Poisson boundaries of II_1 factors

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Abstract

We introduce Poisson boundaries of II_1 factors with respect to density operators that give the traces. The Poisson boundary is a von Neumann algebra that contains the II_1 factor and is a particular example of the boundary of a unital completely positive map as introduced by Izumi. Studying the inclusion of the II_1 factor into its boundary we develop a number of notions, such as double ergodicity and entropy, that can be seen as natural analogues of results regarding the Poisson boundaries introduced by Furstenberg. We use the techniques developed to answer a problem of Popa by showing that all finite factors satisfy the MV-property. We also extend a result of Nevo by showing that property (T) factors give rise to an entropy gap.

1. Introduction

Given a locally compact group G and a probability measure $\mu \in \mathrm{Prob}(G)$, the associated (left) random walk on G is the Markov chain on G whose transition probabilities are given by the measures $\mu * \delta_x$. The Markov operator associated to this random walk is given by

$$\mathcal{P}_\mu(f)(x) = \int f(gx)\, d\mu(g),$$

where f is a continuous function on G with compact support. The Markov operator extends to a contraction on $L^\infty(G)$, which is unital and (completely) positive. A function $f \in L^\infty(G)$ is μ-harmonic if $\mathcal{P}_\mu(f) = f$. We let $\mathrm{Har}(G, \mu)$ denote the Banach space of μ-harmonic functions. The Furstenberg-Poisson boundary [Fur63b] of G with respect to μ is a certain G-probability space (B, ζ), such that we have a natural positivity preserving isometric G-equivariant identification of $L^\infty(B, \zeta)$ with $\mathrm{Har}(G, \mu)$ via a Poisson transform.

An actual construction of the Poisson boundary (B, ζ), which is often described as a quotient of the path space corresponding to the stationary σ-algebra, is less important to us here as its existence, and indeed, up to isomorphisms of G-spaces, it is the unique G-probability space such that $L^\infty(B, \zeta)$ is isomorphic, as an operator G-space, to $\mathrm{Har}(G, \mu)$.

Under natural conditions on the measure μ, the boundary (B, ζ) possesses a number of remarkable properties. It is an amenable G-space [Zim78], it is doubly ergodic with isometric coefficients [Kai92] [GW16], and it is strongly asymptotically transitive [Jaw94, Jaw95]. The boundary has therefore become a powerful tool for studying rigidity properties for groups and their probability measure preserving actions [Mar75, Zim80, BS06, BM02, BF20].

In light of the successful application of the Poisson boundary to rigidity properties in group theory, Alain Connes suggested (see [Jon00]) that developing a theory of the Poisson boundary

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in the setting of operator algebras would be the first step toward studying his rigidity conjecture [Con82], which states that two property (T) icc groups have isomorphic group von Neumann algebras if and only if the groups themselves are isomorphic. Further evidence for this can be seen by the significant role that Poisson boundaries play in [CP17, CP13, Pet15], where a related rigidity conjecture of Connes was investigated.

Poisson boundaries can more generally be defined using any Markov operator associated to a random walk. Markov operators are particular examples of normal unital completely positive (u.c.p.) maps on von Neumann algebras, and motivated by defining Poisson boundaries for discrete quantum groups, Izumi in [Izu02, Izu04] was able to define a non-commutative Poisson boundary associated to any normal u.c.p. map on a general von Neumann algebra. Specifically, if $\mathcal{M}$ is a von Neumann algebra and $\phi : \mathcal{M} \to \mathcal{M}$ is a normal u.c.p. map, then we let $\mathrm{Har}(\phi) = \{x \in \mathcal{M} \mid \phi(x) = x\}$ denote the space of ϕ-harmonic operators. Izumi showed that there exists a (unique up to isomorphism) von Neumann algebra $\mathcal{B}_\phi$ such that, as operator systems, $\mathrm{Har}(\phi)$ and $\mathcal{B}_\phi$ can be identified via a Poisson transform $\mathcal{P} : \mathcal{B}_\phi \to \mathrm{Har}(\phi)$. The existence of this boundary follows by showing that $\mathrm{Har}(\phi)$ can be realized as the range of a u.c.p. idempotent on $\mathcal{M}$ and then applying a theorem of Choi and Effros. Alternatively, the existence of the boundary follows by considering the minimal dilation of ϕ [Izu12]. We include in the appendix of this paper an elementary proof based on this perspective.

There is a well-known dictionary between many analytic notions in group theory and those in von Neumann algebras. For example, states on $\mathcal{B}(L^2(M))$ correspond to states on $\ell^\infty\Gamma$, normal Hilbert M-bimodules correspond to unitary representations, etc., [Con76b, Section 2] [Con80]. This allows one to develop notions such as amenability, property (T), etc., in the setting of finite von Neumann algebras. While Izumi's boundary gives a satisfactory non-commutative analogue of the Poisson boundary associated to a general random walk, still missing is an appropriate notion of a non-commutative Poisson boundary analogous to the group setting.

The main goal of this article is to introduce a theory of Poisson boundaries for finite von Neumann algebras that we believe will fill the role envisioned by Connes. If M is a finite von Neumann algebra with a normal faithful trace τ, and if $\varphi \in \mathcal{B}(L^2(M,\tau))_*$ is a normal state such that $\varphi_{|M} = \tau$, then we will view φ as the distribution of a "non-commutative random walk" on M. To each distribution we associate a corresponding "convolution operator", which is a normal u.c.p. map $\mathcal{P}_\varphi : \mathcal{B}(L^2(M,\tau)) \to \mathcal{B}(L^2(M,\tau))$, such that $M \subset \mathrm{Har}(\mathcal{P}_\varphi)$. We then define the Poisson boundary of M with respect to φ to be Izumi's non-commutative boundary $\mathcal{B}_\varphi$ associated to $\mathcal{P}_\varphi$; more precisely the boundary is really the inclusion of von Neumann algebras $M \subset \mathcal{B}_\varphi$, together with the Poisson transform $\mathcal{P} : \mathcal{B}_\varphi \to \mathrm{Har}(\mathcal{P}_\varphi)$.

Poisson boundaries of groups give rise to natural Poisson boundaries of group von Neumann algebras. Indeed, as was already noticed by Izumi in [Izu12], if Γ is a countable discrete group and $\mu \in \mathrm{Prob}(\Gamma)$, then the non-commutative boundary of the u.c.p. map $\phi_\mu : \mathcal{B}(\ell^2\Gamma) \to \mathcal{B}(\ell^2\Gamma)$ given by $\phi_\mu(T) = \int \rho_\gamma T \rho_\gamma^* \, d\mu(\gamma)$ is naturally isomorphic to the von Neumann crossed-product $L^\infty(B,\zeta) \rtimes \Gamma$ where (B,ζ) is the Poisson boundary of (Γ,μ). Thus, many of the results we obtain are not merely analogues, but are actually generalizations of results from the theory of random walks on groups.

If M is a finite factor, then under natural conditions on the distribution φ, e.g., that its "support" should generate M, we show that the boundary $\mathcal{B}_\varphi$ is amenable/injective (Proposition 2.4), and that the inclusion $M \subset \mathcal{B}_\varphi$ is "ergodic", i.e., $M' \cap \mathcal{B}_\varphi = \mathbf{C}$ (Proposition 2.7). We use techniques of Foguel [Fog75] to obtain equivalent characterizations for when the boundary is trivial (Theorem 2.10). The double ergodicity result of Kaimanovich [Kai92] is more subtle, as unlike in the case for groups, there is no natural "diagonal" inclusion of M into $\mathcal{B}_\varphi \overline{\otimes} \mathcal{B}_\varphi$. There is however a natural notions of left and right convolution operators, so that we may naturally associate with φ a second u.c.p. map $\mathcal{P}_\varphi^{\mathrm{o}}$ which commutes with $\mathcal{P}_\varphi$ (see Section 3 for the precise definition of $\mathcal{P}_\varphi^{\mathrm{o}}$). We may then show that bi-harmonic operators are constant, a result which is equivalent to double ergodicity in the group setting.

Theorem A (Theorem 3.1 below) *Let M be a finite factor and suppose φ is as above. Then we have*

$$\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) \cap \mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi^{\mathrm{o}}) = \mathbf{C}.$$

Motivated by the question of determining whether or not $L\mathbf{F}_\infty$ is finitely generated, Popa studied in [Pop18] the class of separable II_1 factors M that are tight, i.e., M contains two hyperfinte subfactors $L, R \subset M$ such that L and R^{op} together generate $\mathcal{B}(L^2(M))$. He conjectures in Conjecture 5.1 of [Pop18] that if a factor M has the property that all amplifications M^t are singly generated then M is tight. He also notes that a tight factor M satisfies the MV-property, which states that for any operator $T \in \mathcal{B}(L^2(M))$ the weak closure of the convex hull of $\{u(JvJ)T(Jv^*J)u^* \mid u, v \in \mathcal{U}(M)\}$ intersects the scalars. Popa then asks in Problem 7.4 of [Pop19a] and Problem 6.3 in [Pop19b] if free group factors, or perhaps all finite factors have the MV-property. As a consequence of double ergodicity we are able to answer Popa's problem.

Theorem B (Theorem 3.3 below) *All finite factors have the MV-property.*

Other consequences of double ergodicity are that it allows us to show vanishing cohomology for sub-bimodules of the Poisson boundary (Theorem 3.5), it allows us to generalize rigidity results from [CP13] (Theorem 4.1), and it allows us to extend results of Bader and Shalom [BS06] identifying the Poisson boundary of a tensor product with the tensor product of the Poisson boundaries (Corollary 4.5).

We also introduce analogues of Avez's asymptotic entropy and Furstenberg's μ-entropy in the setting of von Neumann algebras (see Section 5 for these definitions). We show that the triviality of the Poisson boundary is equivalent to the vanishing of the Furstenberg entropy (Corollary 5.15). We also use entropy to extend a result of Nevo [Nev03] to the setting of von Neumann algebras, which shows that property (T) factors give rise to an "entropy gap".

Theorem C (Theorem 6.2 below) *Let M be a II_1 factor with property (T) generated by unitaries $u_1, \ldots, u_n$. Define the state $\varphi \in \mathcal{B}(L^2M)_*$ by $\varphi(T) = \frac{1}{n}\sum_{k=1}^n \langle T\hat{u_k}, \hat{u_k}\rangle$. There exists $c > 0$ such that if $M \subsetneq \mathcal{A}$ is an irreducible inclusion of von Neumann algebras and $\zeta \in \mathcal{A}_*$ is any faithful normal state such that $\zeta_{|M} = \tau$, then $h_\varphi(M \subset \mathcal{A}, \zeta) \geq c$.*

We end with an appendix where we construct Izumi's boundary of a u.c.p. map. Our approach is elementary, and has the advantage that it applies for general C^*-algebras. This level of generality has no doubt been known by experts, but we could not find this in the current literature.

2. Boundaries

2.1 Hyperstates and bimodular u.c.p. maps

Fix a tracial von Neumann algebra (M, τ), and suppose we have an embedding $M \subset \mathcal{A}$ where $\mathcal{A}$ is a C^*-algebra. We say a state $\varphi \in \mathcal{A}^*$ is a τ-hyperstate (or just a hyperstate if τ is fixed) if it extends τ. We denote by $\mathcal{S}_\tau(\mathcal{A})$ the convex set of all hyperstates on $\mathcal{A}$. To each a hyperstate φ we obtain a natural inclusion $L^2(M,\tau) \subset L^2(\mathcal{A},\varphi)$ induced from the map $x\hat{1} \mapsto x1_\varphi$ for $x \in M$. We let $e_M \in \mathcal{B}(L^2(\mathcal{A},\varphi))$ denote the orthogonal projection onto $L^2(M,\tau)$. We may then consider the unital completely positive (u.c.p.) map $\mathcal{P}_\varphi : \mathcal{A} \to \mathcal{B}(L^2(M,\tau))$, defined by

$$\mathcal{P}_\varphi(T) = e_M T e_M, \qquad T \in \mathcal{A}. \tag{1}$$

Note that if $x \in M \subset \mathcal{A}$ then we have $\mathcal{P}_\varphi(x) = x$. We call the map $\mathcal{P}_\varphi$ the Poisson transform (with respect to φ) of the inclusion $M \subset \mathcal{A}$.

The following proposition is inspired from [Con76b, Section 2.2].

Proposition 2.1 *The correspondence $\varphi \mapsto \mathcal{P}_\varphi$ defined by (1) gives a bijective correspondence between hyperstates on M, and u.c.p., M-bimodular maps from $\mathcal{A}$ to $\mathcal{B}(L^2(M,\tau))$. Moreover, if $\mathcal{A}$ is a von Neumann algebra, then $\mathcal{P}_\varphi$ is normal if and only if φ is normal.*

Also, this corresondence is a homeomorphism where the space of hyperstates is endowed with the weak-topology, and the space of u.c.p., M-bimodular maps with the topology of pointwise weak operator topology convergence.*

Proof. First note that if φ is a hyperstate on $\mathcal{A}$, then for all $T \in \mathcal{A}$ we have

$$\varphi(T) = \langle T, \hat{1}\rangle_\varphi = \langle \mathcal{P}_\varphi(T)\hat{1}, \hat{1}\rangle_\tau.$$

From this it follows that the correspondence $\varphi \mapsto \mathcal{P}_\varphi$ is one-to-one. To see that it is onto, suppose that $\mathcal{P} : \mathcal{A} \to \mathcal{B}(L^2(M,\tau))$ is u.c.p. and M-bimodular. We define a state φ on $\mathcal{A}$ by $\varphi(T) = \langle \mathcal{P}(T)\hat{1}, \hat{1}\rangle_\tau$. For all $y \in M$ we then have $\varphi(y) = \langle \mathcal{P}(y)\hat{1}, \hat{1}\rangle_\tau = \tau(y)$, hence φ is a hyperstate. Moreover, if $y, z \in M$, and $T \in \mathcal{A}$ then we have

$$\begin{aligned}\langle \mathcal{P}_\varphi(T)\hat{y}, \hat{z}\rangle_\tau &= \langle \mathcal{P}_\varphi(z^*Ty)\hat{1}, \hat{1}\rangle_\tau \\ &= \varphi(z^*Ty) = \langle \mathcal{P}(T)\hat{y}, \hat{z}\rangle_\tau,\end{aligned} \tag{2}$$

hence, $\mathcal{P}_\varphi = \mathcal{P}$.

It is also easy to check that $\mathcal{P}_\varphi$ is normal if and only if φ is.

To see that this correspondence is a homeomorphism when given the topologies above, suppose that φ is a hyperstate, and φ_α is a net of hyperstates. From (2) and the fact that u.c.p. maps are contractions in norm we see that $\mathcal{P}_{\varphi_\alpha}$ converges in the pointwise ultraweak topology to $\mathcal{P}_\varphi$ if φ_α converges weak* to φ. Conversely, setting $y = z = 1$ in (2) shows that if $\mathcal{P}_{\varphi_\alpha}$ converges in the pointwise ultraweak topology to $\mathcal{P}_\varphi$ then φ_α converges weak* to φ. □

Considering the case $\mathcal{A} = \mathcal{B}(L^2(M,\tau))$ we see that to each hyperstate φ on $\mathcal{B}(L^2(M,\tau))$ we obtain a u.c.p. M-bimodular map $\mathcal{P}_\varphi$ on $\mathcal{B}(L^2(M,\tau))$. In particular, composing such maps gives a type of convolution operation on the space of hyperstates. More generally, if $\mathcal{A}$ is a C^*-algebra, with $M \subset \mathcal{A}$, then for hyperstates $\psi \in \mathcal{A}^*$, and $\varphi \in \mathcal{B}(L^2(M,\tau))^*$ we define the convolution $\varphi * \psi$ to be the unique hyperstate on $\mathcal{A}$ such that

$$\mathcal{P}_{\varphi * \psi} = \mathcal{P}_\varphi \circ \mathcal{P}_\psi. \tag{3}$$

We say that ψ is φ-stationary if we have $\varphi * \psi = \psi$, or equivalently, if $\mathcal{P}_\psi$ maps into the space of $\mathcal{P}_\varphi$-harmonic operators

$$\mathrm{Har}(\mathcal{P}_\varphi) = \mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) = \{T \in \mathcal{B}(L^2(M,\tau)) \mid \mathcal{P}_\varphi(T) = T\}.$$

Lemma 2.2 *For a fixed $\psi \in \mathcal{S}_\tau(\mathcal{A})$ the mapping*

$$\mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau))) \ni \varphi \mapsto \varphi * \psi \in \mathcal{S}_\tau(\mathcal{A})$$

is continuous in the weak**-topology.*

Moreover, if $\varphi \in \mathcal{B}(L^2(M,\tau))_*$ *is a fixed normal hyperstate, then the mapping*

$$\mathcal{S}_\tau(\mathcal{A}) \ni \psi \mapsto \varphi * \psi \in \mathcal{S}_\tau(\mathcal{A})$$

is also weak**-continuous.*

Proof. By Proposition 2.1 the correspondence $\varphi \mapsto \mathcal{P}_\varphi$ is a homeomorphism from the weak*-topology to the topology of pointwise ultraweak convergence, this lemma then follows easily from (3). $\square$

2.2 Poisson boundaries of II_1 factors

Definition 2.3 *Let* $\varphi \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau))$ *be a hyperstate. We define the Poisson boundary* $\mathcal{B}_\varphi$ *of* M *with respect to* φ *to be the noncommutative Poisson boundary of the u.c.p. map* $\mathcal{P}_\varphi$ *as defined by Izumi [Izu02], i.e., the Poisson boundary* $\mathcal{B}_\varphi$ *is a* C^**-algebra (a von Neumann algebra when* φ *is normal) that is isomorphic, as an operator system, to the space of harmonic operators* $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$.

