



Observational Inference on the Delay Time Distribution of Short Gamma-Ray Bursts

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Abstract

The delay time distribution of neutron star mergers provides critical insights into binary evolution processes and the merger rate evolution of compact object binaries. However, current observational constraints on this delay time distribution rely on the small sample of Galactic double neutron stars (with uncertain selection effects), a single multimessenger gravitational wave event, and indirect evidence of neutron star mergers based on *r*-process enrichment. We use a sample of 68 host galaxies of short gamma-ray bursts to place novel constraints on the delay time distribution and leverage this result to infer the merger rate evolution of compact object binaries containing neutron stars. We recover a power-law slope of $\alpha = -1.83^{+0.35}_{-0.39}$ (median and 90% credible interval) with $\alpha < -1.31$ at 99% credibility, a minimum delay time of $t_{\min} = 184^{+67}_{-79}$ Myr with $t_{\min} > 72$ Myr at 99% credibility, and a maximum delay time constrained to $t_{\max} > 7.95$ Gyr at 99% credibility. We find these constraints to be broadly consistent with theoretical expectations, although our recovered power-law slope is substantially steeper than the conventional value of $\alpha = -1$, and our minimum delay time is larger than the typically assumed value of 10 Myr. Pairing this cosmological probe of the fate of compact object binary systems with the Galactic population of double neutron stars will be crucial for understanding the unique selection effects governing both of these populations. In addition to probing a significantly larger redshift regime of neutron star mergers than possible with current gravitational wave detectors, complementing our results with future multimessenger gravitational wave events will also help determine if short gamma-ray bursts ubiquitously result from compact object binary mergers.

Unified Astronomy Thesaurus concepts: Gamma-ray bursts (629); Gravitational wave sources (677); Compact binary stars (283)

1. Introduction

The transient event GW170817 (Abbott et al. 2017a), which was observed in both gravitational waves (GWs) and electromagnetic (EM) waves, firmly established the connection between binary neutron star (BNS) mergers and short gamma-ray bursts (GRBs). The host galaxy of this event, NGC4993, was subsequently identified through broadband EM emission of the ensuing kilonova (Allam et al. 2017; Arcavi et al. 2017; Coulter et al. 2017a; Lipunov et al. 2017; Tanvir et al. 2017; Yang et al. 2017) and eventually, the nonthermal afterglow of the GRB (see Margutti & Chornock 2021, and references therein). Under the paradigm that all short GRBs originate from compact object binary mergers, such host associations can unveil unprecedented information regarding the stellar populations in galaxies that host compact object binary mergers, thereby allowing for novel constraints on the progenitors of these events.

One key aspect of compact object binary mergers that can be illuminated using such host associations is the delay time distribution (DTD). The DTD yields important information

regarding the birth properties of compact object binaries, such as their typical orbital separations at birth and inspiral times. This, in turn, provides critical constraints on evolutionary processes of the progenitor binary stellar system and can inform modeling of massive-star binaries, which form GW sources accessible by current and future observatories.

GW170817 is the only GW event to date where the host galaxy has been confidently established, and a number of fortuitous aspects of the system that produced this signal were paramount in the discovery of its EM counterpart: for example, it was well within the horizon of the GW detectors (Abbott et al. 2018), the GRB jet was aligned $\approx 20^\circ$ from our line of sight and eventually came into view (Margutti & Chornock 2021), and it was in close proximity to a massive galaxy that was prioritized by targeted searches (Coulter et al. 2017b), although wide-field follow-up searches would have found it regardless (Soares-Santos et al. 2017). While the BNS detection rate in the forthcoming observing run of the LIGO–Virgo–Kagra interferometer network will be $\mathcal{O}(1)$ /month (Abbott et al. 2018), the rate of BNS mergers with an EM counterpart detection precise enough for confident host association is more uncertain given the unique observational challenges (e.g., Coughlin et al. 2018; Dichiara et al. 2020; Colombo et al. 2022; Perna et al. 2022). In the next observing run, the BNS range for the LIGO detectors is expected to be 160–190 Mpc (Abbott et al. 2018), and GW mergers are more likely to be detected close to the detector horizon where the sensitive

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volume is the largest. Three-dimensional GW sky localization volumes will be larger at these distance, and even with improved search capabilities from state-of-the-art wide-field telescopes, it is unclear how rapidly multimessenger events with definitive host associations will be accumulated in the coming years.

On the other hand, the number of short GRBs with confident host galaxy associations has increased substantially over the past two decades. Over the next decade, the cumulative number of these host associations may still exceed the number of host associations found via multimessenger observations. GRB monitoring missions such as Swift and Fermi have observed hundreds of short GRBs, with subsequent searches in the X-ray, optical, and radio bands for GRB afterglows leading to the identification of probable host galaxies for dozens of these events over a wide range of redshifts (Berger 2014; Fong et al. 2022; Nugent et al. 2022; O’Connor et al. 2022). Photometric and spectroscopic follow-up observations of the afterglow and the host galaxy can then unveil crucial aspects of the GRB–host connection that encode information about the progenitors of BNS systems themselves, such as the stellar mass and star formation rate (SFR) of the galaxy, its star formation history (SFH), and the galactocentric offset of the GRB (Fong et al. 2013; Fong & Berger 2013).

Here, we consider direct implications of GRB–galaxy connections on compact object binary formation and the evolution of their progenitor stars, leveraging an updated sample of GRB host galaxies and their properties presented in Fong et al. (2022) and Nugent et al. (2022) to place novel constraints on the DTD of short GRBs. This expansive catalog of short GRB hosts is likely to outnumber the hosts of multimessenger GW events for years to come, and depending on the highly uncertain rate of discovery of EM counterparts to GW events, may not be met with multimessenger events until the onset of third-generation GW detectors. Besides probing a much larger cosmological volume than a multimessenger sample, future comparisons of DTD constraints between short GRBs and GWs with EM counterparts could help determine whether the BNS merger paradigm for short GRBs is universal. Furthermore, DTD constraints from such an extragalactic population are complementary to those from compact object binaries observed in the Milky Way and with GWs, allowing for comparisons of these distinct probes to help unravel selection effects that impinge upon the detection of each population individually.

The format of this Letter is as follows: In Section 2, we summarize the catalog of GRB host galaxies used in our analysis and assumptions regarding their SFHs. Our main results are in Section 3, where we describe our inference methodology and present our constraints on the DTD of short GRBs. We show implications of our constraints on the predicted compact object binary merger rate in Section 4. In Section 5 we discuss caveats of our analysis and highlight other aspects of binary stellar evolution and compact object binary formation that will be addressed in future work. We summarize our results in Section 6. Throughout this work we employ a standard Λ CDM cosmology of $H_0 = 69.6 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.286$, $\Omega_\Lambda = 0.714$ (Hinshaw et al. 2003; Bennett et al. 2014), consistent with Nugent et al. (2022). All code and data used in this analysis are available on Zenodo.⁷

2. Catalog of Short GRB Hosts

Recently, Fong et al. (2022) presented an updated catalog of 84 short GRB hosts with broadband photometry and well-studied explosion environments. Combined with the literature, the catalog comprises 542 photometric data points, 42 spectra, and 83 offset measurements, and is comprehensive for Swift short GRBs and neutron star GW merger events discovered in 2005–2021.⁸ This catalog, as well as other recent short GRB catalogs (e.g., O’Connor et al. 2022), have more than doubled the number of short GRBs with confident host associations and also include a number of systems that are found to be “hostless” with no coincident host to deep optical limits, depending on the threshold defining host galaxy association (Berger 2010).

