



On the Parabolic Boundary Harnack Principle

Daniela De Silva¹ · Ovidiu Savin¹

Received: 16 March 2021 / Revised: 16 March 2021 / Accepted: 5 May 2021 /

Published online: 12 November 2021

© Springer Science+Business Media LLC, part of Springer Nature 2021

Abstract

We investigate the parabolic Boundary Harnack Principle by the analytical methods developed in De Silva and Savin (J Differ Equ 3(15):2419–2429, 2020; J Math Eng (in press)). Besides the classical case, we deal with less regular space-time domains, including slit domains.

Keywords Parabolic equations · Harnack inequality · Comparison principle · Regularity · Hölder domains · Slit domains

1 Introduction

1.1 Statement of Main Results

In this paper, we provide direct analytical proofs of the parabolic Boundary Harnack Inequality for both divergence and non-divergence type operators, in several different settings. Our strategy is based on our earlier works [8, 9] where the elliptic counterparts of these results were obtained. In order to state our theorems precisely, we introduce some notation.

We denote by $\Gamma \subset \mathbb{R}^{n+1}$ the graph of a continuous function $g(x', t)$ in the x_n direction,

$$\Gamma := \{x_n = g(x', t)\}, \quad (0, 0) \in \Gamma,$$

while \mathcal{C}_r denotes the cylinder of size r on top of Γ (in the e_n direction) i.e.,

✉ Daniela De Silva
desilva@math.columbia.edu
Ovidiu Savin
savin@math.columbia.edu

¹ Department of Mathematics, Barnard College, Columbia University, New York, NY 10027, USA

$$C_r := \{(x', x_n, t) | x' \in B'_r, \quad t \in (-r^2, r^2), \quad g(x', t) < x_n < g(x', t) + r\}.$$

As usual, $x' = (x_1, \dots, x_{n-1})$, while $B'_r \subset \mathbb{R}^{n-1}$ is the ball of radius r centered at the origin.

We consider solutions $u(x, t)$ to the parabolic equation

$$u_t = Lu \quad \text{in } C_1,$$

where $Lu = \operatorname{tr}(A(x)D^2u)$ or $L(u) = \operatorname{div}(A(x)\nabla u)$, with A satisfying,

$$\lambda I \leq A \leq \Lambda I, \quad 0 < \lambda \leq \Lambda < +\infty.$$

First, we recall the standard boundary Harnack inequalities for parabolic equations in Lipschitz domains (Fig. 1). References to known literature will be provided in the next subsection. Here $g \in C_{x',t}^{\alpha,\beta}$ if

$$|g(x', t) - g(y', s)| \leq C(|x' - y'|^\alpha + |t - s|^\beta),$$

and \bar{E} , \underline{E} are points interior to C_1 at times $t = 1/2$ and $t = -1/2$, respectively,

$$\bar{E} = \left(\left(g\left(0, \frac{1}{2}\right) + \frac{1}{2} \right) e_n, \frac{1}{2} \right), \quad \underline{E} = \left(\left(g\left(0, -\frac{1}{2}\right) + \frac{1}{2} \right) e_n, -\frac{1}{2} \right).$$

Theorem 1.1 ($C^{1,\frac{1}{2}}$ domains) Assume that $g \in C_{x',t}^{1,\frac{1}{2}}$ and u, v are two positive solutions to

$$u_t = Lu, \quad v_t = Lv \quad \text{in } C_1,$$

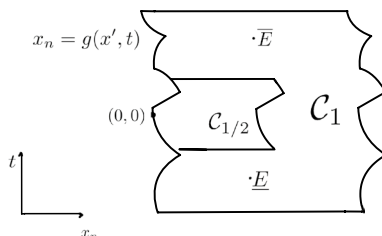
with u vanishing continuously on Γ . Then

$$\frac{u}{v}(x) \leq C \frac{u(\bar{E})}{v(\underline{E})} \quad \text{for all } x \in C_{1/2}, \quad (1.1)$$

with C depending only on $n, \|g\|_{C^{1,1/2}}, \lambda$, and Λ .

In this note, we provide new versions of Theorem 1.1 in more general Hölder domains.

Fig. 1 Theorem 1.1



Theorem 1.2 ($C^{\frac{1}{2}+, \frac{1}{3}+}$ domains) *Theorem 1.1 holds if $g \in C_{x',t}^{\alpha,\beta}$ with $\alpha > 1/2$, $\beta > 1/3$.*

In the case of the heat operator, we may lower further the space regularity of g to any exponent $\alpha > 0$ provided that we have a $1/2$ Hölder modulus of continuity in time (from one-side).

Theorem 1.3 ($C^{\alpha, \frac{1}{2}}$ domains) *Theorem 1.1 holds for the heat equation if $g \in C_{x',t}^{\alpha, \frac{1}{2}}$ with $\alpha > 0$.*

We remark that the only property of the heat equation needed in the proof of Theorem 1.3 is the translation invariance with respect to the x_n, t variables. Hence, the theorem holds also for operators L with coefficients depending only on the x' variable.

Next we state a result in slit domains, that is the case when the equations are satisfied in the complement of a thin set $S \subset \mathbb{R}^{n+1}$ included in a lower dimensional subspace. This case is relevant, for example, in the time-dependent Signorini problem.

Precisely, we assume that S is a closed set and

$$S \subset \{x_n = 0\},$$

and in this case

$$C_r = B_r \times (-r^2, r^2), \quad \bar{E} = (1/2e_n, 1/2), \quad \underline{E} = (1/2e_n, -1/2).$$

With these notation, we state our theorem.

Theorem 1.4 (Thin Parabolic Boundary Harnack) *If u, v are two positive solutions even in the x_n variable,*

$$u_t = Lu \quad v_t = Lv, \quad \text{in } C_1 \setminus S,$$

and u vanishes on S , then (1.1) holds.

