

Stress-Induced High- T_c Superconductivity in Solid Molecular Hydrogen

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Solid molecular hydrogen has been predicted to be metallic and high-temperature superconducting at ultrahigh hydrostatic pressures that push current experimental limits. Meanwhile, little is known about the influence of non-hydrostatic conditions on its electronic properties at extreme pressures where anisotropic stresses are inevitably present and may also be intentionally introduced. Here we show by first-principles calculations that solid molecular hydrogen compressed to multimegabar pressures can sustain large anisotropic compressive or shear stresses which, in turn, cause major crystal symmetry reduction and charge redistribution that accelerate bandgap closure and promote superconductivity relative to pure hydrostatic compression. Our findings highlight a hitherto largely unexplored mechanism for creating superconducting dense hydrogen, with implications for exploring similar phenomena in hydrogen-rich compounds and other molecular crystals.

high pressure | anisotropic stresses | metallic hydrogen | superconductivity | first-principles calculations

Since Wigner and Huntington's pioneering work (1), pressure-induced metallization of hydrogen has attracted immense interest and impelled advances of theoretical and computational methods, ultrahigh-pressure devices, and related measurement and characterization techniques (2, 3). The prediction (4) that metallic hydrogen may be a high-temperature superconductor driven by the BCS-type phonon-mediated mechanism further invigorated the study of this fascinating material and its enigmatic properties. As a prominent thermodynamic quantity, pressure, which is defined as hydrostatic, drives the formation of new material phases through structural transitions and electronic bandgap closure (5). Two main avenues have been pursued for producing metallic hydrogen at high pressure: (i) closure of the electronic bandgap of solid molecular hydrogen and (ii) dissociation of hydrogen molecules to form monatomic metal. These mechanisms have been explored through optical and electrical conductivity techniques over the years (6–16). In addition, at least five molecular phases, labeled I–V, have been documented experimentally to pressures up to 400 GPa and temperatures 300 K (17–21). These very high static pressures can significantly limit, constrain, and challenge full experimental characterization of materials, especially above 300 GPa, leading to ambiguities and controversies in the interpretation of reported results.

Theoretical studies (22–27) of the metallization of solid hydrogen have focused on its formation under hydrostatic pressure. However, conditions inside diamond anvil cells (DACs) at the required pressures (e.g., >300 GPa) can introduce nonuniform deformations and anisotropic stresses in the vicinity of the sample as revealed by direct x-ray imaging study above

300 GPa (28). More recent studies have further explored these effects (29, 30) and have also unveiled unprecedented structural ductility in diamond under large multiaxial deformation modes (31, 32), underscoring complex strain conditions inside DACs under extreme loading conditions. It is therefore instructive to establish the influence of anisotropic stresses on the evolution of structural and electronic properties of molecular hydrogen under nonhydrostatic pressure conditions [More discussion of non-hydrostatic stress states in DACs can be found in the Supporting Information (SI)]. Although gases that are solidified under pressure are often assumed to be hydrostatic over a range of pressures, direct measurements indicate the development of anisotropic strains even well below 100 GPa (33). Moreover, significant pressure gradients were documented even in the earliest optical studies of closing of the bandgap above 200 GPa (2, 7).

Here we report an exploratory study of metallization and superconductivity of solid molecular hydrogen under anisotropic stresses, i.e., at non-hydrostatic pressures. Our first-principles stress-strain calculations establish that despite being a soft and plastic crystal at low pressures (and temperatures) solid hydrogen at megabar pressures can sustain considerable uniaxial compressive and shear stresses, and such anisotropic stresses can have major impact on physical properties. Promi-

Significance Statement

Metallic hydrogen has long been considered the Holy Grail of high pressure physics that holds the intriguing possibility of realizing room-temperature superconductivity. However, the ultrahigh pressures required for hydrogen metallization impose severe experimental challenges and constraints on synthesis and characterization, leaving large ambiguities and controversies in the interpretation of reported results. We show that non-hydrostatic stresses offer a viable route to promote metallization and superconductivity in hydrogen at more accessible static compression conditions (e.g., <300 GPa). Such stress anisotropy is expected to be highly dependent on sample and measurement conditions, which may explain the discrepancies in previously reported experimental results, including variations in the rate of bandgap closure and onset of electrical conductivity in solid molecular hydrogen.