Since M is in the multiplicative domain of $\mathcal{P}_\varphi$, we see that $\mathcal{B}_\varphi$ contains M as a subalgebra. Moreover, note that if we have a C^*-algebra $\mathcal{B}$, an inclusion $M \subseteq \mathcal{B}$ together with a completely positive isometric surjection from $\mathcal{B}$ to $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$, then this induces a completely positive isometric surjection from $\mathcal{B}$ to $\mathcal{B}_\varphi$ which restricts to the identity on M. It's a well-known result of Choi [Cho72] that a completely positive surjective isometry between two C^*-algebras is a $*$-isomorphism. Thus,the Poisson boundary contains M as a subalgebra, and the inclusion $(M \subset \mathcal{B}_\varphi)$ is determined up to isomorphism by the property that there exists a completely positive isometric surjection $\mathcal{P} : \mathcal{B}_\varphi \to \mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$ which restricts to the identity map on M. We will always assume that $\mathcal{P}$ is fixed and we also call $\mathcal{P}$ the Poisson transform.

Given any initial hyperstate $\varphi_0 \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ we may consider the hyperstate given by $\varphi_0 \circ \mathcal{P}$ on $\mathcal{B}_\varphi$. Of particular interest is the state η on $\mathcal{B}_\varphi$ arising from the initial hyperstate $\varphi_0(x) \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ given by $\varphi_0(x) = \langle x\hat{1}, \hat{1}\rangle$, which we call the stationary state on $\mathcal{B}_\varphi$. In this case, using (2) above, it is easy to see that we have $\mathcal{P}_\eta = \mathcal{P}$, and hence $\varphi * \eta = \eta$.

Proposition 2.4 *Let* (M,τ) *be a tracial von Neumann algebra and let* φ *be a fixed hyperstate on* $\mathcal{B}(L^2(M,\tau))$*. Then the Poisson boundary* $\mathcal{B}_\varphi$ *is injective.*

Proof. If we take any accumulation point E of $\left\{\frac{1}{N}\sum_{n=1}^N \mathcal{P}_\varphi^n\right\}_{N\in\mathbf{N}}$ in the topology of pointwise ultraweak convergence, then $E : \mathcal{B}(L^2(M,\tau)) \to \mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$ gives a u.c.p. projection. As $\mathcal{B}_\varphi$ is isomorphic to $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$ as an operator system it then follows that $\mathcal{B}_\varphi$ is injective [CE77, Section 3]. $\square$

The trivial case is when $\varphi_e(x) = \langle x\hat{1}, \hat{1}\rangle_\tau$ in which case we have that $\mathcal{P}_{\varphi_e} = \mathrm{id}$, and the Poisson boundary is nothing but $\mathcal{B}(L^2(M,\tau))$. Note that φ_e gives an identity with respect to convolution. Also note that if $\varphi \in \mathcal{B}(L^2(M,\tau))^*$ is a hyperstate, then we have a description of the space of harmonic operators as:

$$\mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi) = \{T \in \mathcal{B}(L^2(M,\tau)) \mid \varphi(aTb) = \varphi_e(aTb) \text{ for all } a,b \in M\}.$$

Since $\mathcal{P}_\varphi$ is M-bimodular it follows that $\mathcal{P}_\varphi(M') \subset M'$. We say that φ is regular if the restriction of $\mathcal{P}_\varphi$ to M' preserves the canonical trace on M', and we say that φ is generating if M is the largest $*$-subalgebra of $\mathcal{B}(L^2(M,\tau))$ which is contained in $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)),\mathcal{P}_\varphi)$. If φ

is regular, then the conjugate of φ which is given by $\varphi^*(T) = \varphi(JT^*J)$, is again a hyperstate. We'll say that φ is symmetric if it is regular and we have $\varphi^* = \varphi$.

Regular, generating, symmetric hyperstates are easy to find. Suppose (M, τ) is a separable finite von Neumann algebra with a faithful normal trace τ. We consider the unit ball $(M)_1$ of M as a Polish space endowed with the strong operator topology, and suppose we have a σ-finite measure μ on $(M)_1$ such that $\int x^*x \, d\mu(x) = 1$. We obtain a normal hyperstate as

$$\varphi(T) = \int \langle T\widehat{x^*}, \widehat{x^*} \rangle \, d\mu(x) \tag{4}$$

and using (2) we may explicitly compute the Poisson transform $\mathcal{P}_\varphi$ on $\mathcal{B}(L^2(M, \tau))$ as

$$\mathcal{P}_\varphi(T) = \int (Jx^*J)T(JxJ) \, d\mu(x).$$

Proposition 2.5 *Consider φ as given by (4), then*

(i) φ is generating if and only if the support of μ generates M as a von Neumann algebra.

(ii) φ is regular if and only if $\int xx^* \, d\mu(x) = 1$. In this case φ^* is a normal hyperstate.

(iii) If φ is regular then $\mathcal{P}_{\varphi^*}(T) = \int (JxJ)T(Jx^*J) \, d\mu(x)$ and φ is symmetric if $J_*\mu = \mu$, where J is the adjoint operation.

Proof. If the support of μ generates von Neumann algebra $M_0 \subset M$ such that $M_0 \neq M$, then we have $[JxJ, e_{M_0}] = [Jx^*J, e_{M_0}] = 0$ for each x in the support of μ. Hence, $\mathcal{P}_\varphi(T) = \int (JxJ)T(Jx^*J) \, d\mu(x) = T$, for each T in the $*$-algebra generated by M and e_{M_0}. Therefore, φ is not generating. On the other hand, if $T \in \mathrm{Har}(\mathcal{B}(L^2(M, \tau)), \mathcal{P}_\varphi)$ is such that we also have $T^*T, TT^* \in \mathrm{Har}(\mathcal{B}(L^2(M, \tau)), \mathcal{P}_\varphi)$ then for each $a \in M$ we have

$$\begin{aligned}\int \|((JxJ)T - T(JxJ))\hat{a}\|_2^2 \, d\mu(x) \\ = \langle (T^*\mathcal{P}_\varphi(1)T - \mathcal{P}_\varphi(T^*)T - T^*\mathcal{P}_\varphi(T) + \mathcal{P}_\varphi(T^*T))\hat{a}, \hat{a} \rangle = 0,\end{aligned}$$

and by symmetry we also have $\int \|((JxJ)T^* - T^*(JxJ))\hat{a}\|_2^2 \, d\mu(x) = 0$. Hence, $[JxJ, T] = [Jx^*J, T] = 0$ for μ-almost every $x \in (M)_1$. Therefore, if the support of μ generates M as a von Neumann algebra then we then have that $T \in JMJ' = M$, showing that φ is generating, thereby proving (i).

If $y \in M$ then we have $\mathcal{P}_\varphi(JyJ) = \int Jx^*yxJ \, d\mu(x)$. Hence, we see that φ is regular if and only if for all $y \in M$ we have $\tau(y) = \int \tau(x^*yx) \, d\mu(x) = \int \tau(xx^*y) \, d\mu(x)$, which is if and only if $\int xx^* \, d\mu(x) = 1$, thereby proving (ii).

If φ is regular then

$$\begin{aligned}\varphi^*(T) = \varphi(JT^*J) &= \int \langle JT^*J\widehat{x^*}, \widehat{x^*} \rangle \, d\mu(x) \\ &= \int \langle \hat{x}, T^*\hat{x} \rangle \, d\mu(x) = \int \langle T\widehat{x^*}, \widehat{x^*} \rangle \, dJ_*\mu(x).\end{aligned}$$

Therefore, if $J_*\mu = \mu$ then φ is symmetric, thereby proving (iii). □

Given a unital C^*-algebra A, and a u.c.p. map $\mathcal{P} : A \to A$, we denote the set of fixed points of $\mathcal{P}$ by $\mathrm{Har}(A, \mathcal{P})$. That is, $\mathrm{Har}(A, \mathcal{P}) = \{a \in A : \mathcal{P}(a) = a\}$. The following Lemma is well known, see, e.g., [FNW94], [BJKW00, Lemma 3.4], or [CD20, Lemma 3.1] . We include a proof for the convenience of the reader.

Lemma 2.6 *Suppose A is a unital C^*-algebra with a faithful state φ. If $\mathcal{P} : A \to A$ is a u.c.p. map such that $\varphi \circ \mathcal{P} = \varphi$, then $\mathrm{Har}(A, \mathcal{P}) \subset A$ is a C^*-subalgebra.*

Proof. $\mathrm{Har}(A, \mathcal{P})$ is clearly a self-adjoint closed subspace, thus we must show that $\mathrm{Har}(A, \mathcal{P})$ is an algebra. By the polarization identity it is enough to show that $x^*x \in \mathrm{Har}(A, \mathcal{P})$ whenever $x \in \mathrm{Har}(A, \mathcal{P})$. Suppose $x \in \mathrm{Har}(A, \mathcal{P})$. By Kadison's indequality we have $\mathcal{P}(x^*x) - x^*x = \mathcal{P}(x^*x) - \mathcal{P}(x^*)\mathcal{P}(x) \geq 0$. Also, $\varphi(\mathcal{P}(x^*x) - x^*x) = 0$ so that by faithfulness of φ we have $\mathcal{P}(x^*x) = x^*x$. □

Proposition 2.7 *Let M be a finite von Neumann algebra with a normal faithful trace τ. Let $\varphi \in \mathcal{B}(L^2(M, \tau))^*$ be a regular generating hyperstate, and let $\mathcal{B}_\varphi$ be the corresponding Poisson boundary, then $M' \cap \mathcal{B}_\varphi = \mathcal{Z}(M)$. In particular, if φ is also normal and M is a factor then $\mathcal{B}_\varphi$ is also a von Neumann factor.*

Proof. Let $\mathcal{P} : \mathcal{B}_\varphi \to \mathrm{Har}(\mathcal{B}(L^2(M, \tau)), \mathcal{P}_\varphi)$ denote the Poisson transform. If $x \in M' \cap \mathcal{B}_\varphi$, then $\mathcal{P}(x) \in M' \cap \mathcal{B}(L^2(M, \tau)) = JMJ$. Since φ is regular, $\mathcal{P}_\varphi$ preserves the trace when restricted to JMJ. Thus, $\mathrm{Har}(JMJ, \mathcal{P}_\varphi)$ is a von Neumann subalgebra of JMJ by Lemma 2.6. Since φ is generating, M is the largest von Neumann subalgebra of $\mathrm{Har}(\mathcal{B}(L^2(M, \tau))$, and hence $\mathrm{Har}(JMJ, \mathcal{P}_\varphi) \subseteq M$, implying that $\mathrm{Har}(JMJ, \mathcal{P}_\varphi) = \mathcal{Z}(M)$. Therefore, $\mathcal{P}(x) \in \mathrm{Har}(JMJ, \mathcal{P}_\varphi) = \mathcal{Z}(M)$, and hence $x \in \mathcal{Z}(M)$ since $\mathcal{P}$ is injective. □

If φ is a normal hyperstate in $\mathcal{S}_\tau(\mathcal{B}(L^2(M, \tau)))$, then $\mathcal{P}_\varphi : \mathcal{B}(L^2(M, \tau)) \to \mathcal{B}(L^2(M, \tau))$ is a normal map, and hence the dual map $\mathcal{P}_\varphi^*$ preserves the predual of $\mathcal{B}(L^2(M, \tau))$ which we identify with the space of trace-class operators.

We let $A_\varphi \in \mathcal{B}(L^2(M, \tau))$ denote the density operator associated with φ, i.e., A_φ is the unique trace-class operator so that $\varphi(T) = \mathrm{Tr}(A_\varphi T)$ for all $T \in \mathcal{B}(L^2(M, \tau))$. Since φ is positive we have that A_φ is a positive operator. If $P_{\hat{1}}$ denotes the rank one orthogonal projection onto $\mathbf{C}\hat{1}$, then we have $\varphi(T) = \langle \mathcal{P}_\varphi(T)\hat{1}, \hat{1} \rangle = \mathrm{Tr}(\mathcal{P}_\varphi(T)P_{\hat{1}})$, and hence we see that $A_\varphi = \mathcal{P}_\varphi^*(P_{\hat{1}})$. In particular we have that $A_{\varphi^{*n}} = (\mathcal{P}_\varphi^n)^*(P_{\hat{1}})$ for $n \geq 1$.

Proposition 2.8 *Let (M, τ) be a tracial von Neumann algebra and let $\varphi \in \mathcal{S}_\tau(\mathcal{B}(L^2(M, \tau)))$ be a normal hyperstate, then there exists a τ-orthogonal family $\{z_n\}_n$ which gives a partition of the identity as $1 = \sum_n z_n^* z_n$ so that*

$$\mathcal{P}_\varphi(T) = \sum_n (J z_n^* J) T (J z_n J)$$

for all $T \in \mathcal{B}(L^2(M, \tau))$.

Moreover, if $\{\tilde{z}_m\}_m$ is a τ-orthogonal family which gives a partition of the identity as $1 = \sum_n \tilde{z}_n^* \tilde{z}_n$, then the map $\sum_m (J \tilde{z}_m^* J) T (J \tilde{z}_m J)$ agrees with $\mathcal{P}_\varphi$ if and only if for each $t > 0$ we have

$$\mathrm{sp}\{z_n \mid \|z_n\|_2 = t\} = \mathrm{sp}\{\tilde{z}_m \mid \|\tilde{z}_m\|_2 = t\}.$$

Proof. Since A_φ is a positive trace-class operator we may write $A_\varphi = \sum_n a_n P_{y_n}$ where $a_1, a_2, \ldots$ are positive and $\{y_n\}_n$ is an orthonormal family with P_{y_n} denoting the rank one projection onto $\mathbf{C}y_n$. For $T \in \mathcal{B}(L^2(M, \tau))$ we then have

$$\mathrm{Tr}(T A_\varphi) = \sum_n a_n \langle T y_n, y_n \rangle.$$

Taking $T = x^*x \in M$ we have $a_n \|xy_n\|_2^2 \leq \mathrm{Tr}(x^*xA_\varphi) = \|x\|_2^2$, so that $y_n \in M \subset L^2(M,\tau)$ for each n. Hence, for $T \in \mathcal{B}(L^2(M,\tau))$ we have

$$\mathrm{Tr}(\mathcal{P}_\varphi(T)P_{\hat{1}}) = \mathrm{Tr}(TA_\varphi) = \left\langle \sum_n a_n (Jy_nJ)T(Jy_n^*J)\hat{1}, \hat{1} \right\rangle$$
$$= \mathrm{Tr}\left(\left(\sum_n a_n (Jy_nJ)T(Jy_n^*J)\right) P_{\hat{1}}\right).$$

Since $\mathcal{P}_\varphi$ is M-bimodular and since $Jy_nJ \in M'$ it follows that for all $x, y \in M$ we have

$$\mathrm{Tr}(\mathcal{P}_\varphi(T)xP_{\hat{1}}y) = \mathrm{Tr}\left(\left(\sum_n a_n (Jy_nJ)T(Jy_n^*J)\right) xP_{\hat{1}}y\right).$$

In particular, setting $T = y = 1$ we have

$$\tau(x) = \sum_n a_n \tau(y_n^* y_n x),$$

which shows that $\sum_n a_n y_n^* y_n = 1$.

Since the span of operators of the form $xP_{\hat{1}}y$ is dense in the space of trace-class operators it then follows that $\mathcal{P}_\varphi(T) = \sum_n a_n (Jy_nJ)T(Jy_n^*J)$ for all $T \in \mathcal{B}(L^2(M,\tau))$. Setting $z_n = \sqrt{a_n} y_n^*$ then finishes the existence part of the proposition.

