

Geometric flows of G_2 -structures on 3-Sasakian 7-manifolds[☆]

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ABSTRACT

A 3-Sasakian structure on a 7-manifold may be used to define two distinct Einstein metrics: the 3-Sasakian metric and the squashed Einstein metric. Both metrics are induced by nearly parallel G_2 -structures which may also be expressed in terms of the 3-Sasakian structure. Just as Einstein metrics are critical points for the Ricci flow up to rescaling, nearly parallel G_2 -structures provide natural critical points of the (rescaled) geometric flows of G_2 -structures known as the Laplacian flow and Laplacian coflow. We study each of these flows in the 3-Sasakian setting and see that their behaviour is markedly different, particularly regarding the stability of the nearly parallel G_2 -structures. We also compare the behaviour of the flows of G_2 -structures with the (rescaled) Ricci flow.

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1. Introduction

1.1. Nearly parallel G_2 -structures

A G_2 -structure on a 7-manifold is encoded by a 3-form φ satisfying a certain nondegeneracy condition, and such a 3-form determines a Riemannian metric and orientation. One of the most important types of G_2 -structure is a *nearly parallel* G_2 -structure since it defines an Einstein metric with positive scalar curvature, as well as a real Killing spinor [2,7]. Moreover,

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the cone over a 7-manifold with a nearly parallel G_2 -structure admits a conical metric with exceptional holonomy $\text{Spin}(7)$ (and so is Ricci-flat), and thus nearly parallel G_2 -structures are also important in the study of asymptotically conical and conically-singular $\text{Spin}(7)$ manifolds (cf. [14]).

Since the existence of a complete positive Einstein metric will lead to compactness of the underlying manifold by Myers theorem, it is natural to ask which compact 7-manifolds admit nearly parallel G_2 -structures. Though this general question is currently open, an infinite number of examples of such compact 7-manifolds are known, including the 7-sphere, the Aloff–Wallach spaces $N(k, l)$, the Berger space $\text{SO}(5)/\text{SO}(3)$, and the Stiefel manifold $V_{5,2}$ [7]. The largest class of 7-manifolds that are known to admit nearly parallel G_2 -structures are the 3-Sasakian 7-manifolds, which are the focus of this paper.

1.2. Geometric flows

Nearly parallel G_2 -structures are natural to study from the perspective of several geometric flows. Since a nearly parallel G_2 -structure induces a positive Einstein metric, it is natural to evolve its induced metric g by the Ricci flow:

$$\frac{\partial g}{\partial t} = -2\text{Ric}(g). \quad (1.1)$$

The induced metric will define a self-similarly *shrinking* solution to the Ricci flow, and thus a critical point after rescaling. However, a G_2 -structure contains more information than the metric (since the same metric is induced by a whole family of G_2 -structures), so it is worthwhile to examine flows of G_2 -structures relevant to nearly parallel G_2 -structures, and compare and contrast its behaviour to the Ricci flow.

Two such geometric flows of G_2 -structures which have been the most studied, and we shall examine here, are the Laplacian flow (introduced by Bryant [5]) and the Laplacian cflow (first considered in [12]¹).

1.2.1. Laplacian flow

The Laplacian flow evolves the 3-form φ defining the G_2 -structure by its Hodge Laplacian:

$$\frac{\partial \varphi}{\partial t} = \Delta_\varphi \varphi = (dd_\varphi^* + d_\varphi^* d)\varphi. \quad (1.2)$$

(Here, we emphasise the nonlinearity in the formal adjoint d_φ^* of the exterior derivative, since the metric and orientation depend on φ .) The Laplacian flow has received particular attention in the context of *closed* G_2 -structures (when $d\varphi = 0$), where it has many attractive features, particularly with regards to torsion-free G_2 -structures (when $d\varphi = 0$ and $d_\varphi^* \varphi = 0$), which define Ricci-flat metrics with holonomy contained in G_2 . For foundational results and a survey of recent developments in the Laplacian flow for closed G_2 -structures see e.g. [11, 17, 18].

A nearly parallel G_2 -structure defines a self-similarly *expanding* solution to the Laplacian flow (1.2), so can be viewed as a critical point up to rescaling. (We note the immediate difference with the Ricci flow where the induced metric was a shrinker.) A nearly parallel G_2 -structure is, however, not closed but *coclosed*: the defining 3-form φ satisfies $d_\varphi^* \varphi = 0$. Whilst it may seem potentially plausible to study coclosed G_2 -structures using the Laplacian flow (1.2), in fact it is not yet known in general whether this flow even has short time existence starting at a coclosed G_2 -structure. An example situation where it has proved instructive to use the Laplacian flow to study coclosed G_2 -structures can be found in [16].

1.2.2. Laplacian cflow

Currently the best candidate² for studying coclosed G_2 -structures is the Laplacian cflow, which evolves the closed 4-form $\psi = *\varphi$ dual to the 3-form φ defining the G_2 -structure by its Hodge Laplacian:

$$\frac{\partial \psi}{\partial t} = \Delta_\psi \psi = (d_\psi^* d + d d_\psi^*)\psi = dd_\psi^* \psi, \quad (1.3)$$

using the fact that ψ is closed. (The 4-form ψ induces the metric just like φ , but not the orientation, though an orientation can be fixed by the initial choice of G_2 -structure.) The Laplacian cflow preserves the cohomology class $[\psi]$ of ψ , where it may be viewed as the gradient flow of the Hitchin volume functional, and the induced flow of the metric g defined by ψ is

$$\frac{\partial g}{\partial t} = -2\text{Ric}(g) + Q(d\varphi), \quad (1.4)$$

where Q is a quadratic expression in $d\varphi$: see [9, 11] for details. Since Q only depends on first order information on ψ , whereas the Ricci tensor involves second order data, one may view (1.4) as a lower order perturbation of the Ricci flow (1.1).

¹ It should be noted that in [12] the opposite sign for the velocity of the Laplacian cflow is used.

² The Laplacian cflow for coclosed G_2 -structures has many attractive features analogous to the Laplacian flow for closed G_2 -structures, but with the significant difference that the analytic foundations for the Laplacian cflow are currently lacking: see [9, 11] for a discussion of the analytic issues.

However, just as for the Laplacian flow, a nearly parallel G_2 -structure defines a self-similarly *expanding* solution to the Laplacian coflow (1.3), whereas its induced metric defines a *shrinking* for Ricci flow. Hence the “lower order terms” in (1.4) drastically alter the behaviour of the metric flow in this setting.

We should also note that coclosed G_2 -structures satisfy a parametric h-principle (see [6]). Therefore, coclosed G_2 -structures exist on any (compact or non-compact) 7-manifold admitting a G_2 -structure, which just requires the 7-manifold to be oriented and spin, and so the Laplacian coflow can potentially be studied on any oriented spin 7-manifold. By contrast, it is currently not clear how restrictive the closed condition is for a G_2 -structure on a compact manifold.

1.3. 3-Sasakian 7-manifolds

A 3-Sasakian 7-manifold is a Riemannian 7-manifold M so that the metric cone over it is hyperkähler. One can use the 3-Sasakian structure to define two³ distinct nearly parallel G_2 -structures (up to scale), one of which induces the original 3-Sasakian Einstein metric on M , and the other induces the so-called squashed Einstein metric on M . This is most easily seen in the example of the 7-sphere, where the 3-Sasakian metric is the round metric, and the squashed Einstein metric is obtained by rescaling the 3-sphere fibres relative to the 4-sphere base in the Hopf fibration of the 7-sphere.

1.4. Stability

Our primary goal is to study the stability of nearly parallel G_2 -structures on 3-Sasakian 7-manifolds under the Laplacian flow and Laplacian coflow, and to compare the behaviour of these flows to the Ricci flow near their induced Einstein metrics.

For geometric flows, one is primarily interested in the question of dynamical stability of a critical point, i.e. when the flow starting near a critical point will flow back to it. An easier and weaker thing to check is linear stability: whether the critical point is stable for the linearized flow at that point. In some situations, one can infer dynamical stability from linear stability: e.g. for complete positive Einstein metrics in Ricci flow, linear stability plus an integrability assumption implies a weak form of dynamical stability [13].

In the context of nearly parallel G_2 -structures on 7-manifolds M , it was shown in [19] that if the third Betti number $b^3(M) \neq 0$, then under the Ricci flow any Einstein metric induced by a nearly parallel G_2 structure is linearly unstable and therefore dynamically unstable. As 3-Sasakian 7-manifolds M necessarily have $b^3(M) = 0$, this class of examples admitting nearly parallel G_2 -structures is particularly interesting for Ricci flow in light of this result.

In this article, when discussing stability we will always be referring to dynamical stability.

1.5. Main results

On any 3-Sasakian 7-manifold we introduce two disjoint 3-parameter families of coclosed G_2 -structures defined in terms of the 3-Sasakian structure. These families of G_2 -structures each include exactly one of the natural nearly parallel G_2 -structures we discussed above (and their rescalings). We refer the reader to §2 for details.

Our main results concern the behaviour of the Laplacian coflow, the Laplacian flow and the Ricci flow for these families of coclosed G_2 -structures and their induced metrics, which we show are preserved by the flows. (Note, in particular, that the Laplacian flow is shown to preserve the coclosed condition in this setting.)

Our most significant result is for the Laplacian coflow (1.3).

Theorem 1.1. *The Laplacian coflow starting at any initial coclosed G_2 -structure in either of our families converges, after rescaling, to the nearly parallel G_2 -structure in that family. In particular, the nearly parallel G_2 -structures are both stable within their families.*

Comparing the Laplacian flow (1.2) and Laplacian coflow (1.3), one might naively expect them to have similar behaviour as their velocities are Hodge dual. However, in our setting, we have the following, which contrasts sharply with our Laplacian coflow result.

Theorem 1.2. *Both nearly parallel G_2 -structures are unstable sources within their families under the rescaled Laplacian flow, so coclosed G_2 -structure in our families which are not nearly parallel cannot flow to either of them.*

Finally, for the Ricci flow (1.1), we have the following, which differs again from our previous two results.