Q.L., R.J.H., Y.M., and C.C. designed the research; X.S. and C.L. performed the calculations; X.S., C.L., Q.L., R.J.H., Y.M., and C.C. analyzed the data and wrote the paper.

The authors declare no competing interest.

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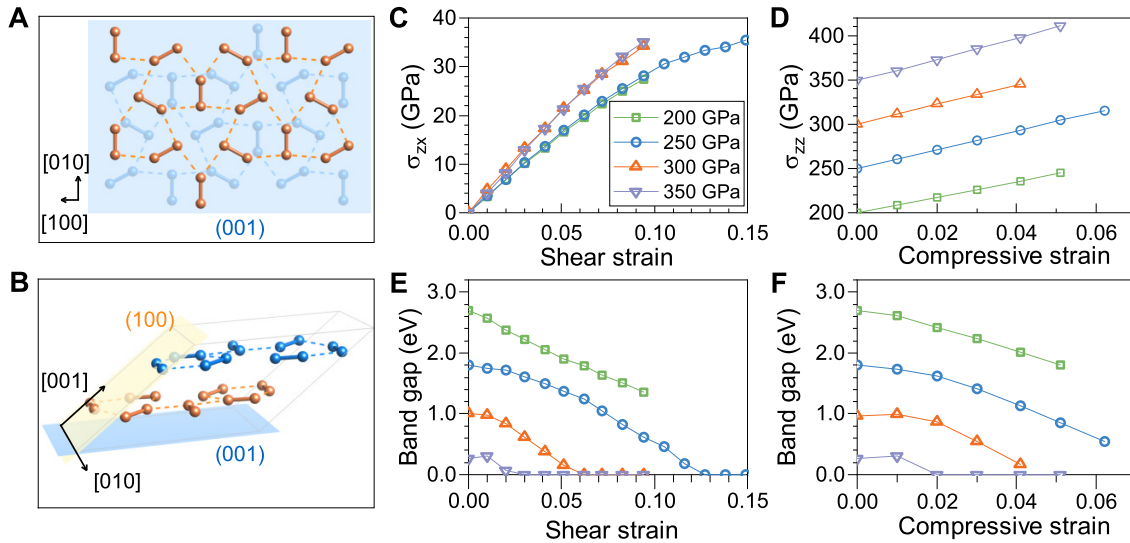


Fig. 1. Structure and stress and bandgap responses of solid molecular hydrogen to anisotropic stresses. (A) Top and (B) polyhedral view of the $C2/c$ phase. The thick solid lines represent the intramolecular covalent bonds of H_2 and the thin dash lines are intermolecular links. The light grey lines outline the unit cell. (C) First-principles determined stable stress responses to shear strains along the (010)[001] direction (the loading path with the steepest descent of E_g) at hydrostatic pressures of 200, 300, and 350 GPa and along the (001)[010] direction (the loading path with the easiest metallization of E_g) at 250 GPa. (D) First-principles determined stable stress responses to compressive strains in the [100] direction at 200, 250, 300, and 350 GPa. (E and F) Shear and compressive strain-modified E_g along the above selected loading paths.

55 nent effects include lattice symmetry reduction and ensuing
56 changes in the electronic band structure that accelerate the
57 bandgap closure, leading to metallization, and enhance the
58 electron-phonon coupling (EPC) giving rise to superconductiv-
59 ity. In particular, we find that shear deformations are highly
60 effective in inducing metallic and superconducting states at
61 much reduced hydrostatic pressures. These phenomena are
62 driven by robust underpinning mechanisms and are therefore
63 expected to remain intact when possible corrections by ther-
64 mal and quantum effects are further considered. The present
65 findings may explain loading-dependent differences in reported
66 experimental results (9–16, 34, 35) and show that introducing
67 controlled non-hydrostatic stresses is a viable route to promote
68 metallization and superconductivity in hydrogen at readily
69 accessible static compression conditions (e.g., <300 GPa).

