

High pressure Cr³⁺ R-line luminescence of zoisite and kyanite: A probe of octahedral site distortion

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Abstract

The Cr³⁺ luminescence spectra of zoisite and kyanite, two geologically important minerals, were studied up to 40 and 20 GPa, respectively, in various pressure media. Cr³⁺ substitutes into the octahedral aluminum sites in both minerals and the R-line luminescence is a particularly sensitive site-specific probe of the octahedral Al site. Unlike many previous studies where Cr³⁺ luminescence was utilized, both these minerals have multiple highly distorted octahedral sites resulting in very large splitting of their R-lines, ~300 cm⁻¹ in zoisite and ~360 cm⁻¹ in kyanite (for reference, ruby is 29 cm⁻¹). For zoisite, the R-line splitting increases as pressure increases and more than triples from its ambient value by 40 GPa, while the R-line splitting in kyanite from the M1 and M2 sites does not change when compressed in a Ne pressure medium up to 20 GPa. We do not observe evidence of any phase transitions in either zoisite or kyanite across the pressure range of these new luminescence measurements. We present some high-pressure luminescence results where kyanite was known to be bridged between the diamond anvils and show how these spectra illustrate the different effect of uniaxial relative to hydrostatic stress on luminescence spectra.

Key words: zoisite, kyanite, high-pressure, luminescence, diamond anvil cell

Introduction

Zoisite [Ca₂Al₃(SiO₄)(Si₂O₇)O(OH)] contains ~2.0 wt.% and is an important source and carrier of water in subduction zones (Dobrzhinetskaya and Green, 2007; Nicholls and Ringwood, 1973; Ono, 1998; Peacock, 1990). For comparison, kyanite [Al₂SiO₅] is a common accessory mineral in eclogites, and can be found both in subduction zones where certain crustal rocks are subducted, and as a breakdown product of aluminous phases (Ono, 1998). The stability field of zoisite in a basalt + H₂O system extends to ~3.1 GPa

at 650 °C (Forneris and Holloway, 2003; Forneris and Holloway, 2004). It has a slightly larger stability field in the CaO-Al₂O₃-SiO₂-H₂O (CASH) system, which ranges from 4.0 GPa at 500 °C to as high as 7.0 GPa at over 1000 °C (Poli and Schmidt, 1998; Schmidt and Poli, 1994); beyond these pressures it breaks down to lawsonite-bearing assemblages. Zoisite has been studied with X-ray and spectroscopic techniques up to 34 GPa at 300 K, and no phase transitions have been reported (Alvaro et al., 2012; Comodi and Zanazzi, 1997; Winkler et al., 1989; Xu et al., 2019), while kyanite has been found to decompose into stishovite and corundum between 5-15 GPa and 1200-1900 K (Ono et al., 2007) and between 14-17.5 GPa and 1272-2273 K (Schmidt et al., 1997). At room temperature, a number of studies of kyanite up to 20 GPa using powder and single crystal X-ray diffraction as well as Raman spectroscopy did not find evidence for any phase transitions (Comodi et al., 1997; Friedrich et al., 2004; Liu et al., 2009; Mernagh and Liu, 1991; Yagi et al., 1998; Yang et al., 1997). However, at pressures >30 GPa and temperatures > 2000 K, other phases of kyanite have been reported to be stable (Ahmed-Zaid and Madon, 1991; Friedrich et al., 2004). New high-pressure phases kyanite-II and kyanite-III have been reported from multi-anvil studies at 13 and 17 GPa, but at temperatures exceeding 2000 K (Zhou et al., 2021; Zhou et al., 2018). Recently, (Gao et al., 2020) compressed kyanite up to 18.6 GPa at room temperature using both Raman and infrared spectroscopy and report that a first order phase transition to kyanite-II (which is likely distinct from the high-pressure/high-temperature kyanite-II polymorph) occurs at 9.7 GPa. These results distinctly differ from previous room temperature, high pressure X-ray and spectroscopic studies on kyanite (Comodi et al., 1997; Friedrich et al., 2004; Liu et al., 2009; Mernagh and Liu, 1991; Yang et al., 1997). Thus, part of the motivation for examining kyanite using luminescence is to determine whether a different probe can resolve any possible phase transitions that may occur under pressure at 300 K in this material.

Cr³⁺ luminescence spectra in the red typically consist of intense relatively sharp emission bands, and in some cases a broad emission band. The sharp emission bands are commonly referred to as R-lines. These emission bands are associated with the spin-forbidden ²E-⁴A₂ transition, with the ⁴A₂ state being

the ground state and the 2E state being split in non-cubic environments (Syassen, 2008). The broad emission band is assigned to the spin-allowed ${}^4T_2-{}^4A_2$ transition. At high crystal field strengths (Burns, 1993; Tanabe and Sugano, 1954), the narrow band ${}^2E-{}^4A_2$ emission dominates the spectrum. For low crystal field strengths (Burns, 1993; Tanabe and Sugano, 1954), the generally broader ${}^4T_2-{}^4A_2$ emission band may dominate the emission spectrum. Both sharp and broad emission bands are observed at intermediate crystal field strengths. The Cr^{3+} luminescence of ruby ($Al_2O_3:Cr$) is well characterized at ambient and high pressures and has been widely used as an *in situ* pressure calibrant for high pressure experiments (Syassen, 2008). This is due to its sharp strong R-line luminescence, lack of high-pressure phase transitions across a wide pressure and temperature range, and chemical inertness towards commonly used pressure media (Syassen, 2008).

Until recently, only a small number of other geologic materials have had their Cr^{3+} luminescence characterized under pressure. These materials were examined for many of the same reasons as ruby and include alexandrite ($BeAl_2O_4$), MgO, spinel ($MgAl_2O_4$) and a range of garnets (Chopelas and Nicol, 1982; Jähren et al., 1992; Jovanic, 2000; Kottke and Williams, 1983; Wamsley and Bray, 1994). Cr^{3+} typically substitutes into the Al^{3+} site and, since it has a slightly larger ionic radius than Al^{3+} , local site distortions can arise from this substitution. Most of the previously minerals characterized by luminescence spectroscopy under high pressures have only one or two unique Cr-bearing octahedral sites that are nearly ideal or only slightly distorted. However, there are a large number of Cr^{3+} luminescence studies on more crystallographically complex natural and synthetic materials at ambient pressure and various temperatures (Gaft et al., 2015 and references therein).

Recently, there has been a growing interest in using Cr^{3+} luminescence as a site-specific probe in geologic materials with moderate to highly distorted Cr-bearing octahedral sites, as well as lower symmetry more chemically complex minerals (Hua et al., 2021; O'Bannon et al., 2018; O'Bannon and Williams, 2016a; O'Bannon and Williams, 2016b; O'Bannon and Williams, 2017; O'Bannon and Williams,

2019). Moreover, due to the difference in ionic radius between Al and Cr, there can be extensive local distortion of the octahedral sites in which Cr³⁺ substitutes that cannot readily be characterized with techniques such as X-ray diffraction. Here, we present new high-pressure Cr³⁺ luminescence results on zoisite up to 40 GPa and kyanite up to 16 GPa. These two minerals have multiple highly distorted Al-sites, and they have R-line splittings of > 300 cm⁻¹ at ambient conditions: a separation far in excess of that of ruby (29 cm⁻¹) and other minerals examined to date by luminescence under compression. As such, the primary goals of this study include determining how luminescence from highly distorted sites is altered by compression, and the degree to which shifts in distortion can be monitored using luminescent spectra.

Methods

Trace element concentrations of the kyanite and zoisite samples were measured with a PhotonMachines Analyte 193H, which is a 193-nm ArF excimer laser system coupled with a ThermoScientific ElementXR single-collect or magnetic sector ICP-MS. The instrument was calibrated with a SRM 610 trace element glass from NIST, and Al was used as the internal standard. All reported concentrations are in ppm and are an average from 4 measurements from different locations on the sample, and the standard deviation of the 4 measurements produces the reported uncertainties. The zoisite sample is from the Mundarara Mine, Arusha Region, Tanzania. The following trace elements were measured: Cr 4860 ± 2104, V 51 ± 22, Mn 91 ± 38, and Fe 6464 ± 2763 ppm; the variability is indicative of some possible compositional zoning within the sample. The kyanite sample is a typical light blue kyanite from an unknown locality which was selected for its strong R-line luminescence and low intensity broad band emission. The following trace elements were measured: Cr 2 ± 1, V 32 ± 2, Fe 194 ± 9 ppm, Mn below detection limit. Our utilization of natural samples is unlikely to produce markedly different local deformations associated with Cr-substitution than are present in synthetic samples: the overarching shift in

lattice parameters associated with Cr-substitution is quite linear from low Cr-concentrations to nearly 60 mol% Cr₂SiO₅ substitution in Cr-kyanites (Hejny et al., 2019). While no comparable data are available for zoisite, clinozoisite appears to show similarly linear lattice parameter trends with Cr-substitution (Nagashima et al., 2009).

High static pressures were generated with either Merrill-Bassett type diamond anvil cells (DACs) (Merrill and Bassett, 1974), symmetric piston-cylinder type DACs, or LLNL type membrane DACs (Jenei et al., 2013). They were equipped with 16 sided type Ia diamond anvils with culets ranging from 500 to 200 μ m in diameter; the gasket material was either stainless steel, spring steel, or rhenium, and the pressure transmitting medium (PTM) was either methanol:ethanol (4:1), Ne, or He, depending on the particular experiment's target pressure. Single crystals of the sample were loaded in the sample chamber with a ruby sphere (Chervin et al., 2001) which was used for pressure calibration using well-known calibrations (Dewaele et al., 2008; Mao et al., 1986). The sample and the ruby were well separated so that luminescence spectra could be collected from the sample with no contamination from ruby luminescence. Ambient pressure low temperature measurements were carried out for peak assignment purposes, with the samples submerged in liquid nitrogen.

