

# Entanglement between Muon and $I > \frac{1}{2}$ Nuclear Spins as a Probe of Charge Environment

Pietro Bonfà<sup>1,\*</sup>, Jonathan Frassinetti<sup>1,2</sup>, John M. Wilkinson<sup>1,3,†</sup>, Giacomo Prando<sup>1,4</sup>, Muhammad Maikudi Isah<sup>1</sup>, Chennan Wang<sup>5</sup>, Tiziana Spina<sup>6</sup>, Boby Joseph<sup>7</sup>, Vesna F. Mitrović<sup>8</sup>, Roberto De Renzi<sup>1</sup>, Stephen J. Blundell<sup>3</sup>, and Samuele Sanna<sup>1,2</sup>

<sup>1</sup>Department of Mathematical, Physical and Computer Sciences, University of Parma, Parco Area delle Scienze 7/A, 43124 Parma, Italy

<sup>2</sup>Department of Physics and Astronomy "A. Righi", University of Bologna and INFN Sezione di Bologna, via Berti Pichat 6/2, 40127 Bologna, Italy

<sup>3</sup>Clarendon Laboratory, Department of Physics, University of Oxford, Parks Road, Oxford OX1 3PU, United Kingdom

<sup>4</sup>Department of Physics, University of Pavia, 27100 Pavia, Italy

<sup>5</sup>Laboratory for Muon Spin Spectroscopy, Paul Scherrer Institute, CH-5232 Villigen, Switzerland

<sup>6</sup>Superconducting Radio Frequency (SRF) Materials and Research Department, Fermilab, Batavia, Illinois 60510, USA

<sup>7</sup>Elettra-Sincrotrone Trieste, S.S. 14-km 163.5, Basovizza, 34149 Trieste, Italy

<sup>8</sup>Department of Physics, Brown University, Providence, 02912 Rhode Island, USA

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We report on the first example of quantum coherence between the spins of muons and quadrupolar nuclei. We reveal that these entangled states are highly sensitive to a local charge environment and thus, can be deployed as a functional quantum sensor of that environment. The quantum coherence effect was observed in vanadium intermetallic compounds which adopt the A15 crystal structure, and whose members include all technologically pertinent superconductors. Furthermore, the extreme sensitivity of the entangled states to the local structural and electronic environments emerges through the quadrupolar interaction with the electric field gradient due to the charge distribution at the nuclear ( $I > \frac{1}{2}$ ) sites. This case study demonstrates that positive muons can be used as a quantum sensing tool to also probe structural and charge-related phenomena in materials, even in the absence of magnetic degrees of freedom.

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Quantum coherence between an implanted positively charged muon and nuclei in a solid was first conclusively demonstrated using muon-spin spectroscopy ( $\mu$ SR) experiments on simple ionic fluorides [1]. The strong hydrogen-like bonding of the implanted positive muon (chemically identified as  $\mu^+$ ) to nearest-neighbor F ions, characterized by a single spin  $1/2$   $^{19}\text{F}$  nuclear isotope, gives rise to a hierarchical separation of the interactions. Typically, dipolar couplings with two nearest-neighbor (NN)  $^{19}\text{F}$  nuclear spins,  $\mathbf{I}_1$  and  $\mathbf{I}_2$ , determine the dominant spin-Hamiltonian of the  $S = 1/2$  muon, whereas all the residual interactions, starting from the next nearest neighbors (NNN), can be ignored to a first approximation. Thanks to the 100% initial muon-spin polarization, a prerogative of  $\mu$ SR, this shows up experimentally as a characteristic coherent spin precession pattern in the muon time-dependent asymmetry, uniquely determined by the geometry of the F- $\mu$ -F bonds. Many fluorinated compounds display this coherent pattern in nonmagnetic phases, including ionic fluorides [2–5], fluopolymers [6,7], and molecular magnets [8]. For these materials, the absence or the fast fluctuations of electronic magnetic moments leave the nuclear spin interactions to determine the dynamics of the muon-spin polarization. This allows a very precise assignment of the muon implantation

