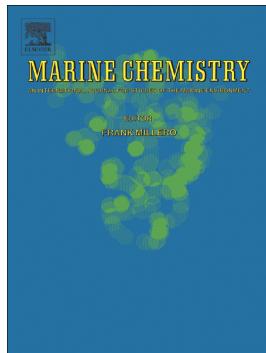


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# ACTINIUM AND RADIUM FLUXES FROM THE SEABED IN THE NORTHEAST PACIFIC BASIN

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## Keywords

Marine sediments, <sup>227</sup>Ac, <sup>228</sup>Ra, <sup>226</sup>Ra, benthic fluxes, distribution coefficient, core incubation, GEOTRACES, and radionuclides.

## Highlights

First study to measure <sup>227</sup>Ac sediment profiles along a cruise tract in Northeast Pacific. The distribution coefficient ( $k_d$ ) for Ac in these sediments is about 7 times higher than that for Ra. These are the first measurements of  $k_d$  for Ac.

Benthic fluxes of <sup>227</sup>Ac and <sup>228</sup>Ra determined by core incubation shows results comparable to predictions from transport models applied to sediment profiles.

Average benthic fluxes of  $^{227}\text{Ac}$ ,  $^{228}\text{Ra}$ , and  $^{226}\text{Ra}$  in the Northeast Pacific are  $90\pm20$ ,  $600\pm200$ , and  $1300\pm200 \text{ dpm m}^{-2}\text{-yr}^{-1}$ , constraining interpretations of water column data.

## Abstract

Five sediment cores were collected along a cruise tract from Hawaii to Alaska in August 2017 (C-Disk-IV cruise) with the objective of characterizing the behavior of  $^{227}\text{Ac}$ ,  $^{228}\text{Ra}$ , and  $^{226}\text{Ra}$  and their fluxes into the overlying water column, information that is essential to the interpretation of the distribution of these tracers in the ocean, for example, as measured on GEOTRACES cruises. Solid phase profiles of these isotopes were measured, and reaction-transport models were applied that incorporated molecular diffusion, bioturbation, sedimentation, distribution coefficients ( $k_d$ ), and the fraction of each isotope released to pore water by parent decay (called F). Fits to these profiles used  $k_d$  values determined in lab experiments for C-Disk-IV sediments. Ra  $k_d$  values ( $1000\text{-}3000 \text{ mL g}^{-1}$ ) agreed with previous estimates for deep-sea sediments, and Ac  $k_d$  values ( $3500\text{-}22000 \text{ mL g}^{-1}$ ) correlated with those for Ra but were about 7 times greater.

Two independent approaches were used to quantify the benthic fluxes of  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  in the Northeast Pacific: (1) use of solid phase profiles with a reaction-transport model, as well as integrated downcore daughter-parent deficiency; and (2) direct measurement of fluxes based on core incubation. The two independent methods agreed within uncertainty, and the average  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  sediment fluxes for the Northeast Pacific are  $90\pm20$  and  $600\pm200 \text{ dpm m}^{-2}\text{-yr}^{-1}$ , respectively. The  $^{226}\text{Ra}$  sediment flux was only determined by the former approach, and the flux

calculated in this study is similar to previous work in the North Pacific, averaging  $1300 \pm 200$  dpm m $^{-2}$ ·yr $^{-1}$ . This is over 2x higher than the water column inventory of  $^{226}\text{Ra}$  in this region (600 dpm m $^{-2}$ ·yr $^{-1}$ ), and indicates the importance of lateral  $^{226}\text{Ra}$  export from the N. Pacific.

The largest  $^{227}\text{Ac}$  and Ra isotope fluxes in the study area are near the center of the Northeast Pacific ( $\sim 37^\circ\text{N}$ ). Smaller  $^{227}\text{Ac}$ ,  $^{228}\text{Ra}$  and  $^{226}\text{Ra}$  fluxes occur north of  $40^\circ\text{N}$ , primarily due to dilution of their Pa and Th ancestors by higher sediment accumulation rates.

## Introduction

$^{227}\text{Ac}$  ( $t_{1/2}=21.77$  y) has shown potential as a tracer for mixing on the basin scale range (Nozaki, 1984; Nozaki et al., 1990; Geibert and Vöge, 2008; Geibert et al., 2002). Nozaki first proposed  $^{227}\text{Ac}$  as a deep-sea tracer in 1984, but for years,  $^{227}\text{Ac}$  went relatively unmeasured due to its low concentrations and difficulties with detection in ocean sediments and waters (Geibert and Vöge, 2008). However, during the last 30 years, better methods have evolved that can measure small concentrations of this isotope, due to development of the Radium Delayed Coincidence Counters (RaDeCC; Moore and Arneborg, 1996; Shaw and Moore, 2002; Dulaiova et al., 2013), and more sensitive mass spectrometers (Levier et al., 2021).

GEOTRACES is an international collaboration whose main objective is to understand the processes that control the distribution and transport of trace elements and isotopes (TEIs) in the oceans (GEOTRACES Planning Group, 2007). One important component for understanding these processes is quantifying transport using naturally occurring radioisotopes, including  $^{228}\text{Ra}$  and  $^{227}\text{Ac}$ . The last 4 US GEOTRACES section cruises (GA03, GP16, GN01, GP15) have

collected Ra and Ac from seawater by pumping large volumes through acrylic fibers impregnated with MnO<sub>2</sub>. This has led to acquisition of a number of data sets for the radium quartet (<sup>223</sup>Ra, <sup>224</sup>Ra, <sup>228</sup>Ra, and <sup>226</sup>Ra) and <sup>227</sup>Ac distribution throughout the oceans. High resolution sampling has occurred in the North Atlantic, North and South Pacific, and the Arctic oceans (Charette et al., 2015, Kipp et al., 2018; Sanial et al., 2018; Hammond et al., in prep; Kemnitz et al., in prep). Applying these isotopes as tracers can be facilitated if their source functions are defined. While <sup>226</sup>Ra and <sup>228</sup>Ra in deep ocean sediments have been well studied at several locations to constrain their source functions, limited data exists for <sup>227</sup>Ac (Cochran and Krishnaswami, 1980; Kadko, 1980; Nozaki et al., 1990; Moore et al., 1996; Rama and Moore, 1996; Kipp et al., 2018; Rutgers van der Loeff et al., 2018). This limits its utility (Ku and Luo, 2009). A few studies have assumed mass balance models (Geibert et al., 2008), but Nozaki et al. (1990) is the only study to date to have directly measured <sup>227</sup>Ac in sediments. Thus, this study aims to improve knowledge of the geochemical behavior of <sup>227</sup>Ac, <sup>226</sup>Ra, and <sup>228</sup>Ra in marine sediments and estimate their fluxes into deep ocean waters.

One objective of this study is to compare methodologies to measure <sup>227</sup>Ac fluxes from marine sediments by using two independent methods: indirectly by modeling solid phase profile activity measurements and directly by core incubations. The former approach utilizes previously published reactive-transport models to predict geochemical behaviors of <sup>227</sup>Ac and other radionuclides (<sup>226</sup>Ra & <sup>228</sup>Ra) in the sediments sampled during this study, while the second approach has not previously been utilized in the deep sea. A second objective is to evaluate spatial variability in sediment fluxes for <sup>227</sup>Ac, <sup>228</sup>Ra, and <sup>226</sup>Ra in a section through the

Northeast Pacific Basin (NEPB) and show how the water column distribution of these radionuclides is related to these fluxes.

## Theory

Cochran and Krishnaswami, (1980) developed a reactive transport model for the behavior of  $^{226}\text{Ra}$  and  $^{228}\text{Ra}$  in marine sediments. This model can be adapted for  $^{227}\text{Ac}$  as well (Nozaki et al., 1990). The model assumes parents of these isotopes are associated with solids, while one fraction of the more mobile daughters may diffuse as solutes or reversibly adsorb to solid phases, and a second fraction of these daughters is tightly bound within solid phases and does not exchange with the fluid phase. This model is idealized, but accounts for the major processes that govern reaction and transport in sediments, including molecular diffusion, radioactive decay, burial, and bioturbation (treated as a diffusive process). Sediments are divided into an upper, bioturbated layer that overlies a deeper layer that is no longer bioturbated. The governing one-dimensional steady state equations are given by:

$$\frac{\partial C}{\partial t} = (D_s + K D_b) \frac{\partial^2 C}{\partial z^2} - S(1 + K) \frac{\partial C}{\partial z} - \lambda(1 + K)C + P = 0 \quad \text{for } 0 \leq z \leq L \quad (1a)$$

$$\frac{\partial C}{\partial t} = D_s \frac{\partial^2 C}{\partial z^2} - S(1 + K) \frac{\partial C}{\partial z} - \lambda(1 + K)C + P = 0 \quad \text{for } z \geq L \quad (1b)$$

where:

$C$  = concentration of dissolved radioisotope in porewater (atoms  $\text{cm}^{-3}$ )

$D_s$  = molecular diffusion corrected for tortuosity ( $\text{cm}^2 \text{yr}^{-1}$ )

$D_b$  = bioturbation ( $\text{cm}^2 \text{yr}^{-1}$ ); assumed constant in the upper layer

$S$  = sedimentation rate ( $\text{cm yr}^{-1}$ )

$\lambda$  = decay constant of daughter ( $\text{yr}^{-1}$ )

$P$  = production rate of daughter atoms per unit of porewater from parent decay (atoms  $\text{cm}^{-3} \text{yr}^{-1}$ )

$K$  = partition coefficient (dimensionless)

$z$  = depth in sediments (cm)

$L$  = thickness of bioturbation zone (cm)

A solution to eq. 1 assumes sedimentation rate, molecular diffusion, and  $K$  are all constant and independent with respect to depth; bioturbation is constant in the upper layer (Cochran and Krishnaswami, 1980; Kadko, 1980). Solid phase  $^{231}\text{Pa}$  and  $^{230}\text{Th}$  profiles in the bioturbated layer (shown later) have been homogenized throughout the upper 6 cm. The production term (P) can be further expanded and expressed as:

$$P = FA_{p1} + fA_{p2} \quad \text{for } 0 \leq z \leq L \quad (2a)$$

$$P = FA_{p1}e^{-\mu(z-L)} + fA_{p2} \quad \text{for } z \geq L \quad (2b)$$

where:

$F$  = fraction of parent isotope ( $^{231}\text{Pa}$  or  $^{230}\text{Th}$ ) decays which recoil daughter products into porewater

$f$  = fraction of grandparent isotope ( $^{235}\text{U}$  or  $^{234}\text{U}$ ) decays which recoil daughter products into porewater

$A_{p1}$  = parent activity in soil multiplied by bulk density and divided by porosity to give equivalent activity per unit of porewater in atoms  $\text{cm}^{-3} \text{ yr}^{-1}$

$A_{p2}$  = grandparent activity defined similarly to  $A_{p1}$ .

$$\mu = \frac{\lambda_p}{S} (\text{cm}^{-1})$$

$\lambda_p$  = decay constant of parent

Note: The grandparent-parent pair for  $A_{p1}$  and  $A_{p2}$  is assumed to be in equilibrium and f is typically minor, due to  $^{234}\text{U}$  and  $^{235}\text{U}$  being within a crystal lattice and having limited ability to recoil daughter products into porewater. The parent and grandparent activities per unit of porewater are calculated from solid phase activity multiplied by solid phase grams per  $\text{cm}^{-3}$  of porewater, based on porosity. Substituting eq. 2 into eq. 1, solving for the concentration, multiplying by  $\lambda_d$ , and expressing concentrations as activities ( $\lambda C = A$ ):

$$A_1(z) = Q_1 e^{\alpha_1 z} + R_1 e^{\delta_1 z} + \frac{f A_{p1} + f A_{p2}}{1+K} \quad \text{for } 0 \leq z \leq L \quad (3a)$$

$$A_2(z) = Q_2 e^{\alpha_2 z} + \frac{f A_{p1} \lambda}{D_1} e^{-\mu(z-L)} + \frac{f A_{p2}}{1+K} \quad \text{for } z \geq L \quad (3b)$$

where Q, R,  $\alpha$ , and  $\delta$  are all constants that depend on  $D_s$ ,  $D_b$ , S, and K (See SI S.1 for full definition of these terms).  $A_{p1}$  is the parent activity of for a specific daughter in  $\text{dpm cm}^{-3}$  of porewater,  $A_{p2}$  is its grandparent activity in  $\text{dpm cm}^{-3}$  of porewater,  $A_1$  is the porewater activity of this daughter in the bioturbated zone in  $\text{dpm cm}^{-3}$ ,  $A_2$  is the porewater activity of this daughter below the bioturbated zone in  $\text{dpm cm}^{-3}$ ,  $A_0$  is the concentration of daughter in porewater at the SWI (sediment-water interface) in  $\text{dpm cm}^{-3}$ , and z is depth in cm. Eq. 3 depends on 4 boundary conditions:

$$1.) \text{ at } z = 0, A_1 = A_0$$

$$2.) \text{ as } z \rightarrow \infty, A_2 \rightarrow \frac{f A_{p2}}{1+K}$$

$$3.) (A_1)_L = (A_2)_L$$

$$4.) D_s \left( \frac{\partial A_1}{\partial z} \right)_L = (D_s + K D_b) \left( \frac{\partial A_2}{\partial z} \right)_L$$

