

Reversible Light-Induced Enhancement of Photoluminescence Lifetime and Intensity in Perovskite-Phase CsPbI₃ Nanocrystals

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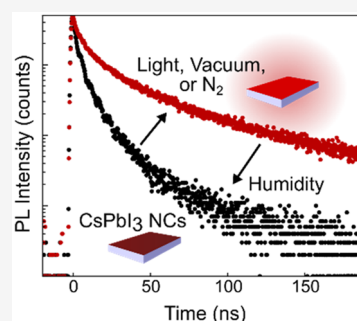
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ABSTRACT: Light-induced changes in photophysical and electronic properties in metal halide perovskites can affect their performance in photovoltaic devices, light-emitting diodes, and other applications. Here we reveal that light induces a slow, reversible enhancement in photoluminescence (PL) lifetime and intensity in films of perovskite-phase CsPbI₃ nanocrystals. When films of CsPbI₃ nanocrystals stored in air are photoexcited, their PL lifetime and intensity increase by as much as a factor of 5 over the course of 20–30 min. Several hours later, without additional light excitation, the initial PL lifetime and intensity return. Placing the films under vacuum or nitrogen for several minutes was also found to increase the PL lifetime and intensity. We propose a model of slow, humidity- and light-sensitive surface states in perovskite-phase CsPbI₃ nanocrystals.



INTRODUCTION

Lead halide perovskite nanocrystals exhibit a variety of potentially useful properties, including size- and composition-tunable band gaps,^{1–4} narrow photoluminescence (PL) emission line widths,^{5–7} near-unity PL quantum yield (QY),^{8–13} superfluorescence,¹⁴ and tolerance to defects.^{15–18} They have been used to fabricate optical detectors,^{19–21} photovoltaic devices (PVs),^{22–34} and light-emitting diodes (LEDs).^{35–41} The performance and design of these devices depend on charge carrier recombination kinetics, which have been the focus of many studies. For example, time-resolved photoluminescence (trPL) spectroscopy measurements have shown that the radiative lifetimes of CsPbX₃ (X = Cl, Br, I) nanocrystals increase with decreasing temperature, which is unique among most semiconductor quantum dots and has been ascribed to a bright triplet state.⁴²

Light-induced changes in the photophysical and electronic properties of metal halide perovskites can also occur. Tsai et al.⁴³ showed that thin films of formamidinium methylammonium cesium lead iodide (FA_{0.7}MA_{0.25}Cs_{0.05}PbI₃) in photovoltaic devices exhibited a 2-fold increase in PL intensity, a PL peak red-shift of about 5 nm, and an increase in open-circuit voltage over the course of about 120 min of 1 sun illumination. This light-induced response was attributed to an elongation of Pb–I bonds, although the details of these results and this mechanism are still under discussion.⁴⁴ The PL intensity and lifetime of metal halide perovskite nanocrystals are also known to be affected by surface chemistry.^{8–10,13,15,45} Most commonly, photoexcitation can decrease the PL, as has been observed from CsPbI₃ nanocrystals under illumination.^{46,47} For CsPbBr₃ nanocrystals under UV illumination, Ruan et al.⁴⁸

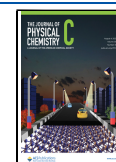
observed an increase in PL intensity and lifetime, and attributed this to self-passivation of halide vacancies. Zhang et al.⁴⁹ showed an increase in PL QY over the course of 5–10 min from CsPbBr₃ nanocrystals exposed to UV-NIR femto-second laser pulses. There have been other reports of increased PL QY and lifetime from CsPbBr₃ nanocrystals following UV exposure, but all in conjunction with additional factors like annealing⁵⁰ or photoinduced morphology changes.⁵¹ Seth et al.⁵² reported significant increases in PL intensity over the course of several hours for CsPbBr₃ and CsPbBr₃I nanocrystals, but with only minor increases in PL lifetime.

Here we show that PL lifetime and intensity of perovskite-phase CsPbI₃ nanocrystal films can slowly increase after an initial photoexcitation for approximately 20–30 min before eventually stabilizing. The PL quantum yield and lifetime can increase during this time by as much as a factor of 5. The increase in PL intensity and lifetime is reversible, and several hours after removing the light source, the initially short lifetime and lower PL intensity return. The nanocrystal films that exhibit this phenomenon are stored in air. By placing the CsPbI₃ nanocrystal films under vacuum or nitrogen in the dark, we found that a similar increase in PL lifetime and intensity was observed. These results point to the presence of slow

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surface states in perovskite-phase CsPbI_3 nanocrystals mediated by humidity and light.

