

1 **The Influence of Boundary Layer Mixing Strength on the**
2 **Evolution of a Baroclinic Cyclone**

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ABSTRACT

7 Sub-grid scale turbulence in numerical weather prediction models is typically handled by
8 a PBL parameterization. These schemes attempt to represent turbulent mixing processes
9 occurring below the resolvable scale of the model grid in the vertical direction, and act upon
10 temperature, moisture, and momentum within the boundary layer. This study varies the
11 PBL mixing strength within 4-km WRF simulations of the 26–29 January 2015 snowstorm
12 to assess the sensitivity of baroclinic cyclones to eddy diffusivity intensity. The bulk critical
13 Richardson number for unstable regimes is varied between 0.0–0.25 within the YSU PBL
14 scheme, as a way of directly altering the depth and magnitude of sub-grid scale turbulent
15 mixing. Results suggest varying the bulk critical Richardson number is similar to select-
16 ing a different PBL parameterization. Differences in boundary layer moisture availability,
17 arising from reduced entrainment of dry, free tropospheric air, lead to variations in the
18 magnitude of latent heat release above the warm frontal region, producing stronger upper-
19 tropospheric downstream ridging in simulations with less PBL mixing. The more amplified
20 flow pattern impedes the northeastward propagation of the surface cyclone and results in a
21 westward shift of precipitation. Additionally, trajectory analysis indicates ascending parcels
22 in the less-mixing simulations condense more water vapor and terminate at a higher po-
23 tential temperature level than ascending parcels in the more-mixing simulations, suggesting
24 stronger latent heat release when PBL mixing is reduced. These results suggest spread
25 within ensemble forecast systems may be improved by perturbing PBL mixing parameters
26 that are not well constrained.

27 1. Introduction

28 The accurate representation of turbulent mixing processes within the planetary boundary
29 layer (PBL) is an important component of numerical weather prediction over a variety of
30 spatial and temporal scales. Several studies have investigated the sensitivity of atmospheric
31 phenomena, including severe weather outbreaks, tropical cyclones (TCs), and baroclinic
32 waves, to PBL and surface layer parameterization schemes and configurations (Adamson
33 et al. 2006; Beare 2007; Plant and Belcher 2007; Kepert 2012; Boutle et al. 2014; Cohen
34 et al. 2015, 2017; Bu et al. 2017). The demonstrated sensitivity to PBL and surface layer
35 processes inspires the current work of identifying pathways through which subgrid-scale mix-
36 ing projects on larger-scale features within full-physics simulations. We seek to understand
37 how and to what degree varying boundary layer mixing strength impacts the development
38 and evolution of an extratropical cyclone using Weather Research and Forecasting (WRF)
39 simulations of the 26–29 January 2015 eastern U.S. snowstorm.

40 It has been well established that boundary and surface layer processes play an important
41 role in baroclinic cyclone evolution (Valdes and Hoskins 1988; Adamson et al. 2006; Beare
42 2007; Plant and Belcher 2007; Boutle et al. 2014). Valdes and Hoskins (1988) investigated
43 the baroclinic instability of a zonal-mean flow and found surface friction can reduce growth
44 rates of baroclinic systems by 50%. Adamson et al. (2006) described how surface friction can
45 influence and dampen baroclinic systems through Ekman pumping and baroclinic potential
46 vorticity (PV) generation. Plant and Belcher (2007) built on the results of Adamson et al.
47 (2006), analyzing the impacts of surface momentum and heat fluxes on the evolution of
48 baroclinic cyclones. The basic features of the PBL-driven damping mechanisms were found
49 to be robust for different frontal structures, surface heat fluxes, and for a range of surface

roughnesses. Beare (2007) and Boutle et al. (2014) performed a series of experiments selectively switching off boundary layer mixing over stable and unstable surface layers within an idealized cyclone. This was achieved by determining, at each timestep, whether the sign of the surface buoyancy flux was negative (stable) or positive (unstable). Their results suggested Ekman pumping and baroclinic PV generation were mostly associated with unstable and stable boundary layers, respectively, and both mechanisms contributed about equally to the dampening of the baroclinic wave.

The impacts of PBL momentum and heat fluxes have been the focus of several baroclinic cyclone studies but the influence of PBL moisture fluxes on baroclinic cyclones has received less attention. Considering the role of condensational heating in extratropical cyclogenesis has been well documented (e.g. Reed et al. 1988; Davis and Emanuel 1988; Kuo et al. 1990; Davis 1992; Davis et al. 1993; Whitaker and Davis 1994; Stoelinga 1996; Brennan and Lackmann 2005), it stands to reason moisture fluxes from the PBL may have a considerable impact on cyclone evolution. Hong and Pan (1996) and Hong et al. (2006) demonstrated changes in PBL mixing strength can impact CAPE and equivalent potential temperature; where shallower, less-intense PBL mixing traps moisture in the lower levels, increasing CAPE and low-level equivalent potential temperature. It is plausible that changes in the mixing of moisture via the PBL may impact the evolution of baroclinic cyclones through augmenting CAPE and low-level equivalent potential temperature fields as latent heat release, particularly around the warm conveyor belt (Carlson 1980; Wernli and Davies 1997; Schemm and Wernli 2014) can substantially impact extratropical cyclogenesis (Reed et al. 1988; Kuo et al. 1990; Stoelinga 1996; Ahmadi-Givi et al. 2004; Joos and Wernli 2012; Binder et al. 2016).

Inspired by the prior literature, one can surmise baroclinic cyclone evolution may be sensitive to the selection of the PBL parameterization scheme. While several studies have

74 addressed the impact of parameterized PBL mixing on baroclinic cyclones, the authors have
75 found no previous literature investigating how variations in PBL mixing strength project on
76 to the baroclinic-cyclone-scale flow, particularly within moist simulations. Our motivations
77 to address this are twofold. First, there is an important concern whether baroclinic cyclone
78 evolution is sensitive to reasonable changes in PBL mixing strength. Second, the impact
79 of PBL mixing strength on the distribution of moisture is an important consideration for
80 determining convective activity in the outer core of TCs (Bu et al. 2017). Therefore, it
81 is hypothesized that latent heating differences arising from variations in low-level moisture
82 distributions, coincident with changes in PBL mixing, can impact the evolution of baroclinic
83 cyclones through diabatic ridge building (Stoelinga 1996). Under these circumstances, this
84 work may help to improve forecast skill through revealing sources of variability in simulated
85 cyclogenesis arising from uncertainty in PBL mixing parameters.

86 In this paper, we investigate how variations in parameterized PBL mixing can impact the
87 evolution of a moist baroclinic cyclone described in Section 2. WRF PBL parameterization
88 schemes and simulation parameters are also discussed in Section 2. Section 3 presents
89 results from various sensitivity experiments in an Eulerian framework and suggests a primary
90 mechanism by which variations in PBL mixing strength impact a baroclinic cyclone. The
91 various experiments are summarized in Table 1 and each simulation is further explained
92 within the text below. Trajectory analyses are presented in Section 4 to further expand
93 upon the findings in Section 3 and a discussion and summary are presented in Section 5.

94 2. Background and methods

95 a. PBL schemes in the Weather Research and Forecasting model

96 PBL parameterizations seek to replicate the effects of subgrid-scale turbulent mixing of
97 heat, momentum, and moisture by calculating vertical diffusion coefficients that are applied
98 to predictive equations for momentum and scalars, such as temperature and moisture, during
99 model runtime. At this writing, the vast majority of PBL schemes are one-dimensional,
100 acting on model columns individually and independently of neighboring columns. There are
101 several PBL mixing parameterization strategies used in numerical weather prediction models
102 and many schemes rely on the concept of eddy diffusivity to find the turbulent vertical flux
103 of a quantity (e.g., heat, momentum, and moisture). Adapting Eq. (3.1) from Holtslag and
104 Boville (1993), vertical turbulent flux can be written as

$$\overline{w'C'} = K_c \frac{\partial C}{\partial z} + N_{NL}, \quad (1)$$

105 where $C \in (q, \theta, u, v)$, K_c is the eddy diffusivity for C , and N_{NL} represents nonlocal terms
106 that vary among schemes.

107 PBL parameterizations can be binned in one of two groups depending on the strategy
108 used to close the turbulence equations and obtain K_c : turbulent kinetic energy (TKE)
109 schemes and bulk K-profile schemes. TKE-based PBL parameterizations, such as the Mel-
110 lor–Yamada–Nakanishi–Niino (MYNN) scheme (Nakanishi and Niino 2009; Olson et al.
111 2019), use variables or gradients of variables vertically adjacent to a given point to de-
112 termine the amount of vertical mixing at that point (Cohen et al. 2015). As a result, TKE
113 schemes develop an eddy diffusivity profile, with contemporary schemes of at least 1.5 order
114 using prognostic equations for TKE to derive eddy diffusivities within a column for heat,

115 momentum, and moisture. K-profile schemes, such as the Yonsei University (YSU) scheme
 116 (Hong et al. 2006), determine the PBL height and impose an empirical eddy diffusivity
 117 profile through the PBL. Some modern schemes, such as the Asymmetric Convective Model
 118 version 2 (ACM2; Pleim 2007), incorporate concepts from both TKE and K-profile closure
 119 approaches, so herein we consider it a third type of scheme.

120 Many PBL schemes featuring a K-profile closure approach, such as ACM2 and YSU, utilize
 121 a critical Richardson number (CRN) to determine PBL height and, consequently, control
 122 the depth and strength of the imposed eddy diffusivity profile. First, a bulk Richardson
 123 number (hereafter BRN) is computed for each grid column. Hong (2010) define the BRN in
 124 the YSU as:

$$\text{BRN}(z) = \frac{g[\theta_v - \theta_s]z}{\theta_{va}U(z)^2}, \quad (2)$$

125 where θ_v is the virtual potential temperature, θ_s is the near-surface virtual potential tem-
 126 perature, θ_{va} is the virtual potential temperature at the lowest model level, and $U(z)$ is the
 127 wind speed at height z . The BRN is calculated from the surface to progressively higher
 128 levels until the CRN is reached. In YSU, a surface thermal excess term is applied to θ_s , in
 129 unstable conditions, and the BRN is recalculated to obtain the final PBL height. ACM2 uses
 130 a similar approach except the BRN is calculated over the entrainment layer only, starting
 131 at the level of neutral buoyancy with respect to rising surface layer air parcels (Pleim 2007)
 132 rather than the surface.

133 Several studies have investigated how changing the CRN in K-profile schemes impacts
 134 model performance (Hong and Pan 1996; Cohen et al. 2017; Bu et al. 2017). Other factors
 135 being equal, a smaller CRN lowers the PBL height and reduces eddy diffusivity through
 136 the PBL when compared to a larger value (Kepert 2012). In the present study, it is shown

in Section 3 that varying the CRN in a single K-profile PBL scheme is useful to achieve variability in PBL mixing strength while simplifying the experiment by limiting the number of variables and isolating the depth and strength of the imposed eddy diffusivity profile as the source of sensitivity within the simulations.

b. 26–29 January 2015 snowstorm and model setup

The 26–29 January snowstorm brought heavy snow to many portions of the northeastern U.S., particularly along the I-95 corridor (Fig. 1). An upper-level trough propagated south-eastward out of Canada and amplified as it approached the U.S. coastline, resulting in the formation of a sub-984-hPa coastal surface cyclone. The intense snowfall in New England and the sharp gradient in snowfall on the western edge of the storm made the event particularly notorious. The tight snowfall gradient contributed to considerable forecast uncertainty regarding precipitation totals along the coast, especially in New York City and surrounding areas of New Jersey.¹ A more detailed review of this event is provided by Greybush et al. (2017) who linked uncertainties in snowfall location to variations in the position of the coastal low.

