

Close encounters of tight binary stars with stellar-mass black holes

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Accepted 2023 January 4. Received 2023 January 3; in original form 2022 November 4

ABSTRACT

Strong dynamical interactions among stars and compact objects are expected in a variety of astrophysical settings, such as star clusters and the disks of active galactic nuclei. Via a suite of three-dimensional hydrodynamics simulations using the moving-mesh code AREPO, we investigate the formation of transient phenomena and their properties in close encounters between an $2 M_{\odot}$ or $20 M_{\odot}$ equal-mass circular binary star and single $20 M_{\odot}$ black hole (BH). Stars can be disrupted by the BH during dynamical interactions, naturally producing electromagnetic transient phenomena. Encounters with impact parameters smaller than the semimajor axis of the initial binary frequently lead to a variety of transients whose electromagnetic signatures are qualitatively different from those of ordinary disruption events involving just two bodies. These include the simultaneous or successive disruptions of both stars and one full disruption of one star accompanied by successive partial disruptions of the other star. On the contrary, when the impact parameter is larger than the semimajor axis of the initial binary, the binary is either simply tidally perturbed or dissociated into bound and unbound single stars ('micro-Hills' mechanism). The dissociation of $20 M_{\odot}$ binaries can produce a runaway star and an active BH moving away from one another. Also, the binary dissociation can either produce an interacting binary with the BH, or a non-interacting, hard binary; both could be candidates of BH high- and low-mass X-ray binaries. Hence, our simulations especially confirm that strong encounters can lead to the formation of the (generally difficult to form) BH low-mass X-ray binaries.

Key words: black hole physics – gravitation – hydrodynamics – stellar dynamics.

1 INTRODUCTION

Dynamical interactions between stars and compact objects in dense environments, such as star clusters, play a very important role in a variety of astrophysical settings. Many dynamical interactions redistribute energy between star cluster members which drives them towards equilibrium. Furthermore, binaries, an energy source in clusters, can be formed or destructed via dynamical interactions (Hut et al. 1992), which determines the cluster's thermodynamic state.

Occasionally, dynamical encounters between astrophysical objects lead them to interact at a close distance. In particular, close encounters involving stars and stellar-mass black holes (BHs) can often create transient phenomena, such as tidal disruption events (TDEs) (Hills 1988; Perets et al. 2016; Lopez et al. 2019a; Kremer et al. 2019; Wang, Perna & Armitage 2021; Kremer et al. 2021, 2022). In those the BH can fully or partially disrupt a star when the closest point of approach is smaller than the so-called tidal disruption radius $r_t \sim (M_{\bullet}/M_{\star})^{1/3} R_{\star}$. Here, M_{\bullet} is the BH mass, M_{\star} the stellar mass, and R_{\star} the stellar radius. In this process, a bright electromagnetic flare can be generated. In particular, encounters involving multiplets (e.g.

binary), which frequently occur near the centre of clusters, can create wider varieties of transients than those between two single objects because of the chaotic nature of interactions (Lopez et al. 2019b; Ryu, Perna & Wang 2022).

Another X-ray source that can form during dynamical interactions between stars and BHs is an X-ray binary (e.g. Kremer et al. 2018b, a), which can show a transient behaviour via, e.g. a disk instability (e.g. King, Kolb & Burderi 1996). A number of BH X-ray binaries has been detected in our Galaxy (Corral-Santana et al. 2016) via radio (e.g. Chomiuk et al. 2013) or X-ray measurements (e.g. Miller-Jones et al. 2015; Shishkovsky et al. 2018). Among these, the detection of two dozens of *BH low-mass X-ray binaries* (BH-LMXBs) (Casares & Jonker 2014) has been puzzling astronomers, and the formation channel of such systems has been debated. In the 'standard' scenario, the orbit of a binary star system shrinks via a common envelope phase before one star collapses to a compact object (e.g. de Kool, van den Heuvel & Pylyser 1987). However, this scenario has been challenged by Podsiadlowski, Rappaport & Han (2003), who suggested that binaries with a very large mass ratio would eventually merge during the common envelope phase. Also dynamically-formed binaries via three-body encounters tend to have comparable-mass companions as a result of a member exchange. Therefore, alternative scenarios have been examined, such as dynamical formation via a number of weak

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encounters (Michaely & Perets 2016) and formation in hierarchical triples (Naoz et al. 2016). In addition to these formation channels, a strong three-body encounter between a star and BH can form such systems. The last possibility adds to the importance of investigating the hydrodynamics of three-body encounters between stars and BHs.

Theoretical investigations of transient phenomena created in multi-body encounters are especially timely in light of the dramatically increase of the number of detectable transients with both ongoing surveys, such as eROSITA¹ and the Zwicky Transient Facility (ZTF),² but especially with the upcoming Vera Rubin Observatory (VRO).³ Despite the impending increase of the number of transient candidates, there have been only a few attempts to perform hydrodynamics calculations for dynamical encounters involving binaries (e.g. McMillan et al. 1991; Goodman & Hernquist 1991; Lopez et al. 2019b). In Ryu et al. (2022) (Paper 1 hereafter), we explored the outcomes of close three-body encounters between a single MS star and a merging (i.e. with lifetime shorter than the Hubble time) binary BH using three-dimensional (3D) smoothed particle hydrodynamics simulations. In particular, we focused on quantifying the impact of the close encounter, frequently resulting in the disruption of the star, on the binary orbit and compared that with what the point-particle approximation predicts. We showed that a single disruption event can change the binary orbital parameters up to 20 per cent of their pre-encounter values, or equivalently change the gravitational-wave-driven merger time-scale up to order unity. This impact is often different from that of a pure scattering. We further showed that the accretion rates of both BHs that have undergone a disruption event are typically super-Eddington with modulations on a time-scale of the binary orbital period.

We continue our investigation on three-body encounters between stars and stellar-mass BHs using hydrodynamics simulations. In particular, in this paper we focus on possible transient formation in close three-body encounters between a binary star and a single BH, using the moving-mesh hydrodynamics code AREPO (Springel 2010). This type of encounters can naturally happen in dense stellar environments where binaries exist, such as open clusters, which have a binary fraction $\simeq 0.2\text{--}0.7$ (Moe & Di Stefano 2017; Sollima et al. 2010), nuclear star clusters (Fragione, Perna & Loeb 2021), globular clusters (Perets et al. 2016; Kremer et al. 2019), young star clusters (Kremer et al. 2021), the disks of Active Galactic Nuclei (Yang et al. 2022). In addition to varieties of transient phenomena qualitatively different from ordinary TDEs, our numerical investigation shows that the outcome of close encounters between BHs and binary stars can lead to a variety of astrophysical outcomes, including the formation of both BH high- and low-mass X-ray binary and runaway stars with active single BHs, among others.

Our paper is organized as follows: our numerical methods, based on the use of the moving mesh code AREPO are detailed in Section 2. Our results are presented in Section 3. In Section 4, we discuss their astrophysical implications (Section 4.1), and estimate the rate (Section 4.2). We summarize and conclude in Section 5.

2 SIMULATION DETAILS

In this work, we investigate the outcome properties of nearly parabolic encounters between a binary star and a single stellar-mass BH. In particular, we focus on the properties of the surviving

remnants and on the accretion and spin state of the BH following the encounter.

2.1 Numerical methods

To achieve our scientific goal, we perform a suite of 3D hydrodynamic simulations of the close encounters using the moving mesh code AREPO (Springel 2010; Pakmor et al. 2016; Weinberger, Springel & Pakmor 2020). AREPO is a massively parallel gravity and magnetohydrodynamic code, which has been used for many astrophysical problems (e.g. Vogelsberger et al. 2014). It adopts a second-order finite-volume scheme to discretize the hydrodynamic equations on a moving Voronoi mesh, and a tree-particle-mesh method for gravitational interactions. This approach of constructing grids represents a compromise between the two widely used hydrodynamics schemes, that is the Eulerian finite-volume method and the Lagrangian smoothed particle method. As a result, this new approach inherits advantages of both schemes, such as improved shock capturing without introducing an artificial viscosity, and adaptive adjustment of spatial resolution. The gas self-gravity is computed in AREPO using a tree solver and is coupled to the hydrodynamics via a Leapfrog time integration scheme.

We use the HELMHOLTZ equation of state (Timmes & Swesty 2000) which includes the radiation pressure, assuming local thermodynamic equilibrium. We include eight isotopes (n, p, ⁴He, ¹²C, ¹⁴N, ¹⁶O, ²⁰Ne, ²⁴Mg) with a nuclear reaction network (Pakmor et al. 2012).

2.2 Binary stars

The initial state of the stars was taken from somewhat evolved main-sequence (MS) stars (with the core H mass fraction of 0.5) computed using the stellar evolution code MESA (version r22.05.1) (Paxton et al. 2011, 2013, 2015, 2018, 2019; Jermyn et al. 2022). We map the one-dimensional MESA model into a 3D AREPO grid with $N \simeq 3 \times 10^5$ cells, using the profiles of density, pressure and chemical composition (Ohlmann et al. 2017). The single star is first relaxed for five stellar dynamical times $t_{\text{dyn}} = \sqrt{R_{\star}^3/GM_{\star}}$ where R_{\star} and M_{\star} are the radius and mass of the star, respectively. The density profiles of the relaxed stars considered in our simulations are depicted in Fig. 1.

We then relax binaries consisting of two relaxed single stars for $0.5P$ where $P = 2\pi\sqrt{a^3/G(M_1 + M_2)}$ is the period of the binary with the primary mass M_1 , secondary mass M_2 and semimajor axis a . We parameterize the semimajor axis a using an approximate analytic estimate of the Roche lobe radius (Eggleton 1983),

$$\frac{r_{\text{RL}}}{a} = \frac{0.49q^{2/3}}{0.6q^{2/3} + \ln(1 + q^{1/3})}, \quad (1)$$

where r_{RL} is the Roche lobe radius and a is the orbital orbital separation. For $q = 1$ and $r_{\text{RL}} = R_{\star}$, we define $a_{\text{RL}} \equiv a(r_{\text{RL}} = R_{\star})$ as the separation at which both stars fill their Roche lobes. It follows that $a_{\text{RL}} \simeq 2.64R_{\star}$ and $P/t_{\text{dyn}} \simeq 19(a/a_{\text{RL}})^{3/2}$.