70 **Stress responses and bandgap evolution of solid molecular**
71 **hydrogen.** To explore the impact of anisotropic stresses on the
72 transition of solid hydrogen from insulating to metallic and
73 potentially high- T_c superconducting states, we focus on phase
74 III of solid hydrogen, which was discovered in 1988 (2, 6) and
75 has been predicted and observed to be the stable insulating
76 phase at temperatures below 200 K in the pressure range of
77 about 150–400 GPa before potentially turning metallic upon
78 further compression. The extremely high pressures required
79 to achieve metallization has been a formidable challenge to
80 reliable measurements and analysis. Theoretical calculations
81 (17, 36) for phase III have predicted two nearly degenerate
82 thus coexisting structures of $C2/c$ symmetry containing 12
83 and 24 atoms per unit cell, respectively, which are compatible
84 with the low-temperature Raman and infrared (IR) spectra
85 (36). In this work, we focus on the $C2/c$ -12 structure (see
86 Fig. 1 A and B) as an exemplary case study; our studies show
87 that this phase is able to sustain large anisotropic stresses and
88 more easily transform into metallic states. We first performed
89 extensive first-principles calculations to determine the stress
90 responses of the $C2/c$ structure to anisotropic compression

and shear strains in the hydrostatic pressure range of 200–350
GPa, and key representative stress responses are summarized
in Fig. 1 C and D (full stress-strain relations are given in *SI*
Appendix Figs. S1 and S2 and crystal structure information
of representative deformed molecular hydrogen is listed in *SI*
Appendix Table S1). It is seen that the structure can sustain
considerable anisotropic compressive stresses ranging from
30 to 60 GPa that increase with rising hydrostatic pressure
of 200–350 GPa or shear stresses of about 30 GPa that are
insensitive to hydrostatic pressure.

Electronic band-structure calculations show that these non-
hydrostatic stresses that give rise to anisotropic strains lower
the bandgap E_g relative to the equivalent isotropic stresses or
hydrostatic pressures (see *SI Appendix* Fig. S3) as shown in Fig.
1 E and F), stemming from band shifts and splittings caused by
the symmetry reduction of the deformed crystal. In particular,
we find that shear strains ε_{zx} are highly effective in reducing
 E_g ; for instance, a shear stress $\sigma_{zx}=20$ GPa at $\varepsilon_{zx}=0.06$ along
the (010)[001] direction reduces E_g at hydrostatic pressure of
200 GPa to approximately E_g at 250 GPa. Remarkably, a
shear stress $\sigma_{zx}=33$ GPa along the (001)[010] direction causes
a complete bandgap closure at 250 GPa (see Fig. 1E), which
occurs at 370 GPa under hydrostatic conditions. At higher
pressures, bandgap closure occurs at smaller shear strains, as
shown in Fig. 1E for selected cases (see *SI Appendix* Fig. S1
for more comprehensive results).

We also examine the evolution of E_g with excess uniaxial
compressive strain (ε_{zz}) added to compressive strains induced
by hydrostatic pressure. The results show a reduction of E_g at
the rate of about 20 meV/GPa ($\Delta E_g/\Delta \sigma_{zz}=0.9$ eV/45 GPa),
which is greater than 16 meV/GPa ($\Delta E_g/\Delta P=2.4$ eV/150
GPa) of E_g reduction under hydrostatic pressure (see Fig. 1
D and F). The metallization can be realized at 350 GPa by
introducing an additional normal compressive stress of about
23 GPa. Additional results on the compressive stress-strain
and corresponding E_g -strain relations presented in *SI Appendix*