site, now known to be particularly accurate with the help of density functional theory (DFT) *ab initio* simulations of the muon stopping site inside the crystal (a technique which is also known as DFT +  $\mu$  [9–13]). A similar coherent spin behavior has been identified in certain hydrides [14–16] and in metal-organic frameworks [17], where for instance a close association of a proton and the positive muon approximates a muoniated hydrogen molecule,  $\mu\text{H}$ , or possibly, a bonded molecular ion,  $(\mu\text{H})^+$ ,  $(\mu\text{H})^-$ . Notice that  $^1\text{H}$ , like  $^{19}\text{F}$ , is a spin  $I = 1/2$  nucleus hence with zero electric quadrupole moment.

In the case of  $^1\text{H}$ , as for the cases of many other nuclear species, such a coherent pattern is rarely observed in  $\mu$ SR experiments. Much more often a large number of unpolarized nuclear spins give rise to a  $T_2^{-1}$  relaxation process with either Gaussian or Lorentzian line shapes, both the hallmarks of fast decoherence on the timescale of the period of the coherent quantum interference processes. Fluorine is special since it is very electronegative, and it has both a small ionic radius and a large nuclear moment, so that its dipolar coupling to the muon is strong, and consequently several oscillations in any quantum-coherent signal can be observed before all muons have decayed or any nuclear relaxation process has become significant. The special

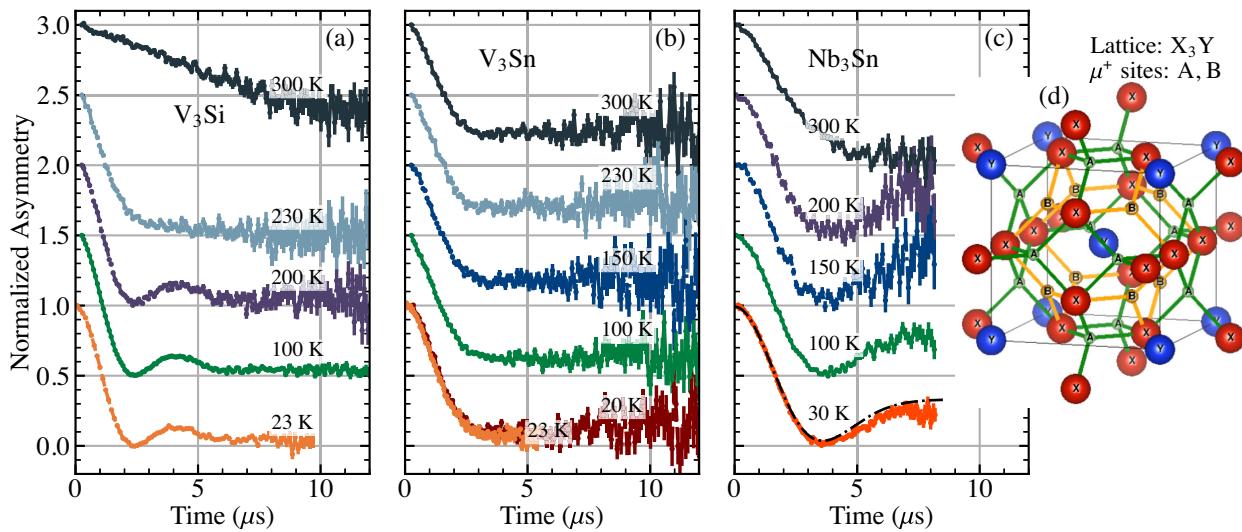


FIG. 1. Experimental results obtained for V<sub>3</sub>Si (a), V<sub>3</sub>Sn (b), and Nb<sub>3</sub>Sn (c) at various temperatures in ZF. The initial asymmetry has been normalized to 1, and the various measurements are shifted along the y axis by multiples of 0.5. The black line in (c) is a best fit to a static Kubo-Toyabe function. In picture (d) the lattice structure of A15 compounds and the candidate muon sites identified in this class of materials are depicted.

