

WEYL LAWS FOR OPEN QUANTUM MAPS

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ABSTRACT. We find Weyl upper bounds for the quantum open baker's map in the semiclassical limit. For the number of eigenvalues in an annulus, we derive the asymptotic upper bound $\mathcal{O}(N^\delta)$ where δ is the dimension of the trapped set of the baker's map and $(2\pi N)^{-1}$ is the semiclassical parameter, which improves upon the previous result of $\mathcal{O}(N^{\delta+\varepsilon})$. Furthermore, we derive a Weyl upper bound with explicit dependence on the inner radius of the annulus for quantum open baker's maps with Gevrey cutoffs.

1. INTRODUCTION

Open quantum maps provide simple finite-dimensional models of open quantum chaos. This makes them especially conducive to numerical experimentation and thus appealing in the study of scattering resonances. They quantize a symplectic relation on a compact phase space. Such relations are toy models for Poincaré sections that arise when considering scattering Hamiltonians with hyperbolic trapped sets. See papers by Nonnenmacher–Sjöstrand–Zworski [NSZ11, NSZ14] for the precise description of the reduction from specific open quantum systems to open quantum maps using Poincaré sections. In this paper, the symplectic relation we consider is the classical baker's map on a 2-torus, which gives rise to the quantum open baker's map. We find a Weyl upper bound for the number of eigenvalues in an annulus.

The quantum open baker's map is an operator on

$$\ell_N^2 = \ell^2(\mathbb{Z}_N), \quad \mathbb{Z}_N = \mathbb{Z}/(N\mathbb{Z})$$

defined by the triple

$$(M, \mathcal{A}, \chi), \quad M \in \mathbb{N}, \quad \mathcal{A} \subset \{0, \dots, M-1\}, \quad \chi \in C_0^\infty((0, 1); [0, 1]). \quad (1.1)$$

Here, M is the *base*, $\mathcal{A}$ is the *alphabet*, and χ is the *cutoff*. Put $N = KM$ where $K \in \mathbb{N}$. Then the quantum open baker's map is given by

$$B_N = \mathcal{F}_N^* \begin{pmatrix} \chi_{N/M} \mathcal{F}_{N/M} \chi_{N/M} & & \\ & \ddots & \\ & & \chi_{N/M} \mathcal{F}_{N/M} \chi_{N/M} \end{pmatrix} I_{\mathcal{A}, M}$$

where $\mathcal{F}_N$ is the unitary discrete Fourier transform, $I_{\mathcal{A}, M}$ is a diagonal matrix whose (j, j) -th entry is equal to one if $\lfloor j/K \rfloor \in \mathcal{A}$ and zero otherwise, and $\chi_{N/M}$ is a discretized

smooth cutoff function. For example, for the triple $(3, \{0, 2\}, \chi)$ and $N = 3K$, the corresponding quantum open baker's map is then

$$B_N = \mathcal{F}_N^* \begin{pmatrix} \chi_K \mathcal{F}_K \chi_K & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \chi_K \mathcal{F}_K \chi_K \end{pmatrix}$$

Define the canonical relation on the torus $\mathbb{T}_{x,\xi}^2$ by

$$\begin{aligned} \mathfrak{K}_{M,\mathcal{A}} : (y, \eta) &\mapsto (x, \xi) = \left(My - a, \frac{\eta + a}{M} \right), \\ (y, \eta) &\in \left(\frac{a}{M}, \frac{a+1}{M} \right) \times (0, 1), \quad a \in \mathcal{A}. \end{aligned} \tag{1.2}$$

Then the corresponding semiclassical Fourier integral operator is given by

$$\begin{aligned} \mathcal{U}_h &:= \sum_{a \in \mathcal{A}} \mathcal{U}_h^a \quad \text{where} \\ \mathcal{U}_h^a v(x) &= \frac{M}{2\pi h} \int_{\mathbb{R}^2} e^{\frac{i}{h}((x+a-My)\theta + xa/M)} \chi(M\theta) \chi(My - a) u(y) dy d\theta. \end{aligned}$$

The quantum open baker's map can then be seen as the discrete analogue of this Fourier integral operator with the corresponding semiclassical parameter $(2\pi N)^{-1}$. For a rigorous analysis of the analogy, see papers of Degli Esposti–Nonnenmacher–Winn [ENW06] and Nonnenmacher–Zworski [NZ06]. Heuristics can be found in earlier works of Balázs–Voros [BV89] and Saraceno–Voros [SV94]. In view of this analogy, one would then expect that forward in time propagation by B_N would lead to localization in frequency space to the Cantor set and backward propagation by B_N would lead to localization in physical space to the Cantor set. Indeed, Fig. 1 demonstrates this property numerically.

Following the above observations, we should expect the eigenfunctions to be localized in frequency near the $(M, \mathcal{A})$ -Cantor set provided that the eigenvalues are not too small (see Fig. 2). The maximum number of eigenfunctions that can be packed into such a region in phase space should then be on the order of N^δ where

$$\delta = \frac{\log |\mathcal{A}|}{\log M} \tag{1.3}$$

is the dimension of the Cantor set. Our result uses such localization properties to provide rigorous upper bounds to the number of eigenvalues of B_N above a threshold. More specifically, consider the eigenvalue counting function

$$\mathcal{N}_N(\nu) = |\text{Spec}(B_N) \cap \{|\lambda| \geq M^{-\nu}\}|, \tag{1.4}$$

defined for $\nu \geq 0$, where the eigenvalues are counted with multiplicities. Then we have the following Weyl upper bound:

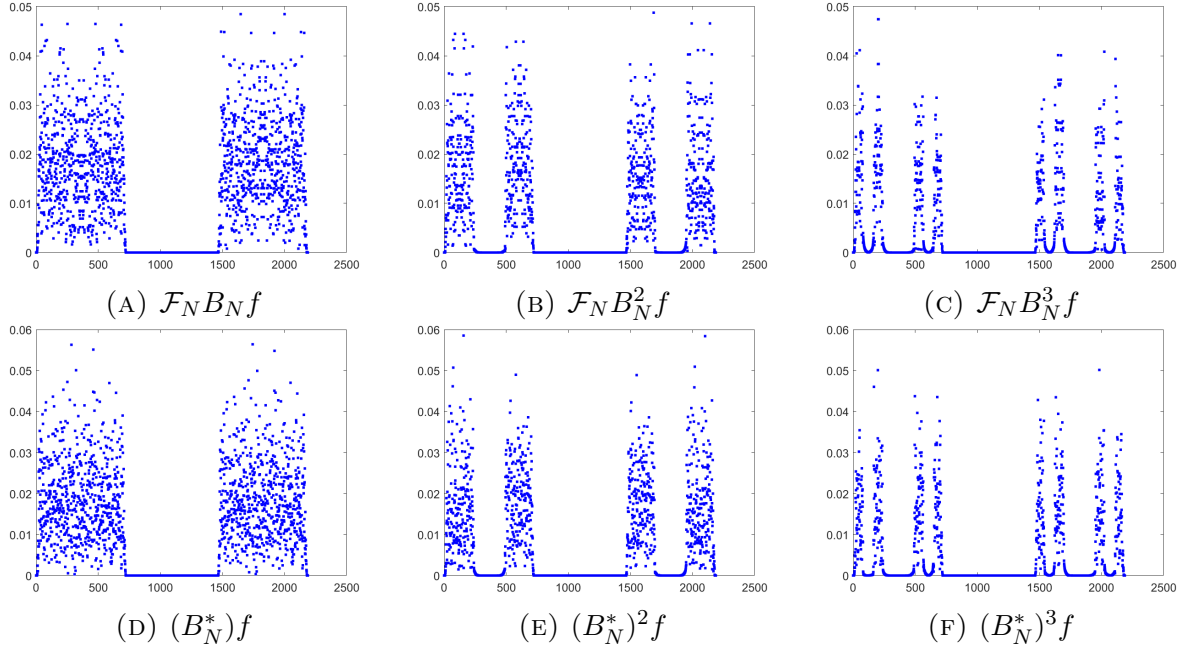


FIGURE 1. A demonstration of the localizing properties of B_N for $M = 3$, $\mathcal{A} = \{0, 2\}$, and $N = 3^7$. An ℓ_N^2 normalized vector f was chosen uniformly at random. Plots (A) – (C) are the frequency side of forward propagation, and (D) – (F) are the spatial side of backward propagation. Each figure plots the absolute value of the indicated vector as a map from $\mathbb{Z}_N \rightarrow \mathbb{R}$