We performed this binary relaxation process for every binary with different orbital parameters (six different binaries in total). The semimajor axis and the eccentricity of the relaxed binaries differ by less than 1 per cent of their initial values.

Note that we performed convergence tests with different number of cells ($N = 2.5 \times 10^5, 5 \times 10^5, 10^6$ and 2×10^6) for a few cases in which we compared the outcomes and a few key quantities such as the semimajor axis and eccentricity of the final binary and BH ejection speed. We found that all the simulations showed converged results.

¹<https://erosita.mpe.mpg.de>

²<https://www.ztf.caltech.edu>

³<https://www.lsst.org>

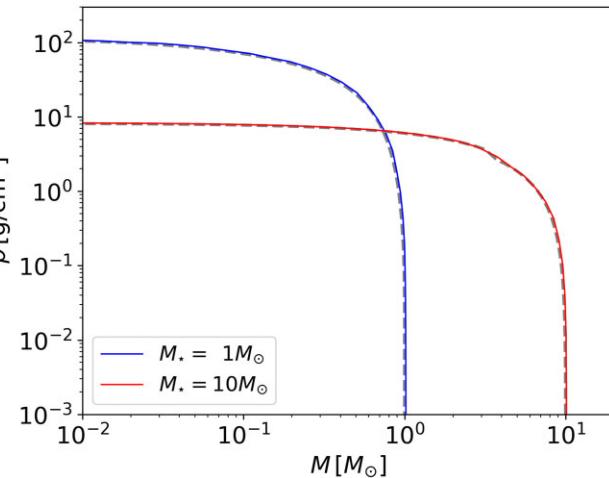


Figure 1. The radial density profile of the MS stars with $M_\star = 1 M_\odot$ (blue) and $10 M_\odot$ (red) relaxed for five stellar dynamical times, as a function of mass. The dashed lines indicate the profiles for the MESA models, which are just sitting below the solid lines.

2.3 Black holes

We model the BH using an initially non-rotating sink particle which interacts via gravity with gas and grows in mass via accretion of gas. The gravitational softening length of the BH is set to be the same as the minimum softening length of the cells of the stars. At every time step, we assume that the accretion rate is determined by the average inward radial mass flux towards the BH. More precisely, accretion occurs in four steps:

(i) *Cell identification for mass flux calculation*: the code identifies cells around the BH within $1.5 \times 10^4 r_g$ where $r_g = GM_\bullet/c^2$ is the gravitational radius of the BH. We choose a somewhat large search radius to ensure the sampling of a large enough number of cells. However, as explained in ii), we compensate for the large search radius by adopting an inverse-distance weight function.

(ii) *Accretion calculation*: to account for the fact that gas closer to the BH is expected to contribute more to the accretion rate, we estimate an average radial mass flux of the identified cells for given position and velocity of the BH as,

$$F^r = \sum_i m_i W(r_i, h_i) v_i^r, \quad (2)$$

where m_i is the mass of the selected neighbouring cells, $W(r_i, h_i)$ the (inverse-distance weighted) spline Kernel (Monaghan & Lattanzio 1985) (equation 4 in Springel 2005), r_i the distance from the BH, h_i the softening length of the BH, v_i^r the velocity of the gas relative to the BH in the radial direction. The effective surface area A is determined using the weighted average volume V of the cells with volume V_i , $V = \sum_i V_i W(r_i, h_i) / \sum_i W(r_i, h_i)$, giving $A = 3^{2/3} (4\pi)^{1/3} V^{2/3}$ (assuming spherical geometry). Finally, the accretion rate is $\dot{M} = \text{Max}(-F^r A, 0)$. At every time step with step size Δt where the BH's position and velocity are updated, the BH grows in mass by an amount of $\Delta M = (1 - \eta) \dot{M} \Delta t$ where η is the radiative efficiency. We assume $\eta = 0$ in our simulations.

(iii) *Mass subtraction and back reaction*: to conserve the mass, the same amount of mass is subtracted from a cell closest to the BH among those with a mass of at least $> 5 \Delta M$ and density $> 10^{-5} \text{ g cm}^{-3}$ and bound to the BH. Although the minimum mass is chosen mostly to ensure to avoid a very small mass of the cell after the mass subtraction, this particular choice should not affect our results

because the masses of the subtracted cells are typically much above the minimum mass. The lower bound of the density is to ensure that the mass is subtracted from 'real' gas, not from a vacuum cell. Then, in order to properly take into account the back-reaction of accretion, we subtract the average momentum of the accreted mass,

$$\mathbf{P} = \Delta M \frac{\sum_i W(r_i, h_i) \mathbf{v}_i}{\sum_i W(r_i, h_i)}, \quad (3)$$

where \mathbf{v}_i is the velocity vector of the i^{th} neighbouring cell identified at the step i , from the 'mass-losing' cell, and add the same amount to the BH.

(iv) *Spin evolution*: accretion of gas can increase the BH spin. We evolve the BH spin due to mass accretion following Paper 1, who adopted the formalism by Fanidakis et al. (2011). The only significant difference is how the direction of the accreted angular momentum is calculated: Paper 1, which performed smoothed particle hydrodynamics simulations, tracked the angular momentum vector simply by cumulatively adding the momentum of the accreted particles, but in this work, we make use of the average angular momentum (measured in the BH frame) of the selected neighbouring cells.

We should note that our simulations do not include the radiation feedback produced by accretion. This may contribute to create significant outflows and regulate the subsequent accretion, especially at super-Eddington accretion rates (e.g. Sadowski et al. 2014). While our simulations reveal a radiation pressure gradient built up in the optically thick gas near the BH, which can push gas away from the BH, its impact is found to be small. Given the purpose of this work, namely, to identify all possible transient types during star-binary/BH three-body encounters, and properly classify the outcomes, we defer improvements of the treatment of accretion and its feedback on surrounding gas to future work specifically dedicated to studying the outcome observables.

2.4 Initial parameters

Throughout the paper, quantities with the subscript containing $b(\star) - \bullet$ refer to those related to the orbit between a binary (single star) and the BH. We consider a parabolic encounter with eccentricity $1 - e_{b-\bullet} = 10^{-5}$. Parabolic orbits are a reasonable assumption for encounters in star clusters given that the typical eccentricity of two-body encounters in a cluster with velocity dispersion σ is $|1 - e| \simeq 10^{-4} (\sigma/15 \text{ km s}^{-1})^2 (M_\bullet/40 M_\odot)^{-2/3} (M_\star/1 M_\odot)^{-1/3} (R_\star/1 R_\odot)$ for $M_\bullet \gg M_\star$.⁴ The distance between the binary's centre of mass and the BH at the first closest approach $r_{p,b-\bullet}$ is parameterized using the impact parameter β , i.e. $r_{p,b-\bullet} = 0.5 \beta r_t$. Here, r_t is defined as $[M_\bullet/(M_1 + M_2)]^{1/3} a(1 + e)$ where a and e are the binary semimajor axis and eccentricity, respectively.

For the mass of the BH we choose $20 M_\odot$ (c.f., Spera & Mapelli 2017). To study the impact of binary mass on outcomes, we consider a low-mass case, where we choose $M_b = M_1 + M_2 = 2 M_\odot$, ($M_{1,2} = 1 M_\odot$) and a high mass case where we chose $M_b = M_1 + M_2 = 20 M_\odot$, ($M_{1,2} = 10 M_\odot$). We consider three semi-major axes: $a/a_{\text{RL}} = 2, 6$, and 9 , and assume the orbit to be circular initially. Massive stars are commonly found in such close binary systems (Sana et al. 2012). For lower-mass stars such tight systems are rare (Duchêne &

⁴The situation is somewhat different in AGN disks, where during the early times, when prograde orbits are in the process of being damped, there is a much higher likelihood of encounters being hyperbolic (Secunda et al. 2021). However, as time goes on, gentler encounters are expected.

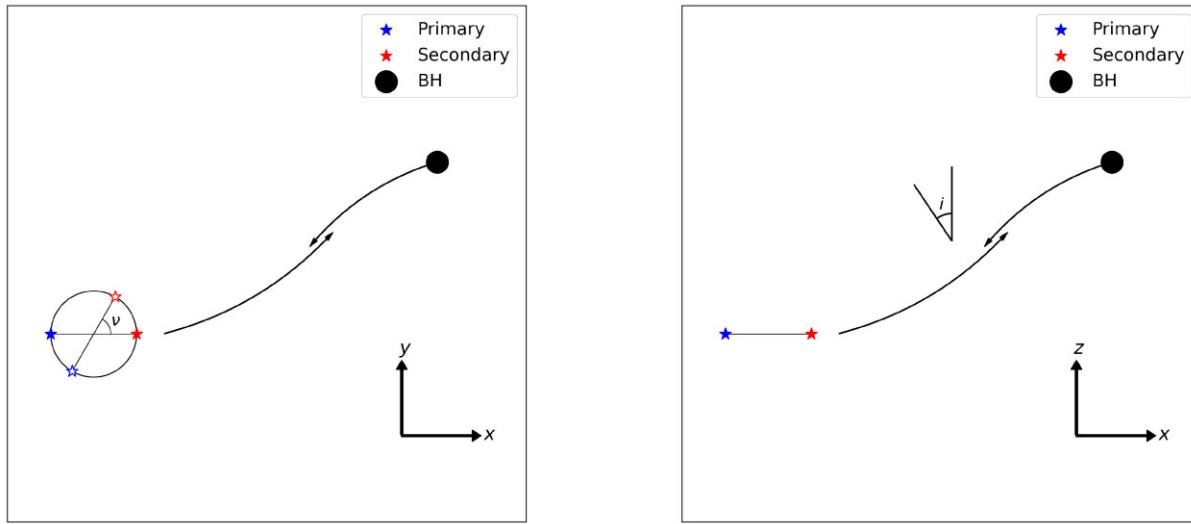


Figure 2. Schematic diagrams for the initial configuration of the stellar binary (blue and red solid stars) and BH (black circle) for a prograde case with the inclination angle $i < 90^\circ$ and the phase angle $v = 0^\circ$, projected onto the $x - y$ plane (left) and the $x - z$ plane. The arrows indicate the instantaneous direction of motion. The open stars, on the same circle with the closed stars, indicate the case with $v > 0^\circ$.