The assumption that u, v are even in the x_n variable can be removed provided that $C_1 \setminus S$ contains a ball of radius σ centered on $\{x_n = 0\}$, and the constant C in estimate (1.1) depends on σ .

We remark that in Theorems 1.2–1.4 whenever the boundary of the domain contains non-regular points for the Dirichlet problem, the statement that u vanishes on it is interpreted in the sense that u is the limit of a sequence of continuous subsolutions which vanish on it.

1.2 Known Literature

For the last 50 years, the boundary Harnack principle has played an essential role in analysis and PDEs in a variety of contexts. The available literature on this topic is

very rich and we collect here only the crucial results, making no attempt to discuss the countless important applications of this fundamental tool.

1.2.1 Elliptic Case

In the elliptic context, the classical Boundary Harnack Principle, that is the case when g is Lipschitz continuous, states the following. Here the notation is the same as above, with u, v, g independent on t .

Theorem 1.5 *Let $u, v > 0$ satisfy $Lu = Lv = 0$ in C_1 and vanish continuously on Γ . Assume u, v are normalized so that $u(e_n/2) = v(e_n/2) = 1$, then*

$$C^{-1} \leq \frac{u}{v} \leq C, \quad \text{in } C_{1/2}, \quad (1.2)$$

with C depending on n, λ, Λ , and the norm of g .

The case when $L = \Delta$ first appears in [1, 7, 18, 26]. Operators in divergence form were then considered in [6], while the case of operator in non-divergence form was treated in [10]. The same result for operators in divergence form was extended also to the so-called NTA domains in [17]. The case of Hölder domains and L in divergence form was addressed with probabilistic techniques in [2, 3], and an analytic proof was then provided in [13]. For Hölder domains and operators L in non-divergence form, it is necessary that the domain is $C^{0,\alpha}$ with $\alpha > 1/2$ or that it satisfies a uniform density property, and this was first established again using a probabilistic approach [5].

In [8, 9], we presented a unified analytic proof the Boundary Harnack Principle that does not make use of the Green's function and which holds for both operators in non-divergence and in divergence form. The idea is to find an “almost positivity property” of a solution, which can be iterated from scale 1 to all smaller scales (some similar ideas were also used in [20, 23] to treat non-divergence equations with unbounded drift). This strategy successfully applies to other similar situations like that of Hölder domains, NTA domains, and to the case of slit domains, providing a unified approach to a large class of results.

1.2.2 Parabolic Case

For parabolic equations, the situation is more complicated, essentially due to the evolution nature of the latter which is reflected in a time-lag in the Harnack Principle. For operators in divergence form, the parabolic boundary Harnack principle in Theorem 1.1 is due to [11, 19, 24]. In the case of operators in non-divergence form in cylinders with C^2 cross sections, Theorem 1.1 was settled in [15], where the author also derived a Carleson estimate (see Lemma 2.6) in Lipschitz domains. The statement of Theorem 1.1 in Lipschitz domain was later obtained in [12], which is (to the authors knowledge) the first instance in which a boundary Harnack type result in Lipschitz domains is obtained without the aid of Green's functions (and it is probably the inspiration for the later works in the elliptic context [20, 23]). In

[16], Theorem 1.1 was also shown to hold for unbounded parabolically Reifenberg flat domains. In the context of time-independent Hölder domains, a result in the spirit of Theorem 1.2 was obtained via probabilistic techniques in [4]. The result in Theorem 1.3 is completely novel. Concerning slit domains, in the case when S is the subgraph of a parabolic Lipschitz graph, the thin-version Theorem 1.4 was established by [22]. Again, our strategy provides a unified approach for a variety of contexts.

1.3 Organization of the Paper

The paper is organized as follows. In Sect. 2, after recalling some standard results, we provide the proof of Theorems 1.1 and 1.4. The key “almost positivity” property to be iterated from scale 1 to all smaller scales is obtained in Lemma 2.5. The following section deals with Hölder domains and the proof of Theorem 1.2, which relies on the same strategy as Theorem 1.1, though the proof of the Carleson estimate in the Hölder setting requires a more involved argument similar to the one in the proof of Lemma 2.5. Section 4 contains the proof of Theorem 1.3, which is based on refined versions of the weak Harnack inequality (see Lemmas 4.2–4.4).

2 Proof of Theorems 1.1 and 1.4

In this section, we provide the proof of the classical result Theorem 1.1 and the novel result Theorem 1.4. We start by collecting standard known Harnack type inequalities. In the divergence setting, these results are due to [21], while in the non-divergence setting they follow from [25].

2.1 Weak Harnack Inequality

Denote by

$$Q_r := (-r, r)^n \times (-r^2, 0], \quad Q_r(x_0, t_0) := (x_0, t_0) + Q_r,$$

the parabolic cubes of size r . The parabolic boundary of Q_r is denoted by $\partial_p Q_r$ and is given by

$$\partial_p Q_r := (\partial(-r, r)^n \times (-r^2, 0)) \cup ((-r, r)^n \times \{-r^2\}).$$

Similarly,

$$Q'_r := (-r, r)^{n-1} \times (-r^2, 0], \quad Q'_r(x'_0, t_0) := (x'_0, t_0) + Q'_r.$$

Our main tools in establishing the boundary Harnack inequalities are the standard weak Harnack estimates. We recall the parabolic versions which as mentioned in the introduction differ from the elliptic counterparts due to the time-lag.

Theorem 2.1 (Supersolution) *If*

$$u_t \geq Lu \quad \text{and} \quad u \geq 0 \quad \text{in} \quad Q_1, \quad u(0, 0) = 1,$$

then

$$\int_{Q_{\frac{1}{2}}(0, -\frac{1}{2})} u^p \, dx dt \leq C,$$

for some $p > 0$ small, C large universal (i.e., dependent on n, λ, Λ).