F- $\mu$ -F case was very recently revisited by some of us [18], showing the role of the rest of the nuclear spins (NNN and beyond) in the slow decoherence process of F- $\mu$ -F. This Letter implies that the very well-known F- $\mu$ -F effect, confined until now among the technicalities of the muon spectroscopy, displays all the features of a very high accuracy *quantum sensor* that can be exploited for microscopic detection of important physical phenomena [19] and can be finely controlled by electromagnetic excitation [5]. Unfortunately, until now, the sensor has been available only for F<sup>-</sup>—and, much more seldom, for H<sup>-</sup>—containing materials, insensitive to quadrupolar interaction.

In the present Letter we demonstrate the same surprising type of quantum coherence due to the entanglement of the muon spin with NN quadrupolar nuclear spin ( $I > 1/2$ ), and we establish that this quantum coherence can provide a muon spectroscopy based quantum sensor of local charge-related phenomena. We show this phenomenon in three intermetallic compounds, Nb<sub>3</sub>Sn, V<sub>3</sub>Si, and V<sub>3</sub>Sn, which belong to the A15 cubic phases (Pm $\bar{3}$ n, group number 223), whose members include several technologically dominant conventional superconductors [20]. In stark contrast to the well-studied  $I = 1/2$  case of <sup>19</sup>F and <sup>1</sup>H, the presence of NN nuclei with  $I > 1/2$ , namely  $I = 7/2, 9/2$  of <sup>51</sup>V and <sup>93</sup>Nb respectively, implies the existence of quadrupolar interactions. This has two effects that could potentially spoil the *quantum sensor* concept: first, it was until now unclear that a detectable quantum coherence could nevertheless show up in the muon asymmetry; second, quadrupolar interactions are proportional to the electric field gradient (EFG) at the nucleus in question, not just on the pure geometry of the bonds. EFG tensors are very accurately determined by DFT in bulk materials [21] and

compared with the values measured for instance by NMR [22]. The muon embedding in the crystal alters the bulk EFG in more than one way. We show that the coherent effect survives, and we develop here an accurate model to describe this phenomenon. Our modeling of the coherence entails identifying precisely the muon site and calculating muon perturbed EFG tensors at NN and NNN nuclei. The results show that the observed phenomenon is highly sensitive to small structural and electronic differences among the same A15 family, paving the way to extend the use of muon spectroscopy as a quantum sensing technique for charge-related phenomena well beyond the proof-of-concept stage presented in this Letter.

Zero-field (ZF)  $\mu$ SR temperature scans, using the EMU spectrometer at the ISIS Muon Source and the GPS spectrometer [23] at the Paul Scherrer Institute, have been conducted as a function of temperature. Further details on the experimental methods are provided in the Supplemental Material [24]. Figure 1 shows the  $\mu$ SR spectra (time-dependent spin polarization of the muon ensemble) for all the samples at various representative temperatures. The temperature dependence is relatively weak, except above 200 K, where thermally activated  $\mu^+$  diffusion occurs in V<sub>3</sub>Si [39]. A similar trend is envisaged for Nb<sub>3</sub>Sn, but with slightly higher activation energy, while a small increase in the tail is observed in V<sub>3</sub>Sn at and above 150 K. At low temperature, where the muon is static in the  $\mu$ SR time window, the results are remarkably sample dependent despite all the  $X_3Y$  samples (X = {V, Nb} and Y = {Si, Sn}) being very similar metals, sharing the same A15 cubic lattice structure. The structure is shown in Fig. 1(d), and our samples have a cubic lattice parameter  $a = 4.72 \text{ \AA}, 4.98 \text{ \AA}, 5.29 \text{ \AA}$  for V<sub>3</sub>Si, V<sub>3</sub>Sn, and Nb<sub>3</sub>Sn respectively (see

Ref. [24]), in agreement with previous results [25,40,41]. The nuclei of the  $X$  atoms are closer to the calculated muon sites, as shown in Fig. 1(d) with labels A and B, and all have similar properties:  $^{51}\text{V}$  with 99.8% abundance has spin  $I = 7/2$ , gyromagnetic ratio  $\gamma_V = 70.45 \times 10^6 \text{ rad}/(\text{sT})$  and quadrupole moment  $Q = -0.052(10) \text{ barn}$ , and  $^{93}\text{Nb}$  with 100% abundance has spin  $I = 9/2$ ,  $\gamma_{\text{Nb}} = 65.64 \times 10^6 \text{ rad}/(\text{sT})$  and  $Q = -0.32(2) \text{ barn}$  [42].

