a harmonic homogeneous 1-form of rate 1. By Lemma 5.11, $Jr^2\phi_{1,j}$ is a linear combination of $d_C r^2 = 2rdr$, $d_C(r^2\phi_{0,j'})$ with $\mu_{0,j'} = 12$, and $r^2\phi_{1,j''}$ with $\mu_{1,j''} = 8$. By Proposition 5.12, $d_C((r^2\phi_{1,j''})^\# \lrcorner \omega) = 0$ and $d_C((r^2\phi_{1,j''})^\# \lrcorner \text{Re } \Omega)$ is a multiple of $\text{Im } \Omega$. Since a linear combination of $r^2\phi_{1,j''}$ is closed if and only if it is 0, by Lemma 5.11, $(r^2\phi_{1,j''})^\# \lrcorner \omega$ equals to a linear combination of forms in (1) and (2). $\square$

Proposition 5.14 *A homogeneous 2-form γ with rate $\lambda \in [-3, 1]$ is harmonic if and only if up to linear combinations, either*

- (1) $\gamma = \phi_{2,1}$ with $\lambda = -2$,
- (2) $\gamma = d_C(r^{\lambda+2}\phi_{1,j})$ with $\lambda \in [0, 1]$ and $\mu_{1,j} = (\lambda + 2)(\lambda + 4)$,
- (3) $\gamma = \omega$ with $\lambda = 0$,
- (4) $\gamma = rdr^\# \lrcorner \text{Re } \Omega$ with $\lambda = 1$,
- (5) $\gamma = r^2\phi_{1,j}^\# \lrcorner \text{Re } \Omega$ with $\lambda = 1$, or
- (6)

$$\gamma = (d_C(r^{\lambda+2}\phi_{0,j}))^\# \lrcorner \text{Re } \Omega$$

with $\lambda \in (0, 1]$ and $\mu_{0,j} = (\lambda + 3)^2 - 4 \in (5, 12]$.

Proof There are two ways of decomposing homogeneous harmonic 2-forms. The first way is to decompose it as in Proposition 5.5. The other way is to decompose it into $(2,0)$, $(0,2)$, multiple of ω and primitive $(1,1)$ components. Assume that $\gamma = \gamma' + \gamma'' + \gamma'''$, where γ' is a linear combination of $(2,0)$, $(0,2)$, multiple of ω homogeneous harmonic forms, γ'' is a linear combination of all homogeneous harmonic forms in Proposition 5.5 except the type (7), and γ''' is primitive $(1,1)$ form of type (7) in Proposition 5.5. By Proposition 5.9, $\gamma''' = 0$. It is easy to see that only the type (4) and type (6) components of γ'' in Proposition 5.5 may be non-zero. They correspond to $\phi_{2,1}$ and $d_C(r^{\lambda+2}\phi_{1,j})$.

By the SU(3) structure, the multiple of ω component of γ' equals to a homogeneous harmonic function of rate λ times ω . By Proposition 5.5, it equals to a constant multiple of ω because $\mu_{0,2} > 5$. Still by the SU(3) structure, the (2,0) and (0,2) component must be the contraction of the metric dual of a homogeneous harmonic 1-form of rate λ with $\text{Re } \Omega$. By Proposition 5.10 and Lemma 5.11, it must be a linear combination of $r^2\phi_{1,j}$ with $\mu_{1,j} = 8$, $d_C(r\sqrt{\mu_{0,j'}+4}^{-2}\phi_{0,j'})$ with $\mu_{0,j'} = (\lambda+3)^2 - 4 \in (5, 12]$, and rdr . $\square$

Corollary 5.15 *Suppose that γ_2 and γ_3 are homogeneous 2-form and 3-form with same rate $\lambda \in (-2, 0]$ on C . Then γ_2 and γ_3 are both closed and coclosed if and only if for the 3-form γ defined as $d\theta \wedge \gamma_2 + \gamma_3$ on $C \times \mathbb{S}^1$, up to linear combinations, either*

- (1) $\gamma = \varphi = \text{Re } \Omega + d\theta \wedge \omega$ with $\lambda = 0$,
- (2) $\gamma = d_C((d_C r^2)^\# \lrcorner \varphi) = 6\text{Re } \Omega + 4d\theta \wedge \omega$ with $\lambda = 0$,
- (3) $\gamma = d_C(J(d_C r^2)^\# \lrcorner \varphi) = 6\text{Im } \Omega$ with $\lambda = 0$, or
- (4)

$$\gamma = d_C((d_C(r^{\lambda+2}\phi_{0,j}))^\# \lrcorner \varphi)$$

with $\lambda \in (-1, 0]$ and $\mu_{0,j} = (\lambda+4)^2 - 4 \in (5, 12]$.

Proof Decompose γ_2 and γ_3 as in Proposition 5.5. The closedness implies that the type (5) and type (7) components in Proposition 5.5 vanish. The coclosedness implies that the type (1) and type (2) components also vanish. The type (3) and type (6) components also vanish by the assumption on p and λ . So γ_2 and γ_3 are of type (4) in Proposition 5.5. This implies that they are in the image of d_C acting on homogeneous harmonic 1-forms or 2-forms. The result follows easily from Propositions 5.13 and 5.14. $\square$

6 Doubling Construction of Calabi–Yau Threefolds

This section proves Theorem 1.3. The notations in Theorem 1.3 are used in this section.

Recall that in the setting of Theorem 1.3, M is glued by $V_\pm$. The first goal is to study the operator $d + d^*$ from odd-degree forms to even-degree forms on $\mathbb{S}^1$ times C , $V_\pm$ or M . Let θ be the standard variable on $\mathbb{S}^1$. Then any odd-degree form can be expressed as

$$\begin{aligned} \gamma &= \gamma^1 + \gamma^3 + \gamma^5 + \gamma^7 \\ &= (d\theta \wedge \gamma_0 + \gamma_1) + (d\theta \wedge \gamma_2 + \gamma_3) + (d\theta \wedge \gamma_4 + \gamma_5) + (d\theta \wedge \gamma_6), \end{aligned}$$

where γ_p is a degree p -form on each slice. A direct calculation shows that

$$\begin{aligned} (d_{\mathbb{S}^1 \times C} + d_{\mathbb{S}^1 \times C}^*)\gamma &= \sum_{p=0,2,4,6} \left(d_C \gamma_{p-1} + d_C^* \gamma_{p+1} - \frac{\partial \gamma_p}{\partial \theta} \right) \\ &\quad - \sum_{p=1,3,5} d\theta \wedge \left(d_C \gamma_{p-1} + d_C^* \gamma_{p+1} - \frac{\partial \gamma_p}{\partial \theta} \right), \end{aligned}$$

where d_C and d_C^* mean doing the d and d^* operators on each slice. C may be replaced by $V_{\pm}$ or M .

Similarly, any even-degree form can be expressed as

$$\begin{aligned}\gamma &= \gamma^0 + \gamma^2 + \gamma^4 + \gamma^6 \\ &= (\gamma_0) + (d\theta \wedge \gamma_1 + \gamma_2) + (d\theta \wedge \gamma_3 + \gamma_4) + (d\theta \wedge \gamma_5 + \gamma_6),\end{aligned}$$

Another direct calculation shows that

$$\begin{aligned}(d_{\mathbb{S}^1 \times C} + d_{\mathbb{S}^1 \times C}^*)\gamma &= - \sum_{p=0,2,4,6} d\theta \wedge \left(d_C \gamma_{p-1} + d_C^* \gamma_{p+1} - \frac{\partial \gamma_p}{\partial \theta} \right) \\ &\quad + \sum_{p=1,3,5} \left(d_C \gamma_{p-1} + d_C^* \gamma_{p+1} - \frac{\partial \gamma_p}{\partial \theta} \right).\end{aligned}$$

When γ is $\mathbb{S}^1$ -invariant, then there is no $\frac{\partial}{\partial \theta}$ part. Therefore, it suffices to study $d + d^*$ on C , M or $V_{\pm}$.

Proposition 6.1 *Suppose that $\delta > 0$ is small enough.*

$$\gamma \in W_{-3-\delta, \dots, -3-\delta}^{k,2}(\Lambda^{\text{even}}(V_{\pm}))$$

or

$$\gamma \in W_{-3-\delta, \dots, -3-\delta}^{k,2}(\Lambda^{\text{even}}(M)).$$

If $(d + d^*)\gamma = 0$, then $d\gamma = d^*\gamma = 0$.

Proof Assume that $\gamma \in W_{-3-\delta, \dots, -3-\delta}^{k,2}(\Lambda^{\text{even}}(M))$. Using the definition of the Hodge Laplacian Δ , it is easy to see that $\Delta\gamma_0 = \Delta\gamma_2 = \Delta\gamma_4 = \Delta\gamma_6 = 0$. Near each singular point $x_i \in V_{\pm}^{\text{sing}}$, $\gamma \in W_{-3-\delta}^{k,2}$. By Propositions 5.5, 5.6 and 5.7, there is no 0-form in $\mathcal{K}_i(\lambda_i)$ with rate $\lambda_i \in (-3-\delta, -\delta)$. So by the definition of $\mathcal{P}_i(\lambda_i)$ and Theorem 2.19, $\gamma_0 \in W_{-2+\delta}^{k,2}(\Lambda^2 M)$.

Similarly, by Propositions 5.6, 5.14, and Theorem 2.19, γ_2 can be written as the linear combination of $\phi_{2,1}$, $\log r \phi_{2,1}$ and an element in $W_{-2+\delta}^{k,2}$ near x_i . Remark that the difference between $\phi_{2,1}$ or $\log r \phi_{2,1}$ and the corresponding element in $\mathcal{P}_i(\lambda_i)$ lies in $W_{-2+\delta}^{k,2}$. By Hodge duality on M , γ_4 can be written as the linear combination of $*\phi_{2,1}$, $*\log r \phi_{2,1}$ and an element in $W_{-2+\delta}^{k,2}$ near x_i . However, the equation

$$d\gamma_2 + d^*\gamma_4 = 0$$

implies that the coefficients of the log terms vanish by Proposition 5.2. So near x_i , $d\gamma_p \in W_{-3+\delta}^{k,2}$ for $p = 0, 2, 4$. Globally, $d\gamma_p \in W_{-3+\delta, \dots, -3+\delta}^{k,2}$.

It is easy to see that the boundary term in the integral

$$(d\gamma_p, d\gamma_p)_{L^2(\{r_i \geq r_0\})} - (\gamma_p, d^*d\gamma_p)_{L^2(\{r_i \geq r_0\})}$$

goes to 0 when r_0 goes to 0. Since $d^*d\gamma_p = -d^*d^*\gamma_{p+2} = 0$, it follows that $d\gamma_p = 0$. The $V_{\pm}$ case is similar. $\square$

Recall the definitions of $V_{\pm}$ and $\tilde{V}_{\pm}$ in Sect. 4. In this section, remark that the definitions of V_+ and $\tilde{V}_+$ remain unchanged but the definitions of V_- and $\tilde{V}_-$ have been changed to another copy of V_+ and $\tilde{V}_+$. By definition, $\tilde{V}_{\pm}$ is the small resolution of $V_{\pm}$. Locally, near each point $x \in V_{\pm}^{\text{sing}}$, the neighborhood of x is topologically a cone over $\mathbb{S}^2 \times \mathbb{S}^3$. The corresponding set in $\tilde{V}_{\pm}$ is topologically $\mathbb{S}^2 \times \mathbb{B}^4$, where x is replaced by $\mathbb{S}^2 \times \{0\}$.