Suppose now that $\{\tilde{z}_m\}_m$ is a τ-orthogonal family which gives a partition of the identity $1 = \sum_n \tilde{z}_n^* \tilde{z}_n$, and set $\tilde{\varphi}(T) = \mathrm{Tr}((\sum_n (J\tilde{z}_n^*J)T(J\tilde{z}_nJ))P_{\hat{1}})$. Then, the density matrix $A_{\tilde{\varphi}}$, corresponding to $\tilde{\varphi}$, is given by $A_{\tilde{\varphi}} = \sum_n \tilde{z}_n^* P_{\hat{1}} \tilde{z}_n$. Since $\{\tilde{z}_n\}_n$ forms a τ-orthogonal family it then follows easily that $\tilde{z}_n^*$ is an eigenvector for $A_{\tilde{\varphi}}$, and the corresponding eigenvalue is $\|\tilde{z}_n^*\|_2^2 = \|\tilde{z}_n\|_2^2$.

Using our notation from the first part of the proof of the proposition, we have that $A_\varphi = \sum_n z_n^* P_{\hat{1}} z_n$. By the same argument as above, we get that z_n^* is an eigenvector for A_φ, and the corresponding eigenvalue is $\|z_n^*\|_2^2 = \|z_n\|_2^2$. Note that $\mathcal{P}_\varphi = \mathcal{P}_{\tilde{\varphi}}$ if and only if $A_\varphi = A_{\tilde{\varphi}}$. Since the corresponding density matrices are positive trace class operators, the moreover part of the proposition follow easily from the Spectral Theorem. □

We say that the form $\mathcal{P}_\varphi(T) = \sum_n (Jz_n^*J)T(Jz_nJ)$ (resp. $\varphi(T) = \sum_n \langle T\widehat{z_n^*}, \widehat{z_n^*}\rangle$) is a standard form for $\mathcal{P}_\varphi$ (resp. φ). It follows from Proposition 2.5 that φ is generating if and only if $\{z_n\}_n$ generates M as a von Neumann algebra. We say that φ is strongly generating if the unital algebra (rather than the unital $*$-algebra) generated by $\{z_n\}_n$ is already weakly dense in M. This is the case, for example, if φ is generating and symmetric, since then we have that $\{z_n\}_n = \{z_n^*\}_n$, and hence the unital algebra generated by $\{z_n\}_n$ is already a $*$-algebra.

Proposition 2.9 *Let (M,τ) be a tracial von Neumann algebra and suppose φ is a normal strongly generating hyperstate, then the stationary state $\zeta = \varphi \circ \mathcal{P}$ gives a normal faithful state on the Poisson boundary $\mathcal{B}_\varphi$ such that $\zeta_{|M} = \tau$.*

Proof. By considering the Poisson transform $\mathcal{P}$, it suffices to show that φ is normal and faithful on the operator system $\mathrm{Har}(\mathcal{P}_\varphi)$. Note that here the stationary state is a vector state and hence normality follows. To see that the state is faithful fix $T \in \mathrm{Har}(\mathcal{P}_\varphi)$, with $T \geq 0$ and $\langle T\hat{1}, \hat{1}\rangle = 0$. Let $\mathcal{P}_\varphi(S) = \sum_n (Jz_n^*J)S(Jz_nJ)$ be the standard form of $\mathcal{P}_\varphi$. Since $T \in \mathrm{Har}(\mathcal{P}_\varphi)$, we have that $\mathcal{P}_\varphi^k(T) = T$, for each $k \in \mathbf{N}$. Expanding the standard form gives

$$0 = \langle T\hat{1}, \hat{1}\rangle = \langle P_\varphi^k(T)\hat{1}, \hat{1}\rangle = \sum_{n_1, n_2, \ldots, n_k} \langle T z_{n_1} z_{n_2} \cdots z_{n_k} \hat{1}, z_{n_1} z_{n_2} \cdots z_{n_k} \hat{1}\rangle.$$

We then have $T\hat{m} = 0$ for all m in the unital algebra generated by $\{z_n\}$, and as φ is strongly generating it then follows that $T = 0$. $\square$

We end this section by giving a condition for the boundary to be trivial. We denote the space of trace-class operators on $L^2(M,\tau)$ by $\mathrm{TC}(L^2(M,\tau))$. We also denote the trace-class norm on $\mathrm{TC}(L^2(M,\tau))$ by $\|\cdot\|_{\mathrm{TC}}$. We identify $\mathcal{B}(L^2(M,\tau))$ with $\mathrm{TC}(L^2(M,\tau))^*$ via the pairing $(A,T) \mapsto Tr(AT)$, where $A \in \mathrm{TC}(L^2(M,\tau))$, and $T \in \mathcal{B}(L^2(M,\tau))$.

Theorem 2.10 *Let (M,τ) be a tracial von Neumann algebra and let ψ be a normal hyperstate. set $\varphi = \frac{1}{2}\psi + \frac{1}{2}\langle \cdot \hat{1}, \hat{1}\rangle$ and let $A_n \in \mathrm{TC}(L^2(M,\tau))$ denote the density matrix corresponding to the normal, u.c.p. M-bimodular map $\mathcal{P}_\varphi^n$. Then the following conditions are equivalent*

(i) *For all $x \in M$ we have $\|xA_n - A_n x\|_{\mathrm{TC}} \to 0$.*

(ii) *For all $x \in M$ we have $xA_n - A_n x \to 0$ weakly.*

(iii) $\mathrm{Har}(\mathcal{P}_\varphi) = M$

Proof. The first condition trivially implies the second. To see that the second implies the third suppose for each $x \in M$ we have $xA_n - A_n x \to 0$ weakly as $n \to \infty$. Let $T \in \mathrm{Har}(\mathcal{P}_\varphi)$. Let $x, a, b \in M$. Then taking inner products in $L^2(M,\tau)$ we have

$$\begin{aligned}|\langle (TJxJ - JxJT)a\hat{1}, b\hat{1}\rangle| &= |\langle (b^*Tax^* - x^*b^*Ta)\hat{1}, \hat{1}\rangle| \\ &= |\langle \mathcal{P}_\varphi^n(b^*Tax^* - x^*b^*Ta)\hat{1}, \hat{1}\rangle| = |\mathrm{Tr}(A_n(b^*Tax^* - x^*b^*Ta))| \\ &= |\mathrm{Tr}((x^*A_n - A_n x^*)b^*Ta)| \to 0.\end{aligned}$$

Hence $T \in JMJ' = M$.

To see that the third condition implies the first we adapt the approach of Foguel from [Fog75]. Suppose $\mathrm{Har}(\mathcal{P}_\varphi) = M$. Set $\mathcal{A}_0 = \{A \in \mathrm{TC}(L^2(M,\tau)) \mid \|(\mathcal{P}_\varphi^n)^*(A)\|_{\mathrm{TC}} \to 0\}$. Note that since $(\mathcal{P}_\varphi^n)^*$ is a contraction in the trace-class norm we have that $\mathcal{A}_0$ is a closed subspace.

Since $\varphi = \frac{1}{2}\psi + \frac{1}{2}\langle \cdot\hat{1}, \hat{1}\rangle$ we have $\mathcal{P}_\varphi^* = \frac{1}{2}\mathrm{id} + \frac{1}{2}\mathcal{P}_\psi^*$ and we compute

$$\begin{aligned}(\mathcal{P}_\varphi^n)^*(\mathrm{id} - \mathcal{P}_\varphi^*) &= 2^{-(n+1)}\left(\sum_{k=0}^{n}\binom{n}{k}(\mathcal{P}_\psi^k)^*\right)(\mathrm{id} - \mathcal{P}_\psi^*) \\ &= 2^{-(n+1)}\sum_{k=1}^{n}\left(\binom{n}{k-1} - \binom{n}{k}\right)\mathcal{P}_\psi^*.\end{aligned}$$

We have $\lim_{n\to\infty} 2^{-(n+1)}\sum_{k=1}^n |\binom{n}{k-1} - \binom{n}{k}| = 0$ (see (1.8) in [OS70]) hence $\|(\mathcal{P}_\varphi^n)^*(P_{\hat{1}} - \mathcal{P}_\varphi^*(P_{\hat{1}}))\|_{\mathrm{TC}} \to 0$. Thus $P_{\hat{1}} - \mathcal{P}_\varphi^*(P_{\hat{1}}) \in \mathcal{A}_0$.

Since $\mathcal{P}_\varphi^*$ is M-bimodular we then have that $aP_{\hat{1}}b - \mathcal{P}_\varphi^*(aP_{\hat{1}}b) \in \mathcal{A}_0$ for each $a, b \in M$ and hence $B - \mathcal{P}_\varphi^*(B) \in \mathcal{A}_0$ for all $B \in \mathrm{TC}(L^2(M,\tau))$. If $T \in \mathcal{B}(L^2(M,\tau))$ is such that $\mathrm{Tr}(AT) = 0$ for all $A \in \mathcal{A}_0$, then for all $B \in \mathrm{TC}(L^2(M,\tau))$ we have $\langle B - \mathcal{P}_\varphi^*(B), T\rangle = 0$ so that $T \in \mathrm{Har}(\mathcal{P}_\varphi) = M$. Hence the annihilator of $\mathcal{A}_0$ is contained in M. So the pre-annihilator of M must be contained in $\mathcal{A}_0$. Thus $A \in \mathcal{A}_0$ whenever $\mathrm{Tr}(Ax) = 0$ for all $x \in M$. In particular, we have $xP_{\hat{1}} - P_{\hat{1}}x \in \mathcal{A}_0$ for all $x \in M$, which is equivalent to the fact that $\|xA_n - A_n x\|_{\mathrm{TC}} \to 0$ for each $x \in M$. $\square$

3. Biharmonic operators

If $\varphi \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ is regular and normal then we define $\mathcal{P}_\varphi^{\mathrm{o}}$ to be the u.c.p. map given by $\mathcal{P}_\varphi^{\mathrm{o}} = \mathrm{Ad}(J) \circ \mathcal{P}_{\varphi^*} \circ \mathrm{Ad}(J)$. Note that $\mathcal{P}_\varphi^{\mathrm{o}}$ and $\mathcal{P}_\eta$ commute for any normal hyperstate η. Indeed, if we have standard forms $\mathcal{P}_\varphi(T) = \sum_n (Jz_n^*J)T(Jz_nJ)$ and $\mathcal{P}_\eta(T) = \sum_m (Jy_m^*J)T(Jy_mJ)$ then by Proposition 2.5 we have $\mathcal{P}_\varphi^{\rm o}(T) = \sum_n z_nTz_n^*$ and hence

$$\mathcal{P}_\varphi^{\rm o} \circ \mathcal{P}_\eta(T) = \mathcal{P}_\eta \circ \mathcal{P}_\varphi^{\rm o}(T) = \sum_{n,m} z_n(Jy_m^*J)T(Jy_mJ)z_n^*.$$

The following is a noncommutative analogue of double ergodicity which was established in [Kai92].

Theorem 3.1 *Let (M,τ) be a tracial von Neumann algebra and let φ be a normal regular strongly generating hyperstate. Then*

$$\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) \cap \mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi^{\rm o}) = \mathcal{Z}(M).$$

Proof. We fix a standard form $\mathcal{P}_\varphi(T) = \sum_n (Jz_n^*J)T(Jz_nJ)$, so that we also have $\mathcal{P}_\varphi^{\rm o}(T) = \sum_m z_mTz_m^*$. We identify the Poisson boundary $\mathcal{B}_\varphi$ with $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi)$, and let ζ denote the stationary state on $\mathcal{B}_\varphi$, which is faithful by Proposition 2.9. For $T \in \mathcal{B}_\varphi$ we have

$$\zeta(\mathcal{P}_\varphi^{\rm o}(T)) = \langle \mathcal{P}_\varphi^{\rm o}(T)\hat{1}, \hat{1}\rangle = \langle \mathcal{P}_\varphi(T)\hat{1}, \hat{1}\rangle = \zeta(\mathcal{P}_\varphi(T)) = \zeta(T).$$

By Lemma 2.6 we then have that $B_0 = \mathrm{Har}(\mathcal{B}_\varphi, \mathcal{P}^{\rm o}_{\varphi|\mathcal{B}_\varphi})$ is a von Neumann subalgebra of $\mathcal{B}_\varphi$. If $p \in B_0$ is a projection and $\xi \in L^2(\mathcal{B}_\varphi, \zeta)$ then

$$\sum_n \|pz_n^*p^\perp\xi\|_2^2 = \sum_n \langle z_npz_n^*p^\perp\xi, p^\perp\xi\rangle = 0.$$

We must therefore have $\|pz_n^*p^\perp\xi\|_2 = 0$ for each n, and hence $pz_n^* = pz_n^*p$, for each n. Repeating this argument with roles of p and $p^\perp$ reversed shows that $z_n^*p = pz_n^*p$, so that $p \in M' \cap \mathcal{B}_\varphi$. Since p was an arbitrary projection we then have $B_0 \subset M' \cap \mathcal{B}_\varphi$ and by Proposition 2.7 we have $B_0 = \mathcal{Z}(M)$. $\square$

The previous result allows us to give an analogue of the classical Choquet-Deny theorem [CD60], which states that if Γ is an abelian group and $\mu \in \mathrm{Prob}(\Gamma)$ has support generating Γ then every bounded μ-harmonic function is constant.

Corollary 3.2 (The Choquet-Deny theorem) *Suppose M is an abelian von Neumann algebra and φ is a normal regular strongly generating hyperstate, then*

$$\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) = \mathcal{Z}(M) = M.$$

We will now describe how Theorem 3.1 leads to a positive answer of a recent question by Popa [Pop19a, Problem 7.4] [Pop19b, Problem 6.3].

Theorem 3.3 *Let M be a finite von Neumann algebra with a normal faithful trace τ and let $\mathcal{G} \subset \mathcal{U}(M)$ be a group which generates M as a von Neumann algebra. Then for any operator $T \in \mathcal{B}(L^2(M,\tau))$ the weak closure of the convex hull of $\{u(JvJ)T(Jv^*J)u^* \mid u,v \in \mathcal{G}\}$ intersects $\mathcal{Z}(M)$.*

Proof. We first consider the case when $\mathcal{G}$ is countable. Let $\mu \in \mathrm{Prob}(\mathcal{G})$ be symmetric with full support and define a normal regular symmetric generating hyperstate φ by $\varphi(T) = \int \langle T\hat{u}, \hat{u}\rangle \, d\mu(u)$. The corresponding Poisson transform is then given by $\mathcal{P}_\varphi(T) = \int (JuJ)T(Ju^*J)\, d\mu(u)$, and we may also compute $\mathcal{P}_\varphi^{\rm o}$ as $\mathcal{P}_\varphi^{\rm o}(T) = \int u^*Tu \, d\mu(u)$.

Fix $T \in \mathcal{B}(L^2(M,\tau))$ and let $\mathcal{C} = \overline{co}^{wk}\{u(JvJ)T(Jv^*J)u^* \mid u,v \in \mathcal{G}\}$. Then $\mathcal{C}$ is preserved by both $\mathcal{P}_\varphi$ and $\mathcal{P}_\varphi^{\rm o}$ and hence $\mathcal{C}$ is preserved by any point-ultraweak limit points E and $E^{\rm o}$ of

$\left\{\frac{1}{N}\sum_{n=1}^N \mathcal{P}_\varphi^n\right\}_{N=1}^\infty$ and $\left\{\frac{1}{N}\sum_{n=1}^N (\mathcal{P}_\varphi^{\mathrm{o}})^n\right\}_{N=1}^\infty$ respectively. Since $\mathcal{P}_\varphi$ and $\mathcal{P}_\varphi^{\mathrm{o}}$ commute we have that E and E^{o} commute. Moreover, as $\|\frac{1}{N}\sum_{n=1}^N \mathcal{P}_\varphi^n - \frac{1}{N}\sum_{n=1}^N \mathcal{P}_\varphi^{n+1}\| \leq 2/N$ it follows that $E : \mathcal{B}(L^2(M,\tau)) \to \mathrm{Har}(\mathcal{P}_\varphi)$ and similarly $E^{\mathrm{o}} : \mathcal{B}(L^2(M,\tau)) \to \mathrm{Har}(\mathcal{P}_\varphi^{\mathrm{o}})$. By Theorem 3.1 we then have $E^{\mathrm{o}} \circ E : \mathcal{B}(L^2(M,\tau)) \to \mathcal{Z}(M)$. Hence

$$E^{\mathrm{o}} \circ E(T) \in \mathcal{C} \cap \mathcal{Z}(M).$$

In the general case, if $G < \mathcal{G}$ is a countable subgroup, then let $N \subset M$ be the von Neumann subalgebra generated by G and let $e_N : L^2(M,\tau) \to L^2(N,\tau)$ be the orthogonal projection. If we define φ as above and set $T_G = E^{\mathrm{o}} \circ E(T)$ then we have $T_G \in \mathcal{C}$, $e_N T_G e_N = E^{\mathrm{o}} \circ E(e_N T e_N)$ and viewing $e_N T e_N$ as an operator in $\mathcal{B}(L^2(N,\tau))$ we may apply Theorem 3.1 as above to conclude that $e_N T_G e_N \in \mathcal{Z}(N) \subset \mathcal{B}(L^2(N,\tau))$. If we consider the net $\{T_G\}_G \subset \mathcal{B}(L^2(M,\tau))$ where G varies over all countable subgroups of $\mathcal{G}$, ordered by inclusion, then letting T_0 be any weak limit point of this net we have that $T_0 \in \mathcal{C}$.