Luminescence spectra were collected either using a Horiba LabRAM HR Evolution Raman spectrometer, or a Princeton instruments SP-2500 spectrometer. For the measurements on the Horiba spectrometer, spectra were collected from 650-800 nm (15380-12500 cm^{-1}). The spectrometer focal length of this system is 800 mm. An Olympus BXFM-ILHS microscope with a 50x long working distance objective was used to focus the 532 nm laser beam onto the sample. An 1800 lines/mm grating with a corresponding spectral resolution of $\sim 1 \text{ cm}^{-1}$ (or, equivalently, $\sim 0.05 \text{ nm}$) was utilized. For measurements on the Princeton Instruments spectrometer, luminescence spectra were collected from 680-770 nm. The spectrometer focal length of this system is 500 mm. A long working distance 30X Mitutoyo objective was used to focus the 488 nm laser beam onto the sample. A 1200 lines/mm grating with a corresponding

spectral resolution of $\sim 1.14 \text{ cm}^{-1}$ was utilized. Combinations of Gaussian and Lorentzian functions were fit to the luminescence spectra with either Horiba Labspec6 software or OriginPro software.

Results and Discussion

Zoisite

Zoisite is orthorhombic, so it is not considered a member of the epidote-group (Armbruster et al., 2006; Mills et al., 2009). It is a sorosilicate with mixed SiO_4 and Si_2O_7 groups and is found in metamorphic and pegmatitic rocks. Zoisite was originally named saualpite for the Saualpe locality in Carinthia, Austria where it occurs in eclogites. The first structural model of zoisite was proposed by Ito (1950). This structural model was later confirmed by Fresenko et al. (1955), Dollase (1968), and Ghose and Tsang (1971). Zoisite has one octahedral chain parallel to [010] with two edge-sharing Al octahedra M1,2 (sometimes referred to as M1) and M3 (sometimes referred to as M2). The two metal sites have very different distortions (Figure 1), which can be quantified by calculating the quadratic elongation (QE) and angle variance (AV) described by Robinson et al. (1971). The QE of M1,2 is 1.006 and for M3 it is 1.021, while the AV is 19.4 for M1,2 and 46.4 for M3 using the single-crystal data from Alvaro et al. (2012). These chains are linked by isolated SiO_4 tetrahedra T3 and Si_2O_7 groups T1 and T2, and in between the octahedral chains and the silica tetrahedra are two crystallographically unique 7-fold sites that are occupied by Ca. The hydrogen atom is bonded to O10 of the octahedral chain and hydrogen bonded across to the O4 atom of the adjacent octahedral chain. Neutron diffraction of zoisite has confirmed the configuration of the hydroxyl unit (Camara et al., 2012; Smith et al., 1987).

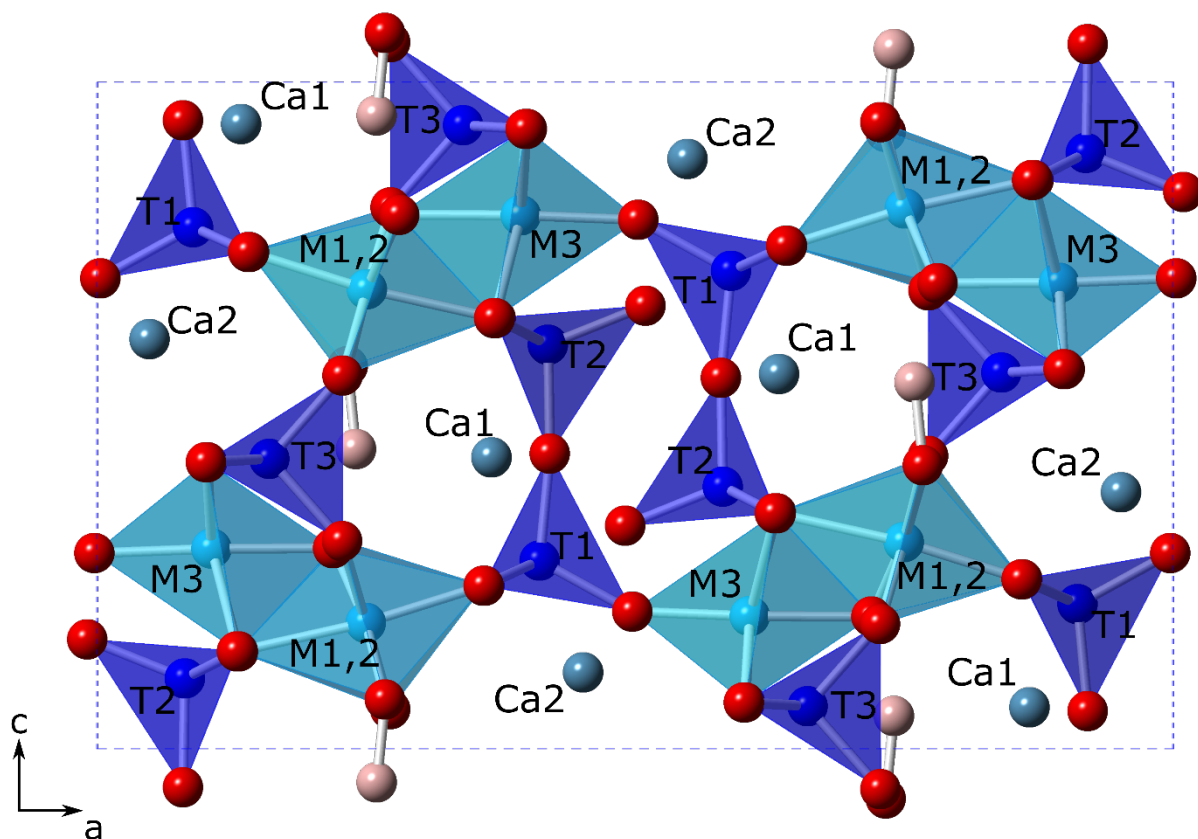


Figure 1. Crystal structure of zoisite showing the isolated SiO_4 site T3 and the Si_2O_7 sites T1 and T2, as well as the two octahedral sites M1,2 and M3, with M3 being highly distorted. Crystal structure diagram generated in CrystalMaker®.

Our 298 and 77 K spectra (Figure 2) are in good agreement with the spectra of a natural Tanzanian zoisite reported by Koziarska et al. (1994). At 77 K, the R_1 peak shifts to 708.6 nm from its 298 K location of 709.3 nm, and we observe a complete quenching of the R_2 peak that is observed at 691.9 nm in our 298 K spectrum. Additionally, several sharp bands that are likely associated with R-line emission from the less distorted M1,2 site are also observed. In particular, there is a relatively intense and sharp emission band that becomes resolvable at 77 K which could be an R-line from V^{2+} emission (Inset Figure 2). Vibronic features and/or paired emission bands (N-Lines) that contribute to the broad, distributed emission on the higher wavelength side of the

R₁ peak are also quenched at 77 K. The relatively high temperature at which we observe quenching of R₂ can be explained by the large R-line splitting in zoisite. The temperature range over which R₂ quenches scales exponentially with an activation energy that is roughly the R₁-R₂ separation, obeying the Arrhenius formula (Weinstein, 1986). Speculatively, there are two emission bands (labelled by tick marks in the inset of Figure 2) which could be the R-lines associated with the less distorted M1,2 site. If this were the case, the R-line separation for the R-lines from this site at 77 K is $\sim 115 \text{ cm}^{-1}$, which would be larger than the expected splitting at room temperature. However, the distortion of the M1,2 site in zoisite (Alvaro et al., 2012) is less than the distortion of the M1 site in spodumene (Arlt and Angel, 2000), while this possible R-line separation in zoisite is greater than observed in spodumene at 77 K, which is 90.2 cm^{-1} (O'Bannon and Williams, 2016b). Thus, the observed separation of these two emission bands is not fully consistent with expectations given the distortion of the site. Nevertheless, the volume of the M1,2 site is compatible with the observed emission peak positions. Therefore, luminescence lifetime measurements would likely be needed to unambiguously distinguish whether these emission bands are Cr³⁺ R-lines. The possible emission band from V²⁺ near 714 nm that is observed at 77 K is difficult to consistently resolve across the pressure range of these measurements, but it appears to shift like the Cr³⁺ R₁ emission band, and its relative intensity also increases under compression.

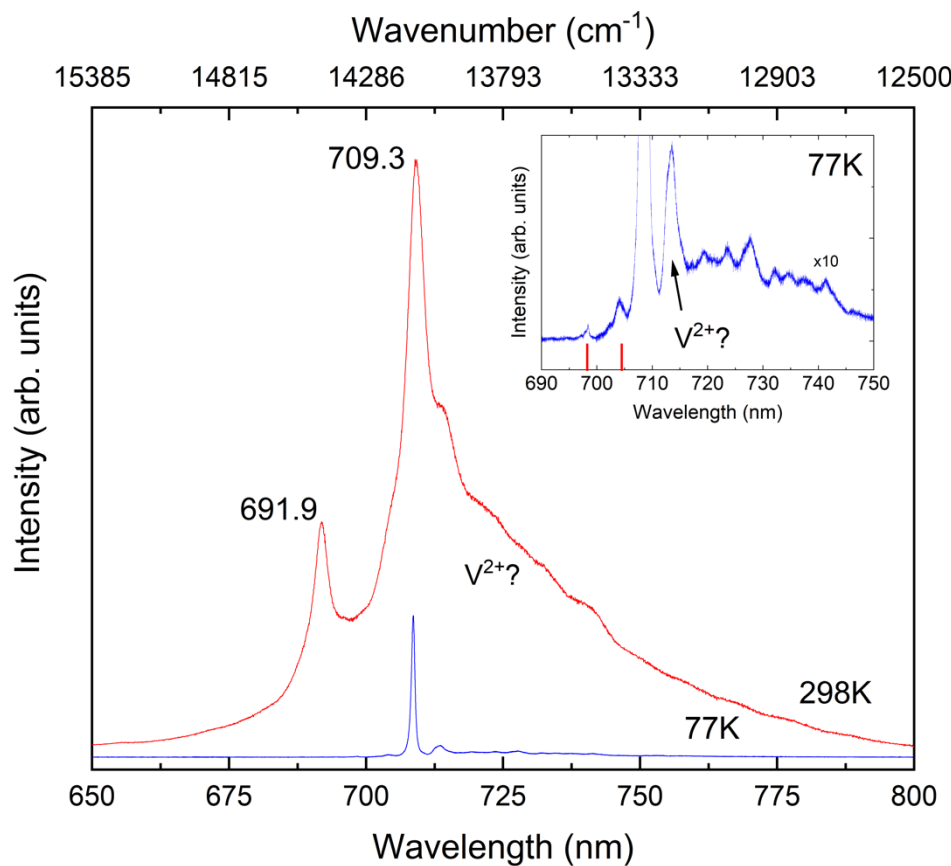


Figure 2. Representative Cr³⁺ luminescence spectra of zoisite at 298 K and 77 K. The inset shows a zoom-in of the 77 K spectrum ($\sim \times 10$), and the vertical lines in the inset mark the possible positions of R-lines originating from the M1,2 site. Possible emission from V²⁺ is also labelled in the inset.

