



On the Metallicities and Kinematics of the Circumgalactic Media of Damped Ly α Systems at $z \sim 2.5^*$

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Abstract

We use medium- and high-resolution spectroscopy of close pairs of quasars to analyze the circumgalactic medium (CGM) surrounding 32 damped Ly α absorption systems (DLAs). The primary quasar sightline in each pair probes an intervening DLA in the redshift range $1.6 < z_{\text{abs}} < 3.5$, such that the secondary sightline probes absorption from Ly α and a large suite of metal-line transitions (including O I, C II, C IV, Si II, and Si IV) in the DLA host galaxy’s CGM at transverse distances $24 \text{ kpc} \leq R_{\perp} \leq 284 \text{ kpc}$. Analysis of Ly α in the CGM sightlines shows an anticorrelation between R_{\perp} and H I column density (N_{HI}) with 99.8% confidence, similar to that observed around luminous galaxies. The incidences of C II and Si II with $N > 10^{13} \text{ cm}^{-2}$ within 100 kpc of DLAs are larger by 2σ than those measured in the CGM of Lyman break galaxies ($C_f(N_{\text{CII}}) > 0.89$ and $C_f(N_{\text{SiII}}) = 0.75^{+0.12}_{-0.17}$). Metallicity constraints derived from ionic ratios for nine CGM systems with negligible ionization corrections and $N_{\text{HI}} > 10^{18.5} \text{ cm}^{-2}$ show a significant degree of scatter (with metallicities/limits across the range $-2.06 \lesssim \log Z/Z_{\odot} \lesssim -0.75$), suggesting inhomogeneity in the metal distribution in these environments. Velocity widths of C IV $\lambda 1548$ and low-ionization metal species in the DLA versus CGM sightlines are strongly ($>2\sigma$) correlated, suggesting that they trace the potential well of the host halo over $R_{\perp} \lesssim 300 \text{ kpc}$ scales. At the same time, velocity centroids for C IV $\lambda 1548$ differ in DLA versus CGM sightlines by $>100 \text{ km s}^{-1}$ for $\sim 50\%$ of velocity components, but few components have velocities that would exceed the escape velocity assuming dark matter host halos of $\geq 10^{12} M_{\odot}$.

Unified Astronomy Thesaurus concepts: Circumgalactic medium (1879); Damped Lyman-alpha systems (349); Quasar absorption-line spectroscopy (1317)

Supporting material: figure set

1. Introduction

The circumgalactic medium (CGM) is defined as the gaseous halo surrounding galaxies that hosts the exchange of gas between large-scale outflows from the host galaxy interstellar medium (ISM), the ambient halo, and accretion from the

intergalactic medium (IGM; Tumlinson et al. 2017). In the last decade, it has become evident that studying the CGM is crucial to fully understanding galaxy evolution (e.g., Peebles et al. 2019). Studies have shown $<25\%$ of the mass expected in $\sim L$ galaxies at $z \sim 0$ is detected in observations of stars or the interstellar medium (ISM; Peebles et al. 2014), potentially making the CGM an opportune place to look for these missing baryons. CGM studies at low redshift have detected the majority of these galactic missing baryons, with just the cool gas mass of the CGM around $\sim L^*$ galaxies estimated to be $\sim (3-6) \times 10^{10} M_{\odot}$ (Prochaska et al. 2011; Werk et al. 2014; Prochaska et al. 2017a). Although the physical origin of these diffuse baryons remains unclear, it is likely that they are deposited in part by winds launched by star formation or active galactic nuclei in the central galaxy (e.g., Bordoloi et al. 2011; Kacprzak et al. 2012; Lan & Mo 2018), as well as by accretion

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of the IGM or recycled wind material (Kereš et al. 2009; Oppenheimer et al. 2010; Rubin et al. 2012). Hence, the detailed properties of the CGM can provide an essential insight into the processes driving the evolution of the galaxies.

A critical epoch to study the CGM is at $z \sim 2-3$, during the peak of cosmic star formation (Storrie-Lombardi & Wolfe 2000; Madau & Dickinson 2014) and supermassive black hole growth (Marconi et al. 2006; Richards et al. 2006). However, the gaseous material that makes up the CGM is diffuse and difficult to detect in emission, especially at higher redshifts (although see, e.g., Arrigoni Battaia et al. 2018; Erb et al. 2018; Cai et al. 2019; O’Sullivan et al. 2020). As a result, the vast majority of high-redshift CGM studies have analyzed its H I and metal content in absorption detected along sightlines to bright background QSOs. For the most part, this work has focused on characterizing the gaseous environments of systems with host galaxies that are bright in the rest-frame UV (i.e., QSOs and Lyman Break Galaxies). These studies have established the masses and extent of the neutral hydrogen overdensities around these systems (e.g., Rakic et al. 2012; Rudie et al. 2012; Prochaska et al. 2013), and have likewise assessed the sizes and masses of their metal components (e.g., Adelberger et al. 2005; Simcoe et al. 2006; Prochaska et al. 2014; Turner et al. 2014; Lau et al. 2016; Rudie et al. 2019a). With the recent advent of VLT/MUSE, it is now also possible to select large samples of Ly α -emitting systems close to background QSO sightlines, enabling similar analyses of the bulk properties of their halos (e.g., Muzahid et al. 2021; Lofthouse et al. 2023).

Studies benefiting from sensitive, high-resolution background QSO spectroscopy have gone beyond assessment of these bulk properties to analyze quantities that provide constraints on the origins and ultimate fate of the circumgalactic material. Detailed analyses of metal-line kinematics along sightlines probing QSO host halos have found that low-ionization transitions trace $\approx 300 \text{ km s}^{-1}$ line widths, which is consistent with material tracing virial motions in these halos, but the robustly-measured asymmetries in these profiles are suggestive of gas tracing large-scale outflows (Lau et al. 2018). Similarly, Rudie et al. (2019a) found that within projected distances (R_{\perp}) $< 100 \text{ kpc}$, the majority of Lyman Break Galaxies (LBGs) exhibit metal-enriched halo gas with velocities which exceed that required to escape the system. In contrast, Turner et al. (2017) found that the absorption kinematics of H I, C IV, and Si IV at larger projected separations (up to $R_{\perp} = 2 \text{ Mpc}$) from LBGs are best explained by large-scale inflow onto their host halos.

Assessment of the metallicity of QSO host halo gas has revealed high levels of overall enrichment ($[M/H] \gtrsim -0.6$) and significantly α -enhanced abundance ratios (Lau et al. 2016; Fossati et al. 2021), which is suggestive that core-collapse supernovae play a dominant role in the enrichment of these environments. A handful of studies have analyzed the metallicities of material both within and well beyond the virial radii of LBGs, uncovering examples of systems as distant as $R_{\perp} = 30 - 110 \text{ kpc}$ that exhibit large scatter in their enrichment levels (e.g., $Z/Z_{\odot} \gtrsim 0.08-0.3$; Simcoe et al. 2006; Crighton et al. 2013; Fumagalli et al. 2017). Several studies (Crighton et al. 2013, 2015; Fumagalli et al. 2016a; Lofthouse et al. 2020) have also now provided clear evidence that at least some high-redshift star-forming galaxy halos are not well-mixed, and can give rise to both near-pristine material ($Z/Z_{\odot} \sim 1/100$) and

mildly subsolar gas along the same background QSO sightline (e.g., with velocity offsets of $\approx 200 \text{ km s}^{-1}$; Crighton et al. 2013). These latter authors in particular used their measurements to argue for the presence of a cold-accretion stream amidst extended and metal-enriched wind ejecta.

Due to their use of continuum or Ly α emission for the identification of foreground galaxy samples, these works have assessed the gaseous environments of halos hosting active galaxies with total halo masses $M_h \gtrsim 10^{11-11.5} M_{\odot}$ (Adelberger et al. 2005; Gawiser et al. 2007; Conroy et al. 2008; Wild et al. 2008; Rakic et al. 2013; White et al. 2012; Font-Ribera et al. 2013; Bielby et al. 2016). An alternative approach is to instead identify high-redshift galaxies from their absorption-line signatures. We pioneered this technique in Rubin et al. (2015), which used spectroscopy of close pairs of QSO sightlines to search for the damped Ly α (DLA) absorption profile associated with galaxies in the foreground. DLAs, defined as absorbers having neutral hydrogen column densities $N_{\text{HI}} \geq 2 \times 10^{20} \text{ cm}^{-2}$ (Wolfe et al. 1986), have been the dominant reservoirs of H I gas since at least $z \sim 5$ (Wolfe et al. 2005). The spatial relationship between DLAs and high-redshift star formation has remained opaque for more than two decades after their discovery; however, studies of DLAs in cosmological simulations have long suggested that they are associated with galaxies spanning a wide range of halo masses ($10^{10} M_{\odot} \lesssim M_h \lesssim 10^{12} M_{\odot}$; Haehnelt et al. 1998; Nagamine et al. 2004; Pontzen et al. 2008; Razoumov et al. 2008; Fumagalli et al. 2011; Cen 2012; Bird et al. 2014; Garratt-Smithson et al. 2021).

Observational searches for the luminous counterparts of DLAs have recently become successful thanks to the advent of near-infrared IFUs on 8–10 m-class telescopes (Péroux et al. 2012; Jorgenson & Wolfe 2014), and in programs targeting DLAs with high metallicities (Fynbo et al. 2010, 2013; Krogager et al. 2012, 2017; Noterdaeme et al. 2012). A meta-analysis of these latter studies conducted by Krogager et al. (2017) concluded that DLAs with detected counterparts typically arise within $R_{\perp} < 25 \text{ kpc}$ of galaxies having star formation rates (SFRs) $\approx 1-30 M_{\odot} \text{ yr}^{-1}$, and that lower-metallicity DLAs are likely associated with host galaxies having luminosities that extend $\gtrsim 2$ magnitudes fainter, and with SFRs as low as $\sim 0.01 M_{\odot} \text{ yr}^{-1}$. More recent follow-up of high-metallicity DLA hosts at $z \sim 4$ with ALMA (Neeleman et al. 2017, 2019; Prochaska et al. 2019) has identified massive high-SFR ($\sim 7-110 M_{\odot} \text{ yr}^{-1}$) counterparts with impact parameters of 10–50 kpc (Neeleman et al. 2019). Most recently, a Keck Cosmic Web Imager (KCWI) study has mapped two DLAs along multiple lines of sight toward a bright gravitationally-lensed background galaxy at $z = 2.7$ (Bordoloi et al. 2022). This work identified Ly α emission that likely arises from the DLA hosts within $R_{\perp} \lesssim 1-2 \text{ kpc}$ of the damped sightlines, and moreover measured the spatial extent of these DLAs to be $\gtrsim 238 \text{ kpc}^2$ (assuming circular geometry with $d = 17 \text{ kpc}$), implying neutral gas masses of $\gtrsim 5.5 \times 10^8 - 1.4 \times 10^9 M_{\odot}$.

Taken together, these studies are suggestive of a scenario in which high-metallicity DLAs arise close to actively star-forming galaxies at high redshift (with SFRs $\gtrsim 1 M_{\odot} \text{ yr}^{-1}$ and halo masses $10^{11} M_{\odot} \lesssim M_h \lesssim 10^{12} M_{\odot}$), while lower-metallicity DLAs likely trace halos with lower star formation rates and masses. This scenario is further corroborated by analyses of the relation between DLA absorption-line widths, metallicities, and host galaxy stellar masses (e.g., Neeleman et al. 2013; Møller et al. 2013; Christensen et al. 2014). This picture is also

Table 1
List of Instruments

Instrument	Telescope	Resolution (\mathcal{R})	Δv (km s ⁻¹) ^a	Wavelength Coverage
MIKE (Blue+Red)	Magellan Clay	35714	8	3350–9500 Å
MIKE-Blue	Magellan Clay	28,000	11	3350–5000 Å
XSHOOTER	VLT UT2	8000	37	3000–25000 Å
MagE	Magellan Clay	5857	51	3100–10000 Å
		4824	62	3100–10000 Å
ESI	Keck II	4545	66	3900–10,900 Å
BOSS	Sloan 2.5 m Telescope	2100	143	3600–10400 Å
GMOS-N	Gemini Telescope	1872	160	3600–9400 Å

Note.

^a Δv is the velocity width of the FWHM resolution element.

broadly consistent with that advocated by analytical work modeling the global distribution function of N_{HI} for DLAs in tandem with their metallicities (Krogager et al. 2020), line widths, and molecular gas content (Theuns 2021). This implies that DLAs are effective signposts for high-redshift galaxies having a broad range of masses and SFRs, and furthermore that the metallicities of DLAs may provide a rough indication of their relative host halo masses (e.g., Wolfe & Prochaska 1998; Ledoux et al. 2006; Neeleman et al. 2013).

In our previous work (Rubin et al. 2015), we searched optical spectroscopy of close pairs of quasars (Findlay et al. 2018) for pairs in which at least one line of sight probed an intervening DLA. Our search yielded a sample of 40 pairs with foreground DLAs having redshifts in the range $1.6 < z_{\text{DLA}} < 3.6$. Our quasar spectroscopy was for the most part obtained at low spectral resolution ($\mathcal{R} \lesssim 2000$), but nevertheless permitted assessment of the covering fraction of optically thick H I and the incidence of strong Si II $\lambda 1526$ and C IV $\lambda \lambda 1548, 1550$ absorption in DLA environments to projected distances $R_{\perp} < 300$ kpc. Since this first work, we have obtained follow-up spectroscopy of a subset of this sample at medium and high spectral resolution ($\mathcal{R} \gtrsim 4000$), enabling detailed assessment of the column densities and kinematics of several ionic species, including Al II, Al III, C II, C IV, Fe II, Mg II, O I, Si II, and Si IV. We present these measurements here, together with a comparison between these CGM properties and those measured in the denser environments surrounding LBGs and QSO hosts (Lau et al. 2016; Rudie et al. 2019a). We assume here that our sample DLAs serve as signposts for nearby star formation, and are located either within the interstellar medium of their galaxy hosts or (more likely) in their “inner” CGM (e.g., Theuns 2021; Stern et al. 2021). Our data set thus offers the unique opportunity to constrain metallicities of both the extended CGM and the ISM/inner CGM material traced by DLAs. Although such comparisons are expected to yield important insight into the origins of circumgalactic gas, they have only been attempted in relatively low-redshift ($z < 1.5$) systems to date (e.g., Péroux et al. 2016; Prochaska et al. 2017a; Kacprzak et al. 2019; Weng et al. 2023).

Our sample selection and data preparation are described in Section 2. We then discuss our methods for measuring H I column densities in the DLA and CGM sightlines in Section 3, and discuss our methods for measuring metal-line column densities and kinematics in Section 4. Section 5 presents the resulting column densities, metallicities, and kinematics measured for our DLA and CGM sightlines. Finally, in Section 6 we combine our results with those in the literature,

and present summaries of the relation between metallicity and R_{\perp} around both DLAs and LBGs, as well as of the relation between metallicity and velocity width. We adopt a Planck Λ CDM cosmology with $\Omega_{\text{M}} = 0.3$, $\Omega_{\Lambda} = 0.70$, and $H_0 = 68$ km s⁻¹ Mpc⁻¹ (Planck Collaboration et al. 2016).

2. Data and Sample Selection

2.1. QSO Pair Sample Selection

Our sample is primarily selected from the Quasars Probing Quasars (QPQ) spectral database (as described in Findlay et al. 2018). The QPQ database contains spectra for 5,627 objects with $z > 2$ which were collected for the purpose of observing pairs of quasars that have close transverse separations on the sky. QPQ targets were initially drawn from low-resolution spectroscopic and photometric surveys that identified sources as quasars, including the SDSS Legacy Survey (2000–2008; York et al. 2000; Gunn et al. 2006), the Baryon Oscillation Spectroscopic Survey (Dawson et al. 2013), and the 2dF QSO Redshift Survey (Croom et al. 2004). These targets were supplemented with a photometrically-selected sample of QSO pair candidates, with photometry measured in SDSS, VST ATLAS (Shanks et al. 2015), and WISE imaging (Wright et al. 2010). Photometrically-identified pairs were followed up with spectroscopy using 2–4 meter-class telescopes, as described in Hennawi et al. (2006, 2010). A subset of these confirmed candidates that were close on the sky (within $\lesssim 30''$) and that have $g \lesssim 21.5$ were then observed with medium- or high-resolution spectrographs, including ESI (Sheinis et al. 2002) on the Keck II telescope, MagE (Marshall et al. 2008) and MIKE (Bernstein et al. 2003) on the Magellan Telescopes, and XSHOOTER on the Very Large Telescope (Vernet et al. 2011). The majority of these high-fidelity spectra were obtained for the purpose of studying the CGM of the foreground QSOs (Prochaska & Hennawi 2009; Lau et al. 2016, 2018). A subset of this sample was targeted specifically for the present study due to the presence of a foreground DLA that was discovered in lower-resolution spectroscopy. A full listing of the telescopes and instruments used to obtain data analyzed in this paper, along with the corresponding spectral coverage and resolution of each instrumental setup, is presented in Table 1.

In an effort to increase our quasar pair sample, we also searched the IGMspec database. IGMspec is a large database that contains 434,686 spectra in the UV, optical, and near-infrared from 16 different surveys (Prochaska 2017). The database includes all the quasars from BOSS DR7 (Abazajian et al. 2009) and DR12 (Alam et al. 2015). Our search yielded eight pairs using the selection criteria described below;

however, none of these sightlines were found to probe foreground DLAs.

From the QPQ and IGMspec databases, we selected only quasar pairs with a maximum transverse proper distance on the sky of $R_{\perp} = 300$ kpc (calculated at the redshift of the foreground QSO). This is much larger than the typical virial radius of massive LBGs at $z \sim 2$ ($R_{\text{virial}} \sim 90$ kpc), and thus this distance criterion allows us to probe the CGM both within and beyond the virial radii of DLA host galaxies at $z \sim 2$. We required that the quasars have redshifts $1.58 < z < 4$ so that their $\text{Ly}\alpha$ transition falls redward of the atmospheric cutoff at 3140 \AA , and so that there is wavelength coverage redward of the $\text{Ly}\alpha$ forest. This initial query yielded 411 QSO pairs.

We then required that at least one QSO in the pair have a medium- or high-resolution spectrum (with $\mathcal{R} \geq 4000$) to enable the metal-line analysis that is described later in Section 4. While the majority of these sightlines were targeted solely due to the presence of a foreground QSO (i.e., for reasons unrelated to the possible presence of a foreground DLA), a subset were targeted after the discovery of a DLA in low-resolution spectroscopy. This latter subsample may be biased toward probing low H I column density DLAs, due to broadening of DLA absorption profiles in low-resolution spectroscopy. However, this effect is small because this marginally biased sample comprises a small fraction of the medium-/high-resolution spectra used in this work. The resulting sample consists of 85 pairs. For each pair, we collected all spectra in each database, including low-resolution spectra if they extended the blue wavelength coverage, to maximize our wavelength search window for DLA signatures. For two pairs in this sample, we also made use of HST WFC3/UVIS grism spectroscopy obtained and reduced as described in Lusso et al. (2018). These data cover $2000 \text{ \AA} < \lambda_{\text{obs}} < 4500 \text{ \AA}$ at a FWHM resolution of $\sim 60 \text{ \AA}$, and therefore can provide useful coverage of the Lyman limit for absorbers discovered at $z < 2.5$ along these sightlines. As described below in Section 3, the HST data was used to improve our constraints on N_{HI} for sightlines with medium or high-resolution optical spectroscopy.

2.2. Continuum Fitting

We fit each quasar continuum using the function `fit_continuum` in the Python package `linetools`¹⁸ (Prochaska et al. 2016), which allows the user to interactively modify a spline fit to the level of the continuum across the spectrum. The typical uncertainty in the continuum level using this method is $\lesssim 10\%$ in the $\text{Ly}\alpha$ forest and $\sim 5\%$ redward of the QSO's $\text{Ly}\alpha$ line (Prochaska et al. 2013).

2.3. Identification of QSO Pairs with Foreground DLAs

We then performed a search for foreground DLAs among these pairs. Initially, we searched each spectrum in a given pair for strong absorption features blueward of the quasar $\text{Ly}\alpha$ emission line. We required that these features meet the following criteria:

1. They appear as a single line with apparent damping wings.
2. The DLA candidate must have a redshift more than 5000 km s^{-1} blueward of the foreground QSO. If the system is redward of that limit, then it may be associated

with the QSO and may not probe the same environment as DLAs that are intervening.

3. There is metal-line absorption present at the same redshift as the putative $\text{Ly}\alpha$ in the same sightline. We searched for metal absorption lines within a $\pm 350 \text{ km s}^{-1}$ window from the transitions Si II $\lambda 1304$, Si IV $\lambda 1526$, O I $\lambda 1302$, and C II $\lambda 1334$ (Wolfe et al. 2005). We chose this velocity window to ensure that we encompass any absorption that could be associated with the DLA. Assuming that DLAs are predominately hosted by halos with masses up to $10^{12} M_{\odot}$, and that the FWHM of the line-of-sight velocity distribution of virialized halo gas is $v_{\text{FWHM}} = 2\sqrt{\ln 2} v_{\text{vir}}$, we expect $v_{\text{FWHM}} \approx 360 \text{ km s}^{-1}$ at the mean redshift of our sample ($\langle z_{\text{abs}} \rangle = 2.45$; Maller & Bullock 2004). Therefore, a search window of $\pm 350 \text{ km s}^{-1}$ fully encompasses the velocity extent of this virialized gas. We note that all of the absorption features in our sample that satisfy the first two criteria also satisfied by this third criterion.