To relate the dissolved concentration to the solid phase, the partition coefficient (K) can be related to the distribution coefficient,  $k_d$  (in units of  $\text{cm}^3 \text{ g}^{-1}$ ) as:

$$K = \frac{k_d(1-\phi)\rho_s}{\phi} \quad (4)$$

where  $\rho_s$  = solid phase density ( $\text{g cm}^{-3}$ ).  $k_d$  relates the mobile solid and dissolved phase concentration  $C'$  ( $\text{dpm g}^{-1}$ ) as:

$$k_d = \frac{C'}{c} \quad (5)$$

The following equations relate the dissolved phase ( $\text{dpm cm}^{-3}$ ) to the solid phase ( $\text{dpm g}^{-1}$ ) by multiplying eq. 3 by  $k_d$  to find the mobile activity (dissolved plus adsorbed) and adding the immobile crystalline activity to obtain the total activity:

$$A_{1-tot}(z) = k_d(A_1) + (1-F)(A_{p1-tot} - A_{p2-tot}) + (1-f)A_{p2-tot} \quad \text{for } 0 \leq z \leq L \quad (6a)$$

$$A_{2-tot}(z) = k_d(A_2) + (1-F)(A_{p1-tot} - A_{p2-tot}) + (1-f)A_{p2-tot} \quad \text{for } z \geq L \quad (6b)$$

where  $A_{1-tot}$  and  $A_{2-tot}$  are the total activities of daughter products in the solid phase for each layer ( $\text{dpm g}^{-1}$ ),  $A_{p1-tot}$  is the total parent activity for a specific daughter ( $^{231}\text{Pa}$  and  $^{230}\text{Th}$ ) in  $\text{dpm g}^{-1}$ ,  $A_{p2-tot}$  is the total grandparent activity for this daughter ( $^{234}\text{U}$  and  $^{235}\text{U}$ ) in  $\text{dpm g}^{-1}$ ,  $(1-F)(A_{p1-tot} - A_{p2-tot})$  is the immobile fraction of  $^{226}\text{Ra}$  or  $^{227}\text{Ac}$  from excess  $^{230}\text{Th}$  or  $^{231}\text{Pa}$  decay ( $\text{dpm g}^{-1}$ ), and  $(1-f)A_{p2-tot}$  is the immobile fraction of  $^{226}\text{Ra}$  or  $^{227}\text{Ac}$  from  $^{234}\text{U}/^{235}\text{U}$  decay (this decay is from within the crystal lattice of sediments).

Lastly, a one-layer model can be derived from eq. 1a if the parent activity and bioturbation are assumed constant throughout the entire profile. Due to its small scale-length, a one-layer model was used for  $^{228}\text{Ra}$  in this study. Two boundary conditions were applied: at  $z=0$ , the concentration of daughter products =  $A_0$  and as  $z \rightarrow \infty$ , the concentration of daughter products  $\rightarrow$

$P/(1+K)$ , where  $P$  is parent activity. Solving for the concentrations with respect to the previous boundary conditions and expressing the concentrations as activities ( $A=\lambda C$ ):

$$A_D(z) = A_0 e^{-\alpha z} + \frac{P}{1+K} (1 - e^{-\alpha z}) \quad (7)$$

$$\alpha = \frac{-S(1+K) + \sqrt{(-S(1+K))^2 + 4(D_s + KD_b)(\lambda_{Ac}(1+K))}}{2(D_s + KD_b)} \quad (8)$$

$A_D$  = concentration of dissolved daughter in porewater (dpm cm<sup>-3</sup>)

$A_0$  = concentration of dissolved daughters in porewater at  $z = 0$  (dpm cm<sup>-3</sup>)

$P$  = activity of parent isotope (<sup>232</sup>Th) per unit volume of pore water (dpm cm<sup>-3</sup>)

For a one-layer model, the following equation relates the dissolved phase (dpm cm<sup>-3</sup>) to the solid phase (dpm g<sup>-1</sup>) by multiplying eq. 7 by  $k_d$  and adding the total parent decay:

$$A'_{Total D}(z) = A_0 k_d e^{-\alpha z} + F A_{p1-tot} (1 - e^{-\alpha z}) + (1 - F) A_{p1-tot} \quad (9)$$

$A'_{Total D}$  = the total activity of daughters in solid phase (dpm g<sup>-1</sup>).

Fluxes of <sup>227</sup>Ac, <sup>226</sup>Ra, or <sup>228</sup>Ra can be calculated by applying Fick's first law to the fitted equations:

$$Flux = -(D_s + D_b K) \phi \frac{dA}{dz} \quad (10)$$

where  $\frac{dA}{dz}$  is the derivative of eq. 6a or 9.

## Study Area

Figure 1 shows the C-Disk-IV (Carbonate Dissolution Kinetics-IV) and US GEOTRACES PMT (GP15) transects in the NEPB. All deep-sea sediment samples were collected during the C-Disk-IV Cruise aboard the R/V *Kilo Moana* in August 2017. Profiles of dissolved  $^{227}\text{Ac}$ ,  $^{231}\text{Pa}$ ,  $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$  in the water column were collected during GP15 cruise aboard R/V *Roger Revelle* between September and November 2018.

Location, depth, surface porosity, sedimentation, and description of sediments for each C-Disk-IV station are summarized in Table 1. Stations 1, 2, and 3 are located within the subtropical gyre, where oligotrophic waters dominate, and organic export to the seafloor is low. The sediments at these stations are characterized by fine-grained, clay material with low organic content. Stations 4 and 5 are in the subarctic gyre, where primary production is high. Their sediments are characterized by larger grain size material with higher biogenic fraction in comparison to stations 1-3. A distinctive fluff layer was observed above the sediments at station 4, indicative of high surface production above this station (Lou et al., 2019).

## Material and Methods

**Sample Collection:** Sediment cores were collected using a multi-coring device (10 cm ID) and sectioned every 1 cm between 0-10 cm and 2 cm between 10-30 cm. Sectioned mud was placed in a cold van after sectioning and later analyzed for porosity at the USC lab by weighing aliquots before and after drying at 50°C. Samples were then dried, gently crushed, and placed in plastic bags until analysis. Weights for sediment samples were corrected for salt contribution, based on water content and salinity.

In the water column, dissolved  $^{227}\text{Ac}$ ,  $^{228}\text{Th}$ , and the radium quartet were collected using a dual-flow path in-situ pump (McLane WRT-LV). Each pump had two filters in parallel: the higher volume flow path had a 1  $\mu\text{m}$  quartz filter (Whatman QMA) and the lower volume flow path had a 0.8  $\mu\text{m}$  polyethersulfone filter (Pall Supor800). This allowed sufficient flow to pass about 1.5  $\text{m}^3$  during a 4-hour pump. Downstream of the filter heads, flow then passed through two grooved acrylic cartridges impregnated with  $\text{MnO}_2$  that sat in series. After collection, cartridges were rinsed with DIW for several minutes to remove sea-salt and then dried to 50-120% moisture using compressed air. Cartridges were re-measured multiple times from a few months to 24 months later to determine  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  based on the release of  $^{219}\text{Rn}$  and  $^{220}\text{Rn}$  progeny.

**Standards and Tracers:** A  $^{229}\text{Th}/^{225}\text{Ac}$  tracer (USC #8)B:  $9.32 \pm 0.12 \text{ dpm g}^{-1}$  was used for measurements of  $^{227}\text{Ac}$  in sediments, calibrated against an  $^{227}\text{Ac}$  standard from Eckert and Zeigler diluted to an activity of 5.90  $\text{dpm g}^{-1}$  on 4/27/2017 (See Kemnitz, 2022).

Standards for gamma counting ( $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$ ) were obtained from the Environmental Protection Agency (SRM-1 diluted pitchblende and SRM-2 diluted monzonite). All Th and Ra daughter products for the  $^{235}\text{U}$  and  $^{232}\text{Th}$  standards are assumed to be in secular equilibrium, with the exception of  $^{210}\text{Pb}$ . Standards were counted in 3.0 cm high geometry, and corrections were made to account for the different sample heights that were used. For low energy gamma regions (i.e., 46 keV:  $^{210}\text{Pb}$ ), a matrix correction of 10% was routinely applied to solid phase samples when liquid solutions were used as standards based on spiking solid samples with high concentrations of known  $^{210}\text{Pb}$ .

$^{210}\text{Pb}$  was standardized using a NIST certified solution (SRM 4337). A few  $^{210}\text{Pb}$  check samples were also run by alpha spectrometry using its  $^{210}\text{Po}$  daughter and a  $^{209}\text{Po}$  standard (USC #65D), calibrated against the NIST  $^{210}\text{Pb}$  solution. The  $^{209}\text{Po}$  standard was corrected for decay to the time of analysis using the recently determined half-life of 125 yr (Colle et al., 2014).

***Distribution Coefficients for Ra and Ac:*** Ac and Ra distribution coefficients ( $k_d$ ) were determined experimentally by taking Ra-free seawater, adding deep-sea sediments, and spiking it with an aged ( $>45$  y)  $^{232}\text{Th}$  (USC #106A: 492 dpm  $\text{g}^{-1}$ ) or  $^{229}\text{Th}$  (USC #89E: 1500 dpm  $\text{g}^{-1}$ ) solution. All daughter products in both spikes are assumed to be in secular equilibrium with their parents.  $^{228}\text{Ra}$  ( $t_{1/2} = 5.75$  y) and  $^{225}\text{Ac}$  ( $t_{1/2} = 10.0$  d) were used to determine Ra and Ac  $k_d$  values.

The procedure for determining Ra  $k_d$  values was as follows: 30 mL of Ra-free seawater was weighed and placed into 50 mL centrifuge tube (some experiments required 80 mL of seawater and used 125 mL PET bottles). Then 1500 dpm of  $^{232}\text{Th}/^{228}\text{Ra}$  was added to the same centrifuge tube, along with enough  $\text{Na}_2\text{CO}_3$  to neutralize acid and bring pH between 7.5-7.8. Next,  $\sim 14$  grams of wet deep-sea sediments were added to tubes and shaken for at least 1 hour, allowing radioisotopes to equilibrate between dissolved and solid phases. After equilibration, the tubes were centrifuged for 5 minutes at 2500 rpm and the seawater was decanted, filtered ( $0.45\ \mu\text{m}$ ), and a weighed aliquot passed through  $\text{MnO}_2$  fibers ( $\sim 1\text{g}$ ) 3 times. Reid et al. (1979) have shown that the absorption efficiency for Ac and Ra onto  $\text{MnO}_2$  fibers are 100% in a single pass at flow rates below  $1\ \text{L min}^{-1}$ . The fibers and wet sediments were then placed in separate polypropylene tubes and measured on HPGe detectors, using gamma rays from  $^{228}\text{Ac}$  daughters (338 and 911

keV) to determine the  $^{228}\text{Ra}$  activity. The final solid phase mass was determined by drying the wet sediments and weighing the final mass (corrected for salt content).

The determination of Ac  $k_d$  values was similar but required additional steps: First, a pure  $^{225}\text{Ac}$  solution was needed. A  $^{229}\text{Th}$  spike (matrix of 3 N  $\text{HNO}_3$ ) was weighed out and diluted with 15 mL of 3 N  $\text{HNO}_3$ .  $^{225}\text{Ac}$  was separated from  $^{229}\text{Th}$  (and  $^{225}\text{Ra}$ ) by using a 2 mL, DGA chromatography cartridge resin from Eichrom Inc. (Martin, 1995). This solution was neutralized, mixed with sediments, and equilibrated as described for Ra, except that Savillex beakers were used instead of a centrifuge tube. The pH was kept between 7.3-7.6, to minimize  $^{225}\text{Ac}$  sorption to the walls of the Savillex beakers. However, about 20% of the  $^{225}\text{Ac}$  did stick to walls of the beaker after one hour of rotating the sediment slurry. This was confirmed by washing Savillex beakers with acid and measuring by HPGC. This  $^{225}\text{Ac}$  loss did not impact our calculations because both sediment and dissolved phases were measured to determine  $k_d$  values.

**Alpha Spectroscopy Procedure:** The procedure developed by Dulaiova et al. (2013) for  $^{227}\text{Ac}$  in rocks was used, with the exception of adding Fe carrier. For these sediments, sufficient Fe is present, so this addition was not needed. Excellent separation and good yields (average of 50%) were observed for Ac in deep-sea sediments, suggesting this method works well for aluminosilicate clay material. Details of this procedure are in SI S.3.

**ICP-MS Procedure:** Uranium, thorium, and protactinium isotopes in sediments were measured by ICP-MS (Element 2, XR type) at Lamont Doherty Earth Observatory using the procedure of Fleisher and Anderson (1991) and Anderson et al. (2012). %Mn was measured at USC by

following the dissolution protocol used for alpha spectroscopy described above, diluting samples in 2% HNO<sub>3</sub> containing In tracer (1 ppb) and analyzing Mn with an Element 2 ICP-MS at mass 55 amu.

**Gamma Spectroscopy Procedures:** Sediment samples (0.5 – 2.0 g) were dried, gently crushed, and placed in polyethylene or polypropylene tubes (tubes were 5 cm in height and 0.5 cm in diameter). Tubes were then placed in a high purity intrinsic germanium well-type detector (HPGe ORTEC, 120 cc active volume), and counted for 2-8 days to determine <sup>210</sup>Pb (46 keV) and <sup>226</sup>Ra (186 keV). <sup>235</sup>U (185.7 keV) interference at 186 keV was removed by using the <sup>238</sup>U activity determined from <sup>234</sup>Th (63 keV) and the natural <sup>235</sup>U/<sup>238</sup>U ratio. For samples near the sediment-water interface, a fraction of the <sup>222</sup>Rn produced in situ should be lost by diffusion into the overlying water column. The scale length (1/e) of the Rn loss was assumed to be 3 cm and 50% of the Rn was assumed to be lost, based on measurements of Rn profiles in equatorial sediments (Hammond et al., 1996). Excess <sup>210</sup>Pb (<sup>210</sup>Pb<sub>ex</sub>) was determined by subtracting the measured <sup>226</sup>Ra (corrected for <sup>222</sup>Rn loss in the upper 5 cm) from total <sup>210</sup>Pb activity and correcting for decay of <sup>210</sup>Pb<sub>ex</sub> between collection and analysis. At depths below the bioturbated layer, <sup>210</sup>Pb and <sup>226</sup>Ra were in good agreement. A more detailed explanation of this procedure is available in SI S.4.

