



Quasiperiodic Microjets Driven by Granular Advection as Observed With High-resolution Imaging at He I 10830 Å

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Abstract

With high-resolution narrowband He I 10830 Å filtergrams from Goode Solar Telescope, we give an extensive analysis for four granule-sized microeruptions which appear as the gentle ejection of material in He I 10830 Å band. The analysis was aided with the EUV data from Atmospheric Imaging Assembly and line-of-sight magnetograms from the Helioseismic and Magnetic Imager on board the Solar Dynamics Observatory. The microeruptions are situated on magnetic polarity inversion lines (PILs), and their roots are accurately traced down to intergranular lanes. Their durations are different: two microeruptions are repetitive microjets, lasting ~ 50 and 27 minutes respectively, while the other two events are singular, lasting ~ 5 minutes. For the two microjets, they are continuous and recurrent in the He I 10830 Å band, and the recurrence is quasiperiodic with a period of ~ 5 minutes. We found that only transient cospatial EUV brightenings are observed for the longer duration microjets and EUV brightenings are absent for the two singular microeruptions. What is essential to the longer duration microjets is that granules with the concentration of a positive magnetic field persistently transport the magnetic field to the PILs, canceling the opposite magnetic flux and making the base of the two microjets and the underlying granules migrate with the speed of ~ 0.25 and 1.0 km s^{-1} . The observations support the scenario of magnetic reconnection for the quasiperiodic microjets and further show that the reconnection continuously generates multitemperature components, especially the cool component with chromospheric temperature. In addition, the ongoing reconnection is modulated by p-mode oscillations inside the Sun.

Unified Astronomy Thesaurus concepts: [Solar mottles \(1074\)](#); [Solar magnetic reconnection \(1504\)](#)

Supporting material: animations

1. Introduction

Even in solar quiet regions, the upper atmosphere is a highly dynamic environment, full of microeruptions, as seen from high-resolution observations. A typical phenomenon is the upward ejection of material, presumably being ejected along magnetic field lines. They were first observed as surges with H α telescopes (Roy 1973). Surges are events of plasma ejection from the chromosphere into the corona, and they were found to be associated with EUV and soft X-ray emission (Schmahl 1981; Schmieder et al. 1988; Svestka et al. 1990; Shen 2021). With the ongoing improvement in spatial and temporal resolutions, and especially observations at multiple wavelengths, we have observed various facets of the upward material ejection in the solar upper atmosphere, and they were called X-ray jet (Shibata et al. 1992), EUV jet (Chae et al. 1999), and H α jet (Wang et al. 1998a; Chae et al. 1998) from the viewing point of a certain wavelength.

In terms of their locations, jets can be classified as polar jets (Wang et al. 1998b; Culhane et al. 2007; Sterling et al. 2015; Panesar et al. 2018; McGlasson et al. 2019), active region jets (Chifor et al. 2008), and quiet region jets (Panesar et al. 2016; McGlasson et al. 2019; Panesar et al. 2020). Some larger EUV/X-ray jets, especially polar jets, can escape into space and appear in the white-light coronagraph, and they are sometimes called coronal jets (St. Cyr et al. 1997; Paraschiv et al. 2015).

Recent observations show that a great number of coronal jets are made by minifilament eruptions that are prepared and triggered by flux cancellation at a magnetic neutral line, especially for those in quiet regions and coronal holes (Panesar et al. 2016, 2018; McGlasson et al. 2019; Panesar et al. 2020). For most solar eruptions that make a coronal jet, an essential component is a long spire that extends out along a far-reaching unipolar magnetic field rooted in an anemone-shaped base. The base depicts a closed-field magnetic configuration over a patch of opposite-polarity magnetic flux. The spire is evidently made by external reconnection of a lobe of the jet's magnetic-anemone base with the ambient oppositely directed far-reaching field. On the other hand, spicules are one class of jet phenomena on the solar limb traditionally observed in narrowband H α images (Beckers 1968), and they may appear as dark mottles, or fibrils (including dynamic fibrils) on the disk (Tsiropoula et al. 1994; Suematsu et al. 1995; Christopoulou et al. 2001). De Pontieu et al. (2007b) found that there are two fundamentally different types of spicules at the limb, type I and type II. The type I spicules evolve on timescales of several minutes while type II spicules occur on timescales of a few tens of seconds during which they are seen to rise and then rapidly disappear. Type I spicules seem to correspond to dynamic fibrils (active regions) and a subset of mottles (quiet Sun), when seen on the solar disk (Hansteen et al. 2006; De Pontieu et al. 2007a; Rouppe van der Voort et al. 2007). Short-lived type II spicules may correspond to rapid blueshifted events on the disk (Langangen et al. 2008). Spicules are also thought to play an important role in heating the upper atmosphere



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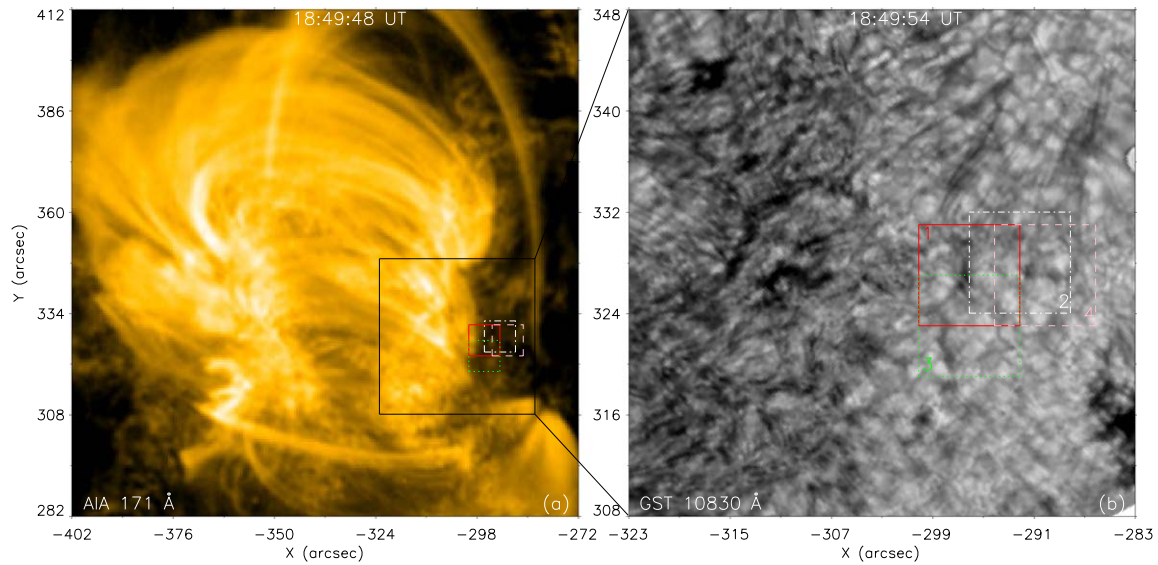