Kraus 2013), but their longer life lifetimes and the fact that they are favoured by the initial mass function may still make possible encounters of these low-mass binaries with a BH relevant to explore. The choice of an initially circular binary orbit may be justified by the fact that tidal circularization time-scales are shorter for more compact binaries. In fact, Meibom & Mathieu (2005) found that binaries in open clusters are circularized out to $P \simeq 8-15$ days, which is comparable to or longer than the binary period considered in this study. Nonetheless, three-body encounters involving eccentric binaries are also possible. Because the effective encounter cross section is greater for eccentric orbits (by a factor $\simeq 1 + e$), the overall encounter rate would be higher, but the impact of other parameters (e.g. semimajor axis, impact parameters, see Section 3.2) on the final outcome would be dominant over that of different initial binary eccentricities.

The binary's angular momentum axis is always along the z axis in our simulations. This configuration defines the mutual inclination angle, which is illustrated in Fig. 2 showing the initial configuration of the stellar binary and BH. We examine the outcomes of encounters with various values of i and β : $i = 0, 30^\circ, 60^\circ, 120^\circ$, and 180° , and $\beta = 1/4, 1/2, 1$, and 2 . However, given the relatively high computational costs, we do not simulate encounters with every combination of i and β . Instead, we simulate the encounters of the intermediate-size binaries ($a/a_{\text{RL}} = 6$) with different combinations of $\beta = 1/4, 1/2, 1$ and 2 , and $i = 30^\circ, 150^\circ$. For the smallest and largest binaries ($a/a_{\text{RL}} = 2$ and 9), we only consider $i = 30^\circ$ and 150° , whereas $\beta = 1/2$. In addition, we further examine the dependence of i on the outcome properties by considering $i = 0, 60^\circ, 120^\circ$ and 180° (for $\beta = 1/2$). Last, we also study the impact of the phase angle v (see Fig. 2) on the encounter outcomes. We define v as the initial angle between the line connecting the two stars in the binary and the x -axis. To simulate encounters with different phase angles ($v = 45^\circ, 90^\circ$, and 135°), we initially place the binary with a different phase angle while all other parameters remains fixed.

The initial separation between the binary and BH is $5r_{\text{t}}$.

We summarize the initial parameters considered in our simulations in Table 1. Each of the models is integrated up to $>25(130)t_p$ for the encounters of the $2(20)M_\odot$ binary which it takes to identify the final

outcomes. Here, $t_p = \sqrt{r_p^3/G(M_\bullet + M_b)}$ is the dynamical time at $r = r_p$. The values of t_p for each model is given in Table 1.

3 RESULTS

3.1 Classification of outcomes

When the binary star and the BH encounter, one or both stars can be fully or partially destroyed by the BH, which produces electromagnetic transient (EMT) phenomena. Thus, it is important to define those events in a more quantitative way to classify the outcomes.

We define a full disruption as the case where the star is completely destroyed and no self-gravitating object is left behind. On the contrary, we define a partial disruption as an event where the star loses more than 1 per cent of its mass and the remnant survives. For other cases where the BH forms a wide and isolated binary whose period is much longer than the simulation duration, we determine the fate of the companion star based on its orbit.⁵ More specifically, to see if the star would undergo another disruption event upon return to the BH, we compare the full (equation 16 in Ryu et al. 2020b) and partial (equation 17 in Ryu et al. 2020b) tidal disruption radii to the pericentre distance.⁶ If the pericentre distance is longer than the

⁵We first compare the orbital period $P_{\text{b},\bullet}$ with the order-of-magnitude estimate of the $1 + 2$ encounter time scale for the binary, $t_{1+2} \simeq [n\Sigma\sigma]^{-1}$ where $n = 10^5 \text{ pc}^{-3}$, Σ is the gravitational focusing cross-section and $\sigma = 15 \text{ km/sec}$ is the velocity dispersion. If $P_{\text{b},\bullet} > t_{1+2}$, the star's orbit would be perturbed before it returned to the BH. Although we choose the specific values of n and σ relevant for typical clusters, the particular choices do not affect the classification because t_{1+2} is many orders of magnitude greater than the period of the largest binary in our simulations.

⁶For the calculation of the partial disruption radius, we define the size of a remnant as the distance from the remnant's centre enclosing 99 per cent of its mass. Note that these analytic formulae for the disruption radii are scaled to match the results of the relativistic simulations for TDEs of realistic MS stars by Ryu et al. (2020a), Ryu et al. (2020b), Ryu et al. (2020c), Ryu et al. (2020d). This means the expression for the partial disruption radius would

Table 1. The initial model parameters. The model name (first column) contains the information of key initial parameters: for the model names of $M(1)a(2)\beta(3)i(4)$, (1) binary mass M_b [M_\odot], (2) the initial semimajor axis of the binary a/a_{RL} , (3) the impact parameter β and (4) the inclination angle i in degrees, and for the model names of $M(1)a(2)v(3)$, (1) binary mass M_b [M_\odot], (2) a/a_{RL} , (3) the initial phase angle v in degrees (see Fig. 2). Note that $a_{RL} \simeq 2.64R_\star$. t_p is the dynamical time at $r = r_p$, in units of hours, and r_p is the pericentre distance of the incoming orbit in hours relative to the BH $r_p = [M_\star/(M_1 + M_2)]^{1/3}(1 + e)a/2$.

Model name	M_b	a/a_{RL}	i	v [°]	β	t_p
Unit	M_\odot	–	°	°	–	hours
$M2a6\beta2i30$	2	6	30	0	2	17
$M2a6\beta1i30$	2	6	30	0	1	6.1
$M2a6\beta1/2i30$	2	6	30	0	1/2	2.2
$M2a6\beta1/4i30$	2	6	30	0	1/4	0.76
$M2a6\beta2i150$	2	6	150	0	2	17
$M2a6\beta1i150$	2	6	150	0	1	6.1
$M2a6\beta1/2i150$	2	6	150	0	1/2	2.2
$M2a6\beta1/4i150$	2	6	150	0	1/4	0.76
$M2a2\beta1/2i30$	2	2	30	0	1/2	0.42
$M2a2\beta1/2i150$	2	2	150	0	1/2	0.42
$M2a9\beta1/2i30$	2	9	30	0	1/2	4.0
$M2a9\beta1/2i150$	2	9	150	0	1/2	4.0
$M2a6\beta1/2i0$	2	6	0	0	1/2	2.2
$M2a6\beta1/2i60$	2	6	60	0	1/2	2.2
$M2a6\beta1/2i120$	2	6	120	0	1/2	2.2
$M2a6\beta1/2i180$	2	6	180	0	1/2	2.2
$M2a6v45$	2	6	30	45	1/2	2.2
$M2a6v90$	2	6	30	90	1/2	2.2
$M2a6v135$	2	6	30	135	1/2	2.2
$M20a6\beta2i30$	20	6	30	0	2	50
$M20a6\beta1i30$	20	6	30	0	1	18
$M20a6\beta1/2i30$	20	6	30	0	1/2	6.3
$M20a6\beta1/4i30$	20	6	30	0	1/4	2.3
$M20a6\beta2i150$	20	6	150	0	2	50
$M20a6\beta1i150$	20	6	150	0	1	18
$M20a6\beta1/2i150$	20	6	150	0	1/2	6.3
$M20a6\beta1/4i150$	20	6	150	0	1/4	2.3
$M20a2\beta1/2i30$	20	2	30	0	1/2	1.2
$M20a2\beta1/2i150$	20	2	150	0	1/2	1.2
$M20a9\beta1/2i30$	20	9	30	0	1/2	12
$M20a9\beta1/2i150$	20	9	150	0	1/2	12
$M20a6\beta1/2i0$	20	6	0	0	1/2	6.4
$M20a6\beta1/2i60$	20	6	60	0	1/2	6.4
$M20a6\beta1/2i120$	20	6	120	0	1/2	6.4
$M20a6\beta1/2i180$	20	6	180	0	1/2	6.4

partial disruption radius, the binary would remain a non-interacting one until encountering another intruder. If the pericentre distance is smaller than the partial disruption radius and larger than the full disruption radius, the star would undergo at least one more partial disruption event. Finally, if the pericentre distance is smaller than the full disruption radius, the star would be totally disrupted in one orbital period.

give a shorter distance than the actual value considering other effects, such as stellar spin and hotter interior, that would make the remnants more subject to a partial disruption. On the contrary, because the expression for the full disruption radius is determined by the core density, we expect that the full disruption radius would be a more robust estimate than the partial disruption radius.

Applying these definitions to the outcomes at the end of simulations, we can categorize our models into two classes, depending on whether EMT phenomena are created during interactions.

(i) *Non-disruptive encounter*: This class corresponds to the case where none of the stars are significantly affected during interactions and survive. So the number of surviving stellar objects is always two. This class happens when the binary and the BH encounter with a large impact parameter, resulting in a perturbation of the binary orbit (without being dissociated) or dissociation of the binary into two single stars. In the case with binary dissociation (models $M2a6\beta2i150$ and $M20a6\beta2i150$), the semimajor axis of the binary does not change much (less than 1 per cent in $M20a6\beta2i150$) or increases (by a factor of 1.5 in $M2a6\beta2i150$). But the binary becomes eccentric ($e \simeq 0.7$ in both models). The orbit of the binary relative to the BH becomes a very eccentric ($e_{b,\bullet} \simeq 0.98 - 0.996$) but more bound ($a_{b,\bullet} \simeq 10 - 27$ AU or $P_{b,\bullet} \simeq 8 - 100$ yr). As a result, the impact parameters of the perturbed binaries' orbits relative to the BH are smaller than those of the initial orbit: $\beta \simeq 0.67 - 0.83$ (using the same definition of r_t in Section 2.4).

On the contrary, the case involving binary dissociation usually happens when $1 \leq \beta \leq 2$ and $a/a_{RL} = 6$ (e.g. models $M2a6\beta1/2i30$ and $M10a6\beta1/2i30$), with no strong dependence on stellar mass and inclination. In all these cases, one star is bound and the other star is unbound. This is effectively the Hills mechanism (Hills 1988) (i.e. the mechanism for the dissociation of a binary by a supermassive BH, resulting in bound and unbound stars) by a stellar-mass BH (or 'micro-Hills mechanism'). The ejection velocity at infinity of the unbound stars is about $50 - 270$ km/sec. On the contrary, the bound stars are on eccentric or nearly parabolic orbits ($e_{\star - \bullet} \simeq 0.6 - 1$) with $a_{\star - \bullet} \simeq 0.3 - 5$ AU (or $P_{\star - \bullet} \simeq 11$ d to 3 yr). Based on the comparison between the tidal radius and the pericentre distance of the bound stars' orbit, there are some cases (e.g. Models $M2a6\beta1i30$ and $M2a6\beta1/4i150$) where the bound stars would undergo at least one more partial disruption event upon return.