Theorem 1.3. *Along the rescaled Ricci flow for our families of metrics, the 3-Sasakian metric is stable, whereas the squashed Einstein metric is a saddle point and so unstable.*

This result again shows that, whilst the Ricci flow and the induced flow of metrics (1.4) from the Laplacian coflow are closely related, their behaviour can be markedly different.

³ In fact, there are three natural nearly parallel G_2 -structures inducing the 3-Sasakian metric, but these are permuted by the symmetries in the 3-Sasakian structure. The same does *not* occur for the squashed Einstein metric.

1.6. Summary

We begin in §2 by discussing background on 3-Sasakian geometry, the nearly parallel G_2 -Structures determined by these geometries, and our geometric flow ansatz. We then study the behaviour of the Laplacian coflow in §3, the Laplacian flow in §4, and the Ricci flow in §5. To do this, we reduce the study of each rescaled flow to the analysis of a nonlinear ODE system for two functions.

2. G_2 -structures on 3-Sasakian 7-manifolds

In this section we recall some of the basics of 3-Sasakian geometry in 7 dimensions and outline its relationship to G_2 geometry. Further details on 3-Sasakian geometry can be found in [3,4]. For information about G_2 -structures, we refer the reader to [10] or [11, pp. 3–50].

2.1. 3-Sasakian 7-manifolds

We first recall the definition of a 3-Sasakian 7-manifold.

Definition 2.1. A complete Riemannian 7-manifold (M^7, g_M) is 3-Sasakian if it has an orthonormal triple of Killing fields $\{E_1, E_2, E_3\}$ satisfying $[E_i, E_j] = 2E_k$ for a cyclic permutation (i, j, k) of $(1, 2, 3)$, such that each E_i defines a Sasakian structure on (M, g_M) .

If (M, g_M) is 3-Sasakian then g_M is Einstein with positive scalar curvature equal to 42 (so M is compact) and there is a locally free action of $SU(2)$ on M whose leaf space N is a 4-dimensional orbifold. Moreover, there is a canonical metric g_N on N , which is anti-self-dual Einstein with positive scalar curvature equal to 48, such that (M, g_M) and (N, g_N) are related by an orbifold Riemannian submersion:

$$\pi : M \rightarrow N. \quad (2.1)$$

Remark 2.2. The simplest example of a 3-Sasakian 7-manifold is the 7-sphere with its constant curvature 1 metric. In this setting, (2.1) just becomes the usual Hopf fibration with $M = S^7$ and $N = S^4$, and $N = S^4$ has its constant curvature 4 metric.

The Levi-Civita connection of (N, g_N) lifts to a connection on the bundle (2.1), and so may be viewed as an $\mathfrak{su}(2)$ -valued 1-form η on M , which can be written as

$$\eta = \sum_{i=1}^3 \eta_i \otimes T_i, \quad (2.2)$$

where η_1, η_2, η_3 are 1-forms on M and $\{T_1, T_2, T_3\}$ is a basis for $\mathfrak{su}(2)$ satisfying $[T_i, T_j] = 2T_k$ for cyclic permutations (i, j, k) of $(1, 2, 3)$. The curvature of η is then an $\mathfrak{su}(2)$ -valued 2-form ω which may be written as

$$\omega = -2 \sum_{i=1}^3 \omega_i \otimes T_i \quad (2.3)$$

for 2-forms $\omega_1, \omega_2, \omega_3$ on M which are, in fact, pullbacks of orthogonal self-dual 2-forms on N since g_N is anti-self-dual Einstein. (The factor of 2 and sign are chosen for convenience.) Moreover, we have that the forms $\omega_1, \omega_2, \omega_3$ are normalized such that

$$\omega_i \wedge \omega_j = 2\delta_{ij}\pi^*\text{vol}_N. \quad (2.4)$$

For later use, we record the following equations satisfied by η and ω where, in each case, (i, j, k) are taken to be a cyclic permutation of $(1, 2, 3)$:

$$d\eta_i = -2\eta_j \wedge \eta_k - 2\omega_i, \quad (2.5)$$

$$d\omega_i = -2\eta_j \wedge \omega_k + 2\eta_k \wedge \omega_j. \quad (2.6)$$

The 3-Sasakian metric g_M on M may then be given in terms of the η_i and g_N as follows:

$$g_M = \eta_1^2 + \eta_2^2 + \eta_3^2 + \pi^*g_N. \quad (2.7)$$

We can scale g_M by any positive constant c and then c^2g_M will still be Einstein with positive scalar curvature. We may also observe the following well-known fact.

Lemma 2.3. *The metric*

$$\tilde{g}_M = \frac{1}{5}(\eta_1^2 + \eta_2^2 + \eta_3^2) + \pi^* g_N \quad (2.8)$$

is Einstein with positive scalar curvature and is known as the squashed Einstein metric on the 3-Sasakian M^7 [7,8].

Remark 2.4. The metric cone on (M, g_M) has holonomy contained in $\mathrm{Sp}(2)$, whereas the metric cone on (M, \tilde{g}_M) (once one scales \tilde{g}_M appropriately) has holonomy $\mathrm{Spin}(7)$. In the first case, the metric cone has the full holonomy $\mathrm{Sp}(2)$ if it is not flat.

2.2. Natural G_2 -structures

We recall that a G_2 -structure on a 7-manifold is determined by a 3-form φ on the manifold satisfying a certain nondegeneracy condition. Such a 3-form determines a metric g_φ and volume form vol_φ , and hence a dual 4-form $\psi = *\varphi$, where $*$ is the Hodge star determined by φ .

Given the data in (2.1), (2.2) and (2.3) above, we may now write down a natural family of G_2 -structures on a 3-Sasakian 7-manifold (M^7, g_M) as follows.

Lemma 2.5. *Given $a_1, a_2, a_3, c > 0$ and $\epsilon \in \{\pm 1\}$, if we let $\mathbf{a} = (a_1, a_2, a_3)$ then the 3-form*

$$\varphi_{\mathbf{a},c,\epsilon} = \epsilon a_1 a_2 a_3 \eta_1 \wedge \eta_2 \wedge \eta_3 - c^2(a_1 \eta_1 \wedge \omega_1 + a_2 \eta_2 \wedge \omega_2 + \epsilon a_3 \eta_3 \wedge \omega_3) \quad (2.9)$$

defines a G_2 -structure on M . Moreover, this G_2 -structure induces the following metric, volume form and dual 4-form:

$$g_{\mathbf{a},c} = a_1^2 \eta_1^2 + a_2^2 \eta_2^2 + a_3^2 \eta_3^2 + c^2 \pi^* g_N; \quad (2.10)$$

$$\mathrm{vol}_{\mathbf{a},c,\epsilon} = \epsilon a_1 a_2 a_3 c^4 \eta_1 \wedge \eta_2 \wedge \eta_3 \wedge \pi^* \mathrm{vol}_N; \quad (2.11)$$

$$\psi_{\mathbf{a},c,\epsilon} = c^4 \pi^* \mathrm{vol}_N - c^2(\epsilon a_2 a_3 \eta_2 \wedge \eta_3 \wedge \omega_1 + \epsilon a_3 a_1 \eta_3 \wedge \eta_1 \wedge \omega_2 + a_1 a_2 \eta_1 \wedge \eta_2 \wedge \omega_3). \quad (2.12)$$

Note that $g_{\mathbf{a},c}$ is independent of ϵ .

This result is an elementary consequence of the fact that $\omega_1, \omega_2, \omega_3$ are the pullbacks of self-dual 2-forms on N satisfying (2.4).

Remark 2.6. Initially, one may allow for $a_1, a_2, a_3 \in \mathbb{R} \setminus \{0\}$. However, φ and $-\varphi$ are the same G_2 -structure up to a change of orientation. Moreover, there are only two possibilities: either a_1, a_2, a_3 all have the same sign, or just two have the same sign. Therefore, we can take a_1, a_2, a_3 to be all positive and use ϵ to account for the two choices.

We now compute the exterior derivatives of $\varphi_{\mathbf{a},c,\epsilon}$ and $\psi_{\mathbf{a},c,\epsilon}$, which together encode all of the information about the torsion of the G_2 -structure.

Lemma 2.7. *Let $\varphi_{\mathbf{a},c,\epsilon}$ and $\psi_{\mathbf{a},c,\epsilon}$ be as in Lemma 2.5. Then:*

$$\begin{aligned} d\varphi_{\mathbf{a},c,\epsilon} &= 4c^2(a_1 + a_2 + \epsilon a_3)\pi^* \mathrm{vol}_N \\ &\quad - 2(\epsilon a_1 a_2 a_3 - c^2 a_1 + c^2 a_2 + \epsilon c^2 a_3)\eta_2 \wedge \eta_3 \wedge \omega_1 \\ &\quad - 2(\epsilon a_1 a_2 a_3 + c^2 a_1 - c^2 a_2 + \epsilon c^2 a_3)\eta_3 \wedge \eta_1 \wedge \omega_2 \\ &\quad - 2(\epsilon a_1 a_2 a_3 + c^2 a_1 + c^2 a_2 - \epsilon c^2 a_3)\eta_1 \wedge \eta_2 \wedge \omega_3; \\ d\psi_{\mathbf{a},c,\epsilon} &= 0. \end{aligned}$$

This result follows quickly from (2.5) and (2.6). Notice in particular that the G_2 -structures are all coclosed.

Remark 2.8. We note the following special cases of our family of G_2 -structures.

- We can always make an overall rescaling so that $c = 1$. (However, we shall see that we will require the freedom to vary the scale c along our flows.)
- Taking $a_1 = a_2 = a_3 = c = 1$ and $\epsilon = 1$ gives a coclosed G_2 -structure inducing the 3-Sasakian metric. It has been referred to as the “canonical” G_2 -structure on a 3-Sasakian 7-manifold (see e.g. [1]).
- Taking $a_1 = a_2 = a_3 = a$ and $c = 1$ gives the family of G_2 -structures considered in [15]. The subfamily where $\epsilon = 1$ was also studied earlier in [7].