Figure 1. Panel (a) depicts an overview of the NOAA AR 11259 as observed by AIA at 171 Å, while panel (b) is a sample He I 10830 Å narrowband filtergram which FOV is given in panel (a) with a black box. Four smaller boxes with different line styles and colors in panels (a) and (b) indicate the places where the four microeruptions are from. Each number corresponds to the number given to each microeruption in this paper.

(Samanta et al. 2019) and many small-scale jets are sources of mass and energy of solar wind (Tian et al. 2014).

Jet activities, either hot or cool, tend to be recurrent from a small localized region (Chae et al. 1999; Jiang et al. 2007; Chifor et al. 2008; Guo et al. 2013; Chandra et al. 2015; Chen et al. 2015). Some authors even proposed a periodic nature for them. Jiang et al. (2007) reported three surges with a recurrent period of about 70 minutes in H α , EUV, and soft X-rays. With wavelet analysis, Li et al. (2015) obtained the periods of 5 and 13 minutes from intensity variations at the base of a series of recurrent EUV jets. With high-resolution narrowband imaging at He I 10830 Å, Wang et al. (2021) reported periodic behavior of a ~ 5 minute period for a small-scale jet. However, in the EUV band and for the same jet, the periodic behavior vanishes and appears as nonperiodic recurrent brightenings. In this regard, high-resolution narrowband imaging at He I 10830 Å is a relatively new window for various jets in the chromosphere.

With the 1.6 m aperture Goode Solar Telescope (GST) at Big Bear Solar Observatory (BBSO), a new facet of solar upper chromosphere was obtained using high-resolution narrowband imaging at He I 10830 Å. Meanwhile, the first direct observations of ultrafine jet-like dynamical events originating in solar intergranular lanes and subsequently lighting up the corona were observed (Ji et al. 2012). With the same set of data, further interesting results have been achieved. For completeness of this paper, they are listed as follows:

1. Zeng et al. (2013) reported the association of a He I 10830 Å surge with the magnetic reconnection driven by advection of a rapidly developing large granule.
2. In a quiet region, Hong et al. (2017) found that EUV emission is mainly concentrated in intergranular lane areas where He I 10830 Å absorption is enhanced.
3. In a EUV moss region, numerous small-scale material ejections shown as enhanced He I 10830 Å absorption are squirted from intergranular lanes periodically for a period of ~ 5 minutes. Furthermore, periodic material ejection was found to be associated with periodic EUV brightenings and magnetoacoustic oscillations with roughly the same period (Hashim et al. 2021; Ji et al. 2021).

Therefore, we need to know how the scale size spectrum of coronal jets extends down to these small-scale subarcsecond material ejections when high-resolution observations are available (Raouafi et al. 2016; Hou et al. 2021). We also need to know the behaviors of all the above-mentioned microeruptions in high-resolution He I 10830 Å imaging observations. In this paper, we extensively analyzed four microeruptions that occurred at four different places with the same set of data. The sizes of these He I 10830 Å microeruptions are comparable with granule sizes. With simultaneous and well-aligned high-resolution images of the photosphere, we are able to pin down these microeruptions in the context of advecting granules and the corresponding magnetic field.

2. Observations

On 2011 July 22, the NOAA Active Region 11259 was observed with narrowband (bandpass: 0.5 Å) imaging at the blue wing (-0.25 Å) of He I 10830 Å and broadband imaging (bandpass: 10 Å) centered at TiO 7057 Å line using the GST and its facilities at the BBSO (Cao et al. 2010; Goode et al. 2010). It is the same set of data analyzed by Ji et al. (2012). The images of He I 10830 Å have a pixel size of $\sim 0''.09$ and a cadence of ~ 14 s. For TiO 7057 Å images, they have a pixel size of $\sim 0''.036$ and a cadence of ~ 15 s. The EUV/UV imaging data taken by the Atmospheric Imaging Assembly (AIA; Lemen et al. 2012) and line-of-sight (LOS) magnetograms taken by the Helioseismic and Magnetic Imager (HMI; Scherrer et al. 2012) on board the Solar Dynamics Observatory (SDO) are used to show the hot components of these microjets and relevant variations of the magnetic field.

The smallest events analyzed in this paper have a footpoint size of $\sim 0''.5$. To understand their physical originations, we have to pin down their locations on the high-resolution photospheric images in the context of granules and intergranular lanes, and a precise coalignment between the images of the photosphere and chromosphere is very necessary. He I 10830 Å is a shallow chromospheric line that allows the photosphere to be visible at most places of the high-resolution narrowband images as a background. This attribute makes

