

Enumerative geometry of the mirror quintic

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Abstract: We evaluate the enumerative invariants of low degree on the mirror quintic threefold.

Mirror symmetry burst upon the mathematical scene with the famous computation made by Candelas, de la Ossa, Green, and Parkes [4] which in modern language proposed to count the number of rational curves of fixed degree on the quintic threefold using a technique from physics. In current language the “counts” are evaluations of Gromov–Witten invariants or Gopakumar–Vafa invariants, and the technique for counting these invariants has been expanded and extended in numerous ways. From the physics point of view, the computation was made on a closely related algebraic variety – the mirror quintic.

In a recent physics paper [10], a new technique was proposed for explicitly evaluating the Gromov–Witten or Gopakumar–Vafa invariants of the mirror quintic itself, not just the quintic. Such explicit evaluations seem rather daunting, since the answer will be a function of 101 variables. One aspect of [10] is to use two variables only and arrive at a more reasonable count.

Upon request of the authors of [10], the present authors worked out the enumerative geometry of the mirror quintic. We have made a conjecture about the Mori cone which comes with a plausibility argument rather than a proof. However, independent of the truth of that conjecture, we are able to explicitly evaluate the Gromov–Witten or Gopakumar–Vafa invariants of low degree for the mirror quintic, involving all 101 variables.

It gives us great pleasure to dedicate this paper to our mentor and friend Herb Clemens. Herb’s work on rational curves on Calabi–Yau threefolds [5, 6, 7] gave an inspiration and a foundation to much of our own work, including this paper.

1. The mirror quintic, its curves, and its divisors

The standard description of the mirror quintic is as a quotient of a quintic hypersurface in \mathbb{P}^4 . Let t_1, \dots, t_5 be homogeneous coordinates. The standard

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equation is then

$$(1) \quad t_1^5 + t_2^5 + t_3^5 + t_4^5 + t_5^5 = 5\psi t_1 t_2 t_3 t_4 t_5.$$

We are taking the quotient of \mathbb{P}^4 by the coordinate-wise action of

$$(2) \quad (\mathbb{Z}_5)^4 = \{(\zeta_1, \zeta_2, \zeta_3, \zeta_4, \zeta_5) \in (\mathbb{Z}_5)^5 \mid \prod_{j=1}^5 \zeta_j = 1\}.$$

The diagonal subgroup acts trivially on \mathbb{P}^4 , so the effective action is by $(\mathbb{Z}_5)^3$.

We will employ an alternate description (due to Batyrev [3]) of the mirror quintic as a singular complete intersection in \mathbb{P}^5 . For this purpose we begin by describing a basis for monomials invariant under the finite group action. Our choice of notation may appear awkward at first, but its utility will become apparent later on.

Let

$$(3) \quad \begin{array}{ll} u = t_1 t_2 t_3 t_4 t_5 & v = t_1^5 \\ w = t_2^5 & x = t_3^5 \\ y = t_4^5 & z = t_5^5. \end{array}$$

Then the equations defining the mirror quintic can be written

$$(4) \quad v + w + x + y + z = 5\psi u$$

$$(5) \quad v w x y z = u^5,$$

which describes the mirror quintic as a $(5, 1)$ complete intersection in \mathbb{P}^5 . We denote this singular model of the mirror quintic by Y_ψ , and refer to the coordinates appearing in (4) and (5) as *Batyrev coordinates*. In fact, this expresses the mirror quintic as a hypersurface (described by (4)) inside a singular toric variety (described by (5)).

There are singularities along lines defined by the simultaneous vanishing of u and two out of $\{v, w, x, y, z\}$ (as well as eq. (4)); there are ten such lines. The transverse singularity at a general point of such a line is A_4 .

There are more complicated singularities at the points defined by the simultaneous vanishing of u and three out of $\{v, w, x, y, z\}$ (and again imposing eq. (4)); there are ten such points. The singularity at each such point takes the local form $\mathbb{C}^3/(\mathbb{Z}_5)^2$, with non-isolated singularities propagating along the coordinate axes in \mathbb{C}^3 having transverse singularity of the form $\mathbb{C}^2/\mathbb{Z}_5 = A_4$.

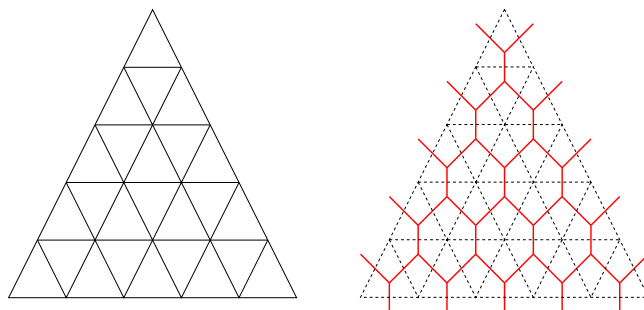


Figure 1: A triangulation specifying a crepant resolution of $\mathbb{C}^3/(\mathbb{Z}_5)^2$, and the corresponding dual graph.

We also observe the feature of five \mathbb{P}^2 's contained within the mirror quintic: each is given by the simultaneous vanishing of u and one of the coordinates $\{v, w, x, y, z\}$, as well as eq. (4).

A concrete procedure for blowing up the singularities was given in [13, Appendix B], but we will use a different birational model described in [14, Section 4].¹ The singularity $\mathbb{C}^3/(\mathbb{Z}_5)^2$ can be described torically, and we display the toric data for the resolution we use in the left half of Figure 1. We will denote the resulting Calabi-Yau threefold by X_ψ . The fully-resolved mirror quintic family is parametrized by $z = \psi^{-5}$, and is smooth for $z \in \mathbb{P}^1 - \{0, 1, \infty\}$. The large radius point is given by $z = 0$, while $z = 1$ is the conifold point and $z = \infty$ is the orbifold point. There is a birational contraction map $\rho : X_\psi \rightarrow Y_\psi$.

The toric data is “dual” to a more geometric description involving curves and surfaces, and we display the corresponding dual graph in the right half of Figure 1. There are six complete divisors within the resolution, represented by hexagons in the dual graph. There are also four incomplete divisors along each edge which represent components of the resolutions of the one-dimensional singular loci, and three incomplete divisors at the vertices of the figure which represent the \mathbb{P}^2 's identified above.

There are also 30 compact curves within the resolution, represented by segments in the dual graph which have both ends meeting a divisor. We can also see portions of other curves which we will describe later.

¹Note that our use of the term “the mirror quintic” is not quite correct – we must specify the birational model, and different models will have different Kähler cones. In this paper, we use the model corresponding to the resolution of singularities specified in [14, Section 4].

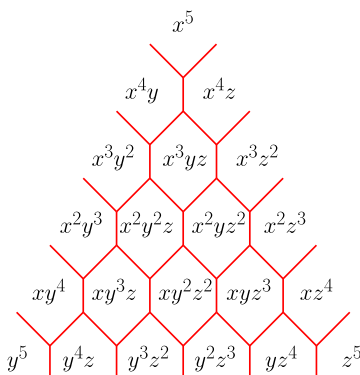


Figure 2: Labeling the divisors with monomials.