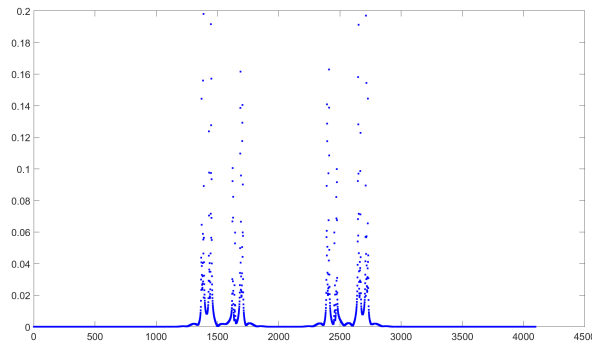


FIGURE 2. The Fourier side of a typical eigenvector. $M = 4$, $\mathcal{A} = \{1, 2\}$, and χ is identically 1 on the Cantor set generated by M and $\mathcal{A}$. Here, $N = 4^6$, and the absolute value of $\mathcal{F}_n v$ is plotted as a function from $\mathbb{Z}_N \rightarrow \mathbb{R}$ where v is eigenvector with the 50th largest eigenvalue at $|\lambda| \approx M^{-0.4869}$.

Theorem 1. *For each $\nu > 0$, we have as $N = KM \rightarrow \infty$,*

$$\mathcal{N}_N(\nu) = \mathcal{O}(N^\delta). \quad (1.5)$$

The proof of the theorem in §4.1 follows the methods used in [DJ17], in which the bound

$$\mathcal{N}_N(\nu) = \mathcal{O}(N^{\delta+\varepsilon})$$

for any $\varepsilon > 0$ was proved. We obtain the ε -improvement in this paper by using tighter propagation estimates and a modified approximate inverse identity (see §3.1 for details).

With stronger assumptions on the decay of the cutoff function χ , an explicit dependence of the upper bound on the depth of the spectrum ν can be extracted. In particular, we consider Gevrey class functions, first introduced in [Gev18] to study regularity of solutions to the heat equation. Given $s \geq 1$, a function $f \in \mathcal{C}^\infty(\mathbb{R})$ is s -Gevrey if for every compact $K \subset \mathbb{R}$, there exists a constant $C_{K,f}$ such that

$$\sup_{x \in K} |\partial^\alpha f(x)| \leq C_{K,f}^{\alpha+1} (\alpha!)^s$$

for all $\alpha \in \mathbb{Z}^+$. For $s = 1$, this is simply the space of real analytic functions, which cannot be compactly supported. However, for every $s > 1$, there exist smooth and compactly supported s -Gevrey functions. For $s > 1$, we write

$$\mathcal{G}_c^s((0, 1)) = \{f \in \mathcal{C}_c^\infty(\mathbb{R}) : f \text{ is } s\text{-Gevrey and } \text{supp } f \subset (0, 1)\}$$

Observe that if $\chi \in \mathcal{G}_c^s((0, 1))$ for some $s > 1$, then

$$|\hat{\chi}(\xi)| \leq C e^{-c|\xi|^{1/s}}, \quad (1.6)$$

for some positive constants C and c . Here, $\hat{\chi}$ denotes the usual Fourier transform given by

$$\hat{\chi}(\varepsilon) = \int_{\mathbb{R}} e^{-2\pi i x \varepsilon} \chi(x) dx. \quad (1.7)$$

So even though we cannot have exponential decay of the Fourier transform that comes with analyticity, we can still get arbitrarily close. Finally, observe that for $\eta > 0$,

$$\int_{\eta}^{\infty} |\hat{\chi}(\xi)| d\xi \leq C e^{-\tilde{c}\eta^{1/s}} \quad (1.8)$$

for some new constant $\tilde{c} < c$. See [Rod93, Chapter 1] for a more detailed account on Gevrey classes and their Fourier decay properties. With this stronger cutoff decay assumption, we then have the following Weyl upper bound:

Theorem 2. *Assume that $\chi \in \mathcal{G}_c^s((0, 1))$ for some $s > 1$, then for all $\nu \geq 1$ and all sufficiently large $N = KM$ where $K \in \mathbb{N}$,*

$$\mathcal{N}_N(\nu) \leq C N^\delta \nu^{(1-\delta)s} \quad (1.9)$$

where the constant C depend only on χ and M .

In the study of quantum chaos, open quantum systems given by the Laplacian on a noncompact Riemannian manifold whose geodesic flow is hyperbolic on the trapped set provide an important mathematical model. In the papers by Nonnenmacher–Sjöstrand–Zworski [NSZ11, NSZ14], the study of resonances for such open quantum system is reduced to the open quantum map, so we should expect our results to run parallel to previous Weyl upper bounds for open quantum systems. We note that in the correspondence between quantum systems and quantum maps, if $\omega \in \mathbb{C}$ with $\text{Im } \omega \leq 0$ is a scattering resonance of the open quantum system, then

$$\lambda = e^{-i\omega \log M} = M^{-i\omega} \quad (1.10)$$

is a corresponding eigenvalue of B_N , which makes sense in view of the fact that B_N can be thought of as a toy model for the time $t = \log M$ propagator of an open quantum system with expansion rate 1. This means that Weyl upper bounds in horizontal strips below the real line should correspond to our Weyl law in an annulus. Weyl upper bounds for resonances of the Laplacian in strips (which corresponds to annuli for the open quantum maps by Eq. (1.10)) were first proved by Sjöstrand [Sjö90]. This was done in the analytic category, which we cannot afford in our case since the cutoff χ is compactly supported. Using Eq. (1.10), the corresponding bound that Sjöstrand found would give $\mathcal{N}_N(\nu) = \mathcal{O}(N^\delta \nu^{1-\delta})$ where δ is the Minkowski dimension of the trapped set. Since we can only assume Gevrey for $s > 1$ in our setting, we see a corresponding loss in our result as we only have $\mathcal{O}(N^\delta \nu^{s(1-\delta)})$. However, we remark that it appears from numerical experiments in §5.2 that if the cutoff is identically 1 near the trapped set, the Sjöstrand bound of $\mathcal{O}(N^\delta \nu^{1-\delta})$ is recovered for ν not too large (Fig. 4), but our methods do not appear to be able to account for this behavior.

Weyl upper bounds for the Laplacian in fixed strips have been proved in various smooth settings by Guillopé–Lin–Zworski [GLZ04], Zworski [Zwo99], Sjöstrand–Zworski [SZ07], Nonnenmacher–Sjöstrand–Zworski [NSZ11, NSZ14], and Datchev–Dyatlov [DD13]. These give the corresponding bounds $\mathcal{N}_N(\nu) = \mathcal{O}(N^\delta)$, which aligns with our result in Theorem 1. Physical microwave experiments on the Weyl law asymptotics have been done by Potzuweit et al. [PWB⁺12], and various numerical experiments can be found in, Lu–Sridhar–Zworski [LSZ03], Borthwick–Weich [BW16], Borthwick [Bor14], and Borthwick–Dyatlov–Weich [DBW19]. The main idea behind deriving the Weyl upper bounds is the localization of the eigenfunctions in phase space (see Proposition 3.1), and this was observed numerically by Keating et al. [KNPS06]. In the setting of Walsh quantization, which uses a modified Fourier transform, the localization was later proved by Keating et al. [KNNS08].

2. OPEN QUANTUM MAPS

In this section, we establish some basic definitions and a general nonstationary phase estimate. We then give a more detailed definition of the quantum open baker's map, and use the nonstationary phase estimate to prove the one-step propagation estimate for the quantum open baker's map B_N that will be iterated in order to get propagation of singularities estimates for long times.