Theorem 2.2 (Subsolution) *If*

$$u_t \leq Lu \quad \text{and} \quad u \geq 0 \quad \text{in} \quad Q_1,$$

then

$$u(0, 0) \leq C(p) \|u\|_{L^p(Q_1)},$$

for any $p > 0$.

The classical (backward) Harnack inequality then reads as follows.

Theorem 2.3 (Harnack inequality) *If*

$$u_t = Lu \quad \text{and} \quad u \geq 0 \quad \text{in} \quad Q_1,$$

then for c small universal (dependent on n, λ, Λ),

$$\min_{Q_{1/2}} u \geq c \max_{Q_{1/2}(0, -\frac{1}{2})} u.$$

Another useful version for the subsolution property is the following measure to pointwise estimate.

Theorem 2.4 (Subsolution) *If*

$$u_t \leq Lu \quad \text{and} \quad 1 \geq u \geq 0 \quad \text{in} \quad Q_2,$$

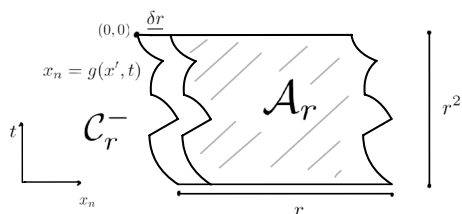
and for some $\delta > 0$,

$$|\{u = 0\} \cap Q_1(0, -1)| \geq \delta.$$

Then

$$u \leq 1 - c(\delta) \quad \text{in} \quad Q_1.$$

Fig. 2 The sets C_r^- and A_r



With these tools at hand, we are ready to provide in the following subsection our proof of the classical result in Theorem 1.1.

2.2 Proof of Theorem 1.1

In what follows, constants depending on n, λ, Λ , and the norm of g , are called universal.

We denote by

$$C_r^- := \{(x', x_n, t) \mid x' \in (-r, r)^{n-1}, \quad t \in (-r^2, 0], \quad g(x', t) < x_n < g(x', t) + r\},$$

the backward-in-time cylinder of size r on top (in the e_n direction) of the graph Γ of g . Also we set,

$$A_r := \{(x, t) \in C_r^- \mid g(x', t) + \delta r \leq x_n < g(x', t) + r\},$$

that is the collection of points in the cylinder C_r^- at height greater or equal than δr on top of Γ , for some $\delta > 0$ small, to be made precise later (Fig. 2).

The key tool for establishing the boundary Harnack estimates is the following iterative lemma. Later, we will apply this lemma for the difference $w = v - cu$ for some sufficiently small constant c , in order to obtain the desired claim in Theorem 1.1.

Lemma 2.5 *There exist universal constants $M, \delta > 0$, such that if w is a solution to*

$$w_t = Lw \quad \text{in } C_r^-,$$

(possibly changing sign) with w^- vanishing continuously on Γ ,

$$w \geq M \quad \text{in } A_r, \tag{2.1}$$

and

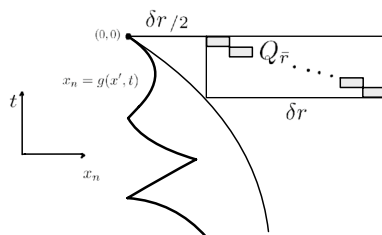
$$w \geq -1 \quad \text{in } C_r^-,$$

then,

$$w \geq Ma \quad \text{in } A_{\frac{r}{2}}, \tag{2.2}$$

and

Fig. 3 Proof of Lemma 2.5



$$w \geq -a \quad \text{in} \quad C_{\frac{r}{2}}^-, \quad (2.3)$$

for some small $a > 0$.

The conclusion can be iterated indefinitely and we obtain that if the hypotheses are satisfied in C_r^- then

$$w > 0 \quad \text{on the line segment} \{(se_n, 0), \quad 0 < s < r\}. \quad (2.4)$$

Proof We start by observing that any point $(x_0, t_0) \in \mathcal{A}_{r/2}$ can be connected through a chain of backwards-in-time adjacent parabolic cubes of size $\bar{r} := c_1 \delta r$ centered at

$$(x_j, t_j) := (x_0 + j\bar{r}e_n, t_0 - j\bar{r}^2),$$

to a last cube $Q_{\bar{r}}(x_m, t_m) \subset \mathcal{A}_r$ (see Fig. 3). Here c_1 is small depending on the $C_{x',t}^{1,1/2}$ norm of g so that

$$Q_{2\bar{r}}(x_j, t_j) \subset C_r^-,$$

and the number m of cubes depends only on c_1 . By Harnack inequality (Theorem 2.3) applied to $w + 1 \geq 0$, using assumption (2.1), we get

$$(w + 1)(x_0, t_0) \geq c^m(M + 1) \implies w \geq 1 \quad \text{in} \quad \mathcal{A}_{r/2},$$

provided that we choose M large depending on c_1 (and independent of δ). Hence (2.2) holds with $a = 1/M$.

To establish (2.3) with this choice of M , a , we first extend $w^- = 0$ in

$$Q'_r \times (\{x_n < g(x', t)\} \cup \{x_n > g(x', t) + \delta r\}),$$

so that w^- is a global subsolution in $Q'_r \times \mathbb{R}$ thanks to assumption (2.1). Then, for each cube $Q_{2\delta r}(x, t)$ satisfying $Q'_{2\delta r}(x', t) \subset Q'_r$, we have

$$|\{w^- = 0\} \cap Q_{2\delta r}(x, t)| \geq \frac{1}{2} |Q_{2\delta r}(x, t)|.$$

This is a consequence of the graph property of Γ . Indeed, for each fixed (x', t) , we consider the 1D line in the e_n direction. Any segment of length $2\delta r$ on this line has at least half of its length either in \mathcal{A}_r or in the complement of C_r^- .