EXPERIMENTAL METHODS

Materials. Lead iodide (PbI_2 , 99.999% (metals basis), Sigma-Aldrich), cesium carbonate (Cs_2CO_3 , 99.9% (metals basis), Sigma-Aldrich), oleylamine ($\text{C}_{18}\text{H}_{35}\text{NH}_2$, OAm, technical grade, 70%, Sigma-Aldrich), oleic acid ($\text{C}_{18}\text{H}_{34}\text{O}_2$, OA, technical grade, 90%, Sigma-Aldrich), 1-octadecene (ODE, $\text{C}_{18}\text{H}_{36}$, technical grade, 90%, Sigma-Aldrich), methyl acetate ($\text{C}_3\text{H}_6\text{O}_2$, MeOAc, 99.5%, anhydrous, Sigma-Aldrich), hexane (>99%, anhydrous, Sigma-Aldrich), and toluene (>99%, anhydrous, Sigma-Aldrich) were used as received without further purification.

CsPbI_3 Nanocrystal Synthesis. Perovskite-phase CsPbI_3 nanocrystals were synthesized by using a procedure similar to Protesescu et al.^{3,53,54} In a three-neck round-bottom flask, 407 mg (1.25 mmol) of Cs_2CO_3 and 1.25 mL of oleic acid were added to 20 mL of octadecene. On a Schlenk line, the mixture was stirred under vacuum for 20 min at room temperature. The temperature was then raised to 100 °C and stirred for 20 min. The flask was purged with nitrogen, and the temperature was raised to 150 °C until the mixture became clear. The temperature was reduced to 130 °C and left until ready to use later in the synthesis.

In a separate flask, 300 mg (0.651 mmol) of PbI_2 was added to 10 mL of octadecene and stirred under vacuum at 110 °C for 20 min. 1.5 mL each of oleylamine and oleic acid were combined, heated to 130 °C, and injected into the flask. The flask was kept stirring under vacuum until vapor was no longer observed to evolve from the mixture. The flask was placed under nitrogen, and the temperature was quickly raised to 165 °C. Once 165 °C was reached, 0.8 mL of the Cs–oleate mixture above was injected, followed by immediate quenching of the reaction with an ice bath.

Nanocrystals were purified by using hexane/MeOAc solvent/antisolvent precipitation, similar to Swarnkar et al.^{32,53,54} MeOAc (25 mL) was added to the crude reaction product (a 2:1 MeOAc to reaction product ratio by volume), and the mixture was centrifuged at 7500 rpm (7232g rcf) for 5 min. The supernatant was discarded, and the precipitated nanocrystals were redispersed in 2 mL of hexane. After the addition of 2 mL of MeOAc, the mixture was immediately centrifuged at 7500 rpm for 5 min. The supernatant was discarded, and the precipitated nanocrystals were redispersed in 4 mL of hexane before finally centrifuging the dispersion at 5000 rpm (3214g rcf) for 2 min to remove insoluble reaction byproducts and poorly capped nanocrystals. The nanocrystals were stored as dispersions in sealed vials in the dark with refrigeration. A typical reaction yields ~40 mg of nanocrystals.

Materials Characterization. For optical characterization, CsPbI_3 nanocrystals were dropcast or spin-coated from hexane or toluene dispersions onto soda lime glass microscope slides (Cardinal Health) to form films with thickness between 100 and 800 nm. The film thickness was measured by using a Keyence VK-X1100 optical profilometer. PL emission spectra were measured by using a Fluorolog-3 FL3–221 spectrofluorometer (Horiba Jobin Yvon) with 450 nm excitation wavelength and slit size of 5 nm for both excitation and emission. Time-resolved PL (trPL) data were measured on the Fluorolog-3 instrument with a FluoroHub-B time-correlated single photon counting (TCSPC) accessory. For trPL, photoexcitation was provided by a 442 nm wavelength

NanoLED pulsed excitation source (pulse duration = 1.2 ns, repetition rate = 0.5–1 MHz). PL measurements under vacuum were taken by using a cryostat chamber (Janis Research Co.) and an Edwards 1.5 vacuum pump (1.1×10^{-3} Torr rated vacuum). UV–vis absorbance spectra were taken by using a Cary 5000 UV–vis–NIR spectrophotometer with an Agilent Technologies diffuse reflectance accessory. Absorbance, A , is the ratio of absorbed radiant energy to incident radiant energy, or $A = 1 - \text{transmittance} - \text{reflectance}$.

X-ray diffraction (XRD) was measured by using a Rigaku R-Axis Spider X-ray diffractometer with an image plate detector and $\text{Cu K}\alpha$ radiation ($\lambda = 1.54$ nm). XRD data were collected for 10 min with 2 deg/s sample rotation, and the X-ray generator operated at 40 kV and 40 mA.