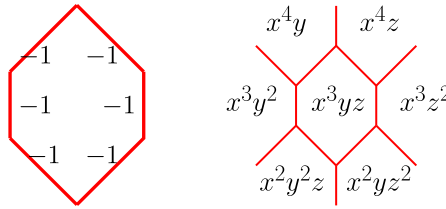
In order to keep track of all of the divisors and curves, we introduce the following notation. Each point with stabilizer $\mathbb{C}^5/(\mathbb{Z}_5)^2$ is associated with three of the variables $\{v, w, x, y, z\}$ and we will use quintic monomials in those three variables to label the corresponding divisors. This is illustrated in Figure 2 for the variables x, y, z . For a given monomial \mathbf{m} , we call the corresponding divisor $D_{\mathbf{m}}$.

Note that every monomial which involves at most three of $\{v, w, x, y, z\}$ corresponds to a divisor, and monomials involving fewer than three of those variables will show up in more than one of the toric diagrams. There are 105 divisors of this form.

Each compact curve in our diagram is the intersection of precisely two compact divisors $D_{\mathbf{m}}$ and $D_{\mathbf{n}}$ and we label the curve by² $\gamma_{\mathbf{m}, \mathbf{n}}$. In addition for each variable \mathbf{s} from our set of variables $\{v, w, x, y, z\}$, we let $\ell_{\mathbf{s}}$ denote a line in $D_{\mathbf{s}^5} \cong \mathbb{P}^2$. (Note that $\ell_{\mathbf{s}}$ is also represented by the curve $D_{\mathbf{s}^5} \cap D_{\mathbf{s}^4 \mathbf{t}}$ for any variable \mathbf{t} distinct from \mathbf{s} .) Finally, given two variables \mathbf{s} and \mathbf{t} from the set $\{v, w, x, y, z\}$, we let $\sigma_{\mathbf{s}, \mathbf{t}}$ be the intersection of $D_{\mathbf{s}^3 \mathbf{t}^2}$ and $D_{\mathbf{s}^2 \mathbf{t}^3}$. We will conjecture below that the 315 curve classes of the form $[\gamma_{\mathbf{m}, \mathbf{n}}]$, $[\ell_{\mathbf{s}}]$, and $[\sigma_{\mathbf{s}, \mathbf{t}}]$ are precisely the generators of all extremal rays and collectively generate the Mori cone of the mirror quintic.

There are a few more curves which should be discussed. For each pair of variables such as $\{y, z\}$, resolving the corresponding one-dimensional singular locus produces four divisors which in the example are $D_{y^4 z}$, $D_{y^3 z^2}$, $D_{y^2 z^3}$, and D_{yz^4} . Each divisor appears over three of the $\mathbb{C}^3/(\mathbb{Z}_5)^2$ points (corresponding

²Our convention is to only use the notation $\gamma_{\mathbf{m}, \mathbf{n}}$ in case the two divisors $D_{\mathbf{m}}$ and $D_{\mathbf{n}}$ are both compact in the inverse image of some $\mathbb{C}^5/(\mathbb{Z}_5)^2$ point.


 Figure 3: A general dP_3 , and a particular dP_3 .

to the triples of variables which contain $\{y, z\}$) and they are labeled the same way in each instance. These divisors are all ruled (by the exceptional curves resolving the non-isolated singularities), and we label the fiber φ of the ruling with the same monomial as that of the divisor (e.g., the ruling on D_{y^4z} is φ_{y^4z}). There are also curves of intersection of adjacent pairs of such divisors, labeled by a pair of monomials, e.g., the intersection of D_{y^4z} and $D_{y^3z^2}$. The only extremal curves among those are the lines ℓ_s, ℓ_t and the “middle” sections $\sigma_{s,t}$.

We determine the relations among the various curve classes by analyzing the structure of the divisors. Each compact divisor of the form $D_{\mathbf{m}}$ where the monomial \mathbf{m} involves three variables (which is represented by a hexagon in the dual graph) is isomorphic to dP_3 , i.e., the blowup of \mathbb{P}^2 in three non-collinear points. Each of the six toric curves on the divisor has self-intersection -1 , as illustrated in the left half of Figure 3. In the right half of that same figure, we have chosen a particular example so as to provide notation for the relations on curve classes that we are about to describe.

A dP_3 has three rulings, each having two singular fibers. Each singular fiber in one of these rulings is the union of two adjacent -1 curves. For example, the ruling whose general fibers run from southwest to northeast in the right side of Figure 3 has special fibers $\gamma_{x^4y, x^3yz} \cup \gamma_{x^3y^2, x^3yz}$ and $\gamma_{x^3yz, x^3z^2} \cup \gamma_{x^3yz, x^2y^2z}$. This leads to a relation among the curve classes which we denote by $R_{x^3yz}^{xy,z}$, namely

$$(6) \quad R_{x^3yz}^{xy,z} = [\gamma_{x^4y, x^3yz}] + [\gamma_{x^3y^2, x^3yz}] - [\gamma_{x^3yz, x^3z^2}] - [\gamma_{x^3yz, x^2y^2z}].$$

The subscript on R denotes the divisor which generated the relation. The superscript identifies the ruling in the following way: the parameter space for the ruling proceeds from the xy edge of the toric diagram to the z vertex of the toric diagram. In this way, we can list all such relations in the xy, z

direction, as follows:

$$\begin{aligned}
 R_{x^3yz}^{xy,z} &= [\gamma_{x^4y,x^3yz}] + [\gamma_{x^3y^2,x^3yz}] - [\gamma_{x^3yz,x^3z^2}] - [\gamma_{x^3yz,x^2yz^2}] \\
 R_{x^2y^2z}^{xy,z} &= [\gamma_{x^3y^2,x^2y^2z}] + [\gamma_{x^2y^3,x^2y^2z}] - [\gamma_{x^1y^2z,x^2yz^2}] - [\gamma_{x^2y^2z,x^2yz^2}] \\
 R_{xy^3z}^{xy,z} &= [\gamma_{x^2y^3,xy^3z}] + [\gamma_{xy^4,xy^3z}] - [\gamma_{xy^3z,xy^2z^2}] - [\gamma_{xy^3z,y^3z^2}] \\
 R_{x^2yz^2}^{xy,z} &= [\gamma_{x^3yz,x^2yz^2}] + [\gamma_{x^2y^2z,x^2yz^2}] - [\gamma_{x^2yz^2,x^2z^3}] - [\gamma_{x^2yz^2,xy^3z^3}] \\
 R_{xy^2z^2}^{xy,z} &= [\gamma_{x^2y^2z,xy^2z^2}] + [\gamma_{xy^3z,xy^2z^2}] - [\gamma_{xy^2z^2,xy^3z^3}] - [\gamma_{xy^2z^2,y^2z^3}] \\
 R_{xyz^3}^{xy,z} &= [\gamma_{x^2yz^2,xyz^3}] + [\gamma_{xy^2z^2,xyz^3}] - [\gamma_{xyz^3,xz^4}] - [\gamma_{xyz^3,yz^4}]
 \end{aligned}
 \tag{7}$$

Similarly, we can relate the fiber in the ruling on one of the “edge” divisors to a degenerate fiber within the toric diagram. The ones in the same xy, z direction are:

$$\begin{aligned}
 R_{x^4y}^{xy,z} &= [\varphi_{x^4y}] - [\gamma_{x^4y,x^4z}] - [\gamma_{x^4y,x^3yz}] \\
 R_{x^3y^2}^{xy,z} &= [\varphi_{x^3y^2}] - [\gamma_{x^3y^2,x^3yz}] - [\gamma_{x^3y^2,x^2yz^2}] \\
 R_{x^2y^3}^{xy,z} &= [\varphi_{x^2y^3}] - [\gamma_{x^2y^3,x^2yz^2}] - [\gamma_{x^2y^3,xy^3z^3}] \\
 R_{xy^4}^{xy,z} &= [\varphi_{xy^4}] - [\gamma_{xy^4,xy^3z}] - [\gamma_{xy^4,y^4z}]
 \end{aligned}
 \tag{8}$$

Other relations of the same type as in eqs. (7) and (8) can be generated by substituting any three variables for x, y, z . This gives 30 versions of eq. (7).