By weak Harnack inequality, Theorem 2.4, as we remove the collection of cubes $Q_{2\delta r}(x, t)$ which are tangent to the parabolic boundary of $Q'_r \times \mathbb{R}$, the norm $\|w^-\|_{L^\infty}$ decays by a factor $1 - c$, $c > 0$ universal. Iterating this for $\sim 1/\delta$ times we find that

$$w^- \leq (1 - c)^{1/\delta} \quad \text{in } C_{r/2}^- \subset Q'_{r/2} \times \mathbb{R}.$$

We choose δ small, so that $w^- \leq a = M^{-1}$ and (2.3) holds. \square

A second ingredient in the proof of Theorem 1.1 is the following Carleson estimate which provides a bound for u in the cylinder $\mathcal{C}_{2/3}$.

Lemma 2.6 (Carleson estimate) *Let u, \bar{E} be as in Theorem 1.1, then*

$$\|u\|_{L^\infty(\mathcal{C}_{2/3})} \leq C u(\bar{E}),$$

with $C > 0$ universal.

Proof The Carleson estimate can be established by similar arguments as in the Lemma 2.5 above. We will use this approach in the case of Hölder domains in the next section. However, for $C^{1, \frac{1}{2}}_{x', t}$ domains, the Carleson estimate is a direct consequence of the weak Harnack inequality.

Indeed, assume that $u(\bar{E}) = 1$. Any point $(x_0, t_0) \in \mathcal{C}_{12/17}$ can be connected to \bar{E} by a chain of forward-in-time adjacent cubes $Q_{r_j}(x_j, t_j)$ included in \mathcal{C}_1 , with r_j proportional to the parabolic distance d_j from (x_j, t_j) to Γ . The number of cubes in this chain is proportional to $|\log d_0|$. By Harnack inequality,

$$u(x_0, t_0) \leq e^{C|\log d_0|} u(\bar{E}) \leq d_0^{-C'}.$$

This means that $\|u\|_{L^p} \leq C$ in $\mathcal{C}_{12/17}$ for some small $p > 0$ universal. The extension of u by 0 in $\Omega_r := ((-r, r)^{n-1} \times (-r^2, r^2)) \times \{x_n \leq g(x', t)\}$ is a subsolution, and now we can apply weak Harnack inequality Theorem 2.2 in cubes $Q_{c_0}(x, t) \subset \mathcal{C}_{12/17} \cup \Omega_r$ for $(x, t) \in \mathcal{C}_{2/3}$ and c_0 small universal, to obtain the desired conclusion. \square

We are now ready to combine the previous two lemmas and obtain the desired Theorem 1.1.

Proof of Theorem 1.1 We assume that $u(\bar{E}) = v(\underline{E}) = 1$ and define $w = C_1 v - c_1 u$. By Harnack inequality applied to v and the Carleson estimate for u , we can choose the constants C_1 large, c_1 small (depending on δ, M) such that w satisfies

$$w \geq -1 \quad \text{in } \mathcal{C}_{2/3}, \quad \text{and} \quad w(x, t) \geq M \quad \text{if } x_n \geq g(x', t) + \delta/4.$$

Then we can apply Lemma 2.5 in cylinders $\mathcal{C}_{1/6}^-$ around any point on $\Gamma \cap \mathcal{C}_{1/2}$, and conclude from (2.4) that $w > 0$ in $\mathcal{C}_{1/2}$. \square

2.3 Proof of Theorem 1.4

The proof is identical to the one of Theorem 1.1 after the appropriate modifications in the definitions of C_r^- and \mathcal{A}_r . Precisely,

$$C_r^- := Q_r \setminus S, \quad \mathcal{A}_r := Q_r \cap \{|x_n| \geq \delta r\}.$$

Lemma 2.5 applies for the difference $w = v - cu$. The hypotheses that u and $w^- = (cu - v)^+$ vanish on S are understood in the sense that each of them is obtained in C_1^- as a pointwise limit of an increasing sequence of continuous subsolutions in Q_1 which vanish on S . Notice that if u_n is such a sequence for u , then $(cu_n - v)^+$ is a corresponding sequence for w^- , (since $v \geq 0$ in C_1^-). Thus, the extensions of u and w^- by 0 on S are subsolutions in Q_1 , and Lemmas 2.5 and 2.6 hold as above. \square

3 Hölder Domains and the Proof of Theorem 1.2

In this section, we prove Theorem 1.2 by extending the arguments of the previous section to Hölder domains. We assume that for some $\alpha > \frac{1}{2}$,

$$[g]_{C^{\alpha, \frac{\alpha}{1+\alpha}}_{x', t}} \leq K, \quad (3.1)$$

for some constant K . Below, constants depending possibly on $n, \lambda, \Lambda, \alpha$ and K are called universal.

We define

$$C_r^- := \{(x', x_n, t) \mid x \in (-r, r)^n, \quad t \in (-r, 0], \quad g(x', t) < x_n < g(x', t) + r\},$$

and notice that here we took the time interval of C_r^- of size r instead of the natural parabolic scaling r^2 that we used in the previous section. This change is due to the fact that the norm of g is no longer left invariant by the parabolic scaling. We also define

$$\mathcal{A}_r := \{x \in C_r^- \mid g(x', t) + r^\beta \leq x_n < g(x', t) + r\},$$

the points in the cylinder C_r^- at height greater or equal than r^β on top of Γ , for some $\beta > 1$ to be made precise later.