Of the 84 short GRB hosts with broadband photometry presented in Fong et al. (2022), 69 were determined to have sufficient observational data (detected in ≥ 3 photometric bands) to model the galactic spectral energy distribution and infer aspects of the host; the detailed properties of this population of short GRB host galaxies are analyzed in Nugent et al. (2022). Approximately 71% of these hosts have spectroscopic redshifts. For the main analysis in this work, we include in our sample NGC4993, the host galaxy of multimessenger event GW170817/GRB170817, and do not include the long-duration GRB060614 and GRB211211A despite that they may be the result of compact object binary mergers; we perform additional fits to the DTD with the inclusion of the host galaxies of these two long-duration GRB events and comment on the impact in Section 5. This totals in 68 host galaxies used in our main analysis, with $\approx 82\%$ of the hosts being star-forming, $\approx 6\%$ transitioning, and $\approx 12\%$ quiescent. The properties of the host galaxies in our sample are modeled consistently as described below.

2.1. Host Galaxy Properties

Properties of the host galaxies, such as host redshift, stellar mass, stellar population ages, metallicity, and SFR, are inferred using the *Prospector* package (Leja et al. 2017; Johnson et al. 2021), as described in Nugent et al. (2022). *Prospector* uses the *dynesty* nested sampler (Speagle 2020) to constrain aspects of galaxies based on available photometric and spectroscopic data. Redshifts are fixed in this inference for galaxies that have spectroscopically determined redshifts, as the redshift measurement uncertainties for these hosts are negligible.

The SFH of each host, $\psi(t)$, is modeled assuming a delayed- τ functional form,⁹

$$\psi(t) = M_F \times \frac{te^{-t/\tau}}{\int_0^{t_{\text{SF}}} t'e^{-t'/\tau} dt'}, \quad (1)$$

⁸ The sample in Fong et al. (2022) also includes long-duration GRB060614 and GRB211211A, which are thought to be the result of neutron star mergers due to the lack of a supernova to deep optical limits and the observation of a kilonova, respectively (Gehrels et al. 2006; Rastinejad et al. 2022).

⁹ A number of other parametric functional forms (Carnall et al. 2019) or nonparametric approaches (Leja et al. 2019) can be used when reconstructing the SFH, sometimes yielding discrepant results. However, parametric fits are used in this work to better establish uniformity in our data set, which has an inconsistent amount and quality of data across host observations. Parametric SFHs are also more commonly used in host galaxy literature, making direct comparisons more seamless. We comment on this systematic uncertainty in Section 5.

⁷ <https://doi.org/10.5281/zenodo.7015221>

where M_F is the total stellar mass formed in the galaxy, t_{SF} is the time at which star formation commenced relative to the observation time, and τ is the e-folding factor. Posterior samples for the parameters M_F , t_{SF} , and τ , as well as the redshift z for hosts that only have photometric redshifts, fully construct the SFH for each `Prospector` posterior sample. We convert Equation (1) into a redshift-dependent SFH, $\psi(z)$, in our DTD parameter hyperlikelihood described in Section 3.1.

2.2. Criteria for Host Associations

The probability of chance coincidence (P_{cc}) metric is used for determining the confidence of a particular host association (Bloom et al. 2002). This metric decreases (i.e., results in a higher likelihood that the GRB is correctly associated with a particular host) with decreasing optical magnitude of the potential host and decreasing angular offset of the GRB, which leads to a higher confidence in host association for GRBs in close proximity to higher-mass galaxies. We use the P_{cc} criteria of Fong et al. (2022) to broadly categorize GRBs in the catalog based on host association confidence: a “Gold Sample” with $P_{cc} \leq 0.02$, a “Silver Sample” with $0.02 < P_{cc} \leq 0.10$, and a “Bronze Sample” with $0.10 < P_{cc} \leq 0.20$.

In this work, our full sample considers the 68 hosts observed in ≥ 3 photometric bands with $P_{cc} \leq 0.20$. We also present results using three subsets of this sample based on their redshift uncertainty and/or P_{cc} : a subset of hosts with spectroscopic redshifts (and therefore negligible redshift uncertainty; 48 hosts), a subset of hosts with $P_{cc} \leq 0.02$ (Gold Sample; 39 hosts), and a subset of hosts in the Gold Sample that also have spectroscopic redshifts (29 hosts). Note that in our main analysis we include GRB170817, which falls under the Gold Sample and has a spectroscopic redshift, and do not include long-duration GRB060614 and GRB211211A; hence the slight differences in the number of hosts in our subsamples compared to Nugent et al. 2022; see their Figure 1. We comment on completeness and potential biases from excluding certain GRB hosts from our sample in Section 5.

3. Delay Time Distributions

The delay time of compact object binaries is defined as the time between the formation of the progenitor stars and the merger of the two compact objects, such that $t_d = t_* + t_{insp}$, where t_* is the time between stellar birth and compact object binary formation, and t_{insp} is the compact object binary inspiral time. The DTD for isolated compact binaries is typically parameterized as a power-law distribution¹⁰ with a minimum delay time t_{min} , which encodes the minimum stellar evolutionary timescales required to form a compact object binary ($\mathcal{O}(10)$ Myr for BNS progenitors) and the minimum orbital separation of a compact object binary at formation, and can include a (somewhat arbitrary) maximum delay time t_{max} that

can be much larger than the Hubble time,

$$p(t_d|\alpha, t_{min}, t_{max}) = \begin{cases} \mathcal{N} t_d^\alpha, & t_{min} \leq t_d \leq t_{max} \\ 0, & \text{otherwise} \end{cases}, \quad (2)$$

where α is the spectral index of the power-law distribution and

$$\mathcal{N} = \begin{cases} \left[\log_{10} \left(\frac{t_{max}}{t_{min}} \right) \right]^{-1}, & \alpha = -1 \\ (1 + \alpha)(t_{max}^{1+\alpha} - t_{min}^{1+\alpha})^{-1}, & \text{otherwise} \end{cases} \quad (3)$$

is the normalization constant for the probability distribution function.

The slope of the DTD for compact object binary mergers is typically assumed to be $\alpha \approx -1$ (e.g., Piran 1992). This assumption stems from the fact that in general $t_{insp} \gg t_*$ and from coupling the equations that govern the GW inspiral with an assumption that the orbital separation distribution of compact object binaries follows the distribution of massive O/B stars, which is approximately flat in log with $dN/da \propto a^{-1}$, where a is the semimajor axis (for example, the best-fit orbital separation distribution from a survey of massive-star binary initial properties in Sana et al. (2012) finds $dN/da \propto a^{-0.83}$). Because the inspiral timescales as $t_{insp} \propto a^4$, it follows that $da/dt_{insp} \propto t_{insp}^{-3/4}$ and, assuming $dN/da \propto a^{-1}$, $dN/dt_{insp} \propto t_{insp}^{-1/4}$. The distribution of inspiral times is thus $dN/dt_{insp} \propto t_{insp}^{-1}$. However, hardening phases during the coevolution of massive-star binaries, such as common envelopes or stable mass transfer phases from a more massive donor star to a less massive compact object accretor, can lead to steeper orbital separation distributions and therefore steeper slopes in the DTD (e.g., Belczynski et al. 2018). The DTD inferred from short GRBs is thus useful in constraining a variety of uncertain aspects of massive-star binary evolution.