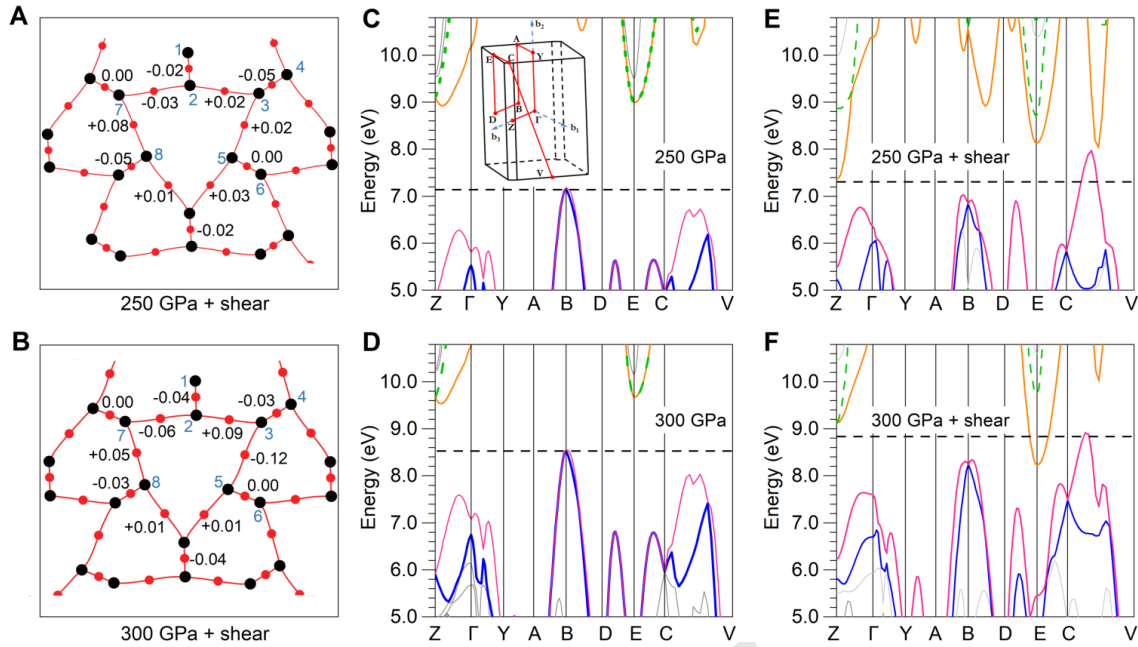


Fig. 2. Electronic properties of shear deformed $C2/c$ hydrogen. (A and B) QTAIM analysis of charge states in the (001) plane of shear deformed $C2/c$ -based structure at 250 GPa and 300 GPa. The large (small) black (red) spheres represent hydrogen atoms and bond critical points (BCPs), respectively. The change of charge density for BCPs is given in the atom unit. (C and D) Electronic band structures at the hydrostatic pressure of 250 GPa and 300 GPa. The inset shows the Brillouin zone path for calculating the band structures. (E and F) Electronic band structures of shear deformed $C2/c$ -based structure at $\sigma_{zx}=35$ GPa ($\varepsilon_{zx}=0.149$) in the (001)[010] direction at 250 GPa and at $\sigma_{zx}=34$ GPa ($\varepsilon_{zx}=0.094$) in the (010)[001] direction at 300 GPa. The black dashed lines mark the Fermi energy. The bands are colored to exhibit band shifts and splits.

Fig. S2 indicate that anisotropic compressive stresses can also play a role but are less effective compared with shear stresses in inducing the metallization.

Shear-induced metallization of molecular hydrogen. To further elucidate the appreciable modifications of electronic properties in non-hydrostatic environments, we performed a quantum theory of atoms in molecules (QTAIM) analysis of changes of charge states in response to stress-induced structural variations. We examine the bonding structures and charge density at bond critical points (BCPs) to evaluate the bonding strength and charge transfer. When hydrostatic pressure increases from 250 GPa to 300 GPa, the charge density values at the BCPs of the hydrogen molecules remain unchanged (2.07 to 2.07 and 2.12 to 2.12) and the values between hydrogen molecules rise by about 0.08 a.u. (0.62 to 0.70, 0.61 to 0.69, and 0.59 to 0.67) (see *SI Appendix* Fig. S4 for more details), implying almost constant bond length of H_2 molecules and increased intermolecular interactions, producing pressure-induced energy-band broadening and E_g decrease (see Fig. 2 C and D). Shear deformation drives intramolecular charge delocalization and differential evolution of hydrogen atomic distance (see *SI Appendix* Fig. S5 a and b) while enhancing intermolecular interactions and breaks the degeneracy of electronic states due to symmetry reduction. The change of charge density at BCPs of the $C2/c$ structure under combined hydrostatic pressure of 250 GPa and shear stress of $\sigma_{zx} = 35$ GPa ($\varepsilon_{zx} = 0.149$) in the (001)[010] direction (see Fig. 2A) shows an in-plane charge transfer from intramolecular to intermolecular regions (the distance of H7-H8 from 1.190 Å to 1.134 Å, H3-H5 from 1.190 Å to 1.171 Å), weakening the intramolecular bonds, elongating the H_2 units, and resulting in hole-type bands (see Fig. 2E). Meanwhile, shear deformation causes the growth of charge distribution

in the interlayer region under shear strains in the (010)[001] direction at 300 GPa (see Fig. 2B) with decreasing interlayer spacing (the distance of H3-H5 from 1.145 Å to 1.260 Å). The antibonding states of H_2 molecules start to be filled, lowering the conduction band minimum, producing the electron-type band crossing at the Fermi energy, and mediating an indirect band gap closure (see Fig. 2F).