The WRF model’s Advanced Research WRF core version 3.7.1, incorporating ERA-Interim (European Centre for Medium-Range Weather Forecasts 2009; Dee et al. 2011) data for initial and boundary conditions, is employed to conduct 72-h simulations of the snowstorm initialized at 0000 UTC 26 January 2015. Three telescoping domains, with horizontal grid spacings of 36, 12, and 4 km (Fig. 2), are employed, and our model physics and vertical resolution are similar to the High-Resolution Rapid Refresh model described

¹The difficult snowfall forecast resulted in an apology issued by a NWS meteorologist at the Mt. Holly, NJ forecast office on their personal Twitter account (Babay 2015) which, to the authors’ knowledge, is an unusual occurrence.

158 in Benjamin et al. (2016) except we use the modified MM5 surface layer scheme (Jiménez
 159 et al. 2012), Noah land surface model (Ek et al. 2003), and a variety of PBL schemes. The
 160 surface layer scheme is responsible for computing the transfer coefficients for momentum and
 161 scalar fluxes from the land surface model to the atmosphere. Consequently, PBL schemes in
 162 WRF rely on the surface fluxes computed from the surface layer to estimate subgrid-scale
 163 mixing and are often paired with unique surface layer schemes. Recent research has shown
 164 considerable variability surrounding how fluxes and frictional effects are calculated within
 165 surface layer schemes available in WRF (Minder et al. 2020). Therefore, to simplify the
 166 experiment and focus our study on the effects of PBL parameterizations, we use a single
 167 surface layer scheme, the modified MM5, due to its compatibility with several PBL schemes.

168 Several sensitivity tests are conducted to examine how robust cyclone evolution sensitivity
 169 is to various physics configurations. A ten-member ensemble is created using the Stochastic
 170 Kinetic Energy Backscatter technique (SKEBS; Berner et al. 2011) to test whether dif-
 171 ferences in cyclone evolution, arising from variations in PBL mixing, are consistent. The
 172 SKEBS scheme assumes that upscale and downscale cascading energy results in forcing for
 173 the resolved flow from unresolved scales with a perturbation amplitude proportional to the
 174 instantaneous dissipation rate. The SKEBS scheme, outlined in Berner et al. (2011) and
 175 implemented in WRF, includes a constant dissipation rate in space and time, allowing the
 176 perturbations to be considered as additive noise to the horizontal velocity (u, v) compo-
 177 nents and potential temperature (θ) fields with prescribed spatial and temporal correlation
 178 (Duda et al. 2016). SKEBS parameters used for the ensemble, such as the magnitude and
 179 scale of the perturbations, follow the recommended values (National Center for Atmospheric
 180 Research, 2015, pp. 5-26).

181 In light of prior work regarding sensitivity to surface friction and heat fluxes, it stands to
182 reason that the variety of methods used to determine PBL mixing may result in differences
183 in cyclone evolution. Simulations of the 26–29 January snowstorm are performed using
184 the MYNN, YSU, and ACM2, representing the three categories of PBL schemes described
185 above. Results, presented in Section 3, indicate considerable diversity among the simulations
186 warranting further study of how baroclinic cyclones are affected by the representation of PBL
187 mixing within a numerical weather prediction model.

188 Trajectories, shown in the Section 4, are computed using LAGRANTO (Sprenger and
189 Wernli 2015) utilizing hourly model output. All model fields shown are using the innermost
190 4-km domain unless specified otherwise. All snowfall accumulations presented herein use
191 a 10:1 snow-to-liquid ratio for consistency. A 10:1 ratio was chosen to approximate the
192 climatological snow-to-liquid ratio for the coastal northeastern U.S. in March (Baxter et al.
193 2005; their Fig. 8); however, we are most interested in the spatial pattern of snowfall for
194 this study rather than the snowfall amounts.

195 **3. Sensitivity experiments and results**

196 *a. Proof of concept and control runs*

197 As a proof of concept experiment, three 72-h simulations of the 26–29 January 2015 snow-
198 storm are performed using the MYNN, YSU, and ACM2 PBL schemes, to establish whether
199 any facet of baroclinic cyclone evolution is sensitive to the selection of PBL parameterization
200 scheme. The mean sea-level pressure (MSLP) centers in all simulations vary somewhat as the
201 systems organize, but the cyclones generally follow a northward track before turning north-
202 eastward off the New England coastline (Fig. 3). While MSLP tracks overlap somewhat,

203 there are subtle yet important differences, particularly in the propagation of the surface
 204 cyclone centers. Concomitant with the variations in track are marked differences in snowfall
 205 footprints among the three simulations. This may be anticipated due to the demonstrated
 206 relationship between the January 26–29 surface cyclone track and snowfall location and in-
 207 tensity by Greybush et al. (2017) (their Fig. 2). The MYNN, which uses a TKE approach
 208 to diagnose subgrid-scale mixing, generates an area of heavy snowfall greater than 61 cm on
 209 Long Island (Fig. 3c), similar to the YSU (Fig. 3a) and ACM2 (Fig. 3b) runs, but produces
 210 less snowfall across the domain elsewhere. The ACM2 and YSU snowfall footprints are more
 211 similar but the YSU snowfall footprint is shifted to the west (Figs. 4a,b), which can be (and
 212 in this case, is) important because the storm tracks run very close to large metropolitan
 213 areas. Discrepancies in forecasted precipitation around densely populated regions, such as
 214 those seen here between the YSU and ACM2 simulations, can lead to substantial variations
 215 in forecasted societal impacts.

216 While the hybrid ACM2 scheme and K-profile YSU scheme differ in many ways, both use
 217 a CRN to diagnose the PBL height and impose a K profile (K_c in Eq. 1) that establishes the
 218 amount of mixing between adjacent vertical grid levels within the PBL. The ACM2 scheme
 219 uses a default CRN of 0.25 for unstable surface layers while the YSU scheme uses a CRN of
 220 0.00 for the same (Pleim 2007; Hong 2010). A CRN of 0.25 allows diagnosis of the PBL top
 221 to be within the capping inversion of a classic convective boundary layer, permitting implicit
 222 mixing and entrainment at the top of the boundary layer. A CRN of 0.00 limits implicit
 223 entrainment; therefore, YSU calculates PBL-top entrainment explicitly using an additional
 224 term under N_{NL} of Eq. 1.

225 An important difference between the two parameterizations, aside from their default CRNs
 226 and the layers over which they calculate the BRN, is how each scheme computes upward and

227 downward mixing. ACM2 is a hybrid, or “asymmetric”, scheme, meaning that upward and
228 downward mixing rates are defined separately: upward mixing is dependent on both local
229 gradients of mixed variables (e.g. heat, momentum, and moisture) and an imposed mixing
230 profile dependent on surface buoyancy and the CRN, whereas downward mixing is related to
231 upward mixing through mass conservation but only uses local gradients to derive downward
232 mixing rates (Pleim 2007). YSU also imposes an eddy diffusivity profile derived from surface
233 fluxes and the CRN, but treats upward and downward mixing equally. While these unique
234 characteristics likely contribute to the variations in storm-total snowfall seen in Fig. 4, we
235 wish to explore whether the differences between these schemes may be approximated by
236 modifying the CRN.

237 *b. Proxy runs*

238 In an attempt to resolve the processes by which these PBL schemes impact a baroclinic
239 cyclone, we seek to reduce the degrees of freedom of the experiment by using a single PBL
240 scheme with modified mixing parameters to capture the variability seen in Fig. 4. Therefore,
241 we conduct another simulation using the YSU scheme, but now with a CRN of 0.25 like
242 ACM2. This simulation will be referred to as “more mixing” and the original YSU run
243 with CRN of 0.00 will be termed “less mixing”. We then investigate whether changing a
244 single mixing parameter yields similar variability seen between the original YSU and ACM2
245 simulations (Fig. 4). Note that raising the CRN to 0.25 permits both implicit and explicit
246 PBL mixing within the YSU due to the PBL-top entrainment term mentioned above. We
247 performed CRN = 0.25 experiments with and without the parameterized entrainment and
248 noted little impact on surface cyclone movement and similar snowfall footprints (not shown).

249 As a consequence, for simplicity, the more-mixing experiments described herein have retained
250 the PBL-top entrainment.

251 Initially, the MSLP fields of the less-mixing and more-mixing simulations evolve similarly,
252 but the locations of the surface low centers begin to diverge around forecast hour 40 (i.e.,
253 1600 UTC 27 January) as the cyclones approach the time of minimum MSLP (Fig. 5a). The
254 surface low pressure center in the more-mixing case begins to move northeastward relative
255 to the surface low pressure center in the less-mixing case (Figs. 5b,c), which appears to
256 stall off the New England coast (Fig. 5c inset), leading to a difference in cyclone center
257 locations on the order of 80 km by 0600 UTC 28 January. This relative orientation persists
258 until approximately 1500 UTC 28 January as the surface cyclones occlude and their respec-
259 tive central pressure minima become less defined. The differences in stalling behavior and
260 cyclone propagation between the two YSU runs, although ostensibly small, contribute to
261 a considerable change in snowfall location and intensity (Fig. 5d) that is impactful, again
262 owing to the storms' proximity to densely populated areas. Note the less-mixing simulation
263 produces more snowfall on the western flank of the cyclone over western New England and
264 eastern New York.

265 Despite the similarities between the more-mixing and ACM2 runs (Fig. 6), the more-
266 mixing simulation does not replicate some other aspects of the ACM2 case, suggesting there
267 are factors other than the CRN that control PBL mixing and how that mixing projects
268 to larger scales. Nevertheless, the more rapid movement of the more-mixing run's cyclone
269 demonstrates that the CRN may be manipulated to provide a reasonable spread in model
270 solutions resembling what can be expected using different PBL physics.

271 The larger CRN of the more-mixing simulation induces stronger eddy diffusivity within
272 the PBL than the less-mixing simulation, which is a direct consequence of it encouraging

273 relatively deeper boundary layers. This can be seen in an area-average profile taken through
 274 the warm sector of the cyclone (Fig. 7a). Here, the warm sector is delineated by first com-
 275 puting the average and standard deviation of the 950–800-hPa potential temperature across
 276 the innermost domain at a given time and using this to convert the field into standardized
 277 anomalies. Next, to create the profiles, we average only those gridpoints possessing a pos-
 278 itive value (i.e. the “warm” values; Fig. 8). This technique divides the domain into two
 279 equal parts at each time for each simulation and seeks to prevent cold sector profiles, which
 280 are predominately unstable in an oceanic cyclone owing to strong cold air advection over
 281 warmer waters, from obscuring differences in the warm-sector PBLs.

282 As anticipated, area-averaged vertical profiles of eddy diffusivity for scalars² indicate
 283 stronger mixing over a deeper depth in the more-mixing simulation (Fig. 7a). The dif-
 284 ferences in eddy diffusivity result, predictably, in higher wind speeds at the surface and
 285 lower wind speeds aloft in the more-mixing simulation (Fig. 7b). Furthermore, the vertical
 286 moisture profile suggests the lowest levels (below 800 m) of the PBL are more moist in the
 287 less-mixing case and drier aloft (Fig. 7c). Reduced mixing limits both entrainment of free
 288 tropospheric air into the PBL and inhibits the upward mixing of moisture away from the
 289 near-surface where moisture content is largest, thereby preserving moisture in the sub-800-m
 290 layer.

