Simultaneous Suspension Control and Energy Harvesting Through Novel Design and Control of a New Nonlinear Energy Harvesting Shock Absorber

Mohammad R. Hajidavalloo , Joel Cosner, Zhaojian Li , Senior Member, IEEE, Wei-Che Tai, and Ziyou Song

Abstract—Simultaneous vibration control and energy harvesting of vehicle suspensions have attracted significant research attention over the past decades. However, existing energy harvesting shock absorbers (EHSAs) compromise suspension performance for high-efficiency energy harvesting and being only responsive to narrow-bandwidth vibrations. In this paper, we propose a new ballscrew-based EHSA design – inerter pendulum vibration absorber (IPVA) – that integrates an electromagnetic rotary EHSA with a nonlinear pendulum vibration absorber. We show that this design simultaneously improves ride comfort and energy harvesting efficiency by exploiting the nonlinear effects of pendulum inertia. To further improve the performance, we develop a novel stochastic linearization model predictive control (SL-MPC) approach in which we employ stochastic linearization to approximate the nonlinear dynamics of EHSA that has superior accuracy compared to standard linearization. In particular, we develop a new stochastic linearization method with guaranteed stabilizability, which is a prerequisite for control designs. This leads to an MPC problem that is much more computationally efficient than the nonlinear MPC counterpart with no major performance degradation. Also, the effect of different road preview configurations on control performance is investigated, which is shown to have a significant impact on the control performance. Extensive simulations are performed to show the superiority of the proposed new nonlinear EHSA and to demonstrate the efficacy of the proposed SL-MPC.

Index Terms—Energy harvesting shock absorber, model predictive control, stochastic linearization.

I. INTRODUCTION

RADITIONAL vehicle suspensions use hydraulic dampers to dissipate undesired vibration energy into heat waste, thereby improving ride comfort. There exists a great potential for harvesting this wasted energy; it is estimated that somewhere between 100 W to 10 kW of power per vehicle can be harvested for an average trip [1], [2]. The growth in hybrid and electric vehicles have further increased the potential impact through smart utilization and management of this harvested energy [3]. Therefore, extensive and increasing research efforts over the past three

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decades have focused on developing energy harvesting shock absorbers (EHSAs) – devices that convert vehicle suspension vibrations into useful electricity. Traditionally, EHSAs use electromagnetic dampers or magnetorheological dampers to recycle this heat waste into useful electricity. Rotary electromagnetic dampers have become popular because of their high conversion efficiency and quick responsiveness [1]. Thanks to their quick responsiveness, electromagnetic dampers are integrated with power electronic circuits to perform damping force control in real time [4], [5] or used as actuators to deliver active force to improve road handling and ride comfort [6].

To drive the rotary electromagnetic (EM) damper, motion conversion mechanisms, such as rack-pinion [7] and ballscrew [8], are required to convert the linear suspension vibrations into angular motion. This operation principle, although being straightforward, has a critical drawback; that is, they require large suspension vibrations to achieve high-efficiency energy harvesting, thereby compromising suspension performance for energy recovery. This drawback has been widely recognized in the literature. Through numerical simulations on a quarter car, Abdelkareem et al. [9] concluded that ride comfort and the harvestable power cannot be optimized at the same time. Through numerical simulations on a quarter car with a traditional EHSA, Casavola et al. [10] showed that a trade-off always exists between road handling and the energy harvesting performance. Huang et al. [11] considered a traditional ball-screw-based EHSA in a quarter car and discovered that ride comfort and the harvested power are conflicting objectives. Guo et al. [12] and Li and Zuo [13] considered a traditional rack-pinion-based EHSA in a quarter car and showed that ride comfort and the harvested power cannot be optimized together.

Therefore, there is a pressing need for fundamentally new EHSA designs to fully realize the potential benefits of harvesting vehicle vibration energy while simultaneously achieving great suspension performance. Recently, Gupta and Tai proposed a nonlinear rack-pinion-based EM damper, known as inerter pendulum vibration absorber (IPVA) [14], [15]. The IPVA consists of a planetary gear set that integrates a rack-pinion-based EM damper and a nonlinear pendulum vibration absorber. It was shown that the nonlinear inertial effects of the pendulum increased the harvested power and energy harvesting bandwidth when subject to harmonic excitation. Although showing promising results, the rack-pinion mechanism is too bulky to fit in a typical vehicle suspension system. Later, Cosner and

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Tai [16] proposed to use a ball-screw in lieu of rack-pinion, which is more compact and suitable for vehicle suspension systems. However, their design was not able to integrate with a rotary EM damper for electricity generation, and only capable of vibration suppression. Nevertheless, they showed that the pendulum absorbed the vibration energy of a suspended platform when subject to white noise excitation, thereby holding promise to achieve energy harvesting and vibration suppression at the same time. In this paper, we propose a new design that integrates the pendulum vibration absorber and a ball-screw-based EM damper. This new ball-screw-based IPVA is integrated with a quarter-car suspension model where we optimize the design parameters to achieve a better tradeoff between the suspension performance (ride comfort) and energy harvesting efficiency of the system when subject to stochastic road excitation.

To further improve the energy harvesting efficiency while maintaining good ride comfort, model predictive control (MPC) is applied to the IPVA-integrated quarter-car suspension model. MPC is a popular tool for solving constrained optimal control problems with the advantage of online implementation relative to methods such as dynamic programming [17]-[20]. Since the IPVA dynamics and the objective function (as defined in the sequel) are nonlinear, a nonlinear MPC (NMPC) can be exploited to solve this problem. However, the NMPC problem is computationally expensive and is difficult for real-time implementations, especially for suspension systems that require very high control frequency. To address this issue, we propose a new MPC framework by exploiting an approximated linear dynamics using the technique of stochastic linearization [21], [22]. The new MPC framework with the stochastically linearized dynamics has comparable control performance with NMPC while requiring significantly less computation power. Moreover, the MPC performance is also investigated with and without the usage of perfect road preview, which can be obtained through recent road information estimation techniques using a single or multiple of vehicles [23], [24]. We show that the usage of road profile preview can greatly improve the performance.

The contributions of this paper include the following. First, we integrate IPVA into automotive suspension systems and optimize the system designs that offer improved ride comfort and energy harvesting efficiency at the same time when compared to the traditional EHSA. Second, we develop a novel stochastic-linearization MPC (SL-MPC) framework by exploiting a stochastically-linearized dynamics based on the nonlinear equation of motion (EOM) of the IPVA. To the best of the authors' knowledge, this is the first time that stochastic linearization is exploited in MPC designs to deal with nonlinear dynamical systems excited by random disturbance signals. Third, we investigate the usage of online estimated road information into the prediction horizon, which we show is able to enhance the performance. Last but not least, extensive simulations are performed to demonstrate the efficacy of the proposed framework. It should be noted that although Chen et al. [2] proposed a nonlinear EHSA that also utilized inertial nonlinearity, their nonlinear EHSA was only efficacious around a resonance peak; that is, it is narrow-banded. Furthermore, their numerical study showed that the maximum energy harvesting efficiency and

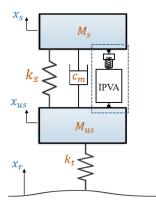


Fig. 1. Schematics of an IPVA-integrated quarter-car model.

worst ride comfort occurred at the same frequency. In other words, energy harvesting and suspension performance are still conflicting objectives in their design.

The rest of this paper is organized as follows. In Section II, the design and modeling of the IPVA-integrated suspension system are introduced. Section III describes the NMPC and SL-MPC designs. Simulations and performance evaluations are presented in Section IV while Section V concludes the paper.

II. ENERGY HARVESTING SYSTEM DESCRIPTION

In this section, we introduce a novel nonlinear EHSA design, along with a linear benchmark model (Section II-A). The governing equations of motion (EOM) for both systems are derived (Section II-B). The optimal parameter designs for both systems are also discussed (Section II-C).