(ii) *Disruptive encounter*: This denotes the case where at least one star is partially or fully disrupted. The most violent case in our simulations associated with disruptions of two stars: full destruction of both stars (e.g. Models $M20a2\beta1/4i30$ and $M20a2\beta1/4i150$) or a full disruption of one star and at least one partial disruption of the other star (e.g. Model $M10a2\beta1/4i30$). As an example, we depict in Fig. 3 the density distribution at four different times in two models where one star is partially disrupted (e.g. Model $M2a6\beta1/2i30$) and two stars are almost instantaneously disrupted (e.g. Model $M20a2\beta1/2i150$).

This classification is also relevant for EMT formation in multi-body encounters with distinctive observational signatures (see Section 3.3). Between the two classes, EMT would definitely be produced in the second class *disruptive encounter*. Nonetheless, there is still the possibility of the formation of EMT in the first class *non-disruptive encounter*: the micro-Hills mechanism produces a hard binary (e.g. SP^* in Table 2). If the orbit continues to shrink via weak encounters with other star, the BH will generate EM flares by accretion of overflowing gas from the star (Section 4.1.2).

We summarize the classification of each model and the properties of the outcomes in Table 2.

3.2 Dependence of outcomes on encounter parameters

In this study, we examine the dependence of outcomes on a few key encounter parameters: M_b (mass of the initial binary), a (semimajor axis of the initial binary), β (impact parameter), i (inclination angle),

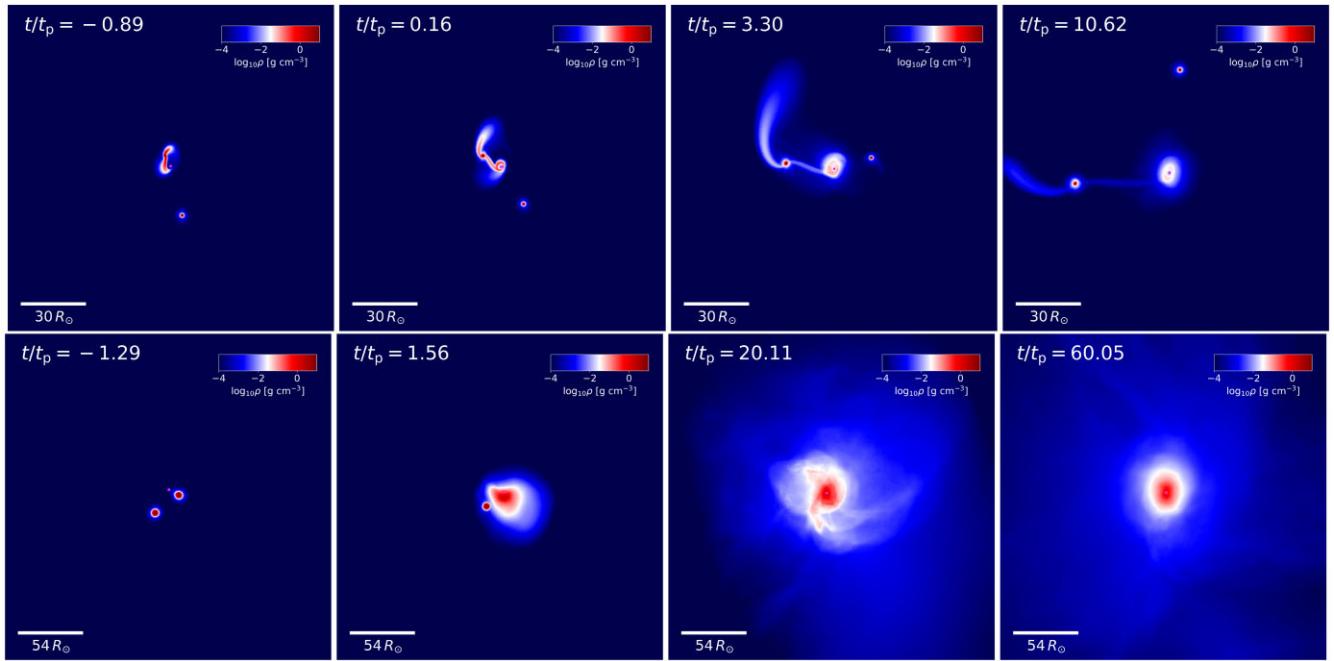


Figure 3. The density distribution in the binary orbital plane in two models (*top*: $M2a6\beta1/2i30$ and *bottom*: $M2a2\beta1/2i150$) at a few different times in units of t_p . The colour bar shows the logarithmic density in g/cm^3 . In the *top* panels, a star is partially disrupted by the BH (magenta dot) at the first pericentre passage, whereas the other star flies away without being significantly perturbed by the BH. In the *bottom* panels, both stars are disrupted by the BH almost simultaneously at the first contact.

and ν (phase angle by varying one parameter at a time, whereas the rest are held fixed). Although our models do not cover the entire parameter space, we can see some clear trends. Our simulations suggest that there are three main parameters which determine the location of the boundary between *non-disruptive* interactions and *disruptive* interactions, namely, the impact parameter β and binary size a/a_{RL} and the phase angle ν .

(i) *Dependence on the impact parameter β* : as one can see from models $M(2, 20)a6\beta(1 - 1/4)i(30, 150)$ in Tables 2, three-body encounters with $\beta > 1$ are *non-disruptive* interactions, whereas those with $\beta \leq 1$ tend to disrupt at least one star. It is not surprising given that β directly determines how close the binary and the BH can approach one another. However, we note that a small impact parameter does not always result in disruption events (e.g. Model $M2a6\beta1i150$ and possibly for a very wide binary, see below). In other words, it is not a sufficient condition but a necessary condition for a disruption event, whereas a large impact parameter is a sufficient condition for *non-disruptive* encounters.

(ii) *Dependence on the semimajor axis a* : the comparison between Models with different a/a_{RL} (e.g. $M2a(2 - 9)\beta1/2i30$) suggests that encounters become more disruptive for smaller a/a_{RL} as long as β is sufficiently small. It is because the interactions move away from the regime of two 1+1 encounters towards the regime of chaotic 2+1 encounter. Stars in large binaries can also be disrupted by the BH [e.g. $M(2, 20)a9\beta1/2i(30, 150)$], but such events become increasingly more like an ordinary TDE by single BH. Interestingly, we found a merger of two partially disrupted stars as an intermediate outcome in one of our models with the smallest a/a_{RL} considered (i.e. Model $M2a2\beta1/2i30$), which is in line with stellar mergers during three-body encounters between stars found in McMillan et al. (1991). Note that the semimajor axis of the initial binary in McMillan et al. (1991) is 4 and $8R_\star$ (cf., $a = 5.3R_\star$ for our smallest binary).

Another consideration related to the dependence of a and β for wide binaries is that if the impact parameter for encounters between a wide binary and a BH is so small that the BH's crossing time across the binary is shorter than the binary period (i.e. the binary is ‘frozen’, whEREAS the BH passes through) and the BH's gravitational force on any of the stars is weaker than the stars' gravitational pull to each other, it may be possible that the BH simply penetrates through the binary in the middle without interacting significantly with any of the stars. This consideration supports that a small impact parameter is a necessary condition for *disruptive* encounters.

(iii) *Dependence on the phase angle ν* : even for a sufficiently small impact parameter and small a , the outcomes can vary depending on the phase angle, as shown in Models with different ν (Models $M2a6\beta1/2i30$ and $M2a6\nu(45 - 135)$). Various outcomes of encounters between a single star and a binary BH for different phase angles was also reported in Paper 1. The statistical likelihood of disruptions are mostly governed by β and a , but it is very important to consider the phase angle for the outcome of an individual encounter case.

(iv) *Dependence on the inclination angle i* : we found no significant dependence on i . The weak dependence on i is at odds with the strong dependence for encounters between a single star and a binary BH found in Paper 1. The comparison is not straightforward given different encounter parameters. However, how chaotic the interactions are (or whether the initial binary is dissociated during encounters) may result in different levels of i –dependence: in encounters between a single star and a BH binary considered in Paper 1, the binary is never dissociated by the star at the first closest approach and the star interacts with the binary mostly once before its disruption or ejection. For that case, the relative velocity between the star and the interacting BH at the first closest approach plays an important role in determining the outcome because that directly governs how long the star is tidally affected by the BH. And the relative velocity is different depending on whether the encounter is

Table 2. The outcomes of each model: (first column) the model name, (second column) the class of encounters, (second column) the class of encounters, *non-disruptive encounter* or *disruptive encounter* (see Section 3.1 for their definitions), and (the remaining columns) the information of outcomes and the BH. The columns for the properties of outcomes and the BH present the type (●) star whose internal structure remains intact, (○) partial disruption, (○) full disruption, $\Delta M/M$ the mass loss in percentile, a the semimajor axis in R_\odot , e the eccentricity, r_p the pericentre distance in R_\odot , P binary orbital period in year and v the ejection velocity at infinity in km/s. v (P) is presented only for unbound (bound) remnants. All the properties of the outcomes presented in the table are measured at the end of the simulation; for the cases where a star forms a bound pair with the BH, we predict their outcomes (circle symbols after \rightarrow) based on their orbit and the disruption (see Section 3.1 for more detailed descriptions). The superscript $*$ indicates interacting binaries found before the end of the simulation. Because of the significant mass loss near pericentre, the orbit of the interacting binaries does not settle until the end of the simulation; the orbital elements are measured at the end of the simulation ($u/\dot{r}_p \geq 80$ for Model $M2a6\beta1/2(i120, 180)$ and ≥ 120 for Model $M2a6\beta1/2(i30)$). We only include v for the single BH. On the contrary, the subscript $*$ indicates potential interacting binary candidates, i.e. non-interacting hard binaries (orbital velocity $\gtrsim 50$ km/sec, c.f., typical velocity dispersion of clusters $\simeq 10 - 15$ km/sec Cohen 1983) (see Section 4.1).