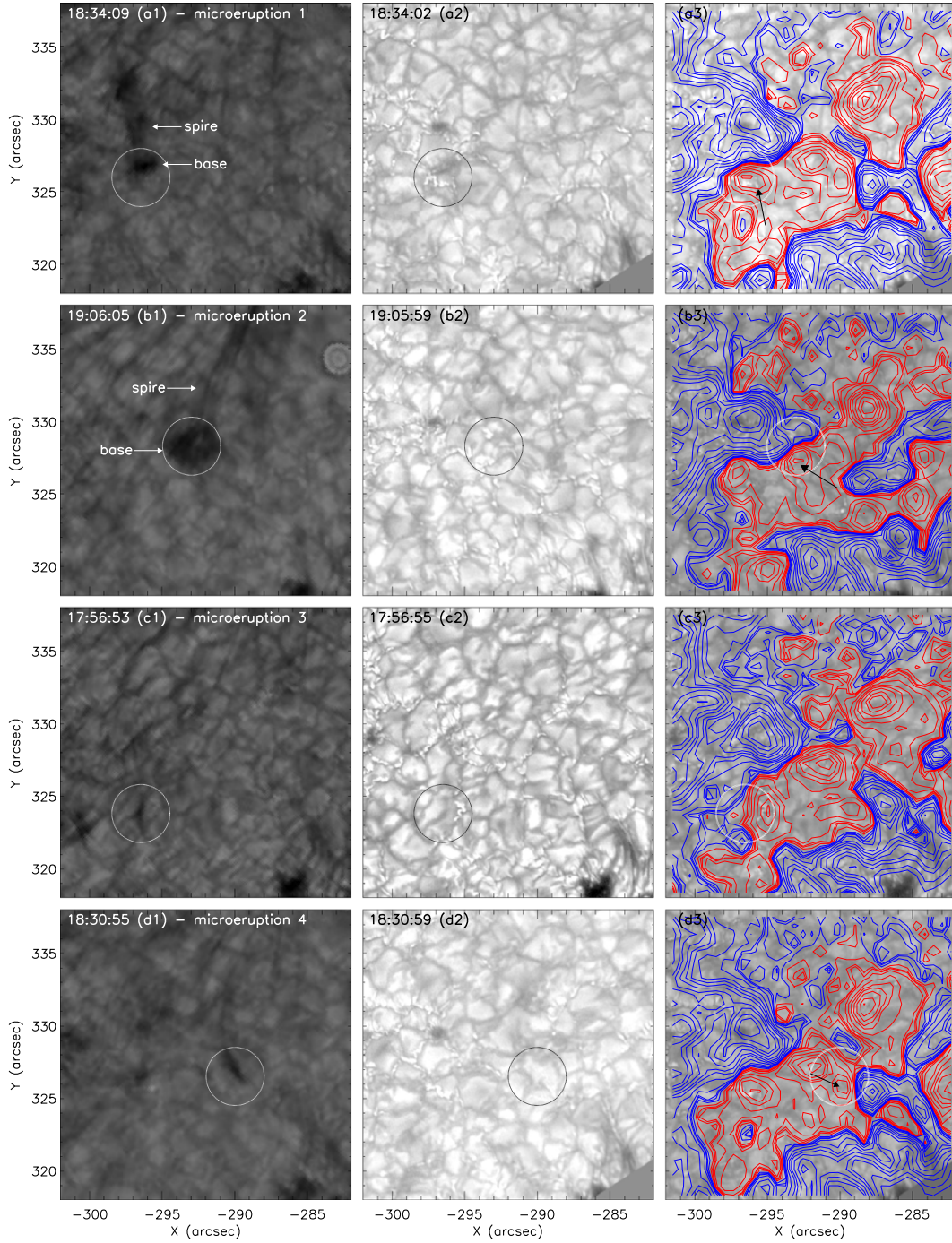


Figure 2. The left column gives four separate snapshots for the four microeruptions (with two microjets in panels (a1) and (b1)), as observed with He I 10830 Å narrowband filtergrams. The middle and right columns are corresponding broadband filtergrams at TiO 7057 Å lines and the same filtergrams overlaid with contours of the LOS magnetic field. The red and blue contours in the magnetograms represent positive and negative LOS magnetic field with the levels of ± 7 , ± 9 , ± 15 , ± 30 , ± 50 , ± 70 , ± 100 , ± 200 , ± 300 , ± 400 G, respectively. The circles in each row have the same position on the maps and are used to relate the four microeruptions with the context of surrounding granules and the magnetic field of the photosphere. The two arrows in panels (a3) and (b3) point to the moving magnetic elements being associated with microjets 1 and 2. The arrow in panel (d3) points to the emerging magnetic element being associated with microeruption 4. An online animation of the figure is available; it lasts for 103 minutes from $\sim 17:35$ UT to 19:18 UT.

(An animation of this figure is available.)

aligning those images with the images of TiO 7057 Å much easier. Meanwhile, the coalignment between GST and SDO data was carried out using sunspots, pores, and bright granules on TiO 7057 Å broadband images and HMI continuum images as the intermediary features. The remaining offset after the coalignment is less than $\sim 0''.5$ (Hong et al. 2017).

3. Results

Figure 1(a) gives an overview of NOAA AR 11259 as observed by AIA 171 Å, with the field of view (FOV) of $130'' \times 130''$. The black box corresponds to the area of the panel (b) with a FOV of $40'' \times 40''$, where a sample He I 10830 Å image is given. The small boxes of different colors