A. System Structure

The IPVA-integrated quarter-car suspension model is shown in Fig. 1, where the mass of a quarter-car body and the unsprung mass (i.e., wheel axle) of a quarter car are represented by M_s and M_{us} , respectively. The interaction between the unsprung mass M_{us} and the ground is characterized by a spring of stiffness k_t (tire stiffness) with its displacement from the equilibrium given by x_{us} . The quarter-car model is excited by the road profile/disturbance x_r . The sprung mass and unsprung mass are connected with a spring of stiffness k_s (suspension stiffness), a viscous damper with damping coefficient c_m , and the IPVA. Note that the mechanical damping c_m is introduced to account for mechanical energy loss due to the ball-screw and generator gears. As shown in Fig. 2(a), the IPVA consists of a ball-screw system with a lead L connected between M_{us} and M_s . It converts the linear oscillations between M_{us} and M_s to the rotation of the screw. The carrier fixed to the screw houses a pendulum of mass m and radius r at a distance of R_p from the screw's axis of rotation. A sun gear is free to rotate with respect to the carrier about the same axis of rotation as the screw and drives the generator. A planet gear fixed to the pendulum rotates and revolves on the sun gear. The housing of the generator is fixed to the sprung mass. Considering the gear ratio g_p between the sun and planet gear to be 1, we have $\psi = \vartheta - \varphi$ as shown in Fig. 2(b), ψ , ϑ , and φ are the angular displacement of the gen-erator rotor, screw, and pendulum, respectively. The generator

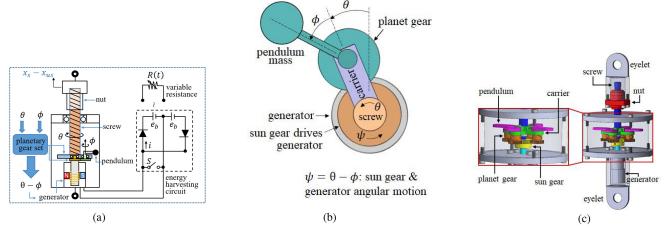


Fig. 2. Schematic for the IPVA. (a) Conceptual design of semi-active inerter pendulum vibration absorber. (b) Top view of pendulum and gear arrangement. (c) Realization of a possible prototype.

is connected with an electricity storage e_b via a pulse width modulated (PWM) step-up chopper to modulate the duty cycle of the generator by switch S, equivalent to a variable resistor R(t) [25]. The corresponding electrical damping coefficient in the generator is denoted by $c_e(t) = M \kappa_t \kappa_e / R(t)$ [26], where *M* is a motion transmission factor that is related to the ball screw lead and gear ratio of the generator, and κ_t and κ_e are the torque and voltage constant of the generator, respectively. Note that the generator's inductance is neglected because the impedance of the inductance is small compared with the resistance considering that the vibration induced by road irregularities is usually in the frequency range of 1–10 Hz [4], [7]. This model is a three degrees of freedom (DOF) system, with the degrees of freedom being the angular displacements of the pendulum (φ) and the screw (ϑ), and the displacement of the unsprung mass x_{us} . Note that the sprung mass displacement x_s is related to ϑ and x_{us} via $x_s - x_{us} = R\vartheta$, where $R = L/2\pi$. Finally, a 3D model for a potential prototype for the proposed IPVA is shown in Fig. 2(c). While the prototype shown consists of four pendulums, the pendulums move synchronously due to the planetary gear system. Hence this prototype is equivalent to a single pendulum system with a quadrupled pendulum mass *m*.

B. Equations of Motion of the EHSA

We next use the Lagrangian method to derive the EOM for the IPVA-integrated quarter-car system shown in Fig. 1. The kinematic relation $x_s - x_{us} = R\vartheta$ relates the suspension deflection $x_s - x_{us}$ to the angular displacement ϑ with $R = L/2\pi$, from which it follows that the total kinetic energy of the system is:

$$T = T_{M_{us}} + T_{M} + T_{c} + T_{p} + T_{r}$$

$$= \frac{1}{2} M_{us} (\dot{x}_{us})^{2} + \frac{1}{2} M_{s} R \vartheta + \dot{x}_{us}^{2} + \frac{1}{2} J \frac{\vartheta^{2}}{2} + \frac{1}{2} m R_{p}^{2} \dot{\vartheta}^{2} + r^{2} \dot{\vartheta} + \dot{\varphi}^{2} + 2 R_{p} r \cos(\varphi) \dot{\vartheta} \dot{\vartheta} + \dot{\varphi}$$

$$+ \frac{1}{2} J_{p} \dot{\vartheta} + \dot{\varphi}^{2} + \frac{1}{2} J_{r} (\dot{\varphi} - \dot{\vartheta})^{2}, \qquad (1)$$

where R_p and r are the distance between the pendulum pivot point and center of the carrier (i.e., half carrier length), and the length of the pendulum, respectively; the parameters J, J_r , and J_p represent the principal moment of inertia (w.r.t. primary rotational axis) of the carrier, generator rotor, and pendulum, respectively. Note that the moment of inertia of the gears and screw are assumed to be negligible. The deformation of the springs contribute to the potential energy, which can be obtained as:

$$V = \frac{1}{2}k_sR^2\vartheta^2 + \frac{1}{2}k_t(x_{us} - x_r)^2 + \frac{1}{2}k_p\varphi^2.$$
 (2)

The last term is related to the torsion spring attached to the pendulum which has a very low stiffness for making sure the stochastic linearized model (discussed in Section III-B) is stable. We use the concept of virtual work to include the nonconservative forces in the system. Since the virtual angular displacement of the rotor is $\delta(\psi) = \delta(\vartheta - \varphi)$, the virtual work on the rotor by the damping torque due to energy harvesting is $\delta W_r = -c_e(\vartheta - \varphi) \cdot \delta(\vartheta - \varphi)$. Here δ represents the variational operator. For the mechanical damper, the virtual displacement of the mass is $R \cdot \delta \vartheta$, so the virtual work is given by $\delta W_m = -c_m R^2 \vartheta \cdot \delta \vartheta$. Therefore, the total virtual work due to non-conservative forces is:

$$\delta W_{nc} = \delta W_r + \delta W_m = -c_e(\dot{\varphi} - \dot{\vartheta}) \cdot \delta \varphi$$
$$+ c_e(\dot{\varphi} - \dot{\vartheta}) \cdot \delta \vartheta - c_m R^2 \dot{\vartheta} \cdot \delta \vartheta. \tag{3}$$

As a result, the viscous damping forces included in the EOM for ϑ and φ are:

$$\vartheta: -c_m R^2 \dot{\vartheta} + c_e (\dot{\varphi} - \dot{\vartheta}),$$

$$\varphi: -c_e (\dot{\varphi} - \dot{\vartheta}).$$

The Lagrange equations are derived as,

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = Q_i, \tag{4}$$

where L = T - V is the Lagrangian, $q_1 = \vartheta$, $q_2 = \varphi$ and $q_3 = x_{us}$ are the degrees of freedom and Q_i is the generalized force

for degree of freedom q_i . After substituting the corresponding terms into (4), the EOM is obtained as:

$$G_{22}\ddot{\vartheta} + G_{24}\ddot{\varphi} + c_m R^2 \dot{\vartheta} - c_e (\dot{\varphi} - \dot{\vartheta}) + kR^2 \vartheta$$

$$- 2mR_p r \dot{\varphi} \dot{\vartheta} \sin(\varphi) - mR_p r \sin(\varphi) \dot{\varphi}^2 + RM_s \ddot{x}_{us} = 0, \quad (5)$$

$$G_{44}\ddot{\varphi}+G_{42}\ddot{\vartheta}+c_e(\dot{\varphi}-\dot{\vartheta})+k_p\varphi+R_pr\sin(\varphi)\dot{\vartheta}^2=0, \quad (6)$$

$$(M_s + M_{us})\ddot{x}_{us} + M_s R \ddot{\vartheta} + k_t (x_{us} - x_r) = 0,$$
 (7)

where

$$G_{22} = M_s R^2 + J + mR_p^2 + mr^2 + 2mR_p r \cos(\varphi) + J_p + J_r,$$

$$G_{24} = G_{42} = mr^2 + mR_p r \cos(\varphi) + J_p - J_r,$$

$$G_{44} = mr^2 + J_p + J_r.$$

For simplicity, the pendulum mass is represented by a point mass and the moment of inertia of the carrier is assumed to be small, resulting in negligible J_p and J.