The next goal is to study the Hodge theory on $V_{\pm}$. It was pioneered by Cheeger [8] using a slightly difference version of weighted analysis and followed by many people including Melrose [52].

Recall that $F_{\infty} = \mathbb{S}^1 \times S_{\pm}$. As in Sect. 6, by Hilbert–Schmidt theorem, assume that $d_{F_{\infty}}\phi_{p-1,j,\infty}$ and $\phi_{p,j,\infty}$ are orthogonal basis for $L^2(\Lambda^p(F_{\infty}))$ satisfying $\Delta_{F_{\infty}}\phi_{p,j,\infty} = \mu_{p,j,\infty}\phi_{p,j,\infty}$. Moreover, let $h_{p,\infty}$ be the p -th Betti number of F_{∞} . Then

Lemma 6.2 *Consider the Hodge Laplacian operator Δ acting on p -forms on $[T, \infty) \times F_{\infty}$.*

$$\mathcal{K}_{\infty}(0) = \text{Span}\{\phi_{p,j,\infty}, t\phi_{p,j,\infty}\}_{j=1}^{h_{p,\infty}} \oplus \text{Span}\{dt \wedge \phi_{p-1,j,\infty}, tdt \wedge \phi_{p-1,j,\infty}\}_{j=1}^{h_{p-1,\infty}}.$$

Proof Any $\gamma \in \mathcal{K}_{\infty}(0)$ can be written as

$$\gamma = dt \wedge \alpha + \beta.$$

Then

$$\Delta_{[T,\infty) \times \mathbb{S}^1 \times S_{\pm}} \gamma = dt \wedge \left(\Delta_{\mathbb{S}^1 \times S_{\pm}} - \frac{\partial^2}{\partial t^2} \right) \alpha + \left(\Delta_{\mathbb{S}^1 \times S_{\pm}} - \frac{\partial^2}{\partial t^2} \right) \beta = 0.$$

Consider the self-adjoint operator $\Delta_{\mathbb{S}^1 \times S_{\pm}}$. Write β as

$$\beta = \sum_{j=h_{p-1,\infty}+1}^{\infty} \beta_{p-1,j}(t) d_{F_{\infty}}\phi_{p-1,j,\infty} + \sum_{j=1}^{\infty} \beta_{p,j}(t) \phi_{p,j,\infty},$$

then

$$\left(\mu_{p,j,\infty} - \frac{d^2}{dt^2} \right) \beta_{p,j} = \left(\mu_{p-1,j,\infty} - \frac{d^2}{dt^2} \right) \beta_{p-1,j} = 0.$$

If $\mu_{p,j,\infty} \neq 0$, $\beta_{p,j}$ is a linear combination of $r^{\pm\sqrt{\mu_{p,j,\infty}}}$. When $\mu_{p,j,\infty} = 0$, $\beta_{p,j}$ is a linear combination of 1 and t . On the other hand, $\beta_{p-1,j}$ is always a linear combination of $r^{\pm\sqrt{\mu_{p-1,j,\infty}}}$. The result for α is similar. $\square$

Lemma 6.3 (Poincaré lemma) *Suppose that $\gamma \in W_{-\delta}^{k,2}(\Lambda^p([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$ is closed for small enough $\delta > 0$. Then there exist $\tilde{\gamma} \in W_{-\delta}^{k+1,2}(\Lambda^{p+1}([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$ and unique constants γ_j such that*

$$\gamma = d\tilde{\gamma} + \sum_{j=1}^{h_{p,\infty}} \gamma_j \phi_{p,j,\infty}.$$

Suppose that γ is a closed form in $W_{\delta}^{k,2}(\Lambda^p([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$ instead. Then there exists $\tilde{\gamma} \in W_{\delta}^{k+1,2}(\Lambda^{p+1}([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$ such that $\gamma = d\tilde{\gamma}$.

Similarly, suppose that γ is a closed form in $W_{-2+\delta}^{k,2}(\Lambda^2(V_{\pm} \cap \{r_x \leq r_0\}))$ for small enough δ and r_0 . Then there exists $\tilde{\gamma} \in W_{-1+\delta}^{k+1,2}(\Lambda^1(V_{\pm} \cap \{r_x \leq r_0\}))$ such that $\gamma = d\tilde{\gamma}$.

Proof Write γ as

$$\begin{aligned} \gamma = dt \wedge & \left(\sum_{j=h_{p-2,\infty}+1}^{\infty} \alpha_{p-2,j}(t) d_{F_{\infty}} \phi_{p-2,j,\infty} + \sum_{j=1}^{\infty} \alpha_{p-1,j}(t) \phi_{p-1,j,\infty} \right) \\ & + \sum_{j=h_{p-1,\infty}+1}^{\infty} \beta_{p-1,j}(t) d_{F_{\infty}} \phi_{p-1,j,\infty} + \sum_{j=1}^{\infty} \beta_{p,j}(t) \phi_{p,j,\infty}, \end{aligned}$$

then

$$\begin{aligned} 0 = d_{[T,\infty) \times F_{\infty}} \gamma \\ = dt \wedge & \left(- \sum_{j=h_{p-1,\infty}+1}^{\infty} \alpha_{p-1,j}(t) d_{F_{\infty}} \phi_{p-1,j,\infty} \right. \\ & + \sum_{j=h_{p-1,\infty}+1}^{\infty} \frac{d\beta_{p-1,j}}{dt}(t) d_{F_{\infty}} \phi_{p-1,j,\infty} + \sum_{j=1}^{\infty} \frac{d\beta_{p,j}}{dt}(t) \phi_{p,j,\infty} \Big) \\ & + \beta_{p,j}(t) d_{F_{\infty}} \phi_{p,j,\infty}. \end{aligned}$$

So $\beta_{p,j}$ are constants. Moreover, they vanish unless $j = 1, 2, \dots, h_{p,\infty}$. Define $\tilde{\alpha}_{p-2,j}(t)$ as

$$-\mu_{p-2,j,\infty} e^{\sqrt{\mu_{p-2,j,\infty}} t} \int_t^{\infty} e^{-2\sqrt{\mu_{p-2,j,\infty}} \tau} \int_T^{\tau} e^{\sqrt{\mu_{p-2,j,\infty}} s} \alpha_{p-2,j}(s) ds d\tau.$$

Define $\tilde{\beta}_{p-2,j}(t)$ as $\frac{1}{\mu_{p-2,j,\infty}} \frac{d}{dt} \tilde{\alpha}_{p-2,j}(t)$. Then it is easy to see that

$$\frac{d}{dt} \tilde{\alpha}_{p-2,j} - \mu_{p-2,j,\infty} \tilde{\beta}_{p-2,j} = 0$$

and

$$-\tilde{\alpha}_{p-2,j} + \frac{d}{dt}\tilde{\beta}_{p-2,j} = \alpha_{p-2,j}.$$

When $\gamma \in W_{-\delta}^{k,2}(\Lambda^p([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$, define $\tilde{\gamma}$ as

$$\begin{aligned} \tilde{\gamma} = dt \wedge & \sum_{j=h_{p-2,\infty}+1}^{\infty} \tilde{\alpha}_{p-2,j}(t)\phi_{p-2,j,\infty} + \sum_{j=h_{p-2,\infty}+1}^{\infty} \tilde{\beta}_{p-2,j}(t)d_{F_{\infty}}\phi_{p-2,j,\infty} \\ & + \sum_{j=1}^{h_{p-1,\infty}} \left(\int_T^t \alpha_{p-1,j}(\tau)d\tau \right) \phi_{p-1,j,\infty} + \sum_{j=h_{p-1,\infty}}^{\infty} \beta_{p-1,j}(t)\phi_{p-1,j,\infty}, \end{aligned}$$

then $\gamma = d\tilde{\gamma} + \sum_{j=1}^{h_{p,\infty}} \beta_{p,j}(t)\phi_{p,j,\infty}$ and $d_{[T,\infty) \times F_{\infty}}^* \tilde{\gamma} = 0$.

When $\gamma \in W_{\delta}^{k,2}(\Lambda^p([T, \infty) \times \mathbb{S}^1 \times S_{\pm}))$, define $\tilde{\gamma}$ as

$$\begin{aligned} \tilde{\gamma} = dt \wedge & \sum_{j=h_{p-2,\infty}+1}^{\infty} \tilde{\alpha}_{p-2,j}(t)\phi_{p-2,j,\infty} + \sum_{j=h_{p-2,\infty}+1}^{\infty} \tilde{\beta}_{p-2,j}(t)d_{F_{\infty}}\phi_{p-2,j,\infty} \\ & + \sum_{j=1}^{h_{p-1,\infty}} \left(- \int_t^{\infty} \alpha_{p-1,j}(\tau)d\tau \right) \phi_{p-1,j,\infty} + \sum_{j=h_{p-1,\infty}}^{\infty} \beta_{p-1,j}(t)\phi_{p-1,j,\infty}, \end{aligned}$$

then $\gamma = d\tilde{\gamma}$ and $d_{[T,\infty) \times F_{\infty}}^* \tilde{\gamma} = 0$.

The estimate on $\tilde{\gamma}$ is standard.

The $x \in V_{\pm}^{\text{sing}}$ case is similar. □

There is a natural map e from the relative deRham cohomology group of $\tilde{V}_{\pm}$ to the absolute deRham cohomology group. By Section 6.4 of [52], the image $e[\mathcal{H}_{\text{dR,rel}}^2(\tilde{V}_{\pm})]$ is isomorphic to the space

$$\frac{\{\gamma \in C_0^{\infty}(\Lambda^2(\tilde{V}_{\pm})), d\gamma = 0\}}{\{\gamma \in C_0^{\infty}(\Lambda^2(\tilde{V}_{\pm})), \gamma = d\gamma', \gamma' \in \cap_{k=1}^{\infty} W_{-\delta}^{k,2}(\Lambda^1(\tilde{V}_{\pm}))\}}$$

for small enough $\delta > 0$.

The next goal is to show that.

Proposition 6.4 *Suppose that $\delta > 0$ is small enough. Then the space*

$$e[\mathcal{H}_{\text{dR,rel}}^2(\tilde{V}_{\pm})] = \frac{\{\gamma \in C_0^{\infty}(\Lambda^2(\tilde{V}_{\pm})), d\gamma = 0\}}{\{\gamma \in C_0^{\infty}(\Lambda^2(\tilde{V}_{\pm})), \gamma = d\gamma', \gamma' \in \cap_{k=1}^{\infty} W_{-\delta}^{k,2}(\Lambda^1(\tilde{V}_{\pm}))\}}$$

is isomorphic to the space $\mathcal{H}_{b,\text{Ho}}^2(V_{\pm})$ defined as

$$\{\gamma \in W_{-3-\delta, \dots, -3-\delta, \delta}^{k,2}(\Lambda^2(V_{\pm})), (d + d^*)\gamma = 0\}.$$

Proof The method in this proof is the combination of the results in Sect. 6.4 of Melrose's book [52]. Suppose that $\gamma \in C_0^\infty(\Lambda^2(\tilde{V}_\pm))$ is closed. Remark that $H^2(\mathbb{S}^2 \times \mathbb{B}^4) = \mathbb{R}$. So $\gamma = d\gamma_x + \gamma_{2,1,x}\phi_{2,1,x}$ on $\mathbb{S}^2 \times \mathbb{B}^4$ corresponding to $x \in V_\pm^{\text{sing}}$, where $\phi_{2,1,x}$ is the pull back of the generator of $H^2(\mathbb{S}^2)$. Choose a cut-off function χ_x which is supported near x and is 1 in a smaller neighborhood. Then

$$\gamma' = \gamma - \sum_{x \in V_\pm^{\text{sing}}} d(\chi_x \gamma_x)$$

equals to $\gamma_{2,1,x}\phi_{2,1,x}$ near x . It is a 2-form on $V_\pm$ if we replace the form $\gamma_{2,1,x}\phi_{2,1,x}$ on $\tilde{V}_\pm$ by $\gamma_{2,1,x}\phi_{2,1,x}$ on $V_\pm$ near x .