Lemma 3.1 *Suppose (3.1) holds for C_r^- and let w be a solution to*

$$w_t = Lw \quad \text{in } C_r^-,$$

for which w^- vanishes on Γ . There exist universal constants $C_0, \beta > 0$ such that if

$$w \geq f(r) \quad \text{on } \mathcal{A}_r,$$

and

$$w \geq -1 \quad \text{on} \quad C_r^-,$$

where

$$f(r) := e^{C_0 r^\gamma}, \quad \gamma := \beta(1 - \frac{1}{\alpha}) < 0,$$

then,

$$w \geq f(\frac{r}{2}) a \quad \text{on} \quad \mathcal{A}_{\frac{r}{2}}, \quad (3.2)$$

and

$$w \geq -a \quad \text{on} \quad C_{\frac{r}{2}}^-, \quad (3.3)$$

for some small $a = a(r) > 0$, as long as $r \leq r_0$ universal.

Proof We adapt the argument of Lemma 2.5 in this case and sketch the details.

We connect a point $(x_0, t_0) \in \mathcal{A}_{r/2}$ (which is not in \mathcal{A}_r) to a point (x_m, t_m) with $x_m = x_0 + r^\beta e_n \in \mathcal{A}_r$ by a chain of adjacent backward-in-time cubes of size $\bar{r} := c_0 r^{\beta/\alpha}$. The number m of cubes depends on r , i.e.,

$$m \sim r^\beta / \bar{r} = c_0^{-1} r^{\beta(1 - \frac{1}{\alpha})} = c_0^{-1} r^\gamma.$$

All the cubes are included in the domain $(t_m := t_0 - m\bar{r}^2)$

$$\{(x - x_0) \cdot e_n \geq 0, \quad t \in [t_m, t_0], \quad (x - x_0)' \in [-\bar{r}, \bar{r}]^{n-1}\},$$

which by (3.1) is included in C_r^- since $m\bar{r}^2 \sim r^\beta \bar{r} = c_0 r^{\beta \frac{\alpha+1}{\alpha}}$, and c_0 is chosen small. Moreover, $Q_{\bar{r}}(x_m, t_m) \subset \mathcal{A}_r$, and Harnack inequality for $w + 1$ implies that

$$w + 1 \geq f(r) e^{-Cm} \geq 2 \quad \text{in} \quad \mathcal{A}_{r/2}, \quad (3.4)$$

where the last inequality is guaranteed if we choose C_0 sufficiently large.

For the second step which bounds w^- we use cylinders of size $2r^\beta$ (instead of $2\delta r$ as before) and get by the same argument as in the Lipschitz case

$$w^- \leq e^{-cr^{1-\beta}} =: a. \quad (3.5)$$

The conclusion follows since in $\mathcal{A}_{r/2}$, $w \geq 1 \geq f(r/2)a$, and in the last inequality we used $1 - \beta < \gamma$, provided that β is chosen sufficiently large. \square

Lemma 3.2 (Carleson estimate) *Let u, \bar{E} be as in Theorem 1.2. Then,*

$$\|u\|_{L^\infty(C_{1/2})} \leq C u(\bar{E}),$$

with C universal.

Proof We apply an iterative argument similar to the one of Lemma 3.1 above.

Assume $u(E) = 1$, and denote by h_Γ the distance in the e_n direction between a point $(x, t) \in C_1$ and Γ

$$h_\Gamma(x, t) := x_n - g(x', t).$$

Any point $(x, t) \in C_{2/3}$ can be connected to \bar{E} by a chain of adjacent forward-in-time cubes included in C_r , so that the size of each cube is proportional to the distance from its center to Γ raised to the power $1/\alpha$. The Hölder continuity of g implies that the number of cubes in this chain is proportional to $(h_\Gamma(x, t))^{1-1/\alpha}$, and by Harnack inequality we find

$$u \leq e^{C_1 h_\Gamma^{1-1/\alpha}} \quad \text{in } C_{2/3}, \quad (3.6)$$

with C_1 universal.

With the same notation as in Lemma 3.1, we wish to prove that if $r \leq r_0$ and

$$u(y, t_y) \geq f(r),$$

for some $(y, t_y) \in C_{1/2}$, then we can find $(z, t_z) \in S$,

$$S := \{(x, t) \mid x' - y' \in (-r, r)^{n-1}, \quad t \in (t_y - r, t_y], \quad 0 < h_\Gamma(x, t) < r^\beta\},$$

such that

$$u(z, t_z) \geq f\left(\frac{r}{2}\right).$$

Since $|(z, t_z) - (y, t_y)| \leq Cr^\alpha$, we see that for r small enough, we can build a convergent sequence of points $(y_k, t_k) \in C_{2/3}$ with $u(y_k, t_k) \geq f(2^{-k}r) \rightarrow \infty$. This is a contradiction if we assume that u vanishes continuously on Γ , and is therefore bounded. If $u = 0$ on Γ is understood in the sense that u is the limit of an increasing sequence of continuous subsolutions which vanish on Γ , then we may apply the argument below to one such subsolution and reach again a contradiction.