Note that for a traditional linear EM damper, the pendulum does not exist and consequently φ , $\dot{\varphi}$, and all nonlinear terms will vanish. The non-conservative force for ϑ remains and the EOM for the linear benchmark is thus:

$$M_{s}R^{2} + J_{r}\vartheta + (c_{m}R^{2} + c_{e})\vartheta + k_{s}R^{2}\vartheta + RM_{s}\ddot{x}_{us} = 0, (8)$$

$$(M_{us} + M_{s})\ddot{x}_{us} + M_{s}R\ddot{\vartheta} + k_{t}(x_{us} - x_{r}) = 0.$$
(9)

C. Optimal Design of IPVA and Linear Benchmark

The performance of the IPVA depends on appropriate choices of design parameters, including R_p , r, and c_e . Towards that end, we define the following dimensionless variables:

$$\eta = \frac{r}{R_p}, \mu_r = \frac{mR_p^2}{M_s R^2}, \xi_e = \frac{c_e}{2\omega_0 M_s R^2}.$$

The maximum of electrical damping coefficient c_e depends on the internal resistance of generator. Considering the maximum electrical damping coefficient in ref. [9], it is assumed that $c_e < 7.2 \, \text{N} \cdot \text{s} \cdot \text{m}$, resulting in $\xi_e < 1$. Furthermore, it is assumed that $\eta < 0.9$ such that the pendulum length is smaller than the carrier radius for compactness and $\mu_r < 0.2$ such that the pendulum mass m < 2.5 kg for a reasonable weight. Then a reasonable choice of design parameters should satisfy the following constraints:

$$0.5 < \eta < 0.9$$
, $0.05 < \mu_r < 0.2$, $\xi_e < 1$. (10)

Note that the suspension spring stiffness is considered as given and not included in the optimization. The main reason is that we envision to fit our EHSA into existing suspension systems with given spring stiffness, according to which one can determine an optimal design and retrofit it to the suspension. Huang *et al.* [11] also fixed the suspension spring stiffness when determining local optimal designs of a traditional ball-screw-based EHSA. To get the optimal values for design parameters, a constrained vector objective optimization problem is defined with the variables R_p ,

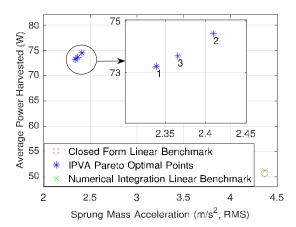


Fig. 3. Pareto optimal points for the proposed IPVA-integrated system (black asterisk) and the linear benchmark on Class-C road (red circle: closed form solution, green cross: numerical integration). $\omega_0 = 14.83$ rad/sec.

r and
$$c_e$$
, i.e.,

$$\min_{c_e, R_p, r} \frac{1}{N} \frac{X^N}{t=0} - c_e(\dot{\varphi}(t) - \dot{\vartheta}(t))^2, \quad X^N = \frac{1}{N} (\ddot{x}_s)^2$$
s.t. (5) - (7), (10).

Here the first element of the objective function vector represents the average harvested power and the second element represents the RMS value of the sprung mass acceleration, which is used to characterize the ride comfort.

The optimal design of the linear benchmark can be similarly obtained by solving the following optimization problem with respect to the variable c_e :

$$\min_{c_{e}} \cdot \frac{1}{N} \sum_{t=0}^{X^{V}} -c_{e} \dot{\vartheta}(t)^{2}, \quad \chi^{V} \quad r \quad \frac{1}{N} (\ddot{x}_{s})^{2}$$
s.t. (8) - (9), $\xi_{e} < 1$. (12)

Note that the optimization problems in (11) and (12) are computationally hard to solve directly. Alternatively, we utilize a grid search method where we discretize the optimization space into grid points and evaluate the performance of each point through Monte-Carlo simulations. More specifically, for each grid point that corresponds to a combination of the parameters to be optimized, we run N (N = 50) simulations for sufficiently long time, each of which is based on one random generated road profile corresponding to the ISO 8608 Class-C and Class-B road [26]. Figs. 3 and 4 summarizes the optimization results where only the Pareto optimal points are shown. Note that instead of combining the harvested energy and ride comfort metrics in a weighted sum, the adopted Monte Carlo method allows us to inspect the Pareto optimal designs to hand-pick the one that achieves the most desirable tradeoff.

Note in Fig. 3, the linear benchmark case, the Pareto front turns out to be a single point. Based on the Pareto front, we choose Point 3 as the "optimal" parameter set since we believe it represents the best trade-off between ride comfort and

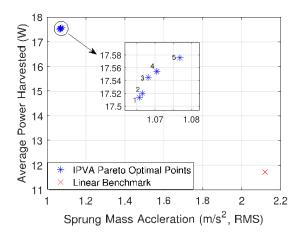


Fig. 4. Pareto optimal points for the proposed IPVA-integrated system (black asterisk) and the linear benchmark on Class-B road.

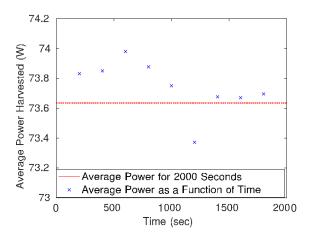


Fig. 5. Average power as a function of time for Pareto point 3 in Fig. 3 (black crosses) and average power over 2000 seconds (red dashed line). ω_0 = 14.83 rad/sec.

harvested energy. This choice corresponds to the following parameters: $R_p = 0.117 \, m$, $r = 0.0897 \, m$, $c_e = 0.225 \, N \cdot s \cdot m$. The optimal design for the linear benchmark corresponds to c_e = $0.225 \,\mathrm{N} \cdot \mathrm{s} \cdot \mathrm{m}$. It can be seen in Fig. 3 that the optimal IPVA system significantly outperforms the optimal linear benchmark design with about 45% increase in the harvested power and 45% reduction in the RMS value of sprung mass acceleration (better ride comfort). This clearly shows that our new nonlinear IPVA design can simultaneously achieve significantly better energy harvesting and ride comfort compared to conventional linear designs. The same arguments are true for the Class-B road analysis shown in Fig. 4 where we can choose any design point from 1 to 5 for the optimal operation of the IPVA, e.g., point 3 corresponds to design parameters $R_p = 0.132 \, m$, r = $0.1012 \, m, c_e = 0.225 \, N \cdot s \cdot m$ and the optimal design for the linear benchmark corresponds to $c_e = 0.225 \, \text{N} \cdot \text{s} \cdot \text{m}$.

To verify that the predictions of the average power associated with the IPVA-integrated system shown in Fig. 3 are weakly stationary, the average power as a function of time is plotted in Fig. 5 for the optimal parameters (Pareto point 3 in Fig. 3).

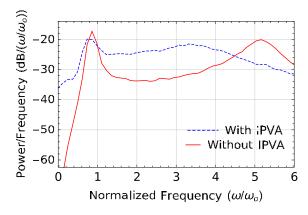


Fig. 6. Power spectral density for instantaneous power with normalized frequency ω/ω_0 ; Linear system: solid line, IPVA-system: dashed line.

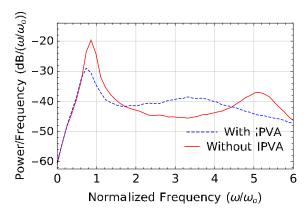


Fig. 7. Power spectral density for sprung mass acceleration with normalized frequency ω/ω_0 ; Linear system: solid line, IPVA-system: dashed line.

It is noted that after about 1,200 seconds of integration time, the average power remains within 0.2% of the average value for an integration time of 2,000 seconds, implying that it is approaching weak stationarity. In addition, Fig. 3 also shows the closed-form solution for RMS acceleration and average power harvested associated with the linear benchmark (see Appendix for the derivations). The solution given by numerical integration of (46) and (43) for 2000 seconds, averaged over 50 realizations, is shown in Fig. 3 as well. It is clear that the closed-form solution and the numerical integration solutions for the linear benchmark are very close, which further confirms that the solution given by numerical integration is very close to stationarity.

In order to further explain the performance improvement with the implementation of the IPVA, the power spectral densities (PSD's) associated with the sprung mass acceleration and instantaneous power for the IPVA-system and the linear system are numerically calculated and compared for the third Pareto point. This is shown in Figs. 6 and 7. Note that the PSD's are a function of normalized frequency ω/ω_0 , where $\omega_0 = \frac{P}{K_s/M_s}$. Furthermore, the natural frequencies associated with linear system were analytically calculated as $\omega_{n1} = .85\omega_0$ and $\omega_{n2} = 5.18\omega_0$, which naturally correspond to the frequencies associated with the PSD peaks of linear system; see Figs. 6 and 7.

Fig. 6 shows that the addition of the IPVA produces a superharmonic peak at about four times the first natural frequency associated with the linear system, while the second natural frequency is nearly destroyed. Specifically, between the first natural frequency ($\omega \approx 0.85\omega_0$) and super-harmonic frequency ($\omega \approx 4 \times 0.85\omega_0$), the power spectral density with the IPVA is significantly larger, leading to greater average power relative to the linear system. Moreover, the acceleration power spectral density for the IPVA-system shown in Fig. 7 displays the same super-harmonic characteristics, with negligible second natural frequency contribution, while the first peak is greatly diminished, and the super-harmonic peak is relatively small compared to second natural frequency contribution associated with the linear system. This naturally leads to lower overall acceleration for the sprung mass.