Model name	Class	Type	$\Delta M/M$	$a[R_\odot]$	e	$r_p[R_\odot]$	$P[\text{yr}]$	v	Type	Outcome 1			Outcome 2		
										$\Delta M/M$	$a[R_\odot]$	e	$r_p[R_\odot]$	$P[\text{yr}]$	v
$M2a6\beta2/30$	Non-disruptive	● ● \rightarrow ●*	-0.21	60	0.83	9.8	0.03	-	●●●●●	-0.12	-	-	-	220	-
$M2a6\beta1/30$	Disruptive	● \rightarrow ○	-0.26	85	0.95	3.9	0.05	-	●●●●	-0.14	-	-	-	170.	-
$M2a6\beta1/2i30$	Disruptive	○ \rightarrow ○	-25	78	0.98	1.3	0.05	-	●●●●	-0.18	-	-	-	180	-
$M2a6\beta1/4i30$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-0.39	-	7400	1.0	24	140
$M2a6\beta2/150$	Non-disruptive	-	-0.16	300	0.93	22	0.36	-	●●●●	-0.24	-	-	-	110	-
$M2a6\beta1/150$	Non-disruptive	●● \rightarrow ●●	-0.34	200	0.99	2.5	0.19	-	●●●●	-0.22	-	-	-	120	-
$M2a6\beta1/2i50$	Non-disruptive	○ \rightarrow ○	-3.1	-	-	-	-	180	● \rightarrow ○	-5.8	63	0.99	0.85	0.03	-
$M2a6\beta1/2/30$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-0.07	-	-	-	310	28
$M2a2\beta1/2i50$	Disruptive	○ \rightarrow ○	-7.4	-	-	-	-	-	●●●●	-23	44	0.98	0.71	0.02	-
$M2a9\beta1/2i30$	Disruptive	○ \rightarrow ○	-10	110	0.98	1.9	0.09	-	●●●●	-0.17	-	-	-	130	-
$M2a9\beta1/4i50$	Non-disruptive	●● \rightarrow ●●	-0.24	300	0.99	3.7	0.35	-	●●●●	-0.21	-	-	-	110	-
$M2a2\beta1/2i40$	Disruptive	○ \rightarrow ○	-76	77	0.98	1.5	0.05	-	●●●●	-0.19	-	-	-	190	-
$M2a2\beta1/2i60$	Disruptive	●● \rightarrow ○	-0.67	170	1.00	0.18	0.15	-	●●●●	-0.14	-	-	-	110	-
$M2a9\beta1/2i120$	Disruptive	●● \rightarrow ○	-0.35	2100	1.0	0.30	6.8	-	●●●●	-0.20	1100	1.0	3.6	2.6	-
$M2a6\beta1/2i80$	Non-disruptive	●● \rightarrow ●●	-0.41	140	0.98	3.6	0.12	-	●●●●	-0.34	-	-	-	130	-
$M2a6\beta45$	Disruptive	●● \rightarrow ●*	-0.35	46	0.93	3.1	0.02	-	●●●●	-0.21	-	-	-	260	-
$M2a6\beta1/2i0$	Non-disruptive	●● \rightarrow ●	-0.19	220	0.96	9.9	0.23	-	●●●●	-0.34	-	-	-	110	-
$M2a6\beta1/2i60$	Disruptive	○ \rightarrow ○	-12	150	0.99	2.3	0.13	-	●●●●	-0.18	-	-	-	91	-
$M2a6\beta1/2i120$	Non-disruptive	●● \rightarrow ●●	-0.55	87	0.62	33	0.05	-	●●●●	-0.42	-	-	-	93	-
$M2a6\beta1/2/30$	Non-disruptive	○ \rightarrow ○	-0.36	220	0.88	26	0.19	-	●●●●	-0.59	-	-	-	97	-
$M2a6\beta1/30$	Disruptive	-	-	-	-	-	-	-	●●●●	-0.93	-	-	-	140	60
$M2a6\beta1/2i30$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	0.08	280	0.86	40	0.27	-
$M2a6\beta1/4i30$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-0.42	2400	0.98	53	8.5	-
$M2a6\beta2/150$	Non-disruptive	-	-	94	0.70	28	0.06	-	●●●●	-0.59	-	-	-	48	-
$M2a6\beta1/1150$	Non-disruptive	●● \rightarrow ●●	-0.65	140	0.87	18	0.09	-	●●●●	-0.59	-	-	-	15	-
$M2a6\beta1/2i30$	Disruptive	○ \rightarrow ○	-	40	0.76	9.5	0.01	-	●●●●	-0.59	-	-	-	230	-
$M2a6\beta1/4i150$	Disruptive	● \rightarrow ●*	-0.56	-	-	-	-	-	●●●●	-0.59	-	-	-	22	-
$M2a6\beta2/1/230$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-0.59	-	-	-	-	-
$M2a6\beta1/2i50$	Disruptive	○ \rightarrow ○	-1.8	-	-	-	-	-	●●●●	-0.41	140	0.69	42	0.09	-
$M2a6\beta1/2i150$	Disruptive	○ \rightarrow ○	-1.5	71	0.82	13	0.04	-	●●●●	-0.65	53	0.98	0.98	0.02	-
$M2a6\beta1/2i60$	Disruptive	○ \rightarrow ○	-1.4	-	-	-	-	-	●●●●	-1.08	-	-	-	170	77
$M2a6\beta1/2i120$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-12	27	0.74	7.0	0.01	-
$M2a6\beta1/2i20$	Disruptive	○ \rightarrow ○	-	-	-	-	-	-	●●●●	-4.4	24	0.48	13	0.01	-

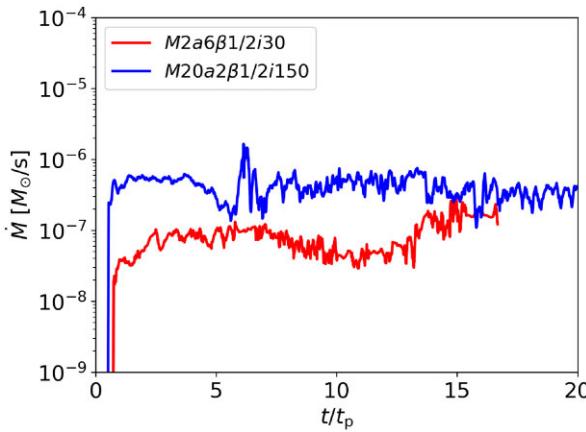


Figure 4. The accretion rate \dot{M} for the two models, $M2a6\beta1/2i30$ (red) and $M20a2\beta1/2i150$ (blue), as a function of t/t_p since the first closest approach. Here, $t_p \simeq 2.2$ (1.2) h for Model $M2a6\beta1/2i30$ ($M20a2\beta1/2i150$).

in a prograde or retrograde direction for given parameters. On the CONTRARY, in the simulations of this study, the binary is always dissociated (except for the two models with $\beta = 2$), followed by more chaotic interactions, which likely removes the memory of the initial approach.

(v) (vi) *Dependence on the binary mass M_b :* The dependence of M_b on the types of outcomes appears to be weak. However, it actually means that the mass (and thus energy) budget for radiation from debris is greater for the disruption of more massive stars (i.e. similar $\Delta M/M$ for larger M_*). Furthermore, the momentum kick given to the BH is greater for encounters with more massive binaries. In particular, the ejection velocity of the BH is $\sim 10 - 30$ km/sec in encounters with the $2 M_\odot$ binary, whereas it is in the range $\sim 50 - 200$ km/sec in encounters with the massive stars of our simulations. Note that the ejection velocity of the unbound stars from the $2 M_\odot$ binary tends to be greater than that for more massive unbound stars, but the degree is not as significant as the ejection velocity of the BH.

3.3 Impact on the black hole properties

The immediate impact of close encounters with binary stars on the BH is momentum kick, growth in mass and, therefore, spin evolution.

The ejection velocity of the BH in most cases is a few tens of km/s, but there are a few cases with an ejection velocity of $\sim 100 - 200$ km/sec (see Table 2). Those velocities are either larger than or comparable to the escape velocity of globular clusters (i.e. a few to 180 km/s; Antonini & Rasio 2016; Gnedin et al. 2002). We find that the only significant correlation of the ejection velocity is with the mass of the binary that encounters the BH, as discussed in Section 3.2: encounters with a more massive binary results in higher ejection velocity of the BH.

If a star is disrupted at least once during encounters, the ejected BH is surrounded by an accretion flow made from stellar debris, as shown in Fig. 3. The gas accretes onto the BH and the BH grows in mass and spins up. Typically, the accretion rate remains nearly the same at super-Eddington $\dot{M} \simeq 10^{-8} - 10^{-6} M_\odot/\text{sec}$ once an accretion disk forms, which is many orders of magnitude higher than typical Roche Lobe overflow rates in interacting BH binary systems (Savonije 1978). This trend is illustrated in Fig. 4 showing the time evolution of the accretion rate for two models, one *single disruption* case (Model $M2a6\beta1/2i30$) and one *double disruption* case (Model

$M20a2\beta1/2i150$). Note that these are the same models shown in Fig. 3.

4 DISCUSSION

4.1 Astrophysical Implications

4.1.1 Varieties of transient phenomena

We have reported that dynamical interactions between binary stars and a single BH can produce a variety of EMT phenomena: a single partial disruption event, a full disruption event, multiple partial, or full disruption events. Even multiple disruption events can be subdivided into two classes: almost instantaneous double disruption events and successive partial disruption events of bound remnants. Furthermore, an interacting star-BH binary (see Section 4.1.2) can form, whereas the other star is fully disrupted (e.g. Model $M20a6\beta1/2i120$) or the other star is ejected via the Hills mechanism (Hills 1988) (or ‘micro-Hills’ mechanism). Dynamical interactions can create ‘collision’-like disruption events where disk formation is prompt. Such varieties constitute the most significant difference from the ordinary TDEs, which are disruptions of a single star by a single BH. In particular, double disruption events can never happen in ordinary TDEs.