2.1. Preliminaries. For $N \in \mathbb{N}$, we have the abelian group

$$\mathbb{Z}_N := \mathbb{Z}/N\mathbb{Z} \simeq \{0, \dots, N-1\},$$

and we have the associated $\ell_N^2 = \ell^2(\mathbb{Z}_N)$ space of functions $u : \mathbb{Z}_N \rightarrow \mathbb{C}$ equipped with the norm

$$\|u\|_{\ell_N^2}^2 = \sum_{j=0}^{N-1} |u(j)|^2.$$

The unitary Fourier transform on ℓ_N^2 is given by

$$\mathcal{F}_N u(j) = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} \exp\left(-\frac{2\pi i j \ell}{N}\right) u(\ell).$$

Given a function $\varphi : [0, 1] \rightarrow \mathbb{C}$, its discretization is a function denoted $\varphi_N \in \ell_N^2$ given by

$$\varphi_N(j) = \varphi\left(\frac{j}{N}\right), \quad j \in 0, \dots, N-1. \quad (2.1)$$

We denote the corresponding Fourier multiplier by

$$\varphi_N^{\mathcal{F}} = \mathcal{F}_N^* \varphi_N \mathcal{F}_N \quad (2.2)$$

For the distance function on $[0, 1]$, we consider the interval with 0 and 1 identified. In particular,

$$d(x, y) = \min\{|x - y|, 1 - |x - y|\}.$$

For the distance between sets, we have the usual

$$d(U, V) = \inf_{x \in U, y \in V} d(x, y).$$

Now we have the following nonstationary phase estimate.

Lemma 2.1. *Fix $\chi \in C_c^\infty((0, 1))$. Assume that $a \in \mathbb{Z}_N$ and*

$$d\left(\frac{a}{N}, 0\right) \geq r$$

for some $r \in (0, 1/2)$. Then

$$\left| \sum_{m=0}^{N-1} \exp\left(\frac{2\pi i a m}{N}\right) \chi\left(\frac{m}{N}\right) \right| \leq N \cdot g_\chi(Nr) \quad (2.3)$$

where

$$g_\chi(x) \leq \begin{cases} C_n x^{-n} & \forall n \text{ if } \varphi \in C_c^\infty((0,1)) \\ C e^{-c x^{1/s}} & \text{if } \varphi \in \mathcal{G}_c^s((0,1)) \end{cases} \quad (2.4)$$

The positive constants C_n , C , and c depend only on the choice of χ .

Proof. The Poisson summation formula gives

$$\sum_{m=0}^{N-1} \exp\left(\frac{2\pi i a m}{N}\right) \chi\left(\frac{m}{N}\right) = N \sum_{\ell \in \mathbb{Z}} \hat{\chi}(N\ell - a).$$

Note that $\hat{\chi}$ is rapidly decaying and by our assumption, $N\ell - a > Nr$. Therefore, for every $n \geq 0$,

$$\begin{aligned} \left| N \sum_{\ell \in \mathbb{Z}} \hat{\chi}(N\ell - a) \right| &\leq CN \sum_{\ell \in \mathbb{Z}} (N\ell - a)^{-n} \\ &\leq 2CN \sum_{\ell \geq 0} (N(r + \ell))^{-n} \\ &\leq C_n N (Nr)^{-n} \end{aligned}$$

where the constant depends only on n . Similarly, if $\chi \in \mathcal{G}_c^s((0,1))$ for some $s > 1$, then in view of Eq. (1.6) and Eq. (1.8),

$$\begin{aligned} \left| N \sum_{\ell \in \mathbb{Z}} \hat{\chi}(N\ell - a) \right| &\leq CN \sum_{\ell \in \mathbb{Z}} \exp\left(-c|N\ell - a|^{\frac{1}{s}}\right) \\ &\leq 2CN \cdot \exp\left(-c(Nr)^{\frac{1}{s}}\right) + 2C \int_{Nr}^{\infty} \exp\left(-c\xi^{\frac{1}{s}}\right) d\xi \\ &\leq \tilde{C} N \exp\left(-\tilde{c}(Nr)^{\frac{1}{s}}\right), \end{aligned}$$

where we interpret the sum as a lower Riemann sum to bound by the integral. Again, all the constants above depend only on χ . Therefore we have both of the desired estimates. $\square$

2.2. One-step propagation. Let the triple $(M, \mathcal{A}, \chi)$ be as in Eq. (1.1), and put $N = KM$ where $K \in \mathbb{N}$. Define the projection $\Pi_a : \ell_N^2 \rightarrow \ell_{N/M}^2$, $a \in \{0, \dots, M-1\}$, by

$$\Pi_a u(j) = u(j + aK), \quad u \in \ell_N^2, j \in \{0, \dots, K-1\}. \quad (2.5)$$

Then the open quantum map $B_N : \ell_N^2 \rightarrow \ell_N^2$ can be written as

$$B_N = \sum_{a \in \mathcal{A}} \mathcal{F}_N^* \Pi_a^* \chi_{N/M} \mathcal{F}_{N/M} \chi_{N/M} \Pi_a.$$

Expanding out the Fourier transforms, we have the formula

$$B_N u(j) = \frac{\sqrt{M}}{N} \sum_{a \in \mathcal{A}} \sum_{m, \ell=0}^{K-1} \exp \left[2\pi i \left(\frac{(j - M\ell)m}{N} + \frac{ja}{M} \right) \right] \chi\left(\frac{m}{K}\right) \chi\left(\frac{\ell}{K}\right) u(\ell + aK). \quad (2.6)$$

It will be useful in the propagation estimates to define the expanding map

$$\Phi = \Phi_{M, \mathcal{A}} : \bigsqcup_{a \in \mathcal{A}} \left(\frac{a}{M}, \frac{a+1}{M} \right) \rightarrow (0, 1)$$

given by

$$\Phi(x) = Mx - a, \quad x \in \left(\frac{a}{M}, \frac{a+1}{M} \right) \quad (2.7)$$

We will obtain estimates in terms of the constants C_n for the propagation of singularities. In particular, we start by showing that by applying B_N once, the resulting function will be roughly microlocalized to $\varkappa_{M, \mathcal{A}}((0, 1)^2)$, and by applying B_N^* once, the resulting function will be roughly microlocalized to $\varkappa_{M, \mathcal{A}}^{-1}((0, 1)^2)$ (of course, this is imprecise since our setting is discrete). This localization behavior is clear in the classical open baker's map $\varkappa_{M, \mathcal{A}}$. In the discrete setting, it is then natural to consider B_N as a matrix consisting of blocks that reflect the classical structure of the baker's map, and each block will be rapidly decaying away from the diagonal, so then we can apply Schur's bound to control the norm. To make precise the above heuristics, we have the following estimate.

Proposition 2.2. *Assume that $\varphi, \psi : [0, 1] \rightarrow [0, 1]$ such that*

$$d(\text{supp } \psi, \Phi^{-1}(\text{supp } \varphi)) \geq r \quad (2.8)$$

where Φ is the expanding map as defined in Eq. (2.7) and r is a small gap satisfying

$$0 < Mr \leq 2d(\text{supp } \chi, 0). \quad (2.9)$$

Let $\psi_N^{\mathcal{F}}$ and $\varphi_N^{\mathcal{F}}$ are Fourier multipliers as defined in (2.2). Then

$$\begin{aligned} \|\varphi_N B_N \psi_N\|_{\ell_N^2 \rightarrow \ell_N^2} &\leq \tilde{g}_\chi(Nr) \\ \|\psi_N^{\mathcal{F}} B_N \varphi_N^{\mathcal{F}}\|_{\ell_N^2 \rightarrow \ell_N^2} &\leq \tilde{g}_\chi(Nr) \end{aligned}$$

where

$$\tilde{g}_\chi(x) \leq \begin{cases} C_n x^{-n} & \forall n \text{ if } \varphi \in C_c^\infty((0, 1)) \\ C e^{-cx^{1/s}} & \text{if } \varphi \in \mathcal{G}_c^s((0, 1)) \end{cases} \quad (2.10)$$

where C_n , c , and C are positive constants depending only on χ .