The oscillatory behavior observed in  $\text{V}_3\text{Si}$  [Fig. 1(a)] is in marked contrast to the other two samples.  $\text{Nb}_3\text{Sn}$  [Fig. 1(c)] resembles the conventional Kubo-Toyabe (KT) relaxation function and empirical KT best fit are shown by the dashed line in the same panel, characterized by a dip and a tail that flattens at  $1/3$  of the initial value.  $\text{V}_3\text{Sn}$  [Fig. 1(b)] could be described by a KT relaxation, with an additional decay of the  $1/3$  tail which has no evident physical origin. The surprisingly slow oscillations observed in  $\text{V}_3\text{Si}$  [Fig. 1(a)] cannot be due to internal fields of electronic origin since all these A15 samples are nonmagnetic. Instead, as we will show, they result from a quantum coherent precession pattern due to the coupling between the muon and nearby  $^{51}\text{V}$  nuclear moments, analogous to the  $\text{F}-\mu-\text{F}$  case, and never reported before for systems containing  $I > 1/2$  nuclear spins.

In order to explain the three precession patterns of Fig. 1 we now consider the microscopic nuclear and electronic degrees of freedom entering the quantum mechanical model of the muon polarization. The model requires the knowledge of three ingredients to reproduce the experimental muon polarization: (i) the muon site, (ii) the perturbation induced by the  $\mu^+$  on the position of the neighboring atoms, and (iii) the perturbation induced by the muon on the EFG at the nuclear sites with spin  $I > 1/2$ . This information allows us to fully define the spin Hamiltonian  $\mathcal{H}$  given by

$$\mathcal{H} = \sum_i^{N_{\text{nuc}}} \frac{\mu_0 \gamma_\mu \gamma_i \hbar^2}{4\pi r_i^3} \mathbf{S}_\mu \cdot \mathbf{D}^i \cdot \mathbf{I}^i + \frac{eQ_i}{2I(2I-1)} \mathbf{I}^i \cdot \mathbf{V}^i \cdot \mathbf{I}^i, \quad (1)$$

where  $\mathbf{S}_\mu$  is the spin of the muon,  $r_i$  is its distance from nucleus  $i$ ,  $\mathbf{I}^i$ , and  $Q_i$  are respectively the spin and the quadrupole moment of nucleus  $i$ ,  $\mathbf{D}$ , and  $\mathbf{V}$  are the dipolar and EFG tensors at nuclear site  $i$ , and other symbols have their standard meaning. All the quantities entering Eq. (1) can be accurately estimated with DFT-based *ab initio* approaches, and we describe below the results that we obtained following the DFT +  $\mu$  procedure.

Two candidate muon sites are present in our A15 compounds and are shown in Fig. 1(d) with labels A and B. Site A corresponds to site T2 in Ref. [39] and is located in the center of the tetrahedron formed by four  $X$  atoms while site B is in the center of the triangle formed by three  $X$  atoms. We find that site B always has higher energy than site A by hundreds of meV (see Ref. [24] for

details) and is therefore omitted from the subsequent analysis. DFT simulations produce, as an additional outcome, the displacements of the atoms surrounding the muon. In all cases, the NN  $X$  atoms are substantially displaced by the muon, and the nearest-neighbor distances increase by 6%, 5%, and 4% respectively in  $\text{V}_3\text{Si}$ ,  $\text{V}_3\text{Sn}$ , and  $\text{Nb}_3\text{Sn}$  (the absolute values are shown in the insets of Fig. 2 against the unperturbed  $\mu-X$  distance and in the Supplemental Material [24]).