This initial search included absorption from DLAs, as well as super Lyman limit systems (SLLS) with $10^{19} \text{ cm}^{-2} \leq N_{\text{HI}} \leq 10^{20.3} \text{ cm}^{-2}$. Once a DLA candidate had been identified, we performed an initial fit of the absorption profile using the XSpecGUI in `linetools` to determine if the column density satisfies the DLA threshold ($N_{\text{HI}} \geq 10^{20.3} \text{ cm}^{-2}$). We assigned a redshift to each DLA that corresponds to the velocity of the peak optical depth of the metal lines. We prioritized lines that arise from low-ionization transitions (i.e., of Si II or C II) and which are not saturated. We then refined our measurement of the H I column density using this redshift, as described below in Section 3.1. The resulting sample included 49 DLAs having $N_{\text{HI}} \geq 10^{20.3} \text{ cm}^{-2}$.

Our final step was to examine the CGM sightlines associated with each of the confirmed DLAs. To ensure precise metal-line analysis, we required the corresponding CGM sightlines to be observed at a resolution $\mathcal{R} > 4000$. If a medium- or high-resolution spectrum for the CGM was not available, then the pair was removed from our sample. The final sample used throughout this paper includes 32 DLA-CGM pairs. Table 2 lists coordinates and redshifts for the QSOs and DLAs in these pairs, as well as the instruments used for spectroscopy of each sightline.

For each DLA-CGM pair, we searched the CGM spectrum for strong $\text{Ly}\alpha$ absorption within $\pm 350 \text{ km s}^{-1}$ of the DLA redshift. We selected the single, strongest $\text{Ly}\alpha$ component that is present within this velocity window. We note that there could be multiple components of $\text{Ly}\alpha$ absorption in the CGM sightline associated with the DLA, so by choosing a single feature we set a lower limit on H I. Some CGM sightlines have multiple $\text{Ly}\alpha$ absorption components of similar strength in this range. In these cases, we included all H I absorption within $\pm 350 \text{ km s}^{-1}$ of the DLA redshift in our column density measurement.

The redshift of the CGM $\text{Ly}\alpha$ absorber was found using the same method described above: we adopted the redshift corresponding to the velocity of the peak optical depth of the metal lines in the CGM sightline. In cases where there are no securely detected CGM metal lines, we estimated the redshift using the $\text{Ly}\alpha$ absorption line. For two CGM sightlines, there was no spectral coverage of H I absorption near the redshift of the DLA. For these systems, we used the DLA redshift as the initial guess to search for associated metal lines.

¹⁸ <https://linetools.readthedocs.io/en/latest/>

Table 2
QSO Observations and DLA Sample

Background QSO ^a	Foreground QSO ^a	$z_{\text{bg}}^{\text{QSO}}$	$z_{\text{fg}}^{\text{QSO}}$	R_{\perp} kpc	Instrument ^b bg/fg	Resolution bg/fg	z_{DLA}	$N_{\text{HI}}^{\text{DLA}}$ $\log(\text{cm}^{-2})$
J 000450.91 – 084452.0	J000450.66 – 084449.6*	3.000	3.000	35	XSHOOTER/XSHOOTER	8000/8000	2.759	20.6 ± 0.2
J 023317.54 – 054230.0	J023316.29 – 054210.8*	2.959	2.629	223	MagE/MagE	5857/5857	2.375	20.45 ± 0.2
J 025049.09 – 025631.7	J025048.86 – 025640.7*	2.844	2.820	79	ESI/ESI	4545/4545	2.571	20.3 ± 0.2
J 025836.62 – 044438.5	J025837.57 – 044426.0*	2.540	2.428	163	MagE/MagE	5857/5857	1.879	20.3 ± 0.2
J 083118.50 + 424728.8*	J083121.58 + 424722.3	3.327	3.011	284	ESI/ESI	4545/4545	2.559	21.15 ± 0.2
J 093225.60 + 092500.2	J093226.34 + 092526.1*	2.602	2.410	237	MagE/XSHOOTER	5857/8000	2.252	20.45 ± 0.2
J 093959.41 + 184757.1	J093959.02 + 184801.7*	2.821	2.727	59	ESI/ESI	4545/4545	2.430	20.5 ± 0.2
J 095543.67 – 012351.5	J095544.29 – 012357.5*	2.844	2.833	89	XSHOOTER/XSHOOTER	8000/8000	2.727	20.65 ± 0.2
J 095723.43 + 622322.9	J095722.78 + 622335.2*	2.257	2.251	111	ESI/ESI	4545/4545	2.143	20.55 ± 0.2
.../BOSS	.../2100
J 102633.21 + 062909.5*	J102633.55 + 062901.5	3.120	2.890	77	MagE/MagE	5857/5857	2.564	21.05 ± 0.2
J 102633.21 + 062909.5	J102633.55 + 062901.5*	3.120	2.890	76	MagE/MagE	5857/5857	2.782	21.2 ± 0.2
J 105644.88 – 005933.4	J105645.25 – 005938.1*	2.132	2.128	62	MagE/MagE	5857/5857	1.967	20.6 ± 0.2
J 111610.69 + 411814.4	J111611.74 + 411821.5*	3.000	3.000	112	ESI/ESI	4545/4545	2.662	20.35 ± 0.2
J 114436.65 + 095904.9*	J114435.54 + 095921.7	3.146	2.974	200	MIKE-Blue/MagE	28,000/5857	2.093	20.35 ± 0.2
J 114958.49 + 430048.4*	J114958.26 + 430041.3	3.273	3.247	60	ESI/ESI	4545/4545	2.777	21.45 ± 0.2
J 115031.14 + 045353.2*	J115031.54 + 045356.8	2.521	2.517	59	MagE/MagE	5857/5857	2.000	21.3 ± 0.2
J 123635.42 + 522057.3	J123635.14 + 522059.0*	2.578	2.571	25	ESI/ESI	4545/4545	2.397	21.0 ± 0.2
J 124025.15 + 432916.5	J124024.93 + 432914.5*	3.264	3.249	24	ESI/ESI	4545/4545	2.979	21.3 ± 0.2
J 124025.15 + 432916.5	J124024.93 + 432914.5*	3.264	3.249	24	ESI/ESI	4545/4545	3.097	20.7 ± 0.2
J 142758.74 – 012136.2	J142758.89 – 012130.4*	2.352	2.279	54	MIKE-Blue/MagE	28,000/4824	1.576	21.1 ± 0.2
...	MagE/...	4824/...
J 142816.51 + 023229.2	J142815.67 + 023243.5*	3.030	3.010	155	XSHOOTER/XSHOOTER	8000/8000	2.626	21.3 ± 0.2
J 152928.37 + 231415.8*	J152929.03 + 231420.0	2.637	2.492	85	ESI/ESI	4545/4545	2.077	20.35 ± 0.2
...	BOSS/...	2100/...
J 154110.40 + 270231.2*	J154110.37 + 270224.8	3.626	3.621	48	ESI/ESI	4545/4545	3.330	20.3 ± 0.2
J 154225.81 + 173323.0	J154226.90 + 173300.5*	3.261	2.782	227	ESI/ESI	4545/4545	2.423	21.4 ± 0.2
J 161302.03 + 080814.3	J161301.69 + 080806.0*	2.386	2.386	84	MagE/MagE	4824/4824	1.617	20.5 ± 0.2
J 162737.25 + 460609.3*	J162738.63 + 460538.4	4.110	3.813	254	ESI/ESI	4545/4545	3.550	20.3 ± 0.2
J 163056.73 + 115250.3*	J163055.96 + 115229.4	3.279	3.257	184	MagE/MagE	5857/5857	3.182	20.3 ± 0.2
J 171946.66 + 254941.1	J171945.87 + 254951.2*	2.172	2.170	125	ESI/ESI	4545/4545	2.019	20.75 ± 0.2
...	GMOSN/GMOSN	1872/1872
J 172524.24 + 303801.0*	J172524.66 + 303803.9	2.647	2.634	50	ESI/ESI	4545/4545	2.508	20.35 ± 0.2
J 210329.25 + 064653.3	J210329.37 + 064650.0*	2.572	2.551	31	MagE/MagE	5857/5857	2.139	20.7 ± 0.2
J 214620.98 – 075303.8*	J214620.68 – 075250.6	2.577	2.112	120	MagE/MagE	5857/5857	1.853	20.45 ± 0.2
J 230044.52 + 015552.1*	J230044.36 + 015541.7	2.951	2.910	86	MagE/MagE	5857/5857	2.730	20.8 ± 0.2

Notes.

^a The QSO marked with a “*” is the DLA sightline.

^b Objects followed by “...” indicate cases in which a low-resolution spectrum is used.

With our DLA-CGM pairs selected, we then measured column densities of H I and column densities and kinematics of several metal ions. Section 3 describes in more detail the methodology that we use to measure the column densities for H I in both the DLA and CGM sightlines. Section 4.1 describes the methods that we use to assess metal-line column densities and kinematics. These measurements are listed in Table 3 and Table 4, respectively.

3. Determining Column Densities of H I

In this section, we describe several complementary approaches that we used to measuring the column density of neutral hydrogen present in each DLA and CGM system. Table 5 lists our CGM N_{HI} measurements and specifies which of the following constraints we used to make this assessment for each sightline pair.

3.1. Damped Ly α Profile Fitting

In all of our DLA sightlines and about a third of our CGM sightlines, we were able to estimate N_{HI} by fitting the characteristic wings of the damped line profile at $\lambda_{\text{rest}} = 1215.67 \text{ \AA}$, as in Wolfe et al. (1986). To fit the H I absorption profile, we used the interactive GUI `XFitDLA` in the Python package `pyigm` (Prochaska et al. 2017). With the redshift obtained as described in Section 2.3, we manually adjusted the N_{HI} and broadening parameters while continuously modifying the continuum level around the absorption line to achieve a close match to the data (as assessed by eye). Examples of three best-fit DLA profiles are shown in Figure 1. We adopted a ± 0.2 dex error for our N_{HI} fits, based on similar analysis from Prochaska et al. (2003a).

There is one case where the CGM H I absorber has $N_{\text{HI}} = 10^{20.2} \text{ cm}^{-2} \pm 0.2$, near the limit of a DLA itself. The sightline with the highest column density is labeled as the DLA

Table 3
Metal-line Column Density Measurements

QSO Pair ^a	z_{DLA}	$\log N_X \text{ (cm}^{-2}\text{)}$		C II	C IV	Fe II	Si II	Si IV	O I
		Al II	Al III						
J0004–0844	2.759	>13.47 <12.07	13.27 ± 0.03 <12.70	>15.05 <13.49	>14.79 >14.06	>14.92 <13.53	15.11 ± 0.13 <12.75	14.00 ± 0.05 13.08 ± 0.21	>15.23 <13.73
J0233–0542	2.375	12.89 ± 0.03 ...	12.90 ± 0.13 <12.43	>14.38 ...	13.94 ± 0.04 ...	14.26 ± 0.06 <13.58	>14.46 <13.18	13.22 ± 0.09
J0250–0256	2.571	<12.05 <12.57	>14.75 13.42 ± 0.08	13.92 ± 0.01 13.66 ± 0.03	14.02 ± 0.01 <12.51	>15.00 <13.19	13.28 ± 0.02 <12.91	>15.10 14.63 ± 0.03
J0258–0444	1.879	>13.61 <12.00	13.30 ± 0.05 <12.52	>15.04 ...	>14.69 12.77 ± 0.30	14.20 ± 0.03 <12.68	14.61 ± 0.49 <14.85	>14.30 13.09 ± 0.16	>15.40 ...
J0831+4247	2.559	>13.77 ...	13.34 ± 0.03 <12.60	>15.16 ...	14.28 ± 0.01 <13.30	>15.05 <13.53	15.80 ± 0.02 <12.89	>14.74 <12.77	>15.53 ...
J0932+0925	2.252	>13.36 ...	13.26 ± 0.05 <12.44	>15.02 ...	14.25 ± 0.01 12.78 ± 0.21	>14.66 <12.60	>15.00 <14.75	13.69 ± 0.02 ...	>15.29 ...
J0939+1848	2.430	>13.09 ...	12.70 ± 0.10 <13.02	>14.67 ...	13.79 ± 0.02 <12.94	14.27 ± 0.01 <12.54	>14.70 <14.48	13.33 ± 0.02 <12.57
J0955–0123	2.727	>14.03 12.91 ± 0.12	13.76 ± 0.04 ...	>15.56 >14.65	>14.77 >14.92	>15.27 <14.01	15.82 ± 0.04 >14.44	>14.35 >14.07	>16.45 14.45 ± 0.40
J0957+6223	2.143	12.80 ± 0.02 <12.57	12.35 ± 0.09 <13.06	>14.35 >14.61	13.84 ± 0.02 <13.87	13.85 ± 0.01 <13.07	14.78 ± 0.09 >14.56	13.42 ± 0.02 >14.42	>14.49 >14.84
J1026+0629	2.564	>13.59 ...	12.74 ± 0.22 13.47 ± 0.07	>15.09 >15.33	>14.99 14.14 ± 0.05	15.08 ± 0.20 >14.74	14.29 ± 0.02 >14.62	>15.87 >15.81
J1026+0629	2.782	>13.47 <12.06	<12.96 <12.81	>14.93 <13.31	13.71 ± 0.15 13.33 ± 0.33	14.45 ± 0.03 <13.09	>14.72 ...	13.57 ± 0.08 <13.19	>15.39 14.15 ± 0.17
J1056–0059	1.967	>13.59 12.95 ± 0.04	13.59 ± 0.02 13.32 ± 0.03	>15.13 >14.89	>15.05 ...	14.48 ± 0.02 13.28 ± 0.03	15.52 ± 0.11 13.67 ± 0.16	>14.52 >14.52	>15.35 ...
J1116+4118	2.662	12.92 ± 0.02 13.03 ± 0.03	>15.18 >14.59	13.73 ± 0.01 >14.81	14.40 ± 0.01 13.52 ± 0.02	15.08 ± 0.04 14.10 ± 0.03	13.24 ± 0.06 13.77 ± 0.03	>15.69 ...
J1144+0959	2.093	>14.08 <11.71	13.61 ± 0.00 <12.68	>15.30 13.60 ± 0.07	>14.97 13.39 ± 0.05	>14.99 <12.71	15.39 ± 0.01 <14.65	>14.44 12.72 ± 0.17
J1149+4300	2.777	>13.33 ...	12.68 ± 0.07 <11.75	>15.00 ...	>14.73 13.12 ± 0.03	>14.64 <11.75	15.27 ± 0.07 <14.24	>14.26 13.15 ± 0.01	>15.34 ...
J1150+0453	2.000	>13.60 <11.95	13.38 ± 0.04	14.31 ± 0.05 13.52 ± 0.06	>15.16 <12.87	15.81 ± 0.05 <13.25	14.06 ± 0.07 12.82 ± 0.11	>15.45 ...
J1236+5220	2.397	>13.48 <11.93	13.06 ± 0.05 <12.49	>15.02 ...	14.32 ± 0.03 13.92 ± 0.03	>14.92 <13.72	15.42 ± 0.06 <13.10	13.88 ± 0.04 13.15 ± 0.08	>15.49 ...
J1240+4329	2.979	>13.74 <12.23	13.08 ± 0.04 <12.69	>15.18 13.07 ± 0.45	14.12 ± 0.02 13.53 ± 0.14	>15.30 ...	>15.90 12.94 ± 0.59	>14.35 13.45 ± 0.05	>15.47 ...
J1240+4329	3.097	>13.59 >13.12	... 13.10 ± 0.10	>15.24 >14.77	14.07 ± 0.02 13.89 ± 0.07	>14.85 <13.69	>15.51 14.44 ± 0.04	13.50 ± 0.03 13.40 ± 0.08	>16.20 >15.24
J1427–0121	1.576	>13.67 12.72 ± 0.01	13.62 ± 0.02 12.90 ± 0.02	>15.09 ...	>14.74 >15.20	>14.85 12.98 ± 0.02	15.42 ± 0.04 14.06 ± 0.02	>14.36 >14.45	>15.52 ...
J1428+0232	2.626	>13.63 <12.43	12.72 ± 0.35 <12.82	>15.14 ...	>14.52 <13.54	>15.04 <13.43	15.66 ± 0.12 <13.04	>14.00 <13.00
J1529+2314	2.077	>13.32 ...	<12.75 13.03 ± 0.14	>15.05 ...	13.35 ± 0.11 ...	14.39 ± 0.01 <13.44	>14.82 <15.31	... >13.61	>15.38 ...
J1541+2702	3.330	>13.24 <12.14	13.69 ± 0.06 12.99 ± 0.07	>15.11 14.15 ± 0.03	>15.26 >15.00	14.01 ± 0.08 ...	14.71 ± 0.32 <13.27	>14.68 >14.30	>14.80 <13.33
J1542+1733	2.423	>13.72 <11.92	13.24 ± 0.03 <12.61	>15.18 <13.85	>14.57 13.54 ± 0.03	>15.09 <12.52	15.74 ± 0.02 <12.89	14.01 ± 0.02 13.53 ± 0.07	>15.54 ...
J1613+0808	1.617	>13.15 12.11 ± 0.09	13.07 ± 0.05 12.80 ± 0.07	>14.73 14.45 ± 0.02	... >15.14	14.44 ± 0.02 12.30 ± 0.17	15.52 ± 0.04 13.83 ± 0.10	13.62 ± 0.04 14.22 ± 0.02	>15.06 ...
J1627+4606	3.550	<12.23 <12.70	... <13.00	13.34 ± 0.04 ...	<13.55 ...	13.09 ± 0.11 <12.98	12.93 ± 0.09 <12.44	... <13.29
J1630+1152	3.182	<12.21 <12.60	>13.90 ...	13.16 ± 0.32 <13.08	13.35 ± 0.23 <13.23	<13.00 <12.77	14.79 ± 0.07 <14.29
J1719+2549	2.019	>13.71 <12.21	13.51 ± 0.02 <12.89	>15.33 ...	>15.27 <13.38	14.50 ± 0.03 <12.97	15.21 ± 0.07 <13.86	>14.45 <13.88
J1725+3038	2.508	>13.20 <12.28	12.51 ± 0.07 <12.72	>14.70 <13.64	13.46 ± 0.04 <13.26	>14.73 <12.85	14.86 ± 0.07 <13.51	13.19 ± 0.02 12.96 ± 0.22	>15.03 <14.06
J2103+0646	2.139	12.96 ± 0.05 <11.98	<12.70 <12.49	>14.81 13.72 ± 0.12	13.35 ± 0.20 13.69 ± 0.07	14.19 ± 0.06 12.63 ± 0.20	14.34 ± 0.09 ...	13.22 ± 0.05 <12.38	>15.20 14.79 ± 0.04
J2146–0753	1.853	>13.16 12.80 ± 0.02	13.32 ± 0.04 13.21 ± 0.03	>14.62 >14.53	14.38 ± 0.02 >14.93	14.31 ± 0.02 13.16 ± 0.02	15.22 ± 0.12 14.30 ± 0.05	13.78 ± 0.07 >14.30	>15.28 14.28 ± 0.07
J2300+0155	2.730	>13.44 ...	<12.91 <12.76	>14.74 <13.34	14.35 ± 0.04 13.32 ± 0.11	>14.66 <13.16	14.73 ± 0.27 <13.20	13.65 ± 0.03 <12.64	>15.21 <13.65

Note.

^a The first row of each pair lists the column densities for the DLA sightline. The second row of each pair lists the column densities for the CGM sightline.

Table 4
DLA and CGM Kinematic Measurements

QSO Pair ^a	R_{\perp} kpc	z_{DLA}	DLA				CGM			
			$\delta v_{\text{weight}}^{\text{low-ion b,c}}$ km s^{-1}	$\Delta v_{90}^{\text{low-ion b,d}}$ km s^{-1}	$\delta v_{\text{weight}}^{1548 \text{ \AA c}}$ km s^{-1}	$\Delta v_{90}^{1548 \text{ \AA d}}$ km s^{-1}	$\delta v_{\text{weight}}^{\text{low-ion b,c}}$ km s^{-1}	$\Delta v_{90}^{\text{low-ion b,d}}$ km s^{-1}	$\delta v_{\text{weight}}^{1548 \text{ \AA c}}$ km s^{-1}	$\Delta v_{90}^{1548 \text{ \AA d}}$ km s^{-1}
J0004–0844	35	2.759	5.4	118	6.7	88	–9.5	88
J0233–0542	223	2.375	–1.6	87
J0250–0256	79	2.571	5.4	111	62.7	81	243.7	81
			109.0	71
J0258–0444	163	1.879	3.7	109	1.9	175	–81.4	175
J0831+4247	284	2.559	5.1	151
J0932+0925	237	2.252	4.6	65	–136.4	65
J0939+1848	59	2.430	0.1	101
J0955–0123	89	2.727	24.6	163	–11.2	193	–36.8	193
J0957+6223	111	2.143	–8.5	101
J1026+0629	77	2.564	21.9	175	122.7	175
			355.4	197	386.6	197
J1026+0629	76	2.782	–6.7	87	–18.7	87	–76.9	87
J1056–0059	62	1.967	51.8	241	–73.3	219
J1116+4118	112	2.662	–19.6	131	4.8	230	90.3	171	68.6	230
J1144+0959	200	2.093	–16.4	139
J1149+4300	60	2.777	14.2	111	36.7	111
J1150+0453	59	2.000	12.2	109	47.9	131	–175.1	109	–152.6	131
J1236+5220	25	2.397	–7.5	111	1.4	111
J1240+4329	24	2.979	–133.4	71	–41.3	71
			143.3	81
J1240+4329	24	3.097	72.7	171	75.0	171
J1427–0121	54	1.576	12.9	82	–30.4	250	181.3	103	150.2	250
			–280.5	104	–298.4	88	972.1	58	946.8	88
J1428+0232	155	2.626
J1529+2314	85	2.077	21.7	211
J1541+2702	48	3.330	0.7	250	131.2	250
J1542+1733	227	2.423	–8.4	101	–18.9	71	–226.0	71
J1613+0808	84	1.617	–0.4	82	70.8	82	78.5	148
			275.4	104	233.1	148
J1627+4606	254	3.550
J1630+1152	184	3.182
J1719+2549	125	2.019	10.3	101
J1725+3038	50	2.508
J2103+0646	31	2.139	–9.4	109	–60.2	65	266.4	65
			156.2	109	62.3	109
J2146–0753	120	1.853	6.1	131	58.6	153	90.4	87	59.3	153
J2300+0155	86	2.730	–98.9	197	126.6	197

Notes.^a The second row below a given QSO pair describes measurements of a secondary velocity component.^b The low-ion transition used for kinematic measurements is chosen to have the highest S/N at the peak of its optical depth profile without being saturated.^c We adopt a 1σ uncertainty for δv_{weight} of 15 km s^{-1} .^d We adopt a 1σ uncertainty for Δv_{90} of 35 km s^{-1} .

sightline, and the other is treated as the CGM sightline. This choice has implications for the interpretation of our results for this pair, so we highlight it in our discussion below (and refer to it as our double-DLA system). Because there is only one such double-DLA in our sample, this does not significantly affect any of the general conclusions made later in this analysis. We indicate the specific CGM sightlines for which we constrain N_{HI} by fitting damping wings with the number 1 in the Method column in Table 5.