2.3. Nearly parallel G_2 -structures

We recall the definition of the distinguished class of G_2 -structures that will be the focus of this paper.

Definition 2.9. A G_2 -structure on a 7-manifold M defined by a 3-form φ with dual 4-form ψ is *nearly parallel* if

$$d\varphi = \lambda\psi$$

for some non-zero constant λ . (A priori λ could be a function on M , but a short argument using $d\psi = 0$ and some representation theory shows that it must in fact be constant.)

A nearly parallel G_2 -structure φ induces an Einstein metric g_φ with positive scalar curvature. If λ is chosen so that the scalar curvature of g_φ is 42, then the cone metric $dr^2 + r^2 g_\varphi$ on $\mathbb{R}^+ \times M$ is Ricci-flat and has holonomy contained in $\text{Spin}(7)$, and φ is *strictly nearly parallel* if the holonomy of this cone metric is $\text{Spin}(7)$. (One should compare this to Remark 2.4.)

We now record the following facts, which follow immediately from (2.5) and (2.6), that show that our family of G_2 -structures contains two nearly parallel G_2 -structures (up to scale).

Lemma 2.10. Take $a_1 = a_2 = a_3 = a$ in $\varphi_{a,c,\epsilon}$.

- If $a = \frac{1}{\sqrt{5}}c$ and $\epsilon = 1$, then the resulting G_2 -structure, which we may write $c^3\varphi^{np}$ with φ^{np} independent of c , is (strictly) nearly parallel and its induced metric is $c^2\tilde{g}_M$.
- If $a = c$ and $\epsilon = -1$, then the resulting G_2 -structure, which we may write $c^3\varphi^{ts}$ where φ^{ts} is independent of c , is nearly parallel and its induced metric is c^2g_M .

Hence, within each branch (determined by $\epsilon \in \{\pm 1\}$) of our family of G_2 -structures, there is one natural critical point (up to scale) for our geometric flows. For $\epsilon = 1$, this is the strictly nearly parallel G_2 -structure φ^{np} inducing the squashed Einstein metric \tilde{g}_M on M , and for $\epsilon = -1$ this is the nearly parallel G_2 -structure φ^{ts} (where “ts” stands for 3-Sasakian) inducing the 3-Sasakian metric g_M .

Remark 2.11. If we take $a = \frac{1}{\sqrt{5}}c$ and $\epsilon = -1$ in Lemma 2.10 then we obtain a coclosed G_2 -structure which induces the Einstein metric $c^2\tilde{g}_M$, but is not nearly parallel. The same occurs when we $a = c$ and $\epsilon = 1$, but now for the Einstein metric c^2g_M : this gives a multiple of the “canonical” G_2 -structure we saw earlier (cf. Remark 2.8).

Remark 2.12. It is worth noting that, by Lemma 2.10 and [6, Examples 1.14 and 1.15], the G_2 -structures defined by $\varphi_{a,c,+1}$ and $\varphi_{a,c,-1}$ cannot be homotopic through G_2 -structures.

2.4. The ansatz

Motivated by Lemma 2.10, we will take our ansatz to be a special case of that of Lemma 2.5 where

$$a_1 = a_2 = a(t), \quad a_3 = b(t) \quad \text{and} \quad c = c(t), \quad (2.13)$$

for positive time-dependent functions $a(t), b(t), c(t)$. These then define 1-parameter families of G_2 3-forms $\varphi_\epsilon(t)$ depending on t , with induced metric $g(t)$, volume form $\text{vol}_\epsilon(t)$ and dual 4-form $\psi_\epsilon(t)$ as follows:

$$\varphi_\epsilon(t) = \epsilon a(t)^2 b(t) \eta_1 \wedge \eta_2 \wedge \eta_3 - a(t)c(t)^2 (\eta_1 \wedge \omega_1 + \eta_2 \wedge \omega_2) - \epsilon b(t)c(t)^2 \eta_3 \wedge \omega_3; \quad (2.14)$$

$$g(t) = a(t)^2 (\eta_1^2 + \eta_2^2) + b(t)^2 \eta_3^2 + c(t)^2 \pi^* g_N; \quad (2.15)$$

$$\text{vol}_\epsilon(t) = \epsilon a(t)^2 b(t) c(t)^4 \eta_1 \wedge \eta_2 \wedge \eta_3 \wedge \pi^* \text{vol}_N; \quad (2.16)$$

$$\psi_\epsilon(t) = c(t)^4 \pi^* \text{vol}_N - \epsilon a(t) b(t) c(t)^2 (\eta_2 \wedge \eta_3 \wedge \omega_1 + \eta_3 \wedge \eta_1 \wedge \omega_2) - a(t)^2 c(t)^2 \eta_1 \wedge \eta_2 \wedge \omega_3. \quad (2.17)$$

We include the subscript ϵ to emphasise the choice of branch given by $\epsilon \in \{\pm 1\}$, as we shall see different behaviour for distinct choices of ϵ , but drop the subscript for $g(t)$ since it is independent of ϵ . We shall make the restriction in (2.13) henceforth in this article.

Remark 2.13. The reader may wonder why we do not simply choose $a = b$ in (2.13) given that this holds for the nearly parallel G_2 -structures in Lemma 2.10. We shall see that the simpler ansatz when $a = b$ is not necessarily preserved along the geometric flows we consider, and so we broaden our study to consider the larger class of 1-parameter families of G_2 -structures given by the condition (2.13). One could also consider curves in the full family of G_2 -structures in Lemma 2.5, but this would be much more challenging and we already exhibit interesting behaviour within the framework provided by (2.13).

For the ansatz, we have the following simplification and slight extension of Lemma 2.7.

Lemma 2.14. Let $\varphi_\epsilon = \varphi_\epsilon(t)$ and $\psi_\epsilon = \psi_\epsilon(t)$ be given by Lemma 2.5 with the conditions in (2.13). Then:

$$\begin{aligned} d\varphi_\epsilon &= 4c^2(2a + \epsilon b)\pi^*\text{vol}_N - 2\epsilon b(a^2 + c^2)\eta_2 \wedge \eta_3 \wedge \omega_1 \\ &\quad - 2\epsilon b(a^2 + c^2)\eta_3 \wedge \eta_1 \wedge \omega_2 - 2\epsilon(a^2b + 2\epsilon ac^2 - bc^2)\eta_1 \wedge \eta_2 \wedge \omega_3; \\ d\psi_\epsilon &= 0. \end{aligned}$$

Moreover, we may write $d\varphi_\epsilon = \tau_0\psi_\epsilon + *\tau_3$ where

$$\tau_0 = \frac{4}{7a^2c^2} \left(4a(a^2 + c^2) + \epsilon b(2a^2 - c^2) \right)$$

and $\tau_3 \wedge \varphi_\epsilon = 0 = \tau_3 \wedge \psi_\epsilon$.

Proof. The formulas for $d\varphi_\epsilon$ and $d\psi_\epsilon$ are immediate from Lemma 2.7. We then compute that

$$\begin{aligned} d\varphi_\epsilon \wedge \varphi_\epsilon &= 4(\epsilon a^2bc^2(2a + \epsilon b) + 2\epsilon abc^2(a^2 + c^2) + bc^2(a^2b + 2\epsilon ac^2 - bc^2))\eta_1 \wedge \eta_2 \wedge \eta_3 \wedge \pi^*\text{vol}_N \\ &= 4\epsilon bc^2(4a(a^2 + c^2) + \epsilon b(2a^2 - c^2))\eta_1 \wedge \eta_2 \wedge \eta_3 \wedge \pi^*\text{vol}_N. \end{aligned}$$

The formula for τ_0 follows. \square

Remark 2.15. Lemma 2.14 shows that, regardless of the choice of $\epsilon \in \{\pm 1\}$, we can always choose initial conditions for our flows of G_2 -structures such that $\tau_0 = 0$ (and necessarily $\tau_3 \neq 0$), even though we are trying to flow to nearly parallel G_2 -structures, which must have $\tau_0 \neq 0$ and $\tau_3 = 0$.

3. Laplacian coflow

We start by studying the Laplacian coflow, which is arguably the natural flow for our ansatz of coclosed G_2 -structures since it manifestly preserves the coclosed condition. We recall that this flow, if it is well-posed and stays within the ansatz, is given by

$$\frac{\partial}{\partial t} \psi_\epsilon(t) = \Delta_{\psi_\epsilon(t)} \psi_\epsilon(t) = dd^*_{\psi_\epsilon(t)} \psi_\epsilon(t), \quad (3.1)$$

for the closed 4-forms $\psi_\epsilon(t)$ in (2.17).

3.1. The flow equations

Since we have that

$$dd^*_{\psi_\epsilon} \psi_\epsilon = d * d\varphi_\epsilon,$$

it is straightforward to compute the right-hand side of (3.1) from Lemma 2.14 as follows.

Lemma 3.1. The Hodge Laplacian of ψ_ϵ in (2.17) is given by:

$$\begin{aligned} \Delta_{\psi_\epsilon} \psi_\epsilon &= 8 \left(2a^2 + b^2 + 2c^2 + \frac{2\epsilon bc^2}{a} - \frac{b^2c^2}{a^2} \right) \pi^*\text{vol}_N \\ &\quad - 4 \left(b^2 + \frac{4\epsilon a^3b}{c^2} + \frac{2a^2b^2}{c^2} + \frac{2\epsilon bc^2}{a} - \frac{b^2c^2}{a^2} \right) (\eta_2 \wedge \eta_3 \wedge \omega_1 + \eta_3 \wedge \eta_1 \wedge \omega_2) \\ &\quad - 4 \left(2a^2 - b^2 + 2c^2 + \frac{4\epsilon a^3b}{c^2} + \frac{2a^2b^2}{c^2} - \frac{2\epsilon bc^2}{a} + \frac{b^2c^2}{a^2} \right) \eta_1 \wedge \eta_2 \wedge \omega_3 \end{aligned} \quad (3.2)$$

In particular, (3.2) is in the same form as (2.17) and so the Laplacian coflow (3.1) is well-defined.