**Core Incubation:** Core incubations were carried out at each station using the methodology from Hammond et al (2004). Briefly, after retrieval of replicate cores from the multi-core device, the spring-loaded arms were released, and rubber stoppers were placed on the bottom of the cores.

Next, a moveable piston was inserted into each core top, which could be advanced as water was withdrawn. Overlying height was adjusted to 12 cm and decreased as samples were taken to measure nutrient fluxes. The incubation plug had a stirring device that continuously stirred the overlying water for the entire incubation period. Results for nutrient fluxes from these cores have been published in Hou et al. (2019).

After the incubation was completed, the remaining water from each core was siphoned out, and all water from the two replicate cores at each site was combined for measurement of  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  (~1.0 L). This water was passed through loose  $\text{MnO}_2$  coated fibers (20 g), at least 3 times at  $< 1.00 \text{ L/min}$  to insure complete absorption of radium, thorium and actinium (Reid et al., 1979).  $\text{MnO}_2$  fibers were then stored at room temp and shipped back to the USC lab. Upon arrival at USC, fibers were washed several times with DIW to remove sea-salt and then dried to 80-120% moisture to be analyzed for its  $^{223}\text{Ra}$  and  $^{224}\text{Ra}$  activity on RaDeCC (Moore and Arnold, 1996; Sun and Torgersen, 1996). Measurements began 7 to 20 days after incubation ended (longer delays for earlier stations). Subsequent measurements were made up to 2 years after collection, allowing  $^{223}\text{Ra}$ ,  $^{224}\text{Ra}$ ,  $^{228}\text{Ra}$ , and  $^{227}\text{Ac}$  activities at the end of the incubation period to be calculated using the Bateman equation fit to the data (Bateman, 1910; See SI S.5 for details). For the long-lived isotopes  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  the approach of Hammond et al. (2004) was used to calculate benthic flux  $J$ :

$$J = \Delta A \times \sum \frac{h}{t} \quad (11)$$

where  $\Delta A$  is the increase in activity ( $\text{dpm m}^{-3}$ ) of  $^{227}\text{Ac}$  or  $^{228}\text{Ra}$  during the incubation (in practice, initial concentrations were negligible),  $t$  is time of incubation period, and  $h$  is height of the overlaying water above the sediment-water-interface (SWI) in the core. Because aliquots of

water were periodically removed to use for other analyses during the incubation,  $h/t$  must be summed for each incubation interval. For short-lived isotopes, the computation is more complex because these isotopes decay during the incubation period. A parameter  $a_i$  was computed ( $m^{-1}$ ) to calculate flux, with  $a_i$  based on water height and time for each incubation step, including radioactive decay (see SI).

## Results and Discussion

Estimates of benthic isotope fluxes have been made based on reaction-transport modeling of profiles of solid phase measurements, integrating profiles of daughter-parent deficiencies, and directly measuring isotope fluxes from incubated cores. Applying some of these approaches required measurements of distribution coefficients and evaluation of bioturbation and sediment accumulation rates. These approaches are detailed in sections below.

**Distribution Coefficients ( $k_d$ ):** One important parameter needed to describe the behavior of radioisotopes in sediments is the partition coefficient (K), which is defined as the ratio of adsorbed solute to dissolved solute and is related to the distribution coefficient  $k_d$  (eq. 4). While some estimates have been made for radium  $k_d$ , (Cochran and Krishnaswami, 1980; Rama and Moore, 1996; Crotwell and Moore, 2003; Colbert and Hammond, 2008; Beck and Cochran, 2013), very little data is available for  $^{227}\text{Ac}$ . The only study to provide K values for  $^{227}\text{Ac}$  was Nozaki et al. (1990), and those values were estimated from Ra K values from an earlier study (Cochran and Krishnaswami, 1980).

Tables 2 and 3 summarize Ra and Ac  $k_d$  values determined experimentally in this study. The range of Ra  $k_d$  values for C-Disk-IV sediments are between 1000-4500  $\text{mL g}^{-1}$ . These values fall within the range of previous Ra  $k_d$  estimates for deep-sea sediments (Beck and Cochran, 2013). For Ac,  $k_d$  values are much larger, but the patterns for different stations are similar. Station 1 shows constant Ra and Ac  $k_d$  values in the upper 3 cm of sediments (1500 & 15,100  $\text{mL g}^{-1}$ ), while other stations tend toward higher values at depth. Station 2 has the highest Ra and Ac  $k_d$  values (3260 & 21,600  $\text{mL g}^{-1}$ ), while results at other stations are about 2x smaller. There is a trend for decreasing  $k_d$  values with larger grain sizes that are visually apparent at stations 4 and 5. Another variable that was tested for Ra  $k_d$  values was the mass-to-water ratio. Deep-sea sediments have an *in-situ* mass-to-water ratio of 0.28  $\text{g mL}^{-1}$  ( $\rho_s=2.50 \text{ g mL}^{-1}$ ,  $\phi=0.90$ ,  $\rho=0.25$ ) that is much larger than experiments. However, there was no trend witnessed with changing mass-to-water ratios on Ra  $k_d$  values (Table 2). Ac  $k_d$  experiments did not vary mass to water ratios, but this variable might be checked in future endeavors.

Stations 1 and 2 are characterized by more clay-like material, and its sediment composition has higher %Mn compared to the northern stations. Ra  $k_d$  values seem to be influenced by %Mn in these sediments, but not as noticeably as Ac (Fig. 2). Previous authors have observed that Ra and other actinides are strongly influenced by adsorption onto Mn-oxides in marine sediments (Kadko, 1980; Kadko et al., 1987; Hayes et al., 2015). %Mn differed at different C-Disk-IV stations, but at each station it was nearly constant in the upper 7.5 cm, and both Ac and Ra  $k_d$  values at each station were also fairly uniform.

Fig. 3 shows Ac vs. Ra  $k_d$  values for these experiments. It indicates that Ac  $k_d$  values are 6.6 times higher than Ra  $k_d$  values. While measurements are limited, this relationship could prove helpful in future modeling of  $^{227}\text{Ac}$  in deep-sea sediments. Previous work has assumed that the Ac  $k_d$  value is close to that of Ra (Nozaki et al., 1990). Additional experiments (data not shown) suggest that good pH control is required, as  $k_d$  for Ac increases at higher pH.

**Profiles of radioisotopes in solid phases:**  $^{231}\text{Pa}$  in these sediments is largely derived from the overlying water column, where it is produced by decay of  $^{235}\text{U}$  and scavenged by sinking particulates (Bacon and Anderson, 1982; Anderson et al., 1983). Based on  $^{238}\text{U}$ , less than 0.1 dpm  $\text{g}^{-1}$  of  $^{231}\text{Pa}$  should be supported by its  $^{235}\text{U}$  parent in sediments (See SI Table S.2 for  $^{238}\text{U}$  activity). Figure 4 displays  $^{231}\text{Pa}$  concentrations in surface sediments vs. latitude in the Pacific Ocean near  $150^\circ\text{W}$ . Both published (Lao et al., 1992) and this study's values are shown. There is a clear trend of increasing  $^{231}\text{Pa}$  concentration moving south from the Aleutians, likely reflecting less dilution of  $^{231}\text{Pa}$  raining from the overlying water column by lower fluxes of detrital sediment (note accumulation rates in Table 1). Data from two regions fall off this trend: one is located near the equatorial Pacific, where a high sedimentation rate due to preservation of  $\text{CaCO}_3$  should dilute the  $^{231}\text{Pa}$  signal, and a second data point is near the Hawaii margin, where water depth is shallower and the  $^{231}\text{Pa}$  rain should be smaller. The profiles of  $^{231}\text{Pa}$  and  $^{230}\text{Th}$  (Fig. 5, 6) suggest bioturbation mixes Pa and Th fairly well throughout the upper 5-8 cm, most likely reflecting the long half-lives of each isotope relative to the mixing time scales.

$^{210}\text{Pb}_{\text{ex}}$  profiles (Fig. 5) indicate varying bioturbation rates ( $D_b$ ) throughout the study area. By fitting an exponential function to the  $^{210}\text{Pb}_{\text{ex}}$  profiles in the top 4 cm (8 cm at sta. 4), assuming constant biodiffusivity ( $D_b$ ) and radioactive decay,  $D_b$  can be calculated at each site (Table 4;

$D_b = \lambda / \mu^2$ ,  $\mu$ =attenuation factor in  $\text{cm}^{-1}$ ). Stations 1, 2 and 3 show low  $D_b$  rates ( $<0.01 \text{ cm}^2 \text{ yr}^{-1}$ ) and stations 4 and 5 show high  $D_b$  rates ( $>0.01 \text{ cm}^2 \text{ yr}^{-1}$ ), likely reflecting the latitude dependence of organic matter rain to the benthic community (Tromp et al., 1995). However, most  $^{210}\text{Pb}_{\text{ex}}$  profiles have an erratic pattern: exponential decay in the upper 3 cm with a slight excess between 4-6 cm at some stations. The excess in the 4-6 cm horizons is evidence of non-local transport (bioturbation that is localized at certain horizons in the sediment profile), which is known to occur in deep-sea sediments (Teal et al., 2008; Smith et al., 1997; Boudreau, 1986). Station 4 has the highest  $^{210}\text{Pb}_{\text{ex}}$  values at depth ( $\sim 5.5 \text{ cm}$ ). Station 2 has a relatively low inventory of  $^{210}\text{Pb}$  in comparison to adjacent stations (Fig. 5), and its possible significance will be discussed later.

$^{227}\text{Ac}$  is less strongly bound to solid phases than  $\text{Ti}$ ,  $\text{Ta}$ , and  $\text{Pb}$ , and can be transported within upper sediments by molecular diffusion and bioturbation, as it is produced by decay of its  $^{231}\text{Pa}$  parent (Nozaki et al., 1990). Consequently, profiles are expected to increase with depth toward secular equilibrium with the parent isotope, a pattern that is generally apparent. The scale length is comparable to the scale length for  $^{210}\text{Pb}_{\text{ex}}$  at each station, which would be expected if bioturbation dominates transport, given the similar half-lives of these isotopes. There is little evidence of an  $^{227}\text{Ac}$  deficiency at depths near 4-6 cm, in association with the non-local mixing of  $^{210}\text{Pb}_{\text{ex}}$ , but downward transport of a few percent of surficial material with low  $^{227}\text{Ac}$  should have less influence on  $^{227}\text{Ac}$  than on  $^{210}\text{Pb}$ .

While the scale lengths for  $^{227}\text{Ac}$  do not vary enormously among the stations, there are two types of  $^{227}\text{Ac}$  profiles found in the North Pacific: One (Sta. 1 and 3) shows a large  $^{227}\text{Ac}$  deficiency relative to  $^{231}\text{Pa}$  in the upper few cm of sediments, and the other (Sta. 2 and 4) shows a small deficiency in the upper few cm of sediments (Fig. 5). At all 5 stations,  $^{227}\text{Ac}$  and  $^{231}\text{Pa}$  activities

in deeper samples show equilibrium (>5cm), which is expected and gives confidence in the  $^{227}\text{Ac}$  analysis procedure used here. Station 5 shows very little  $^{227}\text{Ac}$  deficiency relative to its parent. The cause of differences among these profiles may be the result of differences in the fraction of  $^{227}\text{Ac}$  that is mobile via molecular diffusion. However, at station 5, simple models do not account for many of the observations, and we suspect that non-steady state processes may have influenced distributions at this site. These issues are discussed further in later sections of the paper.

The  $^{226}\text{Ra}$  profiles in the North Pacific are as expected: Most profiles of  $^{226}\text{Ra}$  increase downcore to depths of ~8 cm (Fig. 6), reflecting transport by diffusion and bioturbation, as well as ingrowth toward equilibrium with its  $^{230}\text{Th}$  parent. The longer scale length for  $^{226}\text{Ra}$  reflects its longer half-life (1600 y) and lower  $k_d$ , as compared to  $^{227}\text{Ac}$  and  $^{210}\text{Pb}$ . Below 8 cm,  $^{226}\text{Ra}$  shows more uniform activity until 17-20 cm, where it should approach equilibrium with its parent and decrease in activity as does its parent. Results at stations 2 and 3 establish this behavior, and it is inferred at stations 1 and 4. The  $^{226}\text{Ra}$  profile at station 5 was somewhat erratic, with equilibrium observed between  $^{226}\text{Ra}$  and  $^{230}\text{Th}$  throughout the core, similar to the  $^{227}\text{Ac}/^{231}\text{Pa}$  profile. The absence of a deficiency at this station suggests erosion of overlying sediments may have occurred in the relatively recent past (decades to centuries). This possibility is not supported by pore water profiles of silicic acid at this site (Hou et al., 2019), which do not indicate a disturbance; however, their response times to approach steady state may be only months to a few years.