Define $\mathcal{H}_{b,\text{Ho}}^p(V_\pm)$ as

$$\{\gamma \in W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}(\Lambda^p(V_\pm)), (d + d^*)\gamma = 0\}$$

for $p = 0, 2, 4, 6$ and define $\mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm)$ as

$$\mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm) := \{\gamma \in W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}(\Lambda^{\text{even}}(V_\pm)), (d + d^*)\gamma = 0\}.$$

By Proposition 6.1,

$$\mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm) = \mathcal{H}_{b,\text{Ho}}^0(V_\pm) \oplus \mathcal{H}_{b,\text{Ho}}^2(V_\pm) \oplus \mathcal{H}_{b,\text{Ho}}^4(V_\pm) \oplus \mathcal{H}_{b,\text{Ho}}^6(V_\pm).$$

The L^2 dual of

$$d + d^* : W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^{\text{odd}}(V_\pm)) \rightarrow W_{-3+\delta,\dots,-3+\delta,-\delta}^{k-1,2}(\Lambda^{\text{even}}(V_\pm))$$

is

$$d + d^* : W_{-3-\delta,\dots,-3-\delta,\delta}^{1-k,2}(\Lambda^{\text{even}}(V_\pm)) \rightarrow W_{-4-\delta,\dots,-4-\delta,\delta}^{-k,2}(\Lambda^{\text{odd}}(V_\pm)).$$

The kernel of the dual map is $\mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm)$ by standard elliptic regularity. By the proof of Proposition 6.1, $\mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm) \subset W_{-2+\delta,\dots,-2+\delta,\delta}^{k,2}(\Lambda^{\text{even}}(V_\pm))$. So

$$W_{-3+\delta,\dots,-3+\delta,-\delta}^{k-1,2}(\Lambda^{\text{even}}(V_\pm)) = \mathcal{H}_{b,\text{Ho}}^{\text{even}}(V_\pm) \oplus ((d + d^*)(W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}))$$

by elliptic regularity. Moreover,

$$\begin{aligned} W_{-3+\delta,\dots,-3+\delta,-\delta}^{k-1,2}(\Lambda^2(V_\pm)) &= \mathcal{H}_{b,\text{Ho}}^2(V_\pm) \\ &\oplus dW_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^1(V_\pm)) \oplus d^*W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^3(V_\pm)). \end{aligned}$$

In fact, it suffices to show that the intersection of $dW_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^1(V_\pm))$ and $d^*W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^3(V_\pm))$ is $\{0\}$. Choose any element γ'' in the intersection. It is

harmonic. So by Theorem 2.19 and Lemma 6.2, it can be written as an element in $\mathcal{K}_\infty(0)$ plus an element in $W_\delta^{k-1,2}$ near infinity if δ is small enough. By the closedness and coclosedness, there are no $t\phi_{p,j,\infty}$ and $t dt \wedge \phi_{p-1,j,\infty}$ terms. By the exactness, coexactness and Lemma 6.3, there are no $\phi_{p,j,\infty}$ and $dt \wedge \phi_{p-1,j,\infty}$ terms. So $\gamma'' \in W_\delta^{k-1,2}$ near infinity. Using integration by parts, $\gamma'' = 0$.

There is a natural map $(r, (s_2, s_3)) \rightarrow (s_2, (r, s_3))$ from $C(\mathbb{S}^2 \times \mathbb{S}^3)$ to $\mathbb{S}^2 \times \mathbb{B}^4$ outside the singular point. It induces a map from $W^{k-1,2}(\Lambda^1(\mathbb{S}^2 \times \mathbb{B}^4))$ to $W_{-1-\delta}^{k-1,2}(\Lambda^1(C(\mathbb{S}^2 \times \mathbb{S}^3)))$. Using this map, it is easy to see that the projection of γ' to $\mathcal{H}_{b,\text{Ho}}^2(V_\pm)$ is a well-defined map from $e[\mathcal{H}_{\text{dR},\text{rel}}^2(\tilde{V}_\pm)]$ to $\mathcal{H}_{b,\text{Ho}}^2(V_\pm)$.

In order to show the injectivity, assume that γ is mapped to 0. Then there exists $\gamma''' \in W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^1(V_\pm))$ and $\gamma'''' \in W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^3(V_\pm))$ such that $\gamma' = d\gamma''' + d^*\gamma''''$. So $d\gamma'''$ is both exact and coexact on the end. So $d\gamma''' \in W_\delta^{k-1,2}$ near infinity as before. Using integration by parts, $d^*\gamma'''' = 0$. So $\gamma_{2,1,x}\phi_{2,1,x} = d\gamma'''$ when restricted to $V_\pm \cap \{r_x = r_0\} \simeq \mathbb{S}^2 \times \mathbb{S}^3$ for small enough r_0 . So $\gamma_{2,1,x} = 0$. Therefore $\gamma' = 0$ near x . Since $H^1(\mathbb{S}^2 \times \mathbb{S}^3) = 0$, $\gamma''' = d\gamma_x''''$ near x . So

$$\gamma' = d\left(\gamma''' - \sum_x d(\chi_x \gamma_x''')\right).$$

Therefore,

$$[\gamma] = \left[d\left(\gamma''' - \sum_x d(\chi_x \gamma_x''')\right) + \sum_x d(\chi_x \gamma_x) \right] = 0 \in e[\mathcal{H}_{\text{dR},\text{rel}}^2(\tilde{V}_\pm)].$$

In order to show the surjectivity, pick any form $\tilde{\gamma}$ in $\mathcal{H}_{b,\text{Ho}}^2(V_\pm)$. By Proposition 6.1, $\tilde{\gamma}$ can be written as $\tilde{\gamma} = \tilde{\gamma}_{2,1,x}\phi_{2,1,x} + \tilde{\gamma}'_x$ near each $x \in V_\pm^{\text{sing}}$, where $\tilde{\gamma}'_x \in W_{-2+\delta}^{k,2}$ near x . Since $\tilde{\gamma}'_x$ is closed and is in $W_{-2+\delta}^{k,2}$ near x , by Lemma 6.3 near x , there exists $\tilde{\gamma}''_x$ such that $\tilde{\gamma}'_x = d\tilde{\gamma}''_x$ near x . So $\tilde{\gamma}''' = \tilde{\gamma}_{2,1,x}\phi_{2,1,x} + \tilde{\gamma}'_x - d(\chi_x \tilde{\gamma}''_x)$ is a well defined form on $\tilde{V}_\pm$. On the other hand, by Lemma 6.3, there exists $\tilde{\gamma}''''$ such that $\tilde{\gamma}''' = d\tilde{\gamma}''''$ near infinity. Fix $\chi: \mathbb{R} \rightarrow [0, 1]$ as a smooth function satisfying $\chi(s) = 1$ for $s \leq 1$ and $\chi(s) = 0$ for $s \geq 2$, then $\tilde{\gamma}''' - d((1 - \chi(t_\pm - T + 2))\tilde{\gamma}''') \in C_0^\infty$ and its image approaches $\tilde{\gamma}$ when T goes to infinity. Therefore, the image of $e[\mathcal{H}_{\text{dR},\text{rel}}^2(\tilde{V}_\pm)]$ is dense. Since $\mathcal{H}_{b,\text{Ho}}^2(V_\pm)$ is finite dimensional, the map is in fact surjective. $\square$

Corollary 6.5 Suppose that $\delta > 0$ is small enough and

$$\gamma \in W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}(\Lambda^{\text{even}}(V_\pm)).$$

If $(d + d^*)\gamma = 0$, then $\gamma = 0$.

Proof By Proposition 6.1, $d\gamma_0 = 0$. So γ_0 is a constant. It vanishes because it decays at infinity. By Proposition 6.4, $\gamma_2 = 0$ because as in Proposition 5.38 of [43], the space $e[\mathcal{H}_{\text{dR},\text{rel}}^2(\tilde{V}_\pm)]$ vanishes. By Hodge duality, $\gamma_4 = \gamma_6 = 0$. $\square$

Similarly, it is possible to prove the following:

Proposition 6.6 *Suppose that $\delta > 0$ is small enough. Choose $\chi : \mathbb{R} \rightarrow [0, 1]$ as a smooth function satisfying $\chi(s) = 1$ for $s \leq 1$ and $\chi(s) = 0$ for $s \geq 2$. Define $\chi_\infty = (1 - \chi(t_\pm - T + 1))$. Then the space $\mathcal{H}_{b-\text{abs}}^2(V_\pm)$ defined as*

$$\{\gamma \in \text{Span}\{\chi_\infty \phi_{2,j,\infty}\}_{j=1}^{h_{2,\infty}} \oplus W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}(\Lambda^2(V_\pm)), (d + d^*)\gamma = 0\}$$

is isomorphic to the second absolute deRham cohomology group $H_{\text{dR},\text{abs}}^2(\tilde{V}_\pm)$. Define the space $\mathcal{H}_{b-\text{rel}}^2(V_\pm)$ as

$$\{\gamma \in \text{Span}\{\chi_\infty dt \wedge \phi_{1,j,\infty}\}_{j=1}^{h_{1,\infty}} \oplus W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}(\Lambda^2(V_\pm)), (d + d^*)\gamma = 0\}.$$

Then the space $\mathcal{H}_{eb}^2(V_\pm)$ defined as

$$\{\gamma \in W_{-3-\delta,\dots,-3-\delta,-\delta}^{k,2}(\Lambda^2(V_\pm)), (d + d^*)\gamma = 0\}$$

can be written as $\mathcal{H}_{eb}^2(V_\pm) = \mathcal{H}_{b-\text{abs}}^2(V_\pm) \oplus \mathcal{H}_{b-\text{rel}}^2(V_\pm)$

Proof Remark that

$$H_{\text{dR},\text{abs}}^2(\tilde{V}_\pm) \cong H_{\text{dR},\text{abs}}^2(\tilde{V}_\pm \cap \{t_\pm < T + 1\}),$$

where the isomorphism map is given by restriction. Given any form γ in $H_{\text{dR},\text{abs}}^2(\tilde{V}_\pm \cap \{t_\pm < T + 1\})$, using the fact that

$$H_{\text{dR},\text{abs}}^2(\tilde{V}_\pm \cap \{T < t_\pm < T + 1\}) = H^2([T, T + 1] \times \mathbb{S}^1 \times S_\pm) = H^2(S_\pm),$$

γ can be written as $\gamma = d\gamma_\infty + \phi_\infty$ on $t_\pm \in (T, T + 1)$, where $\phi_\infty \in H^2(S_\pm)$. As in the proof of Proposition 6.4,

$$\gamma' = \gamma - \sum_{x \in V_\pm^{\text{sing}}} d(\chi_x \gamma'_x) - d(\chi_\infty \gamma_\infty)$$

is a 2-form on $V_\pm$. Define γ'' as the projection of γ' to the second component in the decomposition

$$\begin{aligned} W_{-3+\delta,\dots,-3+\delta,-\delta}^{k-1,2}(\Lambda^2(V_\pm)) &= dW_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^1(V_\pm)) \\ &\oplus d^*W_{-2+\delta,\dots,-2+\delta,-\delta}^{k,2}(\Lambda^3(V_\pm)), \end{aligned}$$

as in the proof of Proposition 6.4. Remark that there is no $\mathcal{H}_{b,\text{Ho}}^2(V_\pm)$ component by Corollary 6.5. Then γ'' is closed and coexact. So it is harmonic. By Theorem 2.19 and Lemma 6.2, the leading term of γ'' near the end is the linear combination of $\phi_{2,j,\infty}$,

$t\phi_{2,j,\infty}$, $dt \wedge \phi_{1,j,\infty}$ and $tdt \wedge \phi_{1,j,\infty}$. By the closedness, $t\phi_{2,j,\infty}$ can not appear. By the coclosedness, $tdt \wedge \phi_{1,j,\infty}$ can not appear, too. The coexactness rules out the $dt \wedge \phi_{1,j,\infty}$ terms. So the leading term of γ'' must be the linear combination of $\phi_{2,j,\infty}$. In other words, $\gamma'' \in \mathcal{H}_{b-\text{abs}}^2(V_{\pm})$. It is easy to see that the map from γ to γ'' is a well-defined isomorphism from $\mathcal{H}_{\text{dR,abs}}^2(\tilde{V}_{\pm})$ to $\mathcal{H}_{b-\text{abs}}^2(V_{\pm})$.