To show the existence of the point z , assume for simplicity $y' = 0$, $t_y = 0$, and then $S = C_r^- \setminus \mathcal{A}_r$. Let

$$w := \left(u - \frac{1}{2}e^{C_0 r^\gamma}\right)^+, \quad \text{with } C_0 \gg C_1.$$

By (3.6) we know that

$$w = 0 \quad \text{in } \mathcal{A}_r.$$

If our claim is not satisfied, then we apply Weak Harnack inequality for w in cubes of size $2r^\beta$ repeatedly as in Lemma 3.1. As we move a distance r inside the domain we obtain

$$w \leq f\left(\frac{r}{2}\right) e^{-c_0 r^{1-\beta}} \quad \text{in } C_{r/2}^-. \quad (3.7)$$

In particular

$$\frac{1}{2}f(r) \leq w(y, t) \leq f\left(\frac{r}{2}\right)e^{-c_0 r^{1-\beta}},$$

and we reach a contradiction if r_0 is sufficiently small as long as $1 - \beta < \gamma$ (which is possible because $\alpha > 1/2$). \square

4 Proof of Theorem 1.3

In this section, we assume that (3.1) holds for some $\alpha > 0$ possibly small, and in addition g satisfies a one-sided $C^{1/2}$ bound in the t variable, i.e.,

$$g(x', t + s) - g'(x', t) \geq -Ks^{1/2}, \quad \text{if } s \geq 0. \quad (4.1)$$

We will improve the estimates (3.5), (3.7) of the previous section by applying weak Harnack inequality in parabolic cubes of smaller size $\bar{r} \sim r^{\beta/\alpha}$ (which is the size chosen in the first step to obtain (3.4)) instead of r^β . Then the oscillation of w^- (or w) will decay by a factor $e^{-cr^{1-\beta/\alpha}}$ as we go from C_r^- to $C_{r/2}^-$. However, in cubes of size \bar{r} we can no longer guarantee the uniform measure estimate of the set where $w^- = 0$. To deal with this, we introduce a notion of parabolic capacity for the heat equation. This allows us to diminish the oscillation of w^- more precisely than in the measure estimate of Theorem 2.4.

Definition 4.1 Let E be a closed set. Set,

$$\text{cap}_{Q_1}(E) := \varphi(0, 1)$$

where φ is the solution to the heat equation in $Q_2(0, 1) \setminus (E \cap \overline{Q_1})$ which equals 0 on the parabolic boundary of $Q_2(0, 1)$ and it is equal to 1 in $E \cap \overline{Q_1}$.

The function φ is well defined by the Perron–Wiener–Brelot–Bauer theory (see for example [14]). Similarly, we can define $\text{cap}_{Q_r(x,t)}(E)$ by translating the cube at the origin, and then performing a parabolic rescaling

$$\text{cap}_{Q_r(x,t)}(E) := \text{cap}_{Q_1}(\tilde{E}), \quad \tilde{E} := \{(y, s) | (x + r^2 y, t + rs) \in E\}.$$

We prove here two lemmas about weak Harnack inequality depending on the size of the capacity of E in Q_1 . The first lemma states that a solution to the heat equation in $Q_1 \setminus E$ satisfies the Harnack inequality in measure if E has small capacity (Fig. 4).

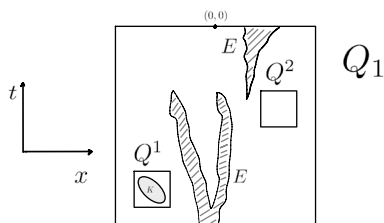
Lemma 4.2 Assume $v \geq 0$ is defined in $Q_1 \setminus E$ and satisfies

$$v_t = \Delta v.$$

Let

$$Q^i := Q_{1/4}(x_i, t_i) \subset Q_1, \quad i = 1, 2$$

Fig. 4 Lemma 4.2



be two cubes of size $1/4$ included in Q_1 , with $t_2 - t_1 \geq 1/4$. Assume that

$$\text{cap}_{Q_1}(E) \leq \delta \quad \text{and} \quad \frac{|\{v \geq 1\} \cap Q^1|}{|Q^1|} \geq 1/2,$$

for some δ small universal. Then

$$\frac{|\{v \geq c_0\} \cap Q^2|}{|Q^2|} \geq 1/2$$

for some c_0 small universal.

Proof Let h be the solution to the heat equation in $Q_1 \setminus K$ with $h = 0$ on the parabolic boundary of Q_1 , and $h = 1$ on $K := \{v \geq 1\} \cap Q^1$. We claim that

$$v \geq h - \varphi \quad \text{in} \quad Q_1 \setminus E,$$

where φ is the function from Definition 4.1. Since both v with $h - \varphi$ solve the heat equation in $Q_1 \setminus (K \cup E)$, it suffices to check the claim on the parabolic boundary of $Q_1 \setminus E$ and on K .

Indeed, $v \geq 0 \geq h - \varphi$ on $\partial_p Q_1$, and $v \geq 1 \geq h - \varphi$ on K . Moreover, $h \leq 1 \leq \varphi$ on E gives $h - \varphi \leq 0$ on E , and since $v \geq 0$ the claim is proved.

The conclusion follows from the inequality above, since by the Weak Harnack inequality, there exists c_0 small universal such that $h \geq 2c_0$ in Q^2 . On the other hand, $\varphi(0, 1) = \text{cap}_{Q_1} E \leq \delta$ implies that $\varphi \leq c_0$ in half the measure of Q^2 provided that δ is chosen sufficiently small. \square

Remark 4.3 We may use cubes Q^i of size σ and with $t_2 - t_1 \geq \sigma^2$, as long as δ and c_0 are allowed to depend on σ as well.

The second lemma states that the weak Harnack inequality holds for a subsolution $v \geq 0$ which vanishes on a set E of positive capacity. It follows directly from the definition of $\text{cap}_{Q_1}(E)$.