The next step is the evaluation of the EFG at the quadrupolar nuclei in each compound. While for ionic materials a point charge approximation may sometimes be sufficient, covalent and metallic systems require more elaborate strategies. Full potential (FP) DFT simulations yield very accurate estimates in materials where the mean field approximation does not break down owing to strong correlation, but are extremely computationally demanding. For this reason, and aiming at providing an easily adoptable approach, we opted for an effective compromise between accuracy and speed using a plane wave basis [43–45] combined with projector augmented wave (PAW) [46] pseudopotentials. A detailed discussion of our strategy and additional comparisons with FP simulations [47] are provided in the Supplemental Material [24]. Notably, this procedure converges much faster than the equivalent technique aimed at the prediction of magnetic contact hyperfine fields at the muon sites [48].

Unsurprisingly, the EFG of the four  $X$  neighbors of the muon is drastically affected by the presence of the interstitial charge. For example, in  $\text{V}_3\text{Si}$  the unperturbed EFG tensor at V nuclei in the pristine material, with  $V_{zz} = 2.2 \times 10^{21} \text{ V/m}^2$  and  $\eta = 0$ , in agreement with the experimental value of  $V_{zz} = 2.37 \times 10^{21} \text{ V/m}^2$ , reduces by almost 1 order of magnitude as a consequence of the presence of the positive impurity and the lattice distortion, in agreement with earlier work [39]. Note that site assignments come with some small uncertainty, and previous investigations that can be compared with experiment [10,18,19,49,50] reveal that a discrepancy of the order of a tenth of Angstrom is to be expected. On the other hand, plane wave based estimations of EFGs are subject to a much larger uncertainty of the order of 30% and  $1.17 \times 10^{21} \text{ V/m}^2$  in relative and absolute terms [51].

Having collected all parameters entering Eq. (1), we proceed to compute the time-dependent muon polarization numerically. For the A15 compounds the internuclear dipolar interactions can be safely neglected [52] thus allowing the adoption of the approach proposed by Celio [26,27] and implemented in the publicly available code UNDI [53], which makes the estimate very quick. Our calculations consider only the effect of the nearest nuclei, but it has recently been shown by some of us [18] how to effectively include the effect of farther nuclei with an appropriate rescaling of a second nearest-neighbor interaction, allowing a substantial reduction of the otherwise