Proof. 1. We computed each entry of $\varphi_N B_N \psi_N$ as an $N \times N$ matrix, $N = KM$ for some $K \in \mathbb{N}$. From the expansion Eq. (2.6) for B_N , we can write

$$\varphi_N B_N \psi_N u(j) = \sum_{a \in \mathcal{A}} \sum_{\ell=0}^{K-1} A_{j\ell}^a u(\ell + aK)$$

where

$$\begin{aligned} A_{j\ell}^a &= \frac{\sqrt{M}}{N} \varphi\left(\frac{j}{N}\right) \exp\left(\frac{2\pi i a j}{M}\right) \chi\left(\frac{\ell}{K}\right) \psi\left(\frac{\ell}{N} + \frac{a}{M}\right) \tilde{A}_{j\ell} \\ \tilde{A}_{j\ell} &= \sum_{m=0}^{K-1} \exp\left(\frac{2\pi i m(j - \ell M)}{N}\right) \chi\left(\frac{m}{K}\right) \end{aligned}$$

Observe that $A_{j\ell}^a$ can be nonzero only when

$$a \in \mathcal{A}, \quad \frac{j}{N} \in \text{supp } \varphi, \quad \frac{\ell}{N} + \frac{a}{M} \in \text{supp } \psi, \quad \frac{\ell}{K} \in \text{supp } \chi. \quad (2.11)$$

For j and ℓ such that condition (2.11) holds, it follows from Eq. (2.8) that

$$d\left(\frac{j - \ell M}{N}, 0\right) \geq \min\{M \cdot r, 2d(\text{supp } \chi, 0)\} \geq M \cdot r. \quad (2.12)$$

Here, note that we crucially used condition (2.9) on r , which controls the case that $\text{supp } \psi$ contains a neighborhood of ak/M and $\text{supp } \Phi^{-1}(\varphi)$ contains a neighborhood of $(a+1)k/M$ or vice versa for some $a \in \mathcal{A}$.

2. We now use Schur's bound (see for instance [Zwo12, §4.5.1]) to bound the operator of $\varphi_N B_N \psi_N$. In particular, it suffices to show that

$$\begin{aligned} \max_{0 \leq j \leq N-1} \sum_{a \in \mathcal{A}} \sum_{\ell=0}^{K-1} |A_{j\ell}^a| &\leq \tilde{g}_\chi(Nr) \\ \max_{\substack{a \in \mathcal{A} \\ 0 \leq \ell \leq N/M-1}} \sum_{j=0}^{N-1} |A_{j\ell}^a| &\leq \tilde{g}_\chi(Nr) \end{aligned}$$

for some $\tilde{g}_\chi$ that satisfies (2.10) in order to conclude $\|\varphi_N B_N \psi_N\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \tilde{g}_\chi(Nr)$.

Let $g = g_{\chi_M}$ as in Proposition 2.1 where $\chi_M(x) = \chi(Mx)$. Then for any $a \in \mathcal{A}$ and $j \in \{0, \dots, N-1\}$, (2.11) and (2.12) means that the conditions of Proposition 2.1 are

satisfied, so

$$\begin{aligned}
\sum_{\ell=0}^{K-1} |A_{j\ell}^a| &\leq \frac{\sqrt{M}}{N} \sum_{\ell=0}^{K-1} |\tilde{A}_{j\ell}| \\
&\leq \sqrt{M} \sum_{\ell: d(\frac{j-\ell M}{N}, 0) \geq Mr} g(j - \ell M) \\
&\leq 2\sqrt{M} \sum_{\ell \geq Nr} g(\ell M)
\end{aligned} \tag{2.13}$$

Similarly, for any ℓ and a , we have that

$$\begin{aligned}
\sum_{j=0}^{N-1} |A_{j\ell}^a| &\leq \frac{\sqrt{M}}{N} \sum_{j=0}^{N-1} |\tilde{A}_{j\ell}| \\
&\leq \sqrt{M} \sum_{j: d(\frac{j-\ell M}{N}, 0) \geq Mr} g(j - \ell M) \\
&\leq 2M^{3/2} \sum_{j \geq NM r} g(j)
\end{aligned} \tag{2.14}$$

3. Now we substitute in the relevant g_χ into the bounds from Step 2 to recover the desired estimates. With no extra assumptions on χ , g decays rapidly. Then it follows from Eq. (2.13) and Eq. (2.14)

$$\begin{aligned}
\sum_{a \in \mathcal{A}} \sum_{\ell=0}^{K-1} |A_{j\ell}^a| &\leq C_n \sum_{\ell \geq Nr} (\ell M)^{-n-1} \leq C_n (Nr)^{-1} \\
\sum_{j=0}^{N-1} |A_{j\ell}^a| &\leq C_n \sum_{j \geq NM r} g(j) \leq C_n (Nr)^{-1}
\end{aligned} \tag{2.15}$$

where C_n can change from line to line but depends only on n , χ , and M . On the other hand, if $\chi \in \mathcal{G}_c^s((0, 1))$, then $g \leq C \exp(-cx^{\frac{1}{s}})$. Therefore, it follows from Eq. (2.13) and Eq. (2.14) that

$$\begin{aligned}
\sum_{a \in \mathcal{A}} \sum_{\ell=0}^{K-1} |A_{j\ell}^a| &\leq C \int_{Nr}^{\infty} \exp\left(-cx^{\frac{1}{s}}\right) dx \leq \tilde{C} \exp\left(-\tilde{c}x^{\frac{1}{s}}\right) \\
\sum_{j=0}^{N-1} |A_{j\ell}^a| &\leq C \int_{NM r}^{\infty} \exp\left(-cx^{\frac{1}{s}}\right) dx \leq \tilde{C} \exp\left(-\tilde{c}x^{\frac{1}{s}}\right)
\end{aligned} \tag{2.16}$$

where the constants depend only on χ and M . Here, the sum can be seen as a lower Riemann sum, which is bounded by the corresponding integral. Therefore, by Schur's estimate, Eq. (2.15) and Eq. (2.16) yield the desired estimates on $\varphi_N B_N \psi_N$. To obtain

the estimates on the Fourier side, we simply have

$$\|\psi_N^{\mathcal{F}} B_N \varphi_N^{\mathcal{F}}\|_{\ell_N^2 \rightarrow \ell_N^2} = \|\mathcal{F}_N^* (\overline{\varphi_N B_N \psi_N})^* \mathcal{F}_N\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \|\varphi_N B_N \psi_N\|_{\ell_N^2 \rightarrow \ell_N^2},$$

which gives the identical bounds for the Fourier side. $\square$

The manifestation of this propagation estimate is clear in Fig. 1. Each time a random function over $\mathbb{Z}_N$ is propagated by B_N , it localizes in frequency space to the next Cantor subset, and similarly propagation by B_N^* yields localization in physical space to the next Cantor subset.

3. PROPAGATION OF SINGULARITIES

Now we are in a position to iteratively apply the one-step propagation estimate Proposition 2.2 to obtain bounds on propagation for a long time. First, we will derive a general estimate for long time propagation. The general estimate will then be applied to the case $\chi \in C_c^\infty((0, 1))$, and then to the Gevrey case $\chi \in \mathcal{G}_c^s((0, 1))$ for $s > 1$.

3.1. Long-time propagation. Let $N = KM$ for $K \in \mathbb{N}$ and let Φ denote the expanding map as defined in the (2.7). Define the fattened Cantor set

$$X_j := \{\Phi^{-j}(x) + y \mod 1 : x \in [0, 1], |y| \leq a_j\} = \Phi^{-j}([0, 1]) + [-a_j, a_j]. \quad (3.1)$$

The gap a_j will be adjusted later. For now, we only need to assume that $a_j > a_{j-1}/M$. Set

$$A_j = (\mathbb{1}_{X_j})_N^{\mathcal{F}}. \quad (3.2)$$

Roughly speaking, A_j is a localizing operator on the Fourier side that localizes $N \cdot a_j$ -close to the j -th discrete Cantor subset in $\mathbb{Z}_N$. We remark that the discrete Cantor subsets of $\mathbb{Z}_N$ are generally not defined in our setting since we do not assume that N is a power of M . We only assume that N is a multiple of M in order to ensure that B_N is well-defined. However, this is not a problem since the fattened Cantor sets are simply defined on the continuum and then discretized.

Note that

$$d(\Phi^{-1}(X_{j-1}), [0, 1] \setminus X_j) \leq a_j - \frac{a_{j-1}}{M}.$$

Define the gap distance by $d_1 = a_1$ and

$$d_j = N \cdot \left(a_j - \frac{a_{j-1}}{M} \right) \quad (3.3)$$

for $j \geq 2$. Therefore, by Proposition 2.2, we have estimates of the form

$$(1 - A_j) B_N A_{j-1} = R_j \quad (3.4)$$

where

$$\|R_j\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \tilde{g}_\chi(d_j).$$

provided that condition (2.9) holds, i.e.

$$0 < \frac{d_j}{N} \leq \frac{2}{M} d(\text{supp } \chi, 0) \quad (3.5)$$

We propagate the estimate (3.4) to obtain long time estimates in the following proposition.