The compact divisor D_{x^3yz} also allows us to exhibit a relation among relations, since the three specified relations among six -1 curves are not linearly independent. The “syzygy” is easily seen to be

$$R_{x^3yz}^{xy,z} + R_{x^3yz}^{yz,x} + R_{x^3yz}^{zx,y} = 0.
 \tag{9}$$

A similar syzygy exists for each compact divisor. We thus conclude that of the 180 relations of type (7) among the 300 γ curves, only 120 are linearly independent.

We can implement these relations as illustrated in Figure 4. In that figure, we have selected a subset of the compact curves (marked with a purple dash) which constitute 4 curves on each compact divisor. In addition, for the unmarked compact curves, we have used a green arrow to designate which divisor’s relations should be used to eliminate that curve from the spanning set. Each compact divisor has two green arrows within it, indicating that two of the relations belonging to that divisor are used in this elimination process. Those relations are linearly independent.

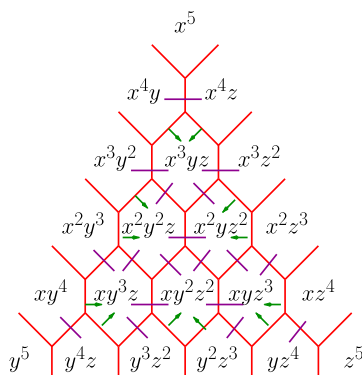


Figure 4: Basis for compact curves.

The divisors on the edges, such as D_{x^4y} also give rise to relations among the γ curves, since three expressions for $[\varphi_{x^4y}]$ are derived from relations of the type in eq. (8). The ones for φ_{x^4y} are

$$\begin{aligned}
 R_{x^4y}^{xy,v} &= [\varphi_{x^4y}] - [\gamma_{x^4y,x^4v}] - [\gamma_{x^4y,x^3yv}] \\
 R_{x^4y}^{xy,w} &= [\varphi_{x^4y}] - [\gamma_{x^4y,x^4w}] - [\gamma_{x^4y,x^3yw}] \\
 R_{x^4y}^{xy,z} &= [\varphi_{x^4y}] - [\gamma_{x^4y,x^4z}] - [\gamma_{x^4y,x^3yz}]
 \end{aligned}
 \tag{10}$$

and by subtraction we obtain two independent relations among the γ curves. Since this can be done for each of 4 divisors along each of 10 edges, there are 80 relations in total of this type. These relations are independent.

Thus, we find a total of 200 linear relations among the 300 compact γ curves, leaving 100 independent classes.

We now consider the divisors on the edges of the toric diagram, which form the resolution of the A_4 loci. We take as an example the sequence of divisors D_{y^4z} , $D_{y^3z^2}$, $D_{y^2z^3}$, D_{yz^4} . The first two divisors in any such sequence are illustrated in Figure 5, in other words, the divisor on the left is D_{y^4z} and the divisor on the right is $D_{y^3z^2}$. To the left of D_{y^4z} is D_{y^5} and to the right of $D_{y^3z^2}$ is $D_{y^2z^3}$. (There is also a third singular fiber in the ruling on each divisor; these are indicated with narrow lines.)

As we mentioned earlier, the divisors D_{y^5} and D_{y^4z} meet on ℓ_y which has self-intersection 1 on D_{y^5} . By the adjunction formula, the self-intersection of this curve on D_{y^4z} (where it is a section of the ruling) must be -3 . It follows that D_{y^4z} is a blowup of the Hirzebruch surface \mathbb{F}_3 , and since there are three singular fibers, the disjoint section $D_{y^4z} \cap D_{y^3z^2}$ must have self-intersection

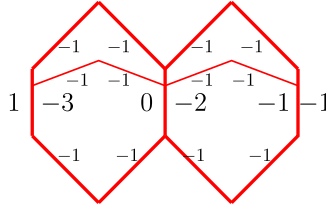


Figure 5: Two adjacent components of the resolution of the A_4 locus.

0 on D_{y^4z} . In fact, we can say more: by the description as a blowup of a Hirzebruch surface³

$$(11) \quad [D_{y^4z} \cap D_{y^3z^2}] = [\ell_y] + 3[\varphi_{y^4z}] - [\gamma_{vy^3z,y^4z}] - [\gamma_{wy^3z,y^4z}] - [\gamma_{xy^3z,y^4z}],$$

which can also be written

$$(12) \quad [D_{y^4z} \cap D_{y^3z^2}] = [\ell_y] + [\gamma_{vy^4,y^4z}] + [\gamma_{wy^4,y^4z}] + [\gamma_{xy^4,y^4z}].$$

Using adjunction again, the self-intersection of $D_{y^4z} \cap D_{y^3z^2}$ on $D_{y^3z^2}$ must be -2 , which tells us that $D_{y^3z^2}$ is the blowup of a Hirzebruch surface \mathbb{F}_2 in three points. The disjoint section $\sigma_{y,z} = D_{y^3z^2} \cap D_{y^2z^3}$ must thus have self-intersection -1 on $D_{y^3z^2}$, and again we get a relation among curve classes:

$$(13) \quad [\sigma_{y,z}] = [D_{y^4z} \cap D_{y^3z^2}] + 2[\varphi_{y^3z^2}] - [\gamma_{vy^2z^2,y^3z^2}] - [\gamma_{wy^2z^2,y^3z^2}] - [\gamma_{xy^2z^2,y^3z^2}],$$

which can also be written

$$(14) \quad [\sigma_{y,z}] = [D_{y^4z} \cap D_{y^3z^2}] + [\gamma_{vy^3z,y^3z^2}] + [\gamma_{wy^3z,y^3z^2}] - [\gamma_{xy^2z^2,y^3z^2}]$$

(a little less naturally this time, since we must single out one of the three singular fibers). Combining these two gives

$$(15) \quad \begin{aligned} [\sigma_{y,z}] &= [\ell_y] + [\gamma_{vy^4,y^4z}] + [\gamma_{wy^4,y^4z}] + [\gamma_{xy^4,y^4z}] \\ &\quad + [\gamma_{vy^3z,y^3z^2}] + [\gamma_{wy^3z,y^3z^2}] - [\gamma_{xy^2z^2,y^3z^2}]. \end{aligned}$$