Representative high-pressure Cr³⁺ luminescence spectra are shown in Figure 3. Three runs were conducted in three different pressure media (4:1 methanol:ethanol, Ne, and He) to a maximum pressure of 40 GPa. We fit one peak under each R-line, while the broad emission band is more difficult to deconvolve since it may be partly due to Cr³⁺, V²⁺ and/or Fe (Koziarska et al., 1994). Since the splitting of the ²E state is produced by deviations from ideal octahedral symmetry, R-line separation is regarded as an indicator of the degree of octahedral distortion. Ruby has an Al-site that is close to an ideal octahedron and has an R-line splitting of $\sim 29 \text{ cm}^{-1}$ (Syassen, 2008). The octahedral sites M1,2 and M3 in zoisite are, in comparison, highly distorted

(Figure 1) and zoisite has an R-line separation of $\sim 340\text{ cm}^{-1}$ (Kozierska et al., 1994). M3 is the more distorted octahedral site in zoisite, so given both the size of the observed splitting and its larger size, it is probable that Cr^{3+} preferentially substitutes into the M3 site over the M1,2 site. This perhaps simplified idea of larger transition-metal cations substituting into the larger, more distorted site is supported by single-crystal X-ray diffraction (Dollase, 1968; Dollase, 1969; Dollase, 1971; Grapes, 1981), Mössbauer spectroscopy (Dollase, 1973), and previous ambient-pressure luminescence spectra (Kozierska et al., 1994) of transition-metal bearing zoisite and zoisite-related phases. The diffraction and Mössbauer studies on a suite of closely related minerals (allanite, epidote, hancockite, piemontite) show that larger cations preferentially substitute into the larger more distorted M3 site, although some substitution into the M1,2 site is also observed. The room temperature luminescence spectrum has strong broad band luminescence (Figure 2), which would likely obscure any low intensity sharp line emission from minor levels of substitution in the less distorted M1,2 site. Yet, it is highly likely that the majority of the Cr^{3+} that has substituted into the zoisite structure (and from which we observe emission) is in the M3 site.

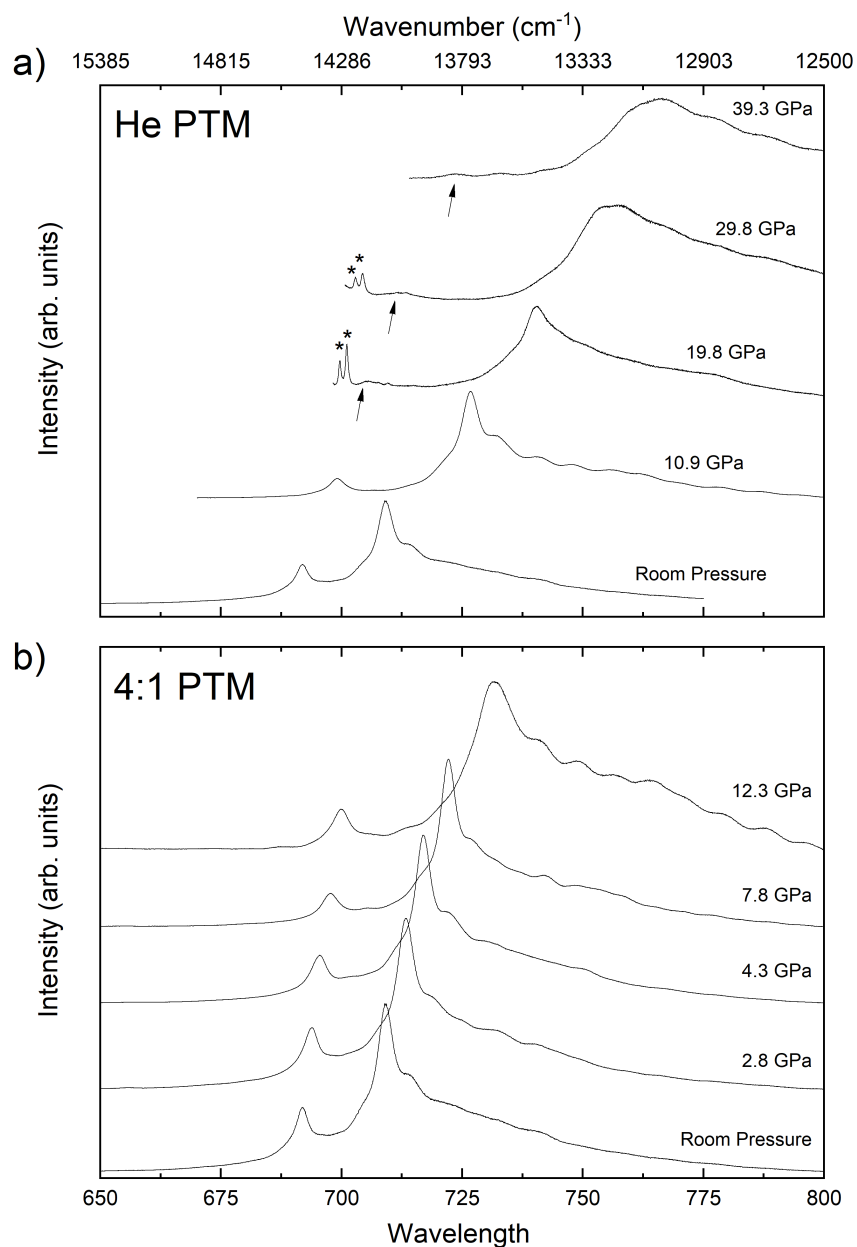


Figure 3. Representative high-pressure luminescence spectra of zoisite (a) He pressure medium and (b) 4:1 methanol ethanol pressure medium. In the He run, ruby emission from the pressure calibrant was also observed in our spectra; these are marked with asterisks, and the arrow indicates the position of the R₂ band of zoisite.

R-line positions as a function of pressure are shown in Figure 4. There is no evidence of a high-pressure phase transition within this pressure range, which is consistent with recent single-crystal X-ray diffraction results (Xu et al., 2019). The decompression data show that the changes

211 that occur on compression are fully reversible with no hysteresis. At ambient conditions, the
212 wavelengths of emission associated with R_1 and R_2 are 709.3 and 691.9 nm. The pressure shift
213 from the He run of R_1 is $-27.2(3) \text{ cm}^{-1}/\text{GPa}$, while for R_2 it is $-12.5(1) \text{ cm}^{-1}/\text{GPa}$; these were
214 determined by non-weighted linear fits. The R_1 shift is the largest known pressure shift of Cr^{3+} R-
215 lines in an oxide reported to date. This is most straightforwardly understood in terms of the
216 compressibility of the M3 site in zoisite. The M3 site in zoisite is more compressible than the
217 octahedral sites in most other minerals in which Cr^{3+} substitutes for Al^{3+} (Alvaro et al., 2012), such
218 as spodumene (Arlt and Angel, 2000), ruby (d'Amour et al., 1978; Finger and Hazen, 1978),
219 clinocllore (Welch and Marshall, 2001), and beryl (Hazen et al., 1986). Because of the differing
220 pressure shifts of R_1 and R_2 , the initially large R-line separation increases even more as pressure
221 increases, and by $\sim 40 \text{ GPa}$ the R-line separation has nearly tripled from its ambient pressure value
222 (Figure 5). Single-crystal X-ray diffraction of zoisite up to 8 GPa confirms that the M3 site becomes
223 more distorted on compression, as manifested by the AV and QE of the M3 site (Alvaro et al.,
224 2012), and our luminescence data indicate that this trend continues to at least 40 GPa .