This section is concluded with a brief discussion on the limitations of the proposed IPVA. The IPVA employs a ball-screw design and has similar limitations with other ball-screw-based EM dampers. It is known that ball-screws have a relatively lower conversion efficiency and relatively higher cost than rack-pinions [1], and may have a risk of structural failure (buckling) [8]. Furthermore, the ball recirculating system and the planetary gear system may have reliability issues due to having complicated transmission mechanisms, which may limit the applicability to heavy-duty vehicles.

III. MODEL PREDICTIVE CONTROLLER DESIGNS

In this section, we present controller designs for the IPVA-integrated suspension system to further improve its performance. Specifically, we first exploit a nonlinear MPC approach based on the derived nonlinear dynamics (Section III-A). Furthermore, we develop a novel stochastic linearized MPC where we use a stochastically linearized model that can accurately approximate the nonlinear dynamics while leading to much more efficient computations (Section III-B).

A. Nonlinear MPC (NMPC) Design

Our NMPC design is based on the derived nonlinear dynamics in (5)–(7). By defining the states $x_1 = \vartheta$, $x_2 = \vartheta$, $x_3 = \varphi$, $x_4 = \dot{\varphi}$, $x_5 = x_{us}$, $x_6 = \dot{x}_{us}$, the control input $u = c_e$ (electrical damping coefficient), and the external disturbance $w = x_r$, the state-space model can be written as:

$$G(x)\dot{x} = F(x, u, w), \tag{13}$$

where

$$G(x) = \begin{bmatrix} ? & 1 & 0 & 0 & 0 & 0 & 0 \\ ? & 0 & G_{22} & 0 & G_{24} & 0 & M_s & R \\ ? & 0 & 0 & 1 & 0 & 0 & 0 \\ ? & 0 & G_{42} & 0 & G_{44} & 0 & 0 \\ ? & 0 & 0 & 0 & 0 & 1 & 0 \\ ? & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & M_s & R & 0 & 0 & 0 & M_s + M_{us} \end{bmatrix}$$

with G_{22} , G_{24} , G_{42} and G_{44} being introduced in (5)–(7); and $\begin{array}{c}
x_2 \\
F_2(x) \\
\hline
P_2 \\
\hline
P_3 \\
\hline
P_4 \\
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P_5 \\
\hline
P_6 \\
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P_7 \\
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P_8 \\
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P_8$

where

$$F_2(x) = -c_m R^2 x_2 + u(x_4 - x_2) - kR^2 x_1 + 2mR_p r x_4 x_2 \sin(x_3) + mR_p r \sin(x_3) x_4^2.$$
 (14)

It can be shown that G(x) is positive definite and therefore invertible. The objectives of the control designs are twofold: 1) Achieve good ride comfort by minimizing the sprung mass acceleration (i.e., $\ddot{x}_s = \ddot{x}_{us} + R\vartheta = \dot{x}_6 + R\dot{x}_2$); and 2) Real-ize efficient energy harvesting by maximizing the regenerated power (i.e., $P = c_e(\vartheta - \dot{\varphi})^2 = u(x_2 - x_4)^2$) by controlling the damping u in real time [2]. Hence, the continuous time economic stage cost function can be defined as

$$I(x, u) = \alpha_1 (\dot{x}_6 + R\dot{x}_2)^2 - \alpha_2 u (x_2 - x_4)^2, \tag{15}$$

where α_1 and α_2 are positive weighting factors to tradeoff the two objectives. As such, the NMPC problem can be formulated as:

$$\min_{U} J = \sum_{k=0}^{N-1} I_{d}(x(k), u(k))$$
s.t. $x(0) = x_{0}, u(k) \mathbb{P} \cup_{d},$

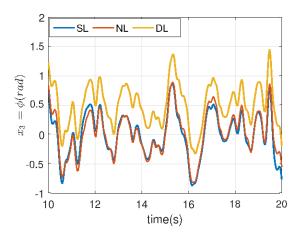
$$x(k+1) = G(x(k))^{-1} F(x(k), u(k), w(k))_{d}, \qquad (16)$$

where U = [u(0), ..., u(N - 1)] is the optimization variable, I_d is the discrete stage cost of I, N is the prediction horizon, and U_d represents the control constraints.

Note that this NMPC problem is computationally heavy and is difficult for onboard implementation due to fast dynamics of suspension systems. Therefore, we next present a sub-optimal MPC that is computationally efficient and thus more suitable for practical uses.

B. Stochastic Linearized MPC (SL-MPC)

As the obtained dynamics of the IPVA-integrated suspension system is inherently nonlinear, the NMPC formulation above is computationally expensive and difficult for online implementations. It would be meritorious if we can find a good linear approximation of the nonlinear dynamics to achieve efficient computations. One option is to linearize the nonlinear system around equilibrium points, referred to as conventional (or deterministic) linearization and it is shown that this approach does not work well for the considered nonlinear system as it produces large prediction error compared to the original nonlinear system (see e.g., Fig. 8). In this subsection, we propose a stochastic linearization approach where we derive a linear approximation such that its system response is statistically close to the nonlinear response when subject to external random excitation. This is



Accuracy of stochastic linearization (SL) approach relative to deterministic linearization (DL).

especially appealing as our EHSA system is indeed subject to random road disturbance. The stochastic linearization approach is a powerful and efficient tool for capturing the complex and random behavior of a nonlinear system [21]. The goal is to obtain a linear system such that its deviation from the original nonlinear system is small in the sense of expectations. To this end, we first write the EOM of the IPVA dynamics (5)-(7) in an alternative form as:

$$M_1\ddot{q} + C_1\dot{q} + K_1q + \Phi(q, \dot{q}, \ddot{q}) = Q(t),$$
 (17)

where we have

$$\Phi_1 = 2mR_p r \cos(\varphi) \ddot{\vartheta} + mR_p r \cos(\varphi) \ddot{\varphi} - 2mR_p r \dot{\varphi} \dot{\vartheta} \sin(\varphi)$$
$$- mR_p r \sin(\varphi) \dot{\varphi}^2,$$

$$\Phi_2 = mR_p r \cos(\varphi) \vartheta + mR_p r \sin(\varphi) \vartheta, \quad \Phi_3 = 0.$$
 (18)

The matrices M_I , C_I and K_I are the linear inertia, damping, and stiffness matrices, respectively; Φ is the collection of nonlinear terms in the equations; and Q(t) is the generalized force. The main idea in SL approach is to find equivalent deterministic inertia, damping, and stiffness matrices (denoted by Me, Ce and K_e, respectively) such that when replaced by the nonlinear terms, the system responses to the random disturbance signal are similar in the statistical sense, that is, the following equivalent stochastic linearized system

$$(M_I + M_e)\ddot{q} + (C_I + C_e)\dot{q} + (K_I + K_e)q = Q(t)$$
 (19)

is close to (17). More specifically, we seek Me, Ce and Ke by solving the following optimization problem:

$$\min_{\mathsf{M}_{\varepsilon}, \mathsf{C}_{\varepsilon}, \mathsf{K}_{\varepsilon}} E\{\varepsilon^{\mathsf{T}} \varepsilon\} \tag{20}$$

where

$$\varepsilon = \Phi(q, \dot{q}, \ddot{q}) - M_e \ddot{q} - C_e \dot{q} - K_e q \qquad (21)$$

is the *n*-dimensional vector difference between the actual nonlinear system and the stochastically linearized system [21], i.e., the difference between (17) and (19).

Assuming the road disturbance is Gaussian and by following the stochastic optimization procedures in [21], one can find the elements of Me, Ce and Ke as:

$$\mathsf{M}_{e,ij} = E \frac{\partial \Phi_i}{\partial \ddot{\mathsf{q}}_j} \,, \tag{22}$$

$$C_{e,ij} = E \qquad QQ_{i} \qquad (23)$$

$$K_{e,ij} = E - \frac{\partial \Phi_i}{\partial q_j}$$
, for $i, j = 1, 2, 3,$ (24)

where $q_1 = \vartheta$, $q_2 = \varphi$, $q_3 = x_{us}$, and Φ_1 , Φ_2 , and Φ_3 are introduced in (18). Solving (22)-(24) can be done analytically or by Monte-Carlo simulation to calculate the expectations. The former poses a challenge due to the complication of nonlinear terms whereas the latter could be done with desired accuracy with adequate number of simulations for any type of nonlinear terms and road disturbance signal. In this paper, we exploit the latter approach to obtain the M_e , C_e , and K_e terms.

The system responses of the stochastic linearized (SL) system and the conventional deterministic linearized (DL) system is compared with the original nonlinear (NL) system, where the results for x_3 on one sample road profile is shown in Fig. 8. It is clear that the stochastic linearization is a much closer representation of the nonlinear dynamics.