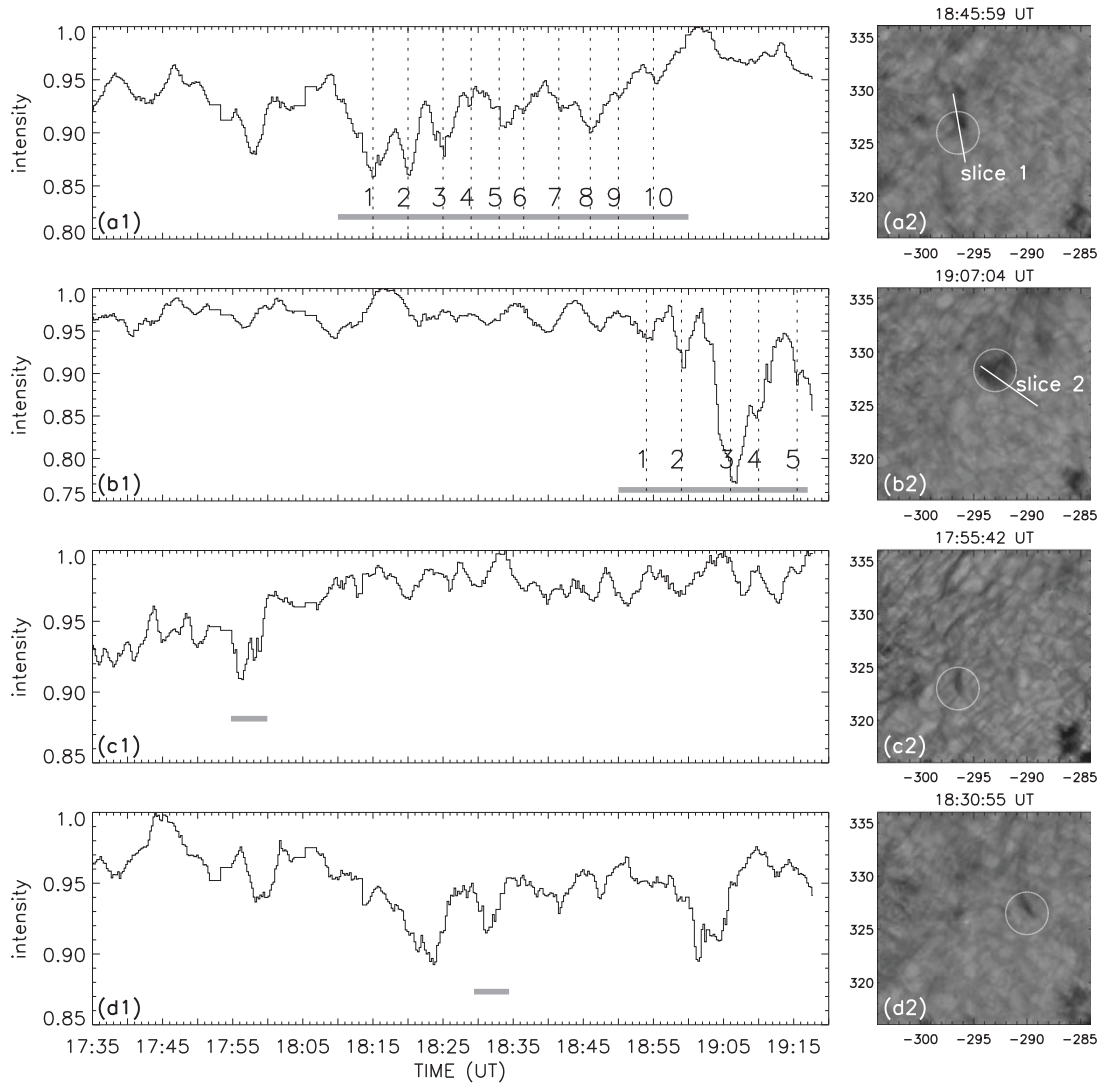


Figure 3. The time profiles (normalized) for He I 10830 Å brightness integrated from the circled areas given on the right column. The gray horizontal lines show the durations of the four microeruptions. Dashed lines in panels (a1) and (b2) give quasiperiodic enhancement of He I 10830 Å absorption during the eruption processes of microjets 1 and 2. Slices 1 and 2 indicate the positions used to make the time–distance diagrams in Figures 6 and 7. An online animation of the figure is available; it lasts for 103 minutes from $\sim 17:35$ UT to 19:18 UT.

(An animation of this figure is available.)

and numbers in panel (b) give the locations of the four microeruptions. It is worthy of noting that a moss region is on the left side of the picture, where many periodic small-scale mass ejections have been studied and they are reported to be powered by magnetoacoustic oscillations (Hashim et al. 2021; Ji et al. 2021). Figure 2 along with its online animation gives detailed information of the four microeruptions on the He I 10830 Å images and their positions (circled area) in the context of granules and LOS magnetograms on the photosphere. The images in the second column show that the locations of the microeruptions are subject to constant advection of granules and have no special appearance when compared with other places. However, the third column makes it apparent that all the microeruptions are located around magnetic polarity inversion lines (PILs). Furthermore, with the help of animation for Figure 2, we can see that all microeruptions squirt out their material from the area of intergranular lanes. Spires and bases in microeruptions 1 and 2 are occasionally visible (as shown in

panels (a1) and (b1)), so the two eruptions can be taken as microjets.

The temporal behaviors for the He I 10830 Å intensity in the circled areas of the two microjets and two microeruptions are given on the left column of Figure 3. On the right column, four snapshots for the two microjets and two microeruptions are given. In this way, an online animation for Figure 3 will enable us to carry out a detailed comparison between their temporal variation and spatial evolution. The gray horizontal lines indicate the duration of the four microeruptions. They can be classified into two categories: long-duration microjets and short-term microeruptions. The two long-duration microjets lasted ~ 50 and 27 minutes respectively. During their existence, they burst periodically for a period of ~ 5 minutes, as can be seen from the oscillations imposed on the light curves. Certainly, the oscillations of the light curves will have the component of the p-mode oscillations of 5 minutes on the photosphere, since He I 10830 Å images contain information from the underlying photosphere. However, when we visually