³If we blowup the Hirzebruch surface \mathbb{F}_n at k points which are contained in a section $\bar{\sigma}$ of the \mathbb{P}^1 -fibration (with fiber $\bar{\varphi}$) having $(\bar{\sigma})_{\mathbb{F}_n}^2 = n \geq 0$, then the blowup contains the proper transform φ of $\bar{\varphi}$, the proper transform σ_∞ of the section $\bar{\sigma}_\infty$ with $(\bar{\sigma}_\infty)_{\mathbb{F}_n}^2 = -n \leq 0$, the proper transform σ of $\bar{\sigma}$, and the exceptional divisors e_1, \dots, e_k . The basic homology relation $[\bar{\sigma}] = [\bar{\sigma}_\infty] + n[\bar{\varphi}]$ on \mathbb{F}_n pulls back to a homology relation $[\sigma] + [e_1] + \dots + [e_k] = [\sigma_\infty] + n[\varphi]$ on the blowup. Note that $[\varphi] - [e_j]$ is represented by an effective (-1) -curve e'_j , and that $[e_j] + [e'_j]$ is homologous to $[\varphi]$.

We thus get 20 additional relations among our 315 curves. However, at most 14 of the 20 relations can be linearly independent, which can be seen from Batyrev's calculation of b_4 of the quintic mirror [3].

The conclusion (from Batyrev's calculation) is that there are precisely 101 independent curve classes generated by our explicit set of 315 curves.

We now consider relations among the divisors. First, for the divisors whose monomials are chosen from $\{x, y, z\}$, we can determine some linear combinations which have intersection number 0 with all 30 compact curves lying over the corresponding singular point. There are three linearly independent combinations, which we label as $D^{x,yz}$, $D^{y,xz}$, and $D^{z,xy}$, and which are described by means of coefficient arrays whose shape matches that of Figure 2, as follows:

$$\begin{aligned}
 D^{y,xz} &= \begin{array}{cccccc} & & & 0 & & \\ & & & 1 & & 0 \\ & & 2 & 1 & 0 & \\ & 3 & 2 & 1 & 0 & \\ 4 & 3 & 2 & 1 & 0 & \\ 5 & 4 & 3 & 2 & 1 & 0 \end{array}, \\
 D^{z,xy} &= \begin{array}{cccccc} & & & 0 & 1 & \\ & & 0 & 1 & 2 & \\ & 0 & 1 & 2 & 3 & \\ 0 & 1 & 2 & 3 & 4 & 5 \end{array}, \\
 D^{x,yz} &= \begin{array}{cccccc} & & & 5 & & \\ & & & 4 & 4 & \\ & & 3 & 3 & 3 & \\ & 2 & 2 & 2 & 2 & \\ 1 & 1 & 1 & 1 & 1 & \\ 0 & 0 & 0 & 0 & 0 & 0 \end{array}.
 \end{aligned}$$

Note that ℓ_x meets D_{x^5} with intersection number -3 , and also meets D_{x^4y} and D_{x^4z} , each with intersection number 1. It follows that $D^{x,yz}$ meets ℓ_x with intersection number -7 , meets ℓ_y once, and meets ℓ_z once.

This calculation suggests how to find a divisor which has intersection number 0 with all curves $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$. Let

$$(16) \quad D^x = \sum_{\substack{\text{monomials } \mathbf{m} \text{ with} \\ \text{at most three variables}}} \nu_x(\mathbf{m}) D_{\mathbf{m}},$$

where $\nu_x(\mathbf{m})$ denotes the exponent of the variable x in the monomial \mathbf{m} . Since the coefficients of D^x in any triangle containing x have the same pattern as $D^{x,yz}$, the intersection with all of the compact curves over such a point is 0. And since the coefficients of D^x are all zero in any triangle not containing x , the intersection number with each of the compact curves over such a point is also 0.

Also, D^x meets ℓ_x at the components D_{x^5} and D_{x^4*} , so the total intersection number is

$$D^x \cdot \ell_x = 5 \cdot (-3) + 4 \cdot 4 \cdot 1 = 1.$$

On the other hand, for any other variable such as y , the only component of D^x with a nonzero coefficient which meets ℓ_y is D_{x^4y} . Thus, $D^x \cdot \ell_y = 1$ as well (and similarly for ℓ_v , ℓ_w , and ℓ_z).

It follows that $D^x \cdot \gamma = 0$, $D^x \cdot \ell = 1$ for all γ curves and lines ℓ , respectively.

Now, the same thing happens for D^v , D^w , D^y , and D^z , so the difference of any two is numerically trivial. This produces four linear relations among the 105 divisors $D_{\mathbf{m}}$, giving a space of dimension 101.

Note that D^x (and the others) can be identified with a hyperplane section H of the original singular model of the mirror quintic.

We would like to write down one more divisor, which has intersection number 1 with all γ curves, and intersection number 0 with all lines ℓ . To this end, we consider the following coefficient array

				16						
				12		12				
			10		9		10			
		10		8		8		10		
	12		9		8		9		12	
16		12		10		10		12		16

The corresponding divisor has intersection number 1 with each compact $\gamma_{\mathbf{m}}$ where \mathbf{m} involves the variables $\{x, y, z\}$. This is verified by checking total intersection number in various sub-diamonds of the coefficient array, such as $16 + 9 - 12 - 12$ which corresponds to intersections with γ_{x^4y,x^4z} .

To make a global version of this, we define a function a on monomials as follows. The nonzero exponents in the monomial \mathbf{m} determine a partition of 5 into at most three elements, and the value of the function only depends on the partition, as follows:

Partition	5	4 + 1	3 + 2	3 + 1 + 1	2 + 2 + 1
$a(\mathbf{m})$	16	12	10	9	8

Then the divisor

$$\mathbb{D} = \sum_{\substack{\text{monomials } \mathbf{m} \text{ with} \\ \text{at most three variables}}} a(\mathbf{m}) D_{\mathbf{m}}$$

has the property that $\mathbb{D} \cdot \gamma = 1$ for all γ . Note that ℓ_x meets D_{x^5} of coefficient 16 with intersection number -3 , and meets 4 divisors of the form D_{x^4*} each of coefficient 12 and with intersection number 1. Thus,

$$\mathbb{D} \cdot \ell_x = 16 \cdot (-3) + 4 \cdot 12 \cdot 1 = 0$$

and similarly for any of the $\mathbb{D} \cdot \ell$.

Having identified the curves $\ell_{\mathbf{t}}$ and $\sigma_{\mathbf{s}, \mathbf{t}}$ for $\mathbf{s}, \mathbf{t} \in \{w, v, x, y, z\}$, and the curves $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$ above, we conjecture that these generate the Mori cone of numerical equivalence classes of effective curves on X_ψ .

Conjecture 1. *The Mori cone M of X_ψ is generated by the classes of the curves $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s}, \mathbf{t}}$, and $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$.*

In the rest of this paper, we assume this conjecture whenever necessary.

If C is a curve on X_ψ , we define the *degree* of C to be the degree of its image $\overline{C} := \rho(C) \subset Y_\psi \subset \mathbb{P}^5$. We will show presently in Lemma 2 that if $C \subset X_\psi$ has degree zero, then its class $[C]$ is in the cone generated by the classes of the curves $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$. Furthermore, in Section 2 we will show that if $C \subset X_\psi$ has degree 1, then its class $[C]$ is in the cone generated by the classes of the curves $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s}, \mathbf{t}}$, and $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$. In other words, the conjecture is true for curves of degree at most 1.