While the SL approach discussed above can generate a linear model that closely approximates the nonlinear dynamics, the optimization in (20) has no guarantees in controllability or stabilizability, which is a prerequisite for control designs. In case that the obtained SL system is not stabilizable, a constrained optimization problem can be formulated to guarantee stabilizability. The process is detailed in Appendix-B where the essential idea is to find stabilizable matrices that are close to the original SL matrices, e.g., in the sense of matrix norm.

With the stochastically linearized system, one can formulate the following SL-MPC problem:

$$\min_{U} \int_{k=0}^{k-1} I_{d}(x(k), u(k))$$

subject to
$$x(0) = x_0$$
, $u(k) \supseteq U_d$,

$$x(k+1) = A_s x(k) + B_s x(k) u(k) + D_s w(k),$$
 (25)

where $U = [u(0); \dots; u(k-1)]$, and A_s , B_s and D_s can be found by inspecting (19) with $x = [q, \dot{q}]^T$. Note that the system model in (25) has a bilinear term, $B_s x(k) u(k)$, where the state and the control variable are cross multiplied. We next follow the

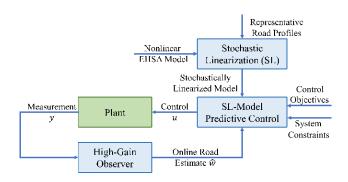


Fig. 9. Flowchart of SL-MPC.

reformulation approach in [27] to transform this bilinear model to a linear one. In specific, we substitute the damping force term, $u(x_4 - x_2)$ by F_d , which follows by changing the control variable from u to F_d and introducing passivity constraints to the model, i.e.,

$$x(k+1) = A_{l}x(k) + B_{l}F_{d}(k) + D_{l}w(k)$$
s.t. $-F_{d}(k)(x_{4}(k) - x_{2}(k)) \le 0$,
$$[F_{d}(k) - c_{\max}(x_{4}(k) - x_{2}(k))] \cdot F_{d}(k) \le 0, \quad (26)$$

where A_I , B_I , and D_I are the induced matrices after the new control formulation. This linear time invariant form will make the MPC problem much more computationally efficient for computations.

We would like to wrap up this section with a flow chart to summarize the work flow, modeling assembly and the control algorithm in SL-MPC. As shown in Fig. 9, our proposed framework starts with the *Stochastic Linearization* (SL) of the Nonlinear EHSA model given the *Representative Road Profile*. This stochastic linearized model is then used for the MPC task by considering the *Control Objectives* (i.e., the trade off between ride comfort, energy harvesting, and vehicle handling) and *System Constraints* (i.e., electrical damping maximum value). These steps form the SL-Model Predictive Control block in the flowchart. The obtained control command is then used for controlling the *Plant*. Lastly, the measurement from plant is used by the *High-Gain Observer* to have an estimate of the road profile which is used by the SL-MPC as the road preview information.

IV. SIMULATION RESULTS

In this section, extensive simulations are presented to investigate the performance of different control designs including the passive IPVA system, NMPC, and SLMPC. Furthermore, we evaluate the control designs in different road preview settings. The road preview refers to the availability of road profile signal data, from time step t to the end of the horizon, t + N - 1, i.e., (w(t), w(t+1), ..., w(t+N-1)), for the MPC task. Specifically, the MPC at each time step t solves a constrained optimization problem with a pre-defined objective and horizon N by predicting the trajectory of the system from time step t to t + N - 1 using the system dynamics (e.g., (16)) and the road profile information (w(t), w(t+1), ..., w(t+N-1)), if available. It is thus clear that a better estimate of the road profile

TABLE I PARAMETERS OF IPVA SYSTEM

Parameters	Physical meanings
M_s , M_{us} , m	mass of sprung, unsprung, and pendulum
L	ball-screw lead
$R = L/(2\pi)$	characteristic length of ball-screw
R_p	distance between the pendulum pivot point and center of the carrier
r	pendulum length
J, J_r, J_p	principal moment of inertia of carrier, generator rotor, and pendulum
k_s , k_l , k_p	stiffness of suspension, tire and pendulum torsion spring

TABLE II SIMULATION PARAMETERS

M_s	M_n	k_t, k_s	k_p	
$250 \ kg$	$35 \ kg$	$150,55 \ kN$	$0.852 \ N$	
J_r	J	J_h	J_p	
$1.21 \times 10^{-4} kg \cdot m^2$	0	0	0	
m	N	$[u_{min}, u_{max}]$	T_s	
$2.5 \ kg$	15	[0, 0.225]	$0.01 \ s$	

will lead to a more accurate prediction, and subsequently better control performance. In the first setting, we assume there is a complete road preview e.g., obtained from recent road information estimation techniques [23], [24]. In the second setting, we exploit an online road profile estimation algorithm, and use the last road disturbance estimation (LRDE) while keeping it constant across the prediction horizon. This is done using the HGO design introduced in Appendix-C. In the third setting, we use a noisy version of the preview with different signal to noise (SNR) ratios in the prediction horizon to evaluate the performance. Three different SNR values are considered which corresponds to road profiles with large noise (SNR 10), moderate noise (SNR 15), and small noise (SNR 20), see Fig. 11. The first setting can be used to evaluate the performance cap while the second and third are realistic settings that has practical implications. The road disturbance signals follow a Class-C (average) road, which are generated following the procedure outlined in [26] and the vehicle speed is considered as 60 mph. The system parameters for simulation are summarized in Table II.

A sample episode of harvested power, dynamic tire load and sprung mass acceleration as functions of time are shown in Fig. 10.

A. Maximizing Harvested Energy

We first examine the case of maximizing the power harvested where we choose α_1 = 0 and α_2 = 1. The results are summarized in Fig. 12, which shows that both NMPC and SLMPC significantly increase the average power harvested across all cases. More specifically, for the perfect preview case, the average power harvested for NMPC and SL-MPC is increased nearly 91% and 60%, respectively. These numbers for the cases where we have an SNR = 20,15,10 or we use the LRDE as the preview for the entire horizon are, respectively, 90% and 56%, 60 % and 49%, 43% and 28%, 24% and 15%. It can be seen that using stochastically linearized model does not cause a major performance degradation but leads to much greater computational efficiency as will be shown later.

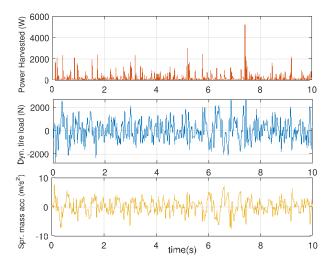


Fig. 10. Power harvested, dynamic tire load and sprung mass acceleration as a function of time for one sample of IPVA operation with SL-MPC control.

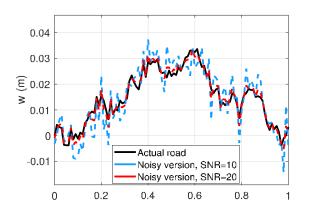


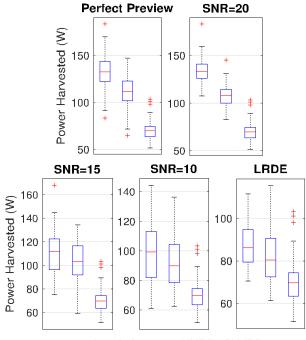
Fig. 11. Actual road profile and the noisy version of road profile with different SNR levels.

An efficiency analysis is used to further quantify the energy harvester performance. To accomplish this, we start with the derivation of mechanical and electrical efficiency associated with the linear system, followed by the IPVA system.

Denote mechanical efficiency as $\eta_{m,l}$, mean input power as $hP_{in}i$, and mean output power as $hP_{out}i$ where $h\cdot i$ denotes statistical expectation. Note that $hP_{out}i = P$ is defined in the appendix. The total input power can be calculated as the sum of the mechanical power input to the ball screw $hP_{in}i$ and power lost due to mechanical damping c_m , $hP_{Lost}i$. Note that the latter is introduced to account for mechanical energy loss due to the ball-screw and generator gears. The total input power to the ball screw is simply the product of the output torque (T) and angular velocity of the generator ($T\vartheta$). The equation of motion for the generator is given as

$$J_r \ddot{\vartheta} + c_e \dot{\vartheta} = T \tag{27}$$

Multiplying the left hand side of (27) by $\dot{\vartheta}$ and taking the statistical expectation will give the mean power input. Furthermore,



All plots from Left to right: NMPC, SLMPC, IPVA.