Transmission electron microscopy (TEM) images were acquired with an FEI Tecnai TEM at 80 kV accelerating voltage. TEM samples were made by dropcasting nanocrystals from a dilute dispersion in hexane onto a continuous carbon-coated copper mesh TEM grid (Electron Microscopy Sciences). High-angle annular dark-field scanning transmission electron microscopy (HAADF-STEM) was performed by using an FEI Titan G2 ChemiSTEM operated at 80 kV with a Cs probe corrector. For HAADF-STEM, the TEM grid was dipped in methyl acetate to remove excess ligand and carbon contamination. Prior to imaging, the TEM grids were loaded into the TEM and left undisturbed in a vacuum for several hours.

RESULTS

Figures 1a and 1b show transmission electron microscopy (TEM) and scanning transmission electron microscopy high angle annular dark field (STEM-HAADF) images of typical perovskite-phase CsPbI_3 nanocrystals used in the study. The nanocrystals have a cuboidal shape with an average edge length of 8.9 ± 2.0 nm. (Additional TEM images and a

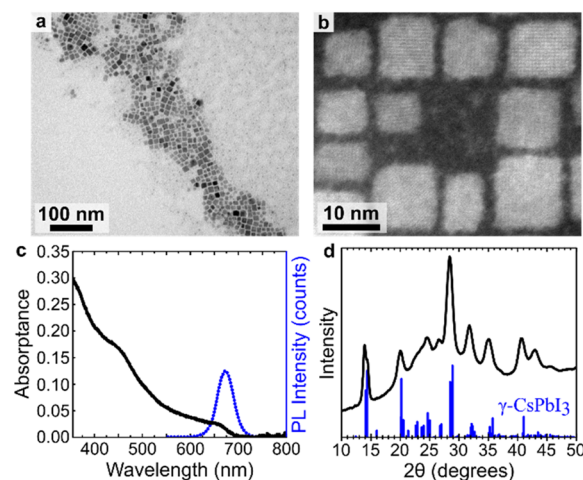


Figure 1. (a) TEM and (b) high-angle annular dark-field (HAADF) STEM images of the CsPbI_3 nanocrystals used in this study. (c) Room temperature absorbance and PL spectra of a CsPbI_3 nanocrystal film on glass ($\lambda_{\text{exc}} = 450$ nm). The absorbance was measured by using a diffuse reflectance integrating sphere to properly account for light scattering and reflection. (d) XRD of CsPbI_3 nanocrystals (black trace) compared to a calculated pattern (blue drop lines) for γ -orthorhombic CsPbI_3 with space group $Pbnm$ (no. 62).^{32,54,58}

histogram of particle sizes are provided as [Supporting Information](#).) Figure 1d shows X-ray diffraction (XRD): the nanocrystals inhabit the perovskite, γ -orthorhombic phase of CsPbI_3 .^{55–59} Room temperature UV–vis absorbance and PL emission spectra from CsPbI_3 nanocrystals deposited on glass are shown in Figure 1c. The band-edge PL peak at 672 nm is typical for γ - CsPbI_3 nanocrystals.^{3,4} PL quantum yields (QYs) of the CsPbI_3 nanocrystals dispersed in hexane range between 25 and 35%.⁵⁰

Figure 2 shows room temperature PL and trPL of a film of CsPbI_3 nanocrystals that had been stored in air. The PL