Multiple observational probes have been explored in constraining aspects of the short GRB/BNS DTD. The redshift distribution of short GRBs has been able to place broad constraints on the DTD (Nakar et al. 2006; Berger et al. 2007; Jeong & Lee 2010; Hao & Yuan 2013; Wanderman & Piran 2015; Anand et al. 2018), with high-redshift GRBs in particular allowing for stronger constraints on the minimum delay time and power-law slope of the DTD (Paterson et al. 2020; Nugent et al. 2022; O’Connor et al. 2022). The binary–host connection interpreted via host galaxy demographics and galaxy scaling relations has also shown promise in constraining BNS kick velocities, delay times, and the properties of galaxies that host BNS mergers and short GRBs (Zheng & Ramirez-Ruiz 2007; Kelley et al. 2010; Fong et al. 2013; Behroozi et al. 2014; Adhikari et al. 2020; Santoliquido et al. 2022). The Milky Way offers a limited sample of BNS systems with well-characterized orbital properties and inspiral times, many of which are much longer than the Hubble time, that are useful for examining the BNS DTD (Beniamini & Piran 2016; Tauris et al. 2017; Vigna-Gómez et al. 2018; Andrews & Zezas 2019; Beniamini & Piran 2019). However, uncertain selection effects inherent to this sample may affect inference of DTD parameters (e.g., Tauris et al. 2017). All these observational probes are complemented by population modeling of compact object binaries that explore variations in binary evolution physical assumptions, and therefore the predicted DTD of compact object binaries (Belczynski et al. 2018; Chruslinska et al. 2018; Broekgaarden et al. 2022; Santoliquido et al. 2022);

¹⁰ Another commonly used functional form for the DTD of short GRBs is a log-normal distribution (Nakar et al. 2006; Berger et al. 2007; Wanderman & Piran 2015), though this functional form is in tension with the growing number of high-redshift short GRBs and is not consistent with the expected DTD from compact object binary mergers based on population modeling predictions and the power-law orbital period distribution of their binary massive-star progenitors.

these predicted compact object populations can be compared and constrained with the population properties and merger rates observed by GW detectors (Abbott et al. 2021a). Though GW observations of compact object binary mergers also hold promise in constraining the DTD of compact object mergers, both through the use of galaxy scaling relations (Safarzadeh & Berger 2019; Safarzadeh et al. 2019b) and multimessenger events whose host galaxies have been identified and characterized (Safarzadeh et al. 2019a), it may require the accumulation of hundreds of host galaxy identifications to achieve strong constraints on the DTD.

3.1. Inferring the DTD Parameters

We follow Safarzadeh et al. (2019a) in constructing our likelihood function for the DTD, which is the key component of our population inference. From the `Prospector` modeling of each host galaxy, we have a set of posterior samples that define the SFH and, for hosts without spectroscopically measured redshifts, the galaxy redshift; a fixed redshift is used for the galaxies that have spectroscopic redshifts. We draw $N_{\text{samp}} = 100$ `Prospector` samples from each host, where we weight our draws with the inverse of the prior to get draws from the likelihood rather than the posterior.¹¹ For a given host galaxy i and likelihood sample j , the expected merger rate at the measured redshift z_i^j is

$$\dot{n}_i^j = \int_{z'=\infty}^{z'=z_i^j} p(t_{\text{lb}}' - t_{\text{d}} | \alpha, t_{\text{min}}, t_{\text{max}}) \lambda \psi_i^j(z') \frac{dt}{dz}(z') dz', \quad (4)$$

where t_{lb}' is the lookback time at redshift z' , λ is the BNS formation efficiency, which we assume is $10^{-5} M_{\odot}^{-1}$ and does not evolve with redshift (see e.g., Broekgaarden et al. 2022), $\psi(z)$ is the SFR as a function of redshift defined in Section 2.1, $dt/dz = [H_0 E(z)(1+z)]^{-1}$, and $E(z) = \sqrt{\Omega_M(1+z)^3 + \Omega_\Lambda(z)}$.

The probability of observing a single short GRB in a host galaxy over a set interval of time Δt follows a Poisson distribution based on the merger rate of Equation (4). Approximating the marginalization over likelihood samples as a discrete sum, the hyperlikelihood for observing a single short GRB in a particular host galaxy given the set of DTD parameters becomes

$$\mathcal{L}(\text{obs}_i | \alpha, t_{\text{min}}, t_{\text{max}}) \approx \frac{\mathcal{A}}{N_{\text{samp}}} \sum_{j=1}^{N_{\text{samp}}} (\dot{n}_i^j \Delta t) e^{-\dot{n}_i^j \Delta t}, \quad (5)$$

where $\mathcal{A}$ is a constant that normalizes the likelihood, and Δt is the fiducial observation time, which we assume to be $\Delta t = 10$ yr, though this choice does not impact results so long as the expected number of events during the observing period satisfies $\dot{n} \Delta t \ll 1$ event, which is the case for all GRB host galaxies in our sample. We note that this formulation of the hyperlikelihood assumes that the DTD parameters are not prone to selection effects, that is, the probability of a detection does not depend on α , t_{min} , and t_{max} . Selection effects may impact the detection of more highly offset (i.e., “hostless”) systems with less luminous counterparts and/or weaker host associations, as well as systems at high redshifts with dimmer

hosts; we discuss potential implications of selection effects further in Section 5.

Assuming the observed short GRB observations are independent, the hyperposterior for the DTD parameters is

$$P(\alpha, t_{\text{min}}, t_{\text{max}} | \text{obs}) \propto \prod_{i=1}^{N_{\text{GRB}}} \mathcal{L}(\text{obs}_i | \alpha, t_{\text{min}}, t_{\text{max}}) \times \pi(\alpha, t_{\text{min}}, t_{\text{max}}), \quad (6)$$

where N_{GRB} is the number of GRB hosts used in our sample, and $\pi(\alpha, t_{\text{min}}, t_{\text{max}})$ is the prior distribution on the DTD parameters. We assume that the priors on α , t_{min} , and t_{max} are independent and uniform on the ranges of $[0, 3]$, $[5 \text{ Myr}, 500 \text{ Myr}]$, and $[3 \text{ Gyr}, 100 \text{ Gyr}]$, respectively.