Such varieties of EMT phenomena and the chaotic nature of three-body interactions indicate that light curves of TDEs or almost collision-like disruptions in three-body interactions should look qualitatively different from those of ordinary TDEs. Detailed quantitative predictions of light curves are beyond the scope of this paper and reserved to future work. Nonetheless, we can still make qualitative predictions for the expected observational signatures of double disruption events. For the case of rapid double full disruption, almost instantaneous disk formation may result in a sudden increase in luminosity to super-Eddington in EUV and X-rays (luminosity $\propto T^4$ -weighted average of temperature $\simeq 10^6 - 10^7$ K), which, if the luminosity $\propto \dot{M}$, may remain roughly constant up to a few months [$\simeq O(1 M_\odot)/O(10^{-6} M_\odot/\text{sec})$].

Furthermore, the total energy budget ($\propto M_*$) is greater than for a single full disruption. However, if super-Eddington accretion causes significant outflows so that some fraction of gas becomes unbound, the total radiated energy would not be as large as expected simply based on the total mass of the disrupted stars, indicating the duration of the burst would be shortened.

For the cases with a full disruption followed by more than one partial disruption event, light curves would reveal a higher peak followed by at least one other less intense burst. If partial disruptions occur successively, the less bright bursts would show a quasi-periodic behaviour with a modulation time-scale \simeq the orbital period of the remnants. In our simulations, the shortest period for multiple partial disruption events is \simeq days. In fact, in two of our models (Models with ** in Table 2, $M20a6\beta1/2i120$ and $M20a6\beta1/2i180$), three-body encounters create a binary system whose eccentricity is low ($e = 0.74$ and 0.48, respectively) for ordinary disruption events where the orbit is typically approximated to be parabolic. We will discuss these low- e binary systems in Section 4.1.2.

4.1.2 Black hole high- and low-mass X-ray binary

Three-body interactions may result in the exchange of a member of the binary with the intruder object (Valtonen & Karttunen 2006) during chaotic interactions and via the micro-Hills mechanism. The member exchange also occurs between the binary star and the BH

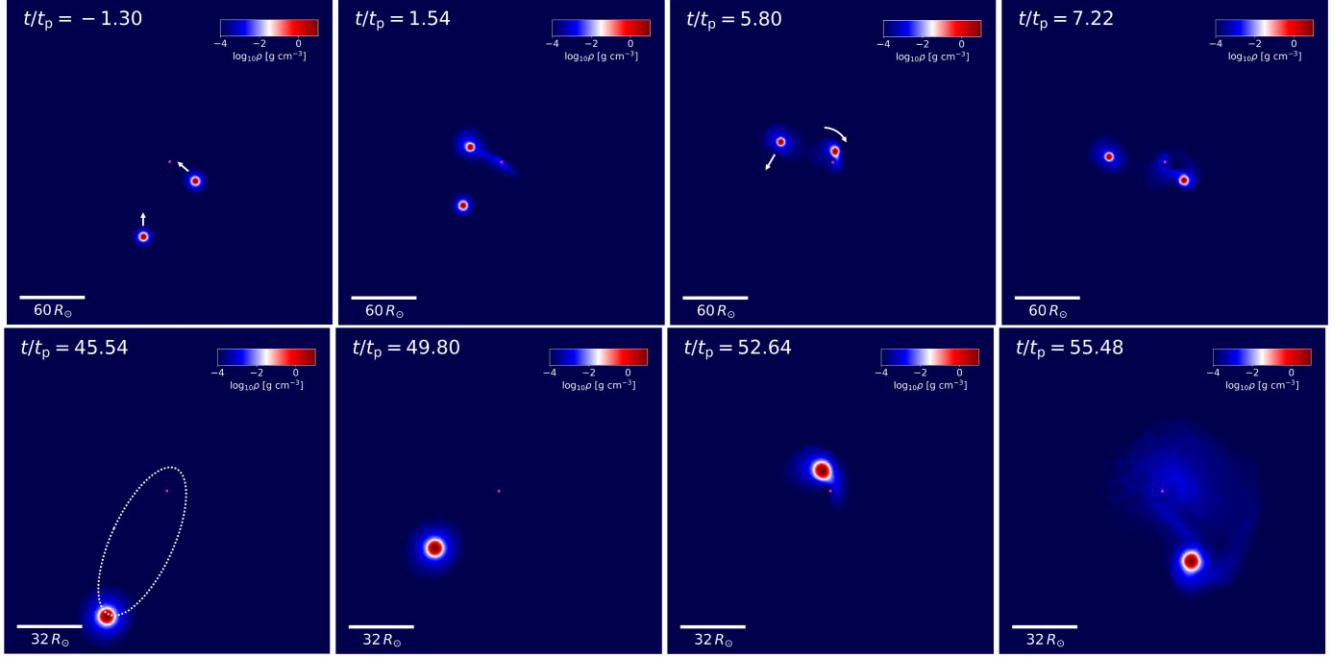


Figure 5. Series of snapshots, centred at the BH, for Model $M20a6\beta1/2i120$ showing a star-BH binary at a few different times in units of $t_p \simeq 2.2$ hours. The colour bar shows the logarithmic density projected in the $x - y$ plane in g/cm^3 . In the lower-left panel, the dotted eccentric circle approximately shows the orbital trajectory of the star relative to the BH.

in our simulations, resulting in a star bound to the BH (e.g. Models $M2a6\beta2i30$ and $M2a6\beta1/2i150$). In some of the models the orbit is nearly parabolic (e.g. Models $M2a2\beta1/2i30$ and $M20a6\beta1/2i30$) so that the bound star would be fully disrupted upon return. This would look exactly like an ordinary TDE. On the contrary, we also found that eccentric star-BH binaries form. In the two Models (Models $M20a6\beta1/2i120$ and $M20a6\beta1/2i180$), we actually simulate multiple quasi-periodic episodes of mass transfer at pericentre in the binary with a $\gtrsim 8 M_\odot$ star. In Fig. 5, we show the density distribution, projected on the $x - y$ plane, in Model $a20a6\beta1/2i120$ at a few different times until an eccentric binary forms ($a \simeq 27 R_\odot$, $e = 0.74$, $P \simeq 3$ d). The repeated mass loss at pericentre leads to quasi-periodic bursts in the mass accretion rate with peak $\dot{M} \simeq 10^{-5} M_\odot/\text{sec}$ on a timescale of the orbital period, as shown in Fig. 6. In that panel, we also show the accretion rate for Model $M20a6\beta1/2i180$ where an eccentric binary with $a \simeq 24 R_\odot$, $e \simeq 0.5$ and $P \simeq 2.5$ days forms. Interestingly, this less eccentric binary does not reveal modulations of the accretion rate. Like the binary in the Model $a20a6\beta1/2i120$, the star is periodically disrupted. But because of not sufficiently large eccentricity and semimajor axis, the mass stripped from the star is simply added to an existing accretion flow around the BH. Our simulations confirm the formation of eccentric high-mass X-ray binaries in three-body encounters.

As shown in Fig. 7, the interacting binaries circularize; both the semimajor axis and eccentricity decrease over time. However, the two circularizing binaries show somewhat different trends in the evolution of a and e . For the wider and more eccentric binary (blue), we see a sudden drop in both parameters (like a step function) whenever the binary undergoes a mass loss episode near pericentre. Furthermore, the absolute time derivatives of a and e increase slightly until the end of the simulation, indicating the circularization accelerates. On the other hand, both parameters for the smaller and more circular binaries (green) decrease like a damped oscillator at a rate that decreases over time. This may indicate that the orbit

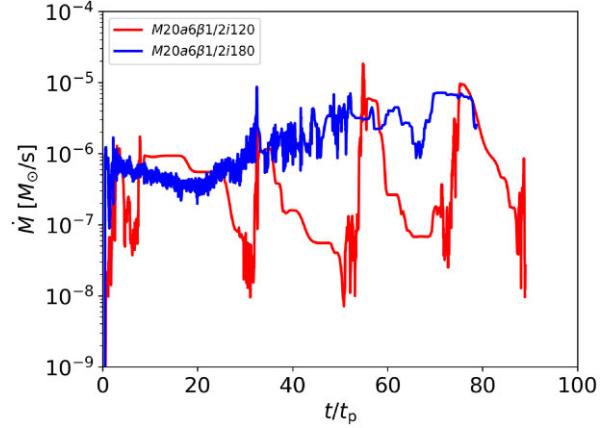


Figure 6. The accretion rate \dot{M} (top) for interacting binaries found in two Models $M20a6\beta1/2i120$ (red) and $M20a6\beta1/2i180$ (blue). The $t_p \simeq 6.4$ hours is the dynamical time at the pericentre distance of the orbit of the initial binary relative to the BH.

evolution of an interacting binary depends on the orbital parameters at the time two companions in the binary start to interact.

We also found eccentric hard binaries consisting of a $\simeq 1 M_\odot$ star and the $20 M_\odot$ BH (Models $M2a6\beta2i30$, $M2a6\beta1i30$, $M2a6v45$ and $M2a6v135$, marked with the superscript \star in Table 2). The orbital period of three of the binaries are 8 - 20 days and the widest one has a period of $\simeq 100$ days. Two of those binaries have pericentre distances ($r_p \simeq 3 - 4 R_\odot$) small enough for close interactions between the star and the BH at pericentre (Models $M2a6\beta1i30$ and $M2a6v45$). Even the two other larger binaries are found to be hard (the orbital velocity at apocentre $\gtrsim 100 \text{ km/sec}$, which is larger than typical velocity dispersion $\simeq 10 - 15 \text{ km/sec}$ of globular clusters Cohen 1983), meaning that unless a strong encounter significantly disrupt