Fig. 12. Box plot of control performance comparison for the case of maximizing the energy harvesting. Red horizontal lines represent the mean and the box heights represent the standard deviations.

the mean dissipated power in the damper is given by

$$hP_{Lost}i = c_m R^2 \dot{\vartheta}^2 . (28)$$

Finally, the sum of the input power and power dissipated is given by

$$D \quad E \qquad D \quad E$$

$$hP_{Total}i = (c_m R^2 + c_e) \quad \dot{\vartheta}^2 + J_r \quad \ddot{\vartheta}\dot{\vartheta} \quad . \tag{29}$$

In order to calculate the second term on the right hand side of (29), $\ddot{\vartheta}$ is first written in terms of state variables according to (8). We then make use of statistical moment equations to complete the computation. The reader is referred to [28] for a detailed explanation of these equations. In this particular case, $\ddot{\vartheta}\dot{\vartheta} = 0$, implying that there is no statistical correlation between the acceleration and velocity of the ball screw. Additionally, we note that $\frac{1}{c_0}hP_{out}i = h\dot{\vartheta}^2i$ and so

$$\eta_{m,l} = \frac{h P_{out} i}{h P_{Total} i} = \frac{c_e}{c_e + c_m R^2}.$$
(30)

As seen in (30), the mechanical efficiency for the linear system is dependent on electrical damping, mechanical damping and the lead of the ball screw through $R = L/2\pi$. It is worth noting that mechanical damping $c_m = 148.32$ Ns/m was chosen such that the linear benchmark has a mechanical efficiency ($\eta_{m,l} \approx 60\%$) similar to the linear EM damper reported in [7]. The same value of c_m is employed in the IPVA system for a fair comparison.

To compute the electrical efficiency it is common to assume $c_e(t) = M \kappa_e \kappa_t / R(t)$, where $R(t) = R_{int} + R_{Load}(t)$ and R_{int} is the internal resistance of the generator [7]. For the

purpose of this paper, $R_{int} = 6.6\Omega$, $M \kappa_e \kappa_t = 11 N \cdot s \cdot m \cdot \Omega$ as directly obtained and reverse engineered from the data given in [7]. The electrical efficiency is then defined as the ratio of output with zero internal resistance to total output power $hP_{out}i$. As such, we can find

$$\eta e = \frac{R_{Load}(t)}{R_{int} + R_{Load}(t)} = \frac{M \kappa_e \kappa_t - R_{int} c_e(t)}{M \kappa_e \kappa_t}.$$
 (31)

Note that (31) applies to the system with IPVA as well. Another remark is that electrical efficiency seems to increase with an decrease in electrical damping or internal resistance. However, an electrical damping coefficient equal to zero corresponds to the case of zero power harvested. Therefore, efficiency is set to zero for cases when $c_e(t) = 0$.

In order to derive the mechanical efficiency of the IPVA system, we choose to take a Lagrangian approach with a holonomic constraint $g(\vartheta, x_{us}, x_s) = R\vartheta + x_{us} - x_s = 0$. To this end, we first rewrite (4) as

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_{i}}\right) - \frac{\partial L}{\partial q_{i}} = Q_{i} + \lambda \frac{\partial g}{\partial q_{i}},$$
(32)

where $q_1 = \vartheta$, $q_2 = \varphi$, $q_3 = x_s$, $q_4 = x_{us}$ and λ is known as a Lagrangian multiplier while it is also the constraint force acting on the IPVA required to impose the constraint g. We are further required to rewrite the Lagrangian in terms of unconstrained potential and kinetic energies. To this end, we find

$$T = T_{M_{us}} + T_{M_s} + T_c + T_p + T_r$$

$$= \frac{1}{2} M_{us} (\dot{x}_{us})^2 + \frac{1}{2} M_s (\dot{x}_s)^2$$

$$+ \frac{1}{2} m R_p^2 \dot{y}^2 + r^2 \dot{y} + \dot{\varphi}^2 + 2R_p r \cos(\varphi) \dot{y} \dot{y} + \dot{\varphi}$$

$$+ \frac{1}{2} J_r (\dot{\varphi} - \dot{y})^2, \qquad (33)$$

and

$$V = \frac{1}{2}k_s(x_s - x_{us})^2 + \frac{1}{2}k_t(x_{us} - x_r)^2.$$
 (34)

Additionally, we rewrite virtual work as

$$\delta W_{nc} = \delta W_r + \delta W_m = -c_e(\dot{\varphi} - \dot{\vartheta}) \cdot \delta \varphi$$
$$+ c_e(\dot{\varphi} - \dot{\vartheta}) \cdot \delta \vartheta - c_m(\dot{x}_s - \dot{x}_{us}) \delta \vartheta. \tag{35}$$

and Q_i corresponds to the coefficient of δq_i in (35). Substituting (35), (33) and (34) into (32) finally allows one to relate constraint force λ to state variables, accelerations and system parameters. Choosing the simplest relationship, corresponding to generalized coordinate $q_3 = x_5$, we have

$$\lambda = - M_s \ddot{x}_s + c_m R \vartheta + k R \vartheta . \tag{36}$$

Note that (36) is just a statement of Newton's second law for the sprung mass which could be expected as λ must be equal and opposite to the constraint force imposed on the sprung mass in order to satisfy the holonomic constraint. Next we sum the power dissipated by the mechanical damper and instantaneous input power to IPVA, where the instantaneous input power is

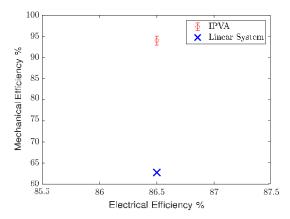


Fig. 13. Mechanical and electrical efficiency for passive IPVA and passive linear system with and error bar indicating standard deviation.

 $\lambda R\dot{\vartheta}$. The sum amounts to

$$P_{Total,IPVA} = c_m R^2 \dot{\vartheta}^2 + \lambda R \dot{\vartheta} = -R M_s \dot{\vartheta} \ddot{x}_s - k R^2 \vartheta \dot{\vartheta}. \tag{37}$$

One remark is that the the sprung mass acceleration \ddot{x}_s can be written in terms of system parameters and state variables and so this dependence is implied rather than explicitly shown in (37). The instantaneous output power going to the harvester is then simply

$$P_H = c_e \varphi - \vartheta^2. \tag{38}$$

To define the mechanical efficiency in this case we choose to compute the average input and output power in each of N=100 realizations in the time domain, followed by the computation of efficiency in each realization and an ensemble average with standard deviation computation. The efficiency in each realization is found to be

$$\eta_{m} = R_{\substack{t=t \\ t=0}}^{R_{t=t_{f}}} P_{H} dt \atop \substack{t=0 \\ t=0}, P_{Total, IPVA} dt},$$
(39)

where $t_f = 5$ as the efficiency value was found to be relatively fixed after 5 seconds. The reason for the averaging is because the calculation of power input to the IPVA may lead to negative values resulting in negative efficiencies. The mechanical efficiency and electrical efficiency for the passive IPVA as well as linear system is shown in Fig. 13.

A remark is given to the interpretation of the mechanical efficiencies shown in Fig. 13. It should be noted that mechanisms that are associated with the pendulums, including the planetary gear system and pendulum bearings, should lead to mechanical energy loss in practice. Therefore, the mechanical efficiency of IPVA should be lower in practice. In this regard, Fig. 13 only shows the maximum mechanical efficiency IPVA can theoretically have provided that the IPVA and linear benchmark have the same mechanical damping coefficient $c_m = 148.32 \text{ Ns/m}$.

This section is concluded by a comparison of the simulated power of IPVA with comparable results reported in the literature; see Table III. It should be noted that only simulation results that consider the ISO 8608 road classes are included in the table for

References	Energy harvesting technology	Model type	Road conditioning	Velocity (km/h)	Dissipated power per
					damper (W)
Guo et al. [12]	Rack-pinion EM ^a damper with MMR ^b	2DOF quarter car	Class C	97	~29-30
Huang et al. [11]	Ball-screw EM damper	SDOF quarter car	Class A	120	~9
			Class B	90	~9.5
			Class C	50	~21
			Class D	30	~24
Ataei et al. [29]	Hybrid Hydraulic-electromagnetic damper	2DOF quarter car	Class C	50	32
Sultoni et al. [30]	Linear EM damper	2DOF quarter car	Class C	50	45
Peng et al. [31]	Hydraulically driven EM damper	2DOF quarter bus	Class B	30	~42
		7DOF full bus	Class C	70	~340
Yu et al. [32]	Rack-pinion EM damper with MMR	7DOF half car	Class C	50	~15
Shi et al. [33]	Linear EM damper with active control	7DOF full car	Class B	80	~85°
Li and Zuo [13]	Rack-pinion EM damper with MMR	2DOF quarter car	Class C	36	~25
Yin et al. [34]	Rotary EM damper with active control	2DOF quarter car	Class B	~120	46
Shi et al. [35]	Linear EM damper with semi-active control	2DOF quarter car	Class C	80	46.57
Tarantini [36]	Ball-screw EM damper	2DOF quarter car	Class C	36	~26
Bao et al. [37]	Hydraulic EHSA	2DOF quarter car	Class C	72	42.5
Chen et al. [2]	Rotary EM damper with inertial nonlinearity & MPC	2DOF quarter car	Class C	not mentioned	11.04
Current Work	Ball-screw-based IPVA with MPC	2DOF quarter car	Class C	90	~70 to ~133
					(perfect preview)

 $\begin{tabular}{l} TABLE III \\ SURVEY OF SIMULATION RESULTS AMONGST RELATED WORKS IN THE FIELD \\ \end{tabular}$

a fair comparison. The reader is referred to [1] for the harvested power of other EHSAs that were obtained experimentally or numerically under other excitation conditions. As shown in Table III, the IPVA system outperformed most of the reported results.