Finally, given $\gamma \in \mathcal{H}_{eb}^2(V_{\pm}) \subset W_{-3+\delta, \dots, -3+\delta, -\delta}^{k-1, 2}(\Lambda^2(V_{\pm}))$, its first component is exact and coclosed. So as in the proof of Proposition 6.4, it belongs to $\mathcal{H}_{b-\text{rel}}^2(V_{\pm})$. On the other hand, its second component belongs to $\mathcal{H}_{b-\text{abs}}^2(V_{\pm})$. $\square$

Proposition 6.7 *Suppose $\delta > 0$ is small enough. Then the space*

$$\mathcal{H}_{\text{Ho}}^2(M) := \{\gamma \in W_{-3-\delta, \dots, -3-\delta}^{k, 2}(\Lambda^2(M)), (d + d^*)\gamma = 0\}.$$

is isomorphic to $H_{\text{dR,abs}}^2(\tilde{M})$, where $\tilde{M}$ is the small resolution of M .

Proof It is proved similarly as Proposition 6.6. Since the manifold is compact, there is no need to do anything near infinity. $\square$

Recall that $\tilde{M}$ is the gluing of $\tilde{V}_+$ and $\tilde{V}_-$ using $t_{\pm} = 2T + 1 - t_{\mp}$. Define t by $t = t_+ - T - \frac{1}{2} = T + \frac{1}{2} - t_-$. Using the fact that $\tilde{M} = \{t < \frac{1}{2}\} \cup \{t > -\frac{1}{2}\}$, there is a long exact sequence for the cohomology groups of $\tilde{M}$, $\{t < \frac{1}{2}\}$, $\{t > -\frac{1}{2}\}$ and $\{|t| < \frac{1}{2}\}$. In particular

$$H^2(\tilde{M}) \rightarrow H^2\left(\left\{t < \frac{1}{2}\right\}\right) \oplus H^2\left(\left\{t > -\frac{1}{2}\right\}\right) \rightarrow H^2\left(\left\{|t| < \frac{1}{2}\right\}\right)$$

is exact. Remark that $\tilde{V}_+$ is isomorphic to $\tilde{V}_-$, so the map from $H^2(\{t > -\frac{1}{2}\})$ to $H^2(\{|t| < \frac{1}{2}\})$ is isomorphic to the map from $H^2(\{t < \frac{1}{2}\})$ to $H^2(\{|t| < \frac{1}{2}\})$. This map is injective by the proof of Proposition 5.38 of [43]. It follows that the long exact sequence is reduced to

$$0 \rightarrow H^1\left(\left\{|t| < \frac{1}{2}\right\}\right) \rightarrow H^2(\tilde{M}) \rightarrow H^2\left(\left\{t < \frac{1}{2}\right\}\right) \rightarrow 0$$

using the fact that $H^1(\{t < \frac{1}{2}\}) \cap H^1(\{t > -\frac{1}{2}\}) = 0$. By Propositions 6.6 and 6.7, it induces a natural map from $\mathcal{H}_{\text{Ho}}^2(M) \simeq H^2(\tilde{M})$ to $\mathcal{H}_{b-\text{abs}}^2(V_+) \simeq H^2(\{t < \frac{1}{2}\})$. Moreover, $\dim \mathcal{H}_{\text{Ho}}^2(M) = \dim \mathcal{H}_{b-\text{abs}}^2(V_+) + 1$ because

$$H^1\left(\left\{|t| < \frac{1}{2}\right\}, \mathbb{R}\right) = H^1(\mathbb{S}^1 \times S_+) = \mathbb{R}.$$

Proposition 6.8 *There exists a linear map from $\mathcal{H}_{\text{Ho}}^{\text{even}}(M)$ to $\mathcal{H}_{eb}^{\text{even}}(V_+)$ such that if $\gamma \in \mathcal{H}_{\text{Ho}}^{\text{even}}(M)$ is mapped to γ' , then*

$$\left(\gamma, \gamma' \chi\left(t_+ - \frac{T}{2}\right)\right)_{W_{-3-\delta, \dots, -3-\delta, -\delta}^{k, 2}} \geq \frac{9}{10} \|e^{-\delta(t+T+\frac{1}{2}x)} \gamma\|_{W_{-3-\delta, \dots, -3-\delta}^{k, 2}}^2$$

for large enough T .

Proof Map 1 to 1 and the volume form $*1$ to $*1$. The estimate is trivial for such components. Using Hodge star, it suffices to define the map for 2-forms. Assume that γ is a 2-form and by normalization, $\|e^{-\delta(t+T+\frac{1}{2})}\gamma\|_{W_{-3-\delta,\dots,-3-\delta}^{k,2}} = 1$.

Since the difference between the asymptotically cylindrical metric and the product metric on the cylinder is $O(e^{-\nu t_{\pm}})$, it is easy to see that

$$\|(d + d^*)_{\infty}\gamma\|_{W^{k-1,2}(|t| < \frac{T}{2})} \leq C e^{-\frac{\nu T}{2} + \frac{3\delta T}{2}},$$

where $(d + d^*)_{\infty}$ is the operator $d + d^*$ defined using the product metric on the cylinder.

Using generalized Fourier series, γ can be written as $dt \wedge \alpha + \beta + \gamma''$ in $|t| < \frac{T}{2}$, where $\alpha \in \mathcal{H}^1(t = 0) \cong \mathbb{R}$, $\beta \in \mathcal{H}^2(t = 0) \cong H^2(S_{\pm})$ and γ'' is an exact form satisfying $\|\gamma''\|_{W^{k,2}(|t| < \frac{1}{2})} \leq C e^{-\frac{\nu T}{4}}$ if both ν and $\frac{\delta}{\nu}$ are small enough. Choose γ''' such that $\|\gamma'''\|_{W^{k+1,2}(|t| < \frac{1}{2})} \leq C e^{-\frac{\nu T}{4}}$ and $d\gamma''' = \gamma''$ when $|t| < \frac{1}{2}$. It is clear that $\gamma - d((1 - \chi(t + \frac{3}{2}))\gamma''')$ induces a form in $W_{-3-\delta,\dots,-3-\delta,-\delta}^{k,2}(V_+)$ which equals to $dt \wedge \alpha + \beta$ when $t_+ > T + 1$. Define $\gamma'_{b-\text{abs}}$ as its $d^*W_{-2+\delta,\dots,-2+\delta,-\delta}^{k+1,2}(\Lambda^3(V_+))$ component. As in the proof of Proposition 6.6, $\gamma'_{b-\text{abs}} \in \mathcal{H}_{b-\text{abs}}^2(V_+)$ and $\gamma'_{b-\text{abs}} - \beta \in W_{\delta}^{k,2}(V_+)$ near infinity. By Propositions 6.6 and 6.7, $\gamma'_{b-\text{abs}}$ is also the image of γ using the restriction map from $H_{\text{dR,abs}}^2(\tilde{M})$ to $H_{\text{dR,abs}}^2(\tilde{V}_+ \cap \{t < \frac{1}{2}\})$.

On the other hand, $*_M\gamma - *_\infty(dt \wedge \alpha + \beta)$ can also be written as $d\gamma''''$ when $|t| < \frac{1}{2}$ for $\|\gamma''''\|_{W^{k+1,2}(|t| < \frac{1}{2})} \leq C e^{-\frac{\nu T}{4}}$. So

$$*_{V_+} \left(*_M\gamma - d \left(\left(1 - \chi \left(t + \frac{3}{2} \right) \right) \gamma'''' \right) \right)$$

also induces a form in $W_{-3-\delta,\dots,-3-\delta,-\delta}^{k,2}(V_+)$ which equals to $*_{V_+} *_\infty(dt \wedge \alpha + \beta)$ when $t_+ > T + 1$. Define $\gamma'_{b-\text{rel}}$ as its $dW_{-2+\delta,\dots,-2+\delta,-\delta}^{k+1,2}(\Lambda^1(V_+))$ component. It belongs to $\mathcal{H}_{b-\text{rel}}^2$ and $\gamma'_{b-\text{abs}} - dt \wedge \alpha \in W_{\delta}^{k,2}$ near infinity.

Define $\gamma' = \gamma'_{b,\text{abs}} + \gamma'_{b-\text{rel}}$. Then $\gamma - d((1 - \chi(t + \frac{3}{2}))\gamma''')$ can be written as $\gamma' + \gamma''''$ for an element $\gamma'''' \in W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}$.

By weighted elliptic estimate and Corollary 6.5,

$$\|\gamma''''\|_{W_{-3-\delta,\dots,-3-\delta,\delta}^{k,2}} \leq C \|(d + d^*)\gamma''''\|_{W_{-4+\delta,\dots,-4+\delta,\delta}^{k-1,2}} \leq C e^{-\frac{\nu T}{8}}.$$

So

$$\begin{aligned} \|\gamma'\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(t_+ < \frac{T}{2})} &\leq \|\gamma\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(t_+ < \frac{T}{2})} + \|\gamma''''\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(t_+ < \frac{T}{2})} \\ &\leq 1 + C e^{-\frac{\nu T}{8}} e^{\delta T}. \end{aligned}$$

Using the fact that γ' is asymptotic to $dt \wedge \alpha + \beta$,

$$\|\gamma'\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(t > \frac{T}{2})} \leq C e^{-\frac{\delta T}{2}} \|\gamma'\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(t < \frac{T}{2})}.$$

In particular,

$$\|\gamma\|_{W^{k,2}(|t| \leq \frac{1}{2})} \leq C e^{\frac{\delta T}{2}}.$$

Using

$$\begin{aligned} & \left\| \left(1 - \chi \left(t + \frac{3}{2} \right) \right) \gamma \right\|_{W_{-3-\delta,\dots,-\delta}^{k,2}(V_-)} \\ & \leq C \left\| (d + d^*)_{V_-} \left(1 - \chi \left(t + \frac{3}{2} \right) \right) \gamma \right\|_{W_{-4-\delta,\dots,-\delta}^{k-1,2}(V_-)}, \end{aligned}$$

it is easy to get the conclusion. $\square$

Proposition 6.9 *For large enough T , and $p = 2, 3$, there exists a linear map*

$$B_T : W_{-2+\delta,\dots,-2+\delta}^{k,2}(\Lambda^p M) \rightarrow W_{-2+\delta,\dots,-2+\delta}^{k,2}(\Lambda^p M)$$

such that $(d + d^)B_T\gamma = d^*\gamma$ and*

$$\|B_T\gamma\|_{W_{-2+\delta,\dots,-2+\delta}^{k,2}(\Lambda^p M)} \leq C e^{3\delta T} \|\gamma\|_{W_{-2+\delta,\dots,-2+\delta}^{k,2}(\Lambda^p M)}.$$