Lemma 2. *The class of any curve in an exceptional divisor for the birational contraction ρ is in the cone generated by the classes of the curves $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s}, \mathbf{t}}$, and $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$.*

In particular, the conjecture is true for curves of degree 0.

Proof. We give separate arguments for the exceptional divisors which contract to a point and to a curve.

The exceptional divisors contracting to a point are all dP_3 s. It is well-known that the Mori cone of dP_3 is generated by the classes of its six exceptional curves of the first kind. But these exceptional curves are precisely the γ curves contained in that dP_3 .

The exceptional divisors contracting to a curve are the components of one of the A_4 resolutions. By symmetry we need only analyze the curves in D_{y^4z} and $D_{y^3z^2}$.

Now D_{y^4z} is identified with the blow up of the Hirzebruch surface \mathbb{F}_3 at three general points, the exceptional curves being identified with the curves γ_{y^4z, y^3zs} for $s \in \{v, w, x\}$. Using the blowup description, the Mori cone is generated by the -3 section, the proper transform of any $+3$ section containing the points being blown up, the proper transform of the fibers, and the exceptional curves. We need only check each curve in turn. The -3 section is identified with the curve ℓ_y . We have already noted that the exceptional curves are all γ curves. The proper transform of the fiber which meets γ_{y^4z, y^3zs} is just γ_{y^4z, y^4s} , also a γ curve. The proper transform of a particular $+3$ section is identified with $D_{y^4z} \cap D_{y^3z^2}$, which is in the cone generated by ℓ_y and γ curves by (12).

Next, we observe that $D_{y^3z^2}$ is similarly identified with the blow up of the Hirzebruch surface \mathbb{F}_2 at three points on a $+2$ section, the exceptional curves being identified with the curves $\gamma_{y^3z^2, y^2z^2s}$ for $s \in \{v, w, x\}$. Using the blowup description, the Mori cone is generated by the -2 section, the proper transform of the unique $+2$ section containing the points being blown up, the proper transform of the fibers, and the exceptional curves. We need only check each curve in turn. The -2 section is identified with the curve $D_{y^4z} \cap D_{y^3z^2}$ already considered above. We have already noted that the exceptional curves are all γ curves. The proper transform of the fiber which meets γ_{y^4z, y^3zs} is just γ_{y^4z, y^4s} , also a γ curve. The proper transform of the $+2$ section is identified with $\sigma_{y,z}$.

We can give a plausibility argument in favor of our conjecture for arbitrary degree as follows. Let C be any irreducible curve on the mirror quintic, and let \overline{C} be its image on the singular model as above. If \overline{C} is contained in the singular locus, then C must lie in one or more of the exceptional divisors of the blowup. This case is handled by Lemma 2. If \overline{C} is not contained in the singular locus then it meets the singular locus in finitely many points. We want to claim – and this is the gap in the argument – that \overline{C} can be deformed away from the singular locus. The family C_t , when lifted to the smooth model, will have a limit C_0 consisting of C together with some curves which are contained in exceptional divisors. As mentioned earlier, those latter curves are in the cone generated by known curves.

We next let $\mathbf{H} = \mathbb{Z}^2 = \mathbb{Z} \cdot \ell \oplus \mathbb{Z} \cdot \gamma$ and project curve classes to H via

$$(17) \quad \pi : M \rightarrow \mathbf{H}, \quad \pi(C) = (C \cdot D^x) \ell + (C \cdot \mathbb{D}) \gamma.$$

The notation for the basis for \mathbf{H} has been chosen so that $\pi(\ell_t) = \ell$ and $\pi(\gamma_{\mathbf{m}_1, \mathbf{m}_2}) = \gamma$. We also observe that if $\pi(\beta) = m\ell + n\gamma$, then $\rho_*(\beta)$ is the class of a degree m curve in $Y_\psi \subset \mathbb{P}^5$.

We compute and record the following curve classes, which are computed from (17) and the definitions of D^x and \mathbb{D} .

$$(18) \quad \pi(\gamma_{t^5, t^4 u}) = \ell$$

$$(19) \quad \pi(\gamma_{t^4 u, t^3 u^2}) = \ell + 3\gamma$$

$$(20) \quad \pi(\sigma_{t, u}) = \ell + 4\gamma$$

where t and u are distinct elements of $\{v, w, x, y, z\}$.

Our interest is in defining and computing Gopakumar-Vafa invariants of classes in \mathbf{H} in order to compare to the calculations of [10]. \square

Lemma 3. *Fix $\delta \in \mathbb{Z}^2$. Then the set of curve classes $\beta \in M$ with $\pi(\beta) = \delta$ is finite.*

Proof. We write $\delta = m\ell + n\gamma$ and may assume $m \geq 0$, $n \geq 0$, and $(m, n) \neq (0, 0)$, otherwise the statement of the lemma is trivial. We can write

$$(21) \quad \beta = \sum m_{\mathbf{t}} \ell^{\mathbf{t}} + \sum n_{i,j} \gamma_{\mathbf{m}_1, \mathbf{m}_2} + \sum p_{\mathbf{s}, \mathbf{t}} \sigma_{\mathbf{s}, \mathbf{t}}$$

with $m_{\mathbf{t}}, n_{i,j}, p_{\mathbf{s}, \mathbf{t}} > 0$, suppressing from the notation the implicit set of 3 variables needed to define $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$. In writing (21), we have adopted the convention that we always use ℓ in place of a class of the form $\gamma_{t^5, t^4 u}$. Then $\pi(\beta) = \delta$ implies $\sum m_{\mathbf{t}} + \sum n_{i,j} + \sum p_{\mathbf{s}, \mathbf{t}} = m$ and $3 \sum n_{i,j} + 4 \sum p_{\mathbf{s}, \mathbf{t}} = n$, which has only finitely many solutions for $m_{\mathbf{t}}, n_{i,j}$ and $p_{\mathbf{s}, \mathbf{t}}$.

Lemma 3 ensures that the two parameter projection of [10] makes mathematical sense: their invariants are simply a sum of finitely many Gopakumar-Vafa invariants.

For $\delta \in \mathbf{H}$ we define

$$(22) \quad n_{m,n}^g = \sum_{\pi(\beta) = m\ell + n\gamma} n_{\beta}^g,$$

which is a finite sum by Lemma 3. In (22), n_{β}^g is the genus g GV invariant associated with the curve class $\beta \in H_2(X_{\psi}, \mathbb{Z})$. The GV invariants will be discussed in more detail in Section 2. \square

2. Enumerative geometry

In this section, we calculate some genus 0 Gromov-Witten invariants. In principle we can algorithmically compute the genus 0 Gromov-Witten invariants for the full 101 parameter model using the toric mirror theorem [8], but we

leave that for future work. It is conceivable that there may be a way to modify the toric mirror theorem to apply directly to our two-parameter family. That would certainly be worth doing.