B. Maximizing Ride Comfort

We then examine the case of maximizing the ride comfort (i.e., minimizing the sprung mass acceleration) where we choose $\alpha_1 = 1$ and $\alpha_2 = 0$. The results are summarized in Fig. 14. It can be seen that with perfect preview and average to high levels of SNRs, both NMPC and SL-MPC can also enhance the ride comfort. More specifically, for the case when we have the perfect preview the average RMS value for the sprung mass acceleration is decreased by 15% and 8 % for the NMPC and SL-MPC, respectively. These numbers for the cases where SNR = 20, 15, 10 are 11% and 7%, 8% and 7%, 1% and 2%, respectively. However, with the LRDE preview, both NMPC and SLMPC fail to outperform the passive IPVA. This is because the passive design is set with the maximum damping, which leads to good ride comfort. So without accurate predictions, MPC finds it difficult to outperform the passive design regarding the ride comfort.

C. Mixed Objective

We next examine the mixed objective case, where both the power harvested and ride comfort are considered. In this case, α_1 is chosen as 1 and α_2 is varied from 0.01 \square 0.1 to observe the trade-off between the power harvested and the ride com-fort. Fig. 15 summarizes the results. For each α_2 , the SLMPC and NMPC designs are simulated 500 times and the average performance are reported. It can be seen that by varying α_2 , different trade-offs between power harvested and ride comfort can be obtained. One can choose an appropriate value that suits best for the design specifications. It can be seen that both NMPC and SL-MPC has multiple parameter settings that offer

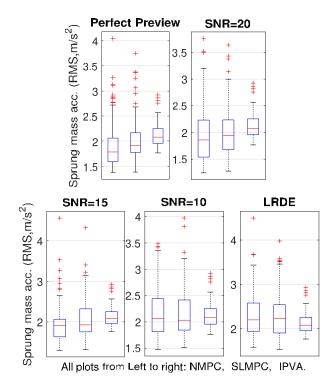


Fig. 14. Box plot of control performance comparison for the case of minimizing vertical acceleration. Red horizontal lines represent the means while the box heights represent the standard deviations.

better trade-offs in both energy harvesting and ride comfort as compared to the passive designs (i.e., in the shaded green areas). As an example, in the perfect preview case for the circle object pointed by the arrow (NMPC), the harvested power and ride comfort are improved by 55% and 9%, respectively, and for the cross shaped object pointed by the arrow (SL-MPC), the improvements are 29% and 4%, respectively.

Lastly, we evaluate the incurred computational complexity of the two MPC approaches. The simulations are done on a PC with

^a EM: electromagnetic, ^b MMR: mechanical motion rectification, ^c The power is an average of 98 kJ over simulation time of 1150 s.

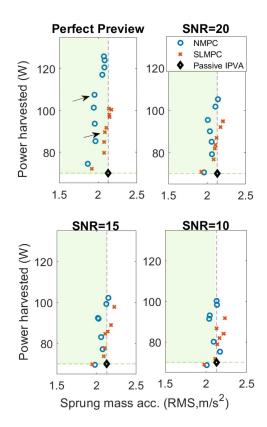


Fig. 15. Control performance comparison for the case with mixed objective.

a 2.5 GHz Intel Core i7-4710HQ CPU with 16 GB of internal memory using MPCTools [38] and CasADi [39]. It is observed that the average computation time per 1000 steps for SL-MPC is 18 seconds, significantly less than that of NMPC, which is 43 seconds. The results show a clear advantage of SLMPC over NMPC in terms of computation complexity, which is a critical factor for online implementations.

D. Vehicle Handling Performance

In this subsection, we investigate the vehicle handling performance of the proposed EHSA. For this goal the stage cost function defined in (15) is modified as follows:

$$I(x, u) = \alpha_1(\dot{x}_6 + R\dot{x}_2)^2 - \alpha_2 u(x_2 - x_4)^2 + \alpha_3(x_5 - w).$$
(40)

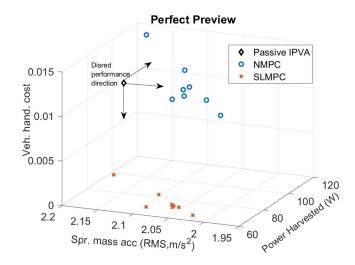


Fig. 16. Control performance comparison for the case with mixed objective.

V. CONCLUSION

In this paper, a nonlinear IPVA was integrated into a quartercar model, and the system with optimized parameters was shown to present simultaneous improvement in vibration control and energy harvesting. To further improve the performance, we investigated two MPC designs: nonlinear MPC (NMPC) and stochastic linearization MPC (SL-MPC), with different settings of road preview. A new SL approach was developed to stochastically linearize systems with guaranteed stabilizability. Extensive simulations were performed which showed that SL-MPC had no major performance degradation and it significantly improved the computational efficiency. Specifically, our simulation results show that the power harvested has the potential to be increased by 60% and the RMS value of the sprung mass acceleration can be reduced by 7% w.r.t the passive case. We showed that road preview has a great impact on control performance. Future research will focus on investigating the stability of SL-MPC as well as developing a prototype system to demonstrate the proposed framework.

APPENDIX

A. Linear System Analytical Solution

In the subsection, we derive the closed-form solution for the average power and acceleration associated with the linear benchmark shown in Fig. 3. Towards that end, (8) and (9) are first transformed into the frequency domain and the ratio of output velocity to input road velocity is determined. Namely,

$$\frac{\dot{\vartheta}(\omega)}{\dot{x}_r(\omega)} = \frac{R\omega^2 k_t M_s}{A + Bi},\tag{41}$$

where

$$A = J_r \omega^2 \omega^2 (M_{us} + M_s) - k_t + R^2 k_s k_t - R^2 \omega^2 k_t M_s + k_s M_s + k_s M_{us} - \omega^2 M_s M_{us},$$

and
$$B = \omega[k_t - \omega^2(M_s + M_{us})](c_e + R^2c_m)$$
 with $i = \sqrt[V]{-1}$.

The road disturbance x_r is modeled as white noise passing through a first-order filter [26] and the corresponding power spectral density, $S_{x_r}(\omega)$, can be written as:

$$S_{x_r}(\omega) = \frac{2\pi G_r V}{\omega^2 + \omega_c^2},\tag{42}$$

where G_r is a road roughness coefficient specific to the class-C road as given by [26], V is the driving speed in m/s and ω_c is the cutoff frequency that keeps the power spectral density bounded at $\omega = 0$. Note that ω_c needs to be significantly smaller than the resonant frequencies of the suspension [40] and it is assumed to be zero when deriving closed-form solutions for the linear benchmark. Otherwise, ω_c is chosen as 0.01 rd/sec. Given $\omega_c = 0$, (42) implies that the time derivative of the road disturbance (i.e., \dot{x}_r) is a white noise with intensity of $2\pi G_r V$. Noting that the instantaneous power is given by $P = c_c \vartheta^2$ and proceeding with the principle of H_2 norm [21], the average power becomes:

$$P = \frac{C_e}{2\pi} \sum_{-\infty}^{\infty} 2\pi G_r V \frac{\dot{\vartheta}}{\dot{x}_r} d\omega = \frac{e_C}{c} \frac{\pi V G_r^t \dot{k}}{\frac{e_R^2 C_m}{m}}$$
(43)