Despite relatively large uncertainties,  $^{228}\text{Ra}$  profiles show a discernible increase with depth toward equilibrium with its parent  $^{232}\text{Th}$  (Fig. 6). Measurement of  $^{228}\text{Ra}$  was based on its  $^{224}\text{Ra}$

progeny detected with gamma spectrometry, which is far less precise at these low concentrations than the ICP analysis of the  $^{232}\text{Th}$  parent. Most profiles have a small  $^{228}\text{Ra}$  deficiency in the upper 2-3 cm of sediments. The much smaller scale length, in comparison to  $^{226}\text{Ra}$ , is due to its much shorter half-life ( $t_{1/2}=5.75$  y). While the uncertainties are large,  $^{228}\text{Ra}$  appears to be in equilibrium with its parent at depth. The activities of  $^{232}\text{Th}$  for these profiles are constant down to 7.5 cm, except station 5, which increases in activity by 7.5 cm.

**$^{230}\text{Th}$  and Accumulation Rates:** S in the upper layer can also be estimated from a mass balance for  $^{230}\text{Th}$  in the mixed layer using logic presented by Bacon (1984) and reviewed by Francois et al. (2004) and Costa et al. (2020). At steady state, the rain of  $^{230}\text{Th}$  to sediments, divided by its production from  $^{234}\text{U}$  in the overlying water column, can be used to define a sediment focusing factor  $g$ :

$$g = \frac{MA_r}{I} \quad (12)$$

where  $M$  = mass accumulation rate ( $\text{g cm}^{-2} \text{ y}^{-1}$ )

$A_r$  =  $^{230}\text{Th}$  in sediment raining to the sea floor and being preserved ( $\text{dpm g}^{-1}$ )

$I$  = production from  $^{234}\text{U}$  in overlying water ( $\text{dpm cm}^{-2} \text{ yr}^{-1}$ )

$g$  = (rain of  $^{230}\text{Th}$  to sediments)/(integrated water column  $^{230}\text{Th}$  production)

The value of  $g$  is near 1.0 in many parts of the ocean, reflecting the relatively short residence time of  $^{230}\text{Th}$  in the water column. However, preferential boundary scavenging of  $^{230}\text{Th}$  in high particle flux areas, sediment winnowing, or lateral inflow of water low in  $^{230}\text{Th}$  may reduce  $g$  below 1.0. Alternatively, lateral inflow of  $^{230}\text{Th}$ -rich water or sediment focusing at a coring site may increase  $g$  above 1.0. The rain of  $^{230}\text{Th}$ , averaged over its lifetime, can also be estimated by

integrating its downcore activity, calculated by summing its decay rate in the bioturbated mixed layer and the integration of its exponential profile based on fits below this layer in Fig. 6. Results for this approach, using cores in this study (Table 5), indicate that  $g$  is often less than 1. Values indicate that either notable sediment winnowing or lateral thorium transport in the water column has been occurring at stations 2, 4, and 5. However, some localized thorium removal is occurring at station 1. It is interesting to note that these stations with low values of  $g$  are also those with much smaller values for  $^{227}\text{Ac}$  F than expected.

An additional factor of importance for these sediments is the relatively thick bioturbated layer that exists, with  $^{230}\text{Th}$  homogenized to depths of 6-7 cm. Rather than using  $^{14}\text{C}$  to correct for the effective age of this layer (Francois et al., 2004), the  $^{230}\text{Th}$  balance can be modified to include its presence and determine the residence time of  $^{230}\text{Th}$  in this layer. The calculation assumes the excess  $^{230}\text{Th}$  leaving the base of the mixed layer equals its input at the top (using eq. 12), minus decay in the mixed layer:

$$MA_m = gI - \lambda\rho L(1 - \phi_{UL})A_m \quad (13a)$$

$$M = \frac{gI}{A_m} - \lambda\rho L(1 - \phi_{UL}) \quad (13b)$$

where  $M$  = mass input to the upper mixed layer ( $\text{g cm}^{-2} \text{ yr}^{-1}$ )

$A_m$  = activity of excess  $^{230}\text{Th}$  in mixed layer ( $\text{dpm g}^{-1}$ )

$\lambda$  = decay constant of  $^{230}\text{Th}$  ( $\text{yr}^{-1}$ )

$\phi_{UL}$  = Upper layer average porosity

The term on the far right of eq. 13a corrects for radioactive decay of  $^{230}\text{Th}$  during passage of sediment through the bioturbated layer, and the equation can be re-arranged to find the mass accumulation rate in this layer using eq. 13b. Below this layer, the LL (Lower Layer) MAR was

derived by multiplying S determined in the LL by  $\rho(1-\phi_{LL})$ , where  $\phi_{LL}$  = Lower layer porosity.

Table 5 summarizes the results for mass accumulation rates (MAR) in the upper and lower layers using the above equations along with  $^{230}\text{Th}$  inventories, sedimentation rates (S), and mixed layer age based on mixed layer mass divided by M. Most stations have similar MAR in the upper and lower layers, suggesting that the thorium dynamics have not changed significantly over the time defined by sampling (about 50 ky). The mixed layer ages here are about 10 to 60 ky.

**Fitting reaction-transport models to solid phase profiles:**  $^{227}\text{Ac}$ ,  $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$  profiles from C-Disk-IV stations 1-4 were fitted with the reaction-transport equations (eq. 6, 9) by using a least squares approach programmed into MATLAB. The model parameters for „best fit“ solid phase  $^{227}\text{Ac}$ ,  $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$  profiles are summarized in Table 4. The mixing depth (L) was between 6-8 cm for stations 1-4 and was chosen based on profiles of the long-lived isotopes  $^{230}\text{Th}$  and  $^{226}\text{Ra}$  (Fig. 6). Based on averages for the mixed layer,  $^{231}\text{Pa}$  and  $^{230}\text{Th}$  activities were considered constant in this layer and decline exponentially thereafter, with a scale length equal to the ratio of sedimentation rate to radioactive decay as defined in equation 2b. A fit to the  $^{226}\text{Ra}$  profile below the mixed layer constrained linear sedimentation rates in this zone.  $^{232}\text{Th}$  activity was assumed constant throughout the each  $^{228}\text{Ra}$  profile, and a one-layer model was used for these profiles as noted above (eq. 9).

Sedimentation rate (S) and the fraction released by parent decay (F) were allowed to be free parameters and were varied in the reaction-transport model for  $^{226}\text{Ra}$  until the „best fit“ was achieved. The model for  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  profiles used S from  $^{226}\text{Ra}$  profiles and only varied F,

the fraction of  $^{227}\text{Ac}$ ,  $^{226}\text{Ra}$ , or  $^{228}\text{Ra}$  that is released into pore water. Krishnaswami and Cochran (1980) predicted this value should be around 50%, since half of all alpha decays from an adsorbed parent should either release daughters into the surrounding pore waters or inject it into the solid on which it is sorbed. The fraction of Ac and Ra ejected into pore waters can either remain in solution or be reversibly adsorbed onto sediment surfaces, but it should remain mobile. The  $^{226}\text{Ra}$  profiles had F values ranging between 50-80%, while  $^{228}\text{Ra}$  had F values between 25-70% (Table 4). The  $^{228}\text{Ra}$  F values are surprisingly high, assuming most of its  $^{232}\text{Th}$  parent is probably contained in crystalline lattices. The  $^{227}\text{Ac}$  F values had a much broader range. Stations 2 & 4 had F values below 20% while the station 3 F value was above 90%. Only station 1 had an F value around 50%. Both  $^{227}\text{Ac}$  and  $^{226}\text{Ra}$  should be primarily produced from their adsorbed parents, and the much lower F values for  $^{227}\text{Ac}$  at stations 2 and 4 are unexpected.

The analytical solution for eq. 1 to find the isotope profiles described by eq. 3a and 3b required the assumption that porosity was constant in each layer. However, changes in porosity with depth do occur; these affect the diffusivity of solutes ( $D_s$ ) and the ratio of sorbed to dissolved solute (K). To evaluate this effect on  $^{227}\text{Ac}$  behavior, the reaction-transport equation (eq. 1) was applied numerically in a MATLAB script. The most rapid porosity change occurs in the uppermost portion of sediments (~0.01m), where the  $^{227}\text{Ac}$  deficiency is observed (See SI S.6 for full details of model). F values and  $^{227}\text{Ac}$  fluxes calculated for the numerical model were slightly (~20%) lower than the reaction-transport model, smaller than the uncertainties, suggesting that changes in porosity in the upper-most centimeters of sediments do not have a major effect on the reaction-transport model fits that assumed porosity is constant.

**Core Incubation:** The  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  concentrations in overlying water collected during core incubations were determined by using RaDeCC to monitor the activity of  $^{223}\text{Ra}$  and  $^{224}\text{Ra}$  in these samples over a 3-year period. As mentioned earlier,  $^{223}\text{Ra}$  is the granddaughter of  $^{227}\text{Ac}$ , with a half-life of 11 days and is much more soluble than its ancestors  $^{227}\text{Th}$  and  $^{227}\text{Ac}$ , producing a higher flux during the incubation. Consequently,  $^{223}\text{Ra}$  on these  $\text{MnO}_2$  fibers was initially in excess over its ancestors but grew into equilibrium with its grandparent activity by 120 days.  $^{224}\text{Ra}$  and  $^{228}\text{Ra}$  show a similar pattern, but over a much longer time scale. Figure 7 shows Bateman equation fits to the  $^{223}\text{Ra}$  and  $^{224}\text{Ra}$  activity for a typical station.

Table 6 summarizes the results for stations 1-4. Results for station 5 indicated unreasonably high fluxes ( $> 500 \text{ dpm m}^{-2} \text{ yr}^{-1}$ ). In part, this can be attributed to a short incubation time that would result in large uncertainty, but it might reflect a problem with filtration or perhaps a non-steady state distribution of its  $^{227}\text{Ac}$  parent progeny in the sediments, as discussed earlier. If an erosional event removed surficial sediments during a time period of perhaps 1 to 20 years prior to core collection, surficial sediments with high  $^{227}\text{Ac}/^{231}\text{Pa}$  would be present, supplying a high, but transient, diffusive flux of  $^{227}\text{Ac}$ .

**Summary of  $^{227}\text{Ac}$  Inputs throughout the NEPB:** Figure 8 summarizes  $^{227}\text{Ac}$  fluxes determined by all the methods discussed in this paper, including calculation of the observed  $^{227}\text{Ac}$  deficiency (See SI S.5 for full details). Briefly, the observed deficiency was calculated by taking the difference of parent-daughter activity downcore and summing the deficiency; this method along with core incubation approach should be the most accurate methodology for determining  $^{227}\text{Ac}$  fluxes, since all variables are known in these two methods. Given the uncertainties, these

approaches agree at most C-Disk-IV stations, as do the fluxes based on the numerical and reaction-transport (analytical) models. Although the incubation-based fluxes are often slightly less precise, they typically agree within about 35% of the mean for the first two approaches. Furthermore, the analytical model results are comparable to the numerical approach. This suggests that a simple model that accounts for major processes governing a radionuclide in marine sediments can describe solid phase  $^{227}\text{Ac}$  profiles fairly well.

To test the accuracy of these  $^{227}\text{Ac}$  benthic inputs, they can be compared to water column inventories of excess  $^{227}\text{Ac}$ , computed along the GP15 transect (See SI Table S-3). The  $^{227}\text{Ac}$  sediment fluxes are an average of the results from the core incubation and observed deficiency methods. Water column fluxes were averaged from GP15 station observations. The computation and interpretation of the water column data is the subject of a companion paper (Kemnitz et al., in prep). However, it is clear that the results are comparable within the uncertainties (~20%). Sediment fluxes for  $^{227}\text{Ac}$  show a local maximum at station 3 and a pronounced drop toward the north (sta. 5). Changes in the water column inventory are less apparent, but clearly are lower at the northern end of the transect. Lateral mixing in the water column is likely to mute spatially varying inputs from sediment, and these patterns will be discussed in the companion paper.

Bioturbation ( $D_b$ ), Fraction released (F), and  $^{231}\text{Pa}$  should be the major controls on  $^{227}\text{Ac}$  fluxes in the NEPB (Nozaki et al., 1990). There is a positive correlation between  $^{227}\text{Ac}$  fluxes and F, while there is little correlation between  $^{227}\text{Ac}$  fluxes and  $k_d$  or  $^{231}\text{Pa}$  (Fig. 9). The large variation in F values in the NEPB masks the correlation of  $^{227}\text{Ac}$  fluxes with the other variables. Despite the low value for F at Station 4, its high flux reflects its large bioturbation rate, which is 50 times greater than any other station in the NEPB (Guinasso and Schink, 1975; Kadko and Heath, 1984;

Nozaki et al., 1990). Perhaps non-local burial of young surface material that should be deficient in  $^{227}\text{Ac}$  has played some role in exhuming deeper material that is less deficient in  $^{227}\text{Ac}$ . At station 2, as noted previously, the inventory of  $^{210}\text{Pb}$  is relatively low near the surface. This might reflect a recent erosion event of the top  $\sim 1$  cm at that site, perhaps a decade ago, so that pore water profiles (Hou et al., 2019) have re-adjusted, but  $^{210}\text{Pb}$  and  $^{227}\text{Ac}$  have not returned to steady state.