291 In an attempt to link the enhanced PBL mixing with reduced near-surface moisture,
 292 accumulated boundary layer mixing ratio (hereafter, ACBLQ) is computed during the model
 293 integrations. ACBLQ is calculated by multiplying the instantaneous tendency in vapor
 294 mixing ratio from the PBL scheme by the timestep (24 s for the innermost domain) and

²The eddy diffusivity of scalars is related to the eddy diffusivity of momentum via the Prandtl number that, while allowed to vary with height in the YSU, produces a similar momentum diffusivity profile shape (not shown) as that seen in Fig. 7a.

binning the result for each grid volume. Consequently, ACBLQ describes the degree to which the PBL scheme is moistening (positive values) or drying a grid volume since model initiation. Fig. 7d indicates the PBL scheme in the less-mixing simulation is moistening the lowest 800-m layer of the warm sector more than the more-mixing run. This suggests the instantaneous moisture differences at 1800 UTC 27 January are associated with PBL moisture tendency variations derived from changes in PBL mixing strength.

c. SKEBS experiment

Motivated by the above experiment, a SKEBS ensemble based on the proxy runs is created to test the robustness of the demonstrated variations in surface cyclone movement. Again, the SKEBS scheme adds stochastic, small-amplitude perturbations to the rotational component of the horizontal wind and potential temperature tendency equations at each timestep (Berner et al. 2011). This additive noise approach essentially creates different versions of the January snowstorm to assess the consistency of the relationship between cyclone track and mixing strength. The ensemble consists of ten members using the same domains in Fig. 2, five being less-mixing (i.e., default YSU) members along with five more-mixing YSU members using a CRN of 0.25.

Each ensemble member from the less-mixing group is paired with its more-mixing counterpart employing the same random seed. The results suggest the SKEBS perturbations add more spread to the model solution than changing the CRN alone (Figs. 9a–c). This may be anticipated given the SKEBS perturbations are added through the entire depth of the model atmosphere.³ However, the MSLP and precipitation difference patterns, highlighted in the

³Although PBL schemes apply vertical mixing through the entire model column, it stands to reason that the largest differences are confined below the boundary layer height where eddy mixing is usually most vigorous.

316 previous experiment, are persistent among the ensemble sets (Figs. 9a–c). The ensemble
 317 mean differences of MSLP low location and accumulated snowfall between the two mixing
 318 regimes (Fig. 9d) is consistent with the difference patterns illustrated by the individual en-
 319 semble sets. The surface cyclones in the more-mixing ensemble move more rapidly toward
 320 the north and east relative to their less-mixing counterparts in each of the five pairs and the
 321 ensemble mean (three sets and the ensemble mean shown for clarity in Figs. 9a–d). These
 322 results suggest the tendency of the more-mixing surface cyclone to move more quickly north
 323 and east is unlikely to be explained by random chance.

324 *d. Latent heat and moisture sensitivity runs*

325 It was hypothesized in Section 1 that moisture content within the PBL could affect the
 326 evolution of a baroclinic cyclone through modifying latent heat release. An experiment is
 327 conducted where latent heating is turned off (NOLH) within the model starting at 0000
 328 UTC 27 January (forecast hour 24) to assess the importance of latent heat release to the
 329 evolution of each cyclone.⁴ The surface cyclones in the NOLH experiment progress along the
 330 east coast of the U.S., in general agreement with the less-mixing/more-mixing experiment,
 331 but fail to deepen to the same degree as in the full-physics simulations. Additionally, the less-
 332 mixing surface cyclone moves *faster* to the north and east than its more-mixing counterpart
 333 (Fig. 10), contrary to the simulations with latent heating (Fig. 5d). This difference in
 334 cyclone progression is consistent with the inverted snowfall difference pattern relative to the
 335 full-physics experiment (Figs. 5d and 10). The NOLH more-mixing simulation produces
 336 more snow around and to the east of the Hudson Valley region, although snowfall difference

⁴We remove latent heating after 24 h into the simulation due to the lack of an organized surface cyclone in runs with no
 latent heating beginning at initialization, making them unsuitable for comparison here.

337 magnitudes are smaller than the full-physics proxy runs due, in part, to the weaker NOLH
338 systems.

339 The results of the NOLH experiment suggest latent heating—and consequently, moisture—
340 play a critical role in how variations in PBL mixing impact baroclinic cyclones. However,
341 removing latent heating drastically alters the evolution and strength of the simulated cyclone
342 to the point where the simulation is no longer representative of the 26–29 January snowstorm.
343 This could be anticipated from Stoelinga (1996), who demonstrated latent heating can have
344 a substantial effect on the low-level circulation field of oceanic cyclones and the subsequent
345 lifecycle of the baroclinic system. Thus, a more surgical approach is needed to preserve the
346 general structure and intensity of the snowstorm while assessing the system’s sensitivity to
347 moisture.

348 To this end, we conduct a YSU experiment (hereafter, MOISTMIX) that increases the
349 CRN used for calculating the eddy diffusivities applied to water species (i.e., vapor and cloud
350 water) to 0.25 while retaining the scheme’s default 0.00 CRN in the determination of eddy
351 diffusivity of momentum and temperature, again whenever and wherever the surface layer
352 is unstable. This is accomplished by solving for heat, moisture, and momentum tendencies
353 using a modified CRN of 0.25 in the YSU. The unstable regime CRN is then adjusted back
354 to 0.00 and the YSU is allowed to recalculate the temperature and momentum tendencies.
355 In other words, MOISTMIX is the less-mixing run except with respect to water vapor
356 and cloud water mixing, for which it mimics the more-mixing simulation. Results from
357 the MOISTMIX experiment (Fig. 11a) indicate that an increase in eddy diffusivity applied
358 to water species has a similar influence on cyclone propagation and snowfall footprint as
359 increased eddy diffusivity across all scalars and momentum (Fig. 5d), with more precipitation
360 on the western flank of the less-moisture-mixing cyclone relative to the faster-propagating

361 MOISTMIX cyclone. Additionally, MOISTMIX closely matches ACM2 in cyclone track,
362 propagation speed, and snowfall footprint (Fig. 11b), further highlighting the importance of
363 moisture mixing by PBL schemes on cyclone evolution.

364 The sensitivity of cyclone propagation to the vertical mixing of moisture in the MOISTMIX
365 experiment suggests that latent heating plays a considerable role in differentiating between
366 the two mixing regime simulations demonstrated earlier (Fig. 5d). As the less-mixing/more-
367 mixing surface cyclones diverge (Figs. 5a–c), higher upper-level heights aloft develop to the
368 north and east of the less-mixing cyclone (Fig. 12a) compared to its more-mixing counter-
369 part. The higher heights are coincident with warmer temperatures in the 400–300-hPa layer
370 (Fig. 12b) and are consistent with stronger upper-level ridging in the less-mixing simulation
371 (Fig. 12a). It is hypothesized that stronger latent heating in the less-mixing cyclone con-
372 tributes to stronger ridge building downstream of the surface cyclone through the diabatic
373 term of the quasigeostrophic height tendency equation, which supports forcing for height
374 rises above maximums in diabatic heating. Although not shown, enhanced ridging is present
375 using middle-tropospheric pressure layers (e.g. 400–600 hPa) to the south-southeast of the
376 green outlined area in Fig. 12a suggesting latent heating differences extend along the warm
377 conveyor belt. The 900–500-hPa thickness lines of the less-mixing run are shifted westward
378 of the more-mixing run (Fig. 12c), indicating warmer air extends further west in the less-
379 mixing run consistent with stronger latent heating in the lower troposphere along the warm
380 conveyor belt.

381 Recall from Figures 5a,c (inset) that the surface cyclones in both simulations have begun
382 to separate concurrent with the differences in 900–500-hPa thickness shown in Figure 12c.
383 The primary consequence of the shift in the thickness field is a reduction in thermal wind
384 strength over the less-mixing surface cyclone, as the strongest gradient in thickness is shifted

385 to the west of the cyclone center (Fig. 12c). Weaker shear above the surface pressure
 386 minimum supports the slower propagation speed of the less-mixing cyclone, seen in Figs. 5a–
 387 c, according to the thermal steering term of Sutcliffe (1947). The Trenberth form of the
 388 quasigeostrophic omega equation (Trenberth 1978) may also be used to explain the variation
 389 in cyclone propagation. A weaker thermal wind field around the less-mixing surface cyclone
 390 results in weaker advection of middle-troposphere relative vorticity by the thermal wind
 391 (not shown) downshear of the surface cyclone, thus reducing both the forcing for upward
 392 vertical motion and propagation of the surface cyclone. Furthermore, enhanced ridging and
 393 associated upper-level divergence amplify the upper-level flow pattern, producing a deeper
 394 trough and a stronger downstream ridge in the less-mixing case (Fig. 12a). These processes,
 395 combined with a shortening of the wavelength of the upper-level baroclinic wave due to the
 396 aforementioned amplification, likely act to slow the northeastward progression of the surface
 397 cyclone.

398 It is hypothesized that the height differences seen in Fig. 12a are a result of stronger latent
 399 heating to the north and east of the surface cyclone in the less-mixing case due to higher
 400 PBL moisture content. A CRN of 0.25 results in stronger mixing through a deeper column,
 401 reducing water vapor content in the low levels of the PBL while increasing the water vapor
 402 content aloft (Figs. 7c,d). The higher moisture values below 800 m in the less-mixing case
 403 (Fig. 7c) suggest weaker PBL mixing effectively preserves PBL moisture, leading to the
 404 higher warm-sector MUCAPE values seen at all times through the simulations (Fig. 13).
 405 This is consistent with the findings of Hong et al. (2006), who found weaker mixing resulted
 406 in higher CAPE and higher moisture values trapped near the surface. Higher water vapor
 407 content and larger MUCAPE may lead to stronger condensational heating through ascent
 408 above the warm front as parcels originating in the PBL are lifted via the warm conveyor

409 belt. The dynamic nature of the warm conveyor belt motivates a Lagrangian approach to
410 analyze how differences in low-level moisture content may impact upper-level flow patterns
411 and, subsequently, surface cyclone propagation, in the next section.

412 **4. Trajectory analysis**

413 Trajectory analysis, using the LAGRANTO (Sprenger and Wernli 2015) software pack-
414 age adapted for WRF, is conducted to investigate the origins of the height differences be-
415 tween the less-mixing/more-mixing simulations (Fig. 12a). LAGRANTO was selected for
416 this work due to its computational efficiency, ease of use for calculating trajectory swarms,
417 and compatibility with WRF. Swarms of 24-h backward trajectories are launched within
418 the less-mixing/more-mixing simulations at 0600 UTC 28 January from the 6–10-km layer
419 surrounding the area of upper-level height differences (green box in Fig. 12a).

420 The majority of parcels experiencing the most ascent in each swarm originate in the warm
421 sector to the east of the storm center and rise to the north of the cyclone center (Figs. 14a,b),
422 resembling the conceptual model of a warm conveyor belt in Schemm and Wernli (2014).
423 A smaller subset of parcels originate to the east of the North Carolina and Virginia coast-
424 line and ascend sharply to the 300-hPa level before transiting to the north of the cyclone,
425 revealing a southern ascent pathway in both simulations. Interestingly, the more-mixing
426 simulation generates more parcels in this southern pathway than the less-mixing simulation.
427 We speculate that stronger PBL mixing reduces CIN and may have led to earlier triggering
428 of convection in the more-mixing simulation similar to the process described in Hong et al.
429 (2006). Aside from the southern ascent pathway differences, ascending parcels follow similar
430 paths in both simulations; however, the ostensibly similar Lagrangian structures belie im-
431 portant differences in the properties of the ascending parcel swarms. Therefore, we compare

the bulk characteristics of each trajectory swarm between the two simulations to investigate the source of the upper-level height differences seen in Fig. 12a.