2.1. Generalities

In lieu of calculating genus 0 Gromov-Witten invariants, we calculate the genus 0 Gopakumar-Vafa invariants [9], which is easier in examples. Let $\beta \in H^2(X_\psi, \mathbb{Z})$. Mathematically, we follow [11] by defining the genus 0 Gopakumar-Vafa invariant n_β^0 as the Donaldson-Thomas invariant of the moduli space of stable sheaves F with $[F] = \beta$ (which is known to be independent of the choice of polarization used to define stability). This definition is consistent with the mathematical definition of higher genus Gopakumar-Vafa invariants n_β^g given later in [12].

The Gromov-Witten invariants N_β^0 are related to the n_β^0 by the Aspinwall-Morrison formula [2, 15]⁴

$$(23) \quad N_\beta^0 = \sum_{k|\beta} \frac{n_{\beta/k}^0}{k^3}$$

so we content ourselves with the calculation of the n_β^0 .

Furthermore, in all of our cases, the stable sheaves are of the form \mathcal{O}_C for rational curves C and so the moduli spaces of sheaves is just the moduli space of curves. Even better, these moduli spaces M_β are smooth, and the Donaldson-Thomas invariant is simply

$$(24) \quad n_\beta^0 = (-1)^{\dim(M_\beta)} e(M_\beta),$$

where $e(M_\beta)$ is the topological euler characteristic.

2.2. Degrees 0 and 1

We now record our results for Gopakumar-Vafa invariants in the classes $m\ell + n\gamma$, with $m = 0, 1$. These are the classes which project to points or lines in

⁴A proof of the multiple cover formula has only been written down for isolated curves. However, in our cases, it can be shown by symplectic techniques that the $\bar{\partial}$ operator can be deformed so that there are only finitely many embedded pseudoholomorphic curves, all isolated, and the number is given by the associated GV invariant, with signed counts corresponding to orientation choices. This is easy for primitive curve classes since the moduli spaces are smooth and projective, but with more care the case of 2γ below can also be handled.

the singular model Y_ψ .

$$\begin{aligned} (25) \quad n_\gamma^0 &= 300 \\ (26) \quad n_{2\gamma}^0 &= -440 \\ (27) \quad n_\ell^0 &= 15 \\ (28) \quad n_{\ell+\gamma}^0 &= -60 \\ (29) \quad n_{\ell+2\gamma}^0 &= 155 \end{aligned}$$

We begin with the classes of the form $n\gamma$. By Lemma 2 and its proof, these curves can be constructed from the γ curves and the fibers of the divisors of the A_4 resolution.

γ . The curves of class γ contained in each rational ruled surface are just the components of the reducible fibers. The curves of class γ contained in each dP_3 are the 6 curves depicted in Figure 3. So we simply count all of these curves.

Each 2-simplex depicted in Figure 1 corresponds to one of the 10 singular points. We count 30 curves of class γ corresponding to each 2-simplex. Thus $n_\gamma^0 = 10 \cdot 30 = 300$.

2γ . Again inspecting Figure 3, we see that any intersecting configuration of two curves of type γ must lie in exactly one of the dP_3 s or the ruled surfaces lying over a singular curve. Since each curve γ has self-intersection -1 in the surface just described, the union $\gamma_1 + \gamma_2$ is a rational curve of self-intersection 0 in a rational surface, hence it moves in a pencil. The corresponding contribution to the GV invariant is $(-1)^1 e(\mathbb{P}^1) = -2$ by (24).

There are 3 such pairs of intersecting -1 curves for each of the 6 dP_3 's over each of the 10 singular points.

There are 4 ruled surfaces over each to the 10 singular curves.

Thus, $n_{2\gamma}^0 = (-2)(10 \cdot 3 \cdot 6 + 10 \cdot 4) = -440$.

We now turn to the curves of the form $\ell + n\gamma$. Since all of these project to lines in Y_ψ , we start by identifying the lines in Y_ψ .

Lemma 4. *Any line in Y_ψ is contained in one of the five \mathbb{P}^2 's defined by the simultaneous vanishing of u and one of the coordinates $\{v, w, x, y, z\}$, as well as eq. (4).*

Proof. The proof is inspired by an argument in [1]. Let $\ell \subset Y_\psi$ be a line and suppose that ℓ is not contained in any of these \mathbb{P}^2 's. If ℓ is not contained in the hyperplane $u = 0$, then it intersects that hyperplane at a point $p \in \ell$. It follows from (5) that each of the five hypersurfaces $\mathbf{t} = 0$ must also contain p .

This is a contradiction, since it would imply that 6 hyperplanes with empty intersection all contain p .

Since u is identically 0 on ℓ , it follows from (5) that at least one of the \mathbf{t} must vanish identically on ℓ , which proves the lemma. \square

Now let C be any irreducible curve of degree 1, and we claim that its class is in the cone generated by the classes of the curves $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s},\mathbf{t}}$, and $\gamma_{\mathbf{m}_1,\mathbf{m}_2}$. In this case $\overline{C} = \rho(C) \subset Y_\psi$ is a line. It follows that C must either be contained in one of the exceptional divisors forming an A_4 resolution or one of the divisors $D_{\mathbf{s}^5}$. We have already shown the claim in the first case. In the second case, $[C] = [\ell_{\mathbf{s}}]$.

ℓ . The classes β which contribute to n_ℓ^0 project to a line in Y_ψ , which must lie in at least one of the five \mathbb{P}^2 's described in Lemma 4. The divisor D_{x^5} is the proper transform of one such \mathbb{P}^2 , and the lines in such a \mathbb{P}^2 have class ℓ . The lines are parametrized by a dual \mathbb{P}^2 and so contribute +3. There are five such \mathbb{P}^2 , giving $5 \cdot 3 = 15$.

We claim that there are no other contributions. The only remaining possibility is for the projection to lie in more than one \mathbb{P}^2 and thus be a singular line. In this case one component of the curve must be a section of one of the 4 ruled surfaces over the singular line. But from (18) and (19) we infer that the Mori cone of D_{x^4y} is generated by ℓ and γ , while the only curve of class ℓ in D_{x^4y} is γ_{x^5,x^4y} , which we have already considered as a line in $D_{x^5} \simeq \mathbb{P}^2$. Similarly, the Mori cone of $D_{x^3y^2}$ is generated by $\ell + 3\gamma$ and γ by (19) and (20), so $D_{x^3y^2}$ contains no curves of class ℓ . Hence, the final answer is $n_\ell^0 = 15$.

$\ell + \gamma$. As above, we must have a component which is a line in one of the 5 \mathbb{P}^2 's, and we look for all ways to attach a curve γ . A typically way is to take the line in D_{x^5} passing through the point $D_{x^5} \cap D_{x^4y} \cap D_{x^4z}$ (there is a pencil of such lines) and attach γ_{x^4y,x^4z} . All other cases are obtained by permutations of the underlying data. Such curves are associated with a partial flag $p \in \mathbb{P}^2$, where p is one of the 10 singular points and \mathbb{P}^2 is one of the 3 \mathbb{P}^2 's containing a given singular point. This gives $n_{\ell+\gamma}^0 = 10 \cdot 3 \cdot (-2) = -60$.