Sampling the likelihood of Equation (5) is computationally prohibitive, as it involves calculating an integral at each step in the parameter space of the DTD distribution. To avoid this, we precompute a regular grid of likelihoods for each host galaxy and create an interpolant of these grids when evaluating the likelihood in Equation (5). This grid contains 20 points for each DTD parameter uniformly spaced across their prior range, resulting in 8000 likelihood interpolants for each host galaxy. We find that this approach provides consistent results compared to directly evaluating the likelihood, as the likelihood surface varies smoothly across the DTD parameter space. For generating posterior distributions of the DTD parameters, we use the `dynesty` nested sampler (Speagle 2020) as implemented in Bilby (Ashton et al. 2019).

We note that this approach only considers host galaxies in which a short GRB has been observed and not galaxies in which no short GRB occurred during the fiducial observing period. Locally, this has been shown to lead to only a mild bias in the inferred DTD parameters (Safarzadeh et al. 2019a).

3.2. DTD Constraints

Our main results are in Figure 1, which shows the constraints on the DTD parameters from our hierarchical inference of Section 3.1. We show the posterior distributions of the DTD parameters using the full population, as well as the three subsets of the population described in Section 2.2.

Considering our full sample of 68 GRBs in the inference, we find that the posterior distribution of DTD parameters significantly deviates from the prior, with a power-law slope of $\alpha = -1.83_{-0.39}^{+0.35}$ and a minimum delay time of $t_{\text{min}} = 184_{-79}^{+67}$ Myr, where we quote the median and 90% symmetric credible interval. These parameters, as expected, exhibit a strong correlation as the bulk of the short GRB sample is consistent with relatively short delay times of $\mathcal{O}(100 \text{ Myr})$; higher values of α (i.e., a shallower power-law slope) correspondingly decrease the inferred values for t_{min} so that the DTD probability distribution function still has enough support at the low end of the distribution to adequately explain the population as a whole.

The constraints on the DTD parameters are fairly consistent when using the full sample of GRB hosts as opposed to the subsamples described in Section 2.2. When only using samples with $P_{\text{cc}} \leq 0.02$ (i.e., Gold Associations) in the inference, the constraints on DTD parameters show little variation compared to the full sample where hosts only satisfy $P_{\text{cc}} \leq 0.20$. This suggests that the DTD constraints are not strongly dependent on the specific criteria used for host associations, namely, the optical magnitude of the host and the offset of the GRB with respect to the host. On the other hand, when considering only

¹¹ All relevant `Prospector` parameters used in this analysis were sampled with uniform priors, making this prior reweighting irrelevant.

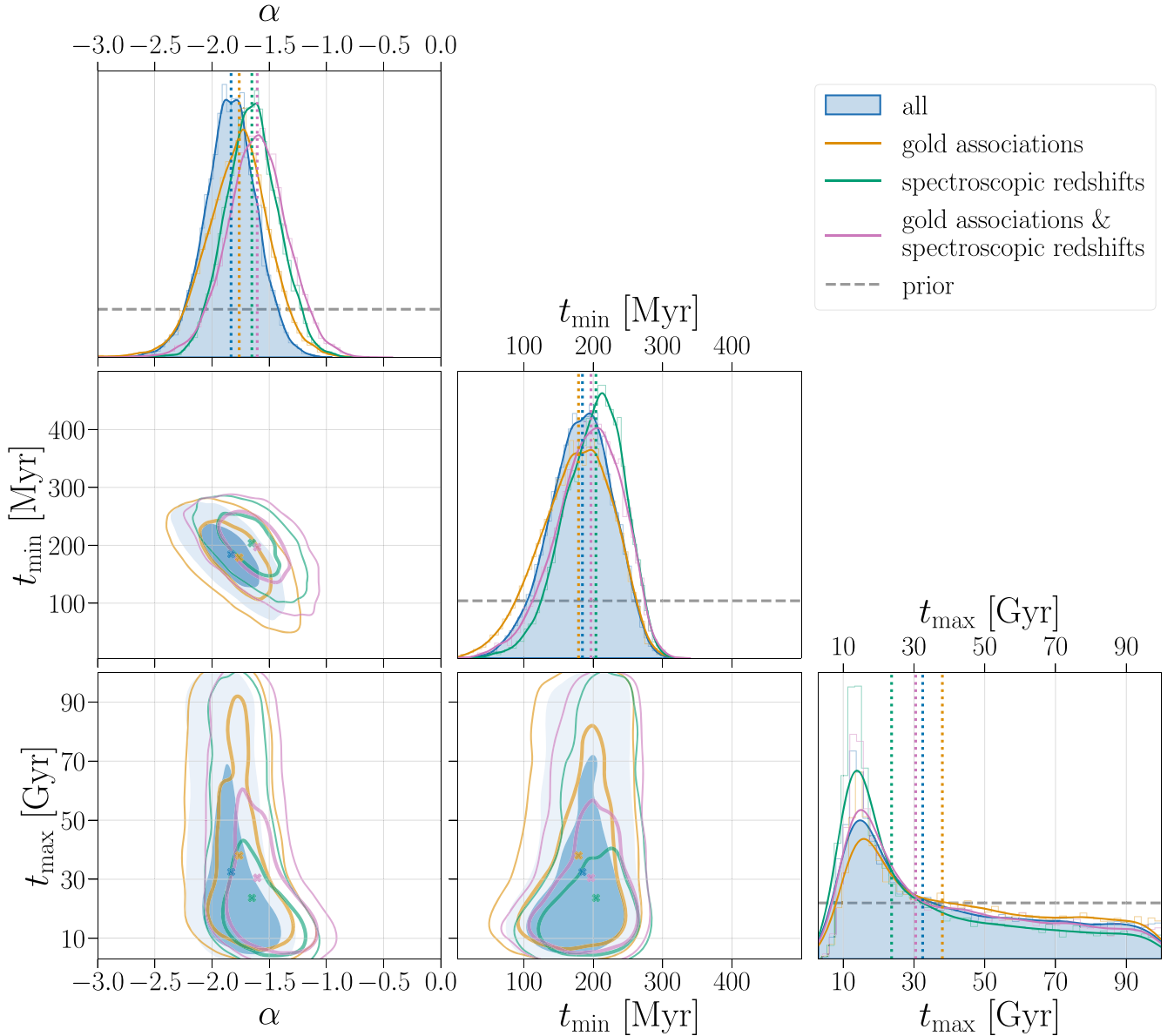


Figure 1. Constraints on DTD parameters when including all host associations with $P_{cc} \leq 0.20$ in the inference (filled blue), as well as three subsets that only include samples that have $P_{cc} \leq 0.02$ (yellow; Gold Association), samples that have spectroscopic redshifts (green), and samples that have spectroscopic redshifts and $P_{cc} \leq 0.02$ (pink). The contours show the 50% and 90% credible regions, with dotted lines marking the median for each marginalized distribution. The gray dashed lines in the panels plotting the marginal distributions show the prior distribution for each parameter, which is uniform for all parameters of the DTD. Both kernel density estimates and binned histograms are shown for the marginal distributions.

associations for which the host galaxy has a spectroscopically measured redshift, we find mild changes in our inferred DTD parameters. In particular, the minimum delay time pushes to slightly larger values and becomes more tightly constrained with a median value that is 20 Myr larger than when all hosts are included, though still within the range of uncertainty. Furthermore, the subsample with spectroscopic redshifts strongly excludes extremely short delay times, with $t_{\min} > 88$ Myr at 99% credibility. The power-law index shifts to slightly shallower slopes as opposed to the full sample, with $\alpha = -1.65^{+0.39}_{-0.38}$. As the spectroscopic sample has precise redshift measurements, the modeled SFHs are much less uncertain compared to the photometric redshift sample, which is the main driver of these tighter constraints. We also note that most of the host galaxies in the spectroscopic redshift sample have redshifts $z < 1$ (Fong et al. 2022), which may lead to mild

differences in our DTD constraints relative to the full sample as discussed in Section 5.