***Summary of  $^{226}\text{Ra}$  Inputs throughout the NEPB:*** Similar to  $^{227}\text{Ac}$  fluxes in this region,  $^{226}\text{Ra}$  fluxes are lower north of  $40^\circ\text{N}$  and then increase below this latitude (Fig. 10). The fluxes for  $^{226}\text{Ra}$  range from 200 - 1700 dpm/m $^2$ -yr throughout C-Disk-IV stations. The highest  $^{226}\text{Ra}$  fluxes were observed between  $25^\circ\text{N}$ - $40^\circ\text{N}$  (Sta.1-3). The  $^{226}\text{Ra}$  fluxes in this region are consistent, averaging around 1350 dpm m $^{-2}$  yr $^{-1}$  (observed deficiency). North of  $40^\circ\text{N}$ ,  $^{226}\text{Ra}$  fluxes decrease toward the Alaskan margins ( $\sim 50 - 200$  dpm m $^{-2}$  yr $^{-1}$ ). Topography steps up a few hundred meters of elevation around  $40^\circ\text{N}$ , which is the boundary of the Mendocino Fracture Zone (MFZ) (Hautala, 2018). North of this boundary,  $^{226}\text{Ra}$  and  $^{227}\text{Ac}$  fluxes are much lower, likely due to the shoaling of bottom depths and increase in sedimentation, factors which should both decrease the  $^{231}\text{Pa}$  and  $^{230}\text{Th}$  concentrations in sediments.

These estimates are in good agreement with the  $^{226}\text{Ra}$  sediment flux data from Cochran and Krisnaswami (1980), which reported  $^{226}\text{Ra}$  fluxes between 800-2000 dpm m $^{-2}$  yr $^{-1}$  for the North Equatorial Pacific. They noted that the local  $^{226}\text{Ra}$  in the overlying water column required a benthic input of only 600 dpm m $^{-2}$  yr $^{-1}$ , indicating that North Equatorial Pacific was a net source of  $^{226}\text{Ra}$  to the world's oceans. For our study area, the  $^{226}\text{Ra}$  sediment fluxes required to explain the average standing crop of  $^{226}\text{Ra}$  in the water column above the C-Disk-IV sites (See SI Table

S-3) are a factor of  $\sim 2$  lower than the measured sediment fluxes. Thus, the NEPB is comparable to the North Equatorial Pacific region in that  $^{226}\text{Ra}$  has a large sediment flux that results in a net source of  $^{226}\text{Ra}$  to other parts of the world's oceans.

***Summary of  $^{228}\text{Ra}$  Inputs throughout the NEPB:*** Lastly,  $^{228}\text{Ra}$  sediment fluxes were measured by core incubation and curve fitting solid phase  $^{228}\text{Ra}$  profiles using eq. 9. At most stations, results are comparable, although the core incubation method is more precise. At station 2, the model is twice the core incubation estimate (Fig. 8). However, the large uncertainty in the model calculation is not inconsistent with the more precise core incubation. The highest  $^{228}\text{Ra}$  flux is located around  $35^\circ\text{N}$  and decreases south and north of this latitude (Fig. 10). Comparing the  $^{228}\text{Ra}$  flux to both  $^{226}\text{Ra}$  and  $^{227}\text{Ac}$  in the NEPB, the highest fluxes occur south of the MFZ ( $\sim 40^\circ\text{N}$ ) and then decrease toward the Alaskan margin. In the center of the NEPB, above oligotrophic waters, fluxes remain high and constant towards the Hawaiian margin. Comparing  $^{228}\text{Ra}$  water column fluxes to the sediment flux, most C-Disk-IV stations agree with each other, except station 3 ( $35^\circ\text{N}$ ) (See SI). At station 3, the water column estimate is 60% higher than the sediment-based approach for the  $^{228}\text{Ra}$  flux. This suggests that at some locations significant horizontal transport may be occurring. Water column distributions of  $^{228}\text{Ra}$  will be discussed in a subsequent article.

### **Summary and Conclusions:**

$^{227}\text{Ac}$ ,  $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$  profiles in sediments were measured and modeled at 5 stations along the C-Disk-IV transect in the Northeast Pacific. The procedure of Dulaiova et al. (2013) was applied to these deep-sea sediments to measure  $^{227}\text{Ac}$  and showed accurate and precise results.

Distribution coefficients ( $k_d$ ) were measured for Ac and Ra in C-Disk-IV surface sediments.

These Ra  $k_d$  values were comparable to previously published work, with values between 1000-3000 mL g<sup>-1</sup>, which agree with Ra  $k_d$  values for similar sediments (Beck and Cochran, 2013).

The Ac  $k_d$  values were between 3500-22000 mL g<sup>-1</sup> for the same sediments. There was a strong positive correlation between Ra and Ac  $k_d$  values, with Ac being almost 6.6 times higher than Ra  $k_d$  values, and both co-varying with the MnO<sub>2</sub> in solid phases. This is the first study to directly measure Ac  $k_d$  values in marine sediments.

Two independent approaches were carried out to quantify the source function of <sup>227</sup>Ac in the Northeast Pacific: Core incubation and solid phase measurements of <sup>227</sup>Ac by alpha spectroscopy. The latter approach used analytical and numerical reaction-transport models to describe the behavior of <sup>227</sup>Ac in sediments. Both models showed comparable results for <sup>227</sup>Ac fluxes when compared to core incubation results. <sup>227</sup>Ac profiles showed a large range in the fraction of mobile <sup>227</sup>Ac released from <sup>231</sup>Pa (called F) in the upper few cm of sediments, versus the expected F value of about 50%. The cause of this wide range is undetermined. At one site (station 2) a relatively low value of excess <sup>210</sup>Pb inventory suggests an erosion event could have occurred in recent decades. At another site that underlies higher productivity waters (station 4), we speculate that non-steady state exhumation of <sup>227</sup>Ac-rich deeper sediment by bioturbation might play a role, but there is little other apparent evidence this is occurring. It is notable that about half of the expected <sup>230</sup>Th is missing from stations 2, 4 and 5, suggesting winnowing might be taking place at these sites, although it may also reflect lateral transport of <sup>230</sup>Th in the water column before it can be scavenged to sediments. Interestingly, these are also sites where the

near-surface  $^{227}\text{Ac}$  deficiency is much smaller than expected, possibly indicating exhumation of older sediments.

Most  $^{226}\text{Ra}$  profiles for C-Disk-IV sediments have a large deficiency relative to  $^{230}\text{Th}$ . At station 5, both  $^{227}\text{Ac}/^{231}\text{Pa}$  and  $^{226}\text{Ra}/^{230}\text{Th}$  profiles show near equilibrium throughout the upper sediments. This is evidence that sediments at this site have been disturbed during the past few decades, although the excess  $^{210}\text{Pb}$  profile shows an exponential decrease through the upper 5 cm. At other C-Disk-IV sites,  $^{226}\text{Ra}$  F values range between 57–85%, in agreement with previous results from deep Pacific sites. Benthic input from the NEP is a net source of  $^{226}\text{Ra}$  to the global ocean, averaging around 1400 dpm  $\text{m}^{-2}\text{-yr}^{-1}$  vs. the 100 dpm  $\text{m}^{-2}\text{-yr}^{-1}$  needed to supply the  $^{226}\text{Ra}$  in the overlying water column inventory, demonstrating the importance of lateral transport of  $^{226}\text{Ra}$  exported from this region.

$^{228}\text{Ra}$  profiles along C-Disk-IV transect also had a modest deficiency relative to its parent,  $^{232}\text{Th}$ . However,  $^{228}\text{Ra}$  profiles were difficult to measure due to high counting uncertainties. F values obtained from  $^{228}\text{Ra}$  profiles were similar to  $^{226}\text{Ra}$  values. Core incubation results provided more precise estimates of benthic input for this isotope.

The largest  $^{227}\text{Ac}$  fluxes observed are located near the center of the Northeast Pacific ( $\sim 35^\circ\text{N}$ ) decreasing somewhat to the south. Smaller fluxes occur north of  $40^\circ\text{N}$  and decrease towards the Alaskan margins.  $^{227}\text{Ac}$  fluxes are influenced by the  $^{231}\text{Pa}$  concentration in sediments, sediment composition, and mass accumulation rate. The largest  $^{226}\text{Ra}$  and  $^{228}\text{Ra}$  fluxes are also located near the center of the Northeast Pacific, and  $^{226}\text{Ra}$  fluxes remain high between Hawaii and  $35^\circ\text{N}$ .

$^{228}\text{Ra}$  fluxes decrease south of 35°N. Both  $^{226}\text{Ra}$  and  $^{228}\text{Ra}$  fluxes decrease north of this latitude, approaching the Alaskan margin.

Finally, we note that a surprisingly large number of the 5 sites we have sampled do not perfectly match the idealized radioisotope distributions we anticipated. Inventories of  $^{230}\text{Th}$  are about half the expected values at stations 2, 4 and 5. Station 5 isotope profiles appear to be non-steady state, with evidence of erosion in recent times.  $^{227}\text{Ac}$  deficiencies in the top 1 cm are often smaller than a simple model predicts, suggesting that older material has been exhausted. This suggests that the NEPB sediments may experience far more disturbance than expected.

### **Declaration of Competing Interest**

Authors have no competing interests at this time.

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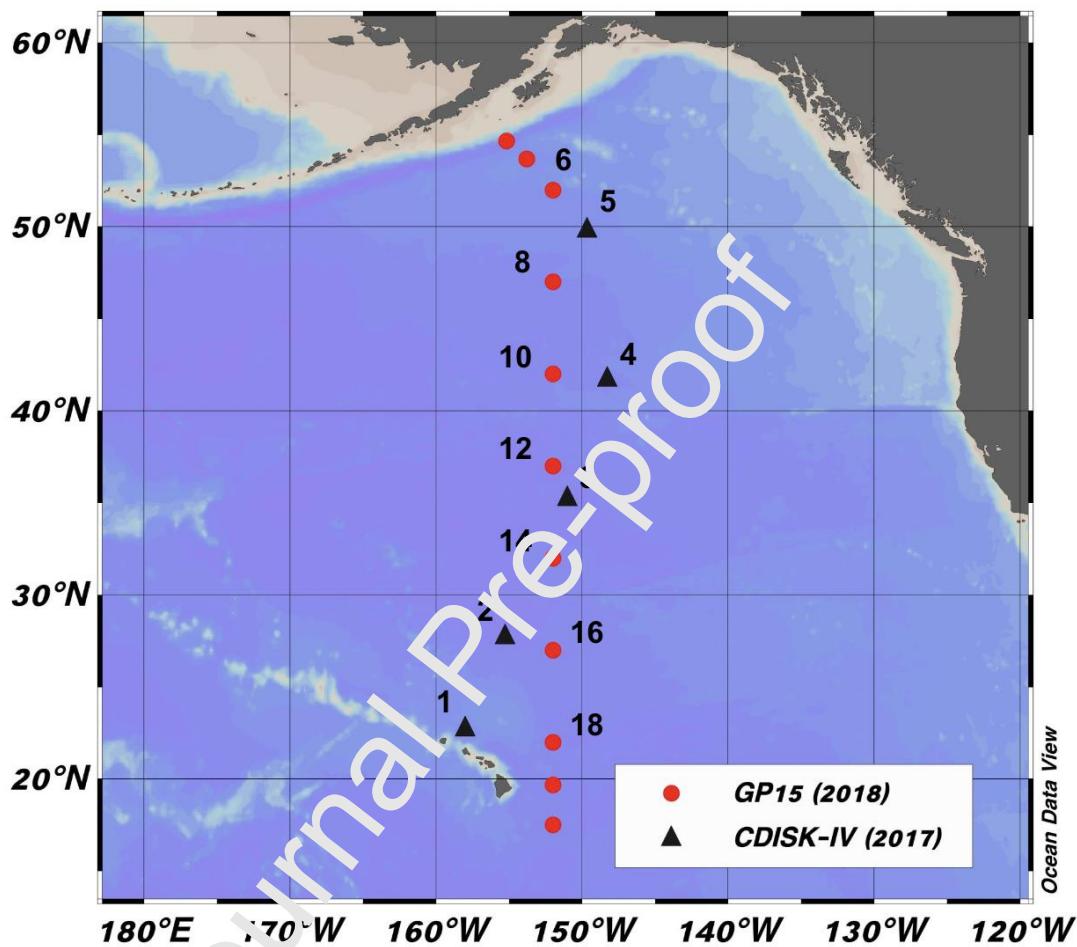
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## Figures and Tables



**Figure 1:** Study Area of the Northeast Pacific Basin (NEPB). Red dots and black triangles refer to stations where  $^{227}\text{Ac}$  was measured in the water column for GEOTRACES PMT (GP15) and C-Disk-IV cruises. Numbers next to symbols refer to station numbers. All sediment isotope profiles (Ra, Th, Ac, Pb, and Po) were made on C-Disk-IV cruise samples.

**Table 1:** Station information for C-Disk-IV cruise. Sediment description is taken from Hou et al. (2019). Porosity was estimated from water loss of sediments for top 3 cm. Density of sediments was assumed  $2.5 \text{ g cm}^{-3}$ . Linear sedimentation rates were found based on fits to the  $^{226}\text{Ra}$  profiles below the mixed layer, as these reflect  $^{230}\text{Th}$  profiles.