Define the annular domain

$$\Omega_\nu := \{M^{-\nu} < |\lambda| < 5\} \subset \mathbb{C}. \quad (3.6)$$

Proposition 3.1. *Let $N = KM$ for some $K \in \mathbb{N}$. Fix a sequence*

$$\{d_j\}_{j=1}^\ell, \quad \ell \leq \frac{\log N}{\log M}$$

such that the condition (3.5) holds. Then there exists a Fourier multiplier

$$A : \ell_N^2 \rightarrow \ell_N^2$$

and families of operators

$$Z(\lambda) : \ell_N^2 \rightarrow \ell_N^2 \quad \mathcal{R}(\lambda) : \ell_N^2 \rightarrow \ell_N^2$$

that satisfy the identity

$$I = Z(\lambda)(B_N - \lambda) + \mathcal{R}(\lambda) + A \quad (3.7)$$

such that

(1) *we have the remainder estimate*

$$\|\mathcal{R}(\lambda)\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \sum_{j=0}^{\ell-1} |\lambda|^{-j-1} \tilde{g}_\chi(d_{\ell-j}) \quad (3.8)$$

where $\tilde{g}_\chi$ is the same as in Eq. (2.10).

(2) *A has rank bounded by*

$$\text{rank } A \leq 2M^{\ell\delta} \left[\frac{N}{M^\ell} + 2 \sum_{k=1}^{\ell} \frac{d_k}{M^{\ell-k}} \right] \quad (3.9)$$

Proof. We obtain the identity from iterating propagation estimate Eq. (3.4). Put

$$a_j = \frac{1}{N} \sum_{k=1}^j \frac{d_k}{M^{j-k}} \quad (3.10)$$

so that

$$d_j = N \cdot \left(a_j - \frac{a_{j-1}}{M} \right).$$

Then we can form the fattened Cantor sets X_j as in Eq. (3.1) with the corresponding Fourier localizers A_j defined in Eq. (3.2). Iterating the estimate Eq. (3.4) ℓ -times, we find

$$\begin{aligned}
(1 - A_\ell)B_N^\ell &= (1 - A_\ell)B_N A_{\ell-1} B_N^{\ell-1} + (1 - A_\ell)B_N(1 - A_{\ell-1})B_N^{\ell-1} \\
&= \sum_{j=0}^{\ell-1} (1 - A_\ell)B_N(1 - A_{\ell-1})B_N \dots (1 - A_{\ell-j})B_N A_{\ell-j-1} B_N^{\ell-j-1} \\
&= E_\ell(B_N - \lambda) \\
&\quad + \sum_{j=0}^{\ell-1} \lambda^{\ell-j-1} (1 - A_\ell)B_N(1 - A_{\ell-1})B_N \dots (1 - A_{\ell-j})B_N A_{\ell-j-1}
\end{aligned}$$

where

$$E_\ell = \sum_{j=0}^{\ell-1} \sum_{k=0}^{\ell-j-2} (1 - A_\ell)B_N(1 - A_{\ell-1})B_N \dots (1 - A_{\ell-j})B_N A_{\ell-j-1} (\lambda^k B^{\ell-j-2-k}).$$

Now, we have the desired approximate inverse identity given by

$$\begin{aligned}
I &= - \left(\sum_{0 \leq k < \ell} \lambda^{-1-k} (I - A_\ell)(B_N)^k \right) (B_N - \lambda) + \lambda^{-\ell} (1 - A_\ell)(B_N)^\ell + A_\ell \\
&= - \left(\sum_{0 \leq k < \ell} \lambda^{-1-k} (I - A_\ell)(B_N)^k \right) (B_N - \lambda) + \lambda^{-\ell} E_\ell(B - \lambda) \\
&\quad + \lambda^{-\ell} \sum_{j=0}^{\ell-1} \lambda^{\ell-j-1} (1 - A_\ell)B_N(1 - A_{\ell-1})B_N \dots (1 - A_{\ell-j})B_N A_{\ell-j-1} + A_\ell \\
&= Z_\ell(B_N - \lambda) + \mathcal{R}_\ell + A_\ell
\end{aligned} \tag{3.11}$$

where

$$\begin{aligned}
Z_\ell &= - \left(\sum_{0 \leq k < \ell} \lambda^{-1-k} (I - A_\ell)(B_N)^k \right) + \lambda^\ell E_\ell \\
\mathcal{R}_\ell &= \sum_{j=0}^{\ell-1} \lambda^{-j-1} B_N^j (1 - A_{\ell-j})B_N A_{\ell-j-1}
\end{aligned}$$

By assumption, condition (3.5) is satisfied, so Eq. (3.4) gives the desired remainder bound

$$\|\mathcal{R}_\ell\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \sum_{j=0}^{\ell-1} |\lambda|^{-j-1} \tilde{g} \left(N \cdot \left(a_{\ell-j} - \frac{a_{\ell-j-1}}{M} \right) \right).$$

To bound the rank of A_ℓ , we observe that $\Phi^{-\ell}([0, 1])$ is the union of $M^{\delta\ell}$ copies of intervals of length $M^{-\ell}$. Then from (3.1), the measure of X_ℓ can be bounded by

$$|X_\ell| \leq M^{\delta\ell}(M^{-\ell} + 2a_\ell).$$

By Eq. (3.10), we then obtain the desired bound

$$\text{rank } A_\ell \leq 2N \cdot |X_\ell| \leq 2M^{\ell\delta} \left[\frac{N}{M^\ell} + 2 \sum_{k=1}^{\ell} \frac{d_k}{M^{\ell-k}} \right].$$

Note that the above inequality holds since $M^{-\ell} \leq 1/N$ by assumption, and the factor of 2 here is merely to account for the discretization. The proposition then follows by putting

$$\begin{aligned} A &= A_\ell \\ Z(\lambda) &= Z_\ell \\ \mathcal{R}(\lambda) &= \mathcal{R}_\ell \end{aligned}$$

□

3.2. Propagation with smooth cutoff. Ultimately, we want to find the asymptotics of the eigenvalue counting function as $N \rightarrow \infty$ for a fixed $\nu > 0$. Therefore, we need some uniform control over the identity Eq. (3.7) for all sufficiently large N at a fixed ν . In particular, we choose d_j so that $\mathcal{R}_\ell$ will be uniformly small for all large N and the rank of A will be on the order N^δ .

In the case that $\chi \in C_c^\infty((0, 1))$, recall that $\tilde{g} = \tilde{g}_\chi$ is rapidly decaying. Then provided that the gaps $d_{\ell-j}$ are chosen so that (3.5) holds, the identity (3.7) holds with the remainder estimate (3.8) given by

$$\|\mathcal{R}_\ell\|_{\ell_N^2 \rightarrow \ell_N^2} \leq C_n \sum_{j=0}^{\ell-1} \lambda^{-j-1} d_{\ell-j}^n \quad (3.12)$$

for constants C_n depending only on χ . Note that for $|\lambda| < 1$, the factor λ^{-j-1} in Eq. (3.12) increases exponentially as j increases. In order for $\mathcal{R}_\ell$ to be small in norm, this growth needs to be tempered by $d_{\ell-j}$. The strategy is to choose $d_{\ell-j}$ in such a way so that the sum in Eq. (3.12) becomes exponentially decreasing in j .