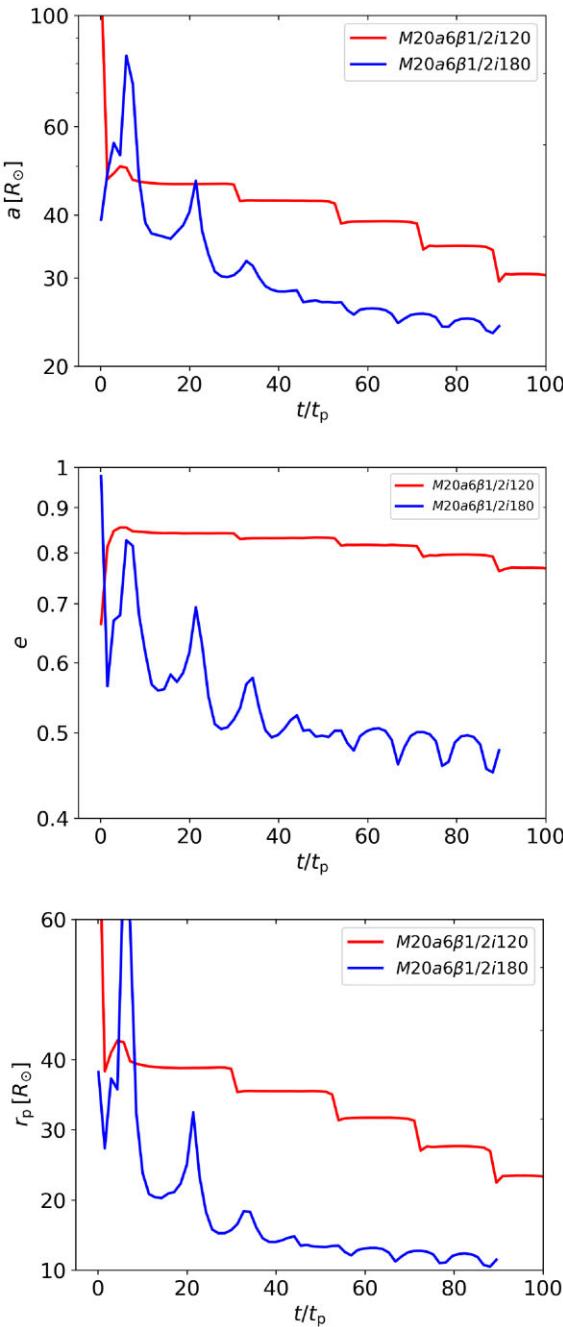


Figure 7. The orbital perimeters, semimajor axis (top), eccentricity (middle) and pericentre distance (bottom) for interacting binaries found in two Models $M20a6\beta1/2i120$ (red) and $M20a6\beta1/2i180$ (blue). $t_p \simeq 6.4$ hours is the dynamical time at the pericentre distance of the orbit of the initial binary relative to the BH.

the binary orbit, they may be hardened further by weak encounters to compact binaries where the star transfers its mass to the BH.

The formation of hard binaries with a $\simeq 1 M_\odot$ star and a $20 M_\odot$ BH may have interesting implications for the formation of detached BH-star binaries (Giesers et al. 2018) and BH-LMXBs (Casares & Jonker 2014). Our simulations confirm the possibility that even a single strong dynamical encounter between a BH and a circular solar-type binary star can result in the formation of a BH-star binary which is already compact enough for mass transfer or is sufficiently hard that it would potentially evolve into an X-ray binary.

Considering that we have not explored the entire range of initial conditions for 3-body encounters, it will be of great interest to extend our simulations to fully identify the parameter space where BH low-mass X-ray binaries can form. We reserve this exploration for future work.

4.1.3 Runaway stars and isolated wandering black holes

Runaway stars are a population of fast-moving O and B type stars at $\gtrsim 30$ km/sec (Blaauw 1961; Stone 1979). Two competing mechanisms for their formation have been suggested: (a) ejection of a star from a binary system when its companion goes off as a supernova explosion (e.g. Zwicky 1957; Blaauw 1961); (b) ejection of a star during dynamical interactions of binaries with other stars in a star cluster (e.g. Poveda, Ruiz & Allen 1967; Gies & Bolton 1986; Ryu, Leigh & Perna 2017) or resulting from the interaction between infalling star clusters and massive bBHs in Galactic Centres (e.g. Capuzzo-Dolcetta & Fragione 2015; Fragione, Capuzzo-Dolcetta & Kroupa 2017).

We find that the three-body encounters involving the relatively tight binary tend to create a rapidly moving ejected star at $v \simeq 50 - 310$ km/sec. Those velocities are large enough to escape globular clusters (the escape velocity of globular clusters \simeq a few to 180 km/sec; Antonini & Rasio 2016; Gnedin et al. 2002).

Interestingly, in the two cases with the initially $20 M_\odot$ binary where one star is fully disrupted and the other star is ejected at a high speed (models $M20a6\beta1/2i30$ and $M20a6\beta1/2i60$), the single BH gains a large momentum, moving at a high speed $\sim 60 - 80$ km/s. These velocities tend to be a factor of 2–3 smaller than the ejection velocity of the unbound star. We attribute this smaller BH velocity to the (not perfect) cancellation of the momentum kicks due to a disruption event and the ejection of a star via the slingshot mechanism. We also estimate that the angle between the unbound star and the BH is around $140^\circ - 180^\circ$. Despite the small sample number, this suggests that the formation of a runaway star via three-body interactions between O-type stars and stellar-mass BHs may be accompanied by the formation of an active BH moving away from the unbound star at a still high speed. This finding has two main implications. First, a detection of a runaway star and an active BH moving away from each other, if the potential formation site can be identified via, e.g. integration of the trajectories backward in time, may serve as strong evidence of close three-body encounters between stars and BHs. Second, if a rapidly moving active single BH is observed, it can be used as a guide to locate a runaway star. Lastly, identifying either transient phenomena or wandering isolated BHs created due to three-body encounters can be used to mutually constrain each other, using the rates of each outcome measured by three-body scattering experiments which carefully take into account hydrodynamic effects.

4.1.4 Long-term evolution of partially disrupted stars

Some of the runaway stars were partially disrupted before being ejected. Upon ejection, not long after the disruption event, the partially disrupted runaway stars are characterized by differential rotation and a hotter thermodynamic state. Ryu et al. (2020c) showed similar features of partial disruption remnants by supermassive BHs. Furthermore, we find that the differential rotation and non-spherical distribution of chemical components during closest passage results in chemical mixing inside the remnants (e.g. an increase in the core H fraction by 23 per cent in a remnant of the initially $1 M_\odot$ star that lost 23 per cent of its mass).

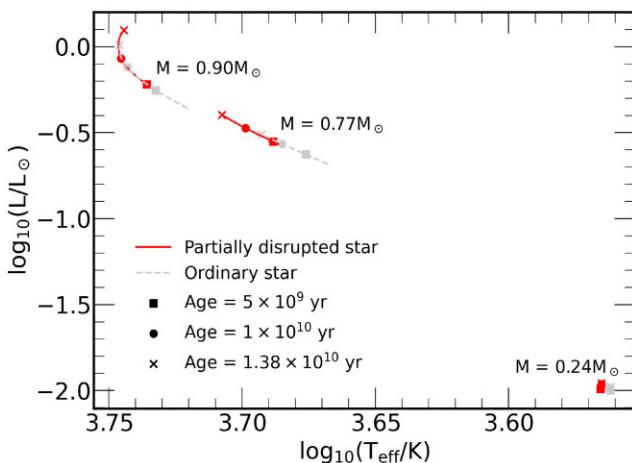


Figure 8. The Hertzsprung–Russell diagram for three partially disrupted stars (red solid lines) which lost $\simeq 10$ per cent ($M = 0.9 M_{\odot}$, Model $M2a9\beta1/2i30$), $\simeq 25$ per cent ($M = 0.77 M_{\odot}$, Model $M2a6\beta1/2i30$) and $\simeq 76$ per cent ($M = 0.24 M_{\odot}$, Model $M2a6\beta1/2i0$) of their mass, respectively. For comparison, we depict the lines for ordinary solar-metallicity stars of the same mass using grey dotted lines. The three markers indicate the age of the star: square (5 Gyr), circle (10 Gyr), and cross (13.8 Gyr).

Although the mass loss of the ejected stars during interactions with the BH is not found to be significant for the encounter parameters considered in this study, severe partial disruptions of a star, followed by its ejection, are in principle possible. Hence, given the possibility of forming runaway stars that underwent a partial disruption, we examined the long-term evolution of three partially disrupted stars: an unbound star with mass loss of 10 per cent, and two bound stars with mass loss of 25 per cent and 76 per cent. As the first attempt for this approach, we only examine the impact of the additional chemical mixing on the structure of the star in this paper. We consider a non-rotating ordinary star with the post-encounter mass in MESA whose chemical composition profile is given by that of a partially disrupted star. We first relax the star until it reaches thermal equilibrium, and evolve it by solving the fully coupled structure and composition equations.

Fig. 8 shows the evolutionary paths of the three partially disrupted stars and ordinary Solar-metallicity stars of the same mass in the Hertzsprung–Russell diagram. The general trend is that the partially disrupted stars, once settled, are hotter and brighter than ordinary stars of the same mass at the same age (at most by a factor of a few) because of a higher He fraction. We should note that we only take the modified chemical composition profile of the partially disrupted stars into account. It is important to include all other thermodynamics quantities properly altered by tidal interactions, such as density, internal energy and angular velocity. Due to the encounter the partially disrupted stars can be spun up beyond their critical rotation rates. At this point, we would then expect the partially disrupted stars to shed angular momentum via mass loss to bring its rotation rate below critical. At this time we have ignored the effect of rotation and rotation-induced mass loss in Fig. 8. We will perform a more systematic study for this evolution with proper modelling of the internal structure, rotation, and mass-loss in our future work.

4.2 Event rate

As an order of magnitude estimate, the differential rate of a single BH encountering a binary star may be estimated as $d\mathcal{R}/dN_{\bullet} \simeq n \Sigma v_{\text{rel}}$.