Neither α nor $t_{\min}$ are strongly correlated with $t_{\max}$, though for $t_{\max} \lesssim 30$ Gyr, decreasing the value of $t_{\max}$ leads to slightly more support for smaller values of $t_{\min}$ and larger values of α (i.e., shallower power-law slopes). The posterior distribution for $t_{\max}$ peaks at ~ 13 Gyr, and the upper limit for $t_{\max}$ rails against the upper bound of our prior. Thus, we cannot make meaningful statements for the upper limit of $t_{\max}$, which translates to a lack of constraints on the maximum orbital separation of compact object binary systems containing a neutron star at formation. Posterior support for values below $t_{\max} \sim 10$ Gyr drops precipitously; using all host galaxies, we constrain $t_{\max} > 7.95$ Gyr at 99% credibility. This constraint is driven by the systems in our population associated with quiescent host galaxies and old stellar populations.

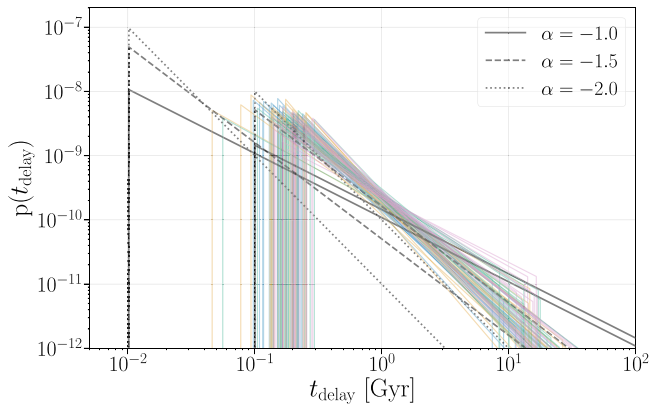


Figure 2. DTDs constructed from random draws of our posterior distribution. The colored lines match the set of samples used in the inference as in Figure 1. The black solid, dashed, and dotted lines show fiducial DTDs with power-law slopes α of -1 , -1.5 , and -2 , respectively, for comparison. We show this set of fiducial distributions for $t_{\min} = 10$ Myr and $t_{\min} = 100$ Myr, with $t_{\max}$ fixed at 100 Gyr.

In Figure 2, we show DTDs constructed from our posterior distributions, as well as DTDs using fiducial values for comparison. As described above, we find a preference for minimum delay times of $t_{\min} \sim 150$ Myr and power-law slopes of $\alpha \sim -1.5$. These constraints can help inform binary evolution modeling of BNS progenitors, as slopes with $\alpha < -1$ may be indicative of hardening phases in the evolution of the progenitor such as mass transfer or common envelope phases (Belczynski et al. 2018), and larger minimum delay times hint at proposed fast-merging BNS channels (e.g., a Case BB common envelope scenario, in which the progenitor of the second-born neutron star proceeds through a second (unstable) mass transfer episode as an evolved naked helium star; see Dewi et al. 2002; Ivanova et al. 2003) may operate inefficiently.

4. Implications for Gravitational Wave Observations

Given our constraints on the parameters of the short GRB DTD, we can construct the expected merger rate evolution of BNS mergers under the assumption that all short GRBs are the result of this class of compact object binary merger. We show the expected merger rate density evolution in Figure 3, assuming that BNSs follow the underlying cosmic SFH of Madau & Fragos (2017). The BNS merger rate density peaks at a redshift of $z \sim 1.6$ – 1.7 . The overall normalization of the merger rate evolution is directly related to the assumed BNS production efficiency, which we take to be fixed at $10^{-5} M_{\odot}^{-1}$. However, this assumption does not affect the shape of the merger rate evolution or the location of its peak. Given this fixed value for the BNS production efficiency, the local merger rate density is predicted to be $\sim 135 \text{ Gpc}^{-3} \text{ yr}^{-1}$, consistent with the current constraints from the LIGO Scientific and Virgo Collaboration (Abbott et al. 2021b).

The inferred BNS merger rate evolution does not vary significantly for the different subsets of samples described in Section 2.2 as the inferred DTD is relatively robust. Though the redshift of the peak merger rate density is well beyond the BNS horizon of current-generation GW detectors operating at design sensitivity (Abbott et al. 2018), it should be within the range of planned third-generation GW detectors (Evans et al. 2021). If inconsistencies are found between the peak of the BNS merger rate density inferred from short GRBs and the peak of the BNS

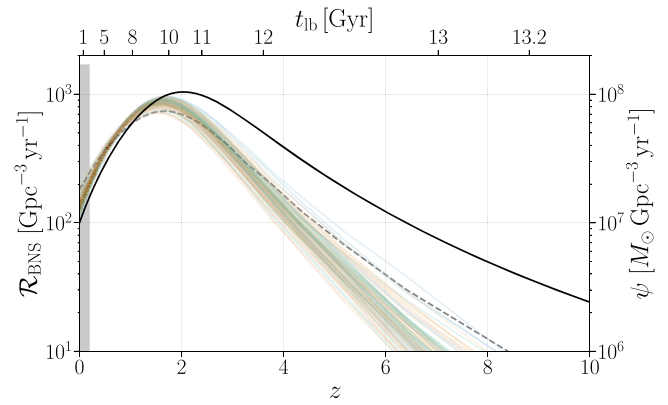


Figure 3. Expected merger rate evolution of BNSs based on our DTD constraints. The black line (right axis) shows the SFR density evolution from Madau & Fragos (2017). The colored lines (left axis) match the set of samples used in the inference as in Figure 1 and show the predicted BNS comoving merger rate density for random draws from our DTD parameter inference, assuming their progenitor stars are born according to the SFR density of Madau & Fragos (2017). The black dashed line shows the expected BNS merger rate evolution assuming fiducial DTD values of $\alpha = -1$, $t_{\min} = 10$ Myr, and $t_{\max} = 14$ Gyr. The shaded gray region shows the local BNS merger rate density constraints from Abbott et al. (2021b).

merger rate observed by future GW detectors, it may indicate an additional mechanism for producing short GRBs is at play or that other transients (e.g., long GRBs) are also caused by compact object binary mergers.

5. Discussion

Although the DTD constraints of short GRBs in this work are not dependent on the specific type of formation mechanism, these constraints can help inform uncertain aspects of massive-star binary evolution within the BNS paradigm for short GRBs. Assuming that the merger time is purely driven by GW inspiral (e.g., in the case of isolated binary evolution), the DTD encodes information about the orbital properties of compact object binary systems at formation. Given the short stellar evolutionary timescales of BNS progenitors prior to the second supernova that forms the compact object binary (~ 10 – 50 Myr; Andrews & Zezas 2019), the delay time is approximately the inspiral time, except for the potential low-end portion of the DTD where a short inspiral time becomes comparable to the stellar evolution timescale.