3.2. Lyman Limit Fitting

For some cases in which there is undamped ($N_{\text{HI}} \leq 10^{18} \text{ cm}^{-2}$) but strong Ly α absorption, we have access to either HST WFC3/UVIS or optical spectroscopic coverage of the flux blueward of the Lyman limit ($\lambda_{\text{rest}} = 912 \text{ \AA}$). For

systems with optical spectral coverage of the Lyman limit, we used the interactive GUI `XFitLLS` from the `pyigm` package to fit the Lyman limits of these systems as described in O’Meara et al. (2013). This program generates a continuum model of the QSO from Telfer et al. (2002), and allows the user to adjust the normalization and power-law tilt of the template to match the QSO continuum redward of the Ly α forest. Any sharp drops in the flux below the QSO’s 912 \AA break may then be modeled as LLSs with optical depth $\tau_{912}^{\text{LL}} \approx N_{\text{HI}}/10^{17.19} \text{ cm}^{-2}$. None of the systems in this work have a strong and clean Lyman limit feature that allows for a direct measurement of N_{HI} . This is due either to strong intervening systems absorbing the continuum close to the Lyman limit at z_{abs} or to the amount of H I in the target absorption system being sufficiently low that it does not produce a detectable Lyman limit break. Therefore, this method

Table 5
DLA and CGM Ionic Ratio Measurements

QSO Pair	R_{\perp} kpc	z_{DLA}	$\log N_{\text{SiIV}}/N_{\text{SiII}}^{\text{DLA}}$	$\{\text{SiII}/\text{HII}\}^{\text{DLA}}$	$\{\text{SiII}/\text{FeII}\}^{\text{DLA}}$	$N_{\text{HI}}^{\text{CGM}}$ $\log(\text{cm}^{-2})$	Method ^a ($N_{\text{HI}}^{\text{CGM}}$)	$\log N_{\text{SiIV}}/N_{\text{SiII}}^{\text{CGM}}$	$\{\text{SiII}/\text{HII}\}^{\text{CGM}}$	$\{\text{OI}/\text{HII}\}^{\text{CGM}}$	$\{\text{OI}/\text{FeII}\}^{\text{CGM}}$
J0004–0844	35	2.759	-1.12 ± 0.14	-1.00 ± 0.24	<0.13	19.10 ± 0.20	1	>0.33	<-1.86	<-2.06	...
J0233–0542	223	2.375	<-1.23	>-1.50	>0.13	13.93 ± 0.10	5	...	<3.74
J0250–0256	79	2.571	<-1.71	>-0.81	>0.91	15.64 ± 0.56	3	...	<2.03	2.30 ± 0.56	>0.87
J0258–0444	163	1.879	>-0.30	-1.20 ± 0.53	0.34 ± 0.50	16.29 ± 1.71	4	>-1.76	<3.05
J0831+4247	284	2.559	>-1.06	-0.86 ± 0.20	<0.69	14.33 ± 0.31	5	...	<3.04
J0932+0925	237	2.252	<-1.31	>-0.96	...	13.98 ± 0.13	5	...	<5.26
J0939+1848	59	2.430	<-1.37	>-1.31	>0.37	16.22 ± 1.78	4	...	<2.75
J0955–0123	89	2.727	>-1.47	-0.34 ± 0.20	<0.49	19.20 ± 0.20	1	...	>-0.27	-1.44 ± 0.45	>-0.80
J0957+6223	111	2.143	-1.36 ± 0.09	-1.28 ± 0.22	0.87 ± 0.09	>0.53
J1026+0629	77	2.564	-0.79 ± 0.20	-1.48 ± 0.28	<0.02	20.20 ± 0.20	1	...	>-0.97	>-1.08	>0.43
J1026+0629	76	2.782	<-1.15	>-1.99	>0.22	15.51 ± 0.79	3	1.95 ± 0.81	>-0.18
J1056–0059	62	1.967	>-1.00	-0.59 ± 0.23	0.98 ± 0.11	19.60 ± 0.20	1	>0.85	-1.44 ± 0.26
J1116+4118	112	2.662	-1.84 ± 0.07	-0.78 ± 0.20	0.62 ± 0.04	20.10 ± 0.20	1	-0.33 ± 0.04	-1.51 ± 0.20
J1144+0959	200	2.093	>-0.94	-0.47 ± 0.20	<0.34	16.40 ± 1.60	4	>-1.93	<2.75
J1149+4300	60	2.777	>-1.01	-1.69 ± 0.21	<0.57	15.27 ± 0.68	3	>-1.09	<3.45
J1150+0453	59	2.000	-1.75 ± 0.09	-1.00 ± 0.21	<0.59	16.37 ± 1.63	4	>-0.43	<1.37
J1236+5220	25	2.397	-1.53 ± 0.07	-1.09 ± 0.21	<0.44	17.20 ± 0.20	2	>0.05	<0.39
J1240+4329	24	2.979	...	>-0.91	...	16.58 ± 1.42	4	0.51 ± 0.60	0.85 ± 1.54
J1240+4329	24	3.097	<-2.01	>-0.70	...	20.00 ± 0.20	1	-1.04 ± 0.09	-1.07 ± 0.20	>-1.45	>0.30
J1427–0121	54	1.576	>-1.06	-1.19 ± 0.20	<0.51	19.70 ± 0.20	1	>0.39	-1.15 ± 0.20
J1428+0232	155	2.626	>-1.66	-1.15 ± 0.23	<0.56	15.46 ± 0.94	3	...	<2.07
J1529+2314	85	2.077	...	>-1.04	>0.37	>-1.70
J1541+2702	48	3.330	>-0.03	-1.10 ± 0.38	0.63 ± 0.33	16.14 ± 0.51	3	>1.03	<1.63	<0.50	...
J1542+1733	227	2.423	-1.73 ± 0.03	-1.17 ± 0.20	<0.59	14.13 ± 1.37	5	>0.64	<3.25
J1613+0808	84	1.617	-1.90 ± 0.06	-0.49 ± 0.20	1.03 ± 0.05	16.70 ± 0.10	2	0.40 ± 0.10	1.62 ± 0.14
J1627+4606	254	3.550	-0.17 ± 0.14	-2.72 ± 0.23	>-0.52	14.36 ± 0.03	5	...	<3.10	<2.24	...
J1630+1152	184	3.182	<-0.35	-2.46 ± 0.31	...	15.53 ± 1.02	3	...	<2.19	<2.07	...
J1719+2549	125	2.019	>-0.77	-1.05 ± 0.21	0.65 ± 0.08	16.19 ± 1.81	4	...	<2.17
J1725+3038	50	2.508	-1.68 ± 0.07	-1.00 ± 0.21	<0.07	15.92 ± 1.08	3	>-0.54	<2.07	<1.45	...
J2103+0646	31	2.139	-1.12 ± 0.11	-1.87 ± 0.22	0.09 ± 0.11	18.85 ± 0.20	1	-0.75 ± 0.20	0.93 ± 0.20
J2146–0753	120	1.853	-1.44 ± 0.14	-0.74 ± 0.23	0.85 ± 0.12	18.55 ± 0.20	1	>0.01	0.24 ± 0.21	-0.96 ± 0.21	-0.11 ± 0.08
J2300+0155	86	2.730	-1.07 ± 0.27	-1.58 ± 0.34	<0.00	14.97 ± 0.48	3	...	<2.72	<1.98	...

Note.

^a Method used to constrain N_{HI} for the CGM sightline as described in Section 3. Method 1 corresponds to the fitting of $\text{Ly}\alpha$ damping wings; method 2 corresponds to model fitting of the Lyman limit covered in HST WFC3/UVIS grism spectroscopy; method 3 corresponds to model fitting of the Lyman limit covered in our optical medium-/high-resolution spectra; method 4 corresponds to bounds established from apparent optical depth measurements for saturated but undamped $\text{Ly}\alpha$ absorption; and method 5 corresponds to summed apparent optical depth column densities assuming that $\text{Ly}\alpha$ is optically thin.

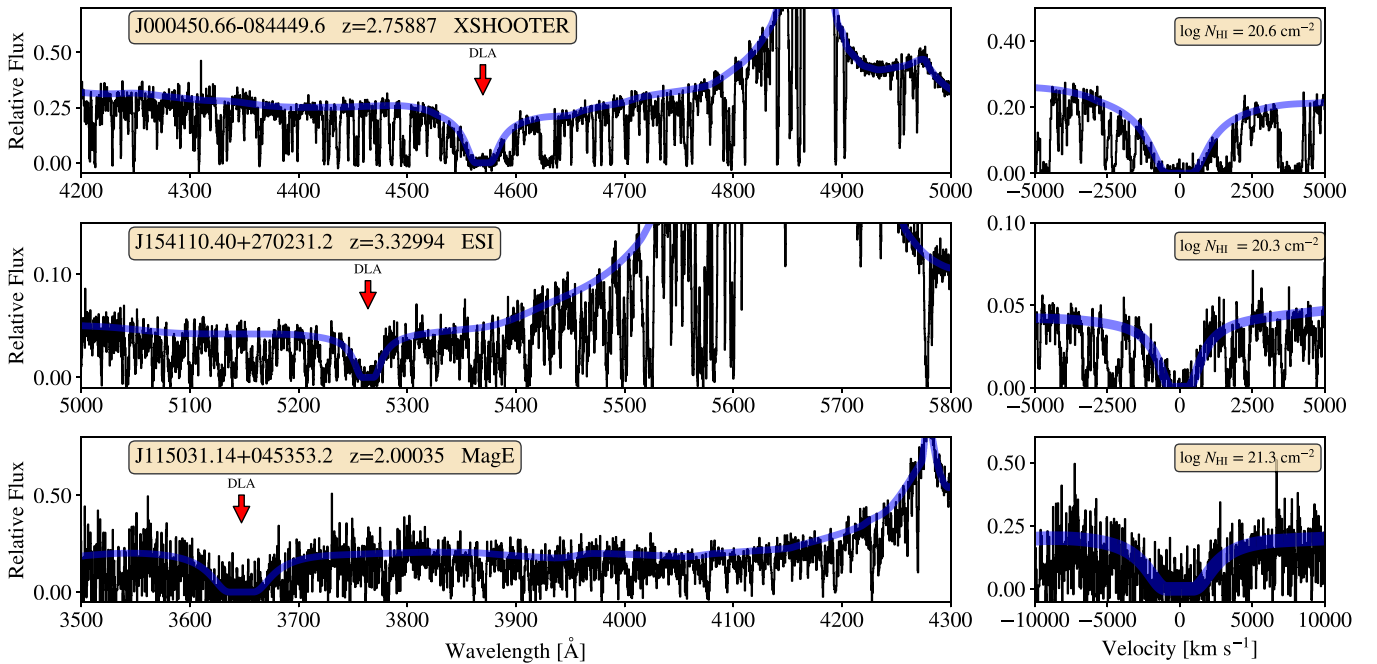


Figure 1. Examples of QSO spectroscopy, the DLA profiles, and continuum fits (in blue) for three systems. The QSO identifier, DLA redshift, and the instrument are given in the top left-hand corner of each panel. The DLA fit and adopted N_{HI} value are shown the subsequent right-hand panels. Velocities are measured relative to the DLA redshift.

allowed for an estimate of the minimum amount of HI that is required to account for the decrease in flux blueward of the DLA’s Lyman limit. Moreover, because there is strong but undamped $\text{Ly}\alpha$ absorption associated with these systems, we also placed an upper bound on their HI columns of $N_{\text{HI}} < 10^{18} \text{ cm}^{-2}$.

Seven CGM sightlines in this work were targeted in the HST WFC3/UVIS grism survey of paired quasars described in Lusso et al. (2018);¹⁹ however, four of our absorbers were not detected in the HST spectroscopy. This was due to either weak HI absorption that limits the detection of a flux decrement at $\lambda_{\text{rest}} = 912 \text{ Å}$, or to strong background absorbers that significantly reduce the flux and signal-to-noise ratio (S/N) near the Lyman limit. One CGM system, in sightline J105644.88-005933.4, has $\text{Ly}\alpha$ damping wings observed in our optical spectroscopy, so we do not make use of the HST coverage to improve our constraints on N_{HI} . The Lyman limit coverage of the grism spectroscopy of the remaining two systems (in sightlines J161302.03+080814.3 at $z_{\text{abs}} = 1.617$ and J123635.42+522057.3 at $z_{\text{abs}} = 2.39691$) were modeled in the same manner as described above, using the XFitLLS GUI. Uncertainties in the value of N_{HI} were determined by perturbing the best-fit value in increments of 0.1 dex and assessing the degree to which each perturbed value was consistent with the data by visual inspection. Using this method, we estimated the error on each N_{HI} measured from these grism spectra to be $\pm 0.2\text{--}0.3$ dex. The CGM sightlines for which we found N_{HI} from Lyman Limit fitting of either HST WFC3/UVIS or optical spectroscopy are indicated with the number 2 or 3 in the Method column in Table 5, respectively.

3.3. Limits on N_{HI}

CGM sightlines for which N_{HI} could not be constrained using the methods described above (but must have $N_{\text{HI}} < 10^{18} \text{ cm}^{-2}$) were treated in one of two ways (described below). If there is a single strong absorption line with some associated metal absorption at the same redshift, then we estimated a lower limit on the HI column density using the apparent optical depth method (Savage & Sembach 1991). We adopted this limit, along with the upper limit $N_{\text{HI}} < 10^{18} \text{ cm}^{-2}$, as conservative bounds on the HI column density assuming that the $\text{Ly}\alpha$ transitions are within the flat region of the curve of growth. We indicate these CGM sightlines with the number 4 in the Method column in Table 5. For CGM systems that have many weak absorption features near $\text{Ly}\alpha$ with no associated metal lines, we assumed that the $\text{Ly}\alpha$ is optically thin and calculated the column density of each absorption feature within $\pm 350 \text{ km s}^{-1}$ of the DLA redshift using the same apparent optical depth method. We then summed the resulting column densities for these features to use as our final estimate of N_{HI} . These sightlines are designated with the number 5 in the Method column in Table 5. Lastly, two CGM sightlines had no spectral coverage of $\text{Ly}\alpha$, and therefore are not used in any HI analysis.

4. Metal-line Profile Analysis

Each metal line in our DLA and CGM sightlines was visually inspected using the interactive GUI XAbsSystem-GUI in the package linetools, which displays multiple transitions for a simultaneous comparison. For all strong transitions in each spectrum, we manually set velocity limits over which we measure the associated absorption line by searching within $\pm 1000 \text{ km s}^{-1}$ of the absorber redshift. This search window was adopted based on the findings of Rudie et al. (2019b), who identified metal-line absorption associated with LBG hosts at relative velocities of up to $\pm 1000 \text{ km s}^{-1}$.

¹⁹ The HST WFC3/UVIS data presented in this paper can be found in Mikulski Archive for Space Telescopes (MAST) at the Space Telescope Science Institute. The specific observations analyzed can be accessed via [10.17909/n7qq-vc30](https://archive.stsci.edu/10.17909/n7qq-vc30).

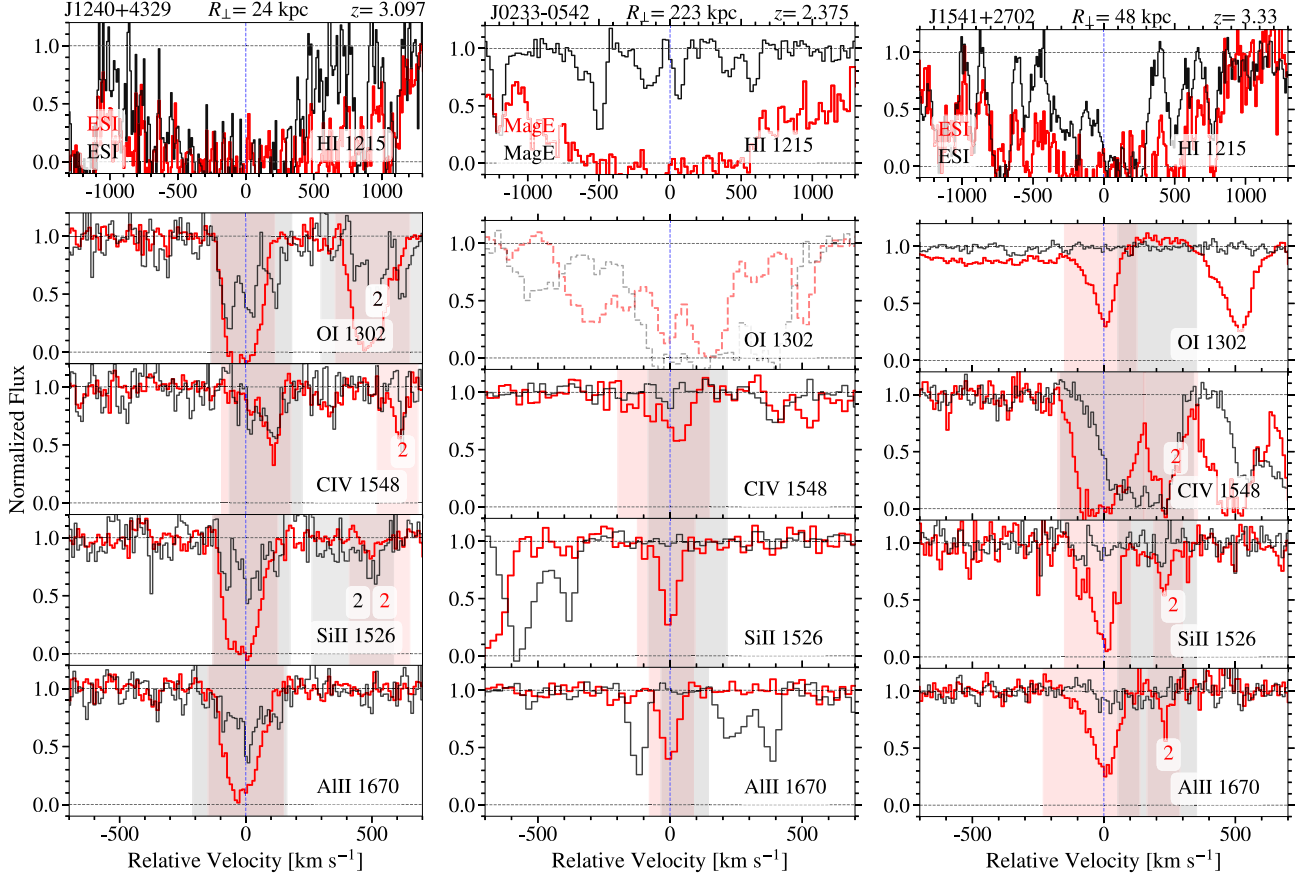


Figure 2. H I and metal-line absorption profiles for three DLA-CGM sightline pairs. H I, O I, C IV, Si II, and Al II transitions are shown (as indicated at the bottom right-hand of each panel). The black histograms show CGM sightlines and the red histograms show the corresponding DLA sightlines. The QSO pair ID is shown at the top of each column, along with the projected distance between the sightlines at the redshift of the host DLA. The blue-dotted line indicates the systemic velocity of the DLA. The shaded areas indicate the placement of the velocity windows used to measure metal-line absorption strength. In the case of profiles with multiple components, the shaded regions are marked “2” or “3” to indicate a second or third component. The instrument used for each sightline is labeled in the bottom left-hand corner of each H I panel. Transitions that are not used in this analysis due to extreme blending are shown with dotted histograms.

We assigned velocity limits for each velocity component in every sightline. In cases in which no absorption is clearly evident, we adopted velocity limits of $\pm 300 \text{ km s}^{-1}$ by default and adjusted the edges of this window to avoid absorption from unassociated systems. In many cases the blending between components associated with our target system is severe, such that they cannot be separated into two distinct absorption lines. In such cases, we separated components only if the flux rises to $>50\%$ of the continuum level between the lines. If a given transition is severely blended with an unassociated absorber such that it could not be separated, then we excluded the line from our analysis. Representative examples of our chosen absorption-line windows, including systems with multiple components, for three quasar pairs are shown in Figure 2.

4.1. Column Densities of Metal Lines

Once the velocity windows had been selected, we used the apparent optical depth method as detailed in Savage & Sembach (1991) to measure the column densities. The optical depth per unit velocity (v) is defined as:

$$\tau(v) = \ln I(v)/I_c = \ln \frac{1}{F(v)}, \quad (1)$$

where I_c is the intensity of the continuum within the set velocity window, $I(v)$ is the absorbed intensity within that

window, and $F(v)$ is the continuum-normalized absorbed flux. Savage & Sembach (1991) used the optical depth to find the total column density, N , as:

$$N = \frac{m_e c}{\pi e^2 f \lambda_0} \sum_i \ln F(v)_i^{-1} \Delta v, \quad (2)$$

where m_e is the mass of an electron, c is the speed of light, e is the elementary charge, f is the oscillator strength of the transition, λ_0 is the rest wavelength of the transition, and Δv is the step in velocity space for each pixel (i) within the velocity window. The error (σ_N) is thus defined as:

$$\sigma_N^2 = \sum_i \left(\frac{m_e c}{\pi e^2 f \lambda_0} \frac{\sigma_{F(v)_i}}{F(v)_i} \Delta v \right)^2, \quad (3)$$

where $\sigma_{F(v)}$ is the error in the spectral flux.