Superconductivity in shear-deformed molecular hydrogen.

There have been extensive theoretical studies of potentially very high- T_c superconductivity in molecular hydrogen induced by hydrostatic compression (37–39). The present work expands such studies by introducing additional tuning mechanisms via anisotropic stresses, which drives metallization of the molecular solid at considerably reduced pressures. We further examine superconductivity of solid molecular hydrogen under shear strains along the easiest metallization paths in the (001)[010] direction at 250 GPa and the (010)[001] direction at 300 GPa. We find enhanced electron-phonon coupling under shear strains. To expose the underlying mechanism, we have calculated the phonon dispersion relations and the Eliashberg spectral function (see Fig. 3). Major contributions to the electron-phonon coupling parameter λ come from the low-frequency branches (acoustic phonons and librations) around 10–20 THz corresponding to the translational and rotational motions of the H_2 molecules. The phonons characterize the centers-of-mass molecular motions and the librations correspond to coupled restricted rotations (40). The distinguishable high-frequency vibrons are coupled stretchy vibrations of the H_2 molecules; notably these give rise to large peaks in the Eliashberg spectral function $\alpha^2F(\omega)$ at some q points but the high phonon frequency ω limits their contributions to the electron-phonon coupling parameter. For the (001)[010] shear loading

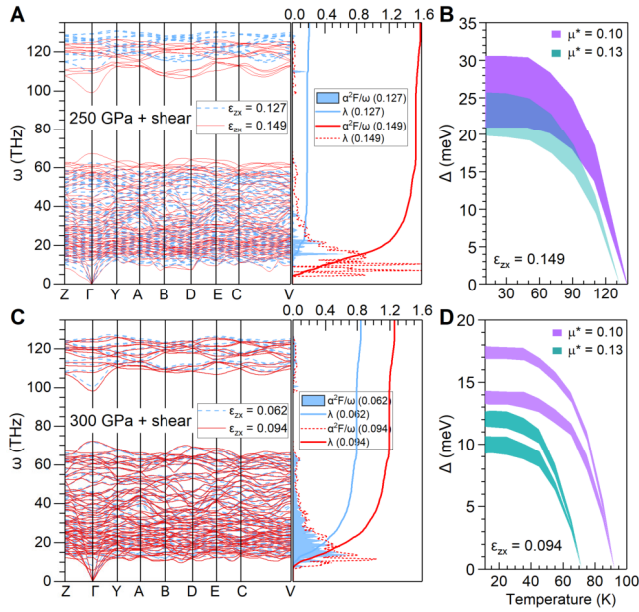


Fig. 3. Shear-induced superconductivity in $C2/c$ hydrogen. (A) Phonon dispersion relations $\omega(k)$ and Eliashberg spectral function $\alpha^2F(\omega)/\omega$ and $\lambda(\omega)$ at $\sigma_{zx}=33$ GPa ($\epsilon_{zx}=0.127$) and $\sigma_{zx}=35$ GPa ($\epsilon_{zx}=0.149$) and (B) the anisotropic superconducting gap (Δ) and superconducting critical temperatures (T_c) at shear strain $\epsilon_{zx}=0.149$ in the (001)[010] direction at 250 GPa. (C) Phonon dispersion relations $\omega(k)$ and Eliashberg spectral function $\alpha^2F(\omega)/\omega$ and $\lambda(\omega)$ at $\sigma_{zx}=25$ GPa ($\epsilon_{zx}=0.062$) and $\sigma_{zx}=34$ GPa ($\epsilon_{zx}=0.094$) and (D) the anisotropic superconducting gap (Δ) and T_c at shear strain $\epsilon_{zx}=0.094$ in the (010)[001] direction at 300 GPa.