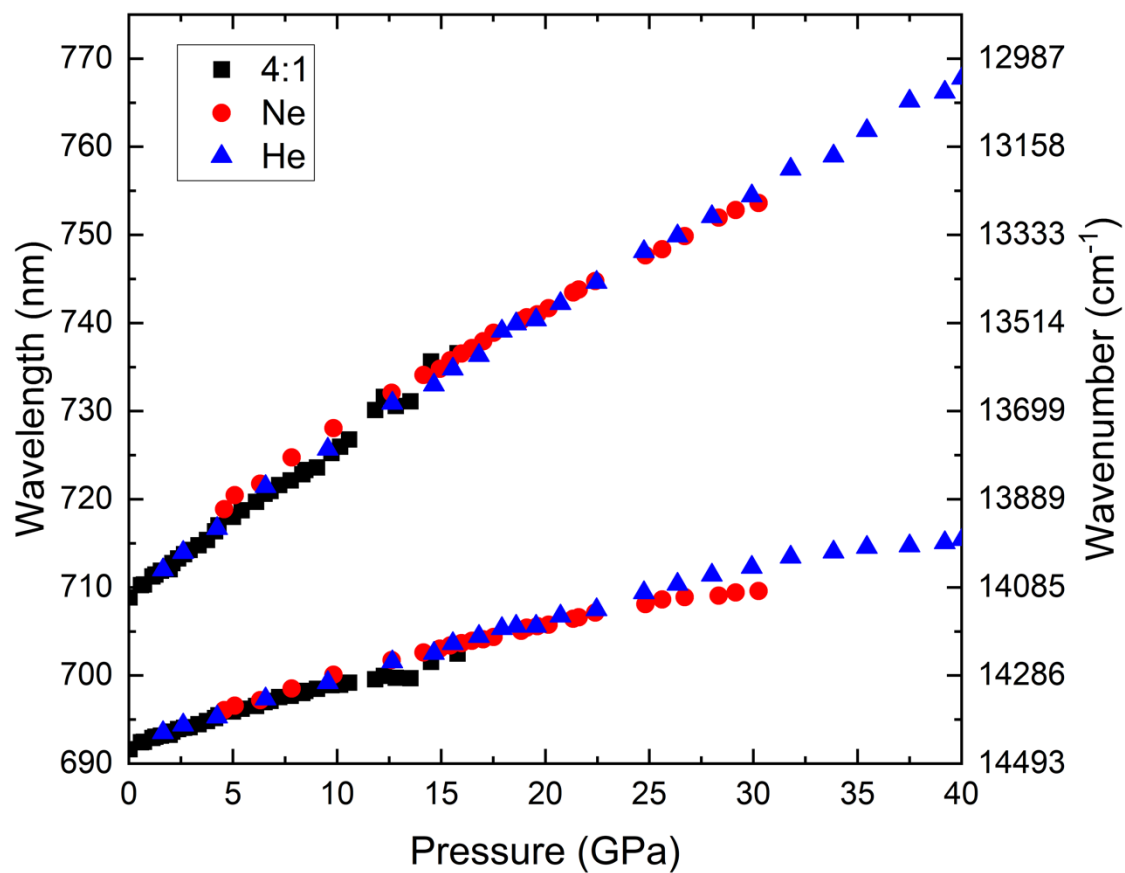


Figure 4. (a) R_1 and R_2 peak positions of zoisite as a function of pressure from three separate runs using different PTM. Error bars are smaller than the symbols.

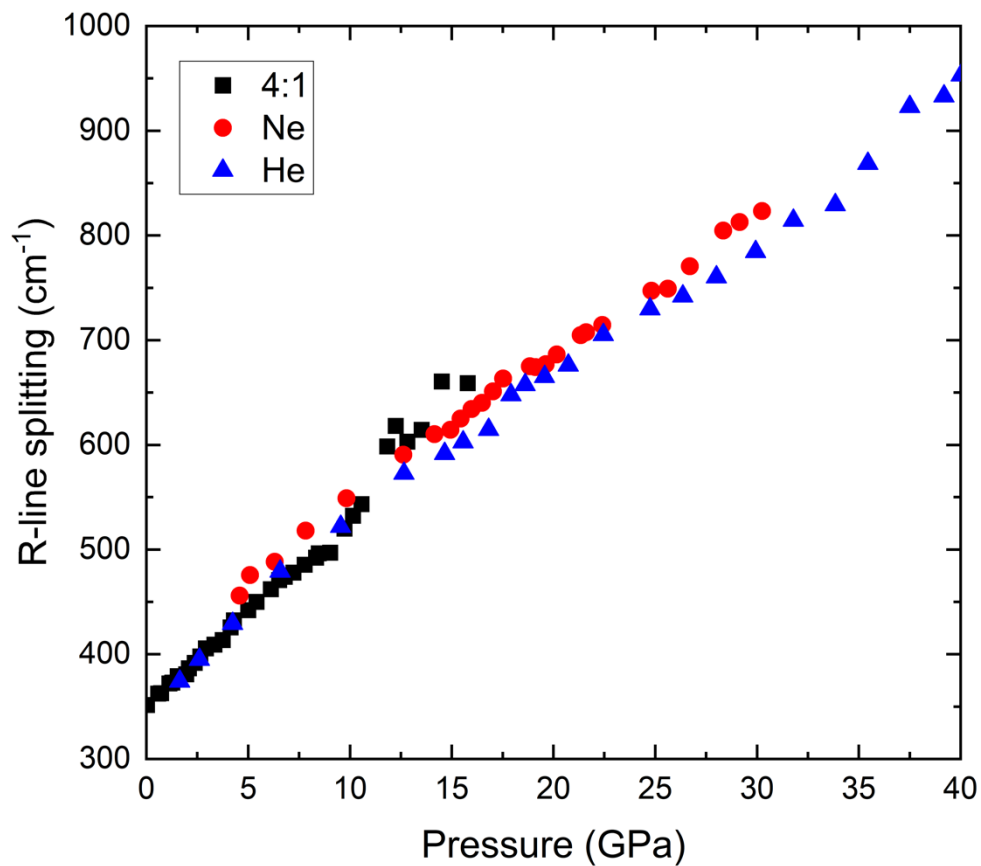


Figure 5. R-line splitting as a function of pressure. Note that the R-line separation nearly triples from its ambient value by 40 GPa. The slight differences in R-line separation between different runs is likely due to the different pressure media that were used. Note the increased scatter above ~10 GPa in the 4:1 methanol:ethanol data.

The results of different PTM are primarily manifested in the R_2 peak positions where changes in slope occur at ~25 GPa in Ne and ~32 GPa in He. It is possible that these changes in slope are indicators of the onset of non-hydrostatic conditions. However, it has been reported that the pressure shift of R_2 in ruby is independent of crystal orientation and non-hydrostatic stresses (Chai and Brown, 1996; Gupta and Shen, 1991), and R_2 may provide a more reliable indicator of the local mean stress. Moreover, the standard deviation in pressure in a He medium at 32 GPa is ~0.15 GPa, while it is ~0.22 GPa in Ne at 25 GPa (Klotz et al., 2009), so these changes in slope may not be occurring at the same differential stress. It is possible that the change in slope

of R_2 is evidence of a change in local compressional mechanism, and/or perhaps be indicative of precursory behavior to a phase transition at pressures > 40 GPa. The intensity of R_2 also decreases dramatically under compression, which increases our uncertainties in the peak position of R_2 above ~ 20 GPa. This decrease in intensity is certainly generated by the increase in R-line separation shown in Figure 5: the amplitude of R_2 relative to R_1 is in accord with Boltzmann's law (e.g., Koziarska et al., 1994), and the increased separation thus suppresses the R_2 amplitude. The anticipated reduction in intensity of R_2 at constant temperature for the factor of 3 increase in splitting across the pressure range of our measurements is about a factor of 20, which is in approximate accord with our observations (Figure 3a).

The broad band emission to the longer wavelength side of the R-lines is more difficult to assign to a particular dopant. This extended, relatively featureless emission shifts to longer wavelengths, and its relative intensity and its width increase at higher pressures (although it is difficult to fit accurately). This broad band emission occurs at wavelengths that are consistent with the $\text{Cr}^{3+} {}^4T_2 \rightarrow {}^4A_2$ transition. It has also been suggested that this broad band emission is due to V^{2+} (Koziarska et al., 1994). Since Cr^{3+} and V^{2+} both have a d^3 configuration, their luminescence spectra are quite similar. The main differences lie in intensity, wavelength, and lifetime of emission associated with these transition metals. This broad band is likely the 4T_2 transition from either Cr^{3+} and/or V^{2+} (although there may also be a vibronic component within the broad band), and it shifts to longer wavelength under compression and persists to the highest pressures of this study, suggesting that the crystal field strength remains weak to intermediate across the pressure range of these measurements. In order to unambiguously assign this broad band emission to Cr^{3+} or V^{2+} or both is not possible from steady state luminescent excitation, as lifetime measurements

267 would be needed.

268 *Kyanite*

269 Kyanite is one of the three Al_2SiO_5 polymorphs found in crustal and upper mantle assemblages,
270 with the other two being sillimanite and andalusite. The pressure-temperature range where these
271 polymorphs are stable span the conditions of regional metamorphism. These polymorphs are important
272 index minerals used in metamorphic petrology to estimate the P - T conditions of rocks. The crystal
273 structure of kyanite was first described by Naray-Szabo et al. (1929) and later by Burnham (1963) who
274 used single-crystal X-ray diffraction. Kyanite is the high-pressure, low temperature polymorph, which has
275 four unique octahedral Al sites and two unique tetrahedral Si sites (Figure 6). There are two groups of
276 octahedra, the chain octahedra (M1 and M2) and the chain-connecting octahedra (M3 and M4). The most
277 distorted metal site (based on the QE and AV) is M3, followed by M2 and M1 which have very similar
278 distortion parameters, and the least distorted and smallest volume site is M4, using ambient single-crystal
279 data from Comodi et al. (1997).