Therefore, we put

$$d_{\ell-j} = \frac{L \cdot M^j}{1.5^j}. \quad (3.13)$$

We will choose $L > 0$. Meanwhile, let the time of propagation be

$$\ell = \left\lfloor \frac{\log N}{\log M} \right\rfloor. \quad (3.14)$$

Since $M \geq 2$, there exists a sufficiently large n so that

$$\frac{1.5^n}{M^{n-\nu}} < \frac{1}{2}. \quad (3.15)$$

Then choose L so that

$$L^n > 4M^\nu C_n \quad (3.16)$$

where the constant C_n is as in Eq. (3.12). Note that the choice of L depends only on χ and M . Next,

$$\frac{d_{\ell-j}}{N} \leq \frac{L}{1.5^j} \frac{M^j}{N} \leq \frac{L}{1.5^\ell}.$$

Therefore for all sufficiently large ℓ (and thus for all sufficiently large N), condition (3.5) will be satisfied. Therefore, Proposition 3.1 applies and we have the remainder estimate

$$\begin{aligned} \|\mathcal{R}_\ell\|_{\ell_N^2 \rightarrow \ell_N^2} &\leq C_n \sum_{j=0}^{\ell-1} \lambda^{-j-1} d_{\ell-j}^{-n} \\ &\leq \frac{C_n M^\nu}{L^n} \sum_{j=0}^{\ell-1} \left(\frac{2^n}{M^{(n-\nu)}} \right)^j \\ &\leq \frac{1}{2} \end{aligned} \quad (3.17)$$

Furthermore, Eq. (3.9) gives the rank bound

$$\text{rank } A \leq 2N^\delta \left(1 + 2 \sum_{k=1}^{\ell} \frac{L}{1.5^{\ell-k}} \right) \leq CN^\delta$$

where C depends only on ν , χ and M . In summary, we have the following corollary of Proposition 3.1:

Corollary 3.2. *Consider the quantum open maps given by the triple $(M, \mathcal{A}, \chi)$ and fix $\nu > 0$. For all sufficiently large $N = KM$ where $K \in \mathbb{N}$ there exists operators A , $Z(\lambda)$, and $\mathcal{R}(\lambda)$ on ℓ_N^2 as in Proposition 3.1 that satisfies the identity Eq. (3.7). Furthermore, they satisfy the remainder estimate*

$$\|\mathcal{R}(\lambda)\|_{\ell_N^2 \rightarrow \ell_N^2} \leq 1/2$$

and the rank bound

$$\text{rank } A \leq CN^\delta$$

where C does not depend on N .

3.3. Propagation with Gevrey cutoff. In the previous section, we fixed some ν and could not have extracted dependence of the rank estimates on ν since we do not know how the constants C_n behave. In particular, the dependence on ν is buried the choice of L in Eq. (3.16). However, if we assume that $\chi \in \mathcal{G}_c^s((0, 1))$, we get more explicit control over the decay of $\tilde{g}$.

For $N \geq \nu^s \geq 1$, put the time of propagation as

$$\ell = \left\lceil \frac{\log \left(\frac{N}{\nu^s} \right)}{\log M} \right\rceil \quad (3.18)$$

With $\chi \in \mathcal{G}_c^s((0, 1))$ for $s > 1$, the remainder bound (3.8) then gives

$$\|\mathcal{R}_\ell\|_{\ell^2 \rightarrow \ell^2} \leq C \sum_{j=0}^{\ell-1} \lambda^{-j-1} e^{-cd_{\ell-j}^{1/s}} = \sum_{j=0}^{\ell-1} e^{\nu(j+1) \log M - cd_{\ell-j}^{1/s}} \quad (3.19)$$

where C and c depend only on χ . Again, the remainder bound holds only if d_j/N is sufficiently small for all j according to Eq. (3.5). This condition will eventually be fulfilled using the choice of propagation time given by Eq. (3.18) and choosing N to be sufficiently large. First, we need to choose the gap distances. We see from Eq. (3.19) that we should put

$$d_{\ell-j}^{1/s} = \left(\frac{\nu \log M + \mu}{c} \right) (j+1), \quad (3.20)$$

where μ is chosen to be sufficiently large so that

$$C \sum_{j=0}^{\infty} e^{-\mu(j+1)} \leq \frac{1}{2}.$$

Then if $d_{\ell-j}$ satisfies (3.5), then the estimate Eq. (3.19) gives the desired remainder bound

$$\|\mathcal{R}_\ell\|_{\ell^2 \rightarrow \ell^2} \leq 1/2.$$

Indeed, note that $d_{\ell-j}$ takes its maximum value at $j = \ell - 1$, so (3.5) is satisfied if

$$\frac{2d(\text{supp } \chi, 0)}{M} \geq N^{-1} d_1 = N^{-1} \left(\frac{\nu \log M + \mu}{c} \left\lceil \frac{\log \left(\frac{N}{\nu^s} \right)}{\log M} \right\rceil \right)^s \quad (3.21)$$

For $1 \leq \nu^s < N$, the above is indeed satisfied for all sufficiently large N , and the threshold depends only on χ , M , and ε . Finally, to estimate the rank of A in Eq. (3.7), we see from Eq. (3.9) that

$$\text{rank } A \leq 2 \frac{N^\delta}{\nu^{s\delta}} \left[\nu^s + 2 \left(\frac{\nu \log M + \mu}{c} \right)^s \sum_{j=0}^{\ell-1} \frac{(j+1)^s}{M^j} \right] \leq C N^\delta \nu^{s(1-\delta)} \quad (3.22)$$

Note analysis above yields the following corollary of Proposition 3.1.

Corollary 3.3. *Consider the quantum open maps given by the triple $(M, \mathcal{A}, \chi)$ where $\chi \in \mathcal{G}_c^s((0, 1))$ for an $s > 1$. Then for all $1 \leq \nu^s < N$, there exists a constant $C_{\chi, M, \varepsilon}$ such that for all $N = KM > C_{\chi, M, \varepsilon}$ where $K \in \mathbb{N}$, there exists operators A , $Z(\lambda)$, and $\mathcal{R}(\lambda)$ on ℓ_N^2 as in Proposition 3.1 that satisfies the identity Eq. (3.7). Furthermore, they satisfy the remainder estimate*

$$\|\mathcal{R}(\lambda)\|_{\ell_N^2 \rightarrow \ell_N^2} \leq 1/2$$

and the rank bound

$$\text{rank } A \leq CN^\delta \nu^{s(1-\delta)}$$

where C does not depend on N or ν .

4. WEYL BOUNDS

Now we proceed to bounding the number of eigenvalues in Ω_ν as defined in (3.6). To do so, we will eventually pass to Jensen's formula from complex analysis:

Lemma 4.1. *Let $f(z)$ be a holomorphic function on a connected open set $\Omega \subset \mathbb{C}$. Let $K \subset \Omega$ be a compact subset. Suppose there exists a constant $L > 0$ and a point $z_0 \in K$ such that*

$$\sup_{z \in \Omega} |f(z)| \leq e^L, \quad |f(z_0)| \geq e^{-L}. \quad (4.1)$$

Then the number of zeros of $f(z)$ in K counted with multiplicity is bounded by

$$|\{z \in K : f(z) = 0\}| \leq CL \quad (4.2)$$

where the constant C depends only on the geometry, i.e. z_0 , Ω , and K .

See [DJ17, Lemma 4.4] for a proof of the lemma.

We want to apply Lemma 4.1 to some expression involving a factor of $\det(B_N - \lambda)$ in the region Ω_ν (defined in (3.6)) in order to count the number of eigenvalues in Ω_ν . To get the lower bound at a point in Ω_ν required in Lemma 4.1, we first modify the approximate inverse identity (3.7) as follows:

$$\begin{aligned} I &= \lambda^{-1}B_N - \lambda^{-1}(B_N - \lambda) \\ &= Z(B_N - \lambda) + \mathcal{R} + \lambda^{-1}AB_N - \lambda^{-1}A(B_N - \lambda) \\ &= (Z - \lambda^{-1}A)(B_N - \lambda) + \mathcal{R} + \lambda^{-1}AB_N \end{aligned} \quad (4.3)$$

where Z and $\mathcal{R}$ depend holomorphically on λ . In either the general cutoff setting and the Gevrey cutoff setting, Corollary 3.2 and Corollary 3.3 both give the bound $\|\mathcal{R}(\lambda)\|_{\ell_N^2 \rightarrow \ell_N^2} \leq 1/2$. Therefore, we can define

$$\begin{aligned} \mathcal{B}_N(\lambda) &:= \lambda^{-1}AB_N(I - \mathcal{R})^{-1} \\ F(\lambda) &:= \det(I - \mathcal{B}_N) \end{aligned} \quad (4.4)$$

Note that $F(\lambda)$ is holomorphic in the annulus Ω_ν . From Eq. (4.3), we have

$$\begin{aligned} F(\lambda) &= \det(I - \mathcal{R}) \det(I - \mathcal{R} - \lambda^{-1}AB_N) \\ &= \det(I - \mathcal{R}) \det(Z - \lambda^{-1}Z) \det(B_N - \lambda). \end{aligned}$$

Therefore if λ is an eigenvalue, it must also be a zero of $F(\lambda)$ considered with multiplicity. Thus it suffices to bound the number of zeros of F .