Proof When $p = 3$, define

$$A_{\pm} : W_{-2+\delta,\dots,-2+\delta,0}^{k,2}(\Lambda^{\text{odd}}(V_{\pm})) \rightarrow W_{-3+\delta,\dots,-3+\delta,0}^{k-1,2}(\Lambda^{\text{even}}(V_{\pm}))$$

and

$$A_T : W_{-2+\delta,\dots,-2+\delta}^{k,2}(\Lambda^{\text{odd}}(M)) \rightarrow W_{-3+\delta,\dots,-3+\delta}^{k-1,2}(\Lambda^{\text{even}}(M))$$

as $A_{\pm}\gamma = e^{\mp\delta t_{\pm}}(d + d^*)(e^{\pm\delta t_{\pm}}\gamma)$ and $A_T\gamma = e^{-\delta t}(d + d^*)(e^{\delta t}\gamma)$. Define the asymptotic kernel as $\chi(t_{\pm} - \frac{T}{2})\text{Ker}A_{\pm}$. The L^2 -dual maps are

$$A_{\pm}^* : W_{-3-\delta,\dots,-3-\delta,0}^{1-k,2}(\Lambda^{\text{even}}(V_{\pm})) \rightarrow W_{-4-\delta,\dots,-4-\delta,0}^{-k,2}(\Lambda^{\text{odd}}(V_{\pm}))$$

and

$$A_T^* : W_{-3-\delta,\dots,-3-\delta}^{1-k,2}(\Lambda^{\text{even}}(M)) \rightarrow W_{-4-\delta,\dots,-4-\delta}^{-k,2}(\Lambda^{\text{odd}}(M))$$

defined as $A_{\pm}^* \gamma = e^{\pm \delta t} (d + d^*) (e^{\mp \delta t} \gamma)$ and $A_T^* \gamma = e^{\delta t} (d + d^*) (e^{-\delta t} \gamma)$. A_T induces a map A_T' from the L^2 complement to the asymptotic kernels of $A_{\pm}$ to the L^2 complement to the asymptotic cokernels. By Proposition 4.2 of [44], A_T' is bijective.

By Proposition 6.8, there exists an injective map from $\mathcal{H}_{\text{Ho}}^2(M)$ to $\mathcal{H}_{eb}^2(V_+)$. However, by Proposition 6.5 and the fact that $H^1(\mathbb{S}^1 \times S_+) = \mathbb{R}$, the dimension of $\mathcal{H}_{b\text{-rel}}^2(V_+)$ is at most 1 but the dimension of $\mathcal{H}_{\text{Ho}}^2(M)$ equals to the dimension of $\mathcal{H}_{b\text{-abs}}^2(V_+)$ plus 1. So the map from $\mathcal{H}_{\text{Ho}}^2(M)$ to $\mathcal{H}_{eb}^2(V_+)$ is bijective. By Hodge duality, the map from $\mathcal{H}_{\text{Ho}}^{\text{even}}(M)$ to $\mathcal{H}_{eb}^{\text{even}}(V_+)$ is also bijective. Therefore, using Proposition 6.8, it is easy to see that the map A_T from the L^2 complement of the asymptotic kernels of $A_{\pm}$ to the L^2 complement of the kernel of A_T^* is also bijective.

For $\gamma \in W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^3 M)$, $e^{-\delta t} d^* \gamma$ is in the L^2 complement of the kernel of A_T^* by Proposition 6.1. So there exists an element γ' in the L^2 complement of the asymptotic kernels of $A_{\pm}$ such that $A_T \gamma' = e^{-\delta t} d^* \gamma$. Let $\gamma'' = e^{\delta t} \gamma'$, then it is easy to see that $d\gamma_3'' + d^* \gamma_5'' = 0$ and $d^* \gamma_3'' + d\gamma_1'' = d^* \gamma$. Using integration by parts, $d^* \gamma_5'' = d\gamma_1'' = 0$. So $B_T \gamma = \gamma_3''$ in this case. The estimate for $B_T \gamma$ follows from Proposition 4.2 of [44].

When $p = 2$, consider the Laplacian operator

$$\Delta : W_{-1+\delta, \dots, -1+\delta}^{k+1,2}(\Lambda^1 M) \rightarrow W_{-3+\delta, \dots, -3+\delta}^{k-1,2}(\Lambda^1 M).$$

The L^2 dual map is

$$\Delta : W_{-3-\delta, \dots, -3-\delta}^{1-k,2}(\Lambda^1 M) \rightarrow W_{-5-\delta, \dots, -3-\delta}^{-1-k,2}(\Lambda^1 M).$$

By Propositions 5.5, 5.6, 5.7, and 5.8, there is no element in $\mathcal{P}_i(\lambda_i)$ for $\text{Re } \lambda_i \in (-4, 0)$. So by standard elliptic regularity and Theorem 2.19, any element in the second kernel also lies in $W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^1 M)$. So using integration by parts, it is closed and coclosed. Therefore, for any element $\gamma \in W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^2 M)$, $d^* \gamma$ is L^2 -perpendicular to the kernel of the second map. So it lies in the image of the first map. Let γ' be its inverse. Then $dd^* \gamma' + d^* d\gamma' = d^* \gamma$. Using integration by parts, $dd^* \gamma' = 0$. Define $B_T \gamma$ as $d\gamma'$, then $(d + d^*) B_T \gamma = d^* \gamma$. Moreover, integration by parts again implies that $\|B_T \gamma\|_{L^2} \leq \|\gamma\|_{L^2} \leq C e^{\delta T} \|\gamma\|_{W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^2 M)}$.

By standard elliptic regularity,

$$\|B_T \gamma\|_{W_{-3, \dots, -3}^{k,2}} \leq C (\|B_T \gamma\|_{L^2} + \|\Delta B_T \gamma\|_{W_{-5, \dots, -5}^{k-2,2}}) \leq C e^{\delta T} \|\gamma\|_{W_{-2+\delta, \dots, -2+\delta}^{k,2}}.$$

So the required estimate is obtained using Theorem 2.19. $\square$

Remark 6.10 The $p = 2$ case of Proposition 6.9 is similar to Theorem A of [38]. The $p = 3$ case of Proposition 6.9 is similar to Proposition 5.40 of [43] in the smooth twisted connected sum case. However, the author is not able to understand the proof of Proposition 5.40 of [43]. This does not affect the main result of [43] because in the smooth case, Proposition 5.40 can be proved using Theorem A of Joyce's paper [38]. In the singular case, it is not possible to find analogue of Theorem A of Joyce's paper

[38] in $p = 3$ case. That is the reason to make full use of the Hodge theory in this paper.

Using the identification $\gamma = d\theta \wedge \gamma_2 + \gamma_3$, B_T can also be viewed as a map from $\mathbb{S}^1$ -invariant 3-form on $\mathbb{S}^1 \times M$ to itself satisfying $(d + d^*)B_T\gamma = d^*\gamma$ and the estimate

$$\|B_T\gamma\|_{W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^3(\mathbb{S}^1 \times M))} \leq Ce^{3\delta T} \|\gamma\|_{W_{-2+\delta, \dots, -2+\delta}^{k,2}(\Lambda^3(\mathbb{S}^1 \times M))}.$$

The next goal is the improvement of the growth rate near each singularity.

Proposition 6.11 *Let $x_i \in V_+^{\text{sing}} \cup V_-^{\text{sing}}$. For simplicity, denote r_{x_i} by r . Then if $\gamma \in W_{-2+\delta}^{k,2}(\Lambda^3(\mathbb{S}^1 \times (V_{\pm} \cap \{r \leq r_{0,x_i}\})))$ is an $\mathbb{S}^1$ -invariant 3-form with $(d + d^*)_{\mathbb{S}^1 \times V_{\pm}}\gamma \in W_{-1+\delta}^{k-1,2}$, then there exist constants $c_{j,1}$, c_2 , c_3 , c_4 and an element $\gamma_{>0} \in W_{\delta}^{k,2}$ such that $\gamma = \gamma_{\leq 0} + \gamma_{>0}$, where*

$$\begin{aligned} \gamma_{\leq 0} = & \sum_{\sqrt{\mu_{0,j}+4}-4 \in (-2+\delta, \delta)} c_{j,1} d((d(r\sqrt{\mu_{0,j}+4}-2)\chi_i\phi_{0,j}))^{\#} \lrcorner \varphi \\ & + \chi_i(c_2 \text{Re } \Omega + c_3 \text{Im } \Omega + c_4 d\theta \wedge \omega). \end{aligned}$$

Moreover,

$$\begin{aligned} & \sum_{\sqrt{\mu_{0,j}+4}-4 \in (-2+\delta, \delta)} |c_{j,1}| + |c_2| + |c_3| + |c_4| + \|\gamma_{>0}\|_{W_{\delta}^{k,2}} \\ & \leq C(\|(d + d^*)\gamma\|_{W_{-1+\delta}^{k-1,2}} + \|\gamma\|_{W_{-2+\delta}^{k,2}}), \end{aligned}$$

where the norm is taking on $\Lambda^*(\mathbb{S}^1 \times (V_{\pm} \cap \{r \leq r_{0,x_i}\}))$.

Roughly speaking, $\gamma_{>0}$ is the component with rate larger than 0 in the expansion of γ , and $c_2\chi_i \text{Re } \Omega + c_3\chi_i \text{Im } \Omega + c_4\chi_i d\theta \wedge \omega$ is the component with rate 0 in the expansion of γ . So it is natural to define $\gamma_{\geq 0}$ as

$$\gamma_{\geq 0} = \gamma_{>0} + c_2\chi_i \text{Re } \Omega + c_3\chi_i \text{Im } \Omega + c_4\chi_i d\theta \wedge \omega,$$

and $\gamma_{<0}$ as

$$\gamma_{<0} = \sum_{\sqrt{\mu_{0,j}+4}-4 \in (-2+\delta, \delta)} c_{j,1} d((d(r\sqrt{\mu_{0,j}+4}-2)\chi_i\phi_{0,j}))^{\#} \lrcorner \varphi.$$

Proof Write γ as $\gamma = d\theta \wedge \gamma_2 + \gamma_3$. Then $(d + d^*)_{V_{\pm}}\gamma_p = 0$ for $p = 2, 3$. Consider the Laplacian operator acting on 2-forms or 3-forms on $V_{\pm}$. By Theorem 2.19, $\gamma = \gamma_{>0} + \gamma_{\leq 0}$, where $\gamma_{>0} \in W_{\delta}^{k,2}$ and $\gamma_{\leq 0}$ is a linear combination of $\chi_i \mathcal{P}_i(\lambda_i)$ for $\text{Re } \lambda_i \in (-2+\delta, \delta)$. By Proposition 5.6, any element in such $\mathcal{K}_i(\lambda_i)$ is a homogeneous harmonic form on C with real rate λ_i . Consider the lowest critical rate $\lambda_i \in (-2+\delta, \delta)$.