Fix $u \in \mathcal{G}$. Then for any countable subgroup $G < \mathcal{G}$, setting $N = G''$ and $\tilde{N} = \langle G, u\rangle''$, we have $e_{\tilde{N}}[u, T_0]e_{\tilde{N}} = [u, e_{\tilde{N}} T_0 e_{\tilde{N}}] = 0$ and hence $e_N[u, T_0]e_N = 0$. If we consider the net of all countable subgroups $G < \mathcal{G}$ ordered by inclusion, then as $\mathcal{G}$ generates M, we have strong operator topology convergence $\lim_{G\to\infty} e_{G''} = 1$. Hence it follows that $[u, T_0] = 0$ and since $u \in \mathcal{G}$ was arbitrary, we have $T_0 \in \mathcal{Z}(M)$. $\square$

Let (M,τ) be a finite von Neumann algebra and $T \in \mathcal{B}(L^2(M,\tau))$. Recall that the distance between T and $\mathcal{Z}(M)$ is defined as $\mathrm{dist}(T, \mathcal{Z}(M)) = \inf\{||T - S|| : S \in \mathcal{Z}(M)\}$. For $T \in \mathcal{B}(L^2(M,\tau))$ we let δ_T denote the derivation given by $\delta_T(x) = [x, T]$.

Corollary 3.4 *Let M be a finite von Neumann algebra, and suppose $T \in \mathcal{B}(L^2(M))$, then*

$$\mathrm{dist}(T, \mathcal{Z}(M)) \leq \|\delta_{T|M'}\| + \|\delta_{T|M}\|.$$

Proof. This follows from the previous theorem since every point $S \in \{u(JvJ)T(Jv^*J)u^* \mid u, v \in \mathcal{U}(M)\}$ satisfies $\mathrm{dist}(T, S) \leq \|\delta_{T|M'}\| + \|\delta_{T|M}\|$. $\square$

As another application of Theorem 3.1 we use Christensen's Theorem [Chr82, Theorem 5.3] to establish the following vanishing cohomology result; the case when $\mathcal{C} = M$ is the celebrated Kadison-Sakai Theorem [Kad66, Sak66].

Theorem 3.5 *Let (M,τ) be a tracial von Neumann algebra and let φ be a normal regular strongly generating hyperstate, suppose $\mathcal{C} \subset \mathcal{B}_\varphi$ is a weakly closed M-bimodule. If $\delta : M \to \mathcal{C}$ is a norm continuous derivation then there exists $c \in \mathcal{C}$ so that $\delta(x) = [x, c]$ for $x \in M$. Moreover, if φ has the form $\varphi(T) = \int \langle T\widehat{u^*}, \widehat{u^*}\rangle \, d\mu(u)$ for some probability measure $\mu \in \mathrm{Prob}(\mathcal{U}(M))$, then c may be chosen so that $\|c\| \leq \|\delta\|$.*

Proof. Identifying $\mathcal{C}$ with its image under the Poisson transform we will view $\mathcal{C}$ as an operator system in $\mathrm{Har}(\mathcal{P}_\varphi) \subset \mathcal{B}(L^2(M,\tau))$. Since $L^2(M,\tau)$ has a cyclic vector for M, Christensen's Theorem [Chr82, Theorem 5.3] shows that $\delta(m) = mT - Tm$ for some $T \in \mathcal{B}(L^2(M,\tau))$. Taking the conditional expectation onto $\mathrm{Har}(\mathcal{P}_\varphi)$, we may assume $T \in \mathrm{Har}(\mathcal{P}_\varphi)$.

We suppose φ is given in standard form $\varphi(T) = \sum_n \langle T\widehat{z_n^*}, \widehat{z_n^*}\rangle$. Note that $z_m\delta(z_m^*) \in \mathcal{C}$, so that

$$T - \mathcal{P}_\varphi^{\mathrm{o}}(T) = \sum_m z_m z_m^* T - \sum_m z_m T z_m^* = \sum_m z_m \delta(z_m^*) \in \mathcal{C}$$

As $\mathcal{P}^{\mathrm{o}}_{\varphi}$ leaves $\mathcal{C}$ invariant (since $\mathcal{C}$ is an M-bimodule), by induction we get that $T-(\mathcal{P}^{\mathrm{o}}_{\varphi})^n(T)\in\mathcal{C}$ for all $n\geq 1$, and hence for $N\geq 1$ we have

$$T-\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}^{\mathrm{o}}_{\varphi})^n(T)\in\mathcal{C}$$

If z is a weak limit point of $\left\{\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}^{\mathrm{o}}_{\varphi})^n(T)\right\}$ then $z\in\mathrm{Har}(\mathcal{P}^{\mathrm{o}}_{\varphi})\cap\mathrm{Har}(\mathcal{P}_{\varphi})$ and so by Theorem 3.1 we have $z\in\mathcal{Z}(M)$. Thus, $T-z\in\mathcal{C}$ implements the derivation.

For the moreover part, note that if φ has the form $\varphi(T)=\int\langle T\widehat{u^*},\widehat{u^*}\rangle\,d\mu(u)$ for some probability measure $\mu\in\mathrm{Prob}(\mathcal{U}(M))$ then

$$\begin{aligned}\|T-z\|&\leq\sup_N\|T-\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}^{\mathrm{o}}_{\varphi})^n(T)\|\\&\leq\sup_n\|T-(\mathcal{P}^{\mathrm{o}}_{\varphi})^n(T)\|\\&=\sup_n\|\int u\delta(u^*)\,d\mu_n\|\leq\|\delta\|,\end{aligned}$$

where μ_n denotes the push forward of $\mu\times\mu\times\cdots\times\mu\in\mathrm{Prob}(\mathcal{U}(M)^n)$ under the multiplication map.

Hence $c=T-z$ implements δ with $\|c\|\leq\|\delta\|$. □

We remark that for a general hyperstate φ, in the proof of the previous theorem we still have $\|T-z\|\leq\|\delta\|_{cb}$, where $\|\delta\|_{cb}$ denotes the completely bounded norm of the derivation δ (see, for instance, [Chr82, Section 2] for the definition of the completely bounded norm). So that in general we may find $c\in\mathcal{C}$ with $\|c\|\leq\|\delta\|_{cb}$.

4. Rigidity for u.c.p. maps on boundaries

The main result in this section is Theorem 4.1, where we generalize [CP13, Theorem 3.2]. We mention several consequences, including a noncommutative version of [BS06, Corollary 3.2], which describes the Poisson boundary of a tensor product as the tensor product of Poisson boundaries.

Theorem 4.1 *Let (M,τ) be a tracial von Neumann algebra, let φ be a normal regular strongly generating hyperstate, and let $\mathcal{B}=\mathcal{B}_{\varphi}$ denote the corresponding boundary. Suppose we have a weakly closed operator system $\mathcal{C}$ such $M\subset\mathcal{C}\subset\mathcal{B}$. Let $\Psi:\mathcal{C}\to\mathcal{B}$ be a normal u.c.p. map such that $\Psi_{|M}=\mathrm{id}$. Then $\Psi=\mathrm{id}$.*

Proof. Let $\mathcal{P}_{\varphi}(T)=\sum_n(Jz_n^*J)T(Jz_nJ)$ denote the standard form of $\mathcal{P}_{\varphi}$ as in Proposition 2.8. Then by Proposition 2.4 we have $\mathcal{P}^{\mathrm{o}}_{\varphi}(T)=\sum_n z_nTz_n^*$. By identifying $\mathcal{C}$ with its image under the Poisson transform we may assume that $\mathcal{C}$ is a weakly closed M-subbimodule of $\mathrm{Har}(\mathcal{P}_{\varphi})$ and $\Psi:\mathcal{C}\to\mathrm{Har}(\mathcal{P}_{\varphi})$ is a normal u.c.p. map such that $\Psi_{|M}=\mathrm{id}$. Note that for $T\in\mathcal{C}$ we have,

$$\begin{aligned}\langle\Psi(T)\hat{1},\hat{1}\rangle&=\langle\mathcal{P}_{\varphi}(\Psi(T))\hat{1},\hat{1}\rangle=\langle\mathcal{P}^{\mathrm{o}}_{\varphi}(\Psi(T))\hat{1},\hat{1}\rangle\\&=\sum_n\langle z_n\Psi(T)z_n^*\hat{1},\hat{1}\rangle=\langle\Psi(\mathcal{P}^{\mathrm{o}}_{\varphi}(T))\hat{1},\hat{1}\rangle,\end{aligned}$$

where the last equality follows from the fact that Ψ is normal and M-bimodular, as M is contained in the multiplicative domain of Ψ. Now, $\langle\Psi(\mathcal{P}^{\mathrm{o}}_{\varphi}(T))\hat{1},\hat{1}\rangle=\langle\Psi(T)\hat{1},\hat{1}\rangle$ for all $T\in\mathcal{C}$ immediately implies that

$$\left\langle \Psi\left(\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi}^{\mathrm{o}})^n(T)\right)\hat{1},\hat{1}\right\rangle = \langle\Psi(T)\hat{1},\hat{1}\rangle \text{for all } T\in\mathcal{C}.$$

Let z be a weak operator topology limit point of $\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi}^{\mathrm{o}})^n(T)$. Then, $z\in\mathcal{Z}(M)$ by Theorem 3.1, so that $\Psi(z)=z$. We then have

$$\langle\Psi(T)\hat{1},\hat{1}\rangle = \langle z\hat{1},\hat{1}\rangle = \langle T\hat{1},\hat{1}\rangle$$

where the last equality follows because z is independent of Ψ. Now, let $a,b\in M$, and $T\in\mathcal{C}$. Then, we have that $b^*Ta\in\mathcal{C}$, and hence by above computation, we get

$$\langle\Psi(T)a\hat{1},b\hat{1}\rangle = \langle\Psi(b^*Ta)\hat{1},\hat{1}\rangle = \langle b^*Ta\hat{1},\hat{1}\rangle = \langle Ta\hat{1},b\hat{1}\rangle.$$

Thus $\Psi(T)=T$. □

Corollary 4.2 *Let M be a finite von Neumann algebra with a normal faithful trace τ, and let φ be a normal regular strongly generating hyperstate. Then, M is a maximal finite von Neumann subalgebra inside $\mathcal{B}_{\varphi}$.*

Proof. Suppose $N\subset\mathcal{B}_{\varphi}$ is a finite von Neumann algebra containing M. Then there exists a normal conditional expectation $E:N\to M$. Hence, by Theorem 4.1, $E(x)=x$ for all $x\in N$, and hence $N=M$. □

Corollary 4.3 *Let M be a II_1 factor, and let φ be a normal regular strongly generating hyperstate. If $\mathcal{B}_{\varphi}\neq M$, then $\mathcal{B}_{\varphi}$ is a type III factor.*

Proof. Note that the stationary state is normal and faithful by Proposition 2.9, and $\mathcal{B}_{\varphi}$ is a factor by Proposition 2.7. We also note that Proposition 2.7 along with von Neumann's bicommutant Theorem shows that $\mathcal{B}_{\varphi}$ is not a type I factor.

Suppose $\mathcal{B}_{\varphi}$ is not a type III factor, then $\mathcal{B}_{\varphi}$ has a semi-finite normal faithful trace Tr. As before, let $\mathcal{P}$ denote the Poisson transform, and let ζ be the normal state on $\mathcal{B}_{\varphi}$ defined by $\zeta(b)=\langle\mathcal{P}(b)\hat{1},\hat{1}\rangle$. Fix $0\leq T\in\mathcal{B}_{\varphi}$ with $Tr(T)<\infty$, and $\zeta(T)\neq 0$. Fix $S\in\mathcal{B}_{\varphi}$ with $S\geq 0$ and $Tr(S)<\infty$. Let z be a ultraweak limit point of $\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi}^{\mathrm{o}})^n(T)$. Then by Theorem 3.1 we have $z\in\mathcal{Z}(M)=\mathbf{C}$ and arguing as in the proof of Theorem 4.1 we have $\zeta(T)=z$. Therefore, $\zeta(T)Tr(S)$ is a limit point of $\{Tr((\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi}^{\mathrm{o}})^n(T))S)\}_{N=1}^{\infty}$. On the other hand, note that for each $N\in\mathbf{N}$, we have that $Tr(\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi}^{\mathrm{o}})^n(T)S)=Tr(T(\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi^*}^{\mathrm{o}})^n(S)))$. Since $|Tr(T(\frac{1}{N}\sum_{n=1}^{N}(\mathcal{P}_{\varphi^*}^{\mathrm{o}})^n(S)))|\leq Tr(T)||S||_{\infty}$, by the above discussion, we then have

$$\zeta(T)Tr(S)\leq Tr(T)||S||_{\infty}.$$

Consider a net of projections $\{S_i\}_{i\in I}$ in $\mathcal{B}_{\varphi}$, such that S_i converges to 1 in the strong operator topology. The above equation then shows that $\zeta(T)Tr(1)\leq Tr(T)<\infty$. As $\zeta(T)\neq 0$ by choice, we get that $Tr(1)<\infty$. Hence $\mathcal{B}_{\varphi}$ is a type II_1 factor and by Corollary 4.2 we have that $\mathcal{B}_{\varphi}=M$. □

Theorem 4.4 *Suppose for each $i\in\{1,2\}$, M_i is a finite von Neumann algebra with normal faithful trace τ_i. Let φ_i be a normal regular strongly generating hyperstate for M_i on $\mathcal{B}(L^2(M_i,\tau_i))$. Then,*

$$\mathrm{Har}(\mathcal{P}_{\varphi_1}\otimes\mathcal{P}_{\varphi_2}) = \mathrm{Har}(\mathcal{P}_{\varphi_1})\,\overline{\otimes}\,\mathrm{Har}(\mathcal{P}_{\varphi_2}).$$

Proof. We clearly have $\mathrm{Har}(\mathcal{P}_{\varphi_1}) \overline{\otimes} \mathrm{Har}(\mathcal{P}_{\varphi_2}) \subset \mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2})$ so we only need to show the reverse inclusion. Note that

$$(\mathcal{P}_{\varphi_1} \otimes \mathrm{id}) \circ (\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2}) = (\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2}) \circ (\mathcal{P}_{\varphi_1} \otimes \mathrm{id}),$$

hence $(\mathcal{P}_{\varphi_1} \otimes \mathrm{id})_{|\mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2})}$ gives a normal ucp map which restricts to the identity on $M_1 \overline{\otimes} M_2$. By Theorem 4.1 we have that $(\mathcal{P}_{\varphi_1} \otimes \mathrm{id})_{|\mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2})}$ is the identity map and hence

$$\mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2}) \subset \mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathrm{id}) = \mathrm{Har}(\mathcal{P}_{\varphi_1}) \overline{\otimes} \mathcal{B}(L^2(M_2)).$$

We similarly have

$$\mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2}) \subset \mathcal{B}(L^2(M_1)) \overline{\otimes} \mathrm{Har}(\mathcal{P}_{\varphi_2}).$$

Since $\mathrm{Har}(\mathcal{P}_{\varphi_1})$ is injective it is semidiscrete [Con76a], and hence has property S_σ of Kraus [Kra83, Theorem 1.9]. We then have

$$\mathrm{Har}(\mathcal{P}_{\varphi_1} \otimes \mathcal{P}_{\varphi_2}) \subset (\mathrm{Har}(\mathcal{P}_{\varphi_1}) \overline{\otimes} \mathcal{B}(L^2(M_2))) \cap (\mathcal{B}(L^2(M_1)) \overline{\otimes} \mathrm{Har}(\mathcal{P}_{\varphi_2})) \subset \mathrm{Har}(\mathcal{P}_{\varphi_1}) \overline{\otimes} \mathrm{Har}(\mathcal{P}_{\varphi_2}).$$

□

Corollary 4.5 *Suppose for each $i \in \{1, 2\}$, M_i is a finite von Neumann algebra with normal faithful trace τ_i. Let φ_i be a normal regular strongly generating hyperstate for M_i on $B(L^2(M_i, \tau_i))$. Then, the identity map on $M_1 \overline{\otimes} M_2$ uniquely extends to a $*$-isomorphism between $\mathcal{B}_{\varphi_1 \otimes \varphi_2}$ and $\mathcal{B}_{\varphi_1} \overline{\otimes} \mathcal{B}_{\varphi_2}$.*

5. Entropy

In this section we introduce noncommutative analogues of Avez's asymptotic entropy [Ave72], and Furstenberg entropy [Fur63a, Section 8].