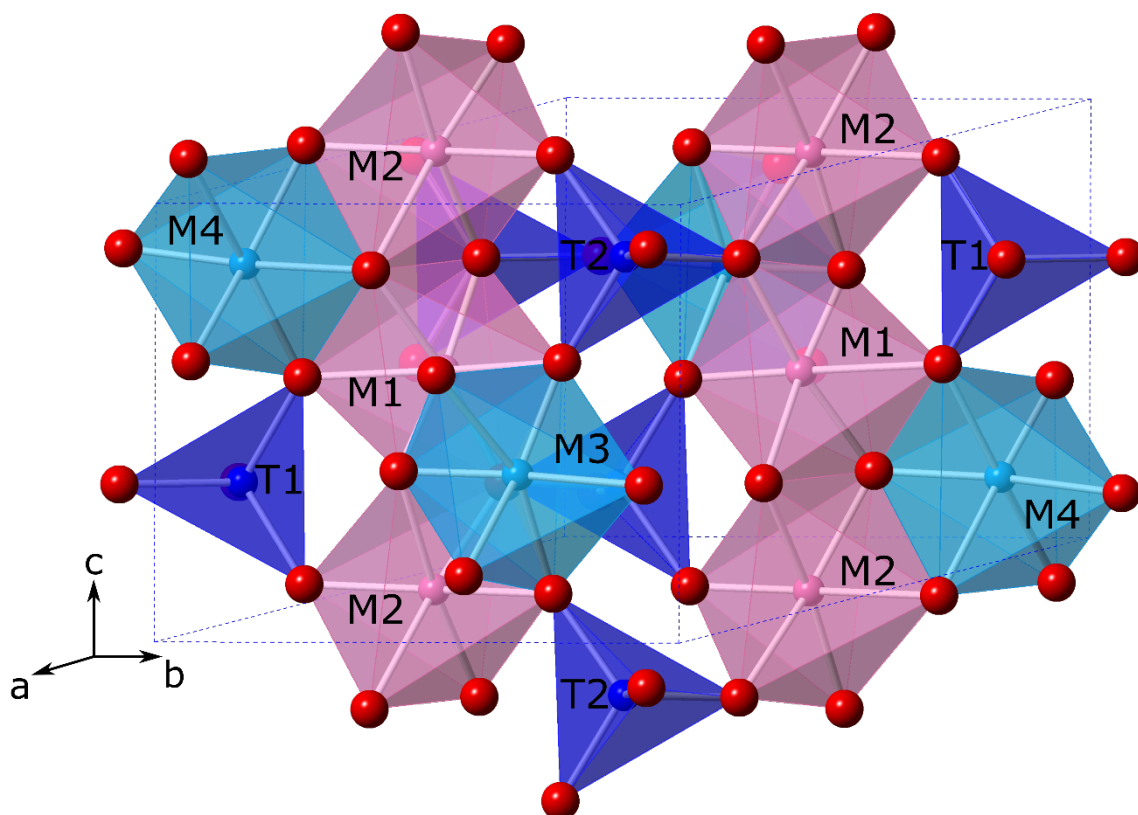


Figure 6. Crystal structure of kyanite. The chain octahedra are M1 and M2 and the chain connecting octahedra are M3 and M4. The two tetrahedral sites are also shown as T1 and T2. Crystal structure diagram generated in CrystalMaker®.

Ambient pressure luminescence spectra collected at 298 and 77 K are in good agreement with previously reported spectra collected at 300 and 26 K (Wojtowicz, 1991), at room temperature and 77 K (Platonov et al., 1998), and at 16 K (Czaja et al., 1995). Our spectra are shown in Figure 7; in order to resolve all of the features in the vicinity of the R-lines and the accompanying broad band emission, different laser powers and collection times were used. At 298 K, two emission bands can be easily deconvolved from each of the sharp emission lines. At 77 K, the two emission bands present at each of the 298 K sharp emission lines are clearly distinct and separated. The two higher intensity components (outer bands) originate from one Cr^{3+} site and are designated R_1 and R_2 , while the two lower intensity components (the inner emission bands) likely originate from a different Cr^{3+} site and are named R_1' and

R_2' . At 298 K the deconvolved peak positions for R_1 and R_2 are 705.9 and 688.5 nm, while for R_1' and R_2' they are 704.7 and 690.3 nm. At 77 K, the R_1 and R_2 bands shift to 704.9 and 687.5 nm while R_1' and R_2' shift to 703.4 and 689.3 nm.

Notably, the measured Cr^{3+} concentration in our kyanite sample is lower than the measured V^{2+} concentration, so it is possible that V^{2+} emission is also observed in our luminescence spectra. Our LA-ICPMS results indicate that the trace element distribution was not homogeneous and that our kyanite sample was chemically zoned. We selected samples that were close to the location where the LA-ICPMS analysis was performed but, given the zonation, it is likely that the samples that were loaded into the DAC had trace element concentrations that differed from the ICPMS analysis spot. However, it has been reported that the Cr concentration in natural kyanite samples with strong narrow emission bands can be very low, with emission observed from samples with less than 0.001 Cr^{3+} per formula unit (Platonov et al. 1998; Gaft et al. 2015). Thus, the quantum efficiency of Cr^{3+} emission in kyanite is likely quite high, in accord with the assumptions of Wojtowicz (1991). The wavelength of the emission bands we observe at both 77 K and 298 K under steady state excitation agree well with previous time resolved luminescence measurements reported by Gaft et al. (2015). Gaft et al. (2015) report that the decay times of the emission bands are consistent with emission from Cr^{3+} . Therefore, in accord with Gaft et al. (2015), all of our assignments assume that Cr^{3+} is the luminescing species responsible for the observed emission bands. We do note, as did Gaft et al. (2015), that a detailed investigation of synthetic kyanites with different Cr and V concentrations is still needed to fully characterize if emission bands from V^{2+} are observed in kyanite.

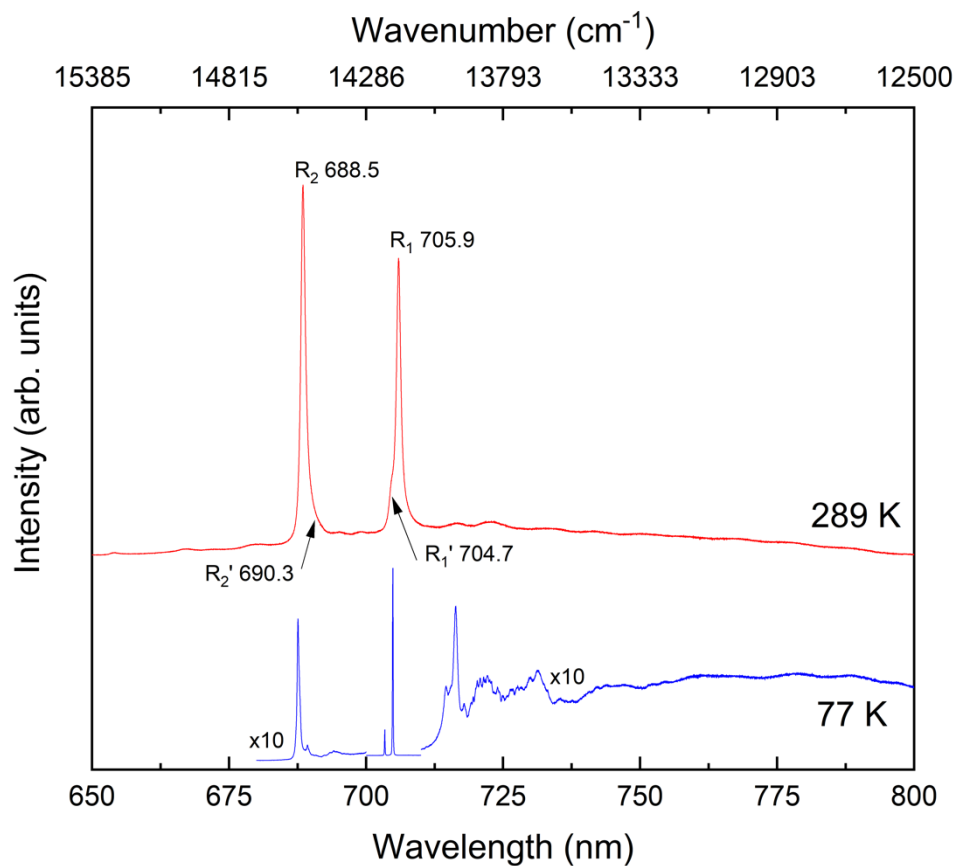


Figure 7. Representative steady state Cr^{3+} luminescence spectra of kyanite 289 K and 77 K. Note that the R_2 emission bands are not entirely quenched in the 77 K spectrum.

Previous assignments of the emission bands have not been straightforward; as discussed above, kyanite has four unique octahedral sites (Figure 6), the chain octahedra (M1 and M2) and the chain-connecting octahedra (M3 and M4). Langer and Seifert (1971) noted that from a geometric point of view, there is no reason why Cr^{3+} would prefer one group of sites over the other. However, based on a comparison between observed quotients of structure factors to those calculated from different models of Cr^{3+} distribution, the best agreement is found when Cr^{3+} does not occupy the M4 site (Langer and Seifert, 1971). Wojtowicz (1991) examined Cr-kyanite using a combination of absorption, steady state, time resolved, and luminescence decay measurements and found that Cr^{3+} emission is observed from 3 of the 4 octahedral sites, with M4 being “dark.” Wojtowicz (1991) proposed that the M4 site could contain Cr^{3+} but that the center might be dark (non-emitting) if the Bartram-Stoneham criterion is fulfilled

(Bartram and Stoneham, 1975). However, Wojtowicz (1991) found that the Bartram-Stoneham criterion was not fulfilled for the M4 site and note that the larger Cr^{3+} ion may not substitute easily into the highest field site (smallest volume site). For comparison, Yeung et al. (1994) conducted a crystal field analysis of the kyanite absorption, emission and excitation spectra and concluded that substitution of Cr^{3+} into the M4 site was unlikely. Three magnetically inequivalent sites were identified in kyanite using electron paramagnetic resonance spectroscopy (Hutton and Troup, 1964), again supporting the presence of one unoccupied site. From a simple crystal structure perspective, M4 is not only the smallest volume site but also the least distorted site ($\text{QE} = 1.013$ and $\text{AV} = 41.9$) meaning Cr^{3+} would be more likely to substitute into the other sites first.