Given this result and (2.17), we may write down the Laplacian coflow (3.1) as the following system of ordinary differential equations for the coefficient functions $a(t)$, $b(t)$, $c(t)$:

$$\begin{aligned}\frac{d}{dt}(c^4) &= 8 \left(2a^2 + b^2 + 2c^2 + \frac{2\epsilon bc^2}{a} - \frac{b^2 c^2}{a^2} \right); \\ \frac{d}{dt}(a^2 c^2) &= 4 \left(2a^2 - b^2 + 2c^2 + \frac{4\epsilon a^3 b}{c^2} + \frac{2a^2 b^2}{c^2} - \frac{2\epsilon bc^2}{a} + \frac{b^2 c^2}{a^2} \right); \\ \frac{d}{dt}(abc^2) &= 4 \left(\epsilon b^2 + \frac{4a^3 b}{c^2} + \frac{2\epsilon a^2 b^2}{c^2} + \frac{2bc^2}{a} - \frac{\epsilon b^2 c^2}{a^2} \right).\end{aligned}$$

We can simplify the analysis of these equations by introducing new variables as follows.

Lemma 3.2. Define

$$X = \frac{a^2}{c^2} \quad \text{and} \quad Y = \frac{ab}{c^2}$$

and introduce a new variable s by

$$\frac{ds}{dt} = \frac{1}{c^2}.$$

If we let $\dot{X} = \frac{dX}{ds}$ and $\dot{Y} = \frac{dY}{ds}$, then the Laplacian coflow equations for (2.17) imply that

$$\dot{X} = \frac{4}{X^2} \left((X+1)Y^2 + 2\epsilon(2X^2 - 2X - 1)XY - 2X^2(2X-1)(X+1) \right); \quad (3.3)$$

$$\dot{Y} = \frac{4Y}{X^2} \left(2(1-X)Y^2 + \epsilon(2X^2 - 3X - 1)Y + 2X(1-2X) \right). \quad (3.4)$$

We note that X and Y are scale-invariant quantities and that solutions to (3.3)–(3.4) give the solutions to the Laplacian coflow (3.1) up to rescaling.

3.2. Critical points and dynamics

To understand the dynamics of the flow (3.3)–(3.4), we need to study its critical points. Some straightforward calculations show the following.

Lemma 3.3. The only critical points for $X, Y > 0$ to the system (3.3)–(3.4) are:

$$X = Y = \frac{1}{5} \quad \text{and} \quad \epsilon = 1 \quad (3.5)$$

and

$$X = Y = 1 \quad \text{and} \quad \epsilon = -1. \quad (3.6)$$

Moreover, if $\epsilon = 1$ the condition $X = Y$ is preserved, but if $\epsilon = -1$ the condition $X = Y$ is not preserved except when $X = Y = 1$.

Remark 3.4. By Lemma 2.10, the critical points (3.5)–(3.6) correspond to the 4-forms ψ^{np} and ψ^{ts} dual to the nearly parallel G_2 -structures φ^{np} and φ^{ts} respectively. Thus, Lemma 3.3 shows that the only critical points for (3.1) up to rescaling are ψ^{np} and ψ^{ts} .

Before considering the general ansatz, we note that if we set $X = Y$ and $\epsilon = 1$ in (3.3)–(3.4) then we obtain:

$$\dot{X} = 4(1 - 5X).$$

Hence, \dot{X} is positive for $X < 1/5$ and negative for $X > 1/5$, which clearly shows the stability along the line $X = Y$ of the critical point (3.5). Thus ψ^{np} is stable within the restricted ansatz (2.13) with $a = b$.

Remark 3.5. Lemma 3.3 shows that the coclosed G_2 -structure with $\epsilon = 1$ inducing the 3-Sasakian metric (up to scale), as well as the one with $\epsilon = -1$ inducing the squashed Einstein metric, have no significance for the Laplacian coflow. It also shows that we need to use the ansatz (2.13) with a and b distinct (i.e. allowing $X \neq Y$) to understand the Laplacian coflow for $\epsilon = -1$.

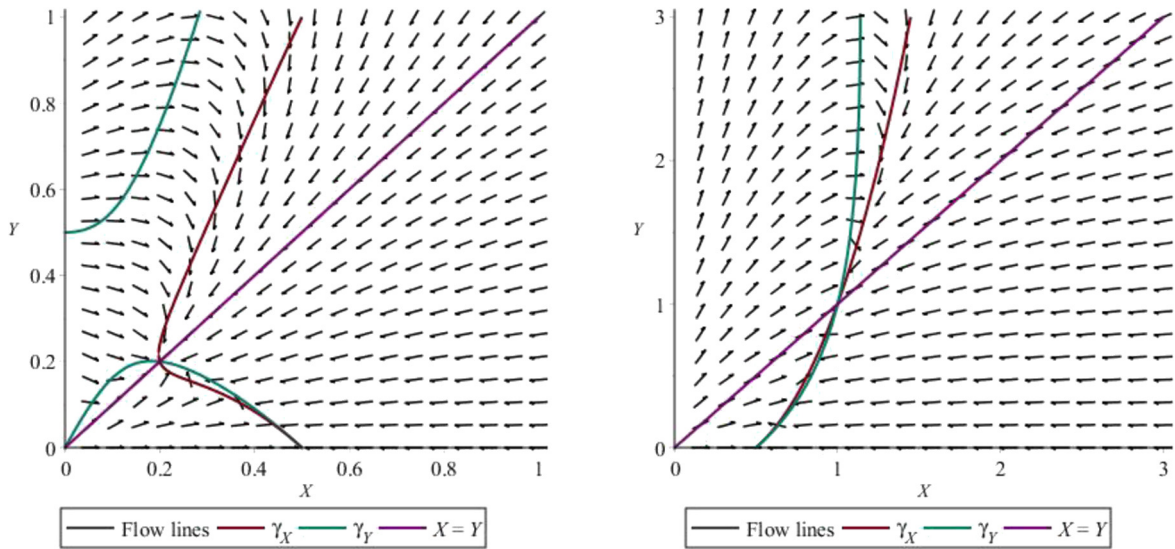


Fig. 3.1. Dynamic plots for Laplacian coflow for $\epsilon = 1$ and $\epsilon = -1$. (For interpretation of the colours in the figure(s), the reader is referred to the web version of this article.)

We provide dynamic plots of the equations (3.3)–(3.4) in Fig. 3.1 for $\epsilon = \pm 1$. In the plots, the curves γ_X and γ_Y across which \dot{X} and \dot{Y} change sign respectively are also shown, along with the line $X = Y$.

Fig. 3.1 indicates that the critical points (3.5) and (3.6), which correspond to ψ^{np} and ψ^{ts} as in Remark 3.4, are both stable. We now show that this is indeed true.

Proposition 3.6. *The 4-forms ψ^{np} and ψ^{ts} dual to the nearly parallel G_2 -structures φ^{np} and φ^{ts} are stable sinks under the Laplacian coflow (3.1), after rescaling.*

Proof. We study the linearization of the flow equations (3.3)–(3.4) at the critical points (3.5) and (3.6) to determine their stability.

At $X = Y = 1/5$ and $\epsilon = 1$, the linearized equations are

$$\dot{X} = -\frac{76}{5}X - \frac{24}{5}Y \quad \text{and} \quad \dot{Y} = -\frac{12}{5}X - \frac{88}{5}Y.$$

(Note that $X = Y$ is preserved by the above system as expected.) The associated 2×2 matrix of coefficients of X, Y in the above equations has two negative eigenvalues ($-64/5$ and -20) and so (3.5) is a stable critical point.

Similarly, at $X = Y = 1$ and $\epsilon = -1$, the linearized equations are

$$\dot{X} = -76X + 24Y \quad \text{and} \quad \dot{Y} = -36X + 8Y,$$

noting that this time $X = Y$ is not preserved. Here, the matrix one obtains again has two negative eigenvalues, which are -4 and -64 , so the critical point (3.6) is stable. \square

Remark 3.7. We see from (3.3)–(3.4) that if we allow $X = 0$ or $Y = 0$ then there are additional critical points:

$$(X, Y) = \left(\frac{1}{2}, 0\right) \quad \text{for } \epsilon = \pm 1.$$

We can see these critical points in Fig. 3.1. We can also consider $(X, Y) = (0, 0)$ to be a degenerate critical point, even though the equations (3.3)–(3.4) are not defined there. We can understand these additional critical points geometrically as follows.

Recall the fibration (2.1) of M^7 over N^4 . The point $(0, 0)$ corresponds to sending the 3-dimensional fibres of (2.1) to zero size (since $a = b = 0$), and so M has collapsed to N (or a point). In this setting, the 4-form ψ_ϵ reduces to simply the volume form of N (or zero if the collapse is to a point).

If we instead view the fibres of (2.1) as circle bundles over S^2 (where E_3 is tangent to the circle direction in the notation of Definition 2.1), then at $(1/2, 0)$ the circle fibres have now collapsed (as $b = 0$ in (2.13)). Since $a \neq 0$ there M^7 has collapsed to a 6-manifold Z which is a 2-sphere bundle over N . This 6-manifold Z is the twistor space of N , and at this critical point it will be endowed with its nearly Kähler metric g_Z , which is an Einstein metric on Z with positive scalar curvature. This Einstein metric g_Z is not Kähler (unlike the standard choice of metric on the twistor space), but instead is related to G_2 geometry as the metric cone over a 6-dimensional nearly Kähler manifold will have holonomy G_2 .