$\ell + 2\gamma$. Again, we must have a component which is a line in one of the 5 \mathbb{P}^2 's. We look for all ways to attach two curves of class γ or one curve of class 2γ . For the first case, we choose one of the five \mathbb{P}^2 's which we identify with some $D_{\mathbf{t}^5}$ and a pair of the 6 singular points contained in that \mathbb{P}^2 . This specifies a line in $D_{\mathbf{t}^5}$ containing the two singular points. Then we glue the curves of type γ which meet each of the singular points. This contributes $5 \cdot 15 = 75$.

We have already seen how the fibers of the ruled surfaces have class 2γ . A typical example is to glue a line $\ell \subset D_{x^5}$ containing a point $p \in \gamma_{x^5,x^4y}$

to a fiber φ_{x^4y} . All other cases are given by permutation. The data $p \in \ell$ is parametrized by the blowup of the dual \mathbb{P}^2 at the point γ_{x^5, x^4y} of this dual \mathbb{P}^2 . This blowup has euler characteristic 4 and contributes +4 to the GV invariant. We have such a pencil for each choice of one of the 5 \mathbb{P}^2 's and one of the 4 singular lines contained in each \mathbb{P}^2 . This contributes $5 \cdot 4 \cdot 4 = 80$.

Combining these cases, we get $n_{\ell+2\gamma}^0 = 75 + 80 = 155$.

2.3. Degree 2

Here we only have partial results, but we note that they already agree with parallel calculations in [10]. We compute the invariants of those curves of class $2\ell + n\gamma$ which lie inside the union of the toric divisors $D_{\mathbf{m}}$, and denote the contribution of these curves by $n_{2\ell+n\gamma}^{0,\text{toric}}$. However, since we do not have an analogue of Lemma 4 at our disposal, it is possible that $n_{2\ell+n\gamma}^{0,\text{toric}} \neq n_{2\ell+n\gamma}^0$.

We record our results.

$$(30) \quad n_{2\ell}^{0,\text{toric}} = -30$$

$$(31) \quad n_{2\ell+\gamma}^{0,\text{toric}} = 150$$

$$(32) \quad n_{2\ell+2\gamma}^{0,\text{toric}} = -500$$

2ℓ . In this class, we clearly have the conics in any of the five toric divisors $D_{\mathbf{s}^5}$, each of which is a \mathbb{P}^2 . Since conics are parametrized by \mathbb{P}^5 , the contribution is $5(-6) = -30$.

The argument that we used in the case $\beta = \ell$ shows that there are no other curves in $\cup D_{\mathbf{m}}$ representing $\beta = 2\ell$.

$2\ell + \gamma$. Analogous to the case of $\ell + \gamma$, our curve must be a connected union of one of the conics of class 2ℓ and a curve γ , to which is associated to one of 30 possible partial flags $p \in \mathbb{P}^2$. The linear system of conics in \mathbb{P}^2 containing a point p is a \mathbb{P}^4 . So the invariant is $30 \cdot 5 = 150$.

$2\ell + 2\gamma$. As in the situation of $\ell + 2\gamma$, we must have a conic $C \subset D_{\mathbf{s}^5}$ as a component. There are then two possibilities. Either C is glued to two distinct γ curves, each meeting $D_{\mathbf{s}^5}$ in one of the 6 singular points that it contains, or C is glued to a fiber of a ruled surface $D_{\mathbf{s}^4\mathbf{t}}$.

In the first case, we choose one of the five \mathbb{P}^2 and one of the 15 pairs of the 6 singular points in that \mathbb{P}^2 at which we glue a γ curve. The linear system of plane conics through a pair of points is a \mathbb{P}^3 . So the net contribution is $5 \cdot 15 \cdot (-4) = -300$.

In the second case, for each of the 5 divisors $D_{\mathbf{s}^5}$, we pick one of the 4 singular lines $\gamma_{\mathbf{s}^5, \mathbf{s}^4\mathbf{t}}$ contained in it. We get a curve of class $2\ell + 2\gamma$ by choosing

a point $p \in \gamma_{\mathbf{s}^5, \mathbf{s}^4 \mathbf{t}}$, taking a conic $C \subset D_{\mathbf{s}^5}$ containing p , and gluing it to the fiber $\phi_{\mathbf{s}^4 \mathbf{t}}$ containing p . This moduli space is a \mathbb{P}^4 bundle over \mathbb{P}^1 , contributing -10 . So the total contribution is $5 \cdot 4 \cdot (-10) = -200$.

Combining these two cases, we get $n_{2\ell+2\gamma}^{0, \text{toric}} = -300 - 200 = -500$.

3. Five parameter projection

In this section, we prove Conjecture 1 in a five-parameter projection. Let us first explain what this means.

It follows from Batyrev's calculation [3] (which used toric constructions to represent curve classes, as we have done) that the curves $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s}, \mathbf{t}}$, and $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$ identified in Conjecture 1 span $H_2(X_\psi, \mathbb{Q})$. Furthermore, the curves $\sigma_{\mathbf{s}, \mathbf{t}}$ are not needed by (15).

We let V be the \mathbb{Q} -vector space with generators $\ell_t, \sigma_{s,t}$, and γ_{m_1, m_2} . Then there is a natural surjective map $r : V \rightarrow H_2(X_\psi, \mathbb{Q})$. The kernel of r is spanned by the relations (7)–(15) together with the analogous relations obtained by substitutions of the variables, described in Section 1.

The group S_5 acts on Y_ψ by permuting the coordinates, and this action lifts naturally to an action on X_ψ . Let W be the \mathbb{Q} -vector space generated by the set of S_5 -orbits of the $\ell_t, \sigma_{s,t}$, and γ_{m_1, m_2} . There is a natural map $n : V \rightarrow W$. Since the set of relations spanning $\ker(r)$ just described is invariant under S_5 , we deduce a natural map

$$(33) \quad \rho : H_2(X_\psi, \mathbb{Q}) \rightarrow \mathbf{N} := W / (n(\ker r)).$$

Let $\mathbf{M} \subset \mathbf{N} \otimes \mathbf{R}$ be the cone spanned by the classes $\rho(\ell_{\mathbf{t}})$, $\rho(\sigma_{\mathbf{s}, \mathbf{t}})$, and $\rho(\gamma_{\mathbf{m}_1, \mathbf{m}_2})$. Recall the Mori cone $M \subset H_2(X_\psi, \mathbb{R})$.

Lemma 5. $\dim(\mathbf{N}) = 5$.