Using the fact that $(d + d^*)\gamma \in W_{-1+\delta}^{k-1,2}$, the element in $\mathcal{K}_i(\lambda_i)$ corresponding to the $\chi_i \mathcal{P}_i(\lambda_i)$ component of $\gamma_{\leq 0}$ must be closed and coclosed. By Corollary 5.15, it must be a linear combination of $d_C((d_C(r^{\lambda_i+2}\phi_{0,j}))^\# \lrcorner \varphi_C)$ with $\lambda_i \in (-1, 0)$ and $\mu_{0,j} = (\lambda_i + 4)^2 - 4 \in (5, 12)$. Using the fact that $\phi_{0,j}$ is pluriharmonic, it is a harmonic function both on C and $V_\pm$ near x_i . So $d_{V_\pm}((d_{V_\pm}(r^{\lambda_i+2}\phi_{0,j}))^\# \lrcorner \varphi_{V_\pm})$ is also closed and coclosed on $V_\pm$ near x_i . So it can be redefined as an element in $\mathcal{P}_i(\lambda_i)$. Then the problem for the second lowest critical $\lambda_i \in (-2 + \delta, \delta)$ is similar. By induction, the problem is reduced to the $\lambda_i = 0$ case. In this case, $Re \Omega$, $Im \Omega$ and $d\theta \wedge \omega$ are in $\mathcal{K}_i(0)$. Moreover, $\phi_{0,j}$ may not be pluriharmonic. However, by choosing $\delta < \nu$, the difference between $\mathcal{P}_i(0)$ and $\mathcal{K}_i(0)$ can be absorbed into $\gamma_{>0}$. $\square$

Recall that $\Theta(\varphi_T)$ is the Hodge dual of φ_T using the metric defined by φ_T . Define $\tilde{\Theta}(\varphi_T)$ as

$$\tilde{\Theta}(\varphi_T) = \Theta(\varphi_T) - \frac{1}{2} \omega_{T,\pm} \wedge \omega_{T,\pm} - d\theta_\pm \wedge Im \Omega_{T,\pm}.$$

It is easy to see that $d\theta(\varphi_T) = d\tilde{\Theta}(\varphi_T)$ and $\tilde{\Theta}(\varphi_T)$ is a form supported in the regions $t_\pm \in [T - 1, T]$. The $W^{k,2}$ norm of $\tilde{\Theta}(\varphi_T)$ is $O(e^{-\nu T})$.

The following lemma is similar to Proposition 10.3.4 of [39]. The proof is omitted.

Lemma 6.12 *Suppose that $\xi \in W_{1+\delta, \dots, 1+\delta}^{k+2,2}$ is a function supported in the $t_\pm < T - 1$ region. It defines a 3-form $\gamma_{<0}$ by*

$$\gamma_{<0} = d((d\xi)^\# \lrcorner \varphi_\pm).$$

Remark that when $t_\pm < T - 1$, $\varphi_T = \varphi_\pm$. Suppose that $\gamma_{\geq 0} \in W_{-\delta, \dots, -\delta}^{k+2,2}$ is a 3-form on M . Define γ as $\gamma_{<0} + \gamma_{\geq 0}$. Then as long as the norms of ξ and $\gamma_{\geq 0}$ in the corresponding spaces are small enough, δ is small enough and T is large enough, the equation

$$(d + d_{\varphi_T}^*)\gamma + *_{\varphi_T} d \left(\left(1 + \frac{1}{3} \langle \gamma_{\geq 0}, \varphi_T \rangle \right) \tilde{\Theta}(\varphi_T) \right) - *_{\varphi_T} d Q_{\varphi_T}(\gamma_{\geq 0}) = 0$$

implies that

$$d\Theta(\varphi_T + \gamma_{\geq 0}) = 0.$$

Remark that Q_{φ_T} in the equation means the non-linear term of Θ defined in Proposition 2.3.

Using the norm on $W_{\delta, \dots, \delta}^{k,2} \oplus (\oplus_{x_i \in V_\pm^{\text{sing}}} (\mathbb{R}\chi_i Re \Omega \oplus \mathbb{R}\chi_i Im \Omega \oplus \mathbb{R}\chi_i d\theta \wedge \omega))$, it is easy to see that $\|*_{\varphi_T} Q_{\varphi_T}(\gamma)\| \leq \|\gamma\|^2$ if $\|\gamma\|$ is small enough. By implicit function theorem, it is possible to find a solution of the equation

$$(d + d_{\varphi_T}^*)\gamma + *_{\varphi_T} d \left(\left(1 + \frac{1}{3} \langle \gamma_{\geq 0}, \varphi_T \rangle \right) \tilde{\Theta}(\varphi_T) \right) - *_{\varphi_T} d Q(\gamma_{\geq 0}) = 0$$

with $\gamma_{\geq 0} \in W_{\delta, \dots, \delta}^{k,2} \oplus (\oplus_{x_i \in V_{\pm}^{\text{sing}}} (\mathbb{R}\chi_i \text{Re } \Omega \oplus \mathbb{R}\chi_i \text{Im } \Omega \oplus \mathbb{R}\chi_i d\theta \wedge \omega))$. So $\varphi_T + \gamma_{\geq 0}$ provides the required $\mathbb{S}^1$ -invariant torsion-free G₂ structure on $\mathbb{S}^1 \times M$, or equivalently, the Calabi–Yau threefold structure on M .

7 The Obstruction of the Singular Twisted Connected Sum Construction

This section attempts to extend the results in the previous section to the singular twisted connected sum case. Even though Theorem 6.11 has its analogue, it turns out that the analogue of Theorem 6.1 is not true. The first goal of this section is to prove the analogue of Theorem 6.11 in order to get familiar with the edge calculus. Then the edge calculus is used to study the obstruction of the analogue of Theorem 6.1.

Let C be the nodal cone as in Example 2.7. Consider the Laplacian operator Δ acting on p -forms on $\mathbb{S}^1 \times C$. Let θ be the standard variable on $\mathbb{S}^1$. Then any p -form can be expressed as

$$\gamma = d\theta \wedge \alpha + \beta,$$

where α is a $(p-1)$ -form on C and β is a p -form on C . A direct calculation shows that

$$\Delta_{\mathbb{S}^1 \times C} \gamma = d\theta \wedge \left(\Delta_C \alpha - \frac{\partial^2}{\partial \theta^2} \alpha \right) + \left(\Delta_C \beta - \frac{\partial^2}{\partial \theta^2} \beta \right),$$

where Δ_C means Δ operator on each slice.

The Laplacian operator is an “edge operator” studied by Mazzeo in [51]. An edge operator L can be characterized by the normal operator $N(L)$ defined in Definition 2.16 of [51] and the indicial operator $I(L)$ defined in Definition 2.18 of [51]. In (5.2) of [51], Mazzeo defined another operator L_0 as the rescaling of the Fourier transform of $N(L)$. When $L = \Delta$, up to a sign, the corresponding operator Δ_0 on $(p-1)$ -forms α and p -forms β on C is given by

$$\Delta_0(\alpha, \beta) = (\Delta_C + 1)(\alpha, \beta),$$

where roughly speaking, $\frac{\partial}{\partial \theta}$ is replaced by $i = \sqrt{-1}$. Up to a sign, the operator $I(\Delta)$ on forms (α, β) on C is given by

$$I(\Delta)(\alpha, \beta) = (\Delta_C \alpha, \Delta_C \beta),$$

where roughly speaking, $\frac{\partial}{\partial \theta}$ is deleted. δ is called critical if it is critical for $I(\Delta)$ as defined in Definition 2.15.

The solution of $\Delta_0 \gamma = 0$ is related to the Bessel function.

Definition 7.1 The Bessel I-function is defined by

$$I_{\mu}(r) = \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m + \mu + 1)} \left(\frac{r}{2}\right)^{2m+\mu}.$$

The Bessel K-function is defined by

$$K_{\mu}(r) = \frac{\pi}{2} \frac{I_{-\mu}(r) - I_{\mu}(r)}{\sin \mu \pi}$$

if $\mu \notin \mathbb{Z}$. When $\mu \in \mathbb{Z}$, the limit $\lim_{\hat{\mu} \rightarrow \mu} K_{\hat{\mu}}(r)$ exists and is defined as $K_{\mu}(r)$. In either cases, $I_{\mu}(r)$ and $K_{\mu}(r)$ are two independent solutions to the modified Bessel equation

$$\left(\left(r \frac{d}{dr} \right)^2 - (r^2 + \mu^2) \right) y = r^2 \frac{d^2 y}{dr^2} + r \frac{dy}{dr} - (r^2 + \mu^2) y = 0.$$

The following proposition is well known.

Proposition 7.2 (1) When r goes to infinity,

$$I_{\mu}(r) = \frac{1}{\sqrt{2\pi r}} e^r \left(1 + O\left(\frac{1}{r}\right) \right),$$

and

$$K_{\mu}(r) = \sqrt{\frac{\pi}{2r}} e^{-r} \left(1 + O\left(\frac{1}{r}\right) \right).$$

(2) When r goes to 0,

$$\lim_{r \rightarrow 0} r^{-\mu} I_{\mu}(r) = \frac{1}{\Gamma(\mu + 1)} \left(\frac{1}{2} \right)^{\mu}.$$

On the other hand, if $\mu > 0$, then

$$\lim_{r \rightarrow 0} r^{\mu} K_{\mu}(r) = \lim_{\tilde{\mu} \rightarrow \mu} \frac{\pi}{2\Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu} \pi} \left(\frac{1}{2} \right)^{-\mu}.$$

If $\mu = 0$, then

$$\lim_{r \rightarrow 0} K_0(r) (\log r)^{-1} = -1.$$

(3) $K'_0(r) = -K_1(r)$.

Proposition 7.3 Suppose that $\gamma \in W_{\delta}^{k,2}(\Lambda^p(C))$ for $\delta > -2$. If $(\Delta + 1)\gamma = 0$, then $\gamma = 0$.

Proof Write γ as the generalized Fourier series

$$\gamma = \sum_{j=1}^{\infty} \gamma_j(r) \hat{\phi}_{p,j}$$

using the forms $\hat{\phi}_{p,j}$ in Proposition 5.4. The equation $(\Delta + 1)\gamma = 0$ is reduced to the ordinary differential equations

$$\left(\left(r \frac{d}{dr} \right)^2 - (r^2 + \hat{\mu}_{p,j}) \right) \gamma_j = 0.$$

However, any linear combination of $I_{\sqrt{\hat{\mu}_{p,j}}}(r)$ and $K_{\sqrt{\hat{\mu}_{p,j}}}(r)$ does not lie in $W_{\delta+2}^{k,2}$. Thus $\gamma = 0$. $\square$

Using the terminology of [51], Proposition 7.3 implies that $\underline{\delta} \leq -2$ for Δ acting on p -forms. An immediate corollary is the following:

Corollary 7.4 Suppose that $\gamma \in W_{\delta}^{k,2}(\Lambda^*(\mathbb{S}^1 \times C))$ for a non-critical $\delta > -2$, then

$$\|\gamma\|_{W_{\delta}^{k,2}(\Lambda^*(\mathbb{S}^1 \times C))} \leq C \|\Delta_{\mathbb{S}^1 \times C} \gamma\|_{W_{\delta-2}^{k-2,2}(\Lambda^*(\mathbb{S}^1 \times C))}.$$

Proof This corollary is essentially due to [51]. As in the proof of Theorem 5.16 of [51], this estimate is obtained from Fourier transform, rescaling, applying the inverse of $\Delta + 1$ and then the inverse Fourier transform. $\square$

The next proposition is the key estimate for the ordinary differential equation involving Bessel functions.