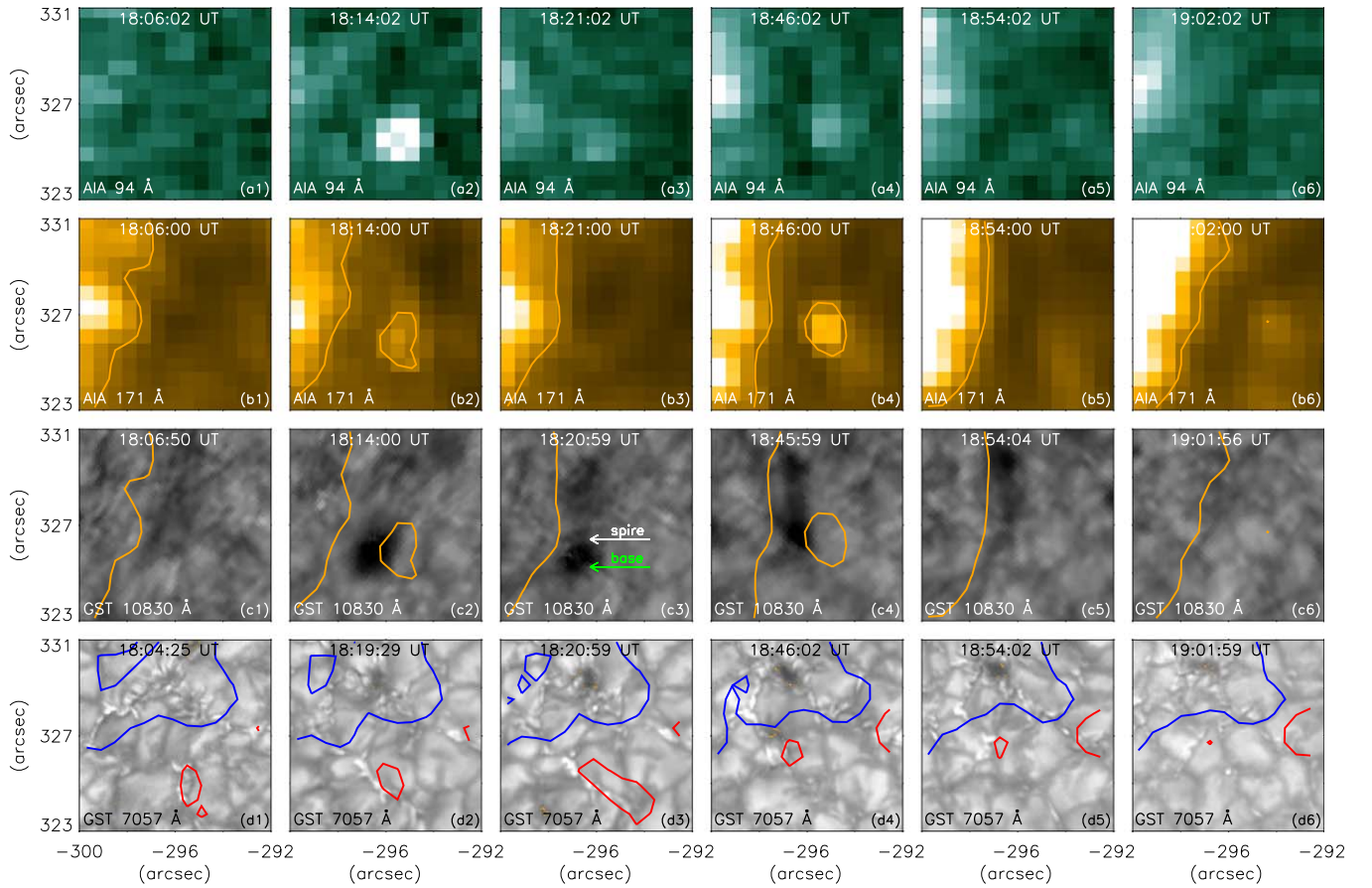


Figure 4. Panels (a1)–(a6) and (b1)–(b6) in the top two rows give the observations for microjet 1 made by AIA at 94 Å and 171 Å. Panels (c1)–(c6) in the third row show the repeated eruptions of microjet 1 as observed with He I 10830 Å filtergrams, overlaid with simultaneous 171 Å contours which have the value of 285 DN pixel^{−1} s^{−1}. The bottom row gives the underlying photosphere as observed with 7057 Å filtergrams. The overlaid red and blue contours in the bottom row represent the LOS magnetic field with the values of +70 and −70 G. White and green arrows in panel (c3) point out the spire and base of the microjet, respectively.

check the movie made for Figure 3, we can see that each burst in the circled area is nicely associated with each valley period on the light curves. In this sense, microeruptions 3 and 4 appear to be singular and last only one period (~ 5 minutes). The light curves for microeruptions 3 and 4 reflect the variation of He I 10830 Å absorption modulated by the p-mode oscillations of 5 minutes on the photosphere.

We compare the spatiotemporal evolution of magnetograms with the animation of Figure 2 and see that the configurations of the PIL associated with microjets 1 and 2 are more stable. It is long and exists there throughout the entire period of the events. There are two moving magnetic elements that continuously push and squeeze the PIL northward (indicated with arrows in panels (a3) and (b3)). The pushing movement continuously squeezes and distorts the adjoining negative magnetic elements.

Seeking their detailed dynamic behaviors associated with the magnetic field and granules, we will focus on the base of the two microjets. Figures 4 and 5 give their temporal evolution as observed in AIA 94 Å and 171 Å, He I 10830 Å and TiO 7057 Å and LOS magnetic field. Microjet 1 has a full size of $\sim 3'' \times 1''$ with a base area of $\sim 1'' \times 1''$. The He I 10830 Å filtergrams in panels (c1) to (c6) in Figure 4 again show that the microjet is recurrent from its initial generating to final fading. Combining with the TiO 7057 Å photosphere images and the contours of the magnetic field (panels (d1)–(d6)), it can be found that the positive magnetic field concentration (red

contours) is decreasing during the whole period of microjet 1. The magnetic concentration is mainly situated in the intergranular lanes and associated with the bright points, which are the important signatures for the stronger magnetic field. Meanwhile, the positive magnetic field and the bright points in intergranular lanes move toward the negative magnetic field (blue contours). The base of the microjet is also moving accordingly and produces several cospatial EUV brightenings at multiple temperatures (panels (a2), (a4), (b2), and (b4)). EUV brightenings are observed as recurrent at the same place. As the microjet gradually disappears, the magnetic field at its root gradually weakens, and the bright points in the lane areas also gradually fade away (see panels (c6) and (d6) of Figure 4).

Through a slice cutting across microjet 1 (slice 1 in Figure 3), we have made its time–distance diagram that is shown in Figure 6. The time span for Figure 6 is from 18:06:00 to 19:02:00 UT (we have no TiO 7057 Å data from 18:04:25 to 18:19:29 UT). From the time–distance image of He I 10830 Å, it can be seen that the microjet squirts its material continuously from $\sim 18:10$ to 19:00 UT. The ejection is quite gentle with a speed of less than 10 km s^{-1} (green lines in Figure 6). Its base moves in accordance with the advection of a corresponding granule. The granules push the intergranular lane moving northward toward the site of negative polarity, as can be seen from the TiO 7057 Å time–distance image. The movement speed of the granules is estimated as 0.25 km s^{-1} . During this process, several patches of recurrent brightenings are shown in