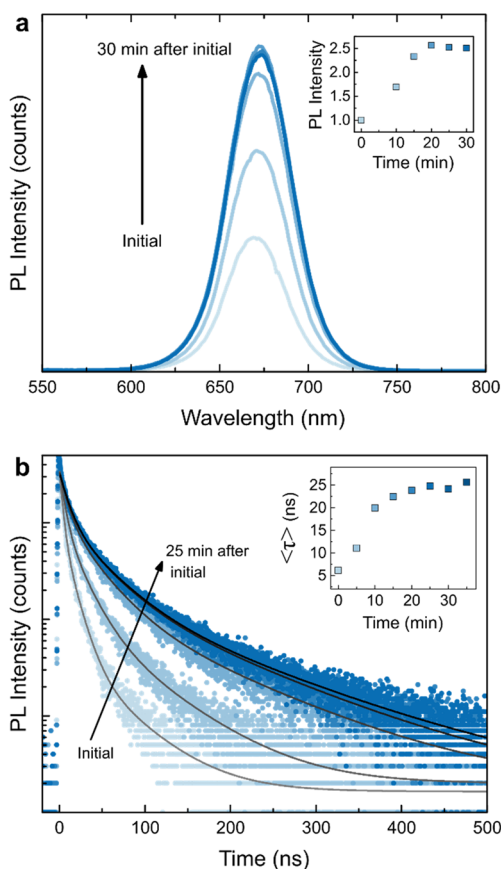


Figure 2. (a) Steady-state PL and (b) trPL of a CsPbI_3 nanocrystal film on glass in air with RH = 45–55% showing the increase in PL intensity and lifetime over time following the initial measurement. (a) PL measurements taken every 5 min over a 30 min period, with darker lines representing later measurements. The excitation wavelength (λ_{exc}) was 450 nm. Inset: time evolution of the integrated PL peak intensity normalized to its initial state after the PL is first measured. (b) trPL measurements taken every 5 min, with darker-colored data points representing later measurements. Each data set was fit with a triexponential decay (gray/black lines, see eq 1). The excitation wavelength (λ_{exc}) was 442 nm. Inset: amplitude average lifetime $\langle\tau\rangle$ of each decay curve vs time after the initial measurement, where $\langle\tau\rangle = \sum a_i\tau_i$. For both sets of measurements, the excitation light was turned on only during measurement.

intensity and lifetime changes significantly for about 20 min before reaching a steady-state response. After leveling off, the PL intensity has increased by nearly a factor of 3, with a slight red-shift of 3 nm and increase in full width at half-maximum (FWHM) of about 3 nm. The PL lifetime can be quantified by fitting the trPL data with a triexponential decay:

$$I(t) = a_1 \exp\left(\frac{-t}{\tau_1}\right) + a_2 \exp\left(\frac{-t}{\tau_2}\right) + a_3 \exp\left(\frac{-t}{\tau_3}\right) \quad (1)$$

Similar to other reports of CsPbX_3 nanocrystals and films,^{50–52,60–62} single- or biexponential decay functions do not accurately fit the data, except when the lifetimes were short—in our case, only during the initial trPL measurement. Once the lifetime increases, the triexponential fit is needed. The amplitude average PL lifetime, $\langle\tau\rangle = \sum a_i\tau_i$,⁶³ of the sample in Figure 2 (shown in the inset of Figure 2b) increased from 6.1 to 24.8 ns. Table 1 provides the fitting parameters used to obtain $\langle\tau\rangle$. The values of $\langle\tau\rangle$ lie within the range of values reported in the literature for CsPbI_3 nanocrystals.^{3,32,60,64}

Table 1. Parameters Used to Fit Eq 1 to the trPL Data in Figure 2b and the Corresponding Amplitude average Lifetime $\langle\tau\rangle$ (Lifetime Values in ns)

	time (min) after light exposure						
	0	5	10	15	20	25	30
τ_1	3.39	5.41	8.15	8.82	9.43	9.77	8.47
a_1	0.74	0.70	0.64	0.65	0.65	0.65	0.62
τ_2	11.9	19.2	31.3	35.2	35.7	38.1	34.4
a_2	0.25	0.27	0.31	0.30	0.29	0.29	0.31
τ_3	49.7	64.2	109	121	118	129	120
a_3	0.01	0.03	0.05	0.05	0.06	0.06	0.07
$\langle\tau\rangle$	6.09	11.05	19.96	22.44	23.83	24.76	24.15

We found that PL intensity and lifetime continued to gradually increase after the initial light exposure during a PL measurement, even without continuous illumination, as shown in Figure 2. The data in Figure 2 were acquired by measuring the PL every 5 min, without any light exposure between measurements. Regardless of when the subsequent measurements were taken, the initial photoexcitation kicks off a gradual, slow change in PL kinetics. This change in PL intensity and lifetime is reversible, however, and when the sample is stored in the dark for several hours, the initial (short) PL lifetime and intensity return. Typically, this would take about 12 h when the nanocrystals are placed in ambient air at about 45–55% relative humidity (RH). When the PL spectra and lifetime are measured again, the photoexcitation during the initial measurement reinitiates a rise in PL intensity and lifetime. XRD (Figure 3) showed that the increased PL intensity and lifetime are not related to any change in crystal structure.

When γ - CsPbI_3 nanocrystals are dispersed in a solvent, they do not show any light-induced changes in PL intensity and lifetime (see Figure S2 in the [Supporting Information](#)). The PL lifetime of CsPbI_3 nanocrystals dispersed in hexane was 13 ns, which is longer than the initial PL lifetime of the films (6 ns) and shorter than the steady-state lifetime observed after photoexcitation (24 ns). One other significant difference between nanocrystals dispersed in a solvent and those dried into films is that the PL rapidly deteriorated in the dispersions, with decreasing lifetime and intensity, most likely due to ligand desorption and degradation as the nanocrystals reverted to their equilibrium yellow, non-emitting orthorhombic δ -phase. As we have previously shown, ligand-stabilized γ - CsPbI_3 nanocrystals are quite stable in air, even under humid conditions.⁵⁴