For our high-resolution spectroscopy (obtained with MIKE), if the absorbed line profile drops below 20% of the flux uncertainty, or if the relevant spectral pixels reach a normalized flux level of <0.05 , then the line was flagged as saturated and our column density estimate was treated as a lower limit. Line saturation is a larger concern for medium-resolution spectroscopy (e.g., from ESI or MagE). For these sightlines, we conservatively defined a line as saturated if the absorbed line profile drops below 50% of the continuum flux

(see Prochaska et al. 2003b). Three- σ_N upper limits are used for non-detections (defined as $N < 3\sigma_N$).

We also investigated the systematic error that is associated with this measurement due to uncertainty in the placement of our velocity windows. We measured the column density of all of the single-component systems after broadening the velocity limits by 50 km s^{-1} on both sides of the line profile. In the case of Si II, we find that this increases our measured column densities by an average of 0.07 dex with a scatter of 0.05 dex. Thus, the choice of a broader velocity window would systematically increase our column density measurements. However, given that the FWHM velocity resolution of our data set is $\lesssim 66 \text{ km s}^{-1}$, our uncertainty in the velocity limit of our absorption profiles is not likely to exceed $\pm 50 \text{ km s}^{-1}$. The implied systematic error in our column densities is therefore $\lesssim 0.1$ dex.

We combined multiple column density constraints for each species as follows: (1) if there is one transition that has yielded a direct measurement of the column density, then that measurement is adopted; (2) if there is more than one detection, then we adopted the mean N value; (3) if there are no direct measurements and one or more transitions are saturated, then the adopted column density is the highest value flagged as a lower limit; and (4) if all of the transitions have yielded upper limits on the column density, then we adopted the lowest upper limit.

Finally, we summed the column densities measured from each separate velocity component associated with a given ion. While we include separate components in this final summation, the velocity components that are kinematically consistent with the primary H I absorber have the largest columns along the line of sight, and therefore dominate these measurements.

4.2. Metal-line Kinematics

We also assessed the kinematic properties of the metal lines in the DLA and CGM sightlines. Our spectral coverage includes singly-, doubly-, and triply-ionized transitions. We focused on the kinematics of singly- and triply-ionized transitions only. We made two kinematic measurements: the flux-weighted velocity centroid (δv_{weight}) and the Δv_{90} velocity width. To estimate the former, we first calculated the flux-weighted wavelength centroid, defined as follows:

$$\lambda_{\text{weight}} = \frac{\sum_i (1 - F_i(\lambda_i)) \lambda_i}{\sum_i (1 - F_i(\lambda_i))}, \quad (4)$$

where $F_i(\lambda_i)$ is the continuum-normalized flux and λ_i is the wavelength at each pixel i within the velocity window for the line. The final δv_{weight} was then calculated using this wavelength relative to the redshift of the associated DLA.

We rely upon δv_{weight} as opposed to the velocity at the peak optical depth (δv_{peak}) for several reasons. First, using δv_{peak} would bias the low-ion kinematics toward $\sim 0 \text{ km s}^{-1}$ because they were used to estimate the redshift of the individual absorption systems (see Section 2.3). Second, many lines are not symmetric about the peak optical depth, such that δv_{peak} probes the velocity of the strongest absorption rather than the average velocity of the absorbing gas. Furthermore, for saturated lines, the velocity at which the optical depth peaks is ambiguous. For most sightlines, the difference in the Si II $\lambda 1526$ δv_{peak} versus δv_{weight} is $\lesssim 50 \text{ km s}^{-1}$, and we find that sightlines for which there is a greater than 50 km s^{-1} difference

in these measures have large widths and strongly asymmetric profiles.

We conducted our δv_{weight} measurements for an unsaturated high-S/N low-ion transition, as well as for C IV $\lambda 1548$. C IV $\lambda 1548$ was selected as representative of the velocity profile of high-ionization material due to its high oscillator strength. We do not report high-ion kinematics for sightlines in which either transition in the C IV doublet is not securely detected or is heavily blended. For sightlines with multiple velocity components, we measured each component's λ_{weight} separately using Equation (4) and computed the corresponding δv_{weight} relative to the redshift of the corresponding DLA. We were able to assess the δv_{weight} of low-ionization material (and C IV-absorbing material) in 31 (30) of our DLA sightlines, seven (six) of which have resolved secondary velocity components. The CGM sightlines have fewer securely detected metal lines, reducing the number of sightlines we could use for kinematic analysis. We measured the δv_{weight} of low-ionization material in eight sightlines, three of which have secondary velocity components. We measured the δv_{weight} of C IV in 20 CGM sightlines, six of which have secondary components.

Our second kinematic measurement is the Δv_{90} velocity width, which was introduced in Prochaska & Wolfe (1997) as a tracer for kinematics of the neutral gas content of DLAs. In that work, the authors analyzed the full absorption profiles of unsaturated low-ionization transitions to assess the bulk neutral gas velocity dispersion, and to ensure that the velocity width is not overestimated due to weak outlying velocity components. We take the same approach for each of our DLA and CGM sightlines. In addition, we assess Δv_{90} on a component-by-component basis for both the low-ionization material in each system and for each C IV $\lambda 1548$ profile (chosen for its high oscillator strength). We make use of these latter (component-by-component) measurements when comparing the kinematics of our sightline pairs in Section 5.5, and report these values in Table 4. We make use of the former Δv_{90} values (measured without component separation) when comparing our sample to global relations in the literature in Section 6.2 and Figure 16.

Prochaska et al. (2008) investigated the artificial broadening associated with Δv_{90} measured from medium-resolution spectra. In that work, they reduced their measured ESI Δv_{90} widths by 20 km s^{-1} and adopted an uncertainty of 20 km s^{-1} . We assume that the artificial broadening of Δv_{90} in our medium-resolution spectra is proportional to what is measured in Prochaska et al. (2008), e.g., a FWHM resolution of 45 km s^{-1} would broaden Δv_{90} by $\sim 20 \text{ km s}^{-1}$, a factor of 0.44 times the FWHM resolution. Using this factor ($0.44 \times \text{FWHM}$ resolution), we estimated the artificial broadening of Δv_{90} in all of the spectra used herein. The measured Δv_{90} widths were then reduced by that estimate to produce the final reported Δv_{90} widths used in the following analysis.

4.2.1. Uncertainties in Kinematic Measurements

The precision of our kinematic measurements depends on the FWHM resolution and S/N of our spectra. To assess the level of uncertainty in our measurements of δv_{weight} and Δv_{90} , we performed a Monte Carlo analysis on mock C IV lines. To be conservative, we use the lowest spectral resolution and S/N among all of our observed sightlines for this analysis, and adopt the resulting uncertainties across our sample.

We first created a mock spectrum with a velocity resolution consistent with that of our ESI data (FWHM $\sim 66 \text{ km s}^{-1}$).

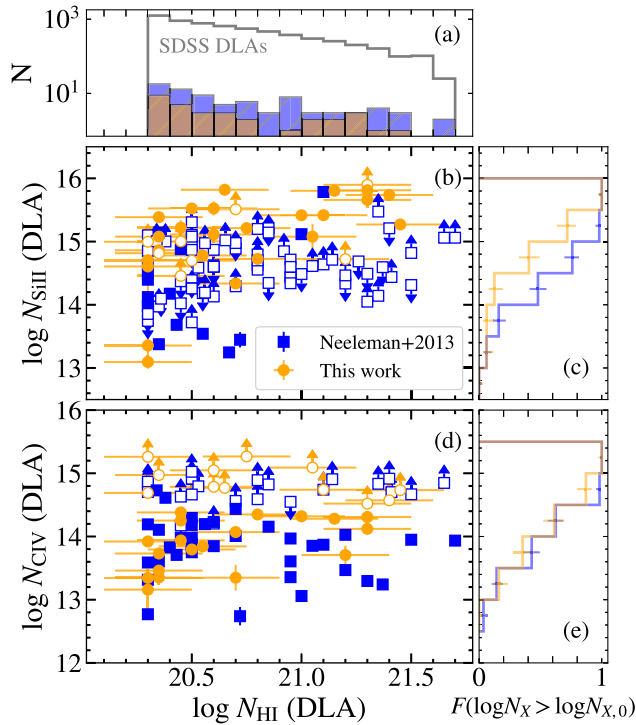


Figure 3. (a) Distribution of N_{HI} for our DLA sample (orange) and the subset of the Neeleman et al. (2013) sample having $z < 3.6$ (blue). The distribution of N_{HI} for DLAs with redshifts $1.6 < z < 3.5$ discovered in SDSS-III DR9 QSO spectra (Noterdaeme et al. 2012) is shown in gray. Bins have a width of $\Delta \log N_{\text{HI}} = 0.1$. All three samples have an average $\log N_{\text{HI}}/\text{cm}^{-2}$ of ≈ 20.7 . (b) N_{SiII} vs. N_{HI} for our sample and that of Neeleman et al. (2013). Open symbols represent limits. (d) N_{CIV} vs. N_{HI} for our sample and that of Neeleman et al. (2013). Panels (c) and (e) show normalized cumulative distributions of N_{SiII} and N_{CIV} , respectively, with a bin width of $\Delta \log N_{\text{SiII,CIV}} = 0.5$. The 1σ Wilson score intervals for these bins are shown as error bars. We include only detections and saturations in these distributions and exclude upper limits. The distributions are similar, which suggests that our sample is representative of typical DLAs around $z \sim 2.5$.

We then added a single, fake CIV line with a column density equal to the minimum column density detection ($\log N_{\text{CIV}} = 13.1 \text{ cm}^{-2}$) in our ESI data set. We adopted the mean Doppler width measured for CIV by (Rudie et al. 2019b), 12.4 km s^{-1} . Because our absorption features likely include unresolved velocity components, we also created mock spectra with two to three of these absorbers at a maximum velocity separation of $\sim 100 \text{ km s}^{-1}$. Finally, we added Gaussian random noise to the mock spectra. We generated 100 realizations of each mock spectrum with a S/N equal to the lowest S/N measured in our observed spectra ($\text{S/N} \sim 20 \text{ pixel}^{-1}$). The standard deviation of the δv_{weight} measurements for our one-component, two-component, and three-component profiles are 8 km s^{-1} , 10 km s^{-1} , and 15 km s^{-1} , respectively. The corresponding values of the dispersion in our Δv_{90} measurements are 10 km s^{-1} , 20 km s^{-1} , and 35 km s^{-1} . We adopted the largest of these values as our 1σ measurement uncertainty for δv_{weight} and Δv_{90} for all sightlines, regardless of their S/N or spectral resolution.

5. Results

5.1. Our DLA Sample as a Representative DLA Population

To better understand whether our DLA sample is representative of random populations of DLAs in this redshift range, we

compare its properties to those of a larger DLA sample from literature. Neeleman et al. (2013) analyzed 100 DLAs observed at high resolution ($\mathcal{R} \sim 40,000$) with $z_{\text{abs}} \sim 1.5\text{--}4$. With these high-quality data, Neeleman et al. (2013) were able to measure precise metal column densities. We restrict our comparison to a subset of the Neeleman et al. (2013) sample that includes 72 DLAs with $z_{\text{abs}} < 3.6$ (i.e., the highest redshift in our DLA sample). Our sample ranges from $z_{\text{abs}} \sim 1.6\text{--}3.5$, with an average absorber redshift of $\langle z_{\text{abs}} \rangle \approx 2.5$. The Neeleman et al. (2013) subset has relatively more systems above $z_{\text{abs}} > 2.9$, yielding an average $\langle z_{\text{abs}} \rangle \approx 2.6$. We perform a two-sample Kolmogorov–Smirnov (K-S) test on the two distributions to test the null hypothesis that the two samples are drawn from the same parent distribution. The maximum absolute difference between the distributions calculated from the two-sample K-S statistic is low ($D_{\text{K-S}} = 0.27$) and has a P -value of 0.06, which suggests that we cannot reject the null hypothesis at a $>95\%$ confidence level. The standard deviation of redshifts for our DLA sample is 0.47, similar to the standard deviation of the Neeleman et al. (2013) subset (0.48). These comparisons suggest that the redshift distributions of these two samples are similar.

We now consider how the physical properties of our sample DLAs relate to those of the parent DLA population during this epoch by comparing their distributions of N_{HI} , N_{SiII} , and N_{CIV} . The latter two ions were selected to be representative of low-ionization and high-ionization metal absorption. We first compare the N_{HI} distribution of our sample to that of a much larger sample of 6132 DLAs with redshifts $1.6 < z < 3.5$ discovered in QSO spectroscopy from the SDSS-III DR9 (Noterdaeme et al. 2012), as well as to that of Neeleman et al. (2013). These comparisons are shown in Figure 3(a). A two-sample K-S test comparing the N_{HI} distribution of our sample and that of Noterdaeme et al. (2012) yields a maximum difference value of $D_{\text{K-S}} = 0.16$ with a P -value of 0.33. The same test comparing our sample and that of Neeleman et al. (2013) yields a maximum difference value of $D_{\text{K-S}} = 0.1$ with a P -value of 0.96. Thus, we find no evidence that either of the two N_{HI} distribution pairs are drawn from different parent populations.

We also compare our sample distributions of N_{SiII} and N_{CIV} with those of Neeleman et al. (2013). N_{HI} and N_{SiII} are similarly correlated in both samples, though we note that the DLAs with the highest N_{HI} values in our sample appear to have higher values of N_{SiII} (shown in Figure 3(b)). These values are, however, consistent with the lower limits on N_{SiII} in the Neeleman et al. (2013) data. Our measurements of Si II column density are somewhat more sensitive than those of Neeleman et al. (2013) because measurements from the latter study relied on Si II $\lambda 1546$ (which is typically saturated in DLA sightlines), whereas we make use of the weaker Si II $\lambda 1808$ when calculating Si II column density. The more highly ionized material, traced by CIV, does not have column densities that are strongly correlated with N_{HI} . However, the distributions of N_{CIV} are similar between both samples (Figure 3(d)). We conclude that the DLAs in our QSO pair sample are representative of typical DLAs at redshifts $1.6 \lesssim z \lesssim 3.5$ from the point of view of column density distributions.

5.2. N_{HI} in DLA Environments

Since DLAs are the dominant reservoirs of neutral gas at $z < 5$, the environments of DLAs can elucidate how HI gas is

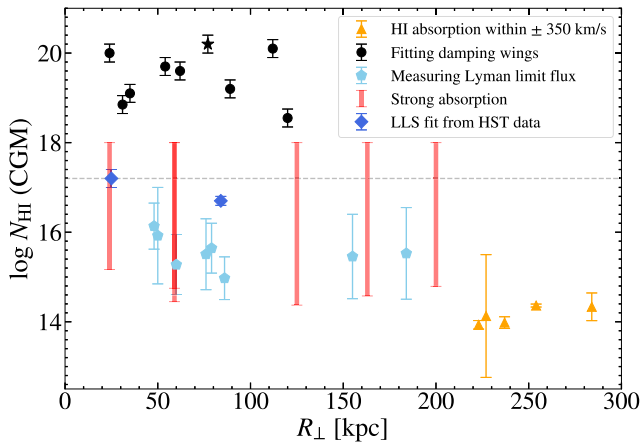


Figure 4. N_{HI} measurements in our CGM sightlines vs. R_{\perp} . Black circles represent sightlines for which we constrain N_{HI} by fitting the damping wings of the Ly α absorption profile. The black star indicates a double-DLA system. Dark blue diamonds indicate N_{HI} measurements obtained from analysis of the Lyman limit observed with HST/WFC3 grism spectroscopy. Light-blue pentagons represent sightlines for which we place an upper limit on N_{HI} by fitting the flux blueward of the system’s Lyman limit. Red bars indicate systems with strong undamped Ly α absorption for which we lack coverage of the Lyman limit. Orange triangles represent systems with no single strong Ly α absorption line. The dashed line shows the limit for optically thick gas, $N_{\text{HI}} \geq 10^{17.2} \text{ cm}^{-2}$. We find that $\log N_{\text{HI}}$ is anticorrelated with R_{\perp} overall, but exhibits significant scatter at $R_{\perp} < 120$ kpc.

distributed in the Universe. We first investigate the absorption strength of H I as a function of distance from our host DLAs. Figure 4 shows $\log N_{\text{HI}}$ measured in each CGM sightline versus transverse distance (R_{\perp}) from the corresponding DLA. All optically thick CGM systems (including those with damped Ly α absorption, indicated in black, as well as one sightline shown as the blue diamond at $R_{\perp} = 25$ kpc) are located within $R_{\perp} < 120$ kpc. However, Figure 4 also includes numerous sightlines within 120 kpc of DLAs that are optically thin, which suggests that neutral gas near DLAs exhibits a wide range of densities. Weak absorption (with $N_{\text{HI}} \sim 10^{14} \text{ cm}^{-2}$, indicated in orange) is only found further than 200 kpc from DLAs, indicating that H I column densities may decrease with increasing R_{\perp} . To evaluate the significance of an anticorrelation between R_{\perp} and H I column density, we calculate the Kendall rank correlation coefficient (τ_K). We caution that six H I absorbers (indicated by the red bars) are on the flat part of the curve of growth, with large error bars that are not accounted for in this calculation. Nevertheless, we find $\tau_K = -0.4$ with a two-sided probability of no correlation of $P = 0.002$, which is indicative of an anticorrelation. Such anticorrelations between H I column density and projected distance are ubiquitous features of CGM sightline samples, including those probing LBG environments (Rakic et al. 2012; Rudie et al. 2012), QSO host environments (Prochaska et al. 2013), and the environments around Ly α emitters (Liang et al. 2021) at $z \sim 2-3$.

We also explore the spatial extent of optically thick ($N_{\text{HI}} \geq 10^{17.2} \text{ cm}^{-2}$) gas in DLA halos, calculating covering fractions within $R_{\perp} < 100$ kpc, at $100 \text{ kpc} < R_{\perp} < 200$ kpc, and at $R_{\perp} > 200$ kpc. The error associated with these covering fractions is estimated by calculating the 1σ Wilson score interval for each bin. The systems marked by red bars in Figure 4, with strong but undamped absorption, have column density constraints that are ambiguous with respect to the optically thick threshold, and are therefore excluded from these

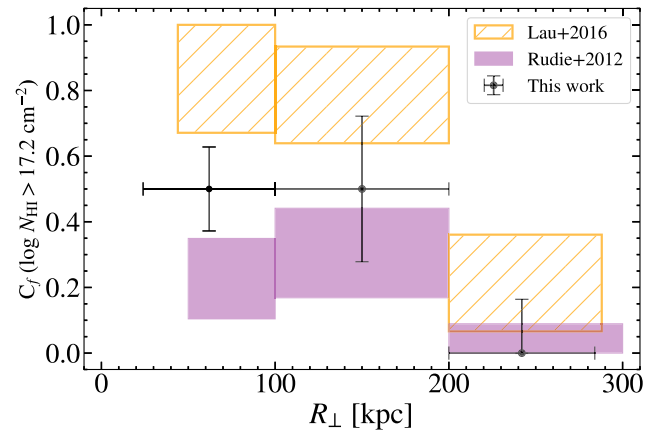


Figure 5. Covering fractions of optically thick H I ($N_{\text{HI}} > 10^{17.2} \text{ cm}^{-2}$) measured in bins of $\Delta R_{\perp} \approx 100$ kpc. The black points represent covering fractions in the CGM of DLAs (this work). The orange boxes represent the covering fraction of optically thick H I measured in the CGM of quasar hosts (Lau et al. 2016). Purple boxes represent the covering fraction of optically thick H I in the CGM of LBGs (Rudie et al. 2012). The vertical error bars represent the 1σ Wilson score interval for each bin. The x-axis error bars indicate the span of R_{\perp} within each bin. We find that DLA halos exhibit more than twice the covering fraction of optically thick H I relative to that measured in LBG halos within $R_{\perp} < 100$ kpc.

calculations. Similarly, we exclude the one blue point which lies at the limit of optically thick gas. The resulting covering fractions (C_f) are shown in Figure 5, with the x-axis error bars indicating the R_{\perp} span of each bin. The halos of DLAs have an incidence of optically thick H I of $50\% \pm 13\%$ and $50\% \pm 22\%$ for bins with $24 \text{ kpc} < R_{\perp} \leq 100 \text{ kpc}$ and $100 \text{ kpc} < R_{\perp} \leq 200 \text{ kpc}$, respectively. At $200 \text{ kpc} < R_{\perp} \leq 284 \text{ kpc}$, we place an upper limit on the covering fraction of $< 16\%$ for optically thick H I. By combining measurements within the first two bins, we find the incidence of optically thick H I to be $50\% \pm 11\%$ within $R_{\perp} < 200$ kpc of DLAs. These findings suggest that the extent of optically thick gas around DLAs is $\lesssim 200$ kpc. Moreover, our finding that DLAs are very rarely detected along both sightlines in our QSO pairs suggests that the total extent of DLAs themselves is likely $\lesssim 50$ kpc (see S. M. Urbano Stawinski et al. 2023, in preparation).