Proposition 7.5 Assume that $\mu > 0$. Suppose that $y \in W_{\delta}^{k,2}((0, 1))$ and $z \in W_{\delta'}^{k-2,2}((0, 1))$ are functions on the interval $(0, 1)$. They vanish in a neighborhood of 1 and they satisfy the equation

$$\left(\left(r \frac{d}{dr} \right)^2 - (n^2 r^2 + \mu^2) \right) y(r) = z(r).$$

- (1) If $-\mu < \delta < \delta' < \mu$, then $y \in L_{\delta'}^2$.
- (2) If $\mu < \delta < \delta'$, then $y \in L_{\delta'}^2$.
- (3) If $-\mu < \delta < \mu < \delta'$ and $n \neq 0$, define $y_{\leq \mu}$ by

$$y_{\leq \mu}(r) = -I_{\mu}(|n|r) \chi(2|n|r) \lim_{\tilde{\mu} \rightarrow \mu} \frac{\Gamma(\tilde{\mu} + 1) \Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu} \pi}{\tilde{\mu} \pi} \int_0^1 K_{\mu}(|n|s) z(s) \frac{ds}{s}.$$

- Define $y_{>\mu}$ as $y - y_{\leq\mu}$, then $y_{>\mu} \in L^2_{\delta'}$.
 (4) If $-\mu < \delta < \mu < \delta'$ and $n = 0$, define $y_{\leq\mu}$ by

$$y_{\leq\mu}(r) = \chi(2r)r^\mu \int_0^1 \left(\int_0^s t^\mu z(t) \frac{dt}{t} \right) s^{-2\mu} \frac{ds}{s}.$$

Define $y_{>\mu}$ as $y - y_{\leq\mu}$, then $y_{>\mu} \in L^2_{\delta'}$.

Proof When $n \neq 0$,

$$y(r) = \lim_{\tilde{\mu} \rightarrow \mu} \frac{\Gamma(\tilde{\mu} + 1)\Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu}\pi}{\tilde{\mu}\pi} \left(-I_\mu(|n|r) \int_r^1 K_\mu(|n|s)z(s) \frac{ds}{s} \right. \\ \left. - K_\mu(|n|r) \int_0^r I_\mu(|n|s)z(s) \frac{ds}{s} \right) + C_1 I_\mu(|n|r) + C_2 K_\mu(|n|r)$$

using the fact that

$$I'_\mu(r)K_\mu(r) - K'_\mu(r)I_\mu(r) = \lim_{\tilde{\mu} \rightarrow \mu} \frac{\tilde{\mu}\pi}{\Gamma(\tilde{\mu} + 1)\Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu}\pi} \frac{1}{r}.$$

When $n = 0$,

$$y(r) = r^\mu \int_0^r \left(\int_0^s t^\mu z(t) \frac{dt}{t} \right) s^{-2\mu} \frac{ds}{s} + C_1 r^\mu + C_2 r^{-\mu}.$$

□

Assume that δ is small enough. Pick $x_i \in V_+^{\text{sing}}$. Denote r_{x_i} by r . Denote r_{0,x_i} by r_0 . Then the following theorem is an analogue of Proposition 6.11:

Theorem 7.6 (1) Suppose that $-2 < \delta' < \delta'' < -1 + \delta$ satisfy $\delta'' - \delta' < \nu$. If

$$\gamma \in W_{\delta'}^{k,2}(\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\})))$$

and

$$(d + d^*)_{\mathbb{S}^1 \times V_+} \gamma \in W_{\delta''-1}^{k-1,2}(\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\}))),$$

then $\gamma \in W_{\delta''}^{k,2}(\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\})))$ and

$$\|\gamma\|_{W_{\delta''}^{k,2}} \leq C \left(\|(d + d^*)\gamma\|_{W_{\delta''-1}^{k-1,2}} + \|\gamma\|_{W_{\delta'}^{k,2}} \right),$$

where the norm is taking on $\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\}))$.

(2) Suppose that $-1 + \delta \leq \delta' < \delta'' < 0$. If

$$\gamma \in W_{\delta'}^{k,2}(\Lambda^3(\mathbb{S}^1 \times (C \cap \{r \leq r_0\})))$$

and

$$(d + d^*)_{\mathbb{S}^1 \times C} \gamma \in W_{\delta''-1}^{k-1,2}(\Lambda^*(\mathbb{S}^1 \times (C \cap \{r \leq r_0\}))),$$

then there exist $\gamma_{\geq \delta''} \in W_{\delta''}^{k,2}$ and constants $c_{n,j}$ such that $\gamma = \gamma_{< \delta''} + \gamma_{\geq \delta''}$, where

$$\gamma_{< \delta''} = \sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta'')} c_{n,j} d \left(\left(d \left(r^{\sqrt{\mu_{0,j}+4}-2} \chi \left(\frac{4|n + \frac{1}{2}|r}{r_0} \right) \phi_{0,j} e^{in\theta} \right) \right)^{\#} \lrcorner \varphi \right).$$

Moreover,

$$\begin{aligned} & \sqrt{\sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta'')} c_{n,j}^2 |n + \frac{1}{2}|^{8-2\sqrt{\mu_{0,j}+4}-2\delta''} + \|\gamma_{\geq \delta''}\|_{W_{\delta''}^{k,2}}^2} \\ & \leq C(\|(d + d^*)\gamma\|_{W_{\delta''-1}^{k-1,2}} + \|\gamma\|_{W_{\delta'}^{k,2}}), \end{aligned}$$

where the norm is taking on $\Lambda^*(\mathbb{S}^1 \times (C \cap \{r \leq r_0\}))$.

(3) Suppose that $-1 + \delta \leq \delta' < 0 < \delta'' < \frac{\delta}{2}$. If

$$\gamma \in W_{\delta'}^{k,2}(\Lambda^3(\mathbb{S}^1 \times (C \cap \{r \leq r_0\})))$$

and

$$(d + d^*)_{\mathbb{S}^1 \times C} \gamma \in W_{\delta''-1}^{k-1,2}(\Lambda^*(\mathbb{S}^1 \times (C \cap \{r \leq r_0\}))),$$

then there exist a form $\gamma_{>0} \in W_{\delta''}^{k,2}$ and constants $c_{n,j,1}$, $c_{n,2}$, $c_{n,3}$, $c_{n,4}$ such that $\gamma = \gamma_{\leq 0} + \gamma_{>0}$, where

$$\begin{aligned} \gamma_{\leq 0} &= \sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta'')} c_{n,j,1} d \left(\left(d \left(r^{\sqrt{\mu_{0,j}+4}-2} \chi \left(\frac{4|n + \frac{1}{2}|r}{r_0} \right) \phi_{0,j} e^{in\theta} \right) \right)^{\#} \lrcorner \varphi \right) \\ &+ \sum_{n=-\infty}^{\infty} \chi \left(\frac{4|n + \frac{1}{2}|r}{r_0} \right) e^{in\theta} (c_{n,2} \operatorname{Re} \Omega + c_{n,3} \operatorname{Im} \Omega + c_{n,4} d\theta \wedge \omega). \end{aligned}$$

Moreover

$$\begin{aligned} & \sqrt{\sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta'')} c_{n,j,1}^2 |n + \frac{1}{2}|^{8-2\sqrt{\mu_{0,j}+4}-2\delta''}} \\ & + \sqrt{\sum_{n=-\infty}^{\infty} (c_{n,2}^2 + c_{n,3}^2 + c_{n,4}^2) |n + \frac{1}{2}|^{-2\delta''} + \|\gamma_{>0}\|_{W_{\delta''}^{k,2}}} \\ & \leq C(\|(d + d^*)\gamma\|_{W_{\delta''-1}^{k-1,2}} + \|\gamma\|_{W_{\delta'}^{k,2}}), \end{aligned}$$

where the norm is taking on $\Lambda^*(\mathbb{S}^1 \times (C \cap \{r \leq r_0\}))$.

(4) Part (2) and (3) are also true if C is replaced by V_+ .

Proof (1) Using the cut-off function, assume that γ is supported in $r < r_0$ and use the metric g_C to define $d + d^*$ and Δ instead of g_{V_+} . Write γ and $\Delta\gamma$ as

$$\gamma = \sum_{n=-\infty}^{\infty} \sum_{j=1}^{\infty} \gamma_{n,j}(r) \hat{\phi}_{3,j} e^{in\theta},$$

and

$$\Delta\gamma = \sum_{n=-\infty}^{\infty} \sum_{j=1}^{\infty} \gamma'_{n,j}(r) \hat{\phi}_{3,j} e^{in\theta}$$

using the forms $\hat{\phi}_{3,j}$ in Proposition 5.4, then the equation is reduced to the ordinary differential equations

$$\left(r \frac{d}{dr}\right)^2 - (\hat{\mu}_{3,j} + n^2 r^2) \gamma_{n,j} = r^2 \gamma'_{n,j}.$$

Remark that the rate of $\hat{\phi}_{3,j}$ is -2 , so $\gamma_{n,j} \in L_{\delta'+2}^2$ and $r^2 \gamma'_{n,j} \in L_{\delta''+2}^2$. By Proposition 7.5, if $\delta' + 2 < \sqrt{\hat{\mu}_{3,j}} < \delta'' + 2$, then $\gamma_{n,j} = \gamma_{n,j,\leq \delta''+2} + \gamma_{n,j,> \delta''+2}$ with $\gamma_{n,j,> \delta''+2} \in L_{\delta''+2}^2$ and

$$\begin{aligned} \gamma_{n,j,\leq \delta''+2} &= -I_{\sqrt{\hat{\mu}_{3,j}}}(|n|r) \chi\left(\frac{|n|r}{r_0}\right) \lim_{\tilde{\mu} \rightarrow \sqrt{\hat{\mu}_{3,j}}} \frac{\Gamma(\tilde{\mu} + 1) \Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu} \pi}{\tilde{\mu} \pi} \\ &\quad \times \int_0^1 K_{\sqrt{\hat{\mu}_{3,j}}}(|n|s) s^2 \gamma'_{n,j}(s) \frac{ds}{s} \end{aligned}$$

if $n \neq 0$, while

$$\gamma_{n,j,\leq \delta''+2} = \chi\left(\frac{r}{r_0}\right) r \sqrt{\hat{\mu}_{3,j}} \int_0^1 \left(\int_0^s t \sqrt{\hat{\mu}_{3,j}+2} \gamma'_{n,j}(t) \frac{dt}{t} \right) s^{-2\sqrt{\hat{\mu}_{3,j}}} \frac{ds}{s}$$

if $n = 0$. Define

$$c_{n,j} = \lim_{\tilde{\mu} \rightarrow \sqrt{\hat{\mu}_{3,j}}} \frac{\Gamma(\tilde{\mu} + 1)\Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu}\pi}{\tilde{\mu}\pi |n|} \int_0^1 K_{\sqrt{\hat{\mu}_{3,j}}}(|n|s) s^2 \gamma'_{n,j}(s) \frac{ds}{s}$$

for $n \neq 0$, then

$$\begin{aligned} |c_{n,j}| &\leq C_{j,1} \sqrt{\int_0^1 K_{\sqrt{\hat{\mu}_{3,j}}}^2(|n|s) s^{2\delta''+4} \frac{ds}{s}} \sqrt{\int_0^1 (s^2 \gamma'_{n,j}(s))^2 s^{-2\delta''-4} \frac{ds}{s}} \\ &\leq C_{j,2} |n|^{-\delta''-2} \|s^2 \gamma'_{n,j}(s)\|_{L_{\delta''+2}^2}, \end{aligned}$$

where

$$C_{j,1} = \lim_{\tilde{\mu} \rightarrow \sqrt{\hat{\mu}_{3,j}}} \frac{\Gamma(\tilde{\mu} + 1)\Gamma(-\tilde{\mu} + 1) \sin \tilde{\mu}\pi}{\tilde{\mu}\pi}$$

and

$$C_{j,2} = C_{j,1} \sqrt{\int_0^\infty K_{\sqrt{\hat{\mu}_{3,j}}}^2(s) s^{2\delta''+4} \frac{ds}{s}}.$$

Remark that replacing $I_{\sqrt{\hat{\mu}_{3,j}}}(|n|r)$ by its leading term $\frac{1}{\Gamma(\sqrt{\hat{\mu}_{3,j}+1})}(\frac{|n|r}{2})^{\sqrt{\hat{\mu}_{3,j}}}$ does not affect the conclusion. By multiplying $(\frac{|n+\frac{1}{2}|}{|n|})^{\sqrt{\hat{\mu}_{3,j}}}$, it can be replaced by $\frac{1}{\Gamma(\sqrt{\hat{\mu}_{3,j}+1})}(\frac{|n+\frac{1}{2}|r}{2})^{\sqrt{\hat{\mu}_{3,j}}}$ so that the $n = 0$ case can be absorbed into the estimate.