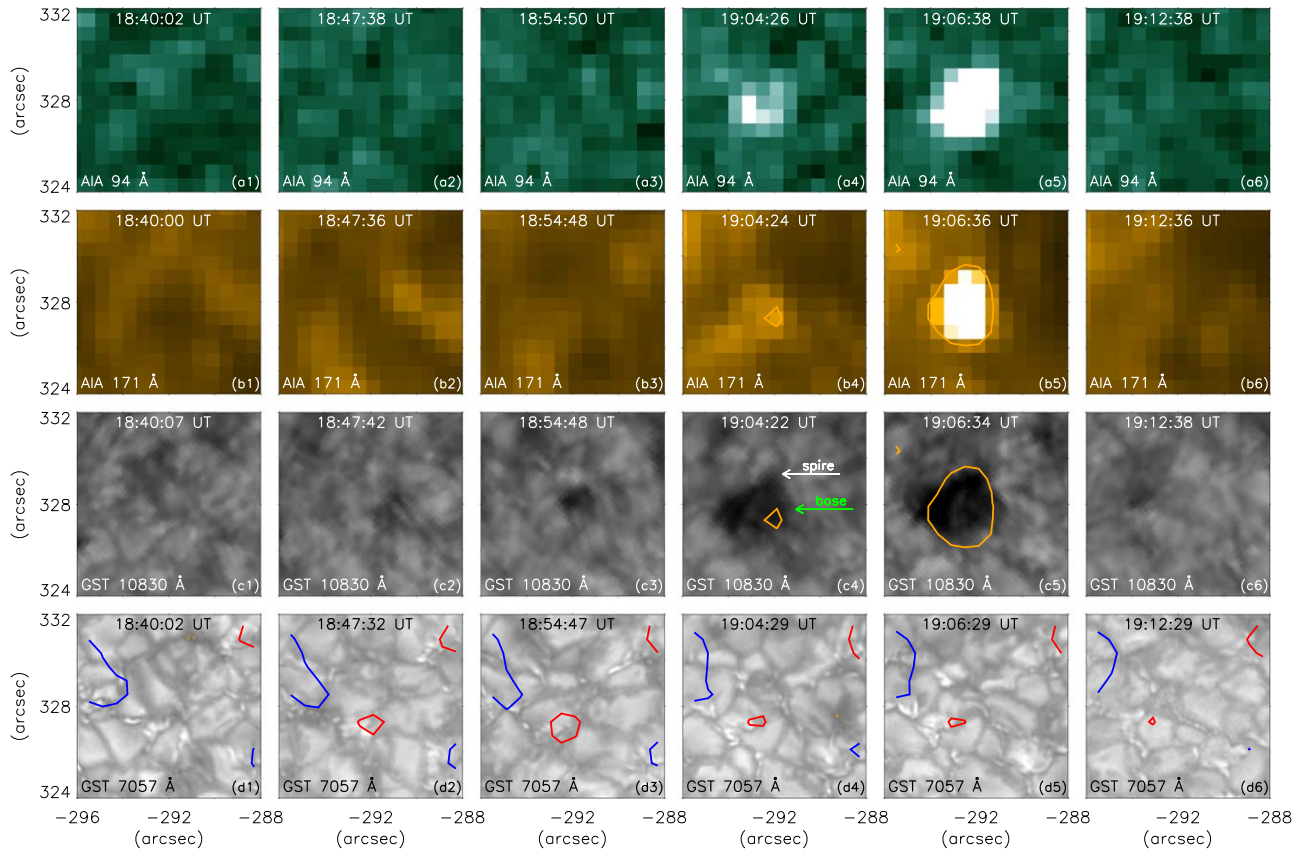


Figure 5. Panels (a1)–(a6) and (b1)–(b6) in the top two rows give the observations for microjet 2 made by AIA at 94 and 171 Å. Panels (c1)–(c6) in the third row show the repeated eruptions of microjet 2 as observed with He I 10830 Å filtergrams, overlaid with simultaneous 171 Å contours of the value of 285 DN pixel^{−1} s^{−1}. The bottom row gives the underlying photosphere as observed with 7057 Å filtergrams. The overlaid red and blue contours in the bottom row represent the LOS magnetic field with the values of +150 and −150 G. The spire and base of the microjet are pointed by white and green arrows in panel (c4), respectively.

the time–distance images of AIA 94 Å and 171 Å. Thus, the generating process of the microjet is well associated with the convective granules which squeeze and cancel the magnetic field of opposite polarity. During this process, the concentration of the magnetic field in the convective granules gradually decreases and the bright points gradually disappear.

The second microjet has a base size of $\sim 3'' \times 3''$. We carried a similar analysis as for the first microjet, and the results are given in Figures 5 and 7 which show that the microjet is associated with the movement of the positive magnetic element. Cospacial multitemperature EUV brightenings at the base occur when the He I 10830 Å absorption of the microjet gets its maximum phase. Through the TiO time–distance image in Figure 5, we again see that the movement of the magnetic field toward opposite polarity coincides with the movement of the intergranular lanes. This behavior is highly similar to that of microjet 1, except that there are no obvious intragranular bright points.

Microeruptions 3 and 4 are singular, lasting ~ 5 minutes. They are tinier with a width well below one arcsec (Table 1). During their appearance, there are no detectable EUV brightenings. Also, there are no detectable magnetic field changes associated with microeruption 3. However, for microeruption 4, it is apparently associated with the emergence of a positive magnetic element (pointed by an arrow in panel d3 of Figure 2). Furthermore, they are still rooted in intergranular lanes and magnetic PILs, being consistent with the corresponding results for microjets 1 and 2.

4. Discussions and Conclusions

With high-resolution He I 10830 Å narrowband images from GST, we give an extensive analysis of four granule-sized microeruptions, two of which are identified as microjets. They all appear as enhanced He I 10830 Å absorption. The investigation of the microjets using high-resolution He I 10830 Å narrowband images represents an extension for the current scale size spectrum of coronal jets down to the smallest spatial scales. We have seen that high-resolution He I 10830 Å imaging is an indispensable observation for various jet phenomena in the chromosphere. We can trace their roots directly to the underlying photosphere and its magnetic field. With the help from the He I 10830 Å data and simultaneous high-resolution images of the photosphere, we report that these microeruptions are rooted in the area of intergranular lanes, which confirms previous results (Ji et al. 2012; Zeng et al. 2013; Wang et al. 2021). Furthermore, they are situated on magnetic PILs.