Lemma 4.4 Assume that $v \geq 0$ in $Q_2(0, 1)$, and

$$\Delta v \geq v_i \text{ in } Q_2(0, 1), \text{ and } v = 0 \text{ in } E \cap \bar{Q}_1.$$

If for some $\delta > 0$

$$\text{cap}_{Q_1}(E) \geq \delta,$$

then

$$v(x, t) \leq (1 - c(\delta))\|v\|_{L^\infty}, \quad (x, t) \in Q_{1/2}(0, 1).$$

Proof Assume $\|v\|_{L^\infty} = 1$. We compare $1 - v$ with φ in $Q_2(0, 1) \setminus (E \cap \overline{Q_1})$ and find $1 - v \geq \varphi$. On the other hand since $\varphi(0, 1) \geq \delta$ and $\varphi = 0$ on the lateral boundary of $Q_2(1, 0) \times [0, 1]$ it follows that φ satisfies the forward Harnack inequality, and $\varphi \geq c\delta$ in $Q_{1/2}(0, 1)$. The same inequality holds for $1 - v$ which gives the desired estimate. \square

Remark 4.5 We may write the conclusion in $[-1/2, 1/2]^n \times [\sigma, 1]$ for any $\sigma > 0$ provided that the constant $c = c(\delta, \sigma)$ depends on σ as well.

We are now ready to provide the proof of Theorem 1.3.

Proof of Theorem 1.3 We only show that the exponent in the estimate (3.5) from the previous section can be improved to

$$w^- \leq e^{-cr^{1-\frac{\beta}{\alpha}}}, \quad (4.2)$$

by the use of the two lemmas above. The rest of the proof remains the same as before. Notice that now $1 - \frac{\beta}{\alpha} < \gamma$ holds simply by choosing $\beta > 1$ and no restriction on range of the Hölder exponent $\alpha > 0$ is needed.

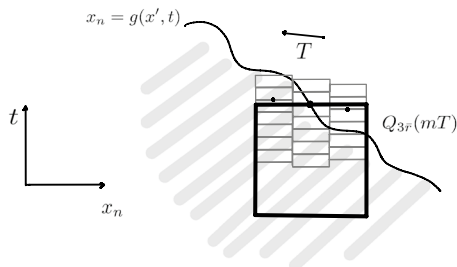
The same argument improves the exponent in (3.7) from $1 - \beta$ to $1 - \beta/\alpha$ in the proof of the Carleson estimate.

We proceed with the proof of (4.2). We set $\bar{r} := r^{\beta/\alpha}$, and by hypothesis, the translation by the vector

$$T := (-\bar{r}e_n, \kappa \bar{r}^2) \in \mathbb{R}^{n+1}$$

maps the complement of \mathcal{C}_1 into itself, provided that $\kappa \leq 1/2$ is small depending on the constant K in (4.1). Thus if we take a cube and then translate it by T , the complement of \mathcal{C}_1 (where $w^- = 0$) “increased” in the translating cube because of (4.1) (see Fig. 5).

Fig. 5 Proof of Theorem 1.3



Decompose the space \mathbb{R}^{n+1} into cubes of size $Q_{\bar{r}}$ in the following way. Take $Q_{\bar{r}}$ centered at the origin and then translate it by a linear combination of the vectors $\bar{r}e_i$, $i < n$, T , and $\bar{r}^2 e_{n+1}$ using integer coefficients. We look at the behavior of w on arrays of cubes translated by multiples of T . Starting with $Q_{\bar{r}}(0)$, we consider $Q_{\bar{r}}(mT)$, with $m \in \mathbb{Z}$. When $m \geq C\bar{r}^{\alpha-1}$, $Q_{3\bar{r}}(mT) \subset \mathcal{A}_r$, and when $m \leq -C\bar{r}^{\alpha-1}$, $Q_{\bar{r}}(mT) \subset E$, where E denotes the complement of \mathcal{C}_1 . Thus, there is an intermediate m_0 where

$$\text{cap}_{Q_{3\bar{r}}(mT)}(E) < \delta \quad \text{if and only if} \quad m \geq m_0.$$

When we decrease m from $C\bar{r}^{\alpha-1}$ to m_0 we may apply Lemma 4.2 in each such $Q_{3\bar{r}}(mT)$. The weak Harnack inequality holds in measure in these cubes (see Remark 4.3, with $\sigma^2 = \kappa/20$), and as in (3.4), (as there are at most $C\bar{r}^{\alpha-1}$ such cubes) we find that

$$\{w + 1 \geq f(r)e^{-C\bar{r}^{\alpha-1}} \geq 2\}$$

in a fixed proportion of each such $Q_{3\bar{r}}(mT)$ with $m \geq m_0$. Thus $w^- = 0$ in a fixed proportion of $Q_{3\bar{r}}(mT)$, and by the weak Harnack inequality

$$\|w^-\|_{L^\infty(Q_{\bar{r}}(x,t))} \leq (1-c)\|w^-\|_{L^\infty(Q_{6\bar{r}}(x,t))} \quad (x,t) = mT + 2\bar{r}^2 e_{n+1}, \quad (4.3)$$

if $m \geq m_0$.

If $m < m_0$ then the capacity of E in $Q_{3\bar{r}}(mT)$ is more than δ . By Lemma 4.4, the inequality above remains valid after possibly relabeling c . We conclude that (4.3) is valid for all cubes centered at $mT + 2\bar{r}^2 e_{n+1}$, and in particular for $Q_{\bar{r}}(2\bar{r}^2 e_{n+1})$.

This argument shows that (4.3) holds in fact at all points $(x,t) \in \mathcal{C}_{3r/4}^-$. Indeed,

$$\text{if } |x_n - g(x', t)| > Cr^\beta \text{ then either } Q_{\bar{r}}(x, t) \subset \mathcal{A}_r \text{ or } Q_{\bar{r}}(x, t) \subset E$$

and (4.3) is satisfied trivially as $w^- = 0$ in $Q_{\bar{r}}(x, t)$. Otherwise, we argue as above by decomposing the space starting with the cube centered at $(x, t) - 2\bar{r}^2 e_{n+1}$ instead of the origin. Notice that $(x, t) \in \mathcal{C}_{3r/4}^-$ and $|x_n - g(x', t)| \leq Cr^\beta$ imply that

$$Q_{3\bar{r}}((x, t) - mT) \subset \mathcal{A}_r \text{ and } Q_{3\bar{r}}((x, t) + mT) \subset E \text{ when } m \sim C\bar{r}^{\alpha-1}$$

and the argument applies as before.