Figure 15a–c depicts the pressure, temperature, and moisture content distributions of each simulation’s backward trajectory swarm through 24 hours. Trajectories from the less-mixing simulation originate at higher pressures and ascend through a greater depth than the more-mixing simulation trajectories across the swarm distribution (Fig. 15a). This ascent is coincident with a larger potential temperature difference between trajectory origination and termination times in the less-mixing simulation (Fig. 15b), implying those parcels experienced stronger diabatic heating than those in the more-mixing run. Furthermore, the less-mixing trajectories originate with significantly higher water vapor content and terminate with roughly similar vapor content as the more-mixing trajectories (Fig. 15c). The larger deficit in specific humidity, along with the greater ascent and warmer termination temperatures, suggest the less-mixing simulation’s parcel trajectories undergo stronger condensational heating during their lifetimes, thereby promoting warmer temperatures aloft (Fig. 12b) and greater upper-tropospheric ridge building downstream of the surface cyclone (Fig. 12a). The more amplified flow pattern accompanies higher MUCAPE and greater low-level water vapor content values in the less-mixing run, contributing in a slower north-eastward progression of the surface cyclone.

5. Discussion and summary

Several studies have demonstrated how boundary layer processes can exert a substantial influence on baroclinic cyclone evolution. Specifically, the fluxes of heat and momentum from the surface impact the strength of Ekman pumping and the production of boundary layer PV, which can influence the system-scale circulation (Adamson et al. 2006; Beare 2007; Plant

455 and Belcher 2007). However, previous studies of PBL processes within baroclinic cyclones
456 have generally used dry, idealized model simulations. An important role for water substance
457 can be anticipated from the tropical cyclone (TC) literature, which indicates PBL mixing
458 can have a substantial impact on TC structure and evolution (Nolan et al. 2009; Kepert
459 2012; Bu et al. 2017). In particular, Bu et al. (2017) demonstrated the vertical mixing of
460 water vapor may influence TC size via modification of outer-core convective activity.

461 This motivated a study concerning the sensitivity of moist baroclinic cyclone evolution to
462 parameterized boundary layer mixing strength. Our principal finding is that variations in
463 eddy diffusion within the PBL can modify cyclone evolution, particularly the propagation
464 of the surface cyclone. The 26–29 January 2015 snowstorm, which dropped heavy snow
465 across densely populated areas along the Northeast coastline and whose snowfall footprint
466 was considerably dependent on cyclone track (Greybush et al. 2017), provides an example
467 of how small changes in model physics can have substantial impacts on 2–3-day forecasts.
468 By comparing simulations of this storm using the MYNN, YSU, and ACM2 PBL schemes,
469 representing 3 approaches to handling the PBL mixing available in the WRF model, we
470 determined that the choice of PBL scheme can influence sensible weather impacts of the
471 baroclinic cyclone by altering precipitation patterns and surface cyclone movement.

472 Both YSU and ACM2 use K-profile assumptions that are sensitive to choice of critical
473 Richardson number (CRN). The schemes use different CRN values when the surface layer is
474 determined to be unstable, reflecting disparate approaches to handling PBL-top entrainment.
475 However, irrespective of the manner in which entrainment is treated, we demonstrated the
476 variations between the YSU and ACM2 simulations with respect to snowfall footprint and
477 surface cyclone movement could largely be reproduced by manipulating the CRN value
478 within the YSU scheme itself to match that employed by ACM2.

Specifically, changing the unstable regime CRN from (the default) 0.00 (less mixing) to 0.25 (more mixing) in the YSU had a sizable impact on the magnitude and depth of the eddy diffusivity generated within the warm sector of the storm. Furthermore, lower CRN values consistently slowed the northeastward progression of the surface cyclone relative to the higher CRN runs, which impacts surface precipitation totals and type. Results from the SKEBS ensemble imply random chance is unlikely to cause these shifts and the special MOISTMIX experiment that limited CRN-associated mixing changes to water species alone demonstrates the primary role of water vapor in modulating the cyclone motion.

The mechanism by which changes in eddy mixing of moisture impact cyclone propagation is evidenced by differences in downstream ridging ahead of the surface cyclone between the less-mixing/more-mixing runs. Simulations with less mixing had stronger ridge building attendant with warmer temperatures in the middle and upper troposphere, amplifying the upper-level flow pattern and slowing the northeastward movement of the system. Reduced values of eddy diffusivity and a shallower PBL contribute to weaker entrainment of drier, free-tropospheric air into the PBL, less ventilation of near-surface moisture out of the PBL, and preservation of higher PBL moisture values within several hundred meters of the surface. This increased moisture content is consistent with higher MUCAPE in the less-mixing simulation and implies more vigorous upward vertical motion transpires when boundary layer parcels are lifted, leading to more condensational heating and contributing to further ridge building. Backward trajectory analysis suggests the strongest ascent occurs within the warm conveyor belt as parcels are lifted over the warm front. Trajectories terminating in the downstream ridge undergo greater ascent and diabatic heating in the less-mixing case, as a result of originating with more moisture and at higher pressures in the troposphere, consistent with the results of (Schäfler and Harnisch 2015) who found drier inflow regions

503 led to lower outflow heights for warm conveyor belt parcels. The enhanced heating both
504 weakens the thermal wind over the less-mixing cyclone and promotes a stronger upper-
505 tropospheric downstream ridge, which combine to inhibiting the northeastward propagation
506 of the less-mixing surface cyclone.

507 The results of this study are intended to demonstrate both the sensitivity of baroclinic
508 cyclones to PBL mixing strength and the predominant pathway through which that sensi-
509 tivity manifests itself. A variety of PBL parameterization schemes exist with little consensus
510 regarding which scheme produces the most realistic mixing in the boundary layer. The sen-
511 sitivity to PBL mixing exhibited by the baroclinic cyclone simulated herein, as well as other
512 phenomena such as tropical cyclones and severe convection discussed in previous studies,
513 motivates future work to better represent uncertainty surrounding PBL mixing strength.
514 Stochastic perturbations of parameters (SPP), such as the CRN in the YSU scheme or other
515 parameters governing vertical mixing that are not well constrained, may be considered as
516 a strategy to address the uncertainty in PBL mixing within ensemble weather prediction
517 systems.

518 Several studies have sought to represent the uncertainties in assumptions within physics
519 parameterization schemes, including PBL schemes, using the SPP approach (Ollinaho et al.
520 2017; Jankov et al. 2017, 2018). Jankov et al. (2017, 2018) evaluated 24-h WRF forecasts
521 using various stochastic perturbation strategies and found SPP can improve model perfor-
522 mance, particularly when combined with other perturbation strategies such as SKEBS and
523 stochastic perturbations of physics tendencies (SPPT). Ollinaho et al. (2017) evaluated an
524 SPP scheme against an SPPT scheme within the ECMWF IFS over 15-day forecast periods
525 and found improved 2-m temperature forecasts over the first couple of days. While only one
526 of these studies limited their SPP strategy to the PBL scheme (Jankov et al. 2018), their col-

lective results suggest SPP schemes can improve model spread, particularly near the surface where SPPT schemes are commonly tapered to zero in order to avoid numerical instabilities (Ollinaho et al. 2017). The present work builds on previous research and hypothesizes a pathway through which perturbations of PBL mixing, specifically mixing of water vapor, may project to larger scales. It is anticipated that understanding the pathways through which physics perturbations may grow upscale and impact synoptic weather features may lead to more informed and targeted applications of physics perturbation methods within numerical weather prediction systems.

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Data availability statement. ERA-Interim data used in this study is available at the NCAR Research Data Archive (<https://doi.org/10.5065/D6CR5RD9>). WRF model source code is available at <https://github.com/wrf-model/WRF>. WRF simulation data is archived at UAlbany and will be made available upon request.

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Experiment Name	Description
ACM2	ACM2 PBL scheme
MYNN	MYNN PBL scheme
Less mixing	Default YSU PBL scheme ($\text{CRN} = 0.00$)
More mixing	YSU PBL scheme with $\text{CRN} = 0.25$
Less mixing SKEBS	five-member less-mixing ensemble with SKEBS perturbations
More mixing SKEBS	five-member more-mixing ensemble with SKEBS perturbations
NOLH less mixing	Same as less mixing but without latent heating
NOLH more mixing	Same as more mixing but without latent heating
MOISTMIX	Same as less mixing but with $\text{CRN} = 0.25$ for moisture mixing

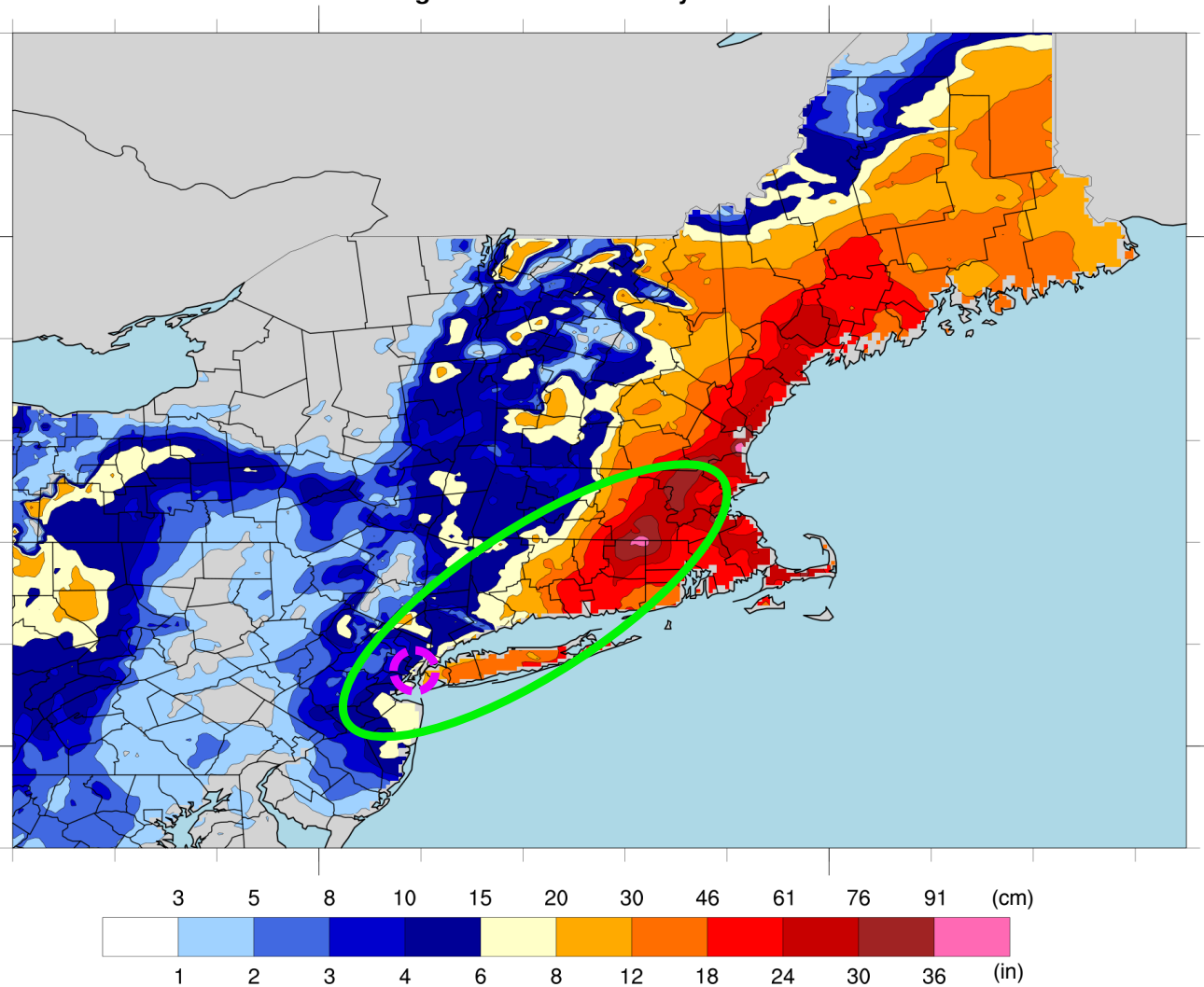
TABLE 1. Summary of experiments

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48-h snowfall accumulation ending 1200 UTC 28 January 2015