Station	Location (°N, °W)	Bottom Depth (m)	Surface Porosity (%)	Sed. Rate (cm kyr <sup>-1</sup> )	Sediment Description
1	22.8, 158.0	4730	86	0.44	uniform light reddish-brown color, tube/worm visible at the sediment surface
2	27.8, 155.2	5640	83	0.11	uniform light brown color, manganese nodules
3	35.3, 151.0	5560	87	0.23	uniform light brown color, manganese nodules, stiff/icky texture of sediments below 2cm
4	41.8, 148.2	4970	89	0.15	very light brown color, 1-3cm fluff layer above sediment, burrows
5	49.8, 149.6	4725	80	0.57	greyish brown color, darker between 5-10cm, grittier below 5cm, worm/tube visible

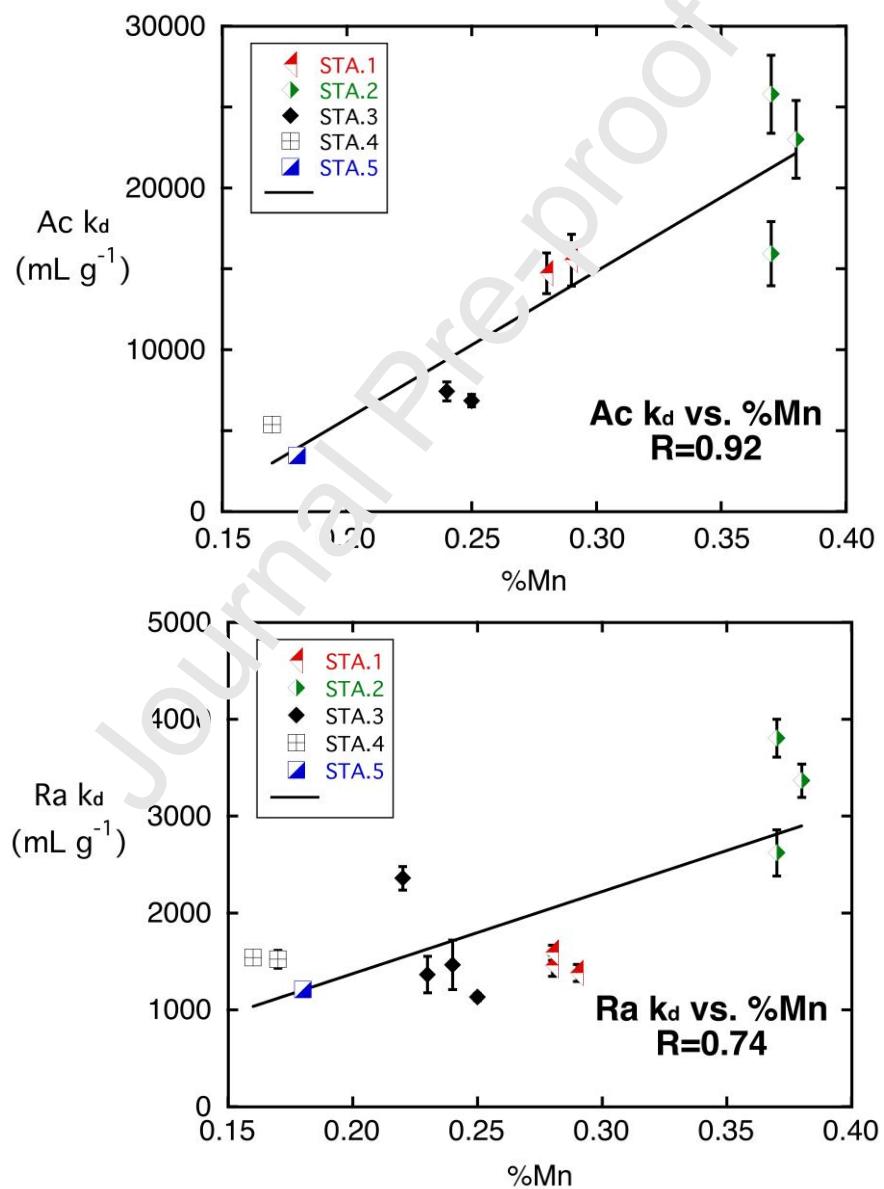
**Table 2:**  $^{228}\text{Ra}$   $k_d$  results using  $^{232}\text{Th}$  spike solution.  $^{228}\text{Ra}$  was measured in the dissolved and solid phase after rotating sediment slurry for at least one hour.

CDISK Station	depth cm	pH	mass:water ratio	$^{228}\text{Ra}$ $k_d$ $\text{mL g}^{-1}$	# measured <i>n</i>
1	0.5	7.6	0.003	1381±88	1
1	1.5	7.7	0.165	1456±77	2
1	2.5	7.9	0.004	1591±78	1
1	17	7.9	0.006	2117±39	1
2	0.5	7.9	0.005	262±169	2
2	1.5	7.6	0.004	3804±195	2
2	2.5	7.9	0.004	3366±171	1
2	17	7.7	0.006	4438±216	1
3	0.5	7.8	0.005	1133±48	2
3	1.5	7.7	0.100	1464±180	2
3	2.5	7.7	0.120	1364±133	2
3	6.5	7.7	0.006	2357±121	1
4	1.5	7.8	0.004	1521±53	3
4	2.5	7.7	0.004	1540±63	1
5	1.5	7.6	0.006	1210±68	1

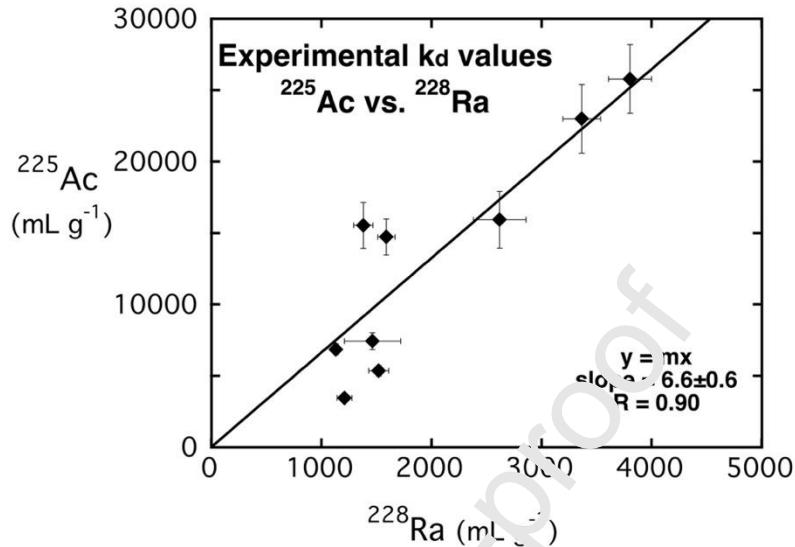
**Table 3:**  $^{225}\text{Ac}$   $k_d$  results using  $^{229}\text{Th}$  spike solution.  $^{225}\text{Ac}$  was measured in the dissolved and solid phases after rotating sediment slurry for one hour.

CDISK Station	depth cm	pH	mass:water ratio	$^{225}\text{Ac}$ $k_d$ $\text{mL g}^{-1}$	# measured <i>n</i>
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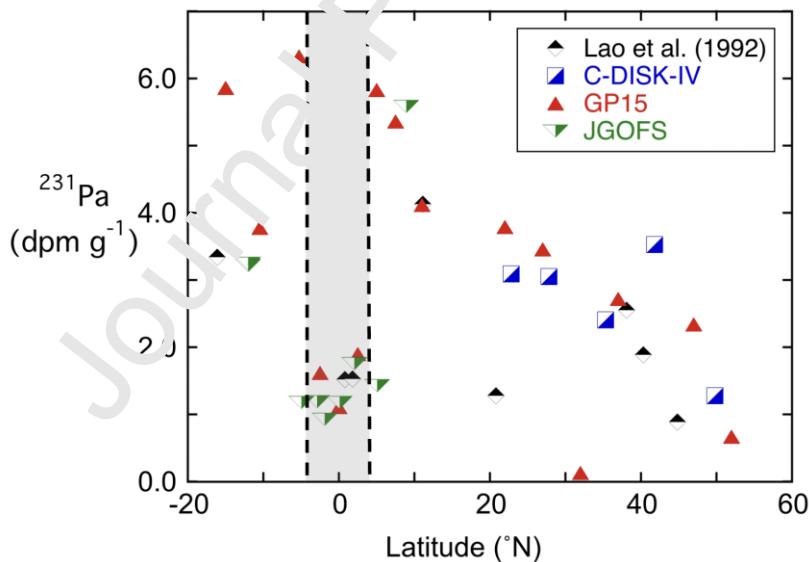
1	0.5	7.4	0.003	15532±1605	3
1	2.5	7.6	0.003	14722±1257	1
2	0.5	7.5	0.003	15931±1984	2
2	1.5	7.6	0.005	25786±2408	2
2	2.5	7.6	0.005	23000±2400	1
3	0.5	7.4	0.003	6846±382	2
3	1.5	7.5	0.003	7433±587	2
4	1.5	7.3	0.004	5354±369	2
5	1.5	7.5	0.003	3453±123	2



**Figure 2:** Ac and Ra  $k_d$  vs. %Mn from aliquots taken from 0-7 cm depth for all C-Disk-IV stations. Spikes of  $^{225}\text{Ac}$  and  $^{228}\text{Ra}$  were added to slurries of seawater and sediments, agitated for an hour, separated, and the activities of water and solids were measured by gamma spectroscopy.

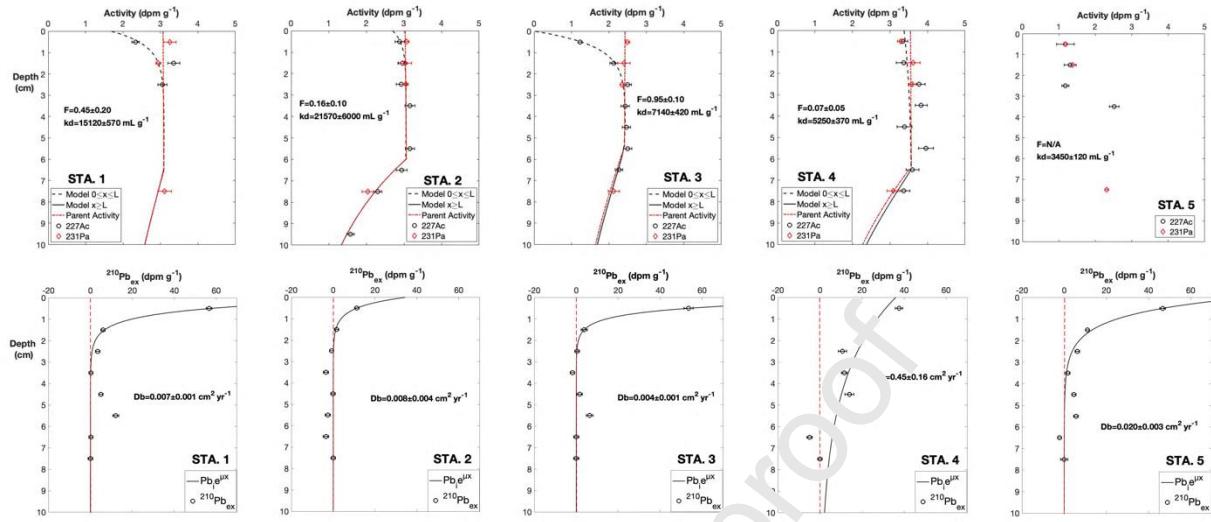


**Figure 3:**  $^{225}\text{Ac}$  vs.  $^{228}\text{Ra}$   $k_d$  values. The regression line is forced through zero.

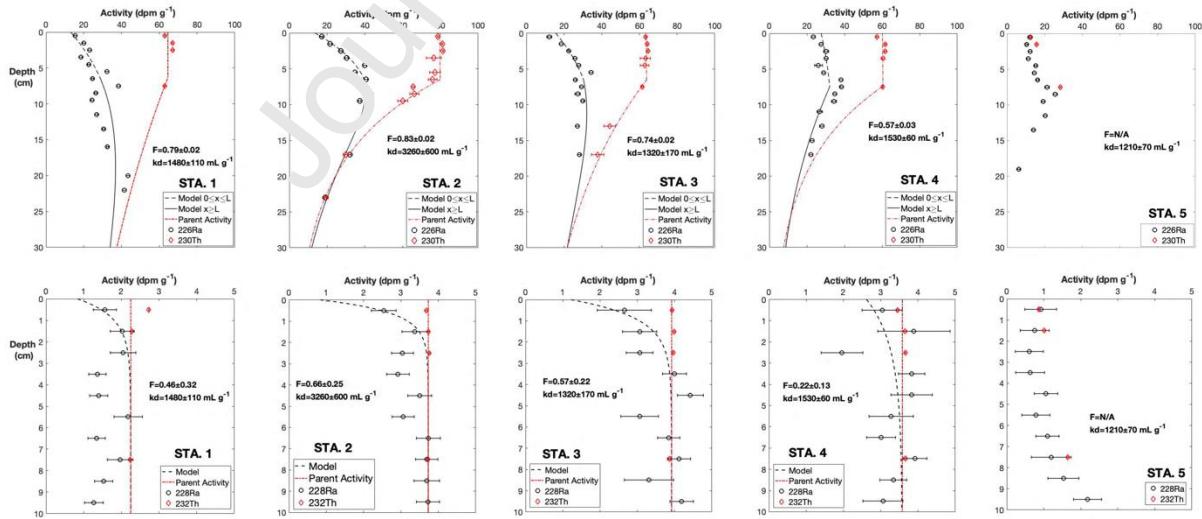


**Figure 4:**  $^{231}\text{Pa}$  activity in surface sediments vs. latitude in the Pacific Ocean (roughly along 150°W). Data taken from Lao et al. (1992), JGOFS (See SI for JGOFS data), and this study (C-Disk-IV and GP15). Equatorial stations (2.5°S - 2.5°N) have lower  $^{231}\text{Pa}$  activities than at surrounding sites because preservation of  $\text{CaCO}_3$  there dilutes the  $^{231}\text{Pa}$  scavenged from the overlying water column. The site from the Hawaiian margin (20°N) has a lower  $^{231}\text{Pa}$  value due to both higher sedimentation rate and shallower depth, which leads to less production of  $^{231}\text{Pa}$  in the overlying water column. The general decrease in  $^{231}\text{Pa}$  concentration northward of the equatorial cores reflects primarily a dilution effect from the northward increase in sediment mass

accumulation rate. The low value at  $\sim 31^\circ\text{N}$  (GP15 Sta 14) is believed to reflect coring of an erosional surface, and the relatively low values near  $10^\circ\text{S}$  may be influenced by downslope transport from adjacent topographic highs.