For the two microjets, we also find that they are caused by advection of granules which transport concentration of positive magnetic field to a relatively stable and long magnetic PIL. The moving positive magnetic field element persistently pushes and distorts the PIL, canceling the magnetic flux of the opposite magnetic field. In quiet regions and coronal holes, Panesar et al. (2016, 2018, 2020) and McGlasson et al. (2019) have found that magnetic flux cancellation at a PIL often makes a PIL-tracing minifilament and triggers its eruption. The minifilament eruption drives external reconnection of the minifilament's

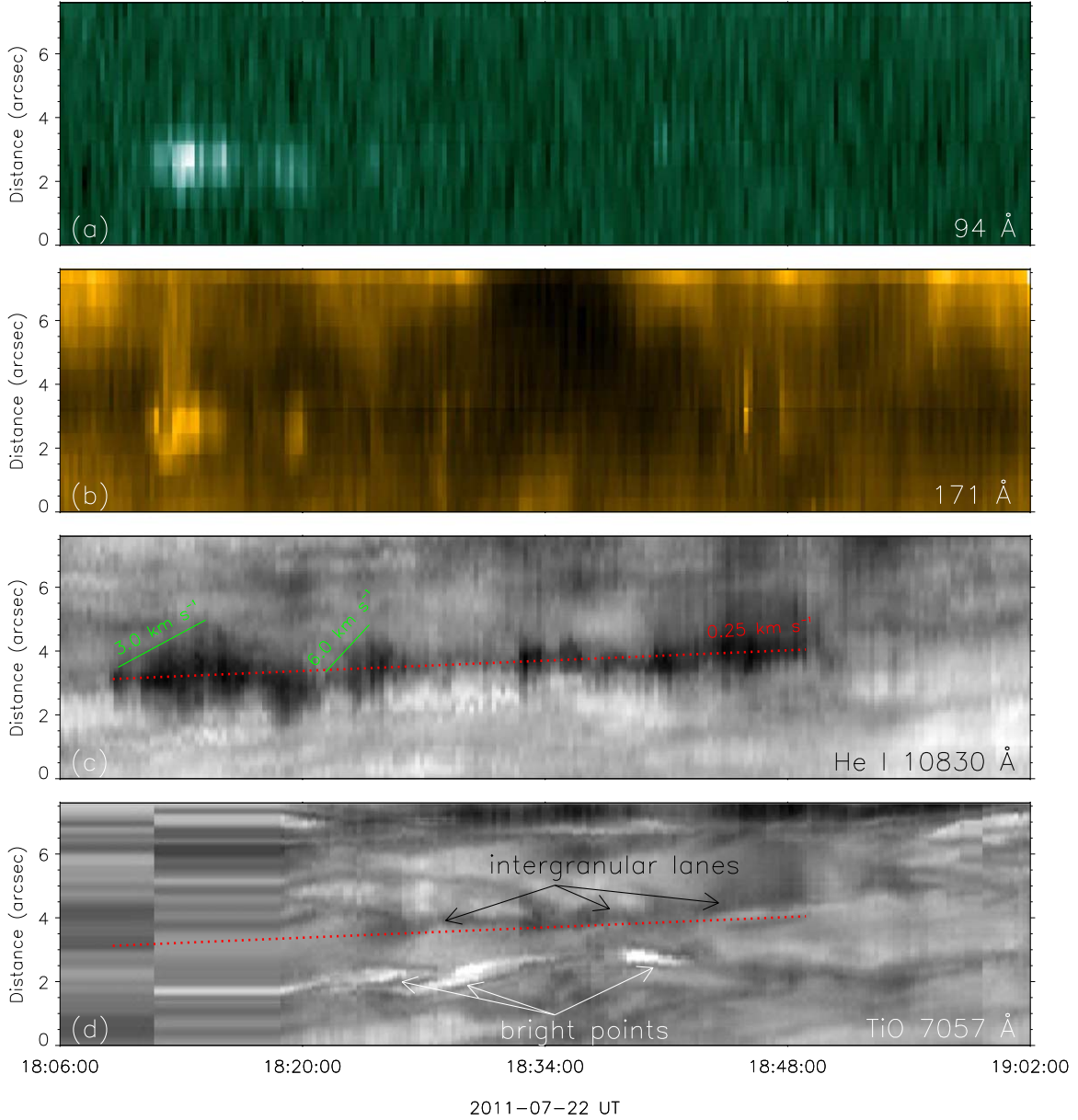


Figure 6. From top to bottom row are the time–distance diagrams at 94, 171, 10830, and 7057 Å for microjet 1, which are obtained by stacking intensity images of slice 1 (Figure 3) of corresponding wavelengths. The two green lines in the third panel give fittings for the speeds of microjet 1 (3.0 km s^{-1} and 6.0 km s^{-1} , respectively), while the red dotted line in the panel gives a fitting to the migration speed (0.25 km s^{-1}) for the base of the microjet. The same red dotted line is overplotted on the fourth panel, showing that the moving speed of the granules is equivalent to the base of the microjet. The dark arrows point out the intergranular lanes and the white arrows point to the bright points.

enveloping magnetic lobe with the far-reaching magnetic field to produce a coronal jet. This phenomenon and model were first discovered and proposed by Sterling et al. (2015). Meanwhile, microeruptions 3 and 4 were more likely confined minifilament eruptions that did not make jets. For a microeruption such as any of those in the present paper, a dark microeruption with or without a discernible microjet spire in He I 10830 Å images, high-resolution $H\alpha$ or EUV images along with high-resolution magnetograms are needed to decide whether the microeruption is the product of a microfilament eruption. Also, many other works have been reported regarding how magnetic flux emergences and cancellations induce various recurring jets (e.g., Chae et al. 1999; Chifor et al. 2008; Zeng et al. 2013; Zhang & Ji 2014; Liu et al. 2016; for a more comprehensive review, see Raouafi et al. 2016).

Therefore, our observation supports the magnetic reconnection scenario for the two microjets.

Coronal jets tend to be recurrent (Jiang et al. 2007; Guo et al. 2013; Cheung et al. 2015; Joshi et al. 2017). In most cases, they are not periodic. Nevertheless, Ning et al. (2004) and Doyle et al. (2006) reported a quasi-period of 3 to 5 minutes for repetitive transition region explosive events. Li et al. (2015) obtained the periods of 5 and 13 minutes from intensity variations at the base of a series of recurrent EUV jets. With high-resolution narrowband imaging at He I 10830 Å, Wang et al. (2021) reported quasiperiodic behavior of a ~ 5 minute period for a small-scale jet. However, in the EUV passbands, the jet only appears as recurrent with no periodicity. Similarly, the long-duration microjets in He I 10830 Å given in this paper are found to be quasiperiodic with a period of ~ 5 minutes. The