3.3. Long-time behaviour

The plots in Fig. 3.1 suggest that, within our ansatz, any initial condition flows to the unique (up to scale) nearly parallel G_2 -structure in the family. We now show that this is indeed the case. For the statement, as in Remark 3.4, we denote by ψ^{np} and ψ^{ts} the duals of the nearly parallel G_2 -structures φ^{np} and φ^{ts} defined in Lemma 2.10, and recall that they induce the squashed Einstein metric and 3-Sasakian metric respectively.

Theorem 3.8. *Let $\psi_\epsilon = \psi_\epsilon(0)$ be a closed 4-form as in (2.17) dual to a G_2 -structure. The solution to the Laplacian coflow (3.1) starting at ψ_ϵ converges, after rescaling, to ψ^{np} if $\epsilon = 1$ and to ψ^{ts} if $\epsilon = -1$, which are the only critical points of the rescaled flow. In particular, the nearly parallel G_2 -structures given by ψ^{np} and ψ^{ts} are stable for (3.1) after rescaling.*

Theorem 3.8 gives Theorem 1.1 in the introduction. The proof of Theorem 3.8 is quite lengthy, so we break it up into several smaller results.

3.3.1. Strategy

Our aim is to prove that, given any initial condition, the flow (3.3)–(3.4) for $(X(s), Y(s))$ exists for all s and converges to the critical point with $X > 0$ and $Y > 0$. From (3.3)–(3.4), long time existence will be guaranteed as long as X, Y remain bounded and X remains bounded away from 0. Moreover, periodic orbits are not possible as we know the Laplacian coflow is a gradient flow. Hence, long time existence will imply convergence to the critical point with $X, Y > 0$ if we additionally know that Y remains bounded away from 0.

To deduce the result for the Laplacian coflow we observe that the evolution equation for c^2 , which determines the parameter s by Lemma 3.2, is:

$$\frac{d}{dt}(c^2) = 8 \left(2X + \frac{Y^2}{X} + 2 + 2\epsilon \frac{Y}{X} - \frac{Y^2}{X^2} \right),$$

so c^2 cannot blow up in finite time because X, Y are bounded and bounded away from zero. In fact, c^2 grows at most linearly, so we can integrate $\frac{ds}{dt} = \frac{1}{c^2}$ to find s as a function of the Laplacian coflow parameter t .

Throughout the proof, we denote the curves where \dot{X} and \dot{Y} change sign, respectively, by γ_X and γ_Y as in Fig. 3.1.

3.3.2. Bounds on X

We start by looking at the behaviour of X .

Lemma 3.9. *The function $X(s)$ is bounded away from zero and can only diverge in the region to the left of γ_X .*

Proof. We first see that X is decreasing whenever it is to the right of γ_X and so X will remain bounded by its initial condition in this region. We then see that X is increasing in the region to the left of γ_X , which is the region containing $X = 0$, and so X is bounded away from zero in this region by its initial condition. \square

Since γ_X meets $Y = X$ (at the critical point), the unbounded part of γ_X lies above the line $Y = X$, i.e. where $Y > X$. Therefore, it suffices to show that Y remains bounded in the region to the left of γ_X where additionally $Y > X$ to deduce that X is bounded everywhere.

3.3.3. Bounds on Y : $\epsilon = 1$

Given our earlier discussion, we now turn to showing that Y remains bounded and bounded away from zero. We start with the case $\epsilon = 1$.

Lemma 3.10. *When $\epsilon = 1$, $Y(s)$ can only diverge in the region above the upper part of γ_Y and can only tend to zero in the region above the lower part of γ_Y but below the line $Y = X$.*

Proof. In this case, $\dot{Y} < 0$ in the region above the lower part of γ_Y and below or to the right of the upper part of γ_Y , so Y remains bounded by its initial value in this region. Note that γ_Y meets the line $X = 0$ at $Y = 0$ and $Y = 1/2$. Hence, in the same region we just considered, we are either to the left of γ_X and so $\dot{X} > 0$, which means we cannot reach $(0, 0)$, or we are to the right of γ_X but also above γ_Y . In this latter region, if we are above the line $Y = X$ we cannot cross it by Lemma 3.3, and so Y remains bounded away from zero here.

We then notice that $\dot{Y} > 0$ in the region bounded by the lower part of γ_Y , in which Y is bounded, and so Y is additionally bounded away from 0 here. \square

We now have our crucial observation for the case $\epsilon = 1$.

Lemma 3.11. When $\epsilon = 1$, $Y(s)$ is decreasing and hence bounded in the region where $Y > X$ and $X \geq 1$.

Proof. In this setting, we can rewrite (3.4) in the following manner:

$$\dot{Y} = \frac{4Y}{X^2} (2(1-X)(Y-X)Y - XY - Y + 2X(1-2X)).$$

Hence, $\dot{Y} < 0$ when $Y > X$ and $X \geq 1$. Therefore, Y will be bounded in this region. \square

Since the only part of the quadrant with $X, Y > 0$ where Y can become unbounded is where $Y > X$ by Lemma 3.10, we deduce that Y can only become unbounded, when $\epsilon = 1$, if X remains in $(0, 1)$. We now show that this is impossible.

Proposition 3.12. For $\epsilon = 1$, there are no solutions $(X(s), Y(s))$ above the upper part of γ_Y with $X(s) \in (0, 1)$ and $Y(s)$ unbounded.

Proof. Suppose not and that we have such a solution. Note that there is a finite $\bar{Y} > 0$ such that for all $X \in (0, 1)$,

$$Y^2 + 2(2X^2 - 2X - 1)XY - 2X^2(2X - 1)(X + 1) > 0 \quad \text{and} \quad (2X^2 - 3X - 1)Y + 2X(1 - 2X) < 0 \quad (3.7)$$

for all $Y > \bar{Y}$. Since Y is increasing and unbounded in the region under consideration, we may assume that $Y(s) > \bar{Y}$. Then, comparing (3.7) and (3.3)–(3.4), we see that

$$\dot{X} \geq \frac{4}{X}Y^2 \quad \text{and} \quad \dot{Y} \leq \frac{8(1-X)}{X^2}Y^3.$$

Hence, we see that

$$\frac{dY}{dX} \leq \frac{2(1-X)}{X}Y.$$

Grönwall's inequality then shows that there are constants C_0, C_1 depending only on the initial conditions so that Y is bounded by $C_0X^2e^{-2X} + C_1$. Since we assumed $X(s) \in (0, 1)$, this is a contradiction. \square

Our results so far show that, when $\epsilon = 1$, both X and Y are bounded and that X is bounded away from zero. To complete the proof of Theorem 3.8 in the $\epsilon = 1$ case we therefore only need the following.

Lemma 3.13. When $\epsilon = 1$, $Y(s)$ is bounded away from zero.

Proof. By Remark 3.7 and Lemma 3.10, we need only consider points near $(X, Y) = (1/2, 0)$ above γ_Y . Linearizing around the critical point $(1/2, 0)$, so writing $X = \frac{1}{2} + \delta X_1$ and $Y = \delta Y_1$ for δ small, we find that the linear term in δ gives

$$\dot{X}_1 + \dot{Y}_1 = -24(X_1 + Y_1) \quad \text{and} \quad \dot{Y}_1 = 0.$$

These equations are degenerate along the line $X_1 + Y_1 = 0$, which is tangent to γ_Y at $(1/2, 0)$. However, γ_Y lies above this line for all points near $(1/2, 0)$, so we may restrict to the region where $X_1 + Y_1 > 0$. Therefore, to leading order, Y_1 remains a non-zero constant and $X_1 + Y_1$ decreases with an exponential rate. Hence, Y will be bounded away from 0, as claimed. \square

We can now put our results together so far.

Proof of Theorem 3.8 for $\epsilon = 1$. Lemma 3.9 shows that X is bounded away from 0 and bounded above if Y is bounded. Lemma 3.10, Lemma 3.11 and Proposition 3.12 show that Y is bounded and hence X is bounded. Finally, Lemma 3.13 shows that Y is bounded away from 0. The observations in §3.3.1 then give the result. \square

3.3.4. Bounds on Y : $\epsilon = -1$

Having proved Theorem 3.8 for $\epsilon = 1$ now move on to the case where $\epsilon = -1$. The arguments here are similar to the $\epsilon = 1$ case, but often easier.

Lemma 3.14. When $\epsilon = -1$, $Y(s)$ can only diverge in the region to the left of γ_Y and above the line $Y = X$, and can only tend to 0 to the right of γ_Y but below the line $Y = X$.

Proof. These are elementary observations given that Y is decreasing to the right of γ_Y and increasing to the left of γ_Y . \square

We now see that the evolution equation (3.4) for Y has a useful feature when $\epsilon = -1$.

Lemma 3.15. When $\epsilon = -1$, there exists a least $\bar{X} \in (1, 2)$ such that $X \leq \bar{X}$ whenever (X, Y) is on γ_Y . Hence, $\dot{Y} \leq 0$ and thus Y is bounded whenever $X \geq \bar{X}$.

Proof. This is an elementary calculation, showing that X takes a maximum value on the curve γ_Y . \square

We deduce from Lemma 3.14 and Lemma 3.15 that the only way Y can become unbounded when $\epsilon = -1$ is if X remains in the interval $(0, \bar{X})$. We now show that this is not possible.

Proposition 3.16. Recall \bar{X} from Lemma 3.15. For $\epsilon = -1$, there are no solutions $(X(s), Y(s))$ to the left of γ_Y and above $Y = X$ with $X(s) \in (0, \bar{X})$ and $Y(s)$ unbounded.

Proof. We suppose, for a contradiction, that there is such a solution. Suppose that the solution enters the part of the region where $X \geq 1$. Then X is strictly increasing, so $X > 1$ for all subsequent times. However, the line γ_Y asymptotes to $X = 1$ as $Y \rightarrow \infty$, so we must eventually have that X is decreasing, which is a contradiction.