We compare these measurements to the results of two other surveys at similar redshifts: one focused on massive quasar host galaxy halos (Lau et al. 2016), and the other assessed halos of Lyman Break Galaxies (LBGs; Rudie et al. 2012). Within 100 kpc, the latter survey implies covering fractions of optically thick material around LBGs of $20^{+15}_{-13}\%$. The C_f we measure around DLAs, $50\% \pm 13\%$, is $\approx 1.5\sigma$ larger than this value. While this offset is not statistically significant, it is nevertheless suggestive that DLA halos may have more uniformly distributed optically thick H I than LBG halos (i.e., it is more likely to find optically thick gas near a DLA than near an LBG). Moreover, this finding may be a natural result of our selection criteria for CGM sightlines, i.e., they must arise close to a region that is already known to have a high neutral column density. Meanwhile, quasar halos have nearly 100% optically thick covering fractions within 100 kpc, $> 1\sigma$ larger than DLA halos. It is therefore even more likely that optically thick material will be found close to quasar host galaxies. Beyond 100 kpc, the error bars on these C_f constraints overlap, such that the covering fractions measured around these three samples are statistically consistent. Previous work has demonstrated that DLAs are clustered to LBGs, with the DLA-LBG correlation

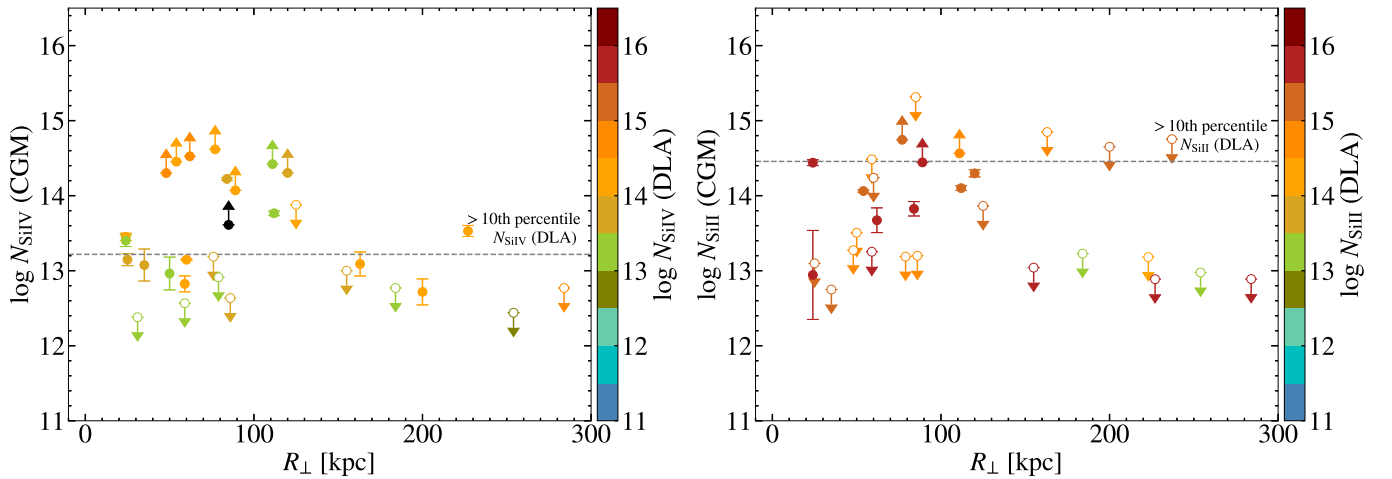


Figure 6. Column densities of Si IV (left-hand panel) and Si II (right-hand panel) measured in our CGM sightlines. Measurements are color-coded according to the corresponding DLA column density for that ion. The black point represents a system for which the DLA sightline has an ambiguous column density. Open symbols indicate that our constraint on the CGM sightline column is an upper limit. The horizontal-dashed lines represent the threshold above which 90% of metal-line column densities for the DLA sightlines fall. The column densities of Si II in the CGM drop below what we measure for the vast majority of DLA sightlines at small R_{\perp} . In contrast, the column densities of Si IV remain similar to those measured in our DLAs up to projected distances of ~ 150 kpc.

length being statistically consistent with that of the LBG-LBG autocorrelation length ($r_0 = 2.81^{+1.4}_{-2.0} h^{-1}\text{Mpc}$; Cooke et al. 2006). While this implies that LBGs and DLAs occupy similar environments, these clustering studies do not sample scales < 400 kpc (as we do here).

5.3. Column Densities and Covering Fractions of Metal Lines

Our medium- and high-resolution spectroscopy uniquely allows us to investigate the properties of metal-enriched halo gas and compare them directly to those observed in the DLA hosts. First, we compare the column densities for different ions measured in the CGM sightlines to column densities of the same ion in the associated DLA sightlines. An illustration of this comparison is shown in Figure 6, which presents the column densities of Si IV and Si II in our CGM sightlines versus R_{\perp} . The colors represent the corresponding DLA column densities for these ions. The horizontal-dashed lines in Figure 6 represent the threshold above which 90% of metal-line column densities for the DLA sightlines fall and can be used to compare individual CGM sightlines to the column densities of the majority of our DLA sample. The vast majority of the DLA column density upper limits are below these thresholds and most lower limits are above them.

First, we note that we measure overall higher DLA column densities of Si II relative to Si IV, which is consistent with previous metal-line absorption studies for DLAs (e.g., Vladilo et al. 2001; Fox et al. 2007; Mas-Ribas et al. 2017) and indicative that DLAs probe overall neutral environments. Second, we find that the CGM column densities of Si II are significantly lower (below the dashed line) than those measured in the DLAs across the full range in R_{\perp} of our sample. By contrast, the CGM column densities of Si IV are similar to those of the DLAs within $R_{\perp} \leq 150$ kpc. Similar patterns are apparent in all of the elements analyzed in this work for which we have access to both low- and high/intermediate-ionization species, including C II, C IV, Al II, and Al III (see Figure 21 in Appendix C). These results imply that high-ionization species observed in DLA sightlines trace halo gas out to distances of ≈ 150 kpc. This finding verifies the results of studies such as Wolfe & Prochaska (2000), who showed that C IV and Si IV

velocity profiles in DLAs are consistent with those arising from halo gas in semianalytic cold dark matter models.

To further investigate the extent of high- (represented by Si IV and C IV), intermediate- (represented by Al III), and low- (represented by Si II, Al II, and C II) ionization gas in DLA halos, we calculate covering fractions for each of these ionic transitions (shown in Figure 7). We require that at least two column density measurements be available in a given bin to compute the corresponding covering fraction. To assess how the column densities in the CGM compare to those of the DLA sightlines and account for the relative abundances of each individual ion, we use a column density threshold set at the 10th percentile value of the column densities measured in the DLA sightlines for each species, i.e., 10% of all DLA metal columns for that species lie below the chosen threshold. Thus, these covering fractions trace the incidence of absorption similar in strength to that observed in DLAs.

We find in general that all high- and intermediate-ionization species have large covering fractions compared to low-ionization species. Within 200 kpc of DLAs, high- and intermediate-ionization species have covering fractions above 40%, while the incidence of low-ionization species never exceeds 30% even at $24 \text{ kpc} < R_{\perp} < 100 \text{ kpc}$. This indicates that the warm ionized material associated with DLAs frequently extends over $\gtrsim 100$ kpc scales, whereas cool photoionized or neutral material seldom exhibits DLA-level absorption strengths across more than $\gtrsim 25$ kpc. We place upper limits on the covering fractions of all intermediate and low ions beyond $R_{\perp} > 200$ kpc (yielding an incidence of $< 32\%$ for Al III, $< 32\%$ for Al II, $< 25\%$ for C II, and $< 20\%$ for Si II); however, these species do exhibit some absorption that is weaker than the corresponding 10th-percentile column density threshold.

Lastly, we investigate the extent of metals in the CGM of DLAs in comparison to other CGM environments at similar redshifts: in quasar halos (Lau et al. 2016) and the halos of LBGs (Rudie et al. 2019a). The results are shown in Figure 8. Here, we measure covering fractions with a threshold of $N > 10^{13} \text{ cm}^{-2}$ for species of carbon and silicon and $N > 10^{12.1} \text{ cm}^{-2}$ for species of aluminum to account for the

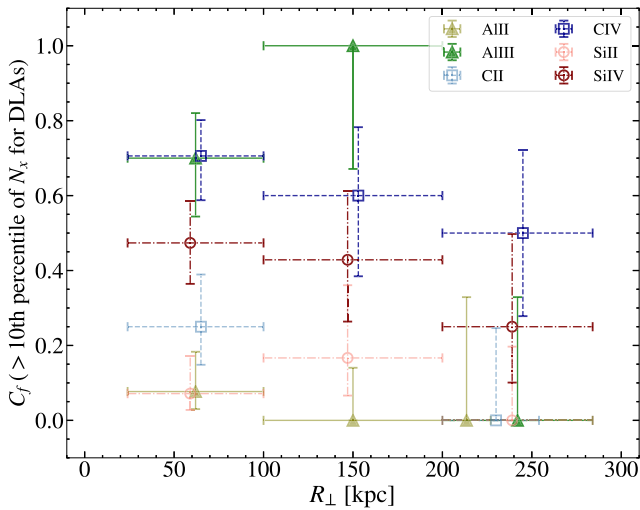


Figure 7. Covering fractions for Al II, Al III, C II, C IV, Si II, and Si IV with the threshold set to encompass the upper 90% of column densities from the DLA sightlines for each species. The x -axis error bars represent the bins used to compute each covering fraction. We reduce the ranges of the first and last bins to indicate the span of R_{\perp} for our sample. In some cases, the bins are sparsely populated, and therefore may not span the full range as represented by the x -axis error bars. The y -axis error bars represent the 1σ Wilson score interval for each covering fraction estimate. We find that (1) the halos of DLAs exhibit a higher incidence of high- and intermediate-ionization gas across all distance bins, and (2) covering fractions decrease as a function of projected distance for most species.

expected differences in the abundance of each element. We only report covering fractions calculated with at least two measurements within each R_{\perp} bin. The y -axis error bars indicate the 1σ Wilson score interval, which accounts for the sample size in each bin. For reference, the DLA-CGM sample has the largest number of measurements (between eight and 18) within 100 kpc, between three and eight measurements at $100 \text{ kpc} < R_{\perp} < 200 \text{ kpc}$, and between two and four measurements at $200 \text{ kpc} < R_{\perp} < 300 \text{ kpc}$. The Rudie et al. (2019a) sample includes between six and eight measurements within 100 kpc, and the Lau et al. (2016) sample includes two measurements within 100 kpc, between three and five measurements at $100 \text{ kpc} < R_{\perp} < 200 \text{ kpc}$, and either four or five measurements at $200 \text{ kpc} < R_{\perp} < 300 \text{ kpc}$.

We find the covering fractions of high-ionization species (Si IV and C IV) around DLAs and LBGs to be similarly high, which is suggestive of a volume-filling medium that extends to comparable projected distances in both environments. In contrast, our sample of DLA halos exhibits higher covering fractions for C II and Si II than the LBG sample by $\sim 2\sigma$, though we caution that the statistical uncertainties are significant. These differences may indicate, e.g., that we are preferentially selecting regions with a higher incidence of neutral (and hence also low-ionization) gas by targeting the CGM of DLAs, and/or that the metallicity of halo material around LBGs is generally lower than that in DLA halos.

Metal covering fractions in quasar halos are either larger than or consistent with those measured around DLAs in all ions. Within $R_{\perp} < 100 \text{ kpc}$, all ionic covering fractions for DLA and QSO halos are consistent within $< 1\sigma$. Within $100 \text{ kpc} < R_{\perp} < 200 \text{ kpc}$, the covering fractions of Si II, C II, and Al III are consistent within $< 1\sigma$. Beyond 200 kpc, all high- and low-ion covering fractions in DLA halos are $\gtrsim 1\sigma$ lower than those measured in QSO halos, which suggests that metal-

enriched gas pervades these latter massive halos to larger impact parameters.

5.4. Metallicity of DLAs and Their Halos

With column densities in hand, we can now provide new constraints on the metallicities of high-redshift DLAs and their associated CGM. In an optimal scenario, the metallicities of absorption-line systems are estimated via photoionization modeling, which can simultaneously constrain the ionization parameter of the system along with its metallicity (e.g., Crighton et al. 2015; Fumagalli et al. 2016b; Prochaska et al. 2017a). However, given that our column density measurements include numerous upper limits with relatively high values (e.g., $N_{\text{SiII}} \lesssim 10^{13-14} \text{ cm}^{-2}$), it is unlikely that photoionization modeling will yield robust metallicity constraints for our data set. Instead, we adopt a simpler approach using ionic ratios to assess the ionization state of each system, as described in Prochaska et al. (2015). We then explore the ratios of low-ionization metal column densities to those of neutral hydrogen, which may be used as a proxy for metallicity in systems with negligible ionization corrections (e.g., Prochaska et al. 2015).

5.4.1. Constraining Metallicities

To estimate the metallicities of our DLA and CGM absorption systems, we make use of a quantity introduced by Prochaska et al. (2015):

$$\{X_i/H_j\} = \log(N_{X_i}/N_{H_j}) - \epsilon_X + \epsilon_H. \quad (5)$$

Here, ϵ is the logarithmic solar abundance for the element X , while i and j represent ionization levels. The bracket notation indicates an ionic ratio of two different elements that ignores ionization corrections. In cases in which the ionization correction is small, we will assume that $\{X_i/H_j\} = [X/H]$. We adopt solar elemental abundances from Asplund et al. (2009).

Previous studies have assessed approximate ionization corrections via the ratio $N_{\text{SiIV}}/N_{\text{SiII}}$. We discuss these ratios in detail for our sample, along with other ionic ratios sensitive to ionization state, in Appendix A. Our measurements of these ratios imply negligible ionization corrections for only a small subset of our CGM sightlines. Therefore, we primarily rely on the ratio $\{\text{OI}/\text{HI}\}$ as a direct indicator of metallicity. For CGM sightlines with $N_{\text{HI}} \gtrsim 10^{19} \text{ cm}^{-2}$, this ionic ratio is insensitive to ionization state due to the similar ionization potentials of H I and O I, the possibility of charge exchange between them (e.g., Field & Steigman 1971; Prochaska et al. 2015), and because oxygen is only weakly depleted by dust (Jenkins 2009). It is commonly assumed that $\{\text{OI}/\text{HI}\} \approx [\text{O}/\text{H}]$ for systems with $N_{\text{HI}} \gtrsim 10^{19} \text{ cm}^{-2}$ (e.g., Crighton et al. 2013). There are six CGM sightlines which have both $N_{\text{HI}} \gtrsim 10^{19} \text{ cm}^{-2}$ and a constraint on $\{\text{OI}/\text{HI}\}$. We also include in the following analysis three more CGM systems with both $N_{\text{HI}} > 10^{18.5} \text{ cm}^{-2}$ and a constraint on $\{\text{SiII}/\text{HI}\}$. Similar to $\{\text{OI}/\text{HI}\}$, $\{\text{SiII}/\text{HI}\}$ can be used to trace $[\text{Si}/\text{H}]$ for mostly neutral systems, although it is somewhat more sensitive to ionization state than $\{\text{OI}/\text{HI}\}$ and overestimates the metallicity as the ionized fraction increases. Since these three systems may have non-negligible ionization corrections, we report these $\{\text{SiII}/\text{HI}\}$ constraints as upper limits on $[\text{Si}/\text{H}]$. For completeness, we also report the ionic ratios $\{\text{SiII}/\text{HI}\}$, $\{\text{CII}/\text{HI}\}$, $\{\text{FeII}/\text{HI}\}$, and $\{\text{OI}/\text{HI}\}$ for all DLAs in our sample in Appendix A. Table 5

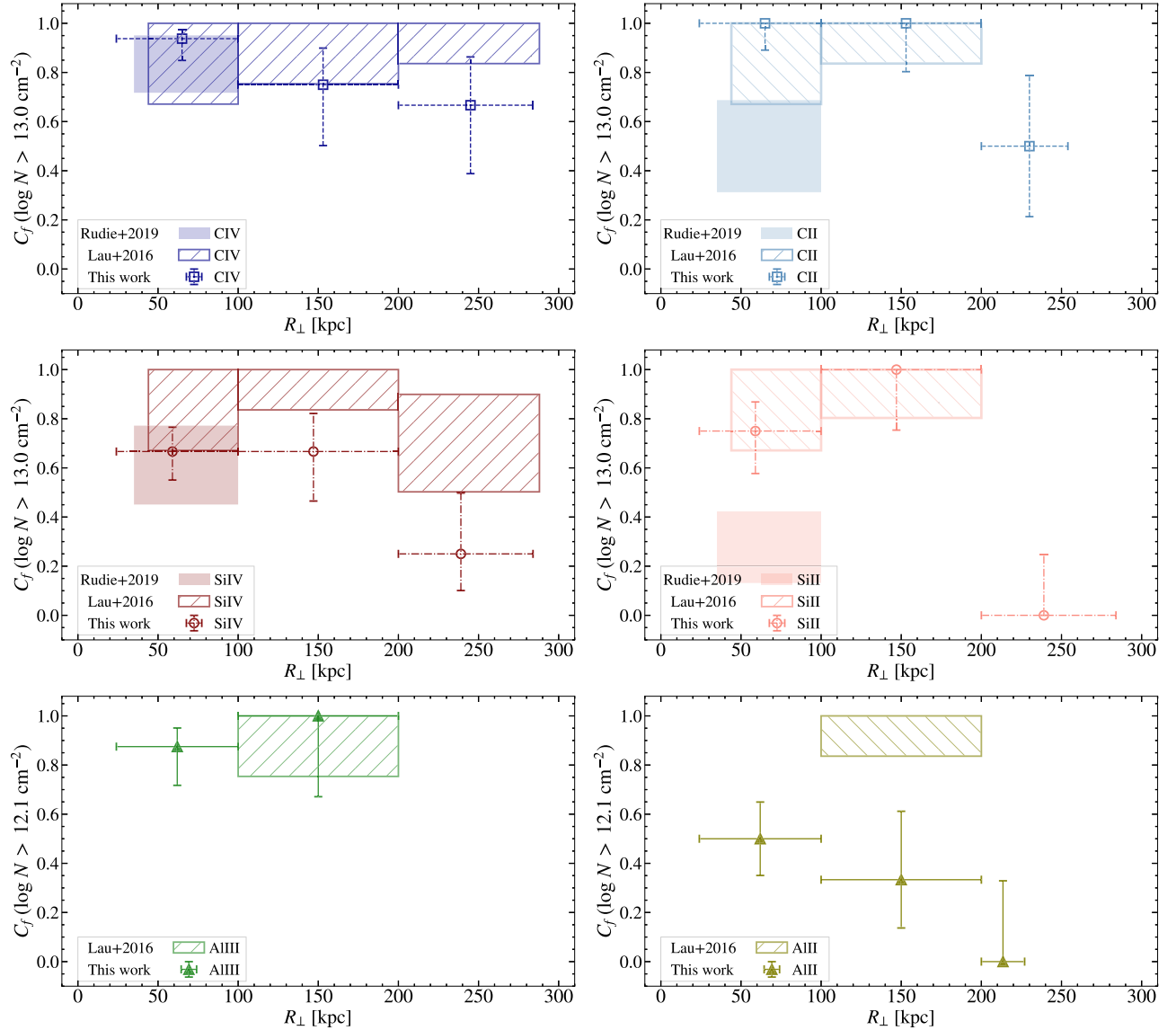


Figure 8. Covering fractions of carbon species (top row), silicon species (middle row), and aluminum species (bottom row) in DLA halos (points with error bars), quasar halos (Lau et al. 2016; hatched boxes), and LBG halos (Rudie et al. 2019a; filled boxes). High-ionization species are shown in the left-hand column and low-ionization species are shown in the right-hand column. Covering fractions are calculated with a threshold of $N > 10^{13} \text{ cm}^{-2}$ for carbon and silicon species and $N > 10^{12.1} \text{ cm}^{-2}$ for aluminum species. Vertical error bars and box heights represent 1σ Wilson score intervals. The x -axis error bars represent the bins used to compute each covering fraction. Although these error bars also show the span of R_\perp of our sample, in some cases the bins are sparsely populated and may not span the full range. The covering fractions in DLA halos are similar to those in LBG halos for high-ionization species and larger by $\sim 2\sigma$ for low-ionization species. DLA halo covering fractions are similar to those measured in QSO halos within 200 kpc (with the exception of covering fractions of Si IV and Al II) but decline relative to the latter beyond this distance.

lists all the ionic ratio measurements for DLA and CGM sightlines in this work.

5.4.2. Metallicity of DLA Halos

Our estimates of the metallicity of individual CGM sightlines, assessed via $\{\text{SiII}/\text{HI}\}$ and $\{\text{OI}/\text{HI}\}$, are shown in Figure 9. The six $\{\text{OI}/\text{HI}\}$ measurements with small ionization corrections and the three $\{\text{SiII}/\text{HI}\}$ metallicity limits are shown with red symbols. Two of these measurements, represented by the red square and the red star, are lower limits with $\{\text{OI}/\text{HI}\} > -1.08$ and -1.45 dex, respectively. The four other sightlines with $\{\text{OI}/\text{HI}\}$ constraints have metallicities ranging from -0.75 dex to at least as low as -2.09 dex. A comparison

of these measurements with CGM metallicities reported in the literature will be discussed in more detail in Section 6.

We also compare our CGM metallicities to the metallicities measured in the corresponding DLA sightlines as a function of R_\perp . For this analysis, we subtract the $\{\text{Xi}/\text{HI}\}$ measured in each DLA sightline from the same ionic ratio measured in its CGM (shown in Figure 10). We find the majority of the points are upper limits, due to the preponderance of upper limits on N_{SiII} and N_{OI} in our CGM sightline sample, and are likely overestimates given the unknown CGM ionization correction.