path at 250 GPa, the in-plane charge transfers from the intramolecular to intermolecular regions, which weakens the covalent diatomic molecular bonds and softening the vibrons, reducing the intermolecular repulsion, and softening the librons (see Fig. 3A). Meanwhile, the hole-type band crosses the Fermi energy along the Brillouin zone path $C \rightarrow V$ (see *SI Appendix* Fig. S6 a and b), producing superconducting critical temperatures (T_c) of 132 K ($\mu^*=0.13$) to 140 K ($\mu^*=0.10$) at shear strain $\sigma_{zx}=0.149$ (see Fig. 3B). For the (010)[001] shear loading path at 300 GPa, the incremental accumulation of interlayer electrons with rising strains contributes to the electron-phonon coupling at $\epsilon_{zx}=0.094$ by initiating new topologies of the Fermi surface (see *SI Appendix* Fig. S6 c and d) and softening phonons (see Fig. 3C). The change produces a two-gap superconducting state stemming from the anisotropic and varied Fermi surface structures with a T_c of 72 K ($\mu^*=0.13$) to 93 K ($\mu^*=0.10$) (see Fig. 3D).

Molecular hydrogen under anisotropic compressive strains.

We also examined $C2/c$ hydrogen under anisotropic compressive strains at 350 GPa (see Fig. 4). The degeneracy of the conduction band remains intact, but the conduction band minimum drops as in the shear-deformed cases (see Fig. 4A and B). In contrast to the hydrostatic and shear-strain cases, the bond lengths of the hydrogen molecules increase here (see *SI Appendix* Fig. S5c), indicating more intramolecular electron transfer to the intermolecular (in-plane and interlayer) regions, in agreement with the charge density analysis of BCPs from QTAIM (see Fig. 4C). These new charge states produce complex band characteristics and Fermi surfaces (see *SI Appendix* Fig. S6 e and f), resulting in a multigap superconducting

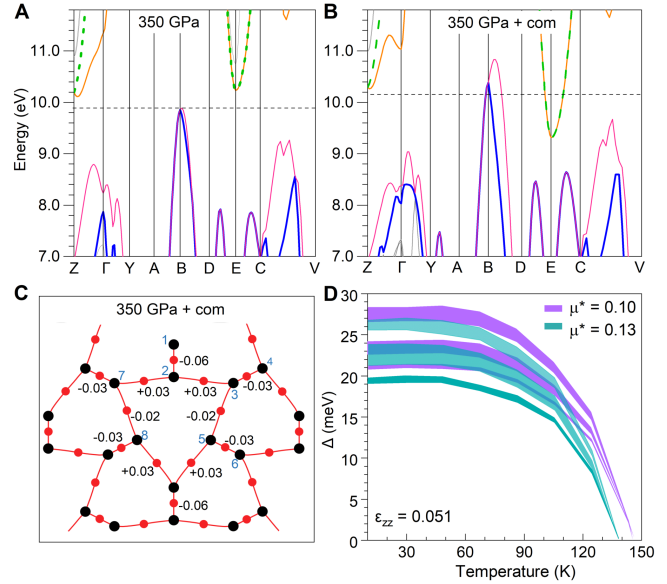


Fig. 4. Electronic properties of $C2/c$ hydrogen under anisotropic strain at 350 GPa. (A and B) Electronic band structures at 350 GPa and at $\sigma_{xx}=\sigma_{yy}=350$ GPa, $\sigma_{zz}=411$ GPa ($\epsilon_{zz}=0.051$) in the [100] direction. The black dashed lines mark the Fermi energy. The bands are colored to exhibit band shifts and splits. (C) The structure and charge density based on QTAIM analysis of $C2/c$ hydrogen under conditions in (B). (D) The corresponding superconducting gap (Δ) and T_c .

state with T_c of 138 K ($\mu^*=0.13$) to 145 K ($\mu^*=0.10$) at $\epsilon_{zz}=0.051$ under the anisotropic compressive stress conditions $\sigma_{xx}=\sigma_{yy}=350$ GPa and $\sigma_{zz}=411$ GPa (see Fig. 4D).