**Figure 5:** Solid phase  $^{231}\text{Pa}$ ,  $^{227}\text{Ac}$ , and  $^{210}\text{Pb}_{\text{ex}}$  profiles for all 5 C-Disk-IV stations. Eq. 6 was fit to the solid phase  $^{227}\text{Ac}$  profile in the upper 3 cm of sediments (Top 5 profiles). Station 5 could not be fit due to equilibrium between  $^{227}\text{Ac}$  and  $^{231}\text{Pa}$  in the upper 3 cm.  $^{210}\text{Pb}_{\text{ex}}$  profiles were calculated by taking the difference of total  $^{210}\text{Pb}$  and  $^{226}\text{Ra}$ . Measured  $^{226}\text{Ra}$  was corrected for *in situ* diffusive loss of  $^{222}\text{Rn}$  from sediments based on deep-sea profiles in the equatorial Pacific (Hammond et al., 1996). An exponential function was fit to the  $^{210}\text{Pb}_{\text{ex}}$  profiles that gave values for initial  $^{210}\text{Pb}_{\text{ex}}$  at the SWI and attenuation ( $\mu$ ).  $\mu$  is the inverse scale length of the exponential curve and it can be used to calculate the bioturbation rate ( $D_b$ ) in units of  $\text{cm}^2 \text{yr}^{-1}$ . Bioturbation is assumed constant throughout the entire  $^{210}\text{Pb}_{\text{ex}}$  profile (Bottom 5 profiles).



**Figure 6:** Solid phase  $^{230}\text{Th}$ ,  $^{226}\text{Ra}$ ,  $^{232}\text{Th}$ , and  $^{228}\text{Ra}$  profiles for all 5 C-Disk-IV stations. Eq. 6 was fit to the entire solid phase  $^{226}\text{Ra}$  profile (Top 5 profiles). Station 5 could not be fit due to equilibrium between  $^{230}\text{Th}$  and  $^{226}\text{Ra}$  in the upper 8 cm. The  $^{228}\text{Ra}$  profiles were fit with eq. 9,

which is a solution to a one-layer model. All parameters such as  $^{232}\text{Th}$  activity and  $D_b$  are assumed constant throughout the entire  $^{228}\text{Ra}$  profile (Bottom 5 Profiles).

**Table 4:** Model parameters for best fit solid phase  $^{227}\text{Ac}$ ,  $^{226}\text{Ra}$ , and  $^{228}\text{Ra}$  profiles (Eq. 6).  $D_b$ ,  $k_d$ ,  $\phi$ ,  $D_m$ ,  $^{230}\text{Th}$ ,  $^{232}\text{Th}$ , and  $^{231}\text{Pa}$  were all measured independently or considered known variables.  $S$  was determined from  $^{226}\text{Ra}$  model parameters. Mixing depth (L) was determined by changes in  $^{230}\text{Th}$  or  $^{231}\text{Pa}$  activity at depth. Porosity ( $\phi$ ) was calculated by averaging the porosity in the upper 3cm of sediments. Only F was allowed to vary in the model until the best fit was achieved.

Sta.	$^{231}\text{Pa}$	$D_b$	$k_d$	L	$\phi$	$D_s$	S	F	$^{227}\text{Ac Flux}$
#	dpm g $^{-1}$	cm $^2$ yr $^{-1}$	mL g $^{-1}$	cm		cm $^2$ yr $^{-1}$	cm kyr $^{-1}$		dpm m $^{-2}$ yr $^{-1}$
1	3.09 $\pm$ 0.12	0.007 $\pm$ 0.001	15120 $\pm$ 570	6.5	0.82	62	0.40	45 $\pm$ 0.20	130 $\pm$ 60
2	3.05 $\pm$ 0.10	0.008 $\pm$ 0.004	21570 $\pm$ 6000	6.5	0.78	56	0.10	11 $\pm$ 0.10	40 $\pm$ 30
3	2.41 $\pm$ 0.09	0.004 $\pm$ 0.001	7140 $\pm$ 420	6.5	0.81	60	0.20	95 $\pm$ 0.10	290 $\pm$ 80
4	3.53 $\pm$ 0.12	0.450 $\pm$ 0.160	5250 $\pm$ 370	7.5	0.87	69	0.11	0.07 $\pm$ 0.05	90 $\pm$ 70
5	1.28 $\pm$ 0.03	0.020 $\pm$ 0.003	3450 $\pm$ 120	N/A	0.79	57	N/A	N/A	N/A

Sta.	$^{230}\text{Th}$	$D_b$	$k_d$	L	$\phi$	$D_s$	S	F	$^{226}\text{Ra Flux}$
#	dpm g $^{-1}$	cm $^2$ yr $^{-1}$	mL g $^{-1}$	cm		cm $^2$ yr $^{-1}$	cm kyr $^{-1}$		dpm m $^{-2}$ yr $^{-1}$
1	64.9 $\pm$ 0.3	0.007 $\pm$ 0.001	1480 $\pm$ 110	6.5	0.80	85	0.40	0.79 $\pm$ 0.02	1690 $\pm$ 490
2	80.3 $\pm$ 0.4	0.008 $\pm$ 0.004	3260 $\pm$ 600	6.5	0.75	75	0.10	0.83 $\pm$ 0.02	1540 $\pm$ 780
3	64.0 $\pm$ 0.3	0.004 $\pm$ 0.001	1320 $\pm$ 170	6.5	0.78	81	0.20	0.74 $\pm$ 0.02	1580 $\pm$ 360
4	60.3 $\pm$ 0.4	0.450 $\pm$ 0.160	1530 $\pm$ 60	7.5	0.82	89	0.17	0.57 $\pm$ 0.03	720 $\pm$ 260
5	14.1 $\pm$ 0.1	0.020 $\pm$ 0.003	1210 $\pm$ 70	N/A	0.76	77	N/A	N/A	N/A

Sta.	$^{232}\text{Th}$	$D_b$	$k_d$	$\phi$	$D_s$	S	F	$^{228}\text{Ra Flux}$
#	dpm g $^{-1}$	cm $^2$ yr $^{-1}$	mL g $^{-1}$		cm $^2$ yr $^{-1}$	cm kyr $^{-1}$		dpm m $^{-2}$ yr $^{-1}$
1	2.25 $\pm$ 0.02	0.008 $\pm$ 0.001	1470 $\pm$ 110	0.82	89	0.40	0.46 $\pm$ 0.32	520 $\pm$ 360
2	3.73 $\pm$ 0.02	0.008 $\pm$ 0.004	5250 $\pm$ 600	0.78	81	0.10	0.66 $\pm$ 0.25	900 $\pm$ 570
3	3.92 $\pm$ 0.02	0.004 $\pm$ 0.001	1320 $\pm$ 170	0.81	87	0.20	0.57 $\pm$ 0.22	1440 $\pm$ 660
4	3.58 $\pm$ 0.02	0.450 $\pm$ 0.160	1530 $\pm$ 60	0.87	100	0.17	0.22 $\pm$ 0.13	700 $\pm$ 390
5	0.93 $\pm$ 0.02	0.020 $\pm$ 0.003	1210 $\pm$ 70	0.79	N/A	N/A	N/A	N/A

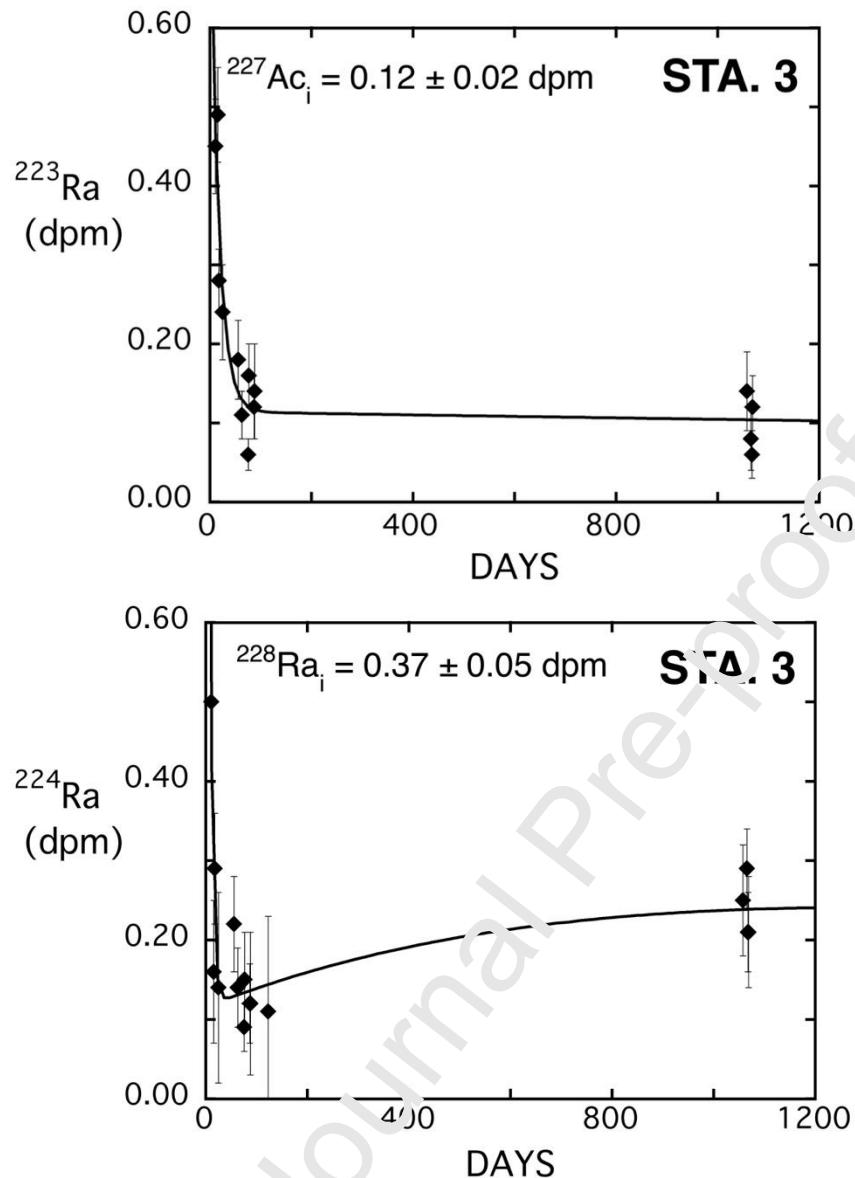
**Table 5:** Mixed layer depth (L),  $^{230}\text{Th}$  activity, inventory of  $^{230}\text{Th}$  for upper- and lower-layer sediments and water column (Inv.), focusing factor (g), mass accumulation rates (MAR),

sedimentation rate (S), and ages are shown for 5 C-Disk-IV profiles. Upper and lower layers, denoted as UL and LL are referenced to layers below or above the mixing depth. Lastly, total inventory of  $^{230}\text{Th}$  in the sediments (Tot. Inv.) and water column (WC) are shown, with g = focusing factor = (sediment inventory)/(water column production).

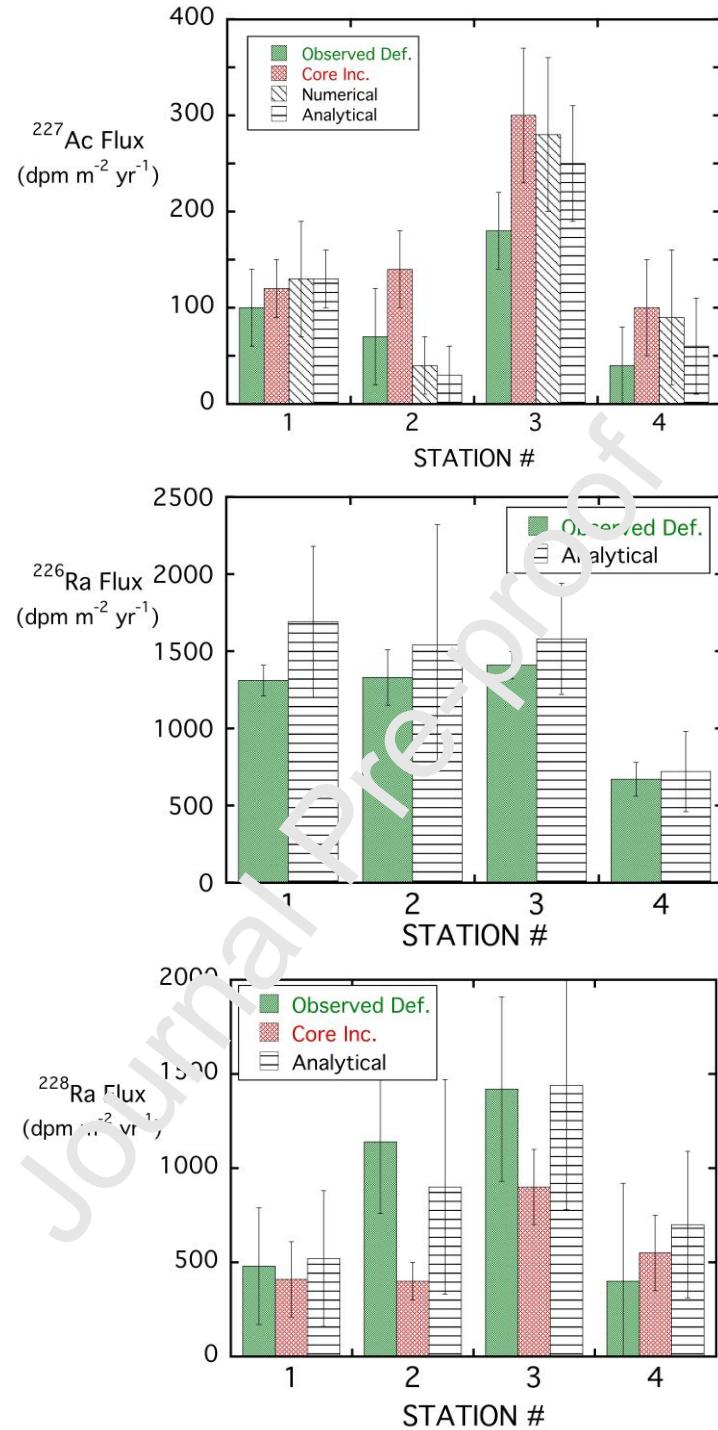
STA. #	L	UL $^{230}\text{Th}$	UL Inv.	LL Inv.	Tot. Inv.	WC Inv.	Focusing Factor	UL MAR	LL MAR	UL S	LL S	UL Age kyr
	cm	dpm g $^{-1}$	dpm cm $^{-2}$	dpm cm $^{-2}$	dpm cm $^{-2}$	dpm cm $^{-2}$	g	g cm $^{-2}$ kyr $^{-1}$	g cm $^{-2}$ kyr $^{-1}$	cm kyr $^{-1}$	cm kyr $^{-1}$	
1	6.5	65	212	1559	1771	1356	1.31	0.221	0.248	0.44	0.40	15
2	6.5	80	322	588	910	1617	0.56	0.066	0.070	0.11	0.10	61
3	6.5	64	231	905	1136	1594	0.71	0.129	0.135	0.23	0.20	28
4	7.5	60	161	496	657	1425	0.46	0.069	0.095	0.15	0.17	39
5	6.5	15	75	543	618	1355	0.46	0.345	0.204	0.57	0.25	11