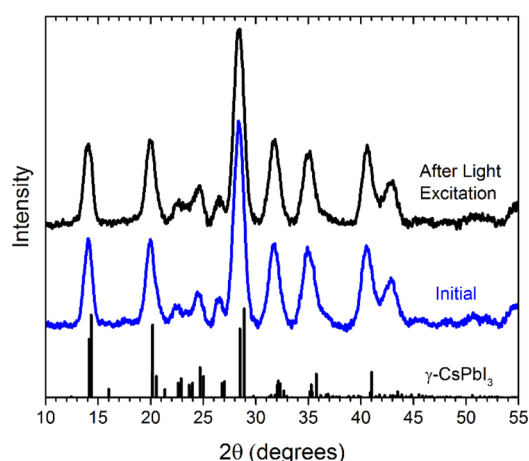


Figure 3. X-ray diffraction from a CsPbI_3 nanocrystal film before (black line) and after (blue line) an induced increase in PL intensity and lifetime by light excitation in air. The black drop lines correspond to calculated diffraction peaks for γ -orthorhombic CsPbI_3 with space group $Pbnm$ (no. 62).⁵⁸

To better understand the reversible light-induced enhancement of PL intensity and lifetime of γ - CsPbI_3 nanocrystals, we measured the PL under a variety of other environmental conditions. The PL intensity and lifetime also increased when the γ - CsPbI_3 nanocrystal film was placed under vacuum in the dark for several minutes. The black trace in Figure 4 shows

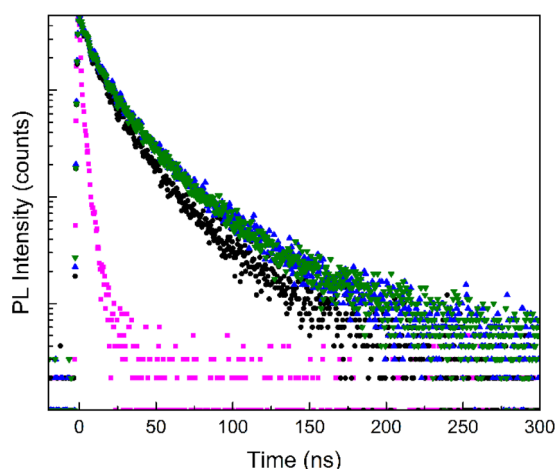


Figure 4. Time-resolved photoluminescence (trPL) of a γ - CsPbI_3 nanocrystal film under vacuum (black, blue, and green data points) compared to a γ - CsPbI_3 nanocrystal film in ambient air (magenta squares). For the film under vacuum, the initial measurement (black circles) was taken after leaving the sample under vacuum for 10 min. The traces shown in blue and green (up and down triangles, respectively) were taken 5 and 10 min after the initial measurement, respectively. The sample is illuminated only during the measurement by the excitation light source; otherwise, the sample is in the dark. The excitation wavelength (λ_{exc}) is 442 nm.

trPL data for a γ - CsPbI_3 nanocrystal film after being held under vacuum for 10 min. The magenta trace in Figure 4 shows the trPL data from a film without this vacuum treatment. The initial PL lifetime is significantly longer for the film placed under vacuum. The green trace in Figure 4 shows the trPL data measured 10 min after the first PL measurement for the film that had been placed under vacuum. There is still an increase

in PL lifetime, but it is less pronounced and reaches steady state after only 5 min. After placing the film back into ambient air, the PL intensity and lifetime returned to their low values after about 12 h. Figure 5 shows a plot of $\langle\tau\rangle$ for a γ - CsPbI_3 nanocrystal film as it changes under vacuum, in air, and after photoexcitation. Both vacuum- and light-induced changes are reversible.

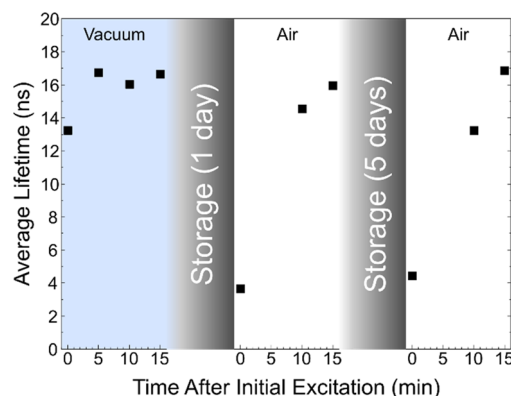


Figure 5. Change in average PL lifetime $\langle\tau\rangle$ of a γ - CsPbI_3 nanocrystal film after light excitation followed by storage in humid air (45–55% RH) at room temperature. The data points labeled “Vacuum” were obtained after the film had been placed under vacuum for 10 min and then measured under vacuum. The trPL data labeled “Air” were taken in ambient air. The sample is illuminated only during the measurement by the excitation light source; otherwise, the sample is in the dark. The excitation wavelength (λ_{exc}) is 442 nm.