We deduce that $X(s) \in (0, 1)$ for all s . We note that there is a finite $\bar{Y} > 0$ such that

$$Y^2 - 2(2X^2 - 2X - 1)XY - 2X^2(2X - 1)(X + 1) > 0$$

$$\text{and } (2X^2 - 3X - 1)Y - 2X(1 - 2X) > 0 \quad (3.8)$$

for all $X \in (0, 1)$ and $Y > \bar{Y}$. Since Y is increasing in the region we are studying and we are assuming it is unbounded, we may restrict to the case where $Y(s) > \bar{Y}$. Comparing (3.8) to (3.3)–(3.4) yields the differential inequalities

$$\dot{X} \geq \frac{4}{X}Y^2 \quad \text{and} \quad \dot{Y} \leq \frac{8(2 - X)Y^3}{X^2}.$$

We deduce that

$$\frac{dY}{dX} \leq \frac{2(2 - X)}{X}Y$$

and so, by Grönwall's inequality, we have that Y is bounded by $X^4 e^{-2X}$ (up to multiplicative factors and additive constants depending only on the initial conditions). Since $X \in (0, 1)$, this forces our required contradiction. \square

To complete the proof in the $\epsilon = -1$ we now only need to show that Y stays away from 0.

Lemma 3.17. When $\epsilon = -1$, $Y(s)$ is bounded away from zero.

Proof. The proof is entirely analogous to that of Lemma 3.13. Remark 3.7 and Lemma 3.14 show that we may restrict attention to points near $(1/2, 0)$ to the right of γ_Y . Writing $X = \frac{1}{2} + \delta X_1$ and $Y = \delta Y_1$ for δ small, we find that the linear term in δ gives

$$\dot{X}_1 - \dot{Y}_1 = -24(X_1 - Y_1) \quad \text{and} \quad \dot{Y}_1 = 0.$$

Note that the line $X_1 = Y_1$ is tangent to γ_Y at $(1/2, 0)$ and that γ_Y lies above this line. Therefore, we need only consider $X_1 - Y_1 > 0$ and see that, to leading order, Y_1 remains a non-zero constant whilst $X_1 - Y_1$ exponentially decreases. Hence, Y will be bounded away from 0. \square

We may now complete the proof of Theorem 3.8.

Proof of Theorem 3.8 for $\epsilon = -1$. We first see that Lemma 3.9 shows that X is bounded away from 0 and bounded if Y is bounded. Lemmas 3.14 and 3.15, together with Proposition 3.16, show that Y is bounded, and thus X is also bounded. Finally, Lemma 3.17 shows that Y is bounded away from zero, which completes the proof by the discussion in §3.3.1. \square

4. Laplacian flow

We now consider the Laplacian flow for our family of G_2 -structures in (2.14). We recall that this flow, if it is well-posed and stays within the ansatz, is given by

$$\frac{\partial}{\partial t} \varphi_\epsilon(t) = \Delta_{\varphi_\epsilon(t)} \varphi_\epsilon(t) = d_{\varphi(t)}^* d \varphi_\epsilon(t), \quad (4.1)$$

for the coclosed 3-forms $\varphi_\epsilon(t)$ in (2.14). Here we have to be particularly mindful that the coclosed condition may not be preserved by the flow, let alone the ansatz.

4.1. The flow equations

We first observe that the Hodge Laplacian of the 3-form defining the G_2 -structure follows immediately from that of the 4-form by taking the Hodge star:

$$dd^*_{\varphi_\epsilon} \varphi_\epsilon = *\Delta_{\psi_\epsilon} \psi_\epsilon.$$

To compute the Hodge star we use the following relations:

$$\begin{aligned} *\varphi_\epsilon \pi^* \text{vol}_N &= \frac{\epsilon a^2 b}{c^4} \eta_1 \wedge \eta_2 \wedge \eta_3; \\ *\varphi_\epsilon (\eta_2 \wedge \eta_3 \wedge \omega_1 + \eta_3 \wedge \eta_1 \wedge \omega_2) &= \frac{\epsilon}{b} (\eta_1 \wedge \omega_1 + \eta_2 \wedge \omega_2); \\ *\varphi_\epsilon (\eta_1 \wedge \eta_2 \wedge \omega_1) &= \frac{\epsilon b}{a^2} \eta_3 \wedge \omega_3. \end{aligned}$$

We can now use Lemma 3.1 to find an expression for the Hodge Laplacian of φ_ϵ , which we need to consider the Laplacian flow (4.1).

Lemma 4.1. *The Hodge Laplacian of φ_ϵ in (2.14) is given by:*

$$\begin{aligned} \Delta_{\varphi_\epsilon} \varphi_\epsilon &= \frac{8\epsilon a^2 b}{c^2} \left(\frac{2a^2}{c^2} + \frac{b^2}{c^2} + 2 + \frac{2\epsilon b}{a} - \frac{b^2}{a^2} \right) \eta_1 \wedge \eta_2 \wedge \eta_3 \\ &\quad - 4 \left(\epsilon b + \frac{4a^3}{c^2} + \frac{2\epsilon a^2 b}{c^2} + \frac{2c^2}{a} - \frac{\epsilon b c^2}{a^2} \right) (\eta_1 \wedge \omega_1 + \eta_2 \wedge \omega_2) \\ &\quad - 4\epsilon b \left(2 - \frac{b^2}{a^2} + \frac{2c^2}{a^2} + \frac{4\epsilon ab}{c^2} + \frac{2b^2}{c^2} - \frac{2\epsilon b c^2}{a^3} + \frac{b^2 c^2}{a^4} \right) \eta_3 \wedge \omega_3 \end{aligned} \quad (4.2)$$

In particular, (4.2) is coclosed and in the same form as (2.14), so the Laplacian flow (4.1) is well-defined.

Remark 4.2. We observe that Lemma 4.1 shows that the coclosed condition on the G_2 -structure is preserved along the Laplacian flow in this situation.

As a consequence of Lemma 4.1 and (2.14), we can write down the Laplacian flow (4.1) as a system of ODEs for $a(t)$, $b(t)$, $c(t)$:

$$\begin{aligned} \frac{d}{dt}(a^2 b) &= \frac{8a^2 b}{c^2} \left(\frac{2a^2}{c^2} + \frac{b^2}{c^2} + 2 + \frac{2\epsilon b}{a} - \frac{b^2}{a^2} \right); \\ \frac{d}{dt}(a c^2) &= 4 \left(\epsilon b + \frac{4a^3}{c^2} + \frac{2\epsilon a^2 b}{c^2} + \frac{2c^2}{a} - \frac{\epsilon b c^2}{a^2} \right); \\ \frac{d}{dt}(b c^2) &= 4b \left(2 - \frac{b^2}{a^2} + \frac{2c^2}{a^2} + \frac{4\epsilon ab}{c^2} + \frac{2b^2}{c^2} - \frac{2\epsilon b c^2}{a^3} + \frac{b^2 c^2}{a^4} \right). \end{aligned}$$

Just as for the Laplacian coflow, it turns out that the Laplacian flow is easier to analyze if we introduce scale-invariant quantities and a new time parameter. The resulting equations we obtain then describe the rescaled Laplacian flow.

Lemma 4.3. *Define*

$$X = \frac{a^2}{c^2} \quad \text{and} \quad Y = \frac{ab}{c^2}$$

and introduce a new variable s by

$$\frac{ds}{dt} = \frac{1}{c^2}.$$

If we let $\dot{X} = \frac{dX}{ds}$ and $\dot{Y} = \frac{dY}{ds}$, then the Laplacian flow equations (4.1) imply that

$$\dot{X} = \frac{4}{X^2} \left(2X^2(2X-1)(X+1) + 2\epsilon(1+2X-2X^2)XY - (1+X)Y^2 \right); \quad (4.3)$$

$$\dot{Y} = \frac{4Y}{X^2} \left(2X(2X-1) + \epsilon(1+3X-2X^2)Y + 2(X-1)Y^2 \right). \quad (4.4)$$

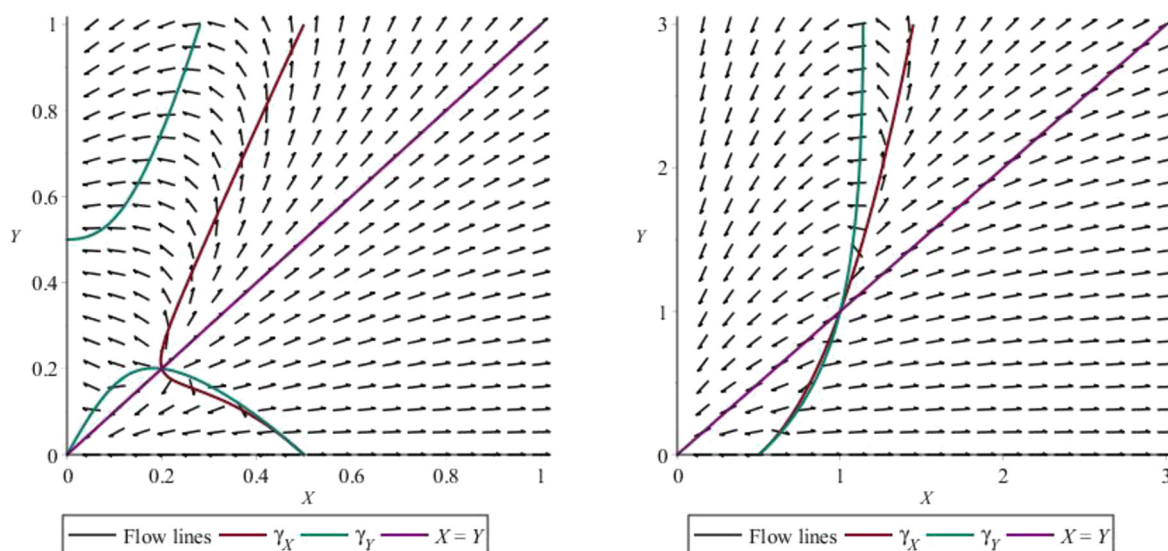


Fig. 4.1. Dynamic plots for Laplacian flow for $\epsilon = 1$ and $\epsilon = -1$.