4.1. Proof of Theorem 1. By Corollary 3.2, we see that for all $\lambda \in \Omega_\nu$,

$$|F(\lambda)| \leq (\|\mathcal{B}_N\|_{\ell_N^2 \rightarrow \ell_N^2} + 1)^{\text{rank } \mathcal{B}_N} \leq (2M^\nu + 1)^{\text{rank } A} \leq e^{CN^\delta} \quad (4.5)$$

where the constant C does not depend on N . Now we want to find a lower bound on $F(\lambda)$ at a single point. Observe that at $\lambda = 4$,

$$\|\mathcal{B}_N(4)\|_{\ell_N^2 \rightarrow \ell_N^2} \leq \frac{1}{2},$$

and so

$$\begin{aligned} |F(4)|^{-1} &= |\det((I - \mathcal{B}_N(4))^{-1})| \\ &= |\det(I + \mathcal{B}_N(4)(I - \mathcal{B}_N(4))^{-1})| \\ &\leq \|I + \mathcal{B}_N(4)(I - \mathcal{B}_N(4))^{-1}\|_{\ell_N^2 \rightarrow \ell_N^2}^{\text{rank } \mathcal{B}_N} \\ &\leq e^{CN^\delta} \end{aligned} \quad (4.6)$$

where again the constant C does not depend on N . Therefore, Theorem 1 follows from applying Lemma 4.1 to Eq. (4.6) and Eq. (4.5).

4.2. Proof of Theorem 2. We modify the definition of the domain Ω_ν slightly to ensure the geometry scales correctly later. Take

$$\Omega_\nu = \{z : M^{-\nu} \leq |z| \leq e^{5\nu}\}.$$

Clearly, counting zeros of $F(\lambda)$ for $\lambda \in \Omega_\nu$ suffices, and as long as $e^{5\nu} \geq 4$, we will be able to find a lower bound at a single point of $F(\lambda)$.

By Corollary 3.3, for all $|\lambda| \geq M^{-\nu}$,

$$|F(\lambda)| \leq (2M^\nu + 1)^{\text{rank } A} \leq M^{C\nu \cdot N^\delta \nu^{s(1-\delta)}}. \quad (4.7)$$

where the constant C is independent of N and ν . For a lower bound on large λ , note that for all $\lambda \geq 4$,

$$\|\mathcal{B}(\lambda)\|_{\ell^2 \rightarrow \ell^2} \leq \frac{1}{2},$$

and thus for such λ ,

$$\begin{aligned} |F(\lambda)|^{-1} &= |\det((I - \mathcal{B}(\lambda))^{-1})| \\ &= |\det(I + \mathcal{B}(\lambda)(I - \mathcal{B}(\lambda))^{-1})| \\ &\leq M^{C\nu \cdot N^\delta \nu^{s(1-\delta)}}, \end{aligned} \tag{4.8}$$

where again the constant does not depend on N or ν . The domain Ω_ν in which we wish to upper bound the number of zeros varies with ν , and the constant in Lemma 4.1 depends on the geometry of the domain. Therefore, in order to capture the dependence on ν , consider the function

$$\tilde{F}(\omega) := F(e^\omega). \tag{4.9}$$

In particular, the number of zeros of $F(\lambda)$ for $\lambda \in \Omega_\nu$ is the same as the number of zeros of $\tilde{G}(\omega)$ for

$$\omega \in \{a + bi : a \in [-\nu \log M, 5\nu], b \in [2\pi k, 2\pi(k+1))\}$$

for any $k \in \mathbb{Z}$. Let

$$\tilde{\Omega}_\nu = \{a + bi : a \in [\nu \log M, 5\nu], b \in [0, 2\pi\nu)\}$$

Let $\mathcal{N}(\nu)$ denote the number of zeros of $F(\lambda)$ for $\lambda \in \Omega_\nu$ and let $\tilde{\mathcal{N}}(\nu)$ denote the number of zeros of $\tilde{F}(\omega)$ for $\omega \in \tilde{\Omega}_\nu$. Note that $\frac{1}{\nu}\tilde{\Omega}_\nu$ is the same domain independent of ν . Then applying Lemma 4.1 with Eq. (4.7) and Eq. (4.8), where the latter is taken at the point $e^{4\nu}$, then for all sufficiently large ν , we have the upper bound

$$\begin{aligned} \mathcal{N}(\nu) &\leq \frac{C}{\nu} \tilde{\mathcal{N}}(\nu) \\ &= \frac{C}{\nu} |\{\text{zeros of } \tilde{F}(\nu\omega) \text{ for } \omega \in \frac{1}{\nu}\tilde{\Omega}_\nu\}| \\ &\leq CN^\delta \nu^{s(1-\delta)}, \end{aligned}$$

where the constant C does not depend on N or ν . This concludes the proof.

5. NUMERICAL DISCUSSION

In this section, we look at how the Weyl upper bounds derived in this paper perform against numerical data. All plots were made using MATLAB, version R2021b.

We use the same smooth cutoff function as in [DJ17] and observe that it is 2-Gevrey. The cutoff is constructed as follows. Let

$$f(x) = \int_{-\infty}^{1.02 \cdot x - 0.01} \exp\left(-\frac{1}{t(1-t)}\right) dt.$$

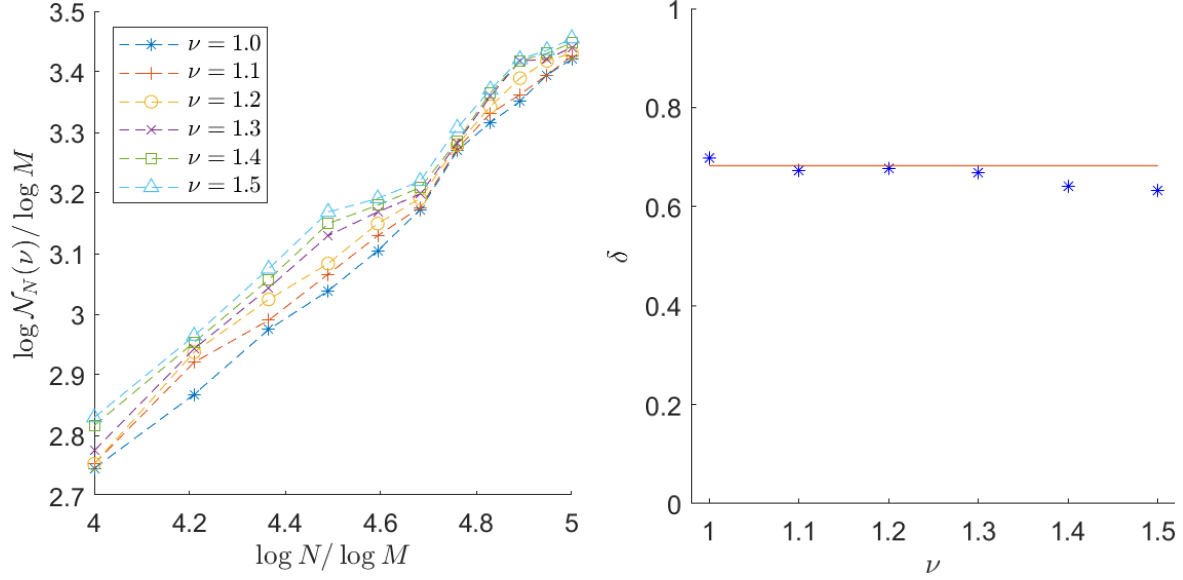


FIGURE 3. $M = 5$, $\mathcal{A} = \{1, 2, 3\}$, $\tau = 0.05$. Left: $\log \mathcal{N}_N(\nu) / \log M$ is plotted against $\log N / \log M$ for $K = N/M = \{125, 175, 225, \dots, 625\}$ given various fixed values of ν . From top to bottom, the corresponding values of ν are 1.5, 1.4, $\dots$, 1.0. Right: the corresponding slope of the linear regression of each curve is plotted against ν , and the red line is at $\delta = \log |\mathcal{A}| / \log M$.

Note that $f(x) = 0$ for $x \leq \frac{0.01}{1.02}$ and $f(x) = 1$ for $x \geq \frac{1.01}{1.02}$. Given a tightness parameter $\tau \in (0, 1/2]$, we then define the cutoff

$$\chi = f\left(\frac{x}{\tau}\right) f\left(\frac{1-x}{\tau}\right).$$

χ is 2-Gevrey and is identically 1 near the interval $[\tau, 1 - \tau]$.