We also found that γ - CsPbI_3 nanocrystal films placed under nitrogen before measuring the PL had a higher PL intensity and longer lifetime. After being stored for 1 h under nitrogen in the dark, the PL lifetime of the film was relatively long (15 ns) (see Figure S3). These films again revert back to lower PL intensity and shorter lifetimes when returned to ambient air (45–55% RH) for more than 12 h. The effects of the vacuum and nitrogen environments indicate that moisture plays a role in these reversible changes in PL intensity and lifetime.

Figure 6 shows the PL lifetime of films stored in air at 15% RH and 55% RH after the PL lifetime had been enhanced by light excitation. After 24 h, the PL lifetime of the film stored in 55% RH had dropped back to its initially short value, while the lifetime of the film stored in 15% RH remained slightly longer. Films took longer to recover their lower lifetime and intensity when stored in air with lower humidity. When the film was stored in nitrogen, it exhibited no decrease in PL lifetime or intensity after 24 h.

CsPbI_3 nanocrystal films were completely unstable under conditions of very high humidity (>90% RH) and turned yellow after about 12 h. Humidity is well-known for its role in destabilizing the black perovskite phase of CsPbI_3 by catalyzing the phase change to the equilibrium yellow, orthorhombic δ -phase.^{54,58,59,65} Figure 7 shows XRD from a CsPbI_3 nanocrystal film that had been stored in air for 12 h at >90% RH and transformed to the equilibrium, non-perovskite δ -orthorhombic crystal phase. It should be noted that these transformed films look completely yellow to the eye but can still exhibit measurable PL from the residual γ -orthorhombic phase in the film. The PL lifetimes of these yellow films vary widely and depend on how much γ - CsPbI_3 is still present.

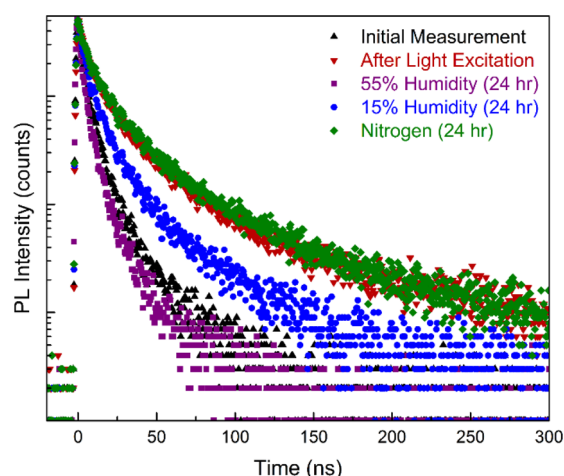


Figure 6. trPL of γ -CsPbI₃ nanocrystal films that had been photoexcited to induce an increase in PL lifetime and then stored under various conditions. Black triangles show the initial trPL of a film that had been stored in air, and red triangles show the trPL after light excitation, which exhibits the longest PL lifetime. The purple squares and blue circles show trPL data for films that were stored at 55 and 15% relative humidity, respectively, for 24 h. When the film was stored under nitrogen (green diamonds), there was no decrease in PL lifetime after 24 h. These measurements were taken in ambient air at 45–55% relative humidity with an excitation wavelength (λ_{exc}) of 442 nm.

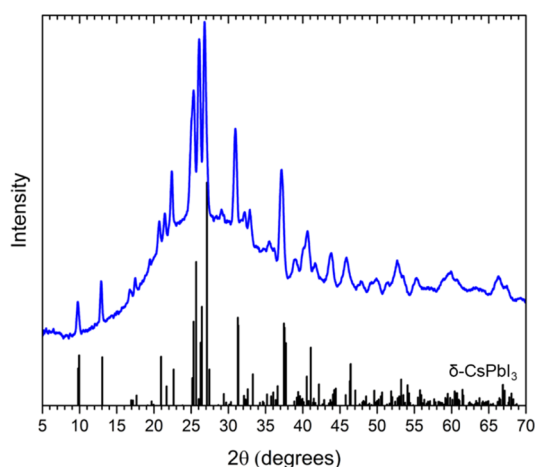


Figure 7. X-ray diffraction from a CsPbI₃ nanocrystal film on glass after 24 h in air at 90% RH (blue line). The nanocrystals have degraded to the yellow δ -orthorhombic phase from the initial black γ -orthorhombic phase. The reference of black drop lines are calculated diffraction peaks for δ -orthorhombic phase CsPbI₃ (PDF #01-074-1970).