Our 77 K spectrum offers additional insight into the nature of the emission bands near 688 nm. At 26 K, Wojtowicz (1991) did not observe these emission bands: they were quenched to the degree that they were not resolved, and they are quite low amplitude (yet resolved) in the spectrum at 16 K of Czaja et al. (1995). At 77 K, they are clearly observed, which is strong evidence that these emission bands can be assigned as the R_2 component of the split ${}^2\text{E}$ emission band. The site nomenclature previously used by different authors in the spectroscopic literature is inconsistent; we use the mineralogically accepted site nomenclature M1-M4 for the four unique octahedral sites in kyanite. Our assignments are based on (1) site volumes, (2) site distortions, (3) intensity of the emission bands, and (4) the assumption that Cr^{3+} does not substitute into the M4 site. We assign the lower wavelength peak near 688 nm and the higher wavelength peak near 705 nm (the strong outer bands in Figure 7) to R_2 and R_1 originating from the M1 site; this is in accord with the preferred assignment of Yeung et al. (1994). The higher wavelength peak near 688 nm and the lower wavelength peak near 705 nm to R_2' and R_1' are assigned as being associated with the M2 site (See figure 7). That the QE of M1 and M2 are identical (M1 and M2 $\text{QE} = 1.016$) and their angle variance is similar (M1 $\text{AV} = 47.7$ and M2 $\text{AV} = 50.8$), coupled with the locations and splittings of the two sets of bands being generally similar, provides support for our assignments. These assignments are

also in accord with those of the absorption spectrum of blue Tanzanian kyanite of Wildner et al. (2013). The interpretation that the splitting of the emission bands is derived from Cr occupancy within two sites, while in accord with the assignments of Yeung et al. (1994), Czaja et al. (1995), Wojtowicz (1991) and Wildner et al. (2013), is at odds with that of Platonov et al. (1998), who attributed the splitting to a superposition of kyanite emission and emission from a corundum-type exsolving solid solution. Notably, the substitutional behavior of Cr in kyanite may shift at very high Cr contents: in synthetic samples with more than a 1:1 Cr/Al ratio, Hejny et al. (2019) observed a preference of Cr for the M3 site. Within more dilute, natural abundance samples (and perhaps within more dilute synthetic samples: Langer and Seifert, 1971), the spectroscopic consensus is that Cr substitution into the M1 and M2 sites predominates. Thus, while emission from Cr-substitution in M1 and M2 clearly dominate the spectrum, there may be evidence of Cr³⁺ emission from the M3 site manifested as the ⁴T₂ emission band (in the 77 K spectrum): a broad feature centered near 770 nm (Czaja et al., 1995), indicating a weaker crystal field at this site. This is in line with crystal structure data which show that M3 is both the largest metal site and the most distorted (QE= 1.019 and AV= 58.0). However, attempts to follow this broad emission band under pressure were unsuccessful.

We observed differences in the luminescence spectra with different orientations of the sample with respect to the polarization direction of the laser excitation. Our kyanite sample displayed the typical elongated columnar crystal habit, which was easily seen in the small single crystals that were cut from the main specimen. Therefore, the *a*-axis was commonly parallel to the incident laser during our measurements and the *b*-*c* plane was perpendicular to the laser probe. In our ambient measurements, the R₁ peak intensity was lower than the R₂ peak intensity when the *c*-axis was perpendicular to the polarization direction of the laser beam. When the *c*-axis was parallel to the polarization direction of the laser, the R₁ peak intensities were greater than the R₂ peak intensities (Figure 8). These differences in the Cr³⁺ luminescence spectra are most likely connected with the different distortions/orientations of the M1

and M2 sites which alternate along the *c*-axis oriented chains. Gaft et al. (2013c & 2015) also noted this polarization behavior in gem quality greenish kyanite and reported the same change in the relative intensities of R_1 and R_2 with different crystallographic orientations relative to the polarized laser probe. They suggest that the luminescence lines with long decay times are polarized in different directions. During our high-pressure measurements, spectra were collected in both orientations at each pressure step.

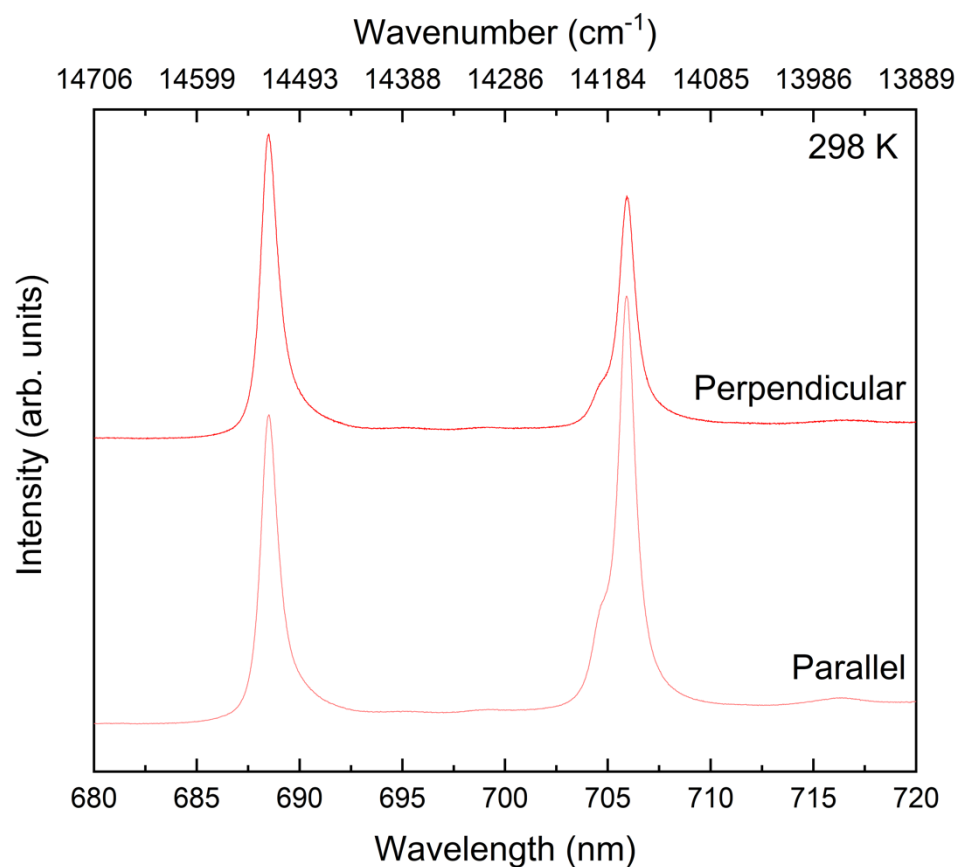


Figure 8. Steady state luminescence spectra of kyanite at ambient conditions showing changes in relative intensities with the *c*-axis of the sample both perpendicular and parallel to the polarization direction of the laser excitation.

Representative high pressure luminescence spectra from the parallel orientation compressed in both a Ne and 4:1 PTM are shown in Figure 9. There are no major differences (beyond relative band intensities and peak widths) in the luminescence spectra, or their pressure dependence, between the two runs in different PTM. In the 4:1 PTM run we collected spectra with the *c*-axis both parallel and

perpendicular to the laser polarization up to a maximum pressure of ~17 GPa. When the c-axis is parallel to the polarization direction of the laser, the two sets of emission bands become well separated above ~1 GPa and can easily be tracked visually. For the perpendicular orientation (not shown), the R_2' component is more difficult to visually track but can be deconvolved across the entire pressure range of these measurements. Attempts to deconvolve the broad, longer wavelength 4T_2 emission bands associated with both the M1 and M2 sites and possibly the M3 site were unsuccessful. The full width at half maximum (FWHM) of R_1 at ambient conditions is $\sim 25\text{ cm}^{-1}$ and in the 4:1 run it remains essentially constant under compression up to ~10 GPa. Above this pressure, the FWHM of R_1 begins to increase, reaching $\sim 46\text{ cm}^{-1}$ at ~16 GPa. In the Ne run, the FWHM of R_1 ranges from $\sim 23\text{--}28\text{ cm}^{-1}$ across the entire pressure range, and no discontinuous changes or changes in slope are observed. This can be seen clearly in Figure 9 by comparing the spectra collected at ~16 GPa in the Ne and 4:1 PTM.

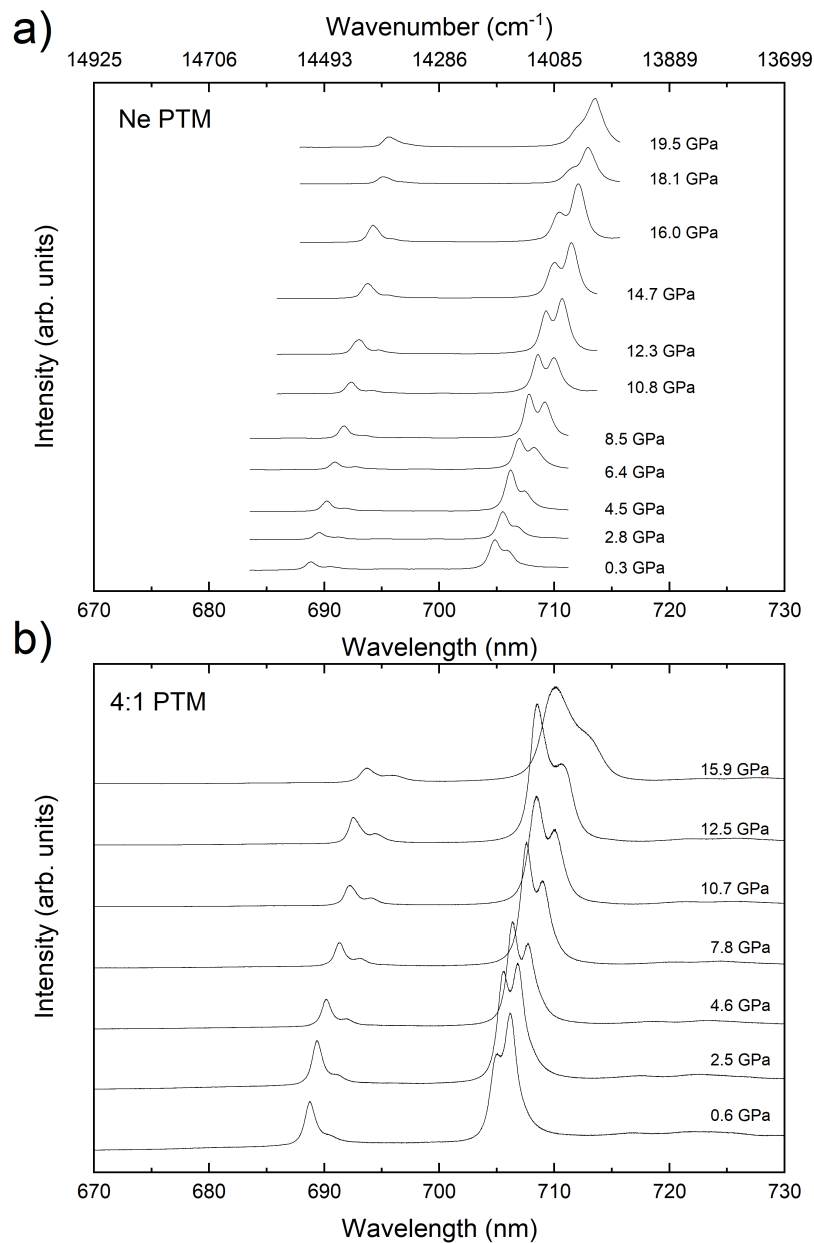


Figure 9. Representative high-pressure Cr³⁺ luminescence spectra of kyanite collected (a) in a Ne PTM with the *c*-axis roughly parallel to the laser polarization and (b) in a 4:1 PTM with the *c*-axis parallel to the laser polarization.