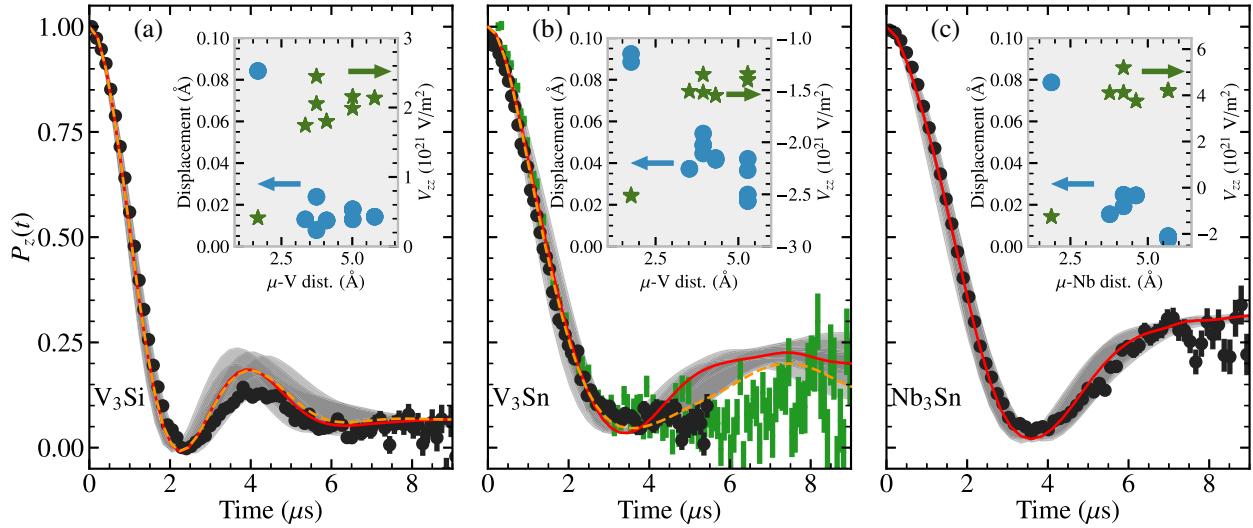


FIG. 2. Comparison between experimental and predicted muon spin polarization obtained using atomic displacements and EFGs from plane wave based DFT calculations. The black dots in panels (a), (b), and (c) are the lowest temperature data collected at PSI for  $V_3Si$ ,  $V_3Sn$ , and  $Nb_3Sn$  respectively (orange and red points of Fig. 1). The green bars in (b) are ISIS results collected at 20 K. A background has been estimated by comparing the asymmetries collected at ISIS and PSI and removed. The red (orange) line in all plots is the depolarization obtained using first principles results from PW (FP) simulations to solve Eq. (1). Shaded area highlights different trends that originate by taking into account typical uncertainties of the DFT-based predictions (see main text). The insets show the perturbation induced by the muon on its  $X$ -type neighbors ( $X = V, Nb$ ). In particular, in the presence of the muon, the displacement of each  $X$  atom from its equilibrium position in the unperturbed lattice, and the values of  $V_{zz}$  at the considered atomic site are reported on the left-hand and right-hand  $y$  axes, respectively, as a function of the unperturbed distance of the considered atom from the  $\mu^+$  interstitial position in a  $3 \times 3 \times 3$  supercell.

exponentially diverging dimension of the Hilbert space. Following Ref. [18], we consider 4 NN and 4 NNN whose positions are homogeneously rescaled by a small amount to compensate for the remaining nuclei (including also the low abundant Sn spins; see the Supplemental Material [24] for details).

The predicted  $\mu$ SR signal obtained fully *ab initio*, i.e., without free parameters, is shown for all samples in Fig. 2 by a red line [plane wave (PW) results] and an orange dashed line (FP results), while the shaded area indicates the uncertainty in the PW based prediction quantified with a reduction or increase of 3% (29%) of  $d_{\mu-X}$  (EFG values). Experimental data acquired at the lowest temperature for each compound are shown for comparison. Perfect agreement is found for  $Nb_3Sn$  [Fig. 2(c)], while for  $V_3Si$  [Fig. 2(a)] a small deviation is observed at about 4  $\mu$ s where the first bump is slightly overestimated, although the experimental result falls inside the shaded area. A small increase of 15 mÅ in the  $\mu$ -V distance allows us to recover perfect agreement (see the Supplemental Material [24]). Remarkably the oscillation (the time position of minima and maxima) is very well reproduced.  $V_3Sn$  is the sample showing worst agreement in the long-time tail. In this case the deviation can be partially attributed to the limits of the PAW approximation in reconstructing the EFG at the V sites. Indeed the FP prediction, which differs from the PW based estimate by 16%, improves the agreement with

the experimental data [54]. These trends demonstrate the exquisite sensitivity of  $\mu$ SR to atomic distances and EFGs.

The striking difference between the muon asymmetries collected in a set of compounds that share the same lattice structure, the same muon site, and similar lattice distortions may appear puzzling at first sight. To address this point, we introduce the simple and analytically solvable case of one muon interacting with a single nucleus of spin  $I$  subject to an axial EFG [55]. In zero field (ZF), the interaction depends on two parameters [56]:

$$\omega_D = \frac{\mu_0 \gamma_\mu \gamma_I \hbar}{4\pi r_I^3}, \quad \omega_Q = \frac{eV_{zz}Q}{4I(2I-1)\hbar}.$$