The DTD can thus be used to place constraints on the minimum separation of BNSs at formation, as well as evolutionary phases in the progenitors that can steepen the compact object binary orbital separation distribution and therefore the DTD (Belczynski et al. 2018; Broekgaarden et al. 2022). For example, given the first percentile of our recovered $t_{\min}$ distribution when using all host galaxies ($t_{\min}^{1\%} = 72$ Myr) and assuming a BNS formation timescale of $t_* = 30$ Myr, the minimum orbital separation at BNS formation is $\approx 1.1 R_{\odot}$ (note that eccentricity in the orbit at BNS formation would expedite the inspiral, leading to larger orbital separations for a given inspiral time). The preference of our results for longer minimum delay times $\gtrsim 100$ Myr may indicate that late-stage hardening phases, such as the Case BB common envelope scenario (e.g., Dewi et al. 2002; Ivanova et al. 2003), may operate inefficiently or be nonexistent. This would have implications for the ability of BNS systems to provide the r -process enrichment observed in low escape velocity environments such as ultra-faint dwarf galaxies (Safarzadeh et al. 2018).

or environments with rapid star formation episodes such as globular clusters (Zevin et al. 2019), as enrichment in these environments requires short delay times of $\mathcal{O}(10 \text{ Myr})$. In addition to informing the distribution of orbital properties at compact object binary formation, delay times can be used to place joint constraints on the properties of a binary prior to the second supernova and the strength of the natal kick that formed the second neutron star (e.g., Andrews et al. 2015; Andrews & Mandel 2019; Tauris et al. 2017). Combining this information with the physical offset of the GRB with respect to its host galaxy can further improve such constraints (e.g., Abbott et al. 2017b; Zevin et al. 2020). We will explore the constraints that this catalog places on BNS progenitor pre-supernova orbital properties and natal kicks based on delay times and offsets in future work.

A significant amount of work has explored constraints on BNS progenitor properties based on the ~ 20 BNS systems observed in the Milky Way (Andrews et al. 2015; Tauris et al. 2017; Andrews & Mandel 2019; Andrews & Zezas 2019). The Galactic BNS population has uncertain selection effects that can be complemented with short GRB host galaxy associations (see below for caveats on the short GRB sample selection). In particular, the shortest delay time BNS systems in the Milky Way may be missed from surveys due to Doppler shifting from the orbit making their detection difficult, as well as their short lifetimes making their existence at any given point in time relatively rare (Tauris et al. 2017). On the other hand, long-delay-time BNS systems in the Milky Way may be hidden from surveys due to their long orbital periods, making their proper motion on the sky minuscule. We find a minimum delay time that is consistent with the shortest delay time systems in the Milky Way. For example, the BNS system J1757-1854 is estimated to have one of the shortest delay times of the population observed in the Milky Way; it will merge in $\sim 70 \text{ Myr}$ and has a characteristic age of $\sim 130 \text{ Myr}$ (Andrews & Mandel 2019) based on the spin-down of the pulsar in the system, giving it a delay time of $\sim 210\text{--}250 \text{ Myr}$ with the inclusion of a $\sim 10\text{--}50 \text{ Myr}$ timescale between stellar birth and BNS formation. This may be an indication that selection effects do not strongly impinge the detection of short delay time systems in the Milky Way, or that certain selection effects against observing short-delay-time systems also affect the short GRB population. The complementarity of these local and cosmological samples may also help us understand selection effects impacting the detection of counterparts to short GRBs, such as the possibility that highly offset short GRBs have fainter afterglows due to the tenuous nature of gas in the outskirts of galaxies (e.g., Berger 2010; Tunnicliffe et al. 2014).

The largest uncertainty in our analysis likely lies in the reconstruction of the SFH. Though the parametric delayed- τ SFH has been shown to be an adequate representation of quiescent galaxies, other parametric or nonparametric approaches may be more physically realistic for SFHs of host galaxies that still have substantial star formation or bursty histories (Carnall et al. 2019; Leja et al. 2019). However, due to the robustness of our results with differing subsets of GRB hosts, we expect our main conclusions to hold despite this systematic uncertainty. In future work, we will explore how the choice of SFH reconstruction affects our results by considering alternative parametric and nonparametric modeling techniques.

The host galaxy of the multimessenger event GW170817/GRB170817, NGC4993, is quiescent with stellar age estimates

of $\gtrsim 3\text{--}10 \text{ Gyr}$ (Blanchard et al. 2017; Im et al. 2017; Levan et al. 2017; Pan et al. 2017; Nugent et al. 2022). Though more quiescent galaxies host short GRBs at low redshifts ($z < 0.5$) than at higher redshifts (Nugent et al. 2022), properties of NGC4993 such as its mass-weighted age and specific SFR make it a mild outlier relative to the other host galaxies used in this work (see Nugent et al. 2022). Despite this, the exclusion of GRB170817 in our analysis has a very minor impact on our DTD constraints, only decreasing (increasing) our median recovered value of $\alpha(t_{\min})$ by 5%, well within the bounds of the uncertainty in these parameters. The DTD constraints in this work are consistent with a long tail of delay times with $t_{\max} > 7.95 \text{ Gyr}$ at high credibility and thus can adequately account for the relatively old stellar population of NGC4993 and the long inferred delay time of GW170817 (Blanchard et al. 2017; Adhikari et al. 2020).

Extremely short delay times of $\lesssim 10 \text{ Myr}$ have been argued for the progenitor of GRB060505, which is spatially associated with an active star-forming region (Ofek et al. 2007). This GRB had a T_{90} of $\sim 4 \text{ s}$, longer than the typical delineation between short and long GRBs of 2 s . However, deep imaging ruled out the presence of a supernova (typically associated with long GRBs) to deep optical limits. Since 10 Myr is approximately the shortest possible evolutionary timescale that a massive-star binary can form a BNS, for a delay time of $\sim 10 \text{ Myr}$ the inspiral timescale would need to be extremely short, requiring a BNS birth semimajor axis of $a_{\text{BNS}} \ll 1 R_{\odot}$, unless the BNS was born with a high eccentricity. Such a scenario is inconsistent with our constraints on the minimum delay time $t_{\min}$; we find $t_{\min} > 72 \text{ Myr}$ at 99% credibility when using all host galaxies in our sample and $t_{\min} > 88 \text{ Myr}$ at 99% credibility when only considering host galaxies with spectroscopically measured redshifts. Given our DTD constraints, this spatial coincidence with a star-forming region may instead be a lucky coincidence following the post-supernova migration of the BNS, or this particular event may be a separate class of long GRBs without a supernova (e.g., Fryer et al. 2006).