4.2. Critical points and dynamics

We observe that (4.3)–(4.4) are just the negative of the equations (3.3)–(3.4) arising from the Laplacian coflow. We therefore have the following.

Lemma 4.4. *The only critical points for $X, Y > 0$ to (4.3)–(4.4) are:*

$$X = Y = \frac{1}{5} \quad \text{and} \quad \epsilon = 1$$

and

$$X = Y = 1 \quad \text{and} \quad \epsilon = -1,$$

and the condition $X = Y$ is preserved for $\epsilon = 1$, but not preserved for $\epsilon = -1$ except when $X = Y = 1$.

The observation that the Laplacian flow equations (4.3)–(4.4) are the negative of the Laplacian coflow equations also implies the following result, based on our stability analysis for the Laplacian coflow, which gives Theorem 1.2 in the Introduction.

Theorem 4.5. *The only critical points of the Laplacian flow (4.1), after rescaling, are the nearly parallel G_2 structures φ^{ts} and φ^{np} inducing the 3-Sasakian and squashed Einstein metrics. Both critical points are unstable sources under the rescaled Laplacian flow.*

For completeness, we provide the dynamic plots in Fig. 4.1 for the Laplacian flow (4.1) for our ansatz with $\epsilon = 1$ and $\epsilon = -1$. We again indicate the curves γ_X and γ_Y where \dot{X} and \dot{Y} change sign, respectively. Of course, the dynamics are simply the opposite of those which appear in the Laplacian coflow plot in Fig. 3.1.

Remark 4.6. Fig. 4.1 shows that several different behaviours are possible for the rescaled Laplacian flow when the initial condition is close to a nearly parallel G_2 -structure. One possibility is flowing to the origin, which corresponds again to the 7-manifold M collapsing to the 4-manifold N through the fibration (2.1). Notice that the 3-dimensional fibres of (2.1) are calibrated by the G_2 -structure φ , i.e. φ restricts to be the volume form on the fibres, and so are associative by definition. The fact that $d\varphi \neq 0$ means that the volume of any compact associative 3-fold is not necessarily topologically determined, unlike for closed G_2 -structures.

5. Ricci flow

In this section we compare the results we have obtained for the Laplacian flow and coflow of our G_2 -structures with the behaviour of the induced metric of the G_2 -structures under the Ricci flow. We recall that this comparison is useful because the nearly parallel G_2 -structures, which are critical points of the rescaled Laplacian flow and coflow that we studied earlier,

induce Einstein metrics which are then critical points of the Ricci flow up to scale. It is also of interest because if we use the Laplacian coflow for coclosed G_2 -structures (1.3), then the induced flow (1.4) on the induced metric of the G_2 -structures is the Ricci flow plus lower order terms determined algebraically by the torsion of the G_2 -structure.

5.1. The flow equations

We wish to study the Ricci flow for our ansatz

$$\frac{\partial}{\partial t} g(t) = -2\text{Ric}(t), \quad (5.1)$$

if it exists. From general theory the Ricci flow will have short time existence starting from our metric ansatz (2.15), though it is not immediately obvious that the ansatz will be preserved.

We begin by computing the Ricci curvature of $g(t)$ from (2.15). (Note that $g(t)$ is independent of $\epsilon \in \{\pm 1\}$.)

Lemma 5.1. *Let $g(t)$ be as in (2.15). The Ricci curvature $\text{Ric}(t)$ of $g(t)$ satisfies*

$$\text{Ric} = 2 \left(2 - \frac{b^2}{a^2} + \frac{2a^4}{c^4} \right) (\eta_1^2 + \eta_2^2) + 2 \left(\frac{b^4}{a^4} + \frac{2b^4}{c^4} \right) \eta_3^2 + 2 \left(6 - \frac{2a^2 + b^2}{c^2} \right) \pi^* g_N.$$

In particular, $\text{Ric}(t)$ is of the same form as (2.15), and so the Ricci flow (5.1) is well-defined.

Proof. Our approach is to use the O'Neill formulas for Riemannian submersions [20].

We let ∇ denote the Levi-Civita connection of the 3-Sasakian metric g_M . Recall the orthonormal Killing fields E_1, E_2, E_3 in Definition 2.1 and let B_0, B_1, B_2, B_3 be local orthonormal vector fields on M which are horizontal for the Riemannian submersion (2.1). It is then straightforward to compute that

$$\nabla_{E_i} E_j = \sum_{k=1}^3 \epsilon_{ijk} E_k, \quad \nabla_{E_i} B_l = 0, \quad \nabla_{B_l} E_i = - \sum_{m=0}^4 \sigma_{ilm} B_m, \quad (\nabla_{B_l} B_m)^\vee = \sum_{i=1}^3 \sigma_{ilm} E_i,$$

where $()^\vee$ indicates the vertical projection with respect to (2.1), ϵ_{ijk} is the sign of the permutation (i, j, k) of $(1, 2, 3)$, and σ_{ilm} is skew-symmetric in l, m satisfying $\sigma_{i0j} = 1$ and $\sigma_{ijk} = \epsilon_{ijk}$ for $i, j, k \in \{1, 2, 3\}$. In particular, we notice that the fibres of the Riemannian submersion (2.1) are totally geodesic, so the O'Neill tensor often denoted T vanishes, and that the other O'Neill tensor, often called A , is horizontally divergence-free.

We then let $\tilde{\nabla}$ denote the Levi-Civita connection of $g = g(t)$ and let

$$\tilde{E}_1 = \frac{E_1}{a}, \quad \tilde{E}_2 = \frac{E_2}{a}, \quad \tilde{E}_3 = \frac{E_3}{b}, \quad \tilde{B}_l = \frac{B_l}{c}.$$

We may then compute that the quantities we need to complete our computation are:

$$\begin{aligned} [\tilde{E}_2, \tilde{E}_3] &= \frac{2}{b} \tilde{E}_1, \quad [\tilde{E}_3, \tilde{E}_1] = \frac{2}{b} \tilde{E}_2, \quad [\tilde{E}_1, \tilde{E}_2] = \frac{2b}{a^2} \tilde{E}_3, \\ \tilde{\nabla}_{\tilde{E}_2} \tilde{E}_3 &= \frac{b}{a^2} \tilde{E}_1, \quad \tilde{\nabla}_{\tilde{E}_3} \tilde{E}_1 = \frac{2a^2 - b^2}{a^2 b} \tilde{E}_2, \quad \tilde{\nabla}_{\tilde{E}_1} \tilde{E}_2 = \frac{b}{a^2} \tilde{E}_3, \quad \tilde{\nabla}_{\tilde{E}_i} \tilde{B}_l = 0, \\ \tilde{\nabla}_{\tilde{B}_l} \tilde{E}_1 &= -\frac{a}{c^2} \sum_{m=0}^4 \sigma_{1lm} \tilde{B}_m, \quad \tilde{\nabla}_{\tilde{B}_l} \tilde{E}_2 = -\frac{a}{c^2} \sum_{m=0}^4 \sigma_{2lm} \tilde{B}_m, \quad \tilde{\nabla}_{\tilde{B}_l} \tilde{E}_3 = -\frac{b}{c^2} \sum_{m=0}^4 \sigma_{3lm} \tilde{B}_m. \end{aligned}$$

Using the fact that the metric on N is Einstein with scalar curvature 48 and the O'Neill formulas, we see that the Ricci tensor of g is diagonal and satisfies:

$$\begin{aligned} \text{Ric}(\tilde{E}_1, \tilde{E}_1) &= \text{Ric}(\tilde{E}_2, \tilde{E}_2) = \frac{4}{a^2} - \frac{2b^2}{a^4} + \frac{4a^2}{c^4}, \\ \text{Ric}(\tilde{E}_3, \tilde{E}_3) &= \frac{2b^2}{a^4} + \frac{4b^2}{c^4}, \\ \text{Ric}(\tilde{B}_l, \tilde{B}_l) &= \frac{12c^2 - 2(2a^2 + b^2)}{c^4} \end{aligned}$$

as desired. \square

Given Lemma 5.1, we can now write down the Ricci flow equations for our ansatz in (2.15):

$$\begin{aligned}\frac{d}{dt}(a^2) &= -4 \left(2 - \frac{b^2}{a^2} + 2 \frac{a^4}{c^4} \right); \\ \frac{d}{dt}(b^2) &= -4 \left(\frac{b^4}{a^4} + 2 \frac{b^4}{c^4} \right); \\ \frac{d}{dt}(c^2) &= -4 \left(6 - 2 \frac{a^2}{c^2} - \frac{b^2}{c^2} \right).\end{aligned}$$

To simplify the analysis of these equations we introduce some scale-invariant quantities and rescale the time parameter to find the following after a short computation.

Lemma 5.2. Define

$$A = \frac{a^2}{c^2} \quad \text{and} \quad B = \frac{b^2}{c^2}$$

and introduce a new variable s by

$$\frac{ds}{dt} = \frac{1}{c^2}.$$

If we let $\dot{A} = \frac{dA}{ds}$ and $\dot{B} = \frac{dB}{ds}$, then the Ricci flow equations for (2.15) imply that

$$\dot{A} = \frac{4(1-A)}{A} (B(1+A) - 2A(1-2A)); \quad (5.2)$$

$$\dot{B} = \frac{4B}{A^2} (2A^2(3-A) - B(1+3A^2)). \quad (5.3)$$

The equations (5.2)–(5.3) describe the rescaled Ricci flow.

5.2. Critical points and dynamics

It is straightforward to find the critical points and observe some basic facts about the dynamics of the flow equations (5.2)–(5.3) as follows.

Lemma 5.3. The only critical points for $A, B > 0$ of the flow equations (5.2)–(5.3) are

$$A = B = \frac{1}{5} \quad \text{and} \quad A = B = 1.$$

Moreover, the lines $A = B$ and $A = 1$ are preserved by the flow.