Figure 3(a) shows the muon polarization as a function of time for various values of  $\omega_Q/\omega_D$  for a single nuclear spin  $I = 7/2$ . This simple model illustrates how, in the two extreme regimes of zero and large quadrupolar splitting, the classical expectation of a single precession frequency is recovered, while, in intermediate regimes, multiple frequencies appear. Similarly, a departure from the semi-classical KT behavior can also be appreciated in the more relevant case of a muon generating an EFG on four tetrahedrally coordinated  $I = 7/2$  nuclei. The polarization as a function of time is obtained numerically in this case and shown in Fig. 3(b). The trend recovers the 1/3 tail of the classical KT limit only in the small and large

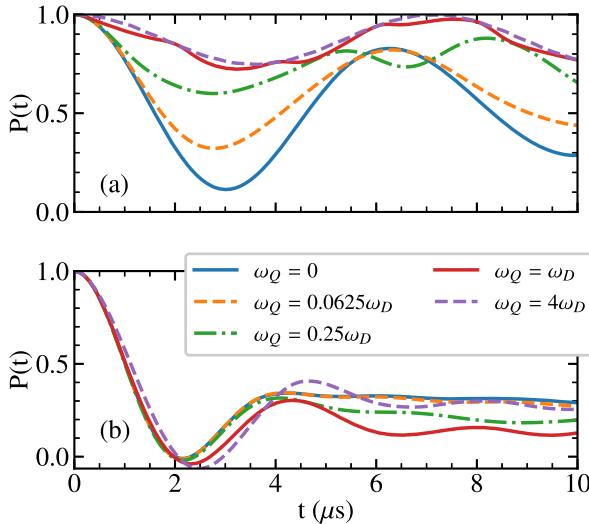


FIG. 3. (a) Time-dependent spin polarization  $P(t)$  for a muon interacting with a single nucleus with spin  $I = 7/2$  subject to an axial EFG for various values of  $\omega_Q/\omega_D$ . (b)  $P(t)$  for a muon tetrahedrally coordinated to four  $I = 7/2$  nuclei, all subject to the EFG generated by the muon itself.

quadrupolar splitting conditions, while substantial deviations happen for the intermediate regime.

While the details of the muon polarizations in  $A15$  compounds are connected to the peculiar composition of intrinsic and muon induced EFGs at the  $X$  sites, yielding to the simulated curves of Fig. 1, the behavior can be qualitatively understood considering the ratio  $|\omega_Q|/\omega_D$  for the NN. Indeed this ratio happens to be about 2.4 for  $Nb_3Sn$ , 2 for  $V_3Sn$ , and 0.2 for  $V_3Si$ , thus qualitatively explaining the deviations from a KT-like trend of the latter two samples.

In conclusion, we have presented the experimental observation of coherent oscillations originating from the interaction between the muon and nuclei with quadrupolar moments. This signal is analogous to what has already been observed in fluorides and other materials containing  $I = 1/2$  nuclei, with higher magnetic moments but zero electric quadrupole moment. An accurate description of the  $\mu$ SR spectra was obtained by solving parameter free spin Hamiltonians that consider the perturbed EFG at nuclear sites surrounding the muon and effectively include all nuclear spins in the system to correctly describe long-time depolarization. In  $\mu$ SR experiments the time evolution of the muon-spin polarization depends dramatically upon the electronic distribution at quadrupolar nuclei coupled to the muon, and an accurate estimation of the perturbed EFG at these sites is crucial for a successful analysis. At the same time, our Letter reveals a high accuracy quantum sensor of local charge environment.

Finally, we have shown that DFT-based simulations can be effectively used to model the charge distribution and how their combination with simple spin Hamiltonians

represents a computationally inexpensive method to accurately predict the  $\mu$ SR spectra of nuclear origin in virtually any crystalline specimen. The strong dependence of the  $\mu$ SR signal on the EFGs and the possibility of estimating quantitatively the perturbation of an interstitial  $\mu^+$  opens the possibility of using positive muons as a quantum sensing tool to probe not only magnetic phenomena in materials but charge-related ones as well.

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\*[pietro.bonfa@unipr.it](mailto:pietro.bonfa@unipr.it)

<sup>†</sup>Present address: ISIS Pulsed Neutron and Muon Facility, Rutherford Appleton Laboratory, Chilton, Oxfordshire OX11 0QX, United Kingdom.

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