5.1 Asymptotic entropy

Let M be a tracial von Neumann algebra with a faithful normal tracial state τ. For a normal hyperstate $\varphi \in \mathcal{S}_\tau(\mathcal{B}(L^2(M, \tau)))$ we define the **entropy of** φ, denoted by $H(\varphi)$, to be the von Neumann entropy of the corresponding density matrix A_φ:

$$H(\varphi) = -\mathrm{Tr}(A_\varphi \log(A_\varphi)).$$

If we have a standard form $\varphi(T) = \sum_n \langle T \widehat{z_n^*}, \widehat{z_n^*} \rangle$ then we may compute this explicitly as

$$H(\varphi) = -\sum_n \|z_n\|_2^2 \log(\|z_n\|_2^2).$$

Theorem 5.1 *If φ and ψ are two normal hyperstates with ψ regular, then*

$$H(\varphi * \psi) \leq H(\varphi) + H(\psi)$$

Proof. Let A_φ and A_ψ be the corresponding density operators and $\mathcal{P}_\varphi$ and $\mathcal{P}_\psi$ be the corresponding u.c.p. M-bimodular maps. Suppose we have the standard forms

$$\varphi(T) = \sum_{i \in I} \langle T \mu_i \hat{a_i^*}, \mu_i \hat{a_i^*} \rangle \text{ with } \mu_i > 0,\ ||a_i^*||_2 = 1, \text{ and } \tau(a_j a_i^*) = 0 \text{ for all } i \neq j \in I.$$

$$\psi(T) = \sum_{j \in J} \langle T \nu_j \hat{c_j^*}, \nu_j \hat{c_j^*} \rangle \text{ with } \nu_j > 0,\ ||c_j^*||_2 = 1, \text{ and } \tau(c_k c_l^*) = 0 \text{ for all } k \neq l \in J.$$

Hence $A_\varphi = \sum_i \mu_i P_{\hat{a}_i}$ and $A_\psi = \sum_j \nu_j P_{\hat{c}_j}$.

Let $b_i = Ja_iJ$ and $d_i = Jc_iJ$ so that

$$\mathcal{P}_\varphi(T) = \sum_i \mu_i b_i T b_i^* \text{ and } \mathcal{P}_\psi(T) = \sum_j \nu_j d_j T d_j^*.$$

Since ψ is regular we have that $\sum_i \nu_i d_i^* d_i = \sum_i \nu_i d_i d_i^* = 1$. Since φ is a hyperstate we have that $\sum_i \mu_i b_i b_i^* = 1$. Now,

$$H(\varphi * \psi) = -\sum_{i,j} Tr[\mu_i \nu_j b_i^* d_j^* P_{\hat{1}} d_j b_i \log(A_{\varphi*\psi})].$$

and

$$b_i^* d_j^* P_{\hat{1}} d_j b_i = \tau(b_i b_i^* d_j^* d_j) P_{\widehat{b_i^* d_j^*}},$$

so that for each k, ℓ we have

$$A_{\varphi*\psi} = \sum_{i,j} \mu_i \nu_j b_i^* d_j^* P_{\hat{1}} d_j b_i \geq \mu_k \nu_\ell \tau(b_k b_k^* d_\ell^* d_\ell) P_{\widehat{b_k^* d_\ell^*}}.$$

As log is operator monotone, for each k, ℓ we then have

$$-\log(A_{\varphi*\psi}) = -\log(\sum_{i,j} \mu_i \nu_j b_i^* d_j^* P_{\hat{1}} d_j b_i) \leq -\log((\mu_k \nu_\ell \tau(b_k b_k^* d_\ell^* d_\ell)) P_{\widehat{b_k^* d_\ell^*}}).$$

Hence,

$$\begin{aligned} H(\varphi * \psi) &\leq -\sum_{i,j} Tr[\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j) P_{\widehat{b_i^* d_j^*}} \log(\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j) P_{\widehat{b_i^* d_j^*}})] \\ &= -\sum_{i,j} Tr[\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j) P_{\widehat{b_i^* d_j^*}} \log(\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j))] \\ &\quad -\sum_{i,j} Tr[\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j) P_{\widehat{b_i^* d_j^*}} \log(P_{\widehat{b_i^* d_j^*}})] \\ &= -\sum_{i,j} \mu_i \nu_j \tau(b_i b_i^* d_j^* d_j) \log(\mu_i \nu_j \tau(b_i b_i^* d_j^* d_j)). \end{aligned}$$

Now define m on $I \times J$ by $m(i,j) = \mu_i \nu_j \tau(b_i b_i^* d_j^* d_j)$. Note that

$$\sum_i m(i,j) = \nu_j \tau(\sum_i \mu_i b_i b_i^* d_j^* d_j) = \nu_j \tau(d_j^* d_j) = \nu_j$$

and

$$\sum_j m(i,j) = \mu_i \tau(\sum_i \nu_j b_i b_i^* d_j^* d_j) = \mu_i \tau(b_i b_i^*) = \mu_i.$$

To finish the proof it then suffices to show

$$H(m) = -\sum_{i,j} m(i,j) \log(m(i,j)) \leq H(\mu) + H(\nu),$$

where $H(\mu) = -\sum_i \mu_i \log(\mu_i)$ and $H(\nu) = -\sum_i \nu_i \log(\nu_i)$. By the remark before Theorem 5.1, a direct calculation yields $H(\mu) = H(\varphi)$ and $H(\nu) = H(\psi)$.

Note that

$$\begin{aligned}
H(m) &= -\sum_{i,j} m(i,j)\log(m(i,j)) \\
&= -\sum_{i,j} \mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\mu_i\tau(b_ib_i^*d_j^*d_j)) - \sum_{i,j}\mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\nu_j) \\
&= -\sum_{i,j} \mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\mu_i) \\
&\quad -\sum_{i,j} \mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\nu_j) - \sum_{i,j}\mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\tau(b_ib_i^*d_j^*d_j)).
\end{aligned}$$

In the last equality above, the first summation is $H(\mu)$, since summing over j we get

$$-\sum_i \mu_i\tau(b_ib_i^*)\log(\mu_i) = -\sum_i \mu_i\log(\mu_i),$$

while the second summation is $H(\nu)$. Hence, all that remains is to show:

$$\sum_{i,j}\mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\tau(b_ib_i^*d_j^*d_j)) \geq 0.$$

Let $\eta(x) = -x\log(x)$ for $x \in [0,1]$. Note that η is concave, and so $\eta(\sum_i \alpha_i x_i) \geq \sum_i \alpha_i\eta(x_i)$ whenever $\alpha_i \geq 0$ and $\sum_i \alpha_i = 1$. So,

$$\begin{aligned}
-\sum_{i,j}\mu_i\nu_j\tau(b_ib_i^*d_j^*d_j)\log(\tau(b_ib_i^*d_j^*d_j)) &= \sum_{i,j}\mu_i\nu_j\eta(\tau(b_ib_i^*d_j^*d_j)) \\
&= \sum_i \mu_i(\sum_j \nu_j\eta(\tau(b_ib_i^*d_j^*d_j))) \\
&\leq \sum_i \mu_i\eta(\sum_j \nu_j\tau(b_ib_i^*d_j^*d_j)) \\
&= \sum_i \mu_i\eta(\tau(b_ib_i^*)) = 0
\end{aligned}$$

□

Corollary 5.2 *If φ is a normal regular hyperstate, then the limit $\lim_{n\to\infty}\frac{H(\varphi^{*n})}{n}$ exists.*

Proof. The sequence $\{H(\varphi^{*n})\}$ is subadditive by Theorem 5.1 and hence the limit exists. □

The **asymptotic entropy** $h(\varphi)$ of a normal regular hyperstate φ is defined to be the limit

$$h(\varphi) = \lim_{n\to\infty}\frac{H(\varphi^{*n})}{n}.$$

5.2 A Furstenberg type entropy

Suppose G is a Polish group and $\mu \in \mathrm{Prob}(G)$. Given a quasi-invariant action $G \overset{a}{\curvearrowright} (X,\nu)$ the corresponding Furstenberg entropy (or μ-entropy) is defined [Fur63a, Section 8] to be

$$h_\mu(a,\nu) = -\iint \log\left(\frac{dg^{-1}\nu}{d\nu}(x)\right)\, d\nu(x)d\mu(g).$$

If we consider the measure space $(G\times X, \nu\times\mu)$ then we have a non-singular map $\pi : G\times X \to G\times X$ given by $\pi(g,x) = (g, g^{-1}x)$, whose Radon-Nikodym derivative is given by

$$\frac{d\pi(\mu\times\nu)}{d(\mu\times\nu)}(x,g) = \frac{dg^{-1}\nu}{d\nu}(x).$$

Recall that for arbitrary positive functions $f, g \in L^1(X,\mu)$ (where (X,μ) is a standard probability space), the relative entropy of the measures $\mu_1 = f d\mu$ and $\mu_2 = g d\mu$, denoted by $S(\mu_1|\mu_2)$, is defined as $S(\mu_1|\mu_2) = \int_X f(\log(f) - \log(g))d\mu$ (see [OP93, Chapter 5]). We may thus rewrite the μ-entropy as a relative entropy

$$h_\mu(a,\nu) = -\iint \log\left(\frac{d\pi(\nu\times\mu)}{d(\nu\times\mu)}(g,x)\right) d(\nu\times\mu) = S((\nu\times\mu)|\pi(\nu\times\mu)).$$

Let (M,τ) be a tracial von Neumann algebra, φ a normal hyperstate for M, and $\mathcal{A}$ a von Neumann algebra, such that $M \subseteq \mathcal{A}$. Let $\zeta \in \mathcal{S}_\tau(\mathcal{A})$ be a normal, faithful hyperstate. Let $\Delta_\zeta : L^2(\mathcal{A},\zeta) \to L^2(\mathcal{A},\zeta)$ be the modular operator corresponding to ζ, and consider the spectral decomposition $\Delta_\zeta = \int_0^\infty \lambda\, dE(\lambda)$. We denote by $\Delta_n = \int_{1/n}^n \lambda d\lambda$, $n \geq 1$, the truncations of the modular operator Δ. We know that Δ_n converges to Δ in the resolvent sense. Throughout this section we denote the one parameter modular automorhism group associated with ζ by $\{\sigma_t^\zeta\}_{t\in\mathbf{R}}$. We also denote the corresponding modular conjugation operator by J, and let $S = J\Delta^{1/2}$. We refer the reader to [Tak03, Chapters VI, VII, VIII] for details regarding Tomita-Takesaki Theory.

Since $\zeta|_M = \tau$, we have a natural inclusion of $L^2(M,\tau)$ in $L^2(\mathcal{A},\zeta)$. Let e denote the orthogonal projection from $L^2(\mathcal{A},\zeta)$ to $L^2(M,\tau)$. The entropy of the inclusion $(M,\tau) \subset (\mathcal{A},\zeta)$ with respect to φ is defined to be

$$h_\varphi(M\subset\mathcal{A},\zeta) = -\int \log(\lambda)\, d\varphi(eE(\lambda)e).$$

The next example shows that $h_\varphi(M\subset\mathcal{A},\zeta)$ can be considered as a generalization of the Furstenberg entropy.

Example 5.3 *If Γ is a discrete group, $\mu \in \mathrm{Prob}(\Gamma)$ and $\Gamma \overset{a}{\curvearrowright} (X,\nu)$ is a quasi-invariant action, then we may consider the state φ on $\mathcal{B}(\ell^2\Gamma)$ given by $\varphi(T) = \int\langle T\delta_\gamma, \delta_\gamma\rangle\, d\mu(\gamma)$, and we may consider the state ζ on $L^\infty(X,\nu)\rtimes\Gamma \subset \mathcal{B}(\ell^2\Gamma\,\overline{\otimes}\,L^2(X,\nu))$ given by $\zeta\left(\sum_{\gamma\in\Gamma} a_\gamma u_\gamma\right) = \int a_e\, d\nu$. Note that a direct computation in this case yields $(\varphi * \zeta)\left(\sum_{\gamma\in\Gamma} a_\gamma u_\gamma\right) = \int a_e\, d(\mu * \nu)$. The modular operator Δ_ζ is then affiliated to the von Neumann algebra $\ell^\infty\Gamma\,\overline{\otimes}\,L^\infty(X,\nu)$, and we may compute this directly as*

$$\Delta_\zeta(\gamma,x) = \frac{d\gamma^{-1}\nu}{d\nu}(x).$$

We also have that the projection e from $\ell^2\Gamma\,\overline{\otimes}\,L^2(X,\nu) \to \ell^2\Gamma$ is given by $\mathrm{id}\otimes\int$. Thus, it follows that the measure $d\varphi(eE(\lambda)e)$ agrees with $d\alpha_*(\mu\times\nu)$, where $\alpha : \Gamma\times X \to \mathbf{R}_{>0}$ is the Radon-Nikodym cocycle, $\alpha(\gamma,x) = \frac{d\gamma^{-1}\nu}{d\nu}(x)$.

In this case we then have

$$\begin{aligned} h_\varphi(L\Gamma \subset L^\infty(X,\nu)\rtimes\Gamma,\zeta) &= -\int \log(\lambda) d\varphi(eE(\lambda)e) \\ &= -\iint \log\left(\frac{d\gamma^{-1}\nu}{d\nu}(x)\right) d(\nu\times\mu) = h_\mu(a,\nu). \end{aligned}$$

Lemma 5.4 *Let $\varphi \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ be a normal hyperstate and write φ in a standard form $\varphi(T) = \sum_n \langle T\widehat{z_n^*}, \widehat{z_n^*}\rangle$. Suppose $\mathcal{A}$ is a von Neumann algebra with $M \subset \mathcal{A}$ and $\zeta \in \mathcal{S}_\tau(\mathcal{A})$ is a normal hyperstate. Then if $h_\varphi(M \subset \mathcal{A}, \zeta) < \infty$ we have that $z_n^* 1_\zeta \in D(\log \Delta_\zeta)$ for each n and*

$$h_\varphi(M \subset \mathcal{A}, \zeta) = -\sum_n \langle \log \Delta_\zeta z_n^* 1_\zeta, z_n^* 1_\zeta\rangle = i \lim_{t\to 0} \frac{1}{t} \sum_n (\zeta(z_n \sigma_t^\zeta(z_n^*)) - 1).$$

Proof. As $\mathcal{A}1_\zeta$ forms a core for S_ζ we get that $z_n^* 1_\zeta \in D(\log(\Delta_\zeta))$. Also, we know that $\lim_{t\to 0} \frac{\Delta_\zeta^{it} - 1}{t}\xi = i\log(\Delta_\zeta)\xi$, for all $\xi \in D(\Delta_\zeta)$. So, we have that

$$\begin{aligned} h_\varphi(M \subset \mathcal{A}, \zeta) &= -\varphi(e\log(\Delta_\zeta)e) = -\sum_n \langle \log \Delta_\zeta z_n^* 1_\zeta, z_n^* 1_\zeta\rangle \\ &= i\sum_n \langle z_n \lim_{t\to 0} \frac{\Delta_\zeta^{it} - 1}{t} z_n^* 1_\zeta, 1_\zeta\rangle = i \lim_{t\to 0} \frac{1}{t} \sum_n (\zeta(z_n \sigma_t^\zeta(z_n^*)) - 1). \end{aligned}$$

□

Example 5.5 *Fix two normal hyperstates $\varphi, \zeta \in \mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ such that φ is regular, and ζ is faithful, and consider the case $\mathcal{A} = \mathcal{B}(L^2(M,\tau))$. Then the density operator A_ζ is injective with dense range and the modular operator on $L^2(\mathcal{B}(L^2(M,\tau)),\zeta)$ is given by $\Delta_\zeta(T1_\zeta) = A_\zeta T A_\zeta^{-1} 1_\zeta$, for $T \in \mathcal{B}(L^2(M,\tau))$ such that $T1_\zeta \in D(\Delta_\zeta)$. In particular note that $\log(\Delta_\zeta)(T1_\zeta) = (\mathrm{Ad}(\log A_\zeta)T)1_\zeta$, where $\mathrm{Ad}(\log A_\zeta)T = (\log A_\zeta)T - T(\log A_\zeta)$.*

We also have that the projection $e : L^2(\mathcal{B}(L^2(M,\tau)),\zeta) \to L^2(M,\tau)$ is given by $e(T1_\zeta) = \mathcal{P}_\zeta(T)\hat{1}$. Therefore, $e\log\Delta_\zeta e x\hat{1} = \mathcal{P}_\zeta(\mathrm{Ad}(\log A_\zeta)x)\hat{1} = \mathcal{P}_\zeta(\mathrm{Ad}(\log A_\zeta))x\hat{1}$. Hence,

$$\begin{aligned} h_\varphi(M \subset \mathcal{B}(L^2(M,\tau)),\zeta) &= \varphi(\mathcal{P}_\zeta(\mathrm{Ad}(\log A_\zeta))) \\ &= \mathrm{Tr}(A_{\varphi*\zeta}\mathrm{Ad}(\log A_\zeta)) \\ &= \mathrm{Tr}(A_{\varphi*\zeta}\log A_\zeta) - \langle \log A_\zeta \hat{1}, \hat{1}\rangle. \end{aligned}$$

Where the last equality follows since φ is regular.