In conclusion, the maximum of w^- is decaying a fixed proportion each time we remove the cubes $Q_{ss6\bar{r}}$ which are tangent to the parabolic boundary of the infinite cylinder in the (x', t) variables

$$\{|x_i| \leq 3/4r, \quad i < n\} \cap \{t \in [-3/4r, 0]\}.$$

Thus

$$w^- \leq e^{-cr\bar{r}^{-1}} \quad \text{in } \mathcal{C}_{r/2}^-$$

as desired, and (4.2) is proved. \square

Declarations

Conflict of interest The authors have no conflicts of interest to declare that are relevant to the content of this article.

References

1. Ancona, A.: Principe de Harnack a la frontiere et theoreme de Fatou pour un operateur elliptique dans un domaine lipschitzien. *Ann. Inst. Fourier* **28**, 169–213 (1978)
2. Banuelos, R., Bass, R.F., Burdzy, K.: Hölder domains and the boundary Harnack principle. *Duke Math. J.* **64**, 195–200 (1991)
3. Bass, R.F., Burdzy, K.: A boundary Harnack principle in twisted Hölder domains. *Ann. Math.* **134**, 253–276 (1991)
4. Bass, R.F., Burdzy, K.: Lifetimes of conditioned diffusions. *Probab. Theory Relat. Fields* **91**, 405–443 (1992)
5. Bass, R.F., Burdzy, K.: The boundary Harnack principle for non-divergence form elliptic operators. *J. Lond. Math. Soc.* **50**(1), 157–169 (1994)
6. Caffarelli, L., Fabes, E., Mortola, S., Salsa, S.: Boundary behavior of non-negative solutions of elliptic operators in divergence form. *Indiana Math. J.* **30**, 621–640 (1981)
7. Dahlberg, B.: On estimates of harmonic measure. *Arch. Rational Mech. Anal.* **65**, 272–288 (1977)
8. De Silva, D., Savin, O.: A short proof of Boundary Harnack inequality. *J. Differ. Equ.* **3**(15), 2419–2429 (2020)
9. De Silva D., Savin O.: On boundary Harnack inequality in Hölder domains. *Math. Eng.* **4** (1), 1–12 (2022)
10. Fabes, E.B., Garofalo, N., Marin-Malave, S., Salsa, S.: Fatou theorems for some nonlinear elliptic equations. *Rev. Mat. Iberoam.* **4**, 227–252 (1988)
11. Fabes, E.B., Garofalo, N., Salsa, S.: A backward Harnack inequality and Fatou theorem for non-negative solutions of parabolic equations. *Ill. J. Math.* **30**(4), 536–565 (1986)
12. Fabes, E.B., Safonov, M.V., Yuan, Y.: Behavior near the boundary of positive solutions of second order parabolic equations II. *Trans. AMS* **351**(12), 4947–4961 (1999)
13. Ferrari, F.: On boundary behavior of harmonic functions in Hölder domains. *J. Fourier Anal. Appl.* **4**(4–5), 447–461 (1988)
14. Friedman, A.: Parabolic equations of the second order. *Trans. Am. Math. Soc.* **93**, 509–530 (1959)
15. Garofalo, N.: Second order parabolic equations in nonvariational form: boundary Harnack principle and comparison theorems for nonnegative solutions. *Ann. Mat. Pura Appl.* **138**, 267–296 (1959)
16. Hofmann, S., Lewis, J.L., Nystrom, K.: Caloric measure in parabolic flat domains. *Duke Math. J.* **122**(2), 281–346 (2004)
17. Jerison, D.S., Kenig, C.E.: Boundary behavior of harmonic functions in non-tangentially accessible domains. *Adv. Math.* **46**, 80–147 (1982)
18. Kemper, J.T.: A boundary Harnack principle for Lipschitz domains and the principle of positive singularities. *Comm. Pure Appl. Math.* **25**, 247–255 (1972)
19. Kemper, J.T.: Temperatures in several variables: kernel functions, representations, and parabolic boundary values. *Trans. Am. Math. Soc.* **167**, 243–262 (1972)
20. Kim, H., Safonov, M.V.: Boundary Harnack principle for second order elliptic equations with unbounded drift Problems in mathematical analysis, No. 61. *J. Math. Sci.* **179**(1), 127–143 (2011)
21. Moser, J.: A Harnack inequality for parabolic differential equations. *Commun. Pure App. Math.* **17**(1), 101–134 (1964)
22. Petrosyan, A., Shi, W.: Parabolic boundary Harnack principles in domains with thin Lipschitz complement. *Anal. PDE* **7**(6), 1421–1463 (2014)
23. Safonov, M.V.: Non-divergence Elliptic Equations of Second Order with Unbounded Drift. *Nonlinear Partial Differential Equations and Related Topics. American Mathematical Society Translational Series 2*, 229, *Advanced Mathematical Science*, 64, pp. 211–232. American Mathematical Society, Providence, RI (2010)
24. Salsa, S.: Some properties of nonnegative solutions of parabolic differential operators. *Ann. Mater. Pura Appl.* **4**(128), 193–206 (1981)

25. Wang, L.: On the regularity theory of fully nonlinear parabolic equations: I. *Commun. Pure Appl. Math.* **45**(1), 27–76 (1992)
26. Wu, J.-M.G.: Comparison of kernel functions, boundary Harnack principle, and relative Fatou theorem on Lipschitz domains. *Ann. Inst. Fourier Grenoble* **28**, 147–167 (1978)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.