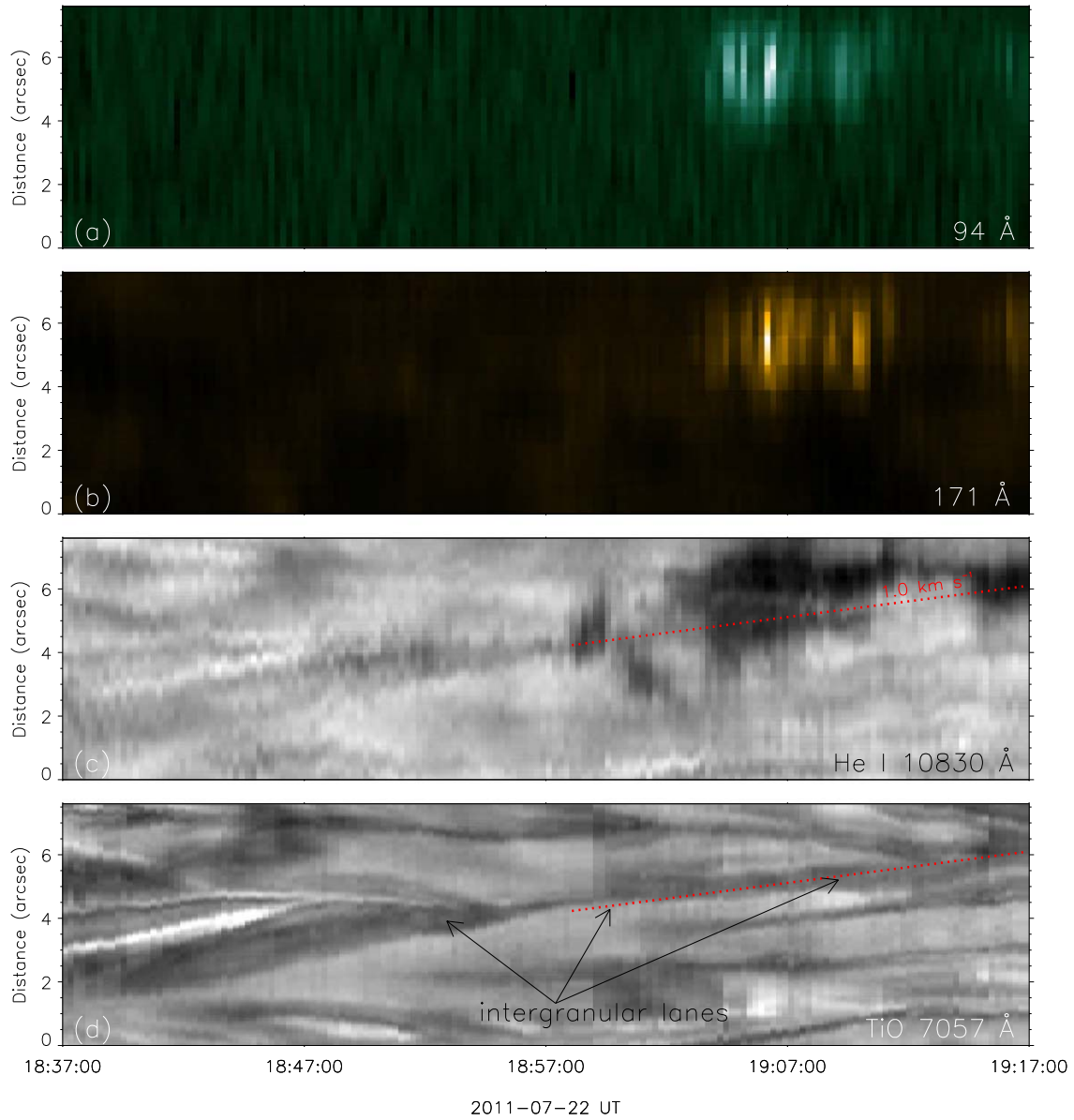


Figure 7. From top to bottom row are the spacetime diagrams at 94, 171, 10830, and 7057 Å for microjet 2, which are obtained by stacking intensity images of slice 2 (Figure 3) of corresponding wavelengths. The red dotted line in the third panel gives a fitting to the migration speed (1.0 km s^{-1}) for the base of the microjet. The same red dotted line is overlotted on the fourth panel, showing that the moving speed is equivalent to the intergranular lanes which are pointed by dark arrows.

Table 1
Parameters of Microeruptions

Microeruptions	Length $\times$ Width (arcsec)	Start Time UT	End Time UT	Duration (min)
1	$\sim 3 \times 1$	$\sim 18:10:00$	$\sim 19:00:00$	~ 50
2	$\sim 3 \times 3$	$\sim 18:50:00$	$\sim 19:17:00$	~ 27
3	$\sim 3 \times 0.5$	$\sim 17:54:45$	$\sim 17:59:56$	~ 5
4	$\sim 3 \times 0.5$	$\sim 18:29:25$	$\sim 18:34:25$	~ 5

material ejections, which are supposed to be in cool chromosphere temperature, are quasiperiodic being occasionally accompanied by nonperiodic yet recurrent multitemperature brightenings in the EUV passbands. From Figures 6 and 7, we can see that EUV brightenings are associated with the intenser cool material ejections, which may indicate a relatively stronger magnetic reconnection. Due to the constant advection

of positive magnetic elements, the cool ejections are actually continuously modulated by the 5 minute oscillations, as can be seen from the two time–distance diagrams (Figures 6 and 7). The observations strongly support a slow magnetic reconnection that is continuously driven by constant advection of granules and the slow magnetic reconnection is modulated by p-mode oscillation.

To explain the quasiperiodic behavior of repetitive transition region explosive events, Chen & Priest (2006) carried out MHD simulations by imposing p-mode oscillations at the bottom boundary. The simulations successfully made the reconnection in the transition region proceed periodically. However, for the microjets reported in this paper, magnetic reconnection is more likely to occur in the solar lower atmosphere. Simulations of magnetic reconnection in the solar lower atmosphere including the modulation of p-mode oscillations are needed. Especially, the magnetic reconnection should be able to produce multithermal components, which

span a large temperature range from 10^4 to 10^6 K, like the simulation given by Ni et al. (2021). Numerical simulations show that, for magnetic reconnection in the lower chromosphere, hot plasma with a temperature of more than 2.0×10^4 K can be produced when plasma β becomes relatively lower (Cheng et al. 2021; Ni et al. 2021). In this way, we can understand why the two singular microeruptions have no EUV counterparts. In the end, we also would like to conclude that, to understand the real nature of these He I 10830 Å microjets, we need further investigations to reveal their relationship with mottles and dynamic fibrils, which have been extensively observed and studied.

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