We recall the following two lemmas from works of D.Petz [Pet86].

Lemma 5.6 *Let Δ_j be positive, self adjoint operators on $\mathcal{H}_j, j = 1,2$. If $T : \mathcal{H}_1 \to \mathcal{H}_2$ is a bounded operator such that:*

- $T(\mathcal{D}(\Delta_1)) \subseteq \mathcal{D}(\Delta_2)$
- $||\Delta_2 T\xi|| \leq ||T|| \cdot ||\Delta_1 \xi||$ *($\xi \in \mathcal{D}(\Delta_1)$),*

then we have for each $t \in [0,1]$, and $\xi \in \mathcal{D}(\Delta_1^t)$,

$$||\Delta_2^t T\xi|| \leq ||T|| \cdot ||\Delta_1^t \xi||$$

Lemma 5.7 *Let Δ be a positive self adjoint operator and $\xi \in \mathcal{D}(\Delta)$. Then:*

$$\lim_{t\to 0+} \frac{||\Delta^{t/2}\xi||^2 - ||\xi||^2}{t}$$

exists. It's finite or $-\infty$ and equals $\int_0^\infty \log\lambda d\langle E_\lambda \xi, \xi\rangle$ where $\int_0^\infty \log\lambda dE_\lambda$ is the spectral resolution of Δ.

Corollary 5.8 $h_\varphi(M\subset\mathcal{A},\zeta) = -\lim\limits_{t\to0+}\dfrac{\sum\limits_{k=1}^{\infty}||\Delta_\zeta^{t/2}ez_n^*\hat{1}||^2-||ez_n^*\hat{1}||^2}{t}$

Lemma 5.9 $h_\varphi(M\subset\mathcal{A},\zeta)\geq 0$

Proof. Let $\mathcal{P}_\zeta(T)=eTe$ for $T\in\mathcal{A}$. Let $\Delta_n=\int_{1/n}^{n}\lambda d\lambda$, $n\geq1$, denote the truncations of the modular operator Δ.
$h_\varphi(M\subset\mathcal{A},\zeta)=\lim\limits_{n\to\infty}\varphi(-e\log\Delta_n e)=-\lim\limits_{n\to\infty}\langle\mathcal{P}_\varphi\circ\mathcal{P}_\zeta(\log\Delta_n)\hat{1},\hat{1}\rangle\geq\lim\limits_{n\to\infty}-\langle\log(\mathcal{P}_\varphi\circ\mathcal{P}_\zeta(\Delta_n))\hat{1},\hat{1}\rangle$ (using the operator Jensen's inequality; recall that log is operator concave).
Notice that $e\Delta_n e\leq e\Delta e=e$. Since $\mathcal{P}_\varphi(e)=e$, we get $\mathcal{P}_\varphi\circ\mathcal{P}_\zeta(\Delta_n)\leq e\leq 1$. As log is operator monotone, we get that $\log(\mathcal{P}_\varphi\circ\mathcal{P}_\zeta(\Delta_n))\leq\log(1)=0$. Hence we are done. $\square$

Theorem 5.10 *Let $\varphi,\psi\in\mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ be two normal hyperstates such that ψ is regular, and suppose $\mathcal{A}$ is a von Neumann algebra with $M\subset\mathcal{A}$, and $\zeta\in S_\tau(\mathcal{A})$ is a normal, faithful hyperstate which is ψ-stationary. Then*

$$h_{\varphi*\psi}(M\subset\mathcal{A},\zeta)=h_\varphi(M\subset\mathcal{A},\zeta)+h_\psi(M\subset\mathcal{A},\zeta).$$

Proof. Suppose we have the standard forms

$$\varphi(T)=\sum_{i\in I}\langle T\mu_i\hat{a_i^*},\mu_i\hat{a_i^*}\rangle \text{ with } \mu_i>0,\ ||a_i^*||_2=1, \text{ and } \tau(a_ja_i^*)=0 \text{ for all } i\neq j\in I.$$
$$\psi(T)=\sum_{j\in J}\langle T\nu_j\hat{b_j^*},\nu_j\hat{b_j^*}\rangle \text{ with } \nu_j>0,\ ||b_j^*||_2=1, \text{ and } \tau(b_kb_l^*)=0 \text{ for all } k\neq l\in J.$$

Let $\mathcal{P}_\varphi$ and $\mathcal{P}_\psi$ be the corresponding u.c.p. maps so that $\mathcal{P}_\varphi(T)=\sum_k\mu_kJa_k^*JTJa_kJ$ and $\mathcal{P}_\psi(T)=\sum_l\nu_lJb_l^*JTJb_lJ$. We shall denote the projection from $L^2(\mathcal{A},\zeta)$ to $L^2(M,\tau)$ by e and Δ_ζ by Δ. We also denote the one parameter modular automorphism group corresponding to ζ by σ_t. We then have

$$\begin{aligned}h_\varphi(M\subset\mathcal{A},\zeta)&=i\lim_{t\to0}\varphi(\frac{e\Delta^{it}e-1}{t})=i\lim_{t\to0}\frac{1}{t}\varphi(e\Delta^{it}e-1)\\&=i\lim_{t\to0}\frac{1}{t}(\sum_k\mu_k\langle(\Delta^{it}-1)a_k^*1_\zeta,a_k^*1_\zeta\rangle)\end{aligned}$$

Similarly,

$$h_\psi(M\subset\mathcal{A},\zeta)=i\lim_{t\to0}\frac{1}{t}(\sum_l\nu_l\langle(\Delta^{it}-1)b_l^*1_\zeta,b_l^*1_\zeta\rangle)$$

and,

$$\begin{aligned}h_{\varphi*\psi}(M\subset\mathcal{A},\zeta)&=i\lim_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle(\Delta^{it}-1)a_k^*b_l^*1_\zeta,a_k^*b_l^*1_\zeta\rangle)\\&=i\lim_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle(b_la_k\sigma_t(a_k^*b_l^*)1_\zeta,1_\zeta\rangle-1)\end{aligned}$$

We shall now show: $\lim\limits_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle b_la_k\sigma_t(a_k^*b_l^*)1_\zeta,1_\zeta\rangle-\sum_{k,l}\mu_k\nu_l\langle b_l\sigma_t(b_l^*)\sigma_t(a_k^*)1_\zeta,1_\zeta\rangle)=0$. Let $y_t=a_k\sigma_t(a_k^*)$. Note that $y_t\to a_ka_k^*$ as $t\to0$, in SOT. We have:

$$\begin{aligned}y_t\sigma_t(b_l^*)-\sigma_t(b_l^*)y_t&=y_t\sigma_t(b_l^*)-y_tb_l^*+y_tb_l^*-\sigma_t(b_l^*)y_t\\&=y_t(\sigma_t(b_l^*)-b_l^*)+(y_tb_l^*-b_l^*y_t)+(b_l^*-\sigma_t(b_l^*))y_t\end{aligned}$$

Now,

$$\begin{aligned}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle(y_tb_l^*-b_l^*y_t)1_\zeta,b_l^*1_\zeta\rangle &= \frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle b_ly_tb_l^*1_\zeta,1_\zeta\rangle-\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle y_t1_\zeta,b_lb_l^*1_\zeta\rangle\\ &= \frac{1}{t}\sum_k\mu_k\langle(\sum_l\nu_lb_ly_tb_l^*)1_\zeta,1_\zeta\rangle-\frac{1}{t}\sum_k\mu_k\langle y_t1_\zeta,1_\zeta\rangle\\ &= \frac{1}{t}\langle y_t1_\zeta,1_\zeta\rangle-\frac{1}{t}\langle y_t1_\zeta,1_\zeta\rangle=0,\end{aligned}$$

where the second to last equality holds by ψ-stationarity of ζ.
Also, $\lim_{t\to0}\frac{1}{t}(y_t(\sigma_t(b_l^*)-b_l^*))$ exists, and hence

$$\lim_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle b_la_k\sigma_t(a_k^*b_l^*)1_\zeta,1_\zeta\rangle-\sum_{k,l}\mu_k\nu_l\langle b_l\sigma_t(b_l^*)\sigma_t(a_k^*)1_\zeta,1_\zeta\rangle)=0.$$

So, we get that

$$\begin{aligned}h_{\varphi*\psi}(M\subset\mathcal{A},\zeta) &= i\lim_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l\langle(b_l\sigma_t(b_l^*)a_k\sigma_t(a_k^*)-1)1_\zeta,1_\zeta\rangle\\ &= i\lim_{t\to0}\frac{1}{t}(\sum_{k,l}\mu_k\nu_l[\langle(b_l\sigma_t(b_l^*)-1)1_\zeta,1_\zeta\rangle\\ &+\langle(a_k\sigma_t(a_k^*)-1)1_\zeta,1_\zeta\rangle\\ &+\langle(a_k\sigma_t(a_k^*)-1)1_\zeta,(b_l\sigma_t(b_l^*)-1)^*1_\zeta\rangle]\end{aligned}$$

The first term equals $h_\varphi(M\subset\mathcal{A},\zeta)$, while second term equals $h_\psi(M\subset\mathcal{A},\zeta)$, and the third term equals zero, as $\lim_{t\to0}\frac{1}{t}(a_k\sigma_t(a_k^*)-1)1_\zeta$ exists, while $\lim_{t\to0}\sum_l\nu_l(b_l\sigma_t(b_l^*)-1)^*1_\zeta=0$. □

Corollary 5.11 *Let $\varphi\in\mathcal{S}_\tau(\mathcal{B}(L^2(M,\tau)))$ be a regular normal hyperstate and suppose $\mathcal{A}$ is a von Neumann algebra with $M\subset\mathcal{A}$, and $\zeta\in S_\tau(\mathcal{A})$ is a faithful φ-stationary hyperstate, then for $n\geq1$ we have*

$$h_{\varphi^{*n}}(M\subset\mathcal{A},\zeta)=nh_\varphi(M\subset\mathcal{A},\zeta).$$

Lemma 5.12 $h_\varphi(M\subset\mathcal{A},\zeta)\leq H(\varphi)$.

Proof. We continue the notation from the proof of Theorem 5.10, so that $\mathcal{P}_\varphi(T)=\sum_k\mu_kb_kTb_k^*$. Let $a_k=Jb_kJ\in M$. It follows from Lemma 5.7 that

$$H(\varphi)=-\lim_{t\to0+}\frac{\sum_{k=1}^\infty\mu_k||A_\varphi^{t/2}a_k^*\hat{1}||^2-||a_k^*\hat{1}||^2}{t}.$$

So by Corollary 5.8 it's enough to show that

$$\lim_{t\to0+}\frac{\sum_{k=1}^\infty\mu_k||A_\varphi^{t/2}a_k^*\hat{1}||^2-||a_k^*\hat{1}||^2}{t}\leq\lim_{t\to0+}\frac{\sum_{k=1}^\infty\mu_k||\Delta_\varphi^{t/2}ea_k^*\hat{1}||^2-||ea_k^*\hat{1}||^2}{t}.$$

So, it's enough to show that

$$||A_\varphi^{t/2}a_k\hat{1}||^2\leq||\Delta_\zeta^{t/2}a_k1_\zeta||^2$$

Define $T : L^2(\mathcal{A},\zeta) \to L^2(M,\tau)$ by $T(a1_\zeta) = \mathcal{P}_\zeta(a)\hat{1}$. Then $||T|| = 1$, as $||T(1_\zeta)|| = 1$ and $||\mathcal{P}_\zeta|| \leq 1$. T takes $\mathcal{D}(\Delta_\zeta)$ into $\mathcal{D}(A_\varphi) = L^2(M,\tau)$. We now denote Δ_ζ by Δ. By Lemma 5.6 it's enough to show:

$$||A_\varphi^{1/2}T\xi|| \leq ||\Delta^{1/2}\xi|| \text{for all } \xi \in \mathcal{D}(\Delta).$$

In fact it's enough to show the above for all vectors in a core for $\mathcal{D}(\Delta)$. Recall that $\mathcal{A}1_\zeta$ forms a core for $\mathcal{D}(\Delta)$. So, we only need to show

$$||A_\varphi^{1/2}Ta1_\zeta|| \leq ||\Delta^{1/2}a1_\zeta|| \text{ for all } a \in \mathcal{A}.$$

To this end, let $a \in \mathcal{A}$. Recall that $S = J\Delta^{1/2}$, so that $\Delta^{1/2} = JS$. We then have

$$\begin{aligned}
||\Delta^{1/2}a1_\zeta||^2 &= \langle \Delta^{1/2}a1_\zeta, \Delta^{1/2}a1_\zeta\rangle = \langle JSa1_\zeta, JSa1_\zeta\rangle \\
&= \langle Ja^*1_\zeta, Ja^*1_\zeta\rangle = \langle a^*1_\zeta, a^*1_\zeta\rangle = \zeta(aa^*) \\
&= \langle \mathcal{P}_\zeta(aa^*)\hat{1}, \hat{1}\rangle
\end{aligned}$$

We also have $\mathcal{P}_\varphi \circ \mathcal{P}_\zeta = \mathcal{P}_\zeta \implies \varphi \circ \mathcal{P}_\zeta = \zeta$. Now:

$$\begin{aligned}
||A_\varphi^{1/2}Ta1_\zeta||^2 &= \langle A_\varphi^{1/2}\mathcal{P}_\zeta(a)\hat{1}, A_\varphi^{1/2}\mathcal{P}_\zeta(a)\hat{1}\rangle = \langle A_\varphi\mathcal{P}_\zeta(a)\hat{1}, \mathcal{P}_\zeta(a)\hat{1}\rangle \\
&= \langle \mathcal{P}_\zeta(a)^*A_\varphi\mathcal{P}_\zeta(a)\hat{1}, \hat{1}\rangle \leq Tr(\mathcal{P}_\zeta(a)^*A_\varphi\mathcal{P}_\zeta(a)) \\
&= Tr(A_\varphi\mathcal{P}_\zeta(a)\mathcal{P}_\zeta(a^*)) \leq Tr(A_\varphi\mathcal{P}_\zeta(aa^*)) \\
&= \langle (\varphi \circ \mathcal{P}_\zeta)(aa^*)\hat{1}, \hat{1}\rangle = \langle \mathcal{P}_\zeta(aa^*)\hat{1}, \hat{1}\rangle \\
&= \zeta(aa^*) = ||\Delta^{1/2}a1_\zeta||^2.
\end{aligned}$$

Hence we are done. ☐

Corollary 5.13 $h_\varphi(M \subset \mathcal{A},\zeta) \leq h(\varphi)$

Proof. By Lemma 5.12, we have that $h_{\varphi^{*n}}(M \subset \mathcal{A},\zeta) \leq H(\varphi^{*n})$. By Corollary 5.11 we have that $h_{\varphi^{*n}}(M \subset \mathcal{A},\zeta) = nh_\varphi(M \subset \mathcal{A},\zeta)$. So we get,

$$h_\varphi(M \subset \mathcal{A},\zeta) \leq \frac{H(\varphi^{*n})}{n} \to h(\varphi).$$

☐

Lemma 5.14 $h_\varphi(M \subset \mathcal{A},\zeta) = 0$ *if and only if there exists a normal* ζ *preserving conditional expectation from* $\mathcal{A}$ *to* M.

Proof. Let φ be a standard form $\varphi(T) = \sum_k \langle T\widehat{a_k^*}, \widehat{a_k^*}\rangle$. Let $\mathcal{E} : \mathcal{A} \to M$ be a normal ζ preserving conditional expectation. Then, we know that $\sigma_t^\zeta(m) = m$ for all $m \in M$, where σ_t^ζ denotes the modular automorphism group corresponding to ζ. Hence,

$$\begin{aligned}
h_\varphi(M \subset \mathcal{A},\zeta) &= i\lim_{t\to 0}\frac{1}{t}\sum_k \langle (\Delta^{it} - 1)a_k^*1_\zeta, a_k^*1_\zeta\rangle \\
&= i\lim_{t\to 0}\frac{1}{t}\sum_k \langle \sigma_t(a_k^*)1_\zeta, a_k^*1_\zeta\rangle - 1 = 0.
\end{aligned}$$

Conversely, suppose $h_\varphi(M \subset \mathcal{A},\zeta) = 0$. This part of the proof is motivated by the proof of Lemma 9.2 in [OP93]. Let $\Delta_\zeta = \Delta$ and let $\Delta = \int_0^\infty \lambda d\lambda$ be it's spectral resolution. Let $\Delta_n = \int_{1/n}^n \lambda d\lambda$, $n \geq 1$ be the truncations. We know that Δ_n converges to Δ in the resolvent sense. As usual, we denote by e the projection from $L^2(\mathcal{A},\zeta)$ to $L^2(M,\tau)$. We have that $e = e\Delta e \geq e\Delta_n e$ for all n.