Likewise, the RMS sprung mass acceleration, $\sigma_{\ddot{x}_s}$, can be found by first deriving

$$\ddot{x}_s(\omega) \qquad i \, k_s R^2 - \omega^2 J_r - \omega \, c_e + c_m R^2 \, \dot{x}_r(\omega)$$

$$\stackrel{t}{=} = k \, \omega - A + B i \qquad (44)$$

where A and B were previously defined. The mean square acceleration can then be determined by integrating the power spectral density for the acceleration over the entire frequency domain, noting that the power spectral density value for \dot{x}_r is $S_{\dot{x}_r}(\omega) = 2\pi G_r V$ and the power spectral density of the sprung mass acceleration is $S_{\ddot{x}_s}(\omega) = \frac{|\ddot{x}_s|}{\hat{x}_r} I^2 S_{\dot{x}_r}(\omega)$. The RMS value is then simply defined as

$$\sigma_{\ddot{x}_{s}} = \int_{-\infty}^{\infty} \frac{Z}{2\pi} \int_{-\infty}^{\infty} S_{\ddot{x}_{s}}(\omega) d\omega = \int_{-\infty}^{\infty} \frac{Z}{2\pi} \int_{-\infty}^{\infty} 2\pi G_{r} V \frac{\ddot{x}_{s}^{2}}{\dot{x}_{r}^{2}} d\omega.$$
(45)

Equation (45) can then be integrated with the known integration formula [21] to obtain

$$\sigma_{\ddot{x}_s} = \frac{p}{\pi V G_r C^{-1} (a_0 + a_1 R^2 + a_2 R^4 + a_3 R^6)}, \quad (46)$$

where $M = M_s + M_{us}$ corresponds to the total mass and $C = R^2 M^2 c_e + R^2 c_m J_r M + R^2 M_s M_{us}$,

$$a_0 = J_r k_t c^2 M + J^3 k_t^2$$

$$a_1 = 2c_e c_m J_r M k_t + c_e^2 M_s M_{us} + k_t^2 M_s - J_r^2 2k_s M k_t$$

$$a_2 = 2c_e c_m k_t M_s M_{us} + c_m^2 J_r k_t M$$

$$+ k_s^2 J_r M^2 - 2k_t k_s J_r M_s M_{us}$$

$$a_3 = M_s M_{us} c_m^2 k_t + k_s^2 M_s M_{us} M.$$

Equations (43) and (46) are then used to plot the closed form solution shown in Fig. 3.

B. Stochastic Linearization With Guaranteed Stability

In this subsection we advance the SL approach so that the linearized model (e.g., (26)) has guaranteed stabilizability. More specifically, we aim at reformulating the problem to find a matrix A_{Is} such that (A_{Is}, B_I) is stabilizable and is close to A_I (e.g., $A_I = A_{Is}$ if (A_I, B_I) is already stabilizable). We next review some definitions regarding stabilizability to place our method in proper context.

Definition 1. Consider a linear time invariant (LTI) system $\dot{x} = Ax + Bu$. The pair (A, B) is called stabilizable if all its uncontrollable eigenvalues are stable.

The following lemma can be used to check the stabilizability of an LTI system.

Lemma 1 [41]. Consider an LTI system $\dot{x} = Ax + Bu$, where A is $n \times n$ and B is $n \times m$. The controllability matrix C is defined as $C = [B, AB, \dots A^{n-1}B]$. If rank C = q < n, then there exists a nonsingular $n \times n$ matrix T such that

$$\hat{P} = T^{-1}AT = \frac{1\hat{A} + 1\hat{A}}{0\hat{A} + 2\hat{A}} = \frac{1\hat{B} + 1\hat{A}}{0\hat{A} + 2\hat{A}} = \frac{1\hat{B} + 1\hat{B}}{0\hat{A} + 1\hat{B}} = \frac{1\hat{B} + 1\hat{B}}{0\hat{A}} = \frac{1\hat{B}}{0\hat{A}} = \frac{$$

where \hat{A}_{11} is $q \times q$, \hat{B}_1 is $q \times m$, and the pair $(\hat{A}_{11}, \hat{B}_1)$ is controllable. The transformation matrix T is called controllability decomposition matrix and can be constructed as follows. Let $T = [X \ Y]$, then X is is an $n \times q$ matrix whose columns span the columns space of C, and Y is an $n \times (n - q)$ matrix whose columns are chosen such that T is nonsingular.

The eigenvalues of \hat{A}_{11} are called the controllable eigenvalues and those of \hat{A}_{22} are called the uncontrollable eigenvalues. Following Lemma 1, if we transform A_I as $\hat{A}_I = T^{-1}A_IT = \begin{bmatrix} A_{I,11} & A_{I,12} \\ 0 & \hat{A}_{I,22} \end{bmatrix}$, then A_I is stabilizable if and only if $A_{I,22}$ is stable.

If this is not the case, then we seek a Hurwitz matrix $\hat{A}_{15,22}$ that is similar to $\hat{A}_{1,22}$ (e.g., in the sense of matrix norms) to replace $\hat{A}_{1,22}$. If $\hat{A}_{1,22}$ is not Hurwitz, we aim for finding the Hurwitz matrix A_{eq} where its second norm difference with $\hat{A}_{1,22}$ is minimized. To this end, using the Lyapunov inequality for verifying the stability of a linear system [42], the optimization problem is such that

min.
$$A_{eq} - \hat{A_{l,22}}^{2}_{eq}^{2}$$

s.t. $A_{eq}^{T}P + PA_{eq} < 0, P > 0.$ (47)

Two decision variables are multiplied together in the first constraint, making the problem a BMI (bilinear matrix inequality) optimization problem, which can be solved using the MATLAB BMIsolver package. After finding A_{eq} , this matrix will replace $A_{l,22}$ in A_l and one can use the matrix T to transform back and get the matrix A_l where the pair (A_l, B_l) is now stabilizable.

C. High Gain Disturbance Observer Design

In this subsection, we show the development of high-gain observer (HGO [43]) to estimate the road disturbance that is exploited in the above MPC formulations in the prediction horizon. Specifically, we rewrite the dynamics of the controlled

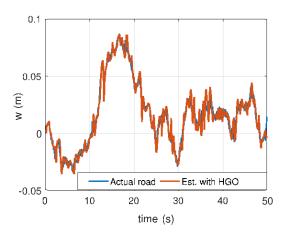


Fig. 17. Road profile estimation using HGO.

energy harvesting system in a compact form based on (13) as follows:

$$\dot{x}_1 = x_2, \dot{x}_2 = f_2(x, u, w),
\dot{x}_3 = x_4, \dot{x}_4 = f_4(x, u, w),
\dot{x}_5 = x_6, \dot{x}_6 = f_6(x, u, w).$$
(48)

where $f_6(x, u, w) = b_1(x, u) + wb_2(x)$.

Eqn. (48) consists of 3 separate sub-dynamics in the normal nonlinear SISO form with $y_1 = x_1$, $y_2 = x_3$ and $y_3 = x_5$ being our measurements. In this regard, in order to design the extended observer, we take $\sigma = f_6(x, u, w)$ as the function to be estimated. Thus, the observer equations are formulated as follows:

$$\dot{\hat{x}}_{1} = \hat{x}_{2} + (\alpha_{1}/\varepsilon_{1})(y_{1} - \hat{x}_{1}),
\dot{\hat{x}}_{2} = \alpha_{2}/\varepsilon_{1}^{2}(y_{1} - \hat{x}_{1}),
\dot{\hat{x}}_{3} = \hat{x}_{4} + (\alpha_{3}/\varepsilon_{2})(y_{2} - \hat{x}_{3}),
\dot{\hat{x}}_{4} = \alpha_{4}/\varepsilon_{2}(y_{2} - \hat{x}_{3}),
\dot{\hat{x}}_{5} = \hat{x}_{6} + (\alpha_{5}/\varepsilon_{3})(y_{3} - \hat{x}_{5}),
\dot{\hat{x}}_{6} = \hat{\sigma} + \alpha_{6}/\varepsilon_{3}^{2}(y_{3} - \hat{x}_{5}),
\hat{\sigma} = \alpha_{7}/\varepsilon_{3}(y_{3} - \hat{x}_{5}).$$
(49)

Here ε_1 , ε_2 and ε_3 are sufficiently small positive constants; α_i and α_{i+1} , i = 1, 2, are chosen such that $s^2 + \alpha_i s + \alpha_{i+1}$ is Hurwitz; and α_5 , α_6 and α_7 are also chosen such that $s^3 + \alpha_5 s^2 + \alpha_6 s + \alpha_7$ is Hurwitz. Then at time t, the disturbance can be estimated as [43]:

$$\hat{w}(t) = \frac{\hat{\sigma}(t) - b_1(\hat{x}(t), u(t))}{b_2(\hat{x}(t))}.$$

A sample of the estimation is shown in Fig. 17.

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