Lemma 5.2 leads us to draw the plot in Fig. 5.1 showing the dynamics of the equations. If we denote the curve where $B(1+A) = 2A(1-2A)$ for $A, B \geq 0$ by γ_A and the curve $2A^2(3-A) = B(1+3A^2)$ for $A, B \geq 0$ by γ_B , these are the curves which, together with the line $A = 1$, determine the sign of \dot{A} and \dot{B} .

The first plot in Fig. 5.1 indicates that there is a stable critical point at $A = B = 1$, where γ_A and $A = 1$ intersect. In the second plot, we focus on the critical point $A = B = 1/5$ where γ_A and γ_B intersect, which appears to be unstable (in fact, a saddle point).

Remark 5.4. The curves γ_A and γ_B also intersect at $A = B = 0$, which is a degenerate critical point in terms of the equations (5.2)–(5.3), since they are not defined there. From the geometric point of view, here M^7 has collapsed to the 4-dimensional orbifold N (or a point) in the fibration (2.1), since $a = b = 0$ in (2.15), and so any flow lines tending to $(0, 0)$ converge to the Einstein metric g_N on N up to scale (or 0). We can avoid the possibility of converging to a point along the rescaled Ricci flow, since the diameter will stay bounded away from 0.

Remark 5.5. One may observe from (5.2)–(5.3) that there are two other degenerate critical points if we allow $A = 0$ or $B = 0$:

$$(A, B) = \left(\frac{1}{2}, 0 \right) \quad \text{and} \quad (A, B) = (1, 0).$$

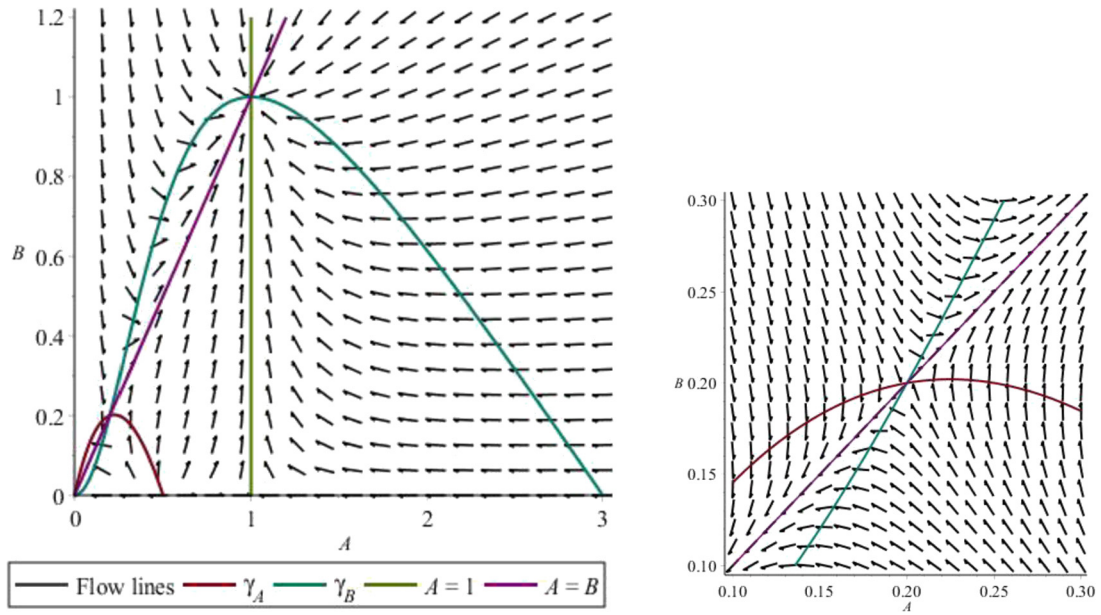


Fig. 5.1. Dynamic plot for Ricci flow.

These critical points correspond to a collapsed situation where the S^1 fibres corresponding to E_3 (in the notation of Definition 2.1) now have zero size since $b = 0$ in (2.13). The 7-manifold M has therefore collapsed to the twistor space Z^6 , which is a 2-sphere bundle over N^4 . The metrics with $(A, B) = (1, 0)$ and $(A, B) = (1/2, 0)$ correspond to two Einstein metrics on the twistor space Z : the canonical Kähler–Einstein and nearly Kähler metric, respectively. It is natural to see these Einstein metrics appear at the boundary of our Ricci flow ansatz as critical points. However, it is interesting to note that the Kähler–Einstein metric on Z does not play a distinguished role in the study of the Laplacian coflow and Laplacian flow, whereas the nearly Kähler metric does.

We now study the dynamics of the flow and show the following, recalling the fibration (2.1) of M^7 over a 4-dimensional base N^4 .

Theorem 5.6. *The only critical points for the Ricci flow (5.1), after rescaling, are the 3-Sasakian metric g_M and the squashed Einstein metric \tilde{g}_M on M^7 . The 3-Sasakian metric is a stable critical point for the rescaled Ricci flow within the ansatz (2.15), whilst the squashed Einstein metric is an unstable critical point which is a saddle point. Moreover, there is an open set of initial metrics in the ansatz (2.15), which can be chosen arbitrarily near \tilde{g}_M , such that they flow either to g_M or to the collapsed limit (even after rescaling) where the 3-dimensional fibres in (2.1) shrink to 0 and the flow converges to the Einstein 4-orbifold (N, g_N) .*

Proof. We recall the curves γ_A , γ_B introduced after Lemma 5.3 and plotted in Fig. 5.1.

For $A > 1$ we have that $\dot{A} < 0$ and hence A remains bounded as long as the flow exists. We also see that $\dot{B} < 0$ whenever B is above the curve γ_B (which passes through $(0, 0)$ and $(3, 0)$) and hence also remains bounded as long as the flow exists. Since

$$\frac{d}{dt}(c^2) = -4(6 - 2A - B),$$

the right-hand side is bounded and so c^2 cannot blow up in finite time. Moreover, c^2 goes to zero at a linear rate in t at most, and so we can integrate $\frac{1}{c^2}$ with respect to t to obtain s in Lemma 5.2.

We linearize (5.2)–(5.3) around $(A, B) = (1, 1)$ to obtain:

$$\dot{A} = -16A \quad \text{and} \quad \dot{B} = -16B.$$

Hence $(1, 1)$ is a stable critical point, which corresponds to the 3-Sasakian metric.

We now observe that if we start above the curve γ_A , which passes through $(0, 0)$ and $(\frac{1}{2}, 0)$, then $\dot{A} > 0$ for $A < 1$, so A is always bounded away from 0 in this region, depending on its initial value. Recalling that $\dot{A} < 0$ for $A > 1$, we deduce that all of the terms in (5.2)–(5.3) are bounded and there can be no periodic orbits in this region as A cannot cross the line $A = 1$ by Lemma 5.3. We also note that $\dot{B} > 0$ when $A \in (1/5, 1)$ and B is near 0 since then (A, B) lies below γ_B , and thus B is bounded away from 0 when $A \in (1/5, 1)$.

Since $(1/5, 1/5)$ lies on γ_A , we deduce that the flow will converge to the critical point at $(1, 1)$ if the initial value of (A, B) lies above or to the right of the curve γ_A and A is greater than $1/5$ initially. This proves the statement about initial conditions near \tilde{g}_M which flow to g_M , since \tilde{g}_M corresponds to $(1/5, 1/5)$.

Now if we specialize to the line $A = B$ which, by Lemma 5.2, is preserved, we see that

$$\dot{A} = 4(1 - A)(5A - 1). \quad (5.4)$$

We see immediately that for $A = B < 1/5$ we have $\dot{A} < 0$, for $A = B \in (1/5, 1)$ we have $\dot{A} > 0$. (For $A = B > 1$ we have $\dot{A} < 0$ again, as expected by the stability of $(1, 1)$.) Thus, the critical point at $(1/5, 1/5)$ is unstable, even within the restricted ansatz when $A = B$. Furthermore, if we linearize the system (5.2)–(5.3) around $(1/5, 1/5)$ we obtain

$$\dot{A} = -\frac{16}{5}A + \frac{96}{5}B \quad \text{and} \quad \dot{B} = \frac{192}{5}A - \frac{112}{5}B,$$

from which it follows that $(1/5, 1/5)$ is a saddle point, as the matrix corresponding to the above dynamical system has one positive and one negative eigenvalue: 16 and $-208/5$.

We see from (5.4) that we have initial conditions for (5.2)–(5.3), even with $A = B$, such that A and B go to 0. In fact, suppose we choose any initial condition in the region \mathcal{R} below γ_A but above γ_B , which means in particular that $A, B \in (0, \frac{1}{5})$. Since the flow lines enter \mathcal{R} vertically from above along γ_A and horizontally from the right along γ_B , no flow lines can leave \mathcal{R} so A, B are bounded and can only reach 0 when $(A, B) = (0, 0)$. In \mathcal{R} , $\dot{A} < 0$ and $\dot{B} < 0$ and there are no periodic orbits. We conclude that all flow lines starting in \mathcal{R} must converge to $(0, 0)$. We again note as in Remark 5.4 that M cannot collapse to a point along the rescaled Ricci flow. The discussion in Remark 5.4 then implies that the point $(0, 0)$ corresponds to the fibres of the fibration (2.1) collapsing, even in the rescaled Ricci flow, so that M^7 collapses to the base N^4 (with its Einstein metric). Since \mathcal{R} is open and has $(1/5, 1/5)$ on its boundary, this completes the proof. \square

Theorem 5.6 proves Theorem 1.3 in the Introduction.

Remark 5.7. By dynamical systems theory, there is a 1-dimensional stable manifold for the squashed Einstein metric within our rescaled Ricci flow ansatz. We can discern this stable manifold from the plots in Fig. 5.1. It might be interesting to understand whether this stable manifold (or, equally, the corresponding unstable manifold) has any geometric significance, e.g. any special curvature properties.

Data availability

No data was used for the research described in the article.

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