So, $(1+t)^{-1} \leq (e\Delta_n e + t)^{-1} \leq e(\Delta_n + t)^{-1}e$ for all n and for all $t > 0$. Taking limits as $n \to \infty$, we get $(1+t)^{-1} \leq e(\Delta + t)^{-1}e$. Now we shall use the following integral representation of log:

$$\log(x) = \int_0^\infty [(1+t)^{-1} - (x+t)^{-1}]dt$$

So that

$$h_\varphi(M \subset \mathcal{A}, \zeta) = -\int_0^\infty \sum_k \langle e[(1+t)^{-1} - (\Delta+t)^{-1}]ea_k^*\hat{1}, a_k^*\hat{1}\rangle dt.$$

From $h_\varphi(M \subset \mathcal{A}, \zeta) = 0$ and the above discussion, we deduce that

$$\begin{aligned}
&\langle e((1+t)^{-1} - (\Delta+t)^{-1})ea_k^*\hat{1}, a_k^*\hat{1}\rangle = 0 \\
&\Rightarrow e((1+t)^{-1} - (\Delta+t)^{-1})ea_k^*\hat{1} = 0 \\
&\Rightarrow (1+t)^{-1}a_k^*\hat{1} = e(\Delta+t)^{-1}ea_k^*\hat{1}
\end{aligned}$$

for almost all $t > 0$, and hence by continuity, for all $t > 0$. We now show that the last relation also holds without the compression e. To this end, note that by differentiating the equation $(1+t)^{-1}a_k^*\hat{1} = e(\Delta+t)^{-1}a_k^*\hat{1}$ with respect to t, we get $(1+t)^{-2}a_k^*\hat{1} = e(\Delta+t)^{-2}ea_k^*\hat{1}$, for all $t > 0$. Therefore, by the following norm calculation in $L^2(\mathcal{A}, \zeta)$ we have

$$\begin{aligned}
||e(\Delta+t)^{-1}ea_k^*\hat{1}||_2^2 &= ||(1+t)^{-1}a_k^*\hat{1}||_2^2 = \langle (1+t)^{-2}a_k^*\hat{1}, a_k^*\hat{1}\rangle \\
&= \langle e(\Delta+t)^{-2}ea_k^*\hat{1}, a_k^*\hat{1}\rangle = \langle (\Delta+t)^{-2}a_k^*\hat{1}, a_k^*\hat{1}\rangle = ||(\Delta+t)^{-1}a_k^*\hat{1}||_2^2.
\end{aligned}$$

So we get that $(1+t)^{-1}a_k^*\hat{1} = (\Delta+t)^{-1}a_k^*\hat{1}$ for all $t > 0$. This implies that $\Delta^{it}a_k^*1_\zeta = a_k^*1_\zeta$, which implies that $\sigma_t^\zeta(a_k^*) = a_k^*$ and hence $\sigma_t^\zeta(m) = m$ for all $m \in M$, as φ is generating. Hence there exists a ζ preserving conditional expectation from $\mathcal{A}$ to M, which is normal, as ζ is normal. □

Corollary 5.15 $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) = M$ *if and only if* $h_\varphi(M \subset \mathcal{B}_\varphi, \zeta) = 0$, *where* $\mathcal{B}_\varphi$ *denotes the Poisson boundary with respect to* φ.

Proof. If $h_\varphi(M \subset \mathcal{B}_\varphi, \zeta) = 0$, then by Lemma 5.14 there exists a normal conditional expectation $\mathcal{E} : \mathcal{B}_\varphi \to M$. By Theorem 4.1, $\mathcal{E} = \mathrm{id}$, which implies that $\mathcal{B}_\varphi = M$, and hence

$$\mathrm{Har}(\mathcal{P}_\varphi) = \mathcal{P}(\mathcal{B}_\varphi) = \mathcal{P}(M) = M.$$

Conversely, if $\mathrm{Har}(\mathcal{B}(L^2M,\tau), \mathcal{P}_\varphi) = M$ then $\Delta_\zeta = I$ and hence $h_\varphi(M \subset \mathcal{B}_\varphi, \zeta) = 0$ □

Corollary 5.16 $\mathrm{Har}(\mathcal{B}(L^2(M,\tau)), \mathcal{P}_\varphi) = M$ *if* $h(\varphi) = 0$.

Proof. Since $0 \leq h_\varphi(M \subset \mathcal{B}_\varphi, \zeta) \leq h(\varphi)$, this result follows from Corollary 5.15. □

6. An entropy gap for property (T) factors

If (M, τ) is a tracial von Neumann algebra, then a Hilbert M-bimodule consists of a Hilbert space $\mathcal{H}$, together with commuting normal representations $L : M \to \mathcal{B}(\mathcal{H})$, $R : M^{\mathrm{op}} \to \mathcal{B}(\mathcal{H})$. We will sometimes simplify notation by writing $x\xi y$ for the vector $L(x)R(y^{\mathrm{op}})\xi$. A vector $\xi \in \mathcal{H}$ is left (resp. right) tracial if $\langle x\xi, \xi\rangle = \tau(x)$ (resp. $\langle \xi x, \xi\rangle = \tau(x)$) for all $x \in M$. A vector is bi-tracial if it is both left and right tracial. A vector $\xi \in \mathcal{H}$ is central if $x\xi = \xi x$ for all $x \in M$. Note that if ξ is a unit central vector then $x \mapsto \langle x\xi, \xi\rangle$ gives a normal trace on M.

The von Neumann algebra M has property (T) if for any sequence of Hilbert bimodules $\mathcal{H}_n$, and $\xi_n \in \mathcal{H}_n$ bi-tracial vectors, such that $\|x\xi_n - \xi_n x\| \to 0$ for all $x \in M$, then we have $\|\xi_n -$

$P_0(\xi_n)\| \to 0$, where P_0 is the projection onto the space of central vectors. This is independent of the normal faithful trace τ [Pop06, Proposition 4.1]. Property (T) was first introduced in the factor case by Connes and Jones [CJ85] where they showed that for an ICC group Γ, the group von Neumann algebra $L\Gamma$ has property (T) if and only if Γ has Kazhdan's property (T) [Kaž67]. Their proof works equally well in the general case when Γ is not necessarily ICC.

We now suppose that M is finitely generated as a von Neumann algebra. Take $\{a_k\}_{k=1}^n \subset M$ a finite generating set such that $\sum_{k=1}^n a_k^* a_k = \sum_{k=1}^n a_k a_k^* = 1$, and let $\mathcal{B}(L^2(M,\tau)) \ni T \mapsto \varphi(T) = \sum_{k=1}^n \langle T\widehat{a_k^*}, \widehat{a_k^*}\rangle$ denote the associated normal regular hyperstate. For a fixed Hilbert bimodule $\mathcal{H}$ we define $\nabla_L, \nabla_R : \mathcal{H} \to \mathcal{H}^{\oplus n}$ by

$$\nabla_L(\xi) = \oplus a_k \xi$$

$$\nabla_R(\xi) = \oplus \xi a_k.$$

Note that we have

$$\|\nabla_L(\xi)\|^2 = \sum_{k=1}^n \|a_k\xi\|^2 = \left\langle \sum_{k=1}^n a_k^* a_k \xi, \xi \right\rangle = \|\xi\|^2,$$

and we similarly have

$$\|\nabla_R(\xi)\|^2 = \left\langle \sum_{k=1}^n \xi a_k a_k^*, \xi \right\rangle = \|\xi\|^2.$$

Thus ∇_L and ∇_R are both isometries. We let T denote the operator given by $T\xi = \sum_{k=1}^n a_k^* \xi a_k$. Note that $T = \nabla_L^* \nabla_R$ and hence T is a contraction.

Suppose now that $M \subset \mathcal{A}$ is an inclusion of von Neumann algebras and $\zeta \in \mathcal{A}_*$ is a faithful normal hyperstate. We may then consider the Hilbert space $L^2(\mathcal{A},\zeta)$ which is naturally a Hilbert M-bimodule where the left action is given by left multiplication $L(x)\hat{a} = \widehat{xa}$, and the right action is given by $R(x^{op}) = JL(x^*)J$. In this case the vector $\hat{1}$ is clearly left tracial, and we also have $Jx^*J\hat{1} = \Delta^{1/2}x\hat{1}$ from which it follows that $\hat{1}$ is also right tracial. If $\xi_0 \in L^2(\mathcal{A},\zeta)$ is a unit M-central vector, then $\tau_0(x) = \langle x\xi_0, \xi_0\rangle$ defines a normal trace on M. We let $s \in \mathcal{Z}(M)$ denote the support of τ_0.

Lemma 6.1 *Let (M,τ), φ, and $(\mathcal{A},\zeta)$ be as given above, then*

$$h_\varphi(M \subset \mathcal{A},\zeta) \geq -2\log\langle T1_\zeta, 1_\zeta\rangle.$$

Proof. Let $\Delta = \int_0^\infty \lambda d\lambda$ be the spectral resolution of the modular operator and let $\Delta_m = \int_{1/m}^m \lambda d\lambda$, $m \geq 1$ be the truncations. Let $\mu_k = \tau(a_k^* a_k)$ and $b_k = \mu_k^{-1/2} a_k$, for $k = 1, 2, \cdots, n$. Note that $\sum_{k=1}^n \mu_k = 1$. Also note that $L_{a_k^*} R_{a_k} 1_\zeta = a_k^* \Delta^{1/2} a_k 1_\zeta$. Now,

$$\begin{aligned}
-2\log\langle T1_\zeta, 1_\zeta\rangle &= -2\log(\sum_{k=1}^n \langle a_k^*\Delta^{1/2}a_k\hat{1},\hat{1}\rangle) = -2\lim_{m\to\infty}\log(\sum_{k=1}^n \mu_k\langle b_k^*\Delta_m^{1/2}b_k\hat{1},\hat{1}\rangle) \\
&\leq -2\lim_{m\to\infty}\sum_{k=1}^n \mu_k\log(\langle b_k^*\Delta_m^{1/2}b_k\hat{1},\hat{1}\rangle) \leq -2\lim_{m\to\infty}\sum_{k=1}^n \mu_k\langle b_k^*\log(\Delta_m^{1/2})b_k\hat{1},\hat{1}\rangle \\
&= -\lim_{m\to\infty}\sum_{k=1}^n \langle a_k^*\log(\Delta_m)a_k\hat{1},\hat{1}\rangle = h_\varphi(M \subset \mathcal{A},\zeta),
\end{aligned}$$

where the second inequality follows from Jensen's operator inequality. □

Theorem 6.2 *Let M be a II_1 factor generated as a von Neumann algebra by $\{a_k\}_{k=1}^n$ such that $\sum_{k=1}^n a_k^* a_k = \sum_{k=1}^n a_k a_k^* = 1$. Let*

$$\mathcal{B}(L^2(M,\tau)) \ni T \mapsto \varphi(T) = \sum_{k=1}^n \langle T\widehat{a_k^*}, \widehat{a_k^*}\rangle$$

denote the associated normal regular hyperstate. If M has property (T), then there exists $c > 0$ such that if $M \subset \mathcal{A}$ is any irreducible inclusion having no normal conditional expectation from $\mathcal{A}$ to M, and if $\zeta \in \mathcal{A}_*$ is any faithful normal hyperstate, then $h_\varphi(M \subset \mathcal{A}, \zeta) \geq c$.

Proof. Suppose M has property (T) and there is a sequence of irreducible inclusions $M \subset \mathcal{A}_m$, and normal faithful hyperstates $\zeta_m \in \mathcal{A}_m$, such that $h_\varphi(M \subset \mathcal{A}_m, \zeta_m) \to 0$. Then by Lemma 6.1 we have that $\langle T1_{\zeta_m}, 1_{\zeta_m}\rangle \to 1$, and hence $\sum_{k=1}^n \|a_k 1_{\zeta_m} - 1_{\zeta_m} a_k\|_2^2 = 2 - 2\langle T1_{\zeta_m}, 1_{\zeta_m}\rangle \to 0$. Since M has property (T) it then follows that for m large enough there exists a unit M-central vector $\xi \in L^2(\mathcal{A}_m, \zeta_m)$. If we let $\tilde{\zeta}$ denote the state on $\mathcal{A}_m$ given by $\tilde{\zeta}(a) = \langle a\xi, \xi\rangle$, then as ξ is M-central we have that $\tilde{\zeta}$ gives an M-hypertrace on $\mathcal{A}_m$. Thus, there exists a corresponding normal conditional expectation form $\mathcal{A}_m$ to M, for all m large enough. □

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7. Appendix: Minimal dilations and boundaries of u.c.p. maps

We include in this appendix a proof of Izumi's result from [Izu02] that for a von Neumann algebra (or even an arbitrary C^*-algebra) A, and a u.c.p. map $\phi : A \to A$, the operator space $\mathrm{Har}(A, \phi)$ has a C^*-algebraic structure. We take the approach in [Izu12] where $\mathrm{Har}(A, \phi)$ is shown to be completely isometric to the $*$-algebra of fixed points associated to a $*$-endomorphism which dialates the u.c.p. map. There are several proofs of the existence of such a dilation, the first proof is by Bhat in [Bha99] in the setting of completely positive semigroups, building on work from [Bha96], [BP94], and [BP95], and then later proofs were given in [BS00], [MS02], and Chapter 8 of [Arv03]. Our reason for including an additional proof is that it is perhaps more elementary than previous proofs, being based on a simple idea of iterating the Stinespring dilation [Sti55].

Lemma 7.1 *If $\mathcal{H}$ and $\mathcal{K}$ are Hilbert spaces, and $V : \mathcal{H} \to \mathcal{K}$ is a partial isometry, then for $A \subset \mathcal{B}(\mathcal{H})$, $B \subset \mathcal{B}(\mathcal{K})$, we have that $V^* *\text{-alg}(VBV^*, A)V = *\text{-alg}(B, V^*AV)$.*

Proof. Using the fact that $V^*V = 1$, this follows easily by induction on the length of alternating products for monomials in VBV^*, and A. □

If $A_0 \subset \mathcal{B}(\mathcal{H}_0)$ is a C^*-algebra, and $\phi : A_0 \to A_0$ is a unital completely positive map, then one can iterate Stinespring's dilation as follows:

Lemma 7.2 *Suppose $A_0 \subset \mathcal{B}(\mathcal{H}_0)$ is a unital C^*-algebra, and $\phi_0 : A_0 \to A_0$ is a unital completely positive map. Then there exists a sequence whose entries consist of:*

(1) a Hilbert space $\mathcal{H}_n$;
(2) an isometry $V_n : \mathcal{H}_{n-1} \to \mathcal{H}_n$;
(3) a unital C^*-algebra $A_n \subset \mathcal{B}(\mathcal{H}_n)$;
(4) a unital representation $\pi_n : A_{n-1} \to \mathcal{B}(\mathcal{H}_n)$, such that $\pi_n(A_{n-1})$, and $V_n A_{n-1} V_n^*$ generate A_n;
(5) a unital completely positive map $\phi_n : A_n \to A_n$;

such that the following relationships are satisfied for each $n \in \mathbf{N}$, $x \in A_{n-1}$:

$$V_n^* \pi_n(x) V_n = \phi_{n-1}(x); \tag{A}$$
$$V_n^* A_n V_n = A_{n-1}; \tag{B}$$
$$\phi_n(\pi_n(x)) = \pi_n(\phi_{n-1}(x)); \tag{C}$$
$$\pi_{n+1}(V_n x V_n^*) = V_{n+1} \pi_n(x) V_{n+1}^*. \tag{D}$$

Moreover, for each $n \in \mathbf{N}$ we have that the central support of $V_n V_n^*$ in A_n'' is 1. Also, if A_0 is a von Neumann algebra and ϕ_0 is normal then A_n will also be a von Neumann algebra and π_n and ϕ_n will be normal for each $n \in \mathbf{N}$.

Proof. We will first construct the objects and show the relationships (A), (B), and (C) by induction, with the base case being vacuous, and we will then show that (D) also holds for all $n \in \mathbf{N}$. So suppose $n \in \mathbf{N}$ and that (A), (B), and (C) hold for all $m < n$, (we leave V_0 undefined).

From the proof of Stinespring's Dilation Theorem we may construct a Hilbert space $\mathcal{H}_n$ by separating and completing the vector space $A_{n-1} \otimes \mathcal{H}_{n-1}$ with respect to the non-negative definite sesquilinear form satisfying

$$\langle a \otimes \xi, b \otimes \eta \rangle = \langle \phi_{n-1}(b^* a)\xi, \eta \rangle,$$

for all $a, b \in A_{n-1}$, $\xi, \eta \in \mathcal{H}_{n-1}$.

We also obtain a partial isometry $V_n : \mathcal{H}_{n-1} \to \mathcal{H}_n$ from the formula

$$V_n(\xi) = 1 \otimes \xi,$$

for $\xi \in \mathcal{H}_{n-1}$.