Proof. As noted above, we can ignore the $\rho(\sigma_{\mathbf{s}, \mathbf{t}})$. Clearly, all of the curves $\ell_{\mathbf{s}}$ are identified by the S_5 action. There are six S_5 orbits of curves $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$:

$$(34) \quad \gamma_{\mathbf{s}^4 \mathbf{t}, \mathbf{s}^4 \mathbf{u}}, \gamma_{\mathbf{s}^4 \mathbf{t}, \mathbf{s}^3 \mathbf{t} \mathbf{u}}, \gamma_{\mathbf{s}^3 \mathbf{t}^2, \mathbf{s}^3 \mathbf{t} \mathbf{u}}, \gamma_{\mathbf{s}^3 \mathbf{t}^2, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}}, \gamma_{\mathbf{s}^3 \mathbf{t} \mathbf{u}, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}}, \gamma_{\mathbf{s}^2 \mathbf{t} \mathbf{u}^2, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}}.$$

However, the third relation in (7) gives $\rho(\gamma_{\mathbf{s}^4 \mathbf{t}, \mathbf{s}^3 \mathbf{t} \mathbf{u}}) = \rho(\gamma_{\mathbf{s}^3 \mathbf{t} \mathbf{u}, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}})$ since the first and last terms cancel in the projection, and similarly the fifth relation in (7) gives $\gamma_{\mathbf{s}^3 \mathbf{t}^2, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}} = \gamma_{\mathbf{s}^2 \mathbf{t} \mathbf{u}^2, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}}$, since the second and third terms cancel in the projection. Thus the $\rho(\gamma_{\mathbf{m}_1, \mathbf{m}_2})$ span a 4-dimensional space. Including the $\rho(\ell_{\mathbf{s}})$, we get $\dim(\mathbf{N}) = 5$. \square

Proposition 6. $\rho(M) = \mathbf{M}$.

Proof. Let $\tau \subset \mathbf{M}$ be the simplicial cone spanned by the $\rho(\gamma_{\mathbf{m}_1, \mathbf{m}_2})$ and $\rho(\ell_{\mathbf{s}})$. We compute the dual cone τ^\vee , expressing the edges in terms of S_5 -invariant divisors on X_ψ .

Claim. Each of these five S_5 -invariant divisors can be chosen to be effective.

We defer the calculation and complete the proof. \square

Let C be an irreducible curve in the mirror quintic. If C is contained in a toric divisor, we already know that $[C]$ is contained in the semigroup spanned by $\ell_{\mathbf{t}}$, $\sigma_{\mathbf{s}, \mathbf{t}}$, and $\gamma_{\mathbf{m}_1, \mathbf{m}_2}$ by our previous calculation of the Mori cone of these toric surfaces. So $\rho([C]) \in \mathbf{M}$.

If C is not contained in any toric divisor, then its intersection with each toric divisor is nonnegative, and in particular its intersection with each of the five S_5 -invariant divisors is nonnegative. But these intersection numbers are simply the coefficients of $\rho(C)$ as it is expressed as a linear combination of the edges of τ . Therefore $\rho(C) \in \tau \subset \mathbf{M}$.

It remains to compute the edges of τ^\vee . We can express the five generators of τ as

$$(35) \quad \rho(\ell_{\mathbf{s}}), \rho(\gamma_{\mathbf{s}^4 \mathbf{t}, \mathbf{s}^4 \mathbf{u}}), \rho(\gamma_{\mathbf{s}^4 \mathbf{t}, \mathbf{s}^3 \mathbf{t} \mathbf{u}}), \rho(\gamma_{\mathbf{s}^3 \mathbf{t}^2, \mathbf{s}^3 \mathbf{t} \mathbf{u}}), \rho(\gamma_{\mathbf{s}^3 \mathbf{t}^2, \mathbf{s}^2 \mathbf{t}^2 \mathbf{u}})$$

We have already constructed an effective toric divisor D^x with $D^x \cdots \ell_{\mathbf{s}} = 1$ and having intersection number zero with each γ curve. Therefore, symmetrizing D^x gives an S_5 -invariant divisor which spans the edge of τ^\vee dual to the face of τ which does not contain $\rho(\ell_{\mathbf{s}})$.

4. Conclusions

We have studied the curves on the quintic mirror threefold, using toric constructions to obtain curves, and using toric surfaces to see the relations among them. Those same toric surfaces allowed us to isolate some of these curves as potentially extremal: they are extremal on the surface that contains them, but not (yet) known to be extremal on the threefold itself. We hope that this approach to studying curves will be useful in other contexts as well.

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References

- [1] A. ALBANO and S. KATZ, *Lines on the Fermat quintic threefold and the infinitesimal generalized Hodge conjecture*, Trans. Amer. Math. Soc. **324** (1991) 353–368. [MR1024767](#)
- [2] P. S. ASPINWALL and D. R. MORRISON, *Topological field theory and rational curves*, Comm. Math. Phys. **151** (1993) 245–262, [arXiv:hep-th/9110048](#). [MR1204770](#)
- [3] V. V. BATYREV, *Dual polyhedra and mirror symmetry for Calabi–Yau hypersurfaces in toric varieties*, J. Algebraic Geom. **3** (1994) 493–535, [arXiv:alg-geom/9310003](#). [MR1269718](#)
- [4] P. CANDELAS, X. C. DE LA OSSA, P. S. GREEN, and L. PARKES, *A pair of Calabi–Yau manifolds as an exactly soluble superconformal theory*, Nucl. Phys. B **359** (1991) 21–74. [MR1115626](#)
- [5] C. H. CLEMENS, *Double solids*, Adv. in Math. **47** (1983) 107–230. [MR0690465](#)
- [6] C. H. CLEMENS, *Homological equivalence, modulo algebraic equivalence, is not finitely generated*, Publ. Math. IHES **58** (1983) 19–38. [MR0720930](#)
- [7] C. H. CLEMENS, *Curves on higher-dimensional complex projective manifolds*, in: Proceedings of the International Congress of Mathematicians, Vols. 1, 2 (Berkeley, Calif., 1986), Amer. Math. Soc., Providence, RI, 1987, pp. 634–640. [MR0934266](#)
- [8] A. GIVENTAL, *A mirror theorem for toric complete intersections*, in: Topological Field Theory, Primitive Forms and Related Topics (Kyoto, 1996), Progr. Math., vol. 160, Birkhäuser Boston, Boston, MA, 1998, pp. 141–175, [arXiv:alg-geom/9701016](#). [MR1653024](#)
- [9] R. GOPAKUMAR and C. VAFA, *M-theory and topological strings. II*, [arXiv:hep-th/9812127](#).
- [10] H. HAYASHI, P. JEFFERSON, H.-C. KIM, K. OHMORI, and C. VAFA, *SCFTs, Holography, and Topological Strings*, [arXiv:1905.00116](#) [hep-th].

- [11] S. KATZ, *Genus zero Gopakumar-Vafa invariants of contractible curves*, J. Differential Geom. **79** (2008) 185–195, [arXiv:math.AG/0601193](#). [MR2420017](#)
- [12] D. MAULIK and Y. TODA, *Gopakumar-Vafa invariants via vanishing cycles*, Invent. Math. **213** (2018) 1017–1097, [arXiv:1610.07303](#) [math.AG]. [MR3842061](#)
- [13] D. R. MORRISON, *Mirror symmetry and rational curves on quintic threefolds: A guide for mathematicians*, J. Amer. Math. Soc. **6** (1993) 223–247, [arXiv:alg-geom/9202004](#). [MR1179538](#)
- [14] D. R. MORRISON, *Geometric aspects of mirror symmetry*, in: Mathematics Unlimited – 2001 and Beyond (B. ENQUIST and W. SCHMID, eds.), Springer-Verlag, 2001, pp. 899–918, [arXiv:math.AG/0007090](#). [MR1852194](#)
- [15] C. VOISIN, *A mathematical proof of a formula of Aspinwall and Morrison*, Compositio Math. **104** (1996) 135–151. [MR1421397](#)

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