The MATLAB function `eig()` was used to compute eigenvalues. We note that column j of $B_{N,\chi}$ is identically zero if $\lfloor j \cdot M/N \rfloor \in \mathcal{A}$. We cut these columns as well as the corresponding rows from the matrix $B_{N,\chi}$ to form an $K|\mathcal{A}| \times K|\mathcal{A}|$ matrix $\tilde{B}_{N,\chi}$ and compute the eigenvalues of the trimmed matrix using MATLAB. The nonzero eigenvalues of $B_{N,\chi}$ are identical to those of $\tilde{B}_{N,\chi}$, so for the sake of counting the number of eigenvalues greater than $M^{-\nu}$, using the trimmed matrix only speeds up the computation.

5.1. Dependence on N . For a fixed ν , the counting function $\mathcal{N}_N(\nu)$ is asymptotically upper bounded by N^δ as $N = KM \rightarrow \infty$. For the numerical experiment in Fig. 3 we plot $\log \mathcal{N}_N(\nu) / \log M$ against $\log N / \log M$ for several different values of ν , and for each ν , we compute the slope of the linear regression. The numerically computed

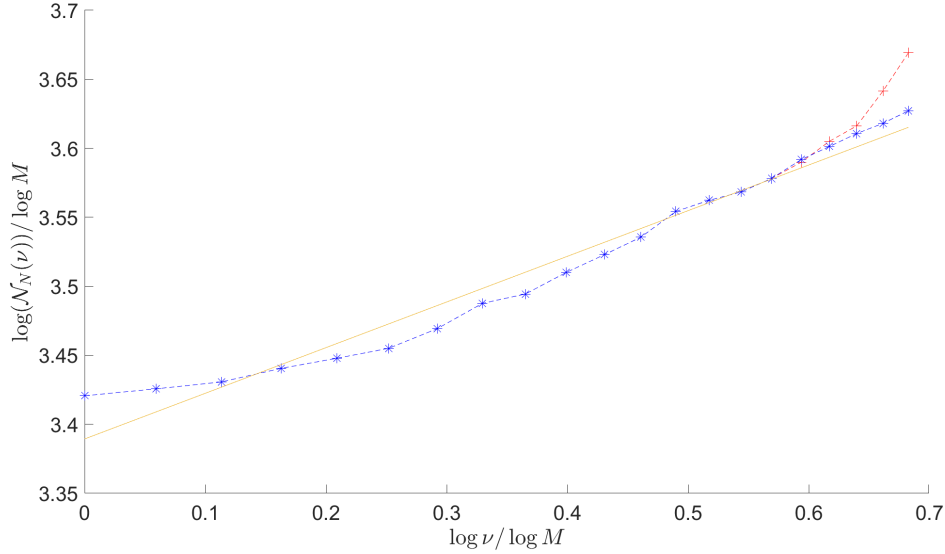


FIGURE 4. $M = 5$, $\mathcal{A} = \{1, 2, 3\}$, $\tau = 0.05$, $N = 5^5$. The blue data set is $\log \mathcal{N}_N(\nu)$ plotted against $\log \nu$ for $\nu = 1.0, 1.1, 1.2, \dots, 3.0$. The red line is the same experiment for the perturbed matrix $\tilde{B}_{N,\chi} + P$ where P is a randomly chosen matrix normalized so that $\|P\|_{\ell^2 \rightarrow \ell^2} \leq 10^{-10}$. The yellow line is the linear regression for the experiment in blue, and has a slope of 0.3308, which is fairly close to $1 - \delta \approx 0.3174$.

slopes are all fairly close to $\delta = \frac{\log |\mathcal{A}|}{\log M}$. Similar numerical results can be obtained for other quantum open baker's maps. This is in numerical agreement with the upper bound derived in this paper, and suggests that there could be matching lower bound, although no such bounds are known.

The numerics depicted in Fig. 3 is fairly stable under perturbations on the order 10^{-5} in the given range of K and ν . In particular, for each $N = K \cdot M$, we also computed the spectrum of $\tilde{B}_{N,\chi} + P$ where P is a random matrix whose entries are i.i.d. random Gaussians, and the whole matrix is normalized so that $\|P\|_{\ell^2 \rightarrow \ell^2} = 10^{-5}$. Running the same experiment as in Fig. 3 with each of the matrices perturbed by a random matrix of norm 10^{-5} , the differences in the resulting slopes are on the order 10^{-3} , which suggests a lack of strong pseudo spectral effects in the range of N and ν of concern.

5.2. Dependence on ν . Now we fix a large $N = KM$ and see how the counting function $\mathcal{N}_N(\nu)$ varies with ν . Since we have the asymptotic upper bound $\sim \nu^{s(1-\delta)}$, we plot $\log \nu / \log M$ against $\log \mathcal{N}_N(\nu) / \log M$. The numerical data for $\mathcal{N}_N(\nu)$ becomes more unstable as ν becomes large. In Fig. 4, we fix $N = 5^5$ and go to largest ν for which a perturbation on the order 10^{-10} yields no discernible difference. The line of best fit as depicted in Fig. 4 has a slope of 0.3308, which is much closer to $1 - \delta \approx 0.3174$ than

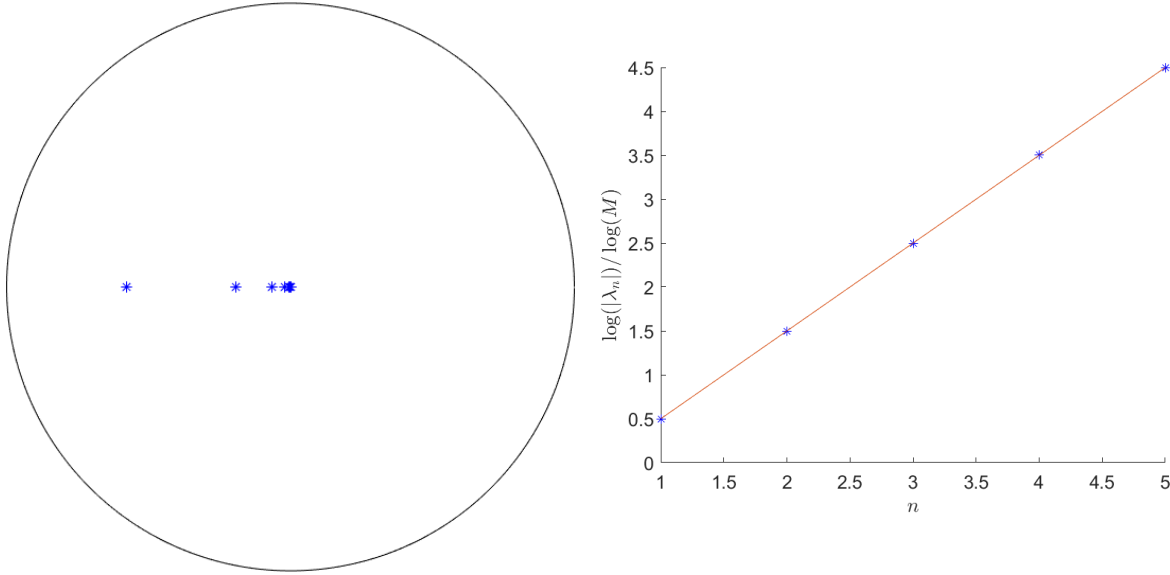


FIGURE 5. $M = 3$, $\mathcal{A} = \{1\}$, $\tau = 0.1$, $N = 3^6$. In particular, this is the degenerate case where the dimension of the trapped set is $\delta = 0$. Left: A plot of the eigenvalues in the complex unit circle. Right: Log of the magnitude of the five largest eigenvalues

to $s(1 - \delta)$. We note that the alphabet and χ we chose is such that χ is identically 1 on the Cantor set associated with the alphabet. In fact, similar experiments with the cutoff identically one near the Cantor set has similar behavior in that $\mathcal{N}_N(\nu)$ behaves like $\nu^{1-\delta}$. However, the Weyl bound $\sim N^\delta \delta^{s(1-\delta)}$ holds for all choices of alphabets and s -Gevrey cutoffs.

In the edge case where we take the alphabet to be a single point and $\delta = 0$, we should see that the magnitude of the first few eigenvalues to decrease exponentially, which would exhibit the $\mathcal{N}_N(\nu) \lesssim \nu^2$ behavior. Indeed, this is what we see in Fig. 5. We remark that essentially the same values for the top eigenvalues is obtained if we take other values of N , which makes sense in light of the fact that the upper bound we derived is independent of N for the degenerate case $\delta = 0$.

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