Here, n is the the binary number density in the vicinity of the BH, $\simeq f_b n_s$ where f_b is the binary fraction, $\simeq 0.1$ for globular clusters (Ji & Bregman 2013; Ivanova et al. 2005; Dalessandro et al. 2011), and n_s is the number density of single stars near the centre. And v_{rel} is the relative velocity between the binary and the BH and Σ is the encounter cross-section. In the gravitational focusing regime ($\sqrt{GM_{\bullet}/r_p} \gg v_{\text{rel}} \simeq \sigma$ where r_p is closest distance between the binary and the BH and σ is the velocity dispersion), $\Sigma \simeq \pi GM_{\bullet}r_p/\sigma^2$. Motivated by the fact that encounters tend to become disruptive when $r_p < a(1 + e)$ and $a \lesssim 9a_{\text{RL}} \simeq 24 R_{\odot}$ (for $M_{\bullet} = 1 M_{\odot}$), we consider $\Sigma \simeq \pi GM_{\bullet}a(1 + e)/\sigma^2$. Then, we find that $d\mathcal{R}/dN_{\bullet}$ is expressed as

$$\begin{aligned} \frac{d\mathcal{R}}{dN_{\bullet}} &\simeq \frac{n\pi GM_{\bullet}a}{\sigma}, \\ &\simeq 10^{-10} \text{ yr}^{-1} \left(\frac{f_b}{0.1} \right) \left(\frac{n_s}{10^5 \text{ pc}^{-3}} \right) \left(\frac{M_{\bullet}}{40 M_{\odot}} \right) \\ &\quad \times \left(\frac{a}{24 R_{\odot}} \right) \left(\frac{\sigma}{15 \text{ km sec}} \right)^{-1}. \end{aligned} \quad (4)$$

Assuming more than tens of single stellar-mass BHs existing in dense clusters at present day (Morscher et al. 2015; Askar, Arca Sedda & Giersz 2018; Kremer et al. 2020)⁷ and $\simeq 150$ globular clusters in Milky Way (Harris 2010), the rate of disruptive three-body encounters per Milky Way-like galaxy is,⁸

$$\mathcal{R} \simeq 10^{-6} \text{ yr}^{-1} \left(\frac{N_{\bullet}}{15000} \right) \left(\frac{f_b}{0.1} \right) \left(\frac{n_s}{10^5 \text{ pc}^{-3}} \right) \left(\frac{M_{\bullet}}{40 M_{\odot}} \right) \\ \times \left(\frac{a}{24 R_{\odot}} \right) \left(\frac{\sigma}{15 \text{ km sec}} \right)^{-1}. \quad (5)$$

We now compare the relative frequency of the three-body encounters for different binary mass M_b . The number of encounters of an equal-mass binary of mass $2M_1$ relative to that of an equal-mass binary of mass $2M_2$ in the lifetime of the binary t_{life} can be expressed as $n_1 \Sigma_1 v_{\text{rel},1} t_{\text{life},1} / [n_2 \Sigma_2 v_{\text{rel},2} t_{\text{life},2}]$.

To start with, let us ignore the impact of mass segregation on the number density of stars with mass. This allows us to assume that $v_{\text{rel},1} = v_{\text{rel},2}$. Adopting the Kroupa stellar mass function (Kroupa 2002) with the cut-off mass $0.5 M_{\odot}$ yields $n_1/n_2 = (m_1/m_2)^{-2.3} (f_{b,1}/f_{b,2})$. The encounter cross-section is again $\Sigma \propto r_p$, giving $\Sigma_1/\Sigma_2 \simeq r_{p,1}/r_{p,2} \simeq a_1/a_2$ where a is the binary semimajor axis. Finally, let us assume that t_{life} is comparable to the lifetime of the star $\propto m^{-3}$.

Combining all of the above we obtain

$$\begin{aligned} \frac{N_1}{N_2} &\simeq \frac{n_1 \Sigma_1 v_{\text{rel},1} t_{\text{life},1}}{n_2 \Sigma_2 v_{\text{rel},2} t_{\text{life},2}} \\ &\simeq \left(\frac{m_1}{m_2} \right)^{-2.3} \left(\frac{f_1}{f_2} \right) \left(\frac{a_1}{a_2} \right) \left(\frac{m_1}{m_2} \right)^{-3} \\ &\simeq \left(\frac{m_1}{m_2} \right)^{-5.3} \left(\frac{f_1}{f_2} \right) \left(\frac{a_1}{a_2} \right). \end{aligned} \quad (6)$$

This suggests that the encounters of a $2 M_{\odot}$ binary (say, $a = 0.05$ AU and $f_1 = 0.1$) with a $20 M_{\odot}$ BH are much more frequent than those of $20 M_{\odot}$ binaries (say, $a = 0.1$ AU and $f_2 = 1$).

⁷Askar et al. (2018) finds that $\simeq 10$ per cent of the retained BH population forms a binary, which is not large enough to significantly affect our order of magnitude estimate for \mathcal{R} .

⁸Note that calculating \mathcal{R} involves an integration of the differential formation rate with several important factors, such star formation history. Thus for a more precise estimate of \mathcal{R} , a more careful consideration of cluster evolution history is required.

If mass segregation is taken into account, the number density of more massive binaries around the BHs (which get to the centre first) would be larger. However, v_{rel} would be smaller for more massive stars by the inverse square-root of the mass ratio. Also, t_{life} would be shortened by the formation time scale of the BH because the encounters considered only occur after BHs form. This means that t_{life} for more massive stars would be shortened even more. For example, the lifetime of a $\simeq 70 M_{\odot}$ star, which would collapse to a $20 M_{\odot}$ BH (fig. 2 in Spera & Mapelli 2017), is $\simeq 3$ Myr. Hence t_{life} for a $20 M_{\odot}$ binary is $20 - 3 = 17$ Myr, whereas that for a $2 M_{\odot}$ binary is $10^4 - 3 \simeq 10^4$ Myr. This suggests that unless the number density for $20 M_{\odot}$ binaries is a few orders of magnitude larger than that for $2 M_{\odot}$ binaries, the encounters between a $2 M_{\odot}$ binary and a $20 M_{\odot}$ BH are more frequent than those between a $20 M_{\odot}$ binary and a BH of the same mass.

However, a close encounter of a BH with more massive binaries (which results in at least one full disruption) is preferentially more likely to be detected since the event is expected to be brighter. The timescale of the peak luminosity will also influence the observability. However, if the timescale is proportional to the peak mass return time, then it has a weak dependence on the stellar mass; only a factor of 2 difference between a $1 M_{\odot}$ star and a $10 M_{\odot}$ star. If the luminosity has the same dependence on the stellar mass as the mass return rate, the luminosity is $\propto m^{0.7}$. Considering the strong dependence of N_1/N_2 on the binary mass, even the number of observable encounters is still likely to be more frequent for encounters involving low-mass stars.

5 SUMMARY AND CONCLUSIONS

In this work, we investigated the outcomes of three-body encounters between a binary star and a BH using the moving-mesh hydrodynamics code AREPO. In particular, we focused on identifying all possible types of transient phenomena produced in the three-body encounters, and studied their properties. We consider a wide range of encounter parameters, i.e. the binary mass, the binary size, the impact parameter, the inclination angle, and the phase angle. However, given the inhomogeneous parameter sampling and the limited number of simulations, our work is not suitable for statistical analysis. Nonetheless, we found some clear qualitative dependence of outcomes and their properties on the system parameters. Our results can be summarized as follows.

- We identified two different types of outcomes:

(i) *Non-disruptive* encounters where neither star is disrupted, which include a weak perturbation of the binary orbit (without being dissociated) and a dissociation of the binary, creating one unbound and one bound stars.

(ii) *Disruptive* encounters where at least one of the stars is fully or partially disrupted. Prompt electromagnetic transient phenomena would be produced in the second class. Although electromagnetic radiation is not expected during encounters of the first type, the formation of hard binaries in that class indicates the possibility of electromagnetic transient phenomena from X-ray binaries at later times if the binaries' orbit continues to shrink via weak encounters.

• The most important factors to determine the parameter space for transient phenomena are the impact parameter, the binary size, and the phase angle. Our simulations suggest that transients are more likely to be created in encounters of smaller binaries with smaller impact parameters. The dependence on the three parameters is:

(i) *Impact parameter*: a large impact parameter is a necessary condition for *non-disruptive* encounters whereas a small impact parameter is a sufficient condition for disruptive encounters.

(ii) *Binary size*: encounters involving an initially smaller binary more likely lead to transient phenomena than non-transient phenomena because the interactions become more chaotic. Disruption events can still happen for a wide binary if the impact parameter is $\simeq a/2$ or interactions become chaotic.

(iii) *Phase angle*: the outcomes can vary, depending on the phase angle even when all other parameters are fixed. The statistical likelihood of transient formation would be mostly governed by the binary size and impact parameter, but the phase angle is an important factor to determine the outcomes of an individual encounter case.

- The transient formation in the three-body encounters has important implications:

(i) *Varieties of transients*: three body-encounters can produce transient phenomena with unique observational features. In particular, multiple disruption events, such as instantaneous double disruption events or a full disruption event followed by repeated partial disruption events can never happen in ordinary tidal disruption events. Furthermore, a nearly collision-like disruption during chaotic interactions would have different light curves than those of ordinary tidal disruption events. Furthermore, the total energy budget available for radiation would be larger.

(ii) *BH X-ray binary*: we found eccentric hard binaries with $\simeq 1 M_{\odot}$ star form via the three-body encounters, two of which have a pericentre distance small enough for the star and BH to interact at pericentre. The formation of these particular star-BH binary systems may serve as evidence for the dynamical formation of BH low-mass X-ray binaries whose formation mechanism remains questionable. On the other hand, we found that a compact nearly equal-mass star-BH binary with mass $40 M_{\odot}$ can form, confirming the formation of high-mass X-ray binaries via three-body interactions between a binary star and a BH.

(iii) *Runaway stars and isolated BHs*: one frequent outcome in our simulations is the formation of single (both $1 M_{\odot}$ and $10 M_{\odot}$) stars ejected at $50 - 270$ km/sec via chaotic interactions or the Hills mechanism (Hills 1988) (which we call 'micro-Hills' mechanism). In particular, the velocities of the ejected stars with $M_{\star} \simeq 10 M_{\odot}$ are comparable to typical velocities of runaway stars ($\gtrsim 30$ km/sec). In the two cases where one star is ejected at a high speed and the other star is fully disrupted, the single BH is also ejected at a velocity of $30 - 80$ km/sec in the opposite direction to the unbound star, becoming active due to the accretion of the surrounding stellar debris. A detection of a runaway star or wandering active BH could mutually constrain the population of each other.

In Paper 1, we investigated the production of transient events in three-body encounters between a single star and a binary BH using hydrodynamics simulations. We found that various types of transients can form in each particular type of encounter, with their own unique observational signatures, such as light curves with periodic modulations. This work further adds to the significance of three-body encounters as a formation channel of transient phenomena.

ACKNOWLEDGEMENTS

The authors are grateful to the anonymous referee for constructive suggestions and comments which helped us improve the manuscript. This research project was conducted using computational resources (and/or scientific computing services) at the Max-Planck Comput-

ing & Data Facility. The authors would like to also thank Stony Brook Research Computing and Cyberinfrastructure, and the Institute for Advanced Computational Science at Stony Brook University for access to the high-performance SeaWulf computing system, which was made possible by a \$1.4M National Science Foundation grant (#1531492). RP acknowledges support by NSF award AST-2006839.

DATA AVAILABILITY

Any data used in this analysis are available on reasonable request from the first author.

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