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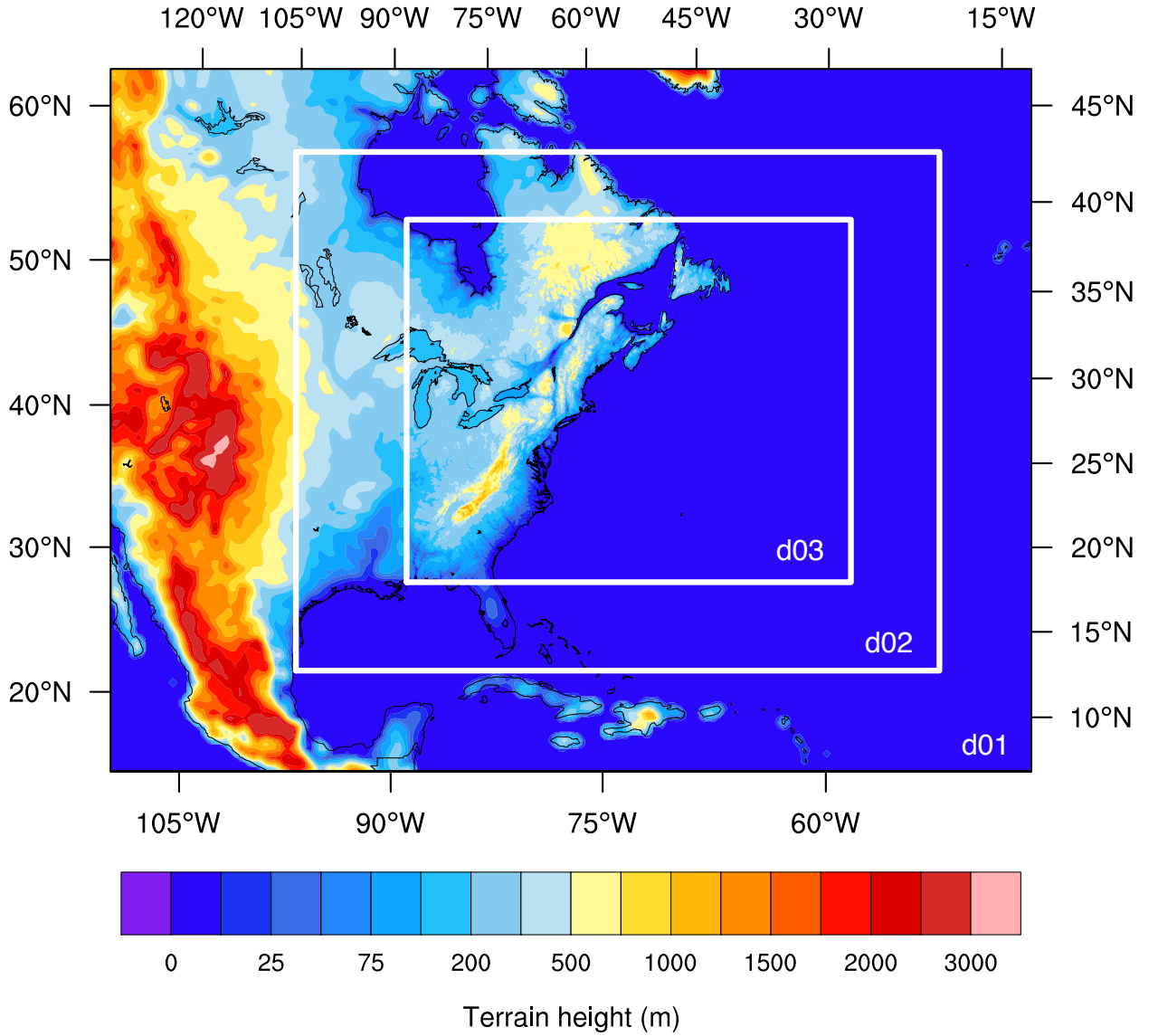


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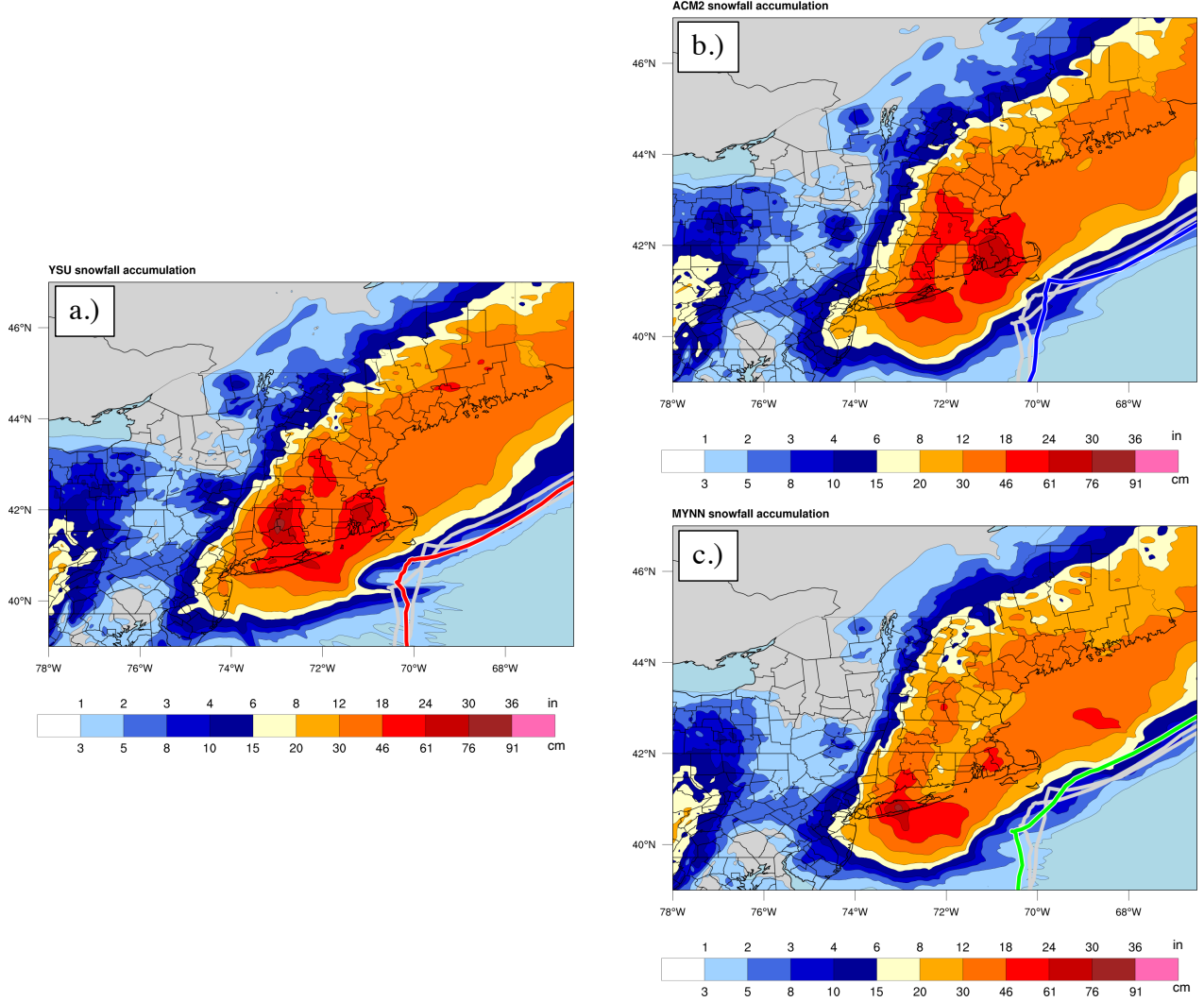


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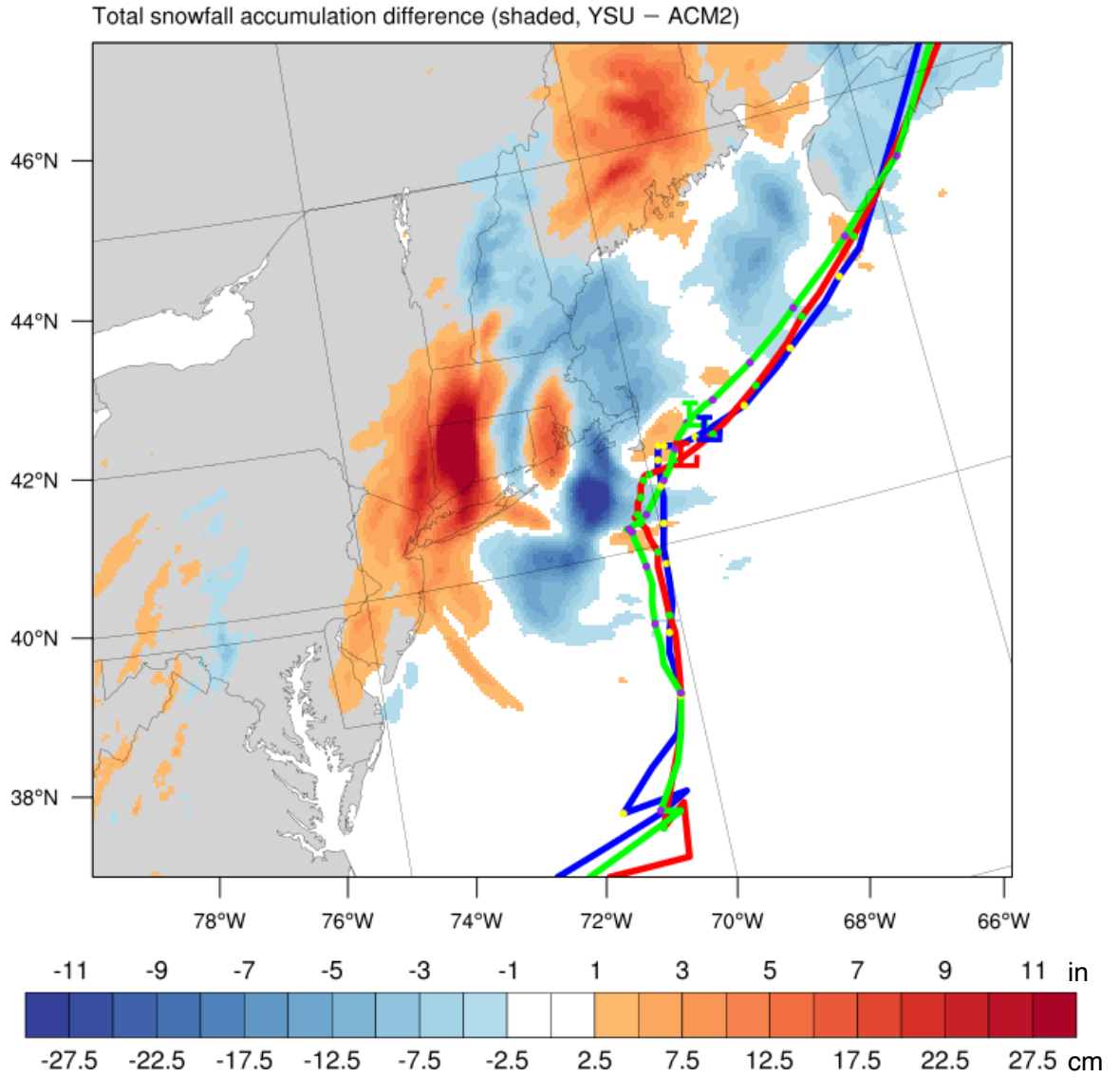