We obtain a representation $\pi_n : A_{n-1} \to \mathcal{B}(\mathcal{H}_n)$ (which is normal when A_0 is a von Neumann algebra and ϕ_0 is normal) from the formula

$$\pi_n(x)(a \otimes \xi) = (xa) \otimes \xi,$$

for $x, a \in A_{n-1}$, $\xi \in \mathcal{H}_{n-1}$. And recall the fundamental relationship $V_n^* \pi_n(x) V_n = \phi_{n-1}(x)$ for all $x \in A_{n-1}$, which establishes (A).

If we let A_n be the C^*-algebra generated by $\pi_n(A_{n-1})$ and $V_n A_{n-1} V_n^*$, then $\pi_n : A_{n-1} \to A_n$, and from Lemma 7.1 we have that $V_n^* A_n V_n$ is generated by $V_n^* \pi_n(A_{n-1}) V_n$ and A_{n-1}. However, $V_n^* \pi_n(A_{n-1}) V_n = \phi_{n-1}(A_{n-1}) \subset A_{n-1}$, hence $V_n^* A_n V_n = A_{n-1}$, establishing (B). Also, when A_0 is a von Neumann algebra and π_n is normal it then follows easily that A_n is then also a von Neumann algebra.

Also note that $\pi_n(A_{n-1}) V_n V_n^* \mathcal{H}_n$ is dense in $\mathcal{H}_n$, and so since $\pi_n(A_{n-1}) \subset A_n$ we have that the central support of $V_n V_n^*$ in A_n'' is 1.

We then define $\phi_n : A_n \to A_n$ by $\phi_n(x) = \pi_n(V_n^* x V_n)$, for $x \in A_n$. This is well defined since $V_n^* A_n V_n = A_{n-1}$, unital, and completely positive. Note that for $x \in A_{n-1}$ we have $\phi_n(\pi_n(x)) =$

$\pi_n(V_n^*\pi_n(x)V_n) = \pi_n(\phi_{n-1}(x))$, establishing (C).

Having established (A), (B), and (C) for all $n \in \mathbf{N}$, we now show that (D) holds as well. For this, notice first that for $a, b \in A_n$, $x \in A_{n-1}$, and $\xi, \eta \in \mathcal{H}_n$ we have

$$\begin{aligned}\langle \pi_{n+1}(V_n x V_n^*)(a \otimes \xi), b \otimes \eta\rangle &= \langle V_n x V_n^* a \otimes \xi, b \otimes \eta\rangle \\ &= \langle \phi_n(b^* V_n x V_n^* a)\xi, \eta\rangle \\ &= \langle \pi_n(V_n^* b^* V_n x V_n^* a V_n)\xi, \eta\rangle \\ &= \langle 1 \otimes \pi_n(x V_n^* a V_n)\xi, b \otimes \eta\rangle.\end{aligned}$$

Setting $x = 1$ and using that $V_{n+1}^*(1 \otimes \zeta) = \zeta$ for each $\zeta \in \mathcal{H}_n$, we see that

$$\begin{aligned}(V_{n+1}V_{n+1}^*)\pi_{n+1}(V_n V_n^*)(a \otimes \xi) &= (V_{n+1}V_{n+1}^*)(1 \otimes \pi_n(V_n^* a V_n)\xi) \\ &= 1 \otimes \pi_n(V_n^* a V_n)\xi \\ &= \pi_{n+1}(V_n V_n^*)(a \otimes \xi),\end{aligned}$$

and hence $\pi_{n+1}(V_n V_n^*) \leq V_{n+1}V_{n+1}^*$. If instead we set $a = 1$ then we have

$$V_{n+1}\pi_n(x)\xi = 1 \otimes \pi_n(x)\xi = \pi_{n+1}(V_n x V_n^*)V_{n+1}\xi,$$

and so $V_{n+1}\pi_n(x) = \pi_{n+1}(V_n x V_n^*)V_{n+1}$. Multiplying on the right by V_{n+1}^* and using that $\pi_n(V_n V_n^*) \leq V_{n+1}V_{n+1}^*$ then gives $V_{n+1}\pi_n(x)V_{n+1}^* = \pi_{n+1}(V_n x V_n^*)$. □

Theorem 7.3 (Bhat [Bha99]) *Let $A_0 \subset \mathcal{B}(\mathcal{H}_0)$ be a unital C^*-algebra, and $\phi_0 : A_0 \to A_0$ a unital completely positive map. Then there exists*

(1) a Hilbert space $\mathcal{K}$;
(2) an isometry $W : \mathcal{H}_0 \to \mathcal{K}$;
(3) a C^*-algebra $B \subset \mathcal{B}(\mathcal{K})$;
(4) a unital $*$-endomorphism $\alpha : B \to B$;

such that $W^*BW = A_0$, and for all $x \in A_0$ we have

$$\phi_0^k(x) = W^*\alpha^k(WxW^*)W.$$

Moreover, we have that the central support of WW^* in B'' is 1, $\alpha^k(WW^*) \leq \alpha^{k+1}(WW^*)$, and for $y \in \mathcal{B}(\mathcal{K})$ we have $y \in B$ if and only if $\alpha^k(WW^*)y\alpha^k(WW^*) \in \alpha^k(WA_0W^*)$ for all $k \geq 0$. Also, if A_0 is a von Neumann algebra and ϕ_0 is normal then B will also be a von Neumann algebra, and α will also be normal.

Proof. Using the notation from Lemma 7.2, we may define a Hilbert space $\mathcal{K}$ as the directed limit of the Hilbert spaces $\mathcal{H}_n$ with respect to the inclusions $V_{n+1} : \mathcal{H}_n \to \mathcal{H}_{n+1}$. We denote by $W_n : \mathcal{H}_n \to \mathcal{K}$ the associated sequence of isometries satisfying $W_{n+1}^*W_n = V_{n+1}$, for $n \in \mathbf{N}$, and we set $P_n = W_nW_n^*$, an increasing sequence of projections.

From (B) we have that $P_{n-1}W_nA_nW_n^*P_{n-1} = W_{n-1}A_{n-1}W_{n-1}^*$, and hence if we define the C^*-algebra $B = \{x \in \mathcal{B}(\mathcal{K}) \mid W_n^*xW_n \in A_n, n \geq 0\}$, then we have $W_n^*BW_n = A_n$, for all $n \geq 0$. Also, if A_0 is a von Neumann algebra, then so is A_n for each $n \in \mathbf{N}$ and from this it follows easily that B is also a von Neumann algebra.

We define the unital $*$-endomorphism $\alpha : B \to B$ (which is normal when A_0 is a von Neumann algebra and ϕ_0 is normal) by the formula

$$\alpha(x) = \lim_{n\to\infty} W_{n+1}\pi_{n+1}(W_n^*xW_n)W_{n+1},$$

where the limit is taken in the strong operator topology. Note that $\alpha(P_n) = P_{n+1} \geq P_n$. From (D) we see that in general, the strong operator topology limit exists in B, and that for $x \in A_n \cong P_n A_\infty P_n$ the limit stabilizes as $\alpha(W_n x W_n^*) = W_{n+1}\pi_{n+1}(x)W_{n+1}^*$.

From (A) we see that for $n \geq 0$, and $x \in A_n$ we have

$$\begin{aligned} P_n\alpha(W_n x W_n^*)P_n &= W_n W_n^* W_{n+1}\pi_{n+1}(x)W_{n+1}^* W_n W_n^* \\ &= W_n V_{n+1}^*\pi_{n+1}(x)V_{n+1}W_n^* \\ &= W_n\phi_n(x)W_n^*. \end{aligned}$$

By induction we then see that also for $k > 1$, and $x \in A_0$ we have

$$\begin{aligned} P_0\alpha^k(W_0 x W_0^*)P_0 &= P_0\alpha^{k-1}(P_0\alpha(W_0 x W_0^*)P_0)P_0 \\ &= P_0\alpha^{k-1}(W_0\phi_0(x)W_0^*)P_0 \\ &= W_0\phi_0^k(x)W_0^*. \end{aligned}$$

By the previous Lemma we have that the central support of P_n in $W_n A_n'' W_n^*$ is P_{n+1}. Hence it follows that the central support of P_0 in B is 1. □

7.1 Poisson boundaries of u.c.p. maps

If $A \subset \mathcal{B}(\mathcal{H})$ is a unital C^*-algebra, and $\phi : A \to A$ a unital completely positive map, then a projection $p \in A$ is said to be coinvariant, if $\{\phi^n(p)\}_n$ defines an increasing sequence of projections which strongly converge to 1 in $\mathcal{B}(\mathcal{H})$, and such that for $y \in \mathcal{B}(\mathcal{H})$ we have $y \in A$ if and only if $\phi^n(p)y\phi^n(p) \in A$ for all $n \geq 0$. Note that for $n \geq 0$, $\phi^n(p)$ is in the multiplicative domain for ϕ, and is again coinvariant. We define $\phi_p : pAp \to pAp$ to be the map $\phi_p(x) = p\phi(x)p$, then ϕ_p is normal unital completely positive. Moreover, we have that $\phi_p^k(x) = p\phi^k(x)p$ for all $x \in pAp$, which can be seen by induction from

$$p\phi^k(x)p = p\phi^{k-1}(p)\phi^k(x)\phi^{k-1}(p)p = p\phi^{k-1}(\phi_p(x))p.$$

Theorem 7.4 (Prunaru [Pru12]) *Let $A \subset \mathcal{B}(\mathcal{H})$ be a unital C^*-algebra, $\phi : A \to A$ a unital completely positive map, and $p \in A$ a coinvariant projection. Then the map $\mathcal{P} : \mathrm{Har}(A, \phi) \to \mathrm{Har}(pAp, \phi_p)$ given by $\mathcal{P}(x) = pxp$ defines a completely positive isometric surjection, between $\mathrm{Har}(A, \phi)$ and $\mathrm{Har}(pAp, \phi_p)$.*

Moreover, if A is a von Neumann algebra and ϕ is normal then $\mathcal{P}$ is also normal.

Proof. First note that $\mathcal{P}$ is well-defined since if $x \in \mathrm{Har}(A, \phi)$ we have

$$\phi_p(pxp) = p\phi(p)x\phi(p)p = pxp.$$

Clearly $\mathcal{P}$ is completely positive (and normal in the case when A is a von Neumann algebra and ϕ is normal).

To see that it is surjective, if $x \in \mathrm{Har}(pAp, \phi_p)$ then consider the sequence $\phi^n(x)$. For each $m, n \geq 0$, we have

$$\phi^m(p)\phi^{m+n}(x)\phi^m(p) = \phi^m(p\phi^n(x)p) = \phi^m(\phi_p^n(x)) = \phi^m(x).$$

It follows that $\{\phi^n(x)\}_n$ is eventually constant for any ξ in the range of $\phi^m(p)$ for any m. Since $\{\phi^n(x)\}_n$ is uniformly bounded and $\{\phi^n(x)\xi\}_n$ converges for a dense subset of $\xi \in \mathcal{H}$ we then have that $\{\phi^n(x)\}_n$ converges in the strong operator topology to an element $y \in \mathcal{B}(\mathcal{H})$ such that $\phi^m(p)y\phi^m(p) = \phi^m(x)$ for each $m \geq 0$. Consequently we have $y \in A$.

In particular, for $m = 0$ we have $pyp = x$. To see that $y \in \mathrm{Har}(A, \phi)$ we use that for all $z \in A$ we have the strong operator topology limit

$$\lim_{n\to\infty} \phi(\phi^n(p) z \phi^n(p)) = \lim_{n\to\infty} \phi^{n+1}(p)\phi(z)\phi^{n+1}(p) = \phi(z),$$

and hence

$$\phi(y) = \lim_{m\to\infty} \phi(\phi^m(p) y \phi^m(p)) = \lim_{m\to\infty} \phi^{m+1}(x) = y.$$

Thus $\mathcal{P}$ is surjective, and since $\phi^n(p)$ converges strongly to 1, and each $\phi^n(p)$ is in the multiplicative domain of ϕ, it follows that if $x \in \mathrm{Har}(A, \phi)$ then $\phi^n(pxp)$ converges strongly to x and hence

$$\|x\| = \lim_{n\to\infty} \|\phi^n(pxp)\| \leq \|pxp\| \leq \|x\|.$$

Thus, $\mathcal{P}$ is also isometric. □

Corollary 7.5 (Izumi [Izu02]) *Let A be a unital C^*-algebra, and $\phi : A \to A$ a unital completely positive map. Then there exists a C^*-algebra B and a completely positive isometric surjection $\mathcal{P} : B \to \mathrm{Har}(A, \phi)$.*

Moreover B and $\mathcal{P}$ are unique in the sense that if $\tilde{B}$ is another C^*-algebra, and $\mathcal{P}_0 : \tilde{B} \to \mathrm{Har}(A, \phi)$ is a completely positive isometric surjection, then $\mathcal{P}^{-1} \circ \mathcal{P}_0$ is an isomorphism.

Also, if A is a von Neumann algebra and ϕ is normal, then B is also a von Neumann algebra and $\mathcal{P}$ is normal.

Proof. Note that we may assume $A \subset \mathcal{B}(\mathcal{H})$. Existence then follows by applying the previous theorem to Bhat's dilation. Uniqueness follows from [Cho74] □

Corollary 7.6 (Choi-Effros [CE77]) *Let A be a unital C^*-algebra and $F \subset A$ an operator system. If $E : A \to F$ is a completely positive map such that $E_{|F} = \mathrm{id}$, then F has a unique C^*-algebraic structure which is given by $x \cdot y = E(xy)$. Moreover, if A is a von Neumann algebra and F is weakly closed then this gives a von Neumann algebraic structure on F.*

Proof. Note that $F \subseteq \mathrm{Har}(A, E)$, as $E_{|F} = \mathrm{id}$. Since the range of E is contained in F, we get $\mathrm{Har}(A, E) = F$.

When A is a C^*-algebra this follows from Corollary 7.5 since $\mathrm{Har}(A, E) = F$. Also note that since $E^n = E$ it follows from the proof of Theorem 7.4 that the product structure coming from the Poisson boundary is given by $x \cdot y = E(xy)$.

If A is a von Neumann algebra and F is weakly closed then F has a predual $F_\perp = \{\varphi \in A_* \mid \varphi(x) = 0, \text{ for all } x \in F\}$ and hence A is isomorphic to a von Neumann algebra by Sakai's theorem. □

Proposition 7.7 *Let A be an abelian C^*-algebra and $\phi : A \to A$ a normal unital completely positive map. Then the Poisson boundary of ϕ is also abelian.*

Proof. Let B be the Poisson boundary of ϕ, and let $\mathcal{P} : B \to \mathrm{Har}(A, \phi)$ be the Poisson transform. If C is a C^*-algebra and $\psi : C \to B$ is a positive map then $\mathcal{P} \circ \psi : C \to \mathrm{Har}(A, \phi) \subset A$ is positive, and since A is abelian it is then completely positive. Hence, ψ is also completely positive. Since every positive map from a C^*-algebra to B is completely positive it then follows that B is abelian. □

Example 7.8 *Let Γ be a discrete group and $\mu \in \mathrm{Prob}(\Gamma)$ a probability measure on Γ such that the support of μ generates Γ. Then on $\ell^\infty\Gamma$ we may consider the normal unital (completely) positive* *map* ϕ_μ *given by* $\phi_\mu(f) = \mu * f$*, where* $\mu * f$ *is the convolution* $(\mu * f)(x) = \int f(g^{-1}x)\, d\mu(g)$*. Then* $\mathrm{Har}(\mu) = \mathrm{Har}(\ell^\infty\Gamma, \phi_\mu)$ *has a unique von Neumann algebraic structure which is abelian by the previous proposition. Notice that* Γ *acts on* $\mathrm{Har}(\mu)$ *by right translation, and since this action preserves positivity it follows from [Cho74] that* Γ *preserves the multiplication structure as well.*

Since the support of μ *generates* Γ*, for a non-negative function* $f \in \mathrm{Har}(\mu)_+$*, we have* $f(e) = 0$ *if and only if* $f = 0$*. Thus we obtain a natural normal faithful state* φ *on* $\mathrm{Har}(\mu)$ *which is given by* $\varphi(f) = f(e)$*.*

Since φ *is* Γ*-equivariant, this extends to a normal u.c.p. map* $\tilde{\varphi} : \ell^\infty\Gamma \rtimes \Gamma \to \ell^\infty\Gamma \rtimes \Gamma$ *such that* $\tilde{\varphi}_{L\Gamma} = \mathrm{id}$*. Note that* $\ell^\infty\Gamma \rtimes \Gamma \cong \mathcal{B}(\ell^2\Gamma)$*. It is an easy exercise to see that the Poisson boundary of* $\tilde{\varphi}$ *is nothing but the crossed product* $\mathrm{Har}(\mu) \rtimes \Gamma$*.*

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