We comment here on a small subset of these sightline pairs that yield interesting constraints on the relative levels of enrichment in the DLA versus CGM sightlines. There are five systems with $\{\text{Xi}/\text{HI}\}_{\text{CGM}} - \{\text{Xi}/\text{HI}\}_{\text{DLA}}$ values less than 0 dex (four with $\{\text{SiII}/\text{HI}\}_{\text{CGM}} - \{\text{SiII}/\text{HI}\}_{\text{DLA}} < 0$, one of

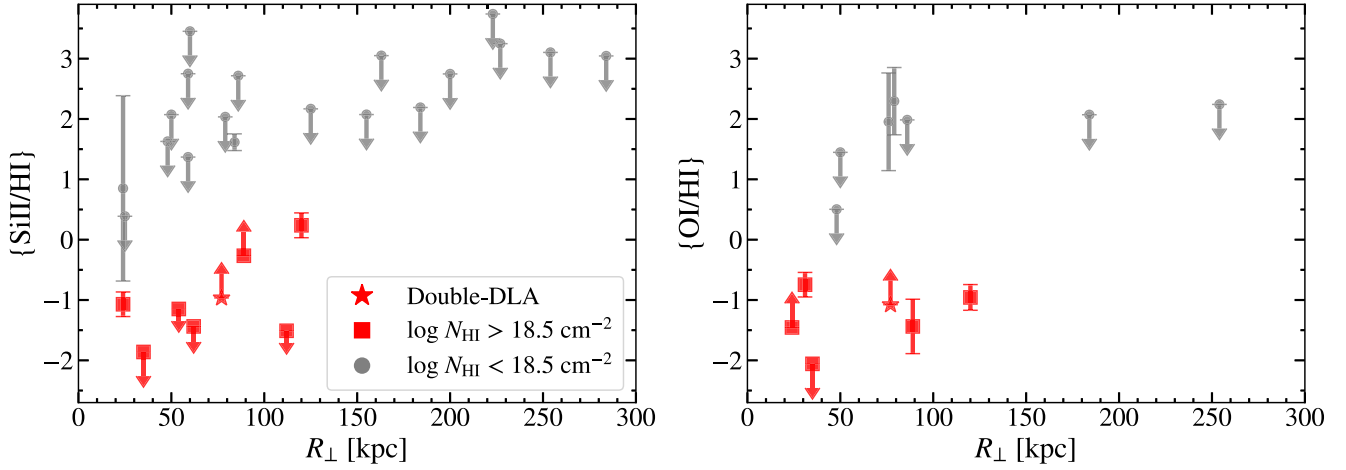


Figure 9. Logarithmic ionic ratios $\{\text{SiII}/\text{HI}\}$ (left-hand panel) and $\{\text{OI}/\text{HI}\}$ (right-hand panel) vs. R_{\perp} for our CGM sightlines. Metallicity constraints for systems with $N_{\text{HI}} > 10^{18.5} \text{ cm}^{-2}$ are represented by red symbols. Measurements shown with gray circles are likely overestimates of $[\text{X}/\text{H}]$. The double-DLA system is shown with a red star.

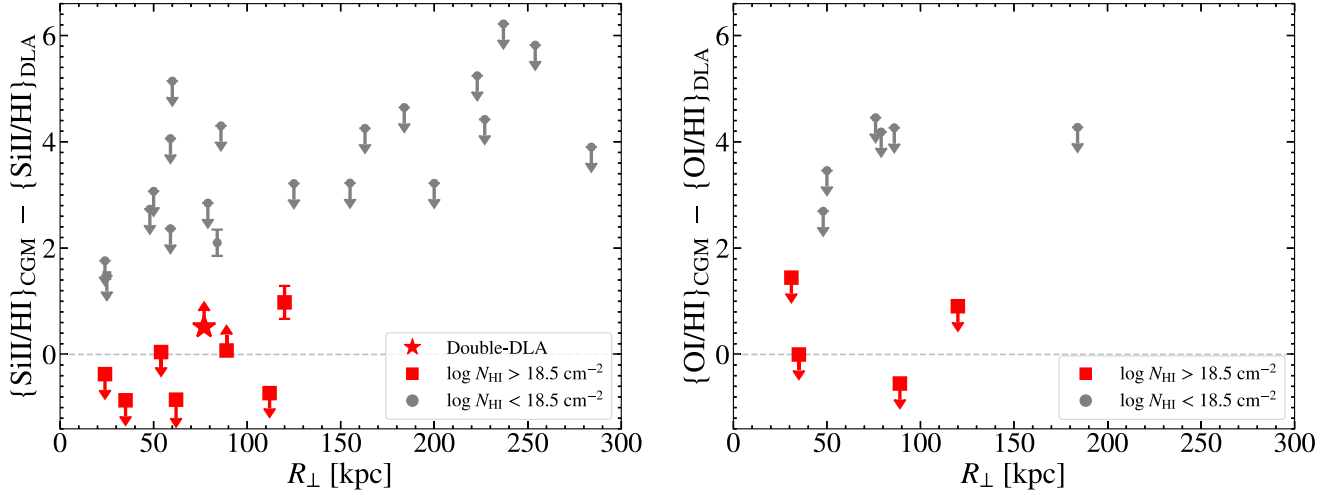


Figure 10. Comparison of metallicity constraints in our DLAs and the associated CGM as a function of R_{\perp} . Offsets that rely on ionic ratio measurements for CGM systems with $N_{\text{HI}} > 10^{18.5} \text{ cm}^{-2}$ are represented by red symbols. Metallicity offsets shown with gray circles are overestimates because our ionic ratios likely overestimate $[\text{X}/\text{H}]$ for the CGM in these systems. The double-DLA system is shown with a red star. All systems for which the CGM metallicity estimate is lower than that of the DLA have optically thick CGM gas. There are four systems for which we can say with confidence that $\{\text{SiII}/\text{HI}\}_{\text{CGM}} - \{\text{SiII}/\text{HI}\}_{\text{DLA}} < 0$ dex and two systems for which $\{\text{OI}/\text{HI}\}_{\text{CGM}} - \{\text{OI}/\text{HI}\}_{\text{DLA}}$ is likely < 0 dex.

which has both $\{\text{SiII}/\text{HI}\}_{\text{CGM}} - \{\text{SiII}/\text{HI}\}_{\text{DLA}}$ and $\{\text{OI}/\text{HI}\}_{\text{CGM}} - \{\text{OI}/\text{HI}\}_{\text{DLA}} < 0$, and an additional sightline with $\{\text{OI}/\text{HI}\}_{\text{CGM}} - \{\text{OI}/\text{HI}\}_{\text{DLA}} < 0$. Each of these systems has optically thick gas in the CGM sightline ($N_{\text{HI,CGM}} > 10^{18.55} \text{ cm}^{-2}$), and all occur within $R_{\perp} < 120$ kpc. In these five cases, we may state with at least $\sim 1\sigma$ confidence that the metallicities in the CGM sightlines are lower than those in the corresponding DLAs by at least -0.2 dex.

There is one sightline with a robust measurement of $\{\text{SiII}/\text{HI}\}_{\text{CGM}} - \{\text{SiII}/\text{HI}\}_{\text{DLA}} > 0.5$ dex. This sightline is the double-DLA, and may therefore probe a different environment than is typical of the other CGM sightlines in our sample. Nonetheless, it is the only system where we are certain the metallicity in the sightline with the lower N_{HI} is higher than in the so-called host DLA.

Taken together, these results point to a significant degree of scatter in the level of enrichment in the CGM at $24 \text{ kpc} < R_{\perp} < 120 \text{ kpc}$ relative to that in the DLA gas in the corresponding galaxy host.

5.4.3. Investigation of α/Fe Ratios in DLAs and Their Halos

In the previous section, we demonstrated that the CGM around DLAs has a wide range of metal enrichment, with robustly-estimated metallicities ranging from as high as -0.75 dex to at least as low as -2.09 dex. We expect that this metal content was originally formed in the interiors of stars and ejected via supernovae (SNe). The comparison of the abundance of α elements to that of Fe is useful in dissecting the specific nucleosynthetic processes that ultimately produced this enriched gas. α elements are produced in massive stars and are ejected by Type II SNe. This process happens on relatively short timescales (10^{6-7} yr). Meanwhile, Fe is produced in both Type II and Type Ia SNe. Type Ia SNe occur on longer timescales, on the order of 10^{8-9} yr (Kobayashi & Nomoto 2009). Once Type Ia SNe begin within a stellar population, the overall α/Fe of the surrounding gas will decrease (Tinsley 1979; Matteucci & Recchi 2001). An intermediate-redshift ($0.1 < z < 1.24$) study from Zahedy et al. (2016) measured the α/Fe ratio in halo gas close to

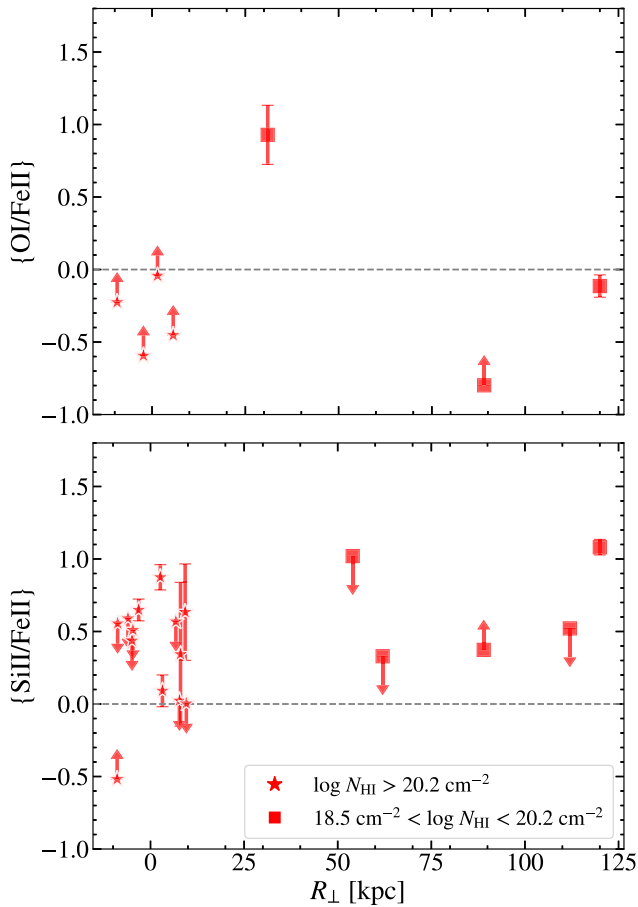


Figure 11. $\{OI/FeII\}$ (top) and $\{SiII/FeII\}$ (bottom) vs. R_{\perp} for sightlines with $\{X/HI\} < -1$ dex. DLAs are represented by stars and LLSs are shown as squares.

galaxies ($R_{\perp} < 60$ kpc) and showed increased α -enrichment in star-forming galaxy halos ($\alpha/Fe = 0.25 \pm 0.21$ dex) compared to that of quiescent galaxy halos ($\alpha/Fe = 0.06 \pm 0.15$ dex). They also found that the α -enrichment increased at larger distances ($R_{\perp} > 60$ kpc), measuring $\alpha/Fe = 0.9 \pm 0.4$ dex around star-forming galaxies and $\alpha/Fe > 0.3$ dex around quiescent galaxies at these distances. They concluded that the higher α -enrichment around star-forming galaxies is a consequence of the presence of young star-forming regions, whereas the higher α -enrichment in the outer halos of quiescent galaxies is suggestive of core-collapse dominated enrichment histories. Thus, this work successfully uses the α/Fe ratio measured in CGM material to trace differences in the stellar populations dominating its enrichment.

When measuring the α/Fe ratios in our sample, we must consider the depletion of Fe due to dust. This depletion scales with metallicity, and therefore has a larger impact on the measured α/Fe ratio in higher metallicity systems. A study of α -enrichment in a larger sample of DLAs found a positive correlation between α/Fe and metallicity for higher metallicity systems ($[X/H] > -1$ dex), which suggests that the depletion of Fe due to dust makes the α/Fe measurement unreliable (Rafelski et al. 2012). Therefore, we limit this analysis to only include sightlines in which we robustly measure a metallicity < -1 dex.

Figure 11 shows two ionic ratios that trace α/Fe : $\{OI/FeII\}$ (top panel) and $\{SiII/FeII\}$ (bottom panel). Higher levels of

ionization tend to elevate $\{SiII/FeII\}$ ratios, so our values may be overestimates of $[Si/Fe]$ for our CGM sightlines. Conversely, $\{OI/FeII\}$ is not very sensitive to ionization corrections; however, most sightlines yield lower limits on this quantity because O I 1302 is saturated in high- N_{HI} sightlines or because Fe II is typically not securely detected for lower- N_{HI} sightlines.

From our constraints on $\{OI/FeII\}$, we find that at least one CGM sightline is α -enriched. There is also one CGM sightline at $R_{\perp} = 120$ kpc for which our constraints imply that the system is not α -enriched, with $\{OI/FeII\} = -0.11 \pm 0.08$ dex. This sightline probes high-metallicity gas at $Z \approx 0.1Z_{\odot}$. The kinematics of this gas (which will be discussed in more detail in Section 5.5) are relatively quiescent, with a Δv_{90} width = 87 km s^{-1} for low-ionization material. Overall from $\{OI/FeII\}$, we find the α -enrichment of this population of DLAs and their halos is ambiguous due to the small sample size and preponderance of limits.

$\{SiII/FeII\}$ ratios yield more detections than $\{OI/FeII\}$ but are overestimates for sightlines with significant ionization. The majority of our CGM sightlines are ionized (implying large ionization corrections to $\{SiII/FeII\}$), and therefore we focus here on the DLA population (red points in the bottom panel of Figure 11). For the low-metallicity DLA population, we find that at least six sightlines are α -enriched, with a median value of $\{SiII/FeII\} = 0.52$ dex among the six detections. This is larger than what was measured in Rafelski et al. (2012). In their examination of DLA abundances at $z > 1.5$, they likewise found low-metallicity DLAs to be mostly α -enriched; however, they reported a mean value of $[\alpha/Fe]$ in their low-metallicity ($[X/H] < -1$ dex) DLA sample of 0.27 ± 0.02 dex.

Together, these results confirm that (1) our low-metallicity DLA population is mostly α -enriched with a median value of $\{SiII/FeII\} = 0.52$ dex; and (2) for CGM sightlines in which we robustly measure α/Fe , we find one sightline is α -enriched and one has an abundance ratio near solar. However, due to uncertainties in both ionization state and the degree of dust depletion across our sample, we cannot comment more generally on the α -enrichment of the CGM of DLAs.

5.5. Kinematics

In this section, we investigate the kinematics of high- and low-ionization gas surrounding DLAs. By necessity, this analysis is limited to sightlines with significantly detected metal-line absorption profiles. As a result, we caution that our conclusions will be biased toward those systems with significant metal content. In many CGM sightlines, there are few detected metal absorption lines, so we choose the strongest transitions as follows: to trace high-ionization gas, we choose C IV $\lambda 1548$ for its high oscillator strength; and to trace low-ionization gas, we choose a low-ion transition with the highest S/N at the peak of the optical depth profile (i.e., the profile with the highest value of $\tau_{\text{peak}}/\langle\sigma_{\tau}\rangle$). We make use of two quantities (fully described in Section 4.2): the flux-weighted velocity centroids (δv_{weight}) measured relative to the DLA redshifts, and the velocity widths measured between the locations where the cumulative optical depth profile reaches 5% of the total integrated optical depth on either side (Δv_{90}). For some sightlines, the C IV $\lambda 1548$ transition is affected by saturation, and in these cases the associated Δv_{90} will likely overestimate the width of 90% of the total line optical depth to some degree. However, because the column densities of C IV in

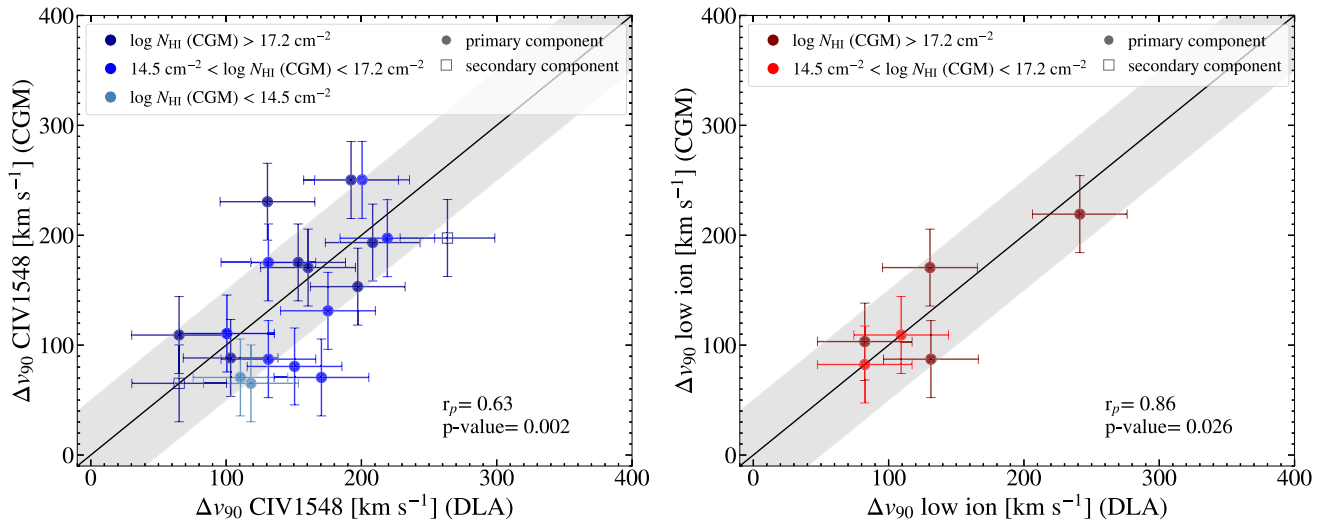


Figure 12. Comparison of Δv_{90} velocity widths for the DLA and CGM sightlines measured from C IV 1548 (left-hand panel) and from an unsaturated low-ionization transition (right-hand panel). Open squares represent secondary velocity components. The black lines show a 1:1 relation, and the gray bars show $\pm 50 \text{ km s}^{-1}$ offsets from this relation in the x - and y - directions. The color of each point is indicative of the H I column density in the corresponding CGM sightline, as shown in the legend. The Δv_{90} velocity widths measured from both C IV and low-ionization transitions closely follow a 1:1 relation. The statistical significance of the correlations between these quantities is indicated with the correlation coefficient (r_p) and the P -value in the bottom right of each plot.

our DLA versus CGM sightlines have overall similar values at $R_{\perp} < 200 \text{ kpc}$ (as shown in Appendix C Figure 21), we posit that saturation effects should not systematically impact our measured DLA velocity widths more than our CGM velocity widths, or vice versa. We adopt uncertainties on these quantities as described in Section 4.2.1.

First, we investigate the difference in the velocity widths for high- and low-ionization gas between DLAs and the corresponding CGM sightlines. We show these results in Figure 12. The colors in Figure 12 represent different bins of CGM H I column density. We include secondary components as squares where there is a detection in both the CGM and DLA sightline. We also include gray bars to show an offset of $\pm 50 \text{ km s}^{-1}$, representing the uncertainty in the x - and y - directions (35 km s^{-1} uncertainty for our Δv_{90} values) added in quadrature. For both high- and low-ionization material, this comparison reveals clear correlations between the DLA and CGM line widths. We perform a Pearson rank correlation test on the two data sets to quantify the strength and significance of a linear correlation of these quantities. For C IV $\lambda 1548$, we find a Pearson correlation coefficient (r_p) of 0.63 with a P -value of 0.2%, which is indicative of a very low probability that these quantities are uncorrelated. For our low-ion Δv_{90} measurements, we find $r_p = 0.86$, which is indicative of a close to 1:1 relation, with a P -value of 2.6%. While these results are suggestive of strong correlations in both cases, we also note a larger degree of scatter in the Δv_{90} measurements of C IV $\lambda 1548$: 35% of our sightline pairs have Δv_{90} values that differ by more than 50 km s^{-1} , which are larger than the uncertainty associated with our Δv_{90} measurements (35 km s^{-1}). For low-ionization gas, none of the Δv_{90} values in our sightline pairs differ by more than 50 km s^{-1} . This distinction may reflect a larger degree of variation in the kinematics of high-ionization gas in nearby sightlines or may be driven by saturation effects in our $\Delta v_{90}(\text{C IV})$ values. We note that systems with differing values of CGM H I column density appear to yield a consistent level of scatter in these quantities. Overall, these results reveal a close correspondence in velocity widths over 24–237 kpc scales. This in turn suggests that these Δv_{90} widths are a

consistent tracer of the potential well of the host halo, regardless of the impact parameter of the sightline.

Both Christensen et al. (2019) and Møller & Christensen (2020) performed a close examination of the Δv_{90} values measured for DLAs as a function of the projected distance from their host galaxies (identified in emission). In particular, Møller & Christensen (2020) measured a gradient of $-0.017 \text{ dex kpc}^{-1}$ in the quantity $\log \Delta v_{90}/\sigma_{\text{em}}$ over an impact parameter range $0 \text{ kpc} < R_{\perp} < 60 \text{ kpc}$, with σ_{em} equal to the velocity width of strong emission lines. These authors demonstrated that this trend is consistent with the projected velocity dispersion profile predicted for a Dehnen (1993) dark matter halo potential model. Moreover, they pointed out that this latter profile flattens at impact parameters $R_{\perp} > 60 \text{ kpc}$. Our finding of a close correspondence between Δv_{90} values over scales of $\gtrsim 100 \text{ kpc}$ is fully consistent with this prediction and may be viewed as further confirmation of the interpretation of Δv_{90} as an effective measure of halo dynamics.