DISCUSSION

Light-induced increases in PL intensity and lifetime have also been observed in the hybrid organic–inorganic perovskites, methylammonium lead iodide (CH₃NH₃PbI₃, MAPbI₃),^{66–71} MAPbBr₃,^{72,73} and MAPb(I_{0.5}Br_{0.5})₃.⁷⁴ Brenes et al.⁶⁹ proposed that photoinduced passivation of traps related to halide vacancies accounted for their observation of a gradual increase in PL from MAPbI films under continuous illumination. In MAPbI, vacancies can be filled by superoxide ions (O₂^{•−}) formed by the interaction of oxygen and photoexcited electrons.^{69,75–78} The light-induced changes that we observe in γ -CsPbI₃ nanocrystal films may also involve

a similar oxygen passivation process; however, a mechanism that only considers oxygen cannot explain the increase in PL lifetime observed when films were placed under vacuum or nitrogen or the significant influence of humidity on the rate of the PL intensity and lifetime changes. In the case of γ -CsPbI₃ nanocrystal films, water must also be involved.

The PL QY and PL lifetime, τ_{PL} , are related to the radiative lifetime, τ_r , and the radiative and nonradiative electron–hole recombination rates, k_r and k_{nr} : $\text{QY} = k_r / (k_r + k_{\text{nr}}) = \tau_{\text{PL}} / \tau_r$.^{79,80} The absorbance profile of the nanocrystal films did not change after illumination (see Figure S4), indicating that the amount of light absorbed by the nanocrystals does not change. Because similar enhancements are observed for both the QY and τ_{PL} , it is the change in nonradiative recombination that is responsible for the increased PL intensity and lifetime of the γ -CsPbI₃ nanocrystal films. We propose that the association of adsorbed water with the CsPbI₃ surface create slow surface states, or traps, that induce nonradiative recombination in the films.

The hydrophobic capping ligands help stabilize the γ -phase by preventing water vapor from reaching the crystal surface.⁵⁴ However, the ligands cannot completely passivate all of the dangling bonds presented by atoms at the nanocrystal surface because of the steric hindrance between neighboring alkyl chains.⁸¹ It is known that water can interact with surface defects on the surface of lead halide perovskites and induce halide vacancies.^{82–84} Liu et al.⁸⁵ showed that oxygen can passivate halide vacancies in CsPbI₂Br and speculated that a decrease in exposed halide vacancies limits the adsorption of water. It is possible that both halide vacancies and water adsorption are involved in the formation of slow surface trap states on γ -CsPbI₃ nanocrystals. Filling of halide vacancies by oxygen and photoexcited electrons could then lead to the PL increase under light excitation that we observed.

According to our model, the rate of nonradiative electron–hole recombination can be reduced by either (1) the removal of adsorbed water and elimination of the trap states under vacuum or nitrogen or (2) the elimination of the surface trap states by a light-induced process involving photoexcited charge carriers, oxygen, and/or adsorbed water. To calculate the characteristic time or rate of deactivation of surface trap states following light excitation in air, we considered these distinct surface traps as dynamic PL quenchers and used kinetic equations to correlate the concentration of active surface traps to the measured PL lifetimes. In the presence of the quencher or surface trap, T, caused by the interaction of water with the perovskite surface, excited electrons can decay by the following steps illustrated in Figure 8:



where S* and S are the excited and ground states, respectively, and k_T is the rate of nonradiative electron–hole recombination associated with the traps. Once the sample has been photoexcited

$$-\frac{d[S^*]}{dt} = k_r[S^*] + k_{\text{nr}}[S^*] + k_T[T][S^*] \quad (4)$$

Solving this differential equation gives

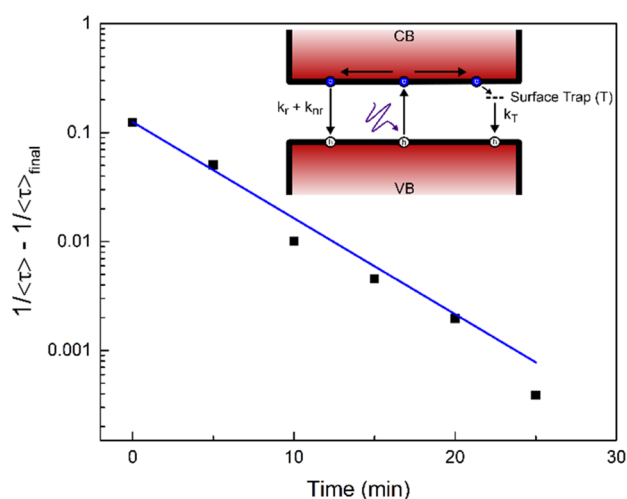


Figure 8. Plot of $1/\langle\tau\rangle - 1/\langle\tau\rangle_{\text{final}}$ vs time after initial photoexcitation for the data shown in Figure 2. $\langle\tau\rangle_{\text{final}}$, the steady-state (maximum) average lifetime of this sample, was taken to be 25 ns based on the values in Table 1. The best fit of eq 10 to the data (blue line) gives a value of the characteristic time for deactivation or elimination of slow surface trap states of $\tau_d = 4.9$ min.