However, a number of other long GRB detections are lending credence to the possibility that some of these events are caused by compact object binary mergers. The discovery of a potential kilonova associated with the minute-long GRB211211A (Rastinejad et al. 2022) hints at BNS mergers causing some fraction of the long GRB population. Similarly, GRB060614 had a long duration (although is also classified as a possible short GRB with extended emission; Lien et al. 2014), and no coincident supernova was detected to deep optical limits (Gehrels et al. 2006). As the properties of the host galaxies of GRB211211A and GRB060614 were also modeled in Nugent et al. (2022), we performed additional analyses with the inclusion of these two hosts in our general short GRB population. These provided nearly identical results. Even though our DTD results are insensitive to the inclusion of these two long GRBs, the number of compact object binary mergers consistent with this picture of extended gamma-ray emission is uncertain and could potentially add more support for short delay times. However, it is unlikely that the majority of long GRBs result from a compact object merger paradigm, as many long GRBs in the local ($z < 0.5$) universe have been followed up extensively by electromagnetic observatories, and GRB211211A is the only one observed thus far to have a plausible kilonova counterpart (Rastinejad et al. 2022).

In addition to some fraction of BNS mergers masquerading as long GRBs, our sample used to constrain the DTD may suffer from other issues of incompleteness. As we rely on the modeling of host galaxies when constraining the DTD, we do not consider short GRBs that do not have a confident host association. Though the properties of short GRB hosts do not seem to deviate strongly as a function of the host-association confidence (see, e.g., Figure 4 of Nugent et al. 2022), neglecting these events may have a potential impact on both the low and high ends of our inferred DTD. For example, GRBs that are highly offset from their hosts may have afterglows with much lower luminosities, making precise localization (and therefore host identification) difficult (Perna et al. 2022). Such systems may have migrated over longer timescales to reach the highly offset locations of the burst and therefore may have longer delay times than the general population. Furthermore, the P_{cc} method for host identification may incorrectly associate a GRB with a faint underlying host rather than a bright host at a larger offset, though Fong et al. (2022) predicted this to be an effect only at the $\lesssim 7\%$ level. On the other hand, if such poorly associated GRBs are instead truly associated with faint galaxies that are below detection limits, we may be excluding additional systems with short delay times as these faint, low-mass galaxies are typically star-forming. Furthermore, though Swift can detect GRBs out to $z \sim 3$, there is likely some fraction of short GRBs that occur beyond this horizon, when the universe was $\lesssim 2$ Gyr old. Short GRBs that occur at these early stages in the history of the universe must have short delay times, and this selection effect may bias the general population in our analysis to longer delay times. This would lead to a larger inferred $t_{\min}$, and, due to the correlation between $t_{\min}$ and α , more negative values of α . However, this population of high-redshift short GRBs is likely small; assuming the SFH from Madau & Fragos (2017), $<10\%$ of stars are born beyond $z = 3$, and the fraction of compact object binary mergers beyond this redshift will be even smaller due to the delay time between formation and merger.

6. Conclusions

We have placed constraints on the delay time distribution (DTD) of short GRBs using the largest catalog of short GRB host associations to date with the inclusion of inference on host galaxy properties. Assuming these transient events result from the merger of BNS systems, the DTD for short GRBs can be directly translated to the DTD of BNS mergers. This allows for predictions of the expected merger rate evolution, and can constrain aspects of massive-star binary evolution physics and compact object binary formation. Our main results are as follows:

1. Based on a catalog of 68 short GRB host galaxies, we constrain the DTD of short GRBs to have a power-law slope steeper than flat-in-log, with a power-law index of $\alpha = -1.83^{+0.35}_{-0.39}$, a minimum delay time of $t_{\min} = 184^{+67}_{-79}$ Myr, and a maximum delay time that is >7.95 Gyr.
2. Using different subsets of the full data set that make cuts based on the host association confidence or whether the host galaxy has a spectroscopically measured redshift, we find our results to be robust. However, when using the spectroscopic redshift sample (which is dominated by $z < 1$ bursts), the DTD pushes to slightly larger values of

the minimum delay time and slightly shallower power-law slopes.

3. Assuming short GRBs are the result of BNS mergers, we construct the expected merger rate evolution of BNS mergers, which we predict to peak at a redshift of $z \sim 1.6\text{--}1.7$.

Constraints on the DTD is the tip of the iceberg of what can be accomplished given a large sample of compact object binary host galaxies. This binary–host sample can enable novel constraints on properties of the compact object binary at birth, and pairing such inference with the physical offsets of short GRBs with respect to their hosts can help unveil supernova mechanisms and the strength of kicks that neutron stars receive at formation. Pairing this sample with multimessenger GW events and the already existing population of BNSs in the Milky Way will help unravel uncertain selection effects and determine the ubiquity of BNS mergers as the cause of short GRBs, all of which we aim to explore in future work. Even when multimessenger BNS observations surpass the number of host-identified short GRBs, they will probe a lower redshift regime, and pairing these samples will be paramount for understanding the binary–host connection and its evolution over cosmic time.

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References

- Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017a, *PhRvL*, **119**, 161101
 Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017b, *ApJL*, **850**, L40
 Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2018, *LRR*, **21**, 3