The pressure shift of the R-lines from both the M1 and M2 sites for the parallel orientation from the Ne and 4:1 PTM runs are shown in Figure 10. All the emission bands show a smooth increase in wavelength as pressure is increased, and no discontinuities or significant changes in the slope are observed up to 16 GPa. There is a slight dip in the trend of the R-lines above ~10 GPa in the 4:1 run, and

there may be a slight change in the slope above this pressure. In the Ne run, the R-lines shift essentially linearly and no discontinuities or changes in slope are observed up to ~20 GPa. In the 4:1 run, the R-line splitting from both the M1 and M2 sites remains essentially constant until about 10 GPa, where the separation of the bands associated with the M1 site starts to increase, while it slightly decreases for those generated by the M2 site (Figure 11). The sudden change in the R-line separation may be an indication of a change in compression mechanism in which the M1 site becomes more distorted and the M2 site becomes less distorted. However, in the run where a Ne PTM was used the R-line splitting for the M1 site shows a very slight increase of $0.34(2) \text{ cm}^{-1}/\text{GPa}$ (which was determined from a non-weighted linear fit) while the R-line splitting for the M2 sites remains essentially constant up to ~20 GPa. It is likely that the observed changes in the kyanite luminescence spectra from the 4:1 run are due to the onset of non-hydrostatic conditions above ~10 GPa. It is well known that 4:1 methanol:ethanol undergoes a glass transition at ~10.5 GPa and that pressure gradients form above this pressure (Klotz et al., 2009). Kyanite appears to be particularly sensitive to the onset of non-hydrostatic conditions. It is clear from the high-pressure run in a Ne PTM that the behavior observed above 10.5 GPa in the 4:1 run is anomalous. The pressure shift of the R-lines as determined from non-weighted linear fits to the Ne data are $R_1 = -7.86(3) \text{ cm}^{-1}/\text{GPa}$, $R_1' = -7.67(4) \text{ cm}^{-1}/\text{GPa}$, $R_2 = -7.51(4) \text{ cm}^{-1}/\text{GPa}$, and $R_2' = -7.59(5) \text{ cm}^{-1}/\text{GPa}$. The relative constancy of the R-line separation under pressure is in general accord with previous single-crystal X-ray results to 4.7 GPa that showed that the quadratic elongation of the M1 and M2 sites either remained constant (Comodi et al., 1997) or weakly decreased under compression (Yang et al., 1997).

Notably, two prior studies conducted reconnaissance investigations of the pressure shift of the R_1 line(s) of kyanite to 14 GPa (Mernagh and Liu, 1991; Parthasarathy, 1997). Both studies tracked only the R_1 emission under pressure. Mernagh and Liu (1991) report that above pressures of 5 GPa a single emission band with a weak shoulder on the high frequency side is observed, while Parthasarathy (1997) report that the two R-lines merge at ~6 GPa with a single broad fluorescence peak observed to 14 GPa.

Parthasarathy (1997) also misassigns the two emission bands near 705 nm as R_1 and R_2 , which is inconsistent with this study and previous studies (Czaja et al., 1995; Wojtowicz, 1991). As discussed above and seen in Figure 9, the two R_1 components remain well separated across the entire pressure range of our new measurements when the c -axis of the sample is parallel or nearly parallel to the laser polarization. When the c -axis was perpendicular to the laser polarization, we also observe a decrease in intensity of the higher frequency component, and by ~10 GPa the emission band is asymmetrical but both components can still be deconvolved. It is possible that the merging of the peaks reported by Mernagh and Liu (1991) and Parthasarathy (1997) is the result of changes in sample orientation between pressure steps.

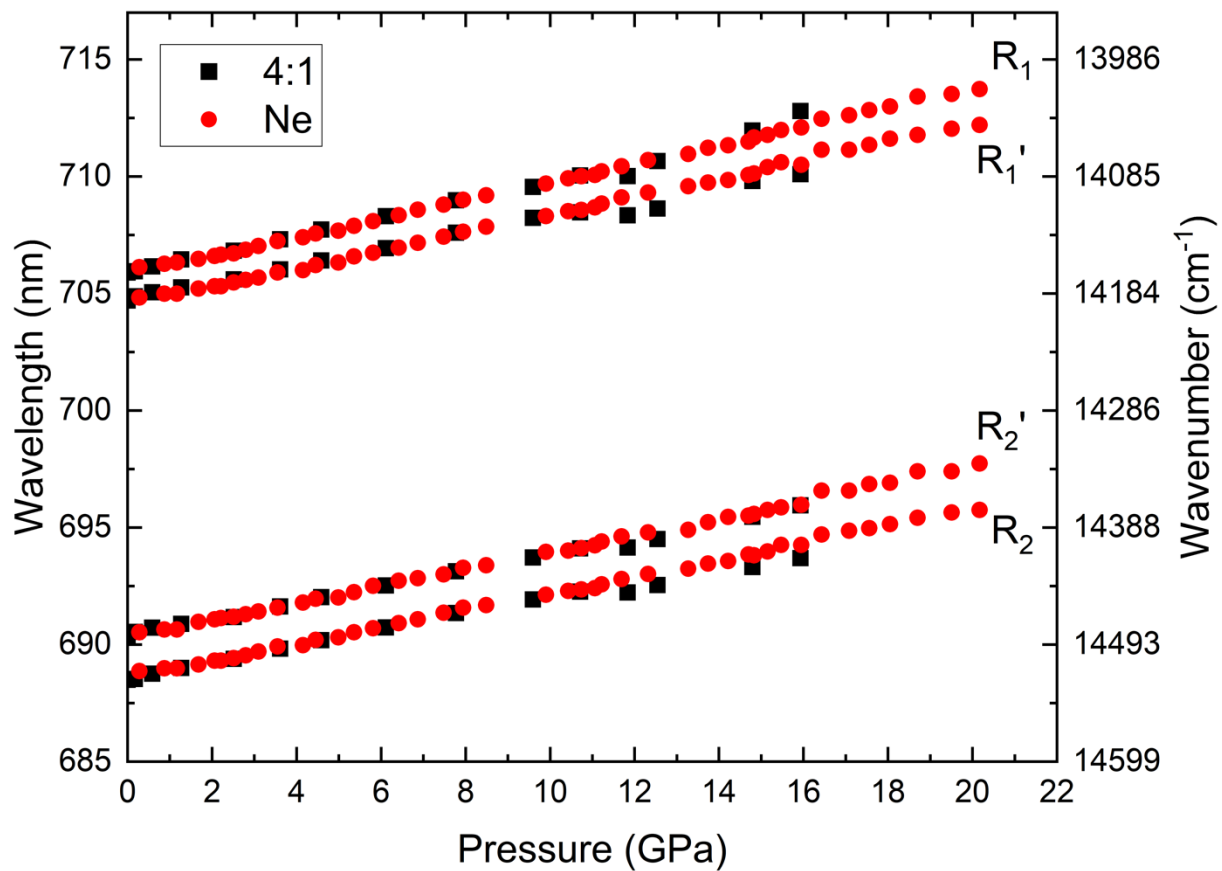


Figure 10. R-line peak positions of kyanite as a function of pressure from two runs; one in a 4:1 PTM and one in a Ne PTM. R_1 and R_2 are associated with the M1 site, and R_1' and R_2' are from the M2 site.

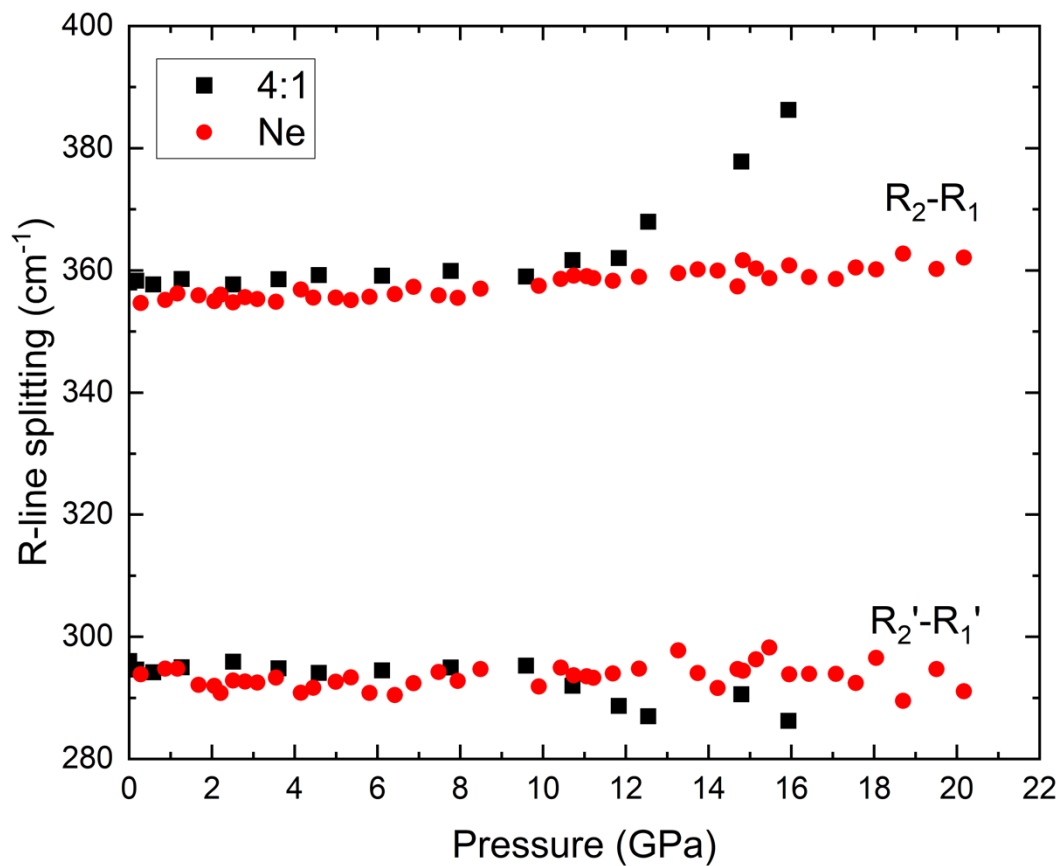


Figure 11. R-line splitting as a function of pressure. For the sample compressed in a 4:1 PTM, the R-line splitting from the M1 site begins to increase above ~10 GPa, and the M2 site shows a slight decrease above 10 GPa. For the sample compressed in a Ne PTM, the R-line splitting changes minimally for both sites up to ~20 GPa.