Discussion

First-principles calculations of stress-strain relations indicate that phase III of solid hydrogen can sustain large anisotropic stresses at megabar pressures, and the accompanying reversible deformations in the elastic region notably alter the crystal symmetry, charge distributions, and bonding behavior. The resulting band shifts, splittings, and enhanced electron-phonon coupling of the phase can dramatically exceed those induced by pure isotropic compression. The effects lead to band gap closure and ensuing superconductivity at significantly reduced compressive stresses compared to hydrostatic conditions. Most notable is shear deformation causing hydrogen to become metallic and superconducting at 250 GPa, which is about 120 GPa (i.e., >30%) lower than that obtained under hydrostatic conditions predicted at the same level of calculation. These results provide specific predictions that could be tested on hydrogen under extreme compression by additional stress conditions that are achievable in the laboratory. Since such stress anisotropy is in general highly sample and P-T path dependent, the present results may help to explain apparent discrepancies in previously reported experimental results, including variations in rate of bandgap closure and onset of electrical conductivity in solid molecular hydrogen (7, 8, 10–12). Furthermore, stress anisotropy in such compressed materials can be both controlled (e.g., by varying the geometry of thin samples and thermal annealing) and can be measured with various spectroscopic, diffraction, and imaging techniques (28, 29, 33).

We now comment on possible impact of quantum and thermal effects on the present findings. It has been proposed that the prototypical quantum crystal, low-density solid ^4He ,

may develop giant nonclassical plasticity (41), but it was also pointed out that the observed phenomenon may be primarily rooted in thermal rather than quantum effects (42), and a later study showed (43) that the stress responses are well described by classical mechanical theory. Moreover, strength enhancement of solid helium under pressure has been documented (44), and analogous results were observed in hydrogen (45). Experimental studies of other low Z systems, such as Li, LiH and LiF, document significant contributions of both thermal (e.g., Debye-Waller) effects and quantum zero-point effects on their lattice dynamics (46, 47). For phase III of hydrogen, anharmonic finite-temperature effects have been shown to reduce the electronic band gaps (48). The quantum nature of the protons and anharmonicity are also predicted to increase the intramolecular distance and enhance superconductivity in molecular hydrogen (49). Thermal effects tend to soften crystals and reduce compressive and shear strengths (50), but probably do not have a large influence on the structural and mechanical behavior at the pressure-temperature regimes of interest here (>200 GPa and <300 K).

In conclusion, this study opens a new avenue and perspective on the metallization and potential high- T_c superconductivity in dense hydrogen under stress conditions possible in high-pressure experiments, most notably static compression studies using diamond anvil cells. The results provide an instructive comparison to many benchmark theoretical studies of molecular hydrogen, i.e., as a classical crystal at the same level of theory under hydrostatic pressures (13, 21–25, 27, 37–39) versus the nonhydrostatic conditions considered in this work. While quantum and thermal effects may modify some predicted quantitative aspects, these effects are not expected to change the trends and alter our principal conclusions about the impact of anisotropic stresses on the electronic properties of hydrogen under pressure. Moreover, the findings have implications for other molecular solids and related systems, which are likely to sustain even higher shear stresses than dense solid hydrogen. These materials include hydride high T_c superconductors (34, 35, 51–54), whose stability ranges and properties may be significantly altered by anisotropic stresses.

Materials and Methods

We have performed the structural relaxation and stress-strain relation calculations using the Vienna Ab-initio Simulation Package (VASP) code, adopting the projector-augmented wave method and the GGA-PBE exchange-correlation functional (XC) with 900 eV cutoff energy and $8 \times 9 \times 10$ Monkhorst-Pack k -mesh. The stress-strain relations are determined with a strain increment of 0.01, where at each step the applied the shear or compressive strain is fixed to determine the shear stress σ_{xz} or compressive stress σ_{zz} , respectively, while the other five independent components of the strain tensors and all atomic positions are simultaneously relaxed. The vdW-DF2 XC has been used to calculate the bandgap E_g , electronic band structure, and Fermi surface of molecular hydrogen implemented in VASP code. We have used the quantum theory of atoms in molecules (QTAIM) analysis implemented in AIM-UC software. The lattice dynamics and electron-phonon coupling calculations were performed using the QUANTUM ESPRESSO code, adopting an energy cutoff of 80 Ry and an $8 \times 8 \times 8$ k mesh and $4 \times 4 \times 4$ q mesh (see SI Appendix Fig. S7 for the convergence test). The anisotropic superconducting gap and T_c are calculated using the open-source electron-phonon Wannier code. More computational details can be found in the Supporting Information.

Data Availability. All study data are included in this article and/or SI Appendix.

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