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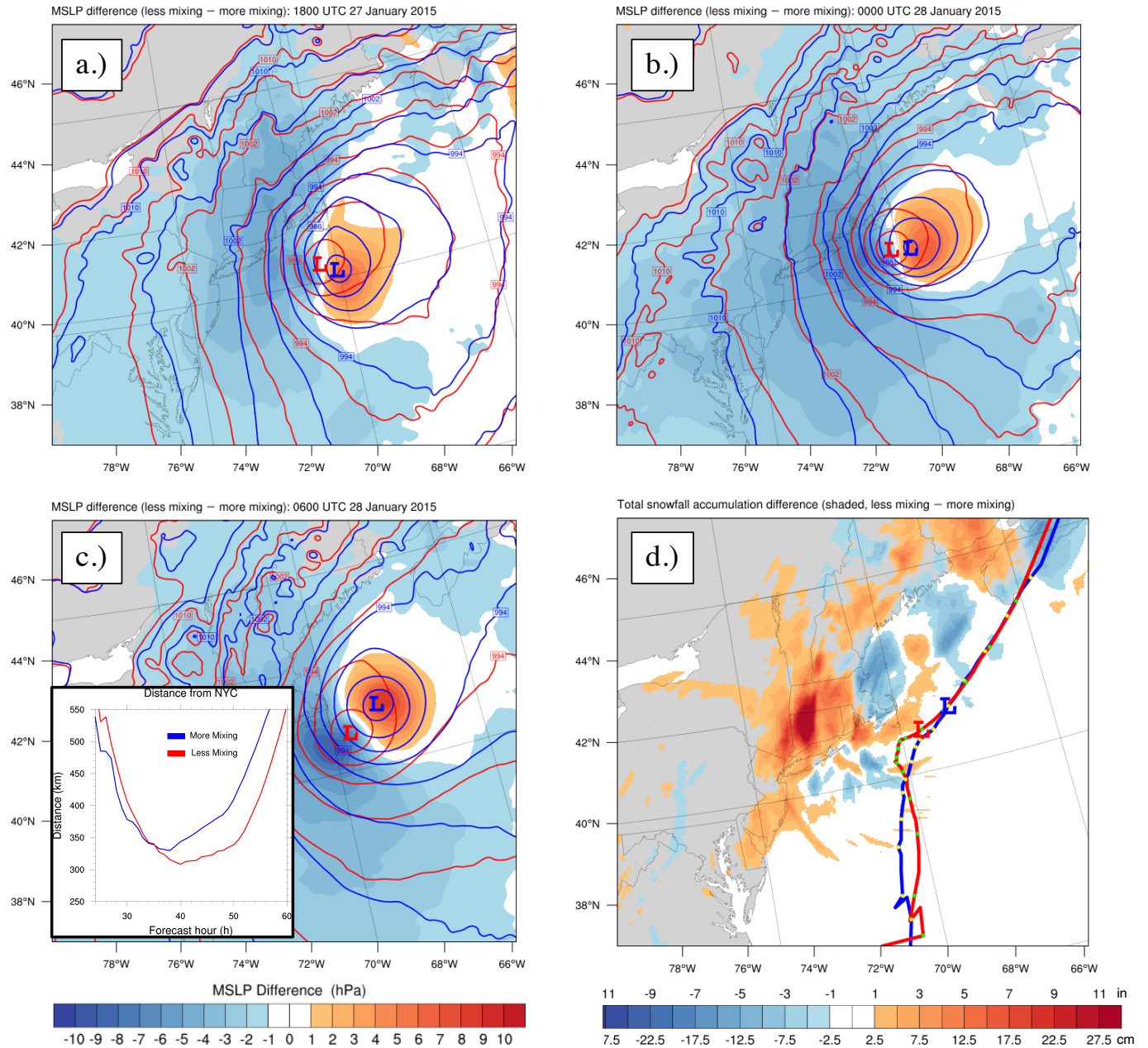


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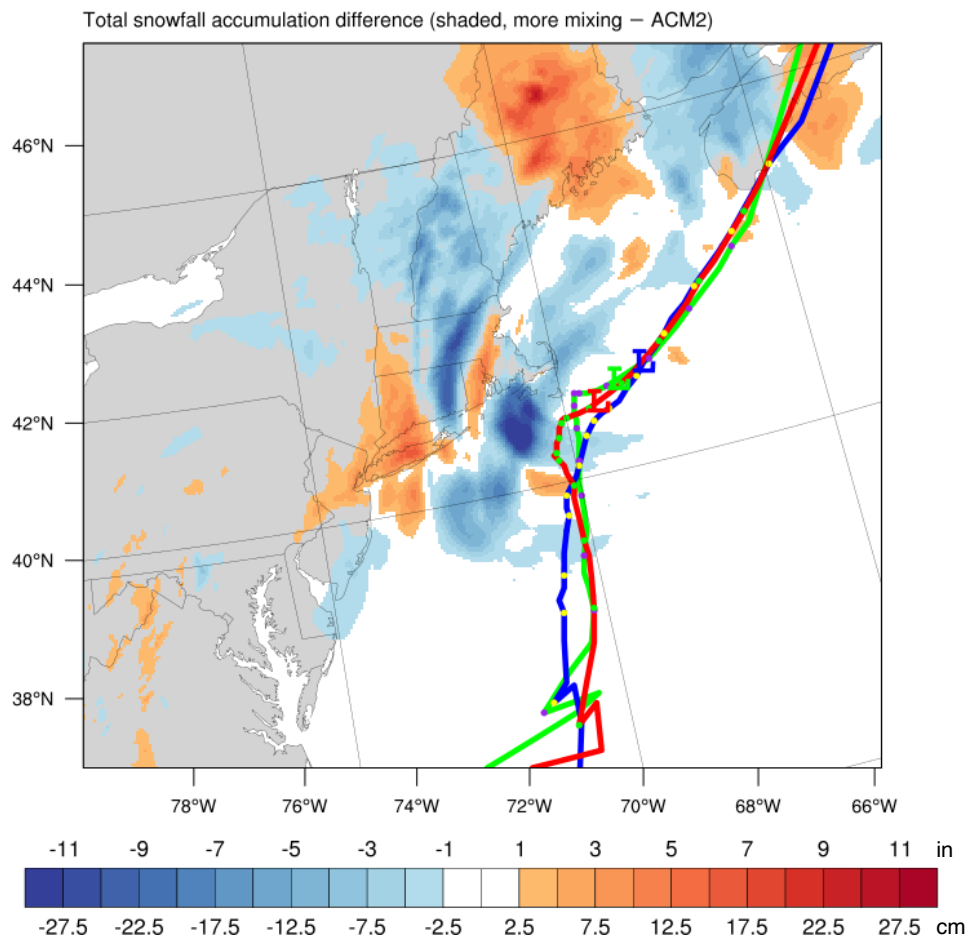
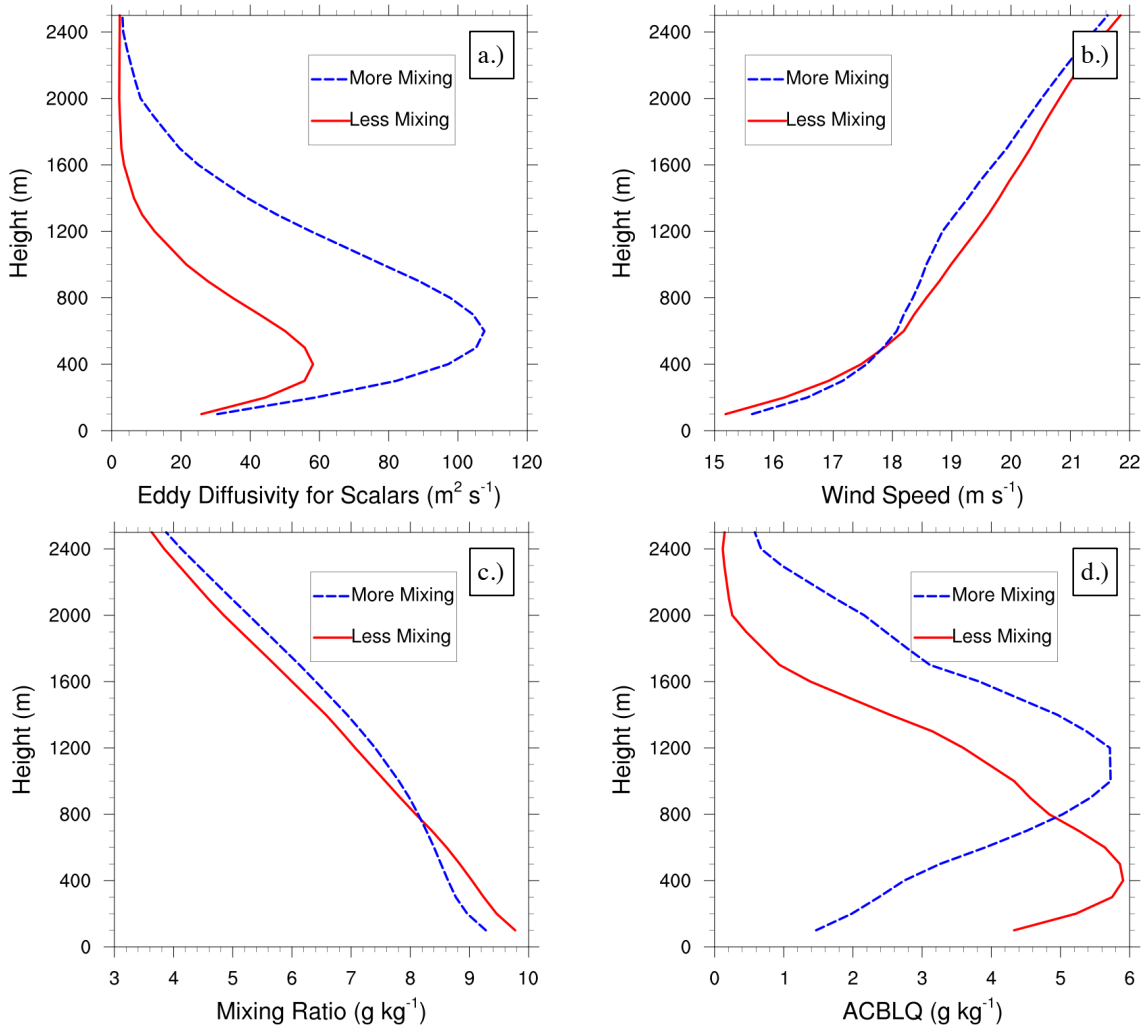