- Abbott, R., Abbott, T. D., Abraham, S., et al. 2021a, arXiv:2111.03606
- Abbott, R., Abbott, T. D., Abraham, S., et al. 2021b, arXiv:2111.03634
- Adhikari, S., Fishbach, M., Holz, D. E., Wechsler, R. H., & Fang, Z. 2020, *ApJ*, **905**, 21
- Allam, S., Annis, J., Berger, E., et al. 2017, GCN, **21530**
- Anand, N., Shahid, M., & Resmi, L. 2018, *MNRAS*, **481**, 4332
- Andrews, J. J., Farr, W. M., Kalogera, V., & Willems, B. 2015, *ApJ*, **801**, 32
- Andrews, J. J., & Mandel, I. 2019, *ApJ*, **880**, L8
- Andrews, J. J., & Zezas, A. 2019, *MNRAS*, **486**, 3213
- Arcavi, I., Howell, D. A., McCully, C., et al. 2017, GCN, **21538**
- Ashton, G., Hübner, M., Lasky, P. D., et al. 2019, *ApJS*, **241**, 27
- Behroozi, P. S., Ramirez-Ruiz, E., & Fryer, C. L. 2014, *ApJ*, **792**, 19
- Belczynski, K., Bulik, T., Olejak, A., et al. 2018, arXiv:1812.10065
- Beniamini, P., & Piran, T. 2016, *MNRAS*, **456**, 4089
- Beniamini, P., & Piran, T. 2019, *MNRAS*, **487**, 4847
- Bennett, C. L., Larson, D., Weiland, J. L., & Hinshaw, G. 2014, *ApJ*, **794**, 2
- Berger, E. 2010, *ApJ*, **722**, 1946
- Berger, E. 2014, *ARA&A*, **52**, 43
- Berger, E., Fox, D. B., Price, P. A., et al. 2007, *ApJ*, **664**, 1000
- Blanchard, P. K., Berger, E., Fong, W., et al. 2017, *ApJ*, **848**, L22
- Bloom, J. S., Kulkarni, S. R., & Djorgovski, S. G. 2002, *AJ*, **123**, 1111
- Broekgaarden, F. S., Berger, E., Stevenson, S., et al. 2022, *MNRAS*, **516**, 5737
- Carnall, A. C., Leja, J., Johnson, B. D., et al. 2019, *ApJ*, **873**, 44
- Chruslinska, M., Belczynski, K., Klencki, J., & Benacquista, M. 2018, *MNRAS*, **474**, 2937
- Colombo, A., Salafia, O. S., Gabrielli, F., et al. 2022, arXiv:2204.07592
- Coughlin, M. W., Tao, D., Chan, M. L., et al. 2018, *MNRAS*, **478**, 692
- Coulter, D. A., Foley, R. J., Kilpatrick, C. D., et al. 2017b, *Sci*, **358**, 1556
- Coulter, D. A., Kilpatrick, C. D., Siebert, M. R., et al. 2017a, GCN, **21529**
- Dewi, J. D., Pols, O. R., Savonije, G. J., & Van Den Heuvel, E. P. 2002, *MNRAS*, **331**, 1027
- Dichiara, S., Troja, E., O'Connor, B., et al. 2020, *MNRAS*, **492**, 5011
- Evans, M., Adhikari, R. X., Afle, C., et al. 2021, arXiv:2109.09882
- Fong, W., & Berger, E. 2013, *ApJ*, **776**, 18
- Fong, W., Berger, E., Chornock, R., et al. 2013, *ApJ*, **769**, 18
- Fong, W., Nugent, A., Dong, Y., et al. 2022, arXiv:2206.01763
- Fryer, C. L., Young, P. A., & Hungerford, A. L. 2006, *ApJ*, **650**, 1028
- Gehrels, N., Norris, J. P., Barthelmy, S. D., et al. 2006, *Natur*, **444**, 1044
- Hao, J. M., & Yuan, Y. F. 2013, *A&A*, **558**, 4
- Hinshaw, G., Spergel, D. N., Verde, L., et al. 2003, *ApJS*, **148**, 135
- Hunter, J. D. 2007, *CSE*, **9**, 99
- Im, M., Yoon, Y., Lee, S.-K., et al. 2017, *ApJL*, **849**, L16
- Ivanova, N., Belczynski, K., Kalogera, V., Rasio, F. A., & Taam, R. E. 2003, *ApJ*, **592**, 475
- Jeong, S., & Lee, C. H. 2010, *JKPS*, **56**, 1619
- Johnson, B. D., Leja, J., Conroy, C., & Speagle, J. S. 2021, *ApJS*, **254**, 22
- Kelley, L. Z., Ramirez-Ruiz, E., Zemp, M., Diemand, J., & Mandel, I. 2010, *ApJL*, **725**, 91
- Leja, J., Carnall, A. C., Johnson, B. D., Conroy, C., & Speagle, J. S. 2019, *ApJ*, **876**, 3
- Leja, J., Johnson, B. D., Conroy, C., van Dokkum, P. G., & Byler, N. 2017, *ApJ*, **837**, 170
- Levan, A. J., Lyman, J. D., Tanvir, N. R., et al. 2017, *ApJ*, **848**, L28
- Lien, A., Sakamoto, T., Barthelmy, S. D., et al. 2014, in Proc. of Swift: 10 Years of Discovery (SWIFT 10), **38**, <http://pos.sissa.it/cgi-bin/reader/conf.cgi?confid=233>
- Lipunov, V., Gorbvskoy, E., Kornilov, V., et al. 2017, GCN, **21570**
- Madau, P., & Fragos, T. 2017, *ApJ*, **840**, 39
- Margutti, R., & Chornock, R. 2021, *ARA&A*, **59**, 155
- McKinney, W. 2010, in Proc. 9th Python in Science Conf., ed. S. van der Walt & J. Millman, 51, doi:10.25080/Majora-92bf1922-00a
- Nakar, E., Gal-Yam, A., & Fox, D. B. 2006, *ApJ*, **650**, 281
- Nugent, A., Fong, W., Dong, Y., et al. 2022, arXiv:2206.01764
- O'Connor, B., Troja, E., Dichiara, S., et al. 2022, *MNRAS*, **515**, 4890
- Ofek, E. O., Cenko, S. B., Gal-Yam, A., et al. 2007, *ApJ*, **662**, 1129
- Oliphant, T. E. 2006, A Guide to NumPy (USA: Trelgol Publishing)
- Pan, Y. C., Kilpatrick, C. D., Simon, J. D., et al. 2017, *ApJL*, **848**, L30
- Paterson, K., Fong, W., Nugent, A., et al. 2020, *ApJ*, **898**, L32
- Pérez, F., & Granger, B. E. 2007, *CSE*, **9**, 21
- Perna, R., Artale, M. C., Wang, Y.-h., et al. 2022, *MNRAS*, **512**, 2654
- Piran, T. 1992, *ApJL*, **389**, 45
- Price-Whelan, A. M., Sipocz, B. M., Günther, H. M., et al. 2018, *AJ*, **156**, 123
- Rastinejad, J. C., Gompertz, B. P., Levan, A. J., et al. 2022, arXiv:2204.10864
- Robitaille, T. P., Tollerud, E. J., Greenfield, P., et al. 2013, *A&A*, **558**, A33
- Safarzadeh, M., & Berger, E. 2019, *ApJ*, **878**, L12
- Safarzadeh, M., Berger, E., Leja, J., & Speagle, J. S. 2019a, *ApJ*, **878**, L14
- Safarzadeh, M., Berger, E., Ng, K. K. Y., et al. 2019b, *ApJ*, **878**, L13
- Safarzadeh, M., Ramirez-Ruiz, E., Andrews, J. J., et al. 2018, *ApJ*, **872**, 105
- Sana, H., De Mink, S. E., De Koter, A., et al. 2012, *Sci*, **337**, 444
- Santoliquido, F., Mapelli, M., Artale, M. C., & Boco, L. 2022, *MNRAS*, **516**, 3297
- Soares-Santos, M., Holz, D. E., Annis, J., et al. 2017, *ApJ*, **848**, L16
- Speagle, J. S. 2020, *MNRAS*, **493**, 3132
- Tanvir, N. R., Levan, A. J., & Vinrouge Collaboration 2017, GCN, **21544**
- Tauris, T. M., Kramer, M., Freire, P. C. C., et al. 2017, *ApJ*, **846**, 170
- Tunnicliffe, R. L., Levan, A. J., Tanvir, N. R., et al. 2014, *MNRAS*, **437**, 1495
- Van Der Walt, S., Colbert, S. C., & Varoquaux, G. 2011, *CSE*, **13**, 22
- Vigna-Gómez, A., Neijssel, C. J., Stevenson, S., et al. 2018, *MNRAS*, **481**, 4009
- Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, *NatMe*, **17**, 261
- Wanderman, D., & Piran, T. 2015, *MNRAS*, **448**, 3026
- Yang, S., Valenti, S., Sand, D., et al. 2017, GCN, **21531**
- Zevin, M., Kelley, L. Z., Nugent, A., et al. 2020, *ApJ*, **904**, 190
- Zevin, M., Kremer, K., Siegel, D. M., et al. 2019, *ApJ*, **886**, 4
- Zheng, Z., & Ramirez-Ruiz, E. 2007, *ApJ*, **655**, 1220