In our previous work (Rubin et al. 2015), we compared the δv_{weight} values for C IV $\lambda 1548$ for 12 DLA-CGM sightline pairs, eight with medium-resolution spectroscopy ($\mathcal{R} \gtrsim 4000$) and four with low-resolution spectroscopy ($\mathcal{R} \lesssim 2400$). We found that the differences in $\delta v_{\text{weight}}^{\text{C IV } \lambda 1548}$ did not exceed 105 km s^{-1} across the full sample, which included sightlines with separations up to $R_{\perp} = 176 \text{ kpc}$. We interpreted this finding as suggestive of strong coherence in C IV absorption over scales of $> 100 \text{ kpc}$. Here, we expand on this analysis with a larger sample of medium- and high-resolution spectroscopy (drawing on 21 absorber pairs, including 19 primary and two secondary components). The results are shown in Figure 13. We split the sample into two bins of DLA metallicity at the median metallicity of the DLA sightlines ($\log Z/Z_{\odot} = -1.07$). The high-metallicity DLAs are more likely to trace more massive halos (with $10^{11.5} M_{\odot} \lesssim M_h \lesssim 10^{12} M_{\odot}$) with $\text{SFR} \gtrsim 1 M_{\odot} \text{ yr}^{-1}$ (Krogager et al. 2017), while lower-metallicity DLAs are associated with lower SFRs and halo masses ($10^{10} M_{\odot} \lesssim M_h \lesssim 10^{11.5} M_{\odot}$; e.g., Bird et al. 2014). We highlight the differences in $\delta v_{\text{weight}}^{\text{C IV } \lambda 1548}$ in our CGM versus DLA sightlines as

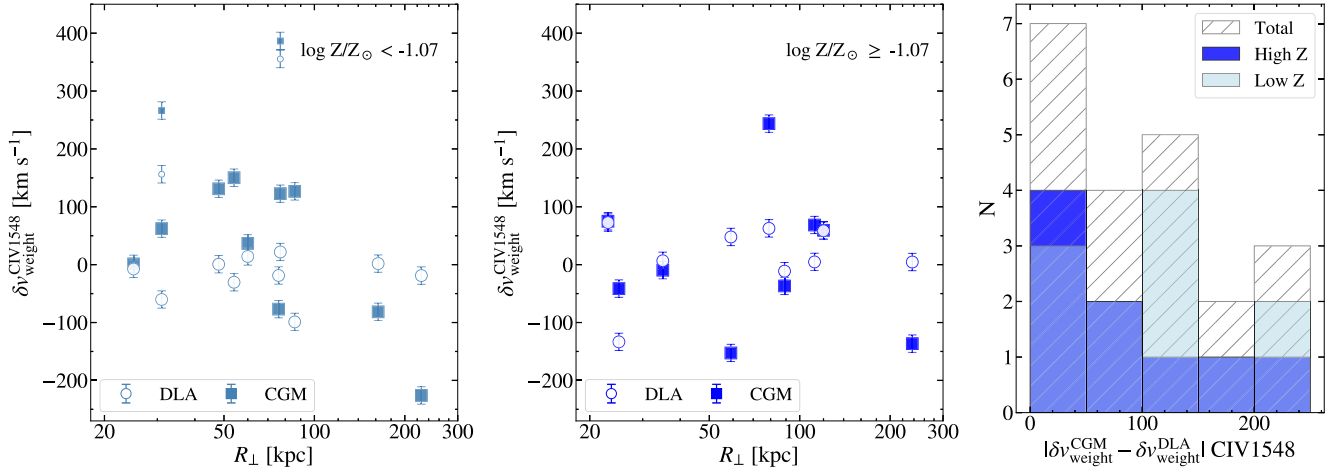


Figure 13. Flux-weighted velocity centroids for C IV ($\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$) for our DLA (open circles) and CGM (filled squares) sightlines plotted vs. R_{\perp} for each sightline pair. For direct comparison between the CGM and DLA $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$, the DLA measurements are placed at the same R_{\perp} as the corresponding CGM sightline. Smaller symbols indicate secondary components. The left-hand panel shows systems with DLA metallicities $\log Z/Z_{\odot} < -1.07$, while the middle panel includes systems with higher DLA metallicities. The right-hand panel shows a histogram of the difference between $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ measured for each DLA and CGM sightline pair with a bin width of 40 km s^{-1} . The hatched histogram shows the full sample, the dark-blue shading shows the high-metallicity sample, and the light-blue shading shows the low-metallicity sample. We find that all systems have a $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ difference of $< 225 \text{ km s}^{-1}$, and 52% have a difference of $< 100 \text{ km s}^{-1}$.

a function of DLA metallicity in the right-hand panel in Figure 13.

Overall, we find no evidence for a correlation between the difference in $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ for the DLA and CGM sightlines and metallicity, separation between the sightlines (R_{\perp}), C IV column density, or H I column density. In addition, this sample yields larger differences in $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ than that analyzed in Rubin et al. (2015). We find that 52% of these sightline pairs have $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ values that differ by $\leq 100 \text{ km s}^{-1}$. Moreover, 86% of our pairs yield differences in $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ of $\leq 200 \text{ km s}^{-1}$. There are three pairs that have differences larger than 200 km s^{-1} . This expanded sample shows clear evidence that the velocities of C IV in the outer halos of DLA hosts are frequently more than 100 km s^{-1} different from that measured in the inner CGM. We use the relations given in Maller & Bullock (2004) to estimate the virial velocities of the DLA host halos (assuming that they have halo masses $M_h \sim 10^{11-12} M_{\odot}$ at $z_{\text{abs}} \sim 2.45$), finding that they span the range $100\text{--}216 \text{ km s}^{-1}$. This suggests that our measured differences in $\delta v_{\text{weight}}^{\text{CIV}\lambda 1548}$ are less than or approaching the virial velocity of the host halos for $\sim 52\%$ of sightlines. Overall, with access to a larger sample, we show there is not a strong coherence ($\lesssim 100 \text{ km s}^{-1}$) in C IV velocity centroids over large scales as seen in our previous work.

Finally, to investigate the possibility of CGM gas escaping the DLA host halos, we compare our measurements of δv_{weight} and Δv_{90} to the radial escape velocities ($v_{\text{esc}} = \sqrt{GM_h/R}$, with $R = R_{\perp}$) of halos with three different values of total mass ($M_h = 10^{10} M_{\odot}$, $10^{11} M_{\odot}$, and $10^{12} M_{\odot}$). The results are shown in Figure 14. The error bars above and below δv_{weight} represent the velocities encompassed by the Δv_{90} interval relative to the line center. Low-ionization gas is represented by circles, and high-ionization gas traced by C IV $\lambda 1548$ is represented by squares. Primary and secondary velocity components are shown with filled and open markers, respectively. As we did in the previous analysis, we divide our sample by the median DLA metallicity, to differentiate between sightlines that likely

trace lower-mass halos ($\lesssim 10^{11.5} M_{\odot}$) and those that are more likely to probe higher-mass halos ($\gtrsim 10^{11.5} M_{\odot}$).

Before interpreting these results, we consider two caveats. First, we note that for any given sightline pair, the R_{\perp} that we measure does not necessarily reflect the true projected distance of the CGM sightline from the center of its host halo because DLAs do not always lie at the centers of their hosts. However, the difference between our measured R_{\perp} and the true R_{\perp} is likely to be small because observational and theoretical studies typically measure DLA-galaxy separations to be $< 25 \text{ kpc}$ (Krogager et al. 2017). Second, we caution that our analysis assesses velocities along the line of sight, rather than the total radial velocity of gas measured with respect to each halo’s center. Our velocities should therefore be interpreted as lower limits of this latter quantity.

Within 100 kpc of the DLAs, we frequently detect velocity components in both high- and low-ionization absorption profiles that have a δv_{weight} , which exceeds the escape velocity for halos with $M_h \leq 10^{11} M_{\odot}$. Moreover, there are two secondary velocity components detected that have a δv_{weight} exceeding the escape velocity of a halo with $M_h = 10^{12} M_{\odot}$. One of these components is detected in the double-DLA sightline, shown in orange in Figure 14, and therefore may trace a different CGM environment, possibly a different halo, from that of the typical isolated DLA. The other high-velocity component, detected in both low-ions and C IV, is at the edge of our search window at 946 km s^{-1} . This gas is likely unbound from the central DLA host halo. Beyond 150 kpc , we detect C IV absorption from a single system at a δv_{weight} that exceeds escape for a halo with $M_h = 10^{12} M_{\odot}$. All three of these velocity components that are detected in excess of the escape velocity for a $M_h = 10^{12} M_{\odot}$ halo are in systems that have low-metallicity DLAs ($\log Z/Z_{\odot} < -1.07$).

Overall, we find no significant correlation between the measured distribution of CGM gas velocities and R_{\perp} or DLA metallicity. We find that 32 of these 35 components are likely to be bound under the assumption that they reside in halos with masses of $\sim 10^{12} M_{\odot}$. The remaining three components are in turn very likely to escape their host halos regardless of their

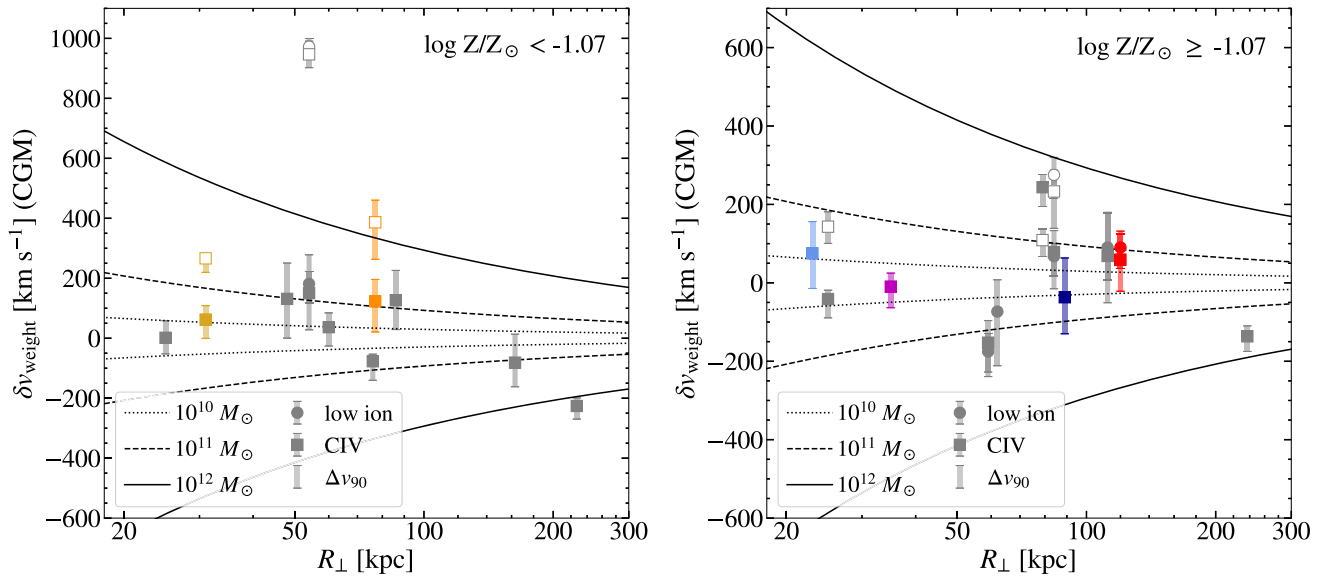


Figure 14. Flux-weighted velocity centroids for low-ions (circles) and C IV (squares) in our CGM sightlines vs. R_{\perp} for systems with low-metallicity DLAs (left-hand panel) and high-metallicity DLAs (right-hand panel). The error bars here represent the velocity range encompassed by the Δv_{90} width relative to the component line center, and therefore show the approximate velocity range exhibited by each absorption component. Primary velocity components are shown with filled markers and secondary velocity components are shown with open markers. The colors highlight systems in which we have robust metallicity constraints mentioned in Section 5.4, which will be used in discussion later in Section 6. The double-DLA system is represented by the orange points in the left-hand panel. Curves represent radial escape velocities for halos with masses $10^{10} M_{\odot}$ (dotted lines), $10^{11} M_{\odot}$ (dashed lines), and $10^{12} M_{\odot}$ (solid lines). Most points lie within these bounds for halos of $10^{11-12} M_{\odot}$. This suggests that if DLAs reside in such massive halos, then the bulk of the absorbing gas traced by both C IV and low-ionization lines will remain bound.

precise dark matter mass (given that they are almost certainly $\lesssim 10^{12} M_{\odot}$). In the case that these DLAs predominately reside in halos with $M_h \sim 10^{11} M_{\odot}$, approximately half (17) of the 35 components in our sample have velocities that exceed that required for escape.

6. Discussion

6.1. Implications for the Metallicity of DLA Halos

We may now place constraints on the metallicities of DLAs and their CGM in the context of other circumgalactic environments at $z \sim 2$. We focus our discussion on the subset of our constraints that we consider to be most robust. As described in Section 5.4, we include six CGM systems for which we can measure $\{\text{O I}/\text{H I}\}$ and which are highly optically thick (i.e., with $N_{\text{H I}} > 10^{18.5} \text{ cm}^{-2}$), such that we may assume $\{\text{O I}/\text{H I}\} \approx \{\text{O}/\text{H}\}$ (e.g., Crighton et al. 2013; Prochaska et al. 2015). We also include three more CGM systems with $N_{\text{H I}} > 10^{18.5} \text{ cm}^{-2}$ for which we constrain metallicity using $\{\text{Si II}/\text{H I}\}$, and report these measurements as upper limits (shown in black in Figure 15). For the corresponding DLA sightlines, we assume $\{\text{Si II}/\text{H I}\} \approx \{\text{Si}/\text{H}\}$ (because O I $\lambda 1302$ is typically saturated in these systems and ionization corrections are likely small). Metallicities for the nine sightline pairs in this subsample are indicated in Figure 15 with colored/black stars for the DLAs and colored/black squares for the CGM (with the point colors pairing CGM systems to the associated DLA).

We also compare our metallicity measurements to CGM metallicities from the literature. Following Figure 4 of Crighton et al. (2013), we indicate the metallicity of ISM absorption measured in a lensed LBG spectrum (cB58; Pettini et al. 2002) with a yellow horizontal bar. We note that oxygen abundances measured from H II region emission from LBGs also fall within this range ($Z/Z_{\odot} \sim 0.4\text{--}0.7$; Strom et al. 2018). We indicate

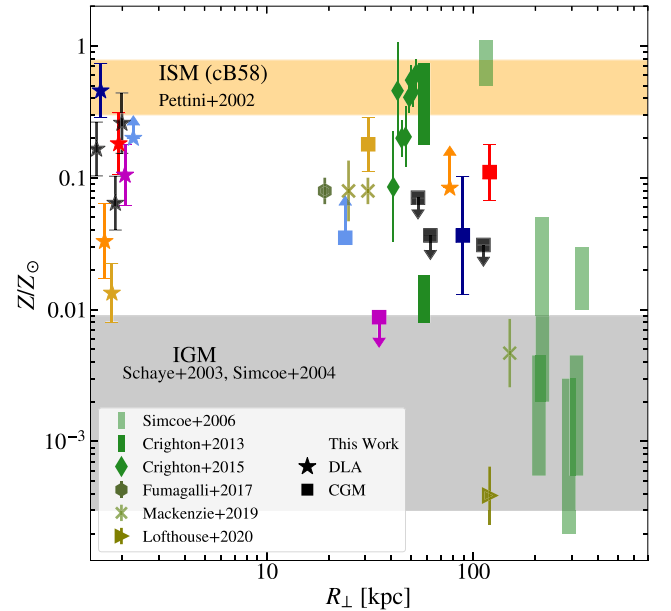


Figure 15. DLA and CGM metallicities vs. R_{\perp} . Metallicities measured along CGM sightlines having $10^{18.5} \text{ cm}^{-2} < N_{\text{H I}} < 10^{20.3} \text{ cm}^{-2}$ are shown as light-blue, dark-blue, purple, red, gold, and black squares. The double-DLA sightlines are represented by orange stars. Black squares show metallicities constrained by $\{\text{Si II}/\text{H I}\}$. The DLA sightlines are shown with stars near $R_{\perp} \approx 2 \text{ kpc}$, and are colored to indicate the corresponding CGM sightline. Green diamonds show metallicities of distinct absorption components detected in the halo of a $\sim 0.2 L^*$ galaxy at $z = 2.5$ (Crighton et al. 2015). The green rectangles indicate metallicity constraints on the CGM of LBGs from Crighton et al. (2013, dark green) and Simcoe et al. (2006, light green). The dark-green hexagon is a DLA located 19.1 kpc from a compact galaxy at $z = 3.25$ (Fumagalli et al. 2017). The green cross-hairs represent the metallicities of DLAs at $3.2 < z < 3.5$ detected close to confirmed LAEs (Mackenzie et al. 2019). The olive green triangle represents a LLS located 120 kpc from a LAE at $z = 3.53$ (Lofthouse et al. 2020). The yellow bar indicates the metallicity measured from interstellar absorption lines in the spectrum of cB58 (Pettini et al. 2002), and the gray bar shows the range in abundances observed in the Ly α forest (Schaye et al. 2003; Simcoe et al. 2004).

the range in abundances measured in the Ly α forest in gray (Schaye et al. 2003; Simcoe et al. 2004). Finally, we include measurements of the metallicity in CGM material detected around sub- L^* systems and LBGs at $z \sim 2.1$ – 3.6 reported in the literature (Simcoe et al. 2006; Crighton et al. 2013, 2015; Fumagalli et al. 2017; Mackenzie et al. 2019; Lofthouse et al. 2020).

In interpreting these results, we first emphasize that our analysis approach cannot reveal order-of-magnitude variations in metallicities along individual sightlines as observed by Crighton et al. (2013, 2015), and Simcoe et al. (2006). Instead, the bulk column densities that we use to compute ionic ratios are dominated by the absorption components with the largest columns along the line of sight—and these dominant components need not arise at the same velocity across all ions. Our metallicities assess the overall level of enrichment integrated along each sightline. Nevertheless, these measurements exhibit a large range of values that are consistent with those observed at much higher spectral resolution.

Looking at these results in detail, we find that three of our CGM sightlines (indicated in red, gold, and orange) exhibit the high metallicities ($\gtrsim 0.05 Z_\odot$) that are observed within $R_\perp \leq 100$ kpc of emission-selected galaxies at $z \sim 2.0$ – 2.5 . At the same time, the host DLAs of these systems have metallicities that are well below that typical of the ISM of LBGs at this epoch. Moreover, two of these DLAs have metallicities lower than that measured in their respective CGM sightline. Meanwhile, five of our CGM sightlines have metallicities lower than those of their respective DLA, including some of the highest metallicity DLAs included in this analysis. One of these CGM sightlines, shown as the purple square at $R_\perp = 35$ kpc, has a metallicity that is consistent with that typical of the IGM (Schaye et al. 2003; Simcoe et al. 2004) and of the CGM of LBGs measured at $R_\perp > 100$ kpc (Simcoe et al. 2006, $\log Z/Z_\odot < -2.09$). The low metallicity of this sightline is likely to be inconsistent with enriched galactic outflows and instead suggests the origin of this gas is from the surrounding IGM. We further discuss the implications of our metallicity measurements for the origins of the CGM material on a system-by-system basis in Appendix B. Overall, under the assumption that high-metallicity DLAs trace higher-mass halos than low-metallicity DLAs, our sample of DLA-CGM metallicities is not indicative of any strong dependence of CGM metallicity on halo mass.

Our findings are consistent with the picture that the CGM is inhomogeneous, containing both enriched (nearing the metallicity of the typical ISM of a LBG) and low-metallicity (near or within the enrichment level of the surrounding IGM) gas. The incidence of higher metallicity (> 0.01 dex) versus low-metallicity (< 0.01 dex) gas along sightlines with $\{\text{OI}/\text{HI}\}$ constraints is high (5:1) and could suggest lower covering fractions for low-metallicity material. Recent cosmological zoom simulations (e.g., Hafen et al. 2019; Stern et al. 2021) also predict a qualitatively inhomogeneous CGM at $z = 2$, and find that its mean metallicity similarly decreases significantly with distance from the host galaxy. Hafen et al. (2019) found that the fractions of the total CGM mass arising from wind material versus from accreted IGM gas are approximately equal within $\lesssim 0.5 R_{\text{vir}}$ for halos with masses $10^{10} M_\odot \lesssim M_h \lesssim 10^{12} M_\odot$, whereas wind material contributes only ~ 30 – 40% of the mass at $\sim R_{\text{vir}}$. This in turn yields a broad distribution of predicted metallicities throughout these environments, which

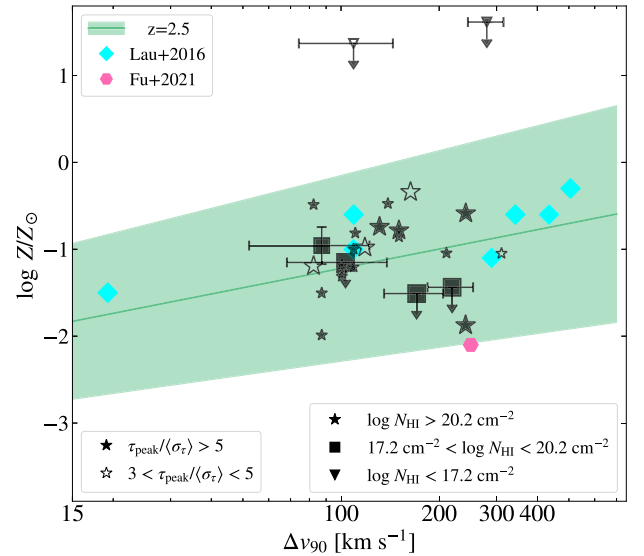


Figure 16. Metallicity vs. velocity width for our DLA (black stars) and CGM (black squares and triangles) sightlines. All systems marked with large squares and stars are also included in Figure 15. Additional DLA and CGM systems from our sample are marked with smaller stars and triangles, respectively. The linear relationship measured by Neeleman et al. (2013) for DLAs at the average redshift of our sample ($z = 2.5$) is shown with a light-green solid line. The width of the colored contour around this line indicates the measured $\pm 1\sigma$ scatter. Filled symbols indicate systems with high-S/N measurements of Δv_{90} (i.e., they have $\tau_{\text{peak}}/\langle\sigma_\tau\rangle > 5$), whereas open symbols indicate lower-S/N assessments ($3 \leq \tau_{\text{peak}}/\langle\sigma_\tau\rangle \leq 5$). We also include measurements of absorption from the CGM of a SMG at $z = 2.674$ (Fu et al. 2021, shown with a pink hexagon) and the CGM of QSOs at $z = 2$ – 3 (Lau et al. 2016, shown with cyan diamonds).

systematically shift to lower enrichment levels at larger radii. Further analysis is required to enable detailed comparisons between these predictions and the metallicities implied by our pencil-beam probes (which are sensitive to gas over a broad range of physical radii, and which may be dominated by the highest-metallicity material along the line of sight). The combined data sets shown in Figure 15 represent substantive observational constraints to motivate such a comparison.