$$[S^*] = Ae^{-(k_r + k_{nr} + k_T[T])t} \quad (5)$$

to provide an expression for the PL lifetime in the presence of surface traps with a concentration of $[T]$:

$$\tau = \frac{1}{k_r + k_{nr} + k_T[T]} \quad (6)$$

In the absence of the surface traps, the PL lifetime (τ_0) is given by

$$\tau_0 = \frac{1}{k_r + k_{nr}} \quad (7)$$

From the ratio of the lifetimes in the absence and presence of the surface traps, a form of the Stern–Volmer equation is obtained:⁸⁰

$$\frac{\tau_0}{\tau} = \frac{k_r + k_{nr} + k_T[T]}{k_r + k_{nr}} = 1 + \tau_0 k_T [T] \quad (8)$$

Rearranging this equation shows that $[T]$ is related to the measured PL lifetimes according to

$$[T] = \frac{1}{k_T} \left(\frac{1}{\tau} - \frac{1}{\tau_0} \right) \quad (9)$$

The concentration of surface trap states $[T]$ is proportional to $1/\tau - 1/\tau_0$. The PL lifetime in the absence of the trap states, τ_0 , can be taken as the steady-state τ_{PL} , or τ_{final} , reached by the sample after photoexcitation. A first-order kinetic equation for the surface trap deactivation or elimination process, with trap state concentration $[T]$ substituted by $1/\tau - 1/\tau_{\text{final}}$, yields

$$\frac{1}{\tau} - \frac{1}{\tau_{\text{final}}} = Ae^{-t/\tau_d} \quad (10)$$

where τ is taken as the amplitude average PL lifetime, $\langle\tau\rangle$, measured at time t after the initial excitation and τ_d is the characteristic time (inverse rate) of the deactivation or elimination of surface trap states caused by the interaction of water with the perovskite surface. Figure 8 shows the PL

lifetime data from Figure 2 replotted to determine the characteristic time for deactivation or elimination of surface trap states, τ_d , which for this sample is 4.9 min. The photoinduced elimination of these trap states is a slow, reversible process. Although slow surface states on metal halide perovskites have not been previously identified, slow surface states are well-known for other semiconductors, such as Ge.⁸⁶ When the nanocrystal film is returned to humid air, the slow surface traps are reactivated by adsorbed water and the PL lifetime reverts back to the initially short values.

CONCLUSIONS

Here, we reveal the presence of humidity-sensitive slow surface states in γ -CsPbI₃ nanocrystals. The PL intensity and lifetime of γ -CsPbI₃ nanocrystals in films can change significantly after light excitation in a PL measurement, increasing by as much as a factor of 5 over the course of 20 min. These changes are reversible in a humid environment. Slow surface states are well known to occur in other semiconductors and are typically associated with reversibly bonded chemical species on the semiconductor surface.⁸⁶ We propose that the slow surface states of the γ -CsPbI₃ nanocrystals affect only the nonradiative electron–hole recombination rates and are photosensitive. Passivation of the surface trap states can occur either through the removal of water interacting with the nanocrystal surface or by a slow light-activated process.

The sensitivity of optical properties to the combination of light and environmental moisture might be related to the fact that the perovskite phase of CsPbI₃ nanocrystals at room temperature is a metastable phase which is quite sensitive to moisture.^{54,58,59} Further study is needed to delineate additional details concerning the reversible enhancement in PL intensity and lifetime of CsPbI₃ nanocrystal films. Light-intensity-dependent studies, for example, might provide more information about the kinetics of the surface interaction process and the role of trap filling. Our results highlight the complex optoelectronic behavior of CsPbI₃ nanocrystals and provide a better understanding of these materials with which to guide strategies for improving the performance and stability of inorganic perovskite devices.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jpcc.2c04305>.

Histogram of nanocrystal size distribution; trPL of nanocrystals dispersed in hexane; trPL of nanocrystal films with and without nitrogen exposure; absorbance spectra of CsPbI₃ nanocrystal films before and after the photoinduced PL increase (PDF)

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Notes

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