Using the fact that

$$(d + d^*)_{\mathbb{S}^1 \times V_+} \gamma \in W_{\delta''-1}^{k-1,2}(\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\}))),$$

it suffices to consider $\hat{\phi}_{3,j}$ such that $\delta' + 2 < \sqrt{\hat{\mu}_{3,j}} < \delta'' + 2$ and $\hat{\phi}_{3,j}$ is both closed and coclosed. By Corollary 5.15, such form does not exist. Therefore, $\gamma_{n,j} = \gamma_{n,j, > \delta''+2} \in L_{\delta''+2}^2$ if $\delta' + 2 < \sqrt{\hat{\mu}_{3,j}} < \delta'' + 2$.

On the other hand, if $\delta' + 2 < \sqrt{\hat{\mu}_{3,j}} < \delta'' + 2$ is not true, by Proposition 7.5, it is also true that $\gamma_{n,j} \in L_{\delta''+2}^2$.

Apply Corollary 7.4 for each term $\gamma_{n,j}(r) \hat{\phi}_{3,j} e^{in\theta}$. Using the fact that the Laplacian of each term are perpendicular to each other in any weighted L^2 -norm, it is easy to see that

$$\|\gamma\|_{L_{\delta''}^2} \leq C \|\Delta_C \gamma\|_{L_{\delta''-2}^2}.$$

By standard elliptic estimate

$$\|\gamma\|_{W_{\delta''}^{k,2}} \leq C \|\Delta_C \gamma\|_{W_{\delta''-2}^{k-2,2}}.$$

(2) The proof is similar to (1).

(3) The proof is similar to (1).

(4) Choose non-critical δ_l so that $\delta' = \delta_1 < \dots < \delta_{N_5} = \delta''$ and $\delta_l - \delta_{l-1} < \nu$. Assume that $\delta_{N_5-1} < 0$. The statement is proved by induction. When $N_5 = 2$, the statement follows from (2) and (3) because $\delta'' - \delta' < \nu$ in this case. Suppose that the statement has been proved for all $N_5 < N_6$. The goal is to prove the statement for $N_5 = N_6$.

By assumption, $\gamma = \gamma_{<\delta_{N_5-1}} + \gamma_{\geq\delta_{N_5-1}}$ with $\gamma_{<\delta_{N_5-1}}$ defined by

$$\sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta_{N_5-1})} c_{n,j} d \left(\left(d \left(r^{\sqrt{\mu_{0,j}+4}-2} \chi \left(\frac{4|n+\frac{1}{2}|r}{r_0} \right) \phi_{0,j} e^{in\theta} \right) \right)^{\#} \lrcorner \varphi \right)$$

using the G_2 structure on $\mathbb{S}^1 \times V_+$, and

$$\begin{aligned} & \sqrt{\sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta_{N_5-1})} c_{n,j}^2 |n + \frac{1}{2}|^{8-2\sqrt{\mu_{0,j}+4}-2\delta_{N_5-1}} + \|\gamma_{\geq\delta_{N_5-1}}\|_{W_{\delta_{N_5-1}}^{k,2}}^2} \\ & \leq C (\|(d + d^*)_{\mathbb{S}^1 \times V_+} \gamma\|_{W_{\delta_{N_5-1}-1}^{k-1,2}} + \|\gamma\|_{W_{\delta'}^{k,2}}), \end{aligned}$$

where the norm is taking on $\Lambda^*(\mathbb{S}^1 \times (V_+ \cap \{r \leq r_0\}))$.

Consider $\gamma'_{<\delta_{N_5-1}}$ defined by

$$\sum_{n=-\infty}^{\infty} \sum_{\sqrt{\mu_{0,j}+4}-4 \in (\delta', \delta_{N_5-1})} c_{n,j} d \left(\left(d \left(r^{\sqrt{\mu_{0,j}+4}-2} \chi \left(\frac{4|n+\frac{1}{2}|r}{r_0} \right) \phi_{0,j} e^{in\theta} \right) \right)^{\#} \lrcorner \varphi \right)$$

using the G_2 structure on $\mathbb{S}^1 \times C$ instead, then

$$\begin{aligned} & (d + d^*)_{\mathbb{S}^1 \times C} (\gamma - \gamma_{<\delta_{N_5-1}} + \gamma'_{<\delta_{N_5-1}}) \\ & = ((d + d^*)_{\mathbb{S}^1 \times C} - (d + d^*)_{\mathbb{S}^1 \times V_+}) (\gamma - \gamma_{<\delta_{N_5-1}}) \\ & \quad + (d + d^*)_{\mathbb{S}^1 \times V_+} \gamma + ((d + d^*)_{\mathbb{S}^1 \times C} \gamma'_{<\delta_{N_5-1}} - (d + d^*)_{\mathbb{S}^1 \times V_+} \gamma_{<\delta_{N_5-1}}). \end{aligned}$$

By Proposition 2.21, $r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}$ is pluriharmonic and therefore harmonic on both C and V_+ . So

$$\begin{aligned} & (d + d^*)d((d(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta}))^\# \lrcorner \varphi) \\ &= \Delta((d(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta}))^\# \lrcorner \varphi) - dd^*((d(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta}))^\# \lrcorner \varphi) \\ &= ((d\Delta(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta}))^\# \lrcorner \varphi) - d * d((d(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta})) \wedge * \varphi) \\ &= n^2((d(r\sqrt{\mu_{0,j}+4}^{-2}\phi_{0,j}e^{in\theta}))^\# \lrcorner \varphi) \end{aligned}$$

on both $\mathbb{S}^1 \times V_+$ and $\mathbb{S}^1 \times C$. By estimates on each terms, it is easy to see that

$$\begin{aligned} & \|(d + d^*)_{\mathbb{S}^1 \times C}(\gamma - \gamma_{<\delta_{N_5-1}} + \gamma'_{<\delta_{N_5-1}})\|_{W_{\delta''-1}^{k-1,2}} \\ & \leq C(\|(d + d^*)_{\mathbb{S}^1 \times V_+} \gamma\|_{W_{\delta''-1}^{k-1,2}} + \|\gamma\|_{W_{\delta'}^{k,2}}). \end{aligned}$$

The induction statement follows from (2) or (3) applied to $\gamma - \gamma_{<\delta_{N_5-1}} + \gamma'_{<\delta_{N_5-1}}$. $\square$

The analogue of Proposition 6.11 has been proved. The next goal is to obtain the analogue of Proposition 6.1. The key point in the proof of Proposition 6.1 is the fact that there is no $\log r \phi_{2,1}$ term in the $\mathbb{S}^1$ -invariant case. However, in the singular twisted connected sum case, it is easy to see that for $n \neq 0$,

$$(d + d^*)_{\mathbb{S}^1 \times C}(e^{in\theta} d\theta \wedge K_0(|n|r)\phi_{2,1} + \frac{i|n|}{n}e^{in\theta} K_1(|n|r)dr \wedge \phi_{2,1}) = 0.$$

Remark that $K_0(|n|r)$ is asymptotic to $-\log(|n|r)$ when $|n|r$ is small. By comparison to the proof of Propositions 6.1 and 7.6, they provide the infinite dimensional obstruction space for the singular twisted connected sum construction.

The leading obstruction terms are

$$e^{in\theta} d\theta \wedge -\log(|n|r)\phi_{2,1} + \frac{i|n|}{n}e^{in\theta}(|n|r)^{-1}dr \wedge \phi_{2,1}.$$

Remark that the decay rates of $e^{in\theta} d\theta \wedge -\log(|n|r)\phi_{2,1}$ are $r^{-2} \log r$ while the decay rates of $\frac{i|n|}{n}e^{in\theta}(|n|r)^{-1}dr \wedge \phi_{2,1}$ are r^{-3} . So the leading obstruction term is an infinite dimensional linear combination of $e^{in\theta}(|n|r)^{-1}dr \wedge \phi_{2,1}$.

The obstruction comes from inverting the error term in some sense. Using Fourier series with respect to the first $\mathbb{S}^1$ factor, M_+ is $\mathbb{S}^1$ -invariant. On the other hand, even though M_- is $\mathbb{S}^1$ -invariant with respect to the second $\mathbb{S}^1$ factor, the deviation of M_- from being $\mathbb{S}^1$ -invariant with respect to the first $\mathbb{S}^1$ factor decays exponentially. Moreover, the decay rate for the $e^{in\theta}$ factor is $O(e^{-|n|t_-})$. The leading term is the $n = \pm 1$ case. So the leading obstruction term is a linear combination of $\cos \theta r^{-1}dr \wedge \phi_{2,1}$ and $\sin \theta r^{-1}dr \wedge \phi_{2,1}$. By rescaling and changing θ by a constant, the leading obstruction term is $\sin \theta r^{-1}dr \wedge \phi_{2,1}$.

Remark that the Calabi–Yau metric on the deformation C_ϵ of C is asymptotically conical with leading error term $\epsilon r^{-1} dr \wedge \phi_{2,1}$ [19], so roughly speaking, the obstruction will be resolved if each slice $\{\theta\} \times C$ of $\mathbb{S}^1 \times C$ is deformed to $C_{\sin \theta}$. It is an analogue to the construction of Li [46]. There are singularities near $\theta = 0$ and $\theta = \pi$. Topologically, when $\theta = 0$ or $\theta = \pi$, the slice is $C(\mathbb{S}^2 \times \mathbb{S}^3)$. The other slices are $\mathbb{B}^3 \times \mathbb{S}^3$. The total space is $C(\mathbb{S}^3 \times \mathbb{S}^3)$ near $\theta = 0$ or $\theta = \pi$. In other words, there are strong evidences that the nodal singularity along $\mathbb{S}^1$ should be replaced by two isolated conical singularities with model $C(\mathbb{S}^3 \times \mathbb{S}^3)$. Such problem will be left for future studies.

Remark that the main tool of Li's construction [46] is Yau's solution [64] of the Calabi conjecture and its non-compact generalizations by Tian-Yau [62] and Hein [35]. In the G_2 case, the analogue of Yau's theorem is not available, so one has to instead find extra structures to reduce the dimension.

Acknowledgements The author would like to thank Edward Witten for introducing him to this problem and for many fruitful discussions. The author is also grateful to the referees for detailed suggestions as well as Jeff Cheeger, Xiuxiong Chen, Sir Simon Donaldson, Lorenzo Foscolo, Mark Haskins, Hans-Joachim Hein, Helmut Hofer, Fanghua Lin, Rafe Mazzeo, Johannes Nordström, Song Sun, Akshay Venkatesh, Jeff Viaclovsky and Ruobing Zhang for helpful conversations. This material is based upon work supported by the National Science Foundation under Grant No. 1638352, as well as support from the S. S. Chern Foundation for Mathematics Research Fund.

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