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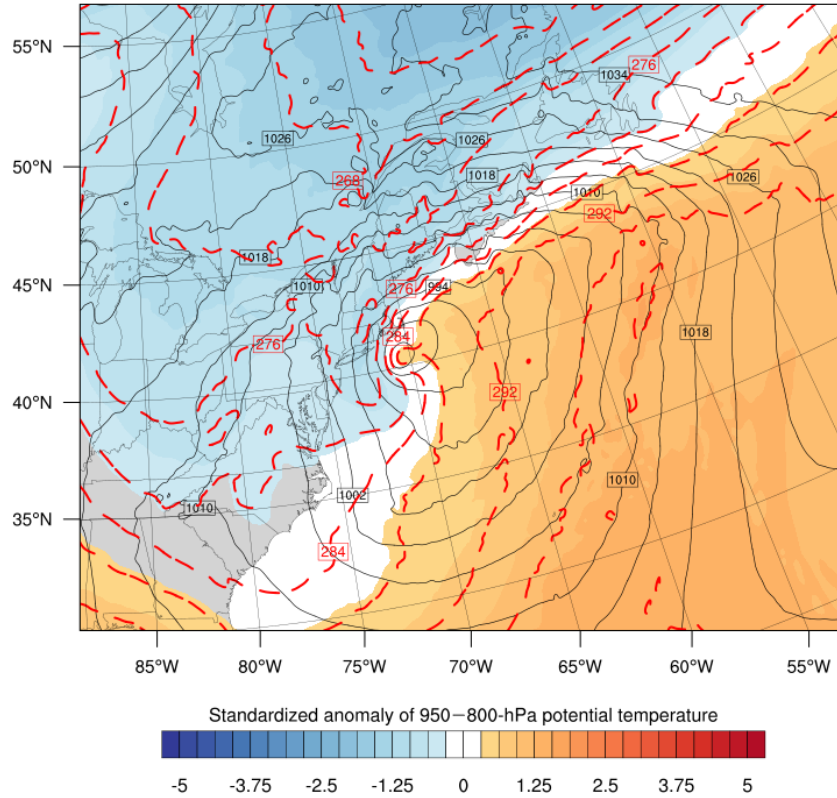


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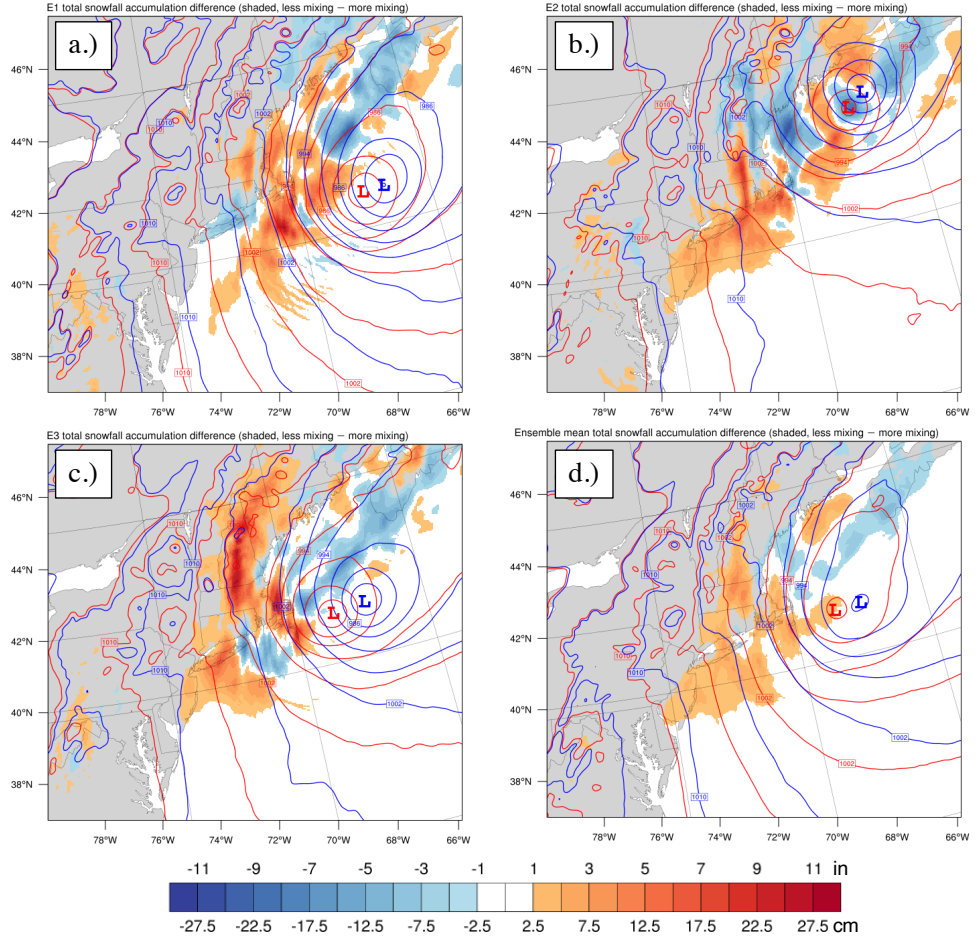


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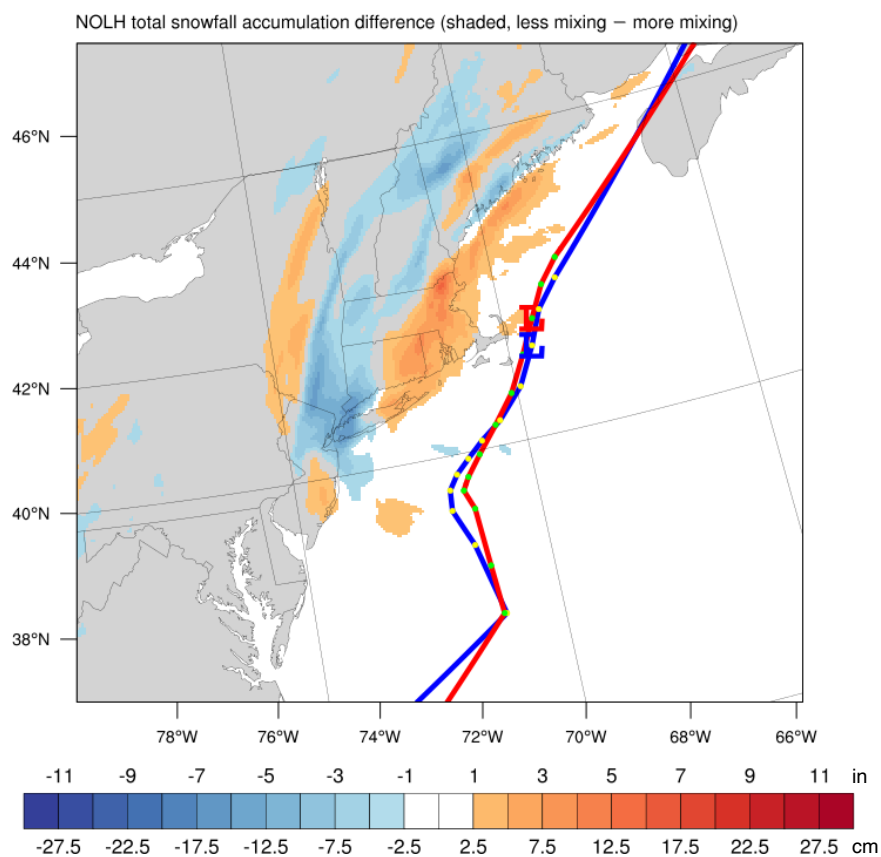


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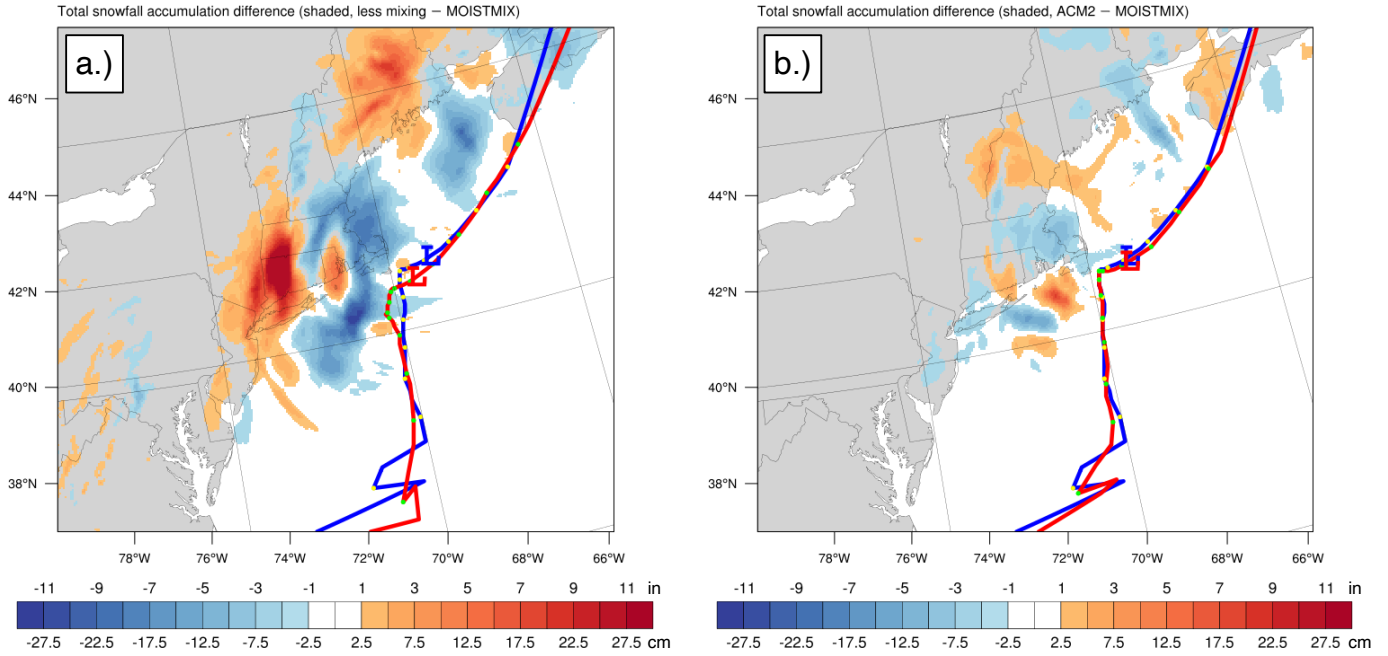


FIG. 11. (a) MSLP center tracks for (red) less-mixing and (blue) MOISTMIX simulations and total snowfall accumulation difference (less mixing–MOISTMIX, cm). Dots are 3-hourly surface cyclone center locations and bold "L"s represent surface low center locations at 0600 UTC 28 January 2015 for each simulation. (b) As in (a) but for (red) ACM2 and (blue) MOISTMIX simulations.