Recently, a phase transition was reported to occur at 9.7 GPa in kyanite (Gao et al., 2020). Gao et al. (2020) used either a KBr pressure medium or silicone oil pressure medium and collected both Raman and infrared spectra up to 18.6 GPa. They used the appearance of new peaks, discontinuous shifts in certain peaks and changes in the pressure-dependent slope of several peaks as evidence for a phase transition at ~9.7 GPa to a phase they termed kyanite-II. The present luminescence study, as well as previous vibrational and X-ray studies at room temperature and pressures up to 30 GPa, have not found evidence for any phase transitions in kyanite (Comodi et al., 1997; Friedrich et al., 2004; Liu et al., 2009; Mernagh and Liu, 1991; Yang et al., 1997). In particular, Mernagh and Liu (1991) conducted Raman spectroscopy to 14 GPa; Liu et al. (2009) performed powder diffraction to 17.5 GPa, and Yagi et al. (1998)

to 10.2 GPa. Friedrich et al. (2004) presented evidence for the persistence of kyanite to 30 GPa, with the possible onset of a high-pressure transition above 30 GPa based on powder X-ray diffraction results. Gao et al. (2020) offer three possible explanations for this discrepancy: (1) different PTM, (2) better spectral resolution than previous studies, and (3) differences in trace element concentrations. Gao et al. (2020) used KBr and silicone oil as their pressure media; each of these media is characterized by non-hydrostatic gradients at low pressure conditions, with KBr supporting pressure gradients at low pressure conditions, and silicone oil showing signs of pressure gradients above ~2.5 GPa (Klotz et al., 2009). Therefore, it is possible that non-hydrostatic conditions may be at the root of this discrepancy. Differences in spectral resolution between Gao et al. (2020) and (Mernagh and Liu, 1991) may be a possibility with respect to Raman results, but it does not explain the non-observance of this transition in high-pressure powder diffraction studies (Friedrich et al., 2004; Liu et al., 2009). Lastly, Gao et al. (2020) postulate that the very low Fe content of their sample can possibly explain the phase transition, but the sample used in this study and the study of Yagi et al., (1998) both had low Fe concentrations, and no evidence of the proposed transition at 9.7 GPa was observed in either study.

The role of non-hydrostatic stress in generating (or inducing) possible phase transitions is complex. If the proposed 9.7 GPa transition reported by Gao et al. (2020) is generated by non-hydrostatic stress, then an extreme scenario would involve bridging of kyanite within the diamond cell. This would introduce a complex stress field characterized by a large uniaxial component. Indeed, Gao et al. (2020) loaded inhomogeneous thickness samples in a KBr pressure medium for Raman mapping of the coexistence of kyanite and the proposed kyanite-II phases, with the rationale being “to generate a larger concentration of stress”. In one of our spectroscopic experiments, the kyanite sample bridged between the anvils, and we observed several new emission bands within this sample alone. In this run, we compressed the sample directly to ~11 GPa and collected a spectrum (Figure 12). The new emission bands in this spectrum indicated that a subset of the Cr³⁺ ions were present within octahedral sites with different

distortions and volumes than the M1 and M2 sites. Within this bridged sample, we increased the pressure to above 15 GPa and collected another spectrum where additional emission bands were observed, indicating that the new stress field had further modified the Cr³⁺ environments. On decompression, by ~9.0 GPa, the sample was no longer bridged. All the spectra collected at ~9.0 GPa and below (within a hydrostatic regime) showed the typical set of emission bands assigned to M1 and M2.

Importantly, uniaxial stress has long been known to generate structural changes that can induce splitting of fluorescence lines (Imbusch et al., 1965). For example, ruby luminescence spectra collected with different uniaxial stress applied along the direction $\theta = 61 \pm 2^\circ$ and $\omega = 8 \pm 2^\circ$ showed splitting due to separation of the Cr³⁺ ions into two σ -sublattices (Monteil et al., 1984). Since our sample was oriented in the DAC when the sample bridged, the applied uniaxial stress was along the a -direction. Thus, while the predominant non-hydrostatic stress would likely be close-to-uniaxial, the full stress tensor in a bridged, oriented sample within a pressure medium is not expected to be simple. The key point here is that Figure 12 demonstrates that particularly extreme non-hydrostatic stress in kyanite can induce reversible structural changes within the first nearest neighbor coordination sphere relative to normal kyanite. We speculate that this reversible behavior under differential stress could be driven by a soft shear mode within kyanite's elastic tensor: however, constraints on the individual elastic constants of kyanite at high pressure do not exist.

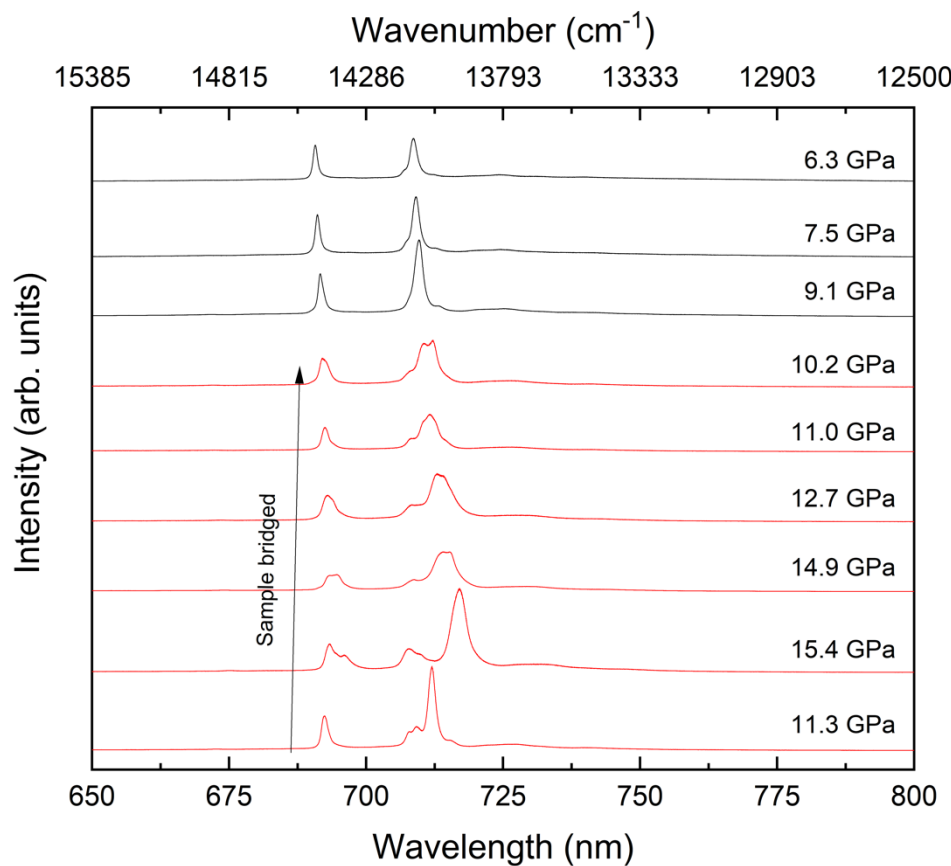


Figure 12. High pressure Cr^{3+} luminescence spectra of kyanite where the sample bridged between the diamond anvils during the experiment. The red spectra (11.3 to ~10.2 GPa, bottom six spectra) were collected while the sample was bridged. In the black spectra (9.1 to 6.3 GPa), the sample was no longer bridged between the diamond anvils and the luminescence spectra show emission bands from only the M1 and M2 sites.

Conclusions

The Cr^{3+} luminescence spectra of zoisite and kyanite have been characterized up to 40 and 20 GPa, respectively. These two minerals each have very large R-line splittings, and for zoisite the separation increases as pressure increases and more than triples from its ambient value by 40 GPa. For comparison, in kyanite the splitting remains virtually constant up to ~20 GPa when compressed in a Ne PTM. The markedly different behavior of the Cr^{3+} substituted octahedral sites (the M3 site in zoisite and the M1 and M2 sites in kyanite) indicate that the zoisite site increases dramatically in its distortion under pressure, while the two sites in kyanite are largely unchanged: each of these results are in accord with substantially

lower pressure single-crystal diffraction results and demonstrate that R-line luminescence is a particularly sensitive probe of the octahedral metal sites. Notably, kyanite appears to be particularly sensitive to the onset of non-hydrostatic conditions. Phase transitions are not observed in either zoisite or kyanite up to the maximum pressure of each study, consistent with most previous studies. Lastly, we illustrate that the stress field of the sample is important for possible structural changes within kyanite, and that non-hydrostatic/nearly uniaxial stress components can produce anomalous luminescence and structural results in this material.

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Author Contributions

E.F.O. prepared and carried out the experiments. E.F.O. carried out the data analysis with input from Q.W. E.F.O. and Q.W. wrote the manuscript.

Competing Interests

The authors declare no competing interests.

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