**Table 6:** Fluxes derived from Bateman equation fits to the core incubation samples measured over 1200 days. These fits provided the  $^{227}\text{Ac}$ ,  $^{223}\text{Ra}$ , and  $^{228}\text{Ra}$  activity in the core-top water at the conclusion of the incubation. Time/height is the summation of time incubated divided by height of the remaining water in the core for all segments of the incubation. Volume indicates the amount of seawater that was passed through the Mn-fibers.  $a_i$  is the effective inverse height needed to calculate flux for a short-lived radioisotope (see SI S 2).  $^{227}\text{Ac}$ ,  $^{223}\text{Ra}$ , and  $^{228}\text{Ra}$  fluxes were calculated by eq. 11.

Sta.	time/height	Vol.	$^{227}\text{Ac}$	$^{223}\text{Ra}$	$^{228}\text{Ra}$	$^{227}\text{Ac}$ Flux	$^{228}\text{Ra}$ Flux	$a_i$	$^{223}\text{Ra}$ Flux
	days m $^{-1}$	L	dpm	dpm	dpm	dpm m $^{-2}$ yr $^{-1}$	dpm m $^{-2}$ yr $^{-1}$	m $^{-1}$	dpm m $^{-2}$ yr $^{-1}$
1	142	1.35	0.07 $\pm$ 0.01	0.83 $\pm$ 0.12	0.22 $\pm$ 0.05	120 $\pm$ 30	410 $\pm$ 200	5.60	2430 $\pm$ 850
2	171	1.29	0.09 $\pm$ 0.03	0.90 $\pm$ 0.22	0.24 $\pm$ 0.17	140 $\pm$ 40	390 $\pm$ 130	6.49	2400 $\pm$ 620
3	128	1.16	0.12 $\pm$ 0.02	0.82 $\pm$ 0.17	0.37 $\pm$ 0.05	280 $\pm$ 60	880 $\pm$ 220	5.49	2870 $\pm$ 560
4	95	1.00	0.03 $\pm$ 0.01	0.41 $\pm$ 0.06	0.1 $\pm$ 0.05	100 $\pm$ 50	650 $\pm$ 220	4.35	2080 $\pm$ 920
5	51	1.00	0.12 $\pm$ 0.02	0.48 $\pm$ 0.06	0.21 $\pm$ 0.05	850 $\pm$ 110	850 $\pm$ 490	2.70	3980 $\pm$ 1200



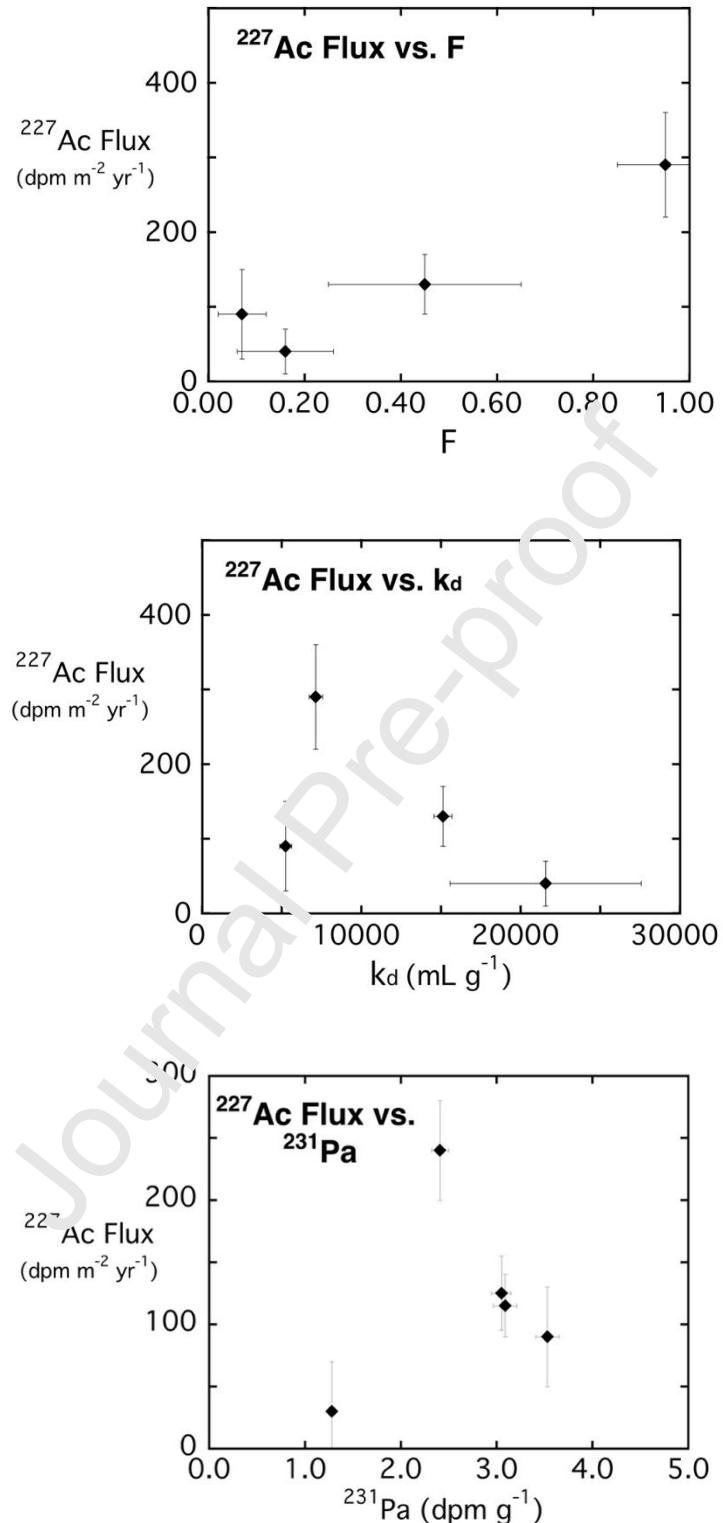
**Figure 7:** Bateman equation fits to determine  $^{223}\text{Ra}$ ,  $^{224}\text{Ra}$ ,  $^{228}\text{Ra}$ , and  $^{227}\text{Ac}$  activities in incubation samples at the conclusion of the incubation. Samples were stored and counted for up to 3 years while  $^{223}\text{Ra}$  and  $^{224}\text{Ra}$  grew into equilibrium with their longer-lived ancestors. Only one station is shown for reference (STA. 3).  $^{223}\text{Ra}$  grows into equilibrium with its grandparent activity,  $^{227}\text{Ac}$ , by 125 days, while ingrowth of  $^{224}\text{Ra}$  toward equilibrium with its  $^{228}\text{Ra}$  parent is still continuing. Samples from most stations had very little excess  $^{224}\text{Ra}$  activity due to long wait times between collection and counting ( $^{224}\text{Ra}$ :  $t_{1/2}=3.4$  days).



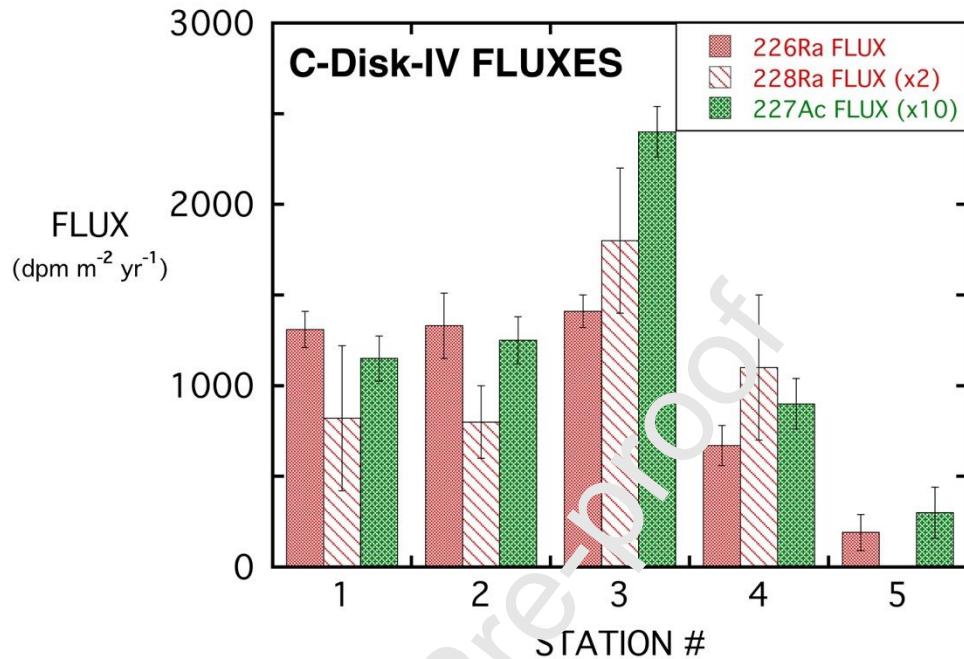
**Figure 8: Top:**  $^{227}\text{Ac}$  flux was determined by all methods described in this manuscript: subtracting  $^{227}\text{Ac}$  and  $^{231}\text{Pa}$  from each sediment horizon in the top 3 cm of sediments (observed deficiency), curve fitting solid phase  $^{227}\text{Ac}$  vs. depth using eq.6 (reaction-transport model), creating a model that accounts for varying porosity downcore (numerical model), and direct measurements via core incubation (core incubation). Observed deficiency and core incubation should give the most accurate results. **Middle:**  $^{226}\text{Ra}$  flux was determined by two methods:

observed deficiency and curve fitting solid phase  $^{226}\text{Ra}$  vs. depth (eq.6). **Bottom:**  $^{228}\text{Ra}$  flux was determined by observed deficiency, core incubation, and curve fitting solid phase  $^{228}\text{Ra}$  vs. depth (eq.9). Note that core incubation results are most precise for this isotope.

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**Fig. 9:**  $^{227}\text{Ac}$  Flux vs.  $F$ ,  $k_d$ , and  $^{231}\text{Pa}$  for C-Disk-IV sediments. The results for  $F$  and  $k_d$  are taken from the reaction-transport model and  $^{231}\text{Pa}$  is  $^{231}\text{Pa}$  activity in the mixed layer.



**Fig. 10:**  $^{226}\text{Ra}$ ,  $^{228}\text{Ra}$ , and  $^{227}\text{Ac}$  Flux vs. Station # for C-Disk-IV sediments. The  $^{228}\text{Ra}$  and  $^{227}\text{Ac}$  fluxes were multiplied x2 and x10 in order to produce comparable scales, since  $^{226}\text{Ra}$  flux was so much larger.

# ACTINIUM AND RADIUM FLUXES FROM THE SEABED IN THE NORTHEAST PACIFIC BASIN

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## Highlights

First study to measure  $^{227}\text{Ac}$  sediment profiles along a cruise tract in Northeast Pacific. The distribution coefficient ( $k_d$ ) for Ac in these sediments is about 7 times higher than that for Ra. These are the first measurements of  $k_d$  for Ac.

Benthic fluxes of  $^{227}\text{Ac}$  and  $^{228}\text{Ra}$  determined by core incubation shows results comparable to predictions from transport models applied to sediment profiles.

Average benthic fluxes of  $^{227}\text{Ac}$ ,  $^{228}\text{Ra}$ , and  $^{226}\text{Ra}$  in the Northeast Pacific are  $90 \pm 20$ ,  $600 \pm 200$ , and  $1300 \pm 200$  dpm  $\text{m}^{-2}\text{-yr}^{-1}$ , constraining interpretations of water column data.