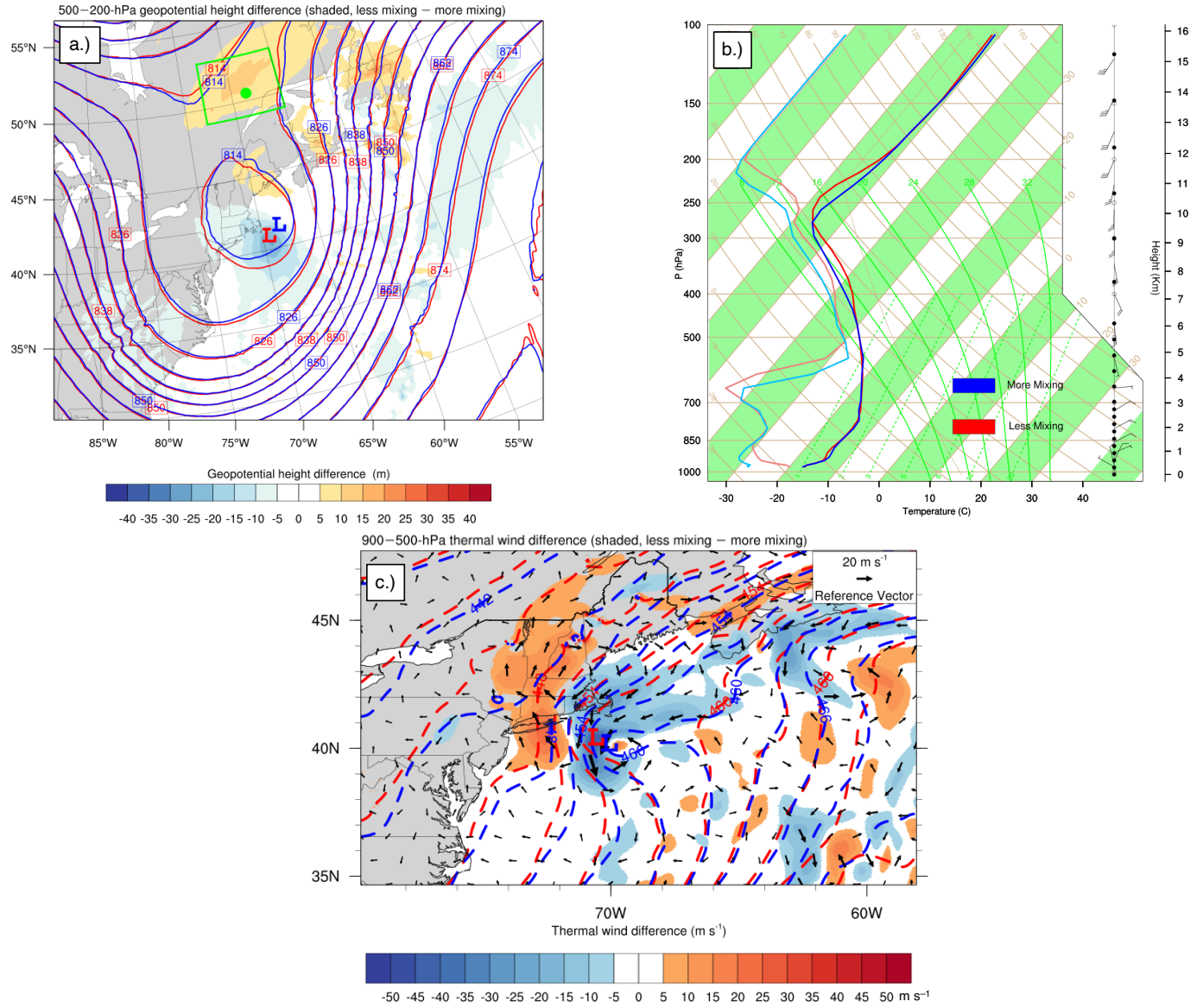


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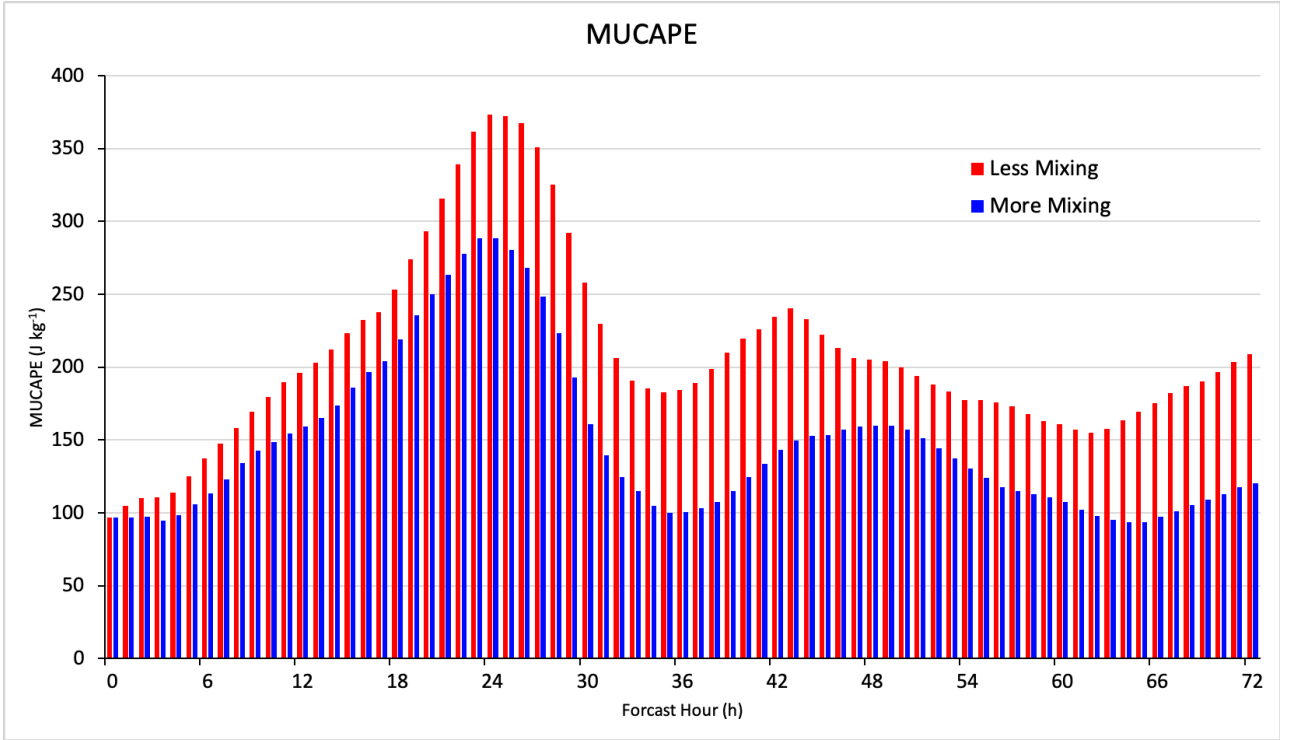


FIG. 13. Warm-sector area-averaged MUCAPE for the (red) less-mixing and (blue) more-mixing simulations.

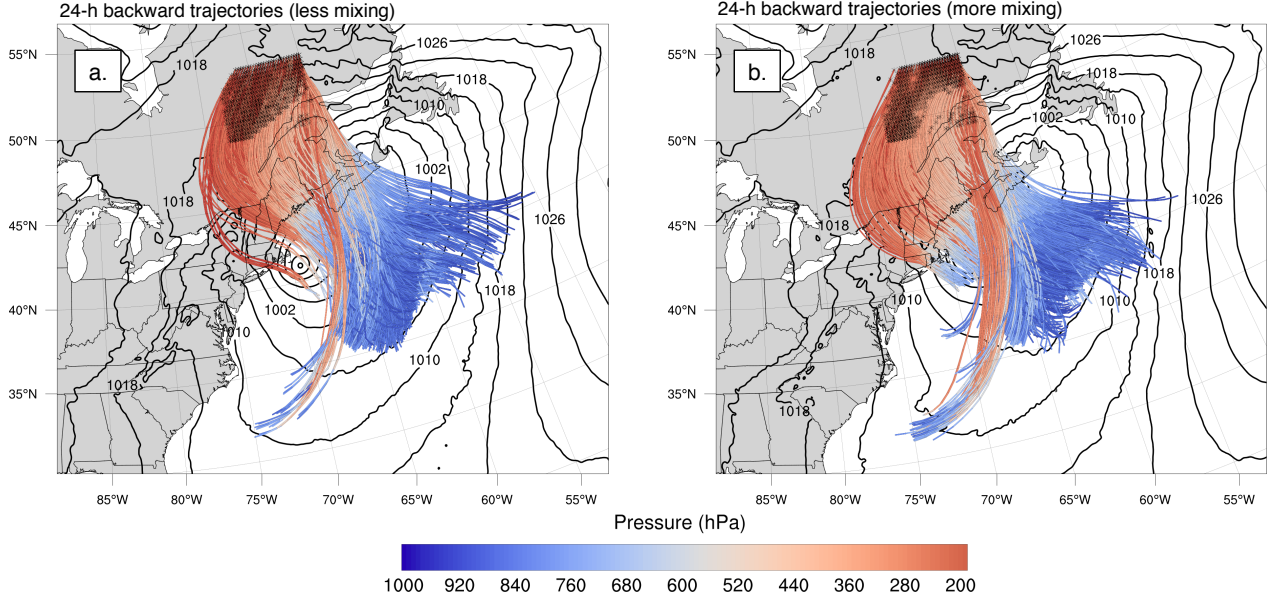


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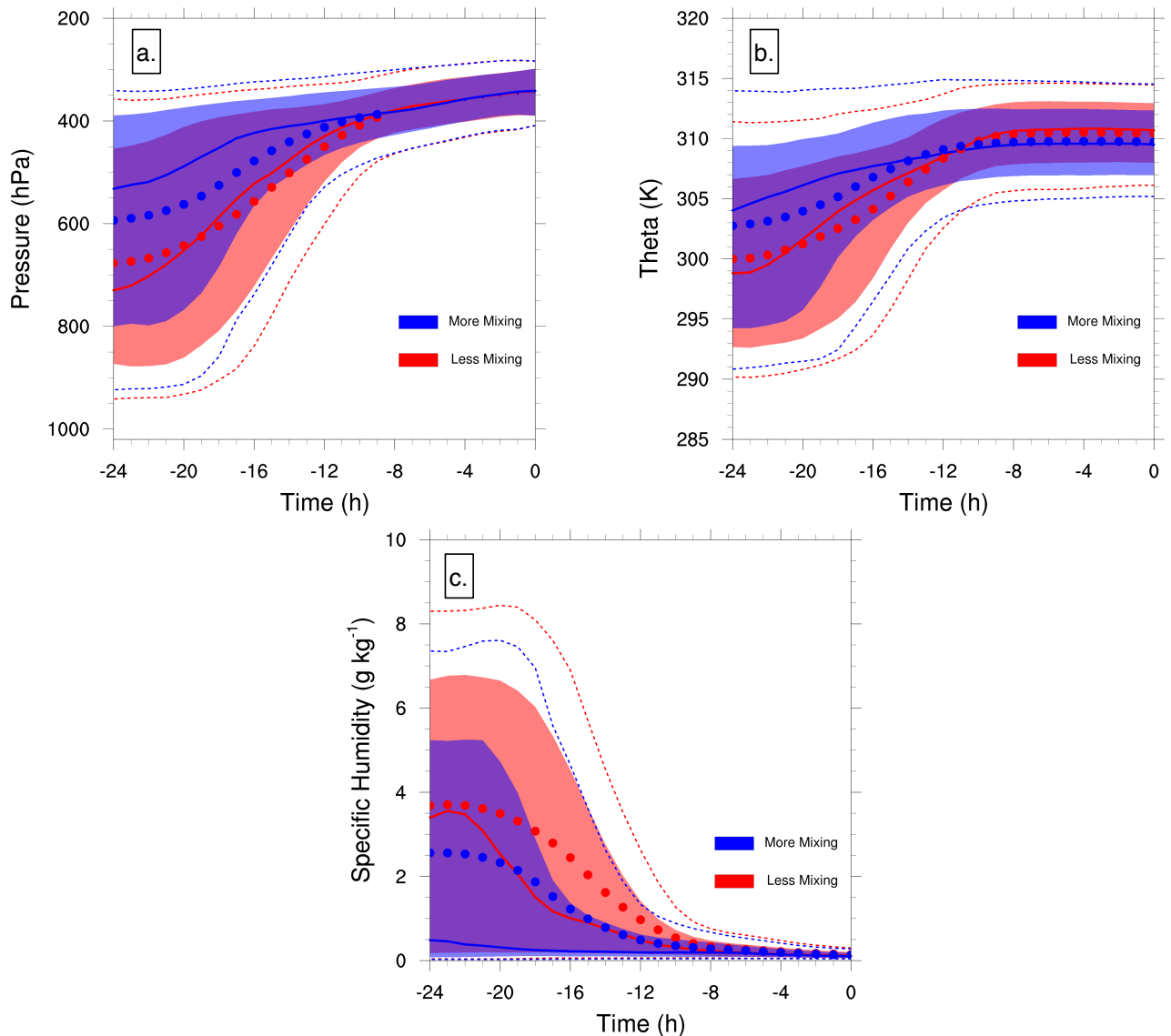


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