6.2. A Global Velocity Width-Metallicity Relation for Absorption-Line Systems

Our unique data set allows us to investigate the Δv_{90} width-metallicity relation in our CGM sightlines and compare them to both the corresponding Δv_{90} for the DLA sightlines and impact parameter. We find the Δv_{90} width in the same way for both sets of sightlines, using an unsaturated low-ion transition with the highest ratio $\tau_{\text{peak}}/\langle\sigma_\tau\rangle$, and here do so without first separating the profiles into distinct velocity components. The majority of low-ion transitions in our CGM sightlines are not significantly detected, so this analysis is limited to six CGM sightlines. Two of these six systems have $N_{\text{HI}} < 10^{17.2} \text{ cm}^{-2}$, so the majority of this subset represent LLSs around DLAs (and one of these CGM sightlines is also a DLA). The Δv_{90} -metallicity distribution of these systems is shown in Figure 16, along with the same measurements for 20 of the DLAs in our sample. We show those systems with robust Δv_{90} measurements (having $\tau_{\text{peak}}/\langle\sigma_\tau\rangle > 5$) with filled symbols, and show those with lower-S/N constraints on Δv_{90} (with $3 \leq \tau_{\text{peak}}/\langle\sigma_\tau\rangle \leq 5$) with open symbols. Figure 16 also indicates the mean relation between Δv_{90} and metallicity for the DLA population as reported by Neeleman et al. (2013) at

$z = 2.5$ (the average redshift of our sample). The shaded region indicates the measured $\pm 1\sigma$ scatter in this relation. We find that the Δv_{90} -metallicity distribution of the DLAs in our sample is overall consistent with the relation fit by Neeleman et al. (2013).

To increase our sample size, we have included other CGM measurements from the literature. We consider measurements from Fu et al. (2021), who analyzed CGM absorption at a transverse separation of 93 kpc from a confirmed submillimeter galaxy (SMG) at $z = 2.674$, as well as measurements from Lau et al. (2016), who studied the CGM within $R_{\perp} < 300$ kpc of several QSO host galaxies at $z \sim 2-3$. For the Lau et al. (2016) sample, we used the metallicities reported in that work, and used their publicly-available spectra to calculate the Δv_{90} width for these systems in the same manner as for our sample. Seven of these QSO-CGM sightlines have low-ionization transitions that are sufficiently strong to yield a robust Δv_{90} measurement. Overall, the DLA-CGM and other CGM systems included here occupy a similar region of this parameter space as the DLA population.

The Δv_{90} -metallicity relation for DLAs is thought to be driven by the well-known relationship between mass and metallicity among star-forming galaxies (Tremonti et al. 2004; Erb et al. 2006; Møller et al. 2013), under the assumption that Δv_{90} is an effective tracer of the dynamical mass of the system (Ledoux et al. 2006; Prochaska et al. 2008). Neeleman et al. (2013) bolstered evidence for this assumption by demonstrating that DLAs exhibit a Δv_{90} -metallicity relation that evolves with cosmic time in a manner similar to the stellar mass-metallicity relation. The sample of Δv_{90} and metallicity measurements for CGM systems that we have assembled permits two novel tests of this picture.

First, we may compare the values of Δv_{90} measured for the QSO-CGM versus DLA-CGM samples. Because the dark matter halos hosting high-redshift QSOs have masses $M_h > 10^{12.5}$ (Wild et al. 2008; White et al. 2012; Font-Ribera et al. 2013), we posit that the associated Δv_{90} values must be overall larger than those exhibited by DLAs and their CGM, if this quantity is indeed an effective tracer of halo mass. We find that the average values of Δv_{90} for the DLA sightlines, DLA-CGM sightlines, and QSO-CGM sightlines are 141 km s^{-1} , 161 km s^{-1} , and 258 km s^{-1} , respectively. This latter value is $97-117 \text{ km s}^{-1}$ larger than the DLA and DLA-CGM sample means. At the same time, we see from Figure 16 that there is some overlap among the Δv_{90} values of these three subsamples. Six of the seven QSO-CGM sightlines (86%) have $\Delta v_{90} > 100 \text{ km s}^{-1}$, and three ($\approx 43\%$) have $\Delta v_{90} > 300 \text{ km s}^{-1}$. Among our DLA and DLA-CGM sightlines, 80% have $\Delta v_{90} > 100 \text{ km s}^{-1}$, and just one ($\approx 4\%$) has $\Delta v_{90} > 300 \text{ km s}^{-1}$. These findings are consistent with the idea that Δv_{90} is indeed correlated with halo mass, and furthermore bolsters previous indications that DLAs arise in halos across a broad mass range, including those as massive as QSO hosts (e.g., Font-Ribera et al. 2012; Mackenzie et al. 2019).

Second, we may use these samples to test for a dependence of Δv_{90} on the location of the background QSO sightline with respect to the halo center. We note first that the absolute differences in the Δv_{90} values measured for five of the six available DLA-CGM sightline pairs are small: these offsets have a median value of 21 km s^{-1} and a maximum value of 44 km s^{-1} . Only one sightline pair exhibits a much larger difference of 197 km s^{-1} . Given that the typical uncertainty associated with our Δv_{90} measurements is 35 km s^{-1} , this suggests that Δv_{90} is likely to be an effective tracer of

dynamical mass, regardless of the location of the background sightline (with the caveat that this paired sightline comparison sample is small).

The suggestion that the Δv_{90} -metallicity relation may change with distance from the host center was previously studied in the literature. As discussed in Section 5.5, the work of Christensen et al. (2019) and Møller & Christensen (2020) investigated this topic, finding a negative correlation in $\log \Delta v_{90}/\sigma_{\text{em}}$ versus projected distance out to 60 kpc. They concluded that the concomitant decline in DLA metallicity with R_{\perp} implies that the Δv_{90} -metallicity relation does not depend on impact parameter. Meanwhile, the analysis by Fu et al. (2021) found a CGM absorber detected at $R_{\perp} = 93$ kpc from a host SMG that deviates from the Δv_{90} -metallicity relation for DLAs by $+230 \text{ km s}^{-1}$ (see Figure 16). Fu et al. (2021) argued that this could be due to the large impact parameter of the sightline relative to typical DLA host galaxy projected separations.

To investigate the nature of the offsets from the Δv_{90} -metallicity relation within our sample, we focus on those systems with secure metallicity measurements (including one DLA-CGM sightline, the one SMG-CGM sightline, and all seven QSO-CGM sightlines). We use Δv_{90} and the system redshifts to calculate the expected DLA metallicities from the relation reported in Neeleman et al. (2013). We then compare the expected value to the measured metallicities by calculating the difference divided by the $\pm 1\sigma$ scatter in the Δv_{90} -metallicity relation ($|Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}}$), and relate that to R_{\perp} for each system. These latter values fall across the range $40 \text{ kpc} < R_{\perp} < 180 \text{ kpc}$.

The median value of $|Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}}$ is 0.02, with the SMG-CGM sightline having the maximum value of 1.08. Excluding this sightline, all eight values of this offset are < 0.6 , and are therefore consistent with the DLA Δv_{90} -metallicity relation within $\pm 1\sigma$. We also find no evidence for a dependence of this offset on R_{\perp} : the four sightlines within 100 kpc have values in the range $0.05 < |Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}} < 1.08$, while at large impact parameters (> 100 kpc) we measure values $0.02 < |Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}} < 0.19$.

While this analysis is limited, it is nevertheless suggestive that there is no clear relation between $|Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}}$ and R_{\perp} over the impact parameter range $40 \text{ kpc} < R_{\perp} < 180 \text{ kpc}$. We interpret this result as being consistent with the conclusions of Christensen et al. (2019) and Møller & Christensen (2020), as discussed above. Their analysis implied that the Δv_{90} -metallicity relation for DLAs is local in nature, meaning that the metallicity measured along the line of sight follows the local value of Δv_{90} (i.e., the local gravitational potential). Our results may similarly suggest that on average, CGM gas detected at much larger impact parameters than previously explored follows the same mean Δv_{90} -metallicity relation. However, a larger sample size is needed to confidently rule out a relation between $|Z/Z_{\odot, \text{observed}} - Z/Z_{\odot, \text{expected}}|/\sigma_{Z/Z_{\odot}}$ and R_{\perp} over large scales.

Our comparison of metallicities between DLA-CGM and QSO-CGM sightlines has strengthened existing lines of evidence, indicating that Δv_{90} is correlated with halo mass. Under this assumption, given that (1) in the foregoing analysis, we have used the measured Δv_{90} width of individual sightlines to calculate the expected metallicity, and that (2) CGM gas metallicities are likely to span a larger range than the metallicities of DLA material, we suggest that any deviation from the Δv_{90} -metallicity relation for CGM sightlines is due to the scatter in metallicities extant in the CGM at fixed halo

mass. While DLAs are likely to be confined to star-forming regions or the “inner” CGM (e.g., Krogager et al. 2017; Stern et al. 2021; Theuns 2021), the cool circumgalactic medium at this epoch is fed by both pristine inflow and metal-enriched outflow (e.g., Crighton et al. 2013, 2015; Fumagalli et al. 2017; Lofthouse et al. 2020), and it is likely that many of these gaseous structures are destroyed by hydrodynamical instabilities before they can mix with the surrounding material (Schaye et al. 2007; Crighton et al. 2015; Gronke et al. 2022). This interpretation could also explain the slight deviation of the Fu et al. (2021) system from the Δv_{90} -metallicity relation. Based on its relatively low metallicity, Fu et al. (2021) argued that this sightline is probing a cold inflowing stream around the host SMG; therefore, while the Δv_{90} width may fall in the range expected for a massive halo typically associated with SMGs, the metallicity in this particular line of sight may scatter away from the Δv_{90} -metallicity relation due to stochasticity in the contents of the CGM.

In summary, we have found that optically thick CGM absorbers around both DLAs and QSO hosts occupy a similar region of Δv_{90} -metallicity parameter space as DLAs themselves. In addition, we find no evidence for a relationship between the offset from the Δv_{90} -metallicity relation and R_{\perp} (though our sample is small). Instead, we propose that circumgalactic systems that are outliers from this relation arise due to poorly-mixed halo material.

7. Conclusion

In previous work, we introduced a technique to study the CGM of high-redshift DLAs (which act as signposts for high-redshift galaxies) in absorption using close quasar pairs. This technique probes gas in the extended CGM in one QSO sightline and permits a direct comparison to the ISM/inner CGM material traced by the DLA in the second sightline. In this work, we have analyzed medium and high-resolution ($R \geq 4000$) spectroscopy for 32 such quasar pairs, each of which has an intervening DLA in the redshift range $1.6 < z_{\text{DLA}} < 3.5$, and which have sightline separations in the range $24 \text{ kpc} < R_{\perp} < 284 \text{ kpc}$. We have reported column densities and assessed kinematics for several ionic species in each sightline pair. We have also performed a novel comparison between the metallicities of DLAs and nearby circumgalactic material.

Here, we summarize the main results of this paper:

1. We find that the HI column densities measured in our CGM sightlines are anticorrelated with R_{\perp} with 99.8% confidence. We also report a high incidence of optically thick HI ($N_{\text{HI}} > 10^{17.2} \text{ cm}^{-2}$) around DLAs, with $C_f = 50 \pm 13\%$ for sightlines within $24 \text{ kpc} \leq R_{\perp} \leq 100 \text{ kpc}$, and $C_f = 50 \pm 22\%$ for sightlines at $100 \text{ kpc} < R_{\perp} \leq 200 \text{ kpc}$. These results suggest both that DLAs are located close to the centers of their host halos and that these systems are located in environments in which neutral gas extends over large scales ($>100 \text{ kpc}$). The CGM of Lyman Break Galaxies (LBGs) at $z \sim 2-3$ exhibits a marginally lower incidence of optically thick HI at $50 \text{ kpc} \leq R_{\perp} \leq 100 \text{ kpc}$ ($C_f = 20^{+15}_{-13}\%$; Rudie et al. 2012).
2. We report covering fractions within 100 kpc-wide R_{\perp} bins for several metal species. C II and C IV yield the largest covering fractions within $R_{\perp} < 100 \text{ kpc}$, with $C_f(N_{\text{CII}} > 10^{13} \text{ cm}^{-2}) > 89\%$ and $C_f(N_{\text{CIV}} > 10^{13} \text{ cm}^{-2}) = 94^{+4}_{-9}\%$.

We likewise find high incidences of singly- and triply-ionized silicon, with $C_f(N_{\text{SiII}} > 10^{13} \text{ cm}^{-2}) = 75^{+12}_{-17}\%$ and $C_f(N_{\text{SiIV}} > 10^{13} \text{ cm}^{-2}) = 67^{+10}_{-12}\%$. Comparing to LBG and QSO-host halos at this epoch, we find that the covering fractions of high-ionization species are similar (consistent within $\pm 1\sigma$) within $R_{\perp} < 100 \text{ kpc}$. However, the covering fractions for C II and Si II around DLAs are larger than around LBGs at $0 \text{ kpc} < R_{\perp} < 100 \text{ kpc}$ by $\sim 2\sigma$. DLA covering fractions for all species drop below those measured in QSO halos by $\gtrsim 1\sigma$ beyond $R_{\perp} > 200 \text{ kpc}$, suggesting that QSO halos have enriched gas that extends to larger impact parameters.

3. We identify, species by species, the thresholds above which 90% of metal-line column densities in our DLA sample lie (i.e., the 10th percentile column density values). We then assess the covering fraction of CGM systems with column densities above these thresholds. At $24 \text{ kpc} < R_{\perp} < 200 \text{ kpc}$, both intermediate- and high-ionization species exhibit covering fractions $>40\%$ relative to the corresponding 10th percentile thresholds; however, even within $<100 \text{ kpc}$, the incidence of low-ionization species does not exceed $C_f = 30\%$. This suggests that the warm material traced by Si IV and C IV associated with DLAs frequently extends over 100–200 kpc scales, whereas cool, photoionized or neutral material seldom exhibits DLA-level absorption strengths across length scales $\gtrsim 30 \text{ kpc}$.
4. We identify nine DLA-CGM systems, all having impact parameters $24 \text{ kpc} < R_{\perp} < 120 \text{ kpc}$, for which the ionic ratios $N_{\text{OI}}/N_{\text{HI}}$ or $N_{\text{SiII}}/N_{\text{HI}}$ yield robust constraints or limits on the CGM metallicity (given our assumption that these ratios are less sensitive to ionization state in systems with $N_{\text{HI}} \gtrsim 10^{18.5} \text{ cm}^{-2}$; e.g., Crighton et al. 2013). These values range from a maximum of $\log Z/Z_{\odot} = -0.75$ (i.e., close to that observed in the ISM of LBGs) to an upper limit $\log Z/Z_{\odot} < -2.06$ (consistent with that observed in the IGM at this epoch). These metallicities are consistent with or lower than those estimated for the associated DLA (using the ionic ratio $N_{\text{SiII}}/N_{\text{HI}}$) in five of these sightline pairs, and are higher than that of the associated DLA in two systems. Overall, we find no evidence for a correlation between the metallicities observed in the DLAs and their associated CGM.
5. The Δv_{90} velocity widths of both low-ionization transitions and the C IV $\lambda 1548$ transition in the DLA versus CGM sightlines lie along a 1:1 relation and are correlated with high statistical significance. Moreover, the low-ionization Δv_{90} values differ by $<40 \text{ km s}^{-1}$ in the preponderance (66%) of sightline pairs. This suggests (1) that metal-line kinematic widths exhibit strong coherence over $R_{\perp} \lesssim 200 \text{ kpc}$ scales and (2) that the Δv_{90} kinematic measure is an effective indicator of the potential well of the host halo, regardless of the location of the sightline relative to the halo center. We find that our DLA-CGM systems, along with several CGM systems drawn from the literature, lie along the same Δv_{90} -metallicity relation as that exhibited by DLAs themselves.
6. Velocity centroids for C IV $\lambda 1548$ differ by $>100 \text{ km s}^{-1}$ for nearly half (48%) of the velocity components across the sightline pairs. However, the vast majority (32/35) of C IV and low-ion component centroids have line-of-sight velocities that are less than the escape velocity of a

putative DLA host dark matter halo with $M_h \approx 10^{12} M_\odot$. If we instead assume a host halo mass of $M_h \approx 10^{11} M_\odot$, then 18 of these 35 components have radial velocities less than the implied escape velocity.

The CGM, while crucial to our understanding of galaxy evolution, is diffuse and difficult to detect at $z \gtrsim 2$. As demonstrated by the foregoing analysis, absorption-line studies remain important tools for assessing the enrichment histories of the gas that feeds star formation during this critical epoch. Ongoing efforts to improve our understanding of the baryonic cycling through circumgalactic environments will include increasing the samples of known galaxy counterparts to strong absorbers at high redshift. One such effort is the MUSE Analysis of Gas around Galaxies (MAGG) survey, which uses VLT/MUSE to search for galaxy counterparts in emission in several bright quasar fields (see Dupuis et al. 2021; Lofthouse et al. 2020; Fossati et al. 2021; Lofthouse et al. 2023). The increasing sample of confirmed galaxy/CGM sightline pairs will lead to a better understanding of the connection between the properties of host galaxy/halo centers (i.e., star formation rates, halo mass) and those of extended halo gas. At the same time, a comparison of the measured CGM properties in these environments to the predictions of state-of-the-art cosmological zoom simulations (e.g., FIRE or FOGGIE; Hopkins et al. 2018; Peebles et al. 2019; Stern et al. 2021) will provide a unique insight into the physics of the baryonic flows feeding such extended gas reservoirs. Ultimately, these combined efforts will aid our understanding of the complex structure of the CGM and

provide an insight into how halo gas may be linked to different phases of galaxy growth.

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The authors wish to recognize and acknowledge the very significant cultural role and reverence that the summit of Maunakea has always had within the indigenous Hawaiian community. We are most fortunate to have the opportunity to conduct observations from this mountain.

Appendix A Constraining Metallicities

A.1. Constraints on Ionization State from Ionic Ratios

To estimate the systematic uncertainty in our metallicity measurements, to be described in Appendix A.2, we must

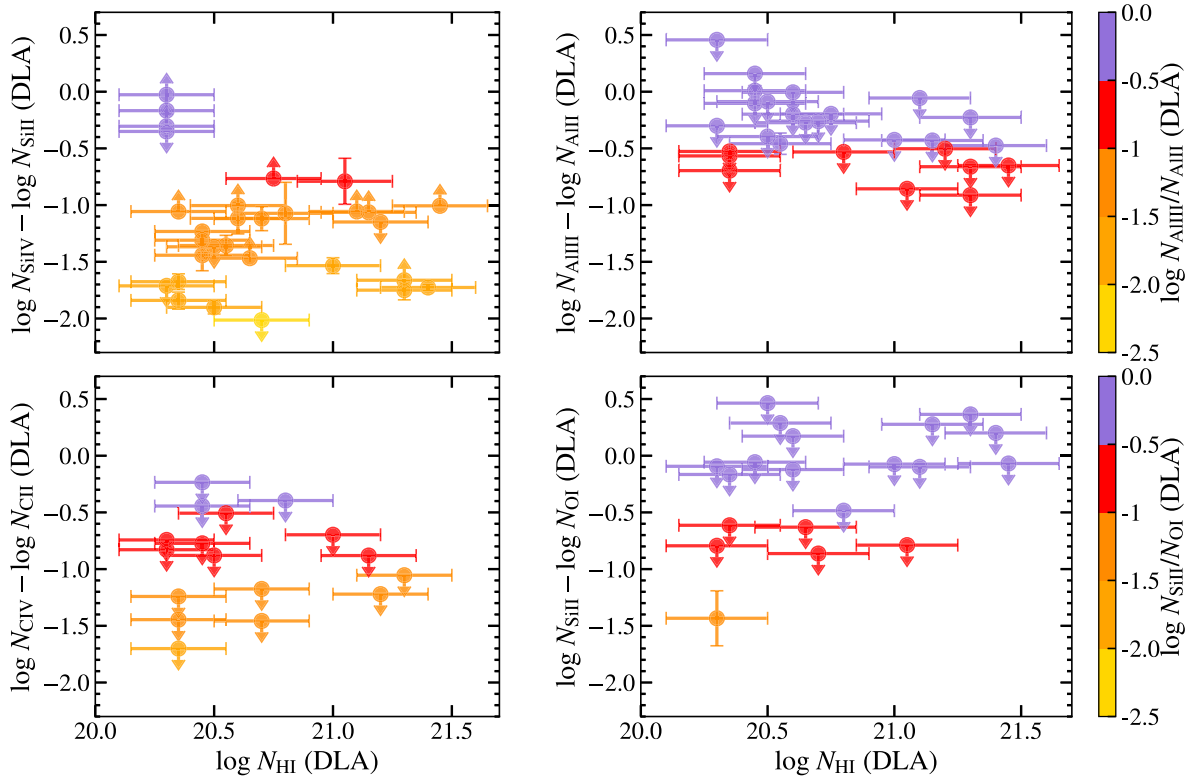


Figure 17. Ionic ratios of $N_{\text{SiIV}}/N_{\text{SiII}}$ (top left-hand panel), $N_{\text{CIV}}/N_{\text{CII}}$ (bottom left-hand panel), $N_{\text{AlIII}}/N_{\text{AlII}}$ (top right-hand panel), and $N_{\text{SiII}}/N_{\text{OI}}$ (bottom right-hand panel) vs. N_{HI} for the DLAs in our sample. The color bar reflects the y-axis value of each point in each panel and will be used to color code these systems in Figure 19. Larger values are indicative of more highly ionized material and imply larger uncertainties in metallicities estimated from ionic ratios. Our constraints on N_{OI} , N_{AlII} , and N_{AlIII} are upper limits in the majority of sightlines, resulting in the vast majority of ionic ratios including these species appearing here as limits. C II is also typically saturated in DLA sightlines, such that many of our $N_{\text{CIV}}/N_{\text{CII}}$ values are upper limits. We therefore rely on ionic ratios calculated from $N_{\text{SiIV}}/N_{\text{SiII}}$ to indicate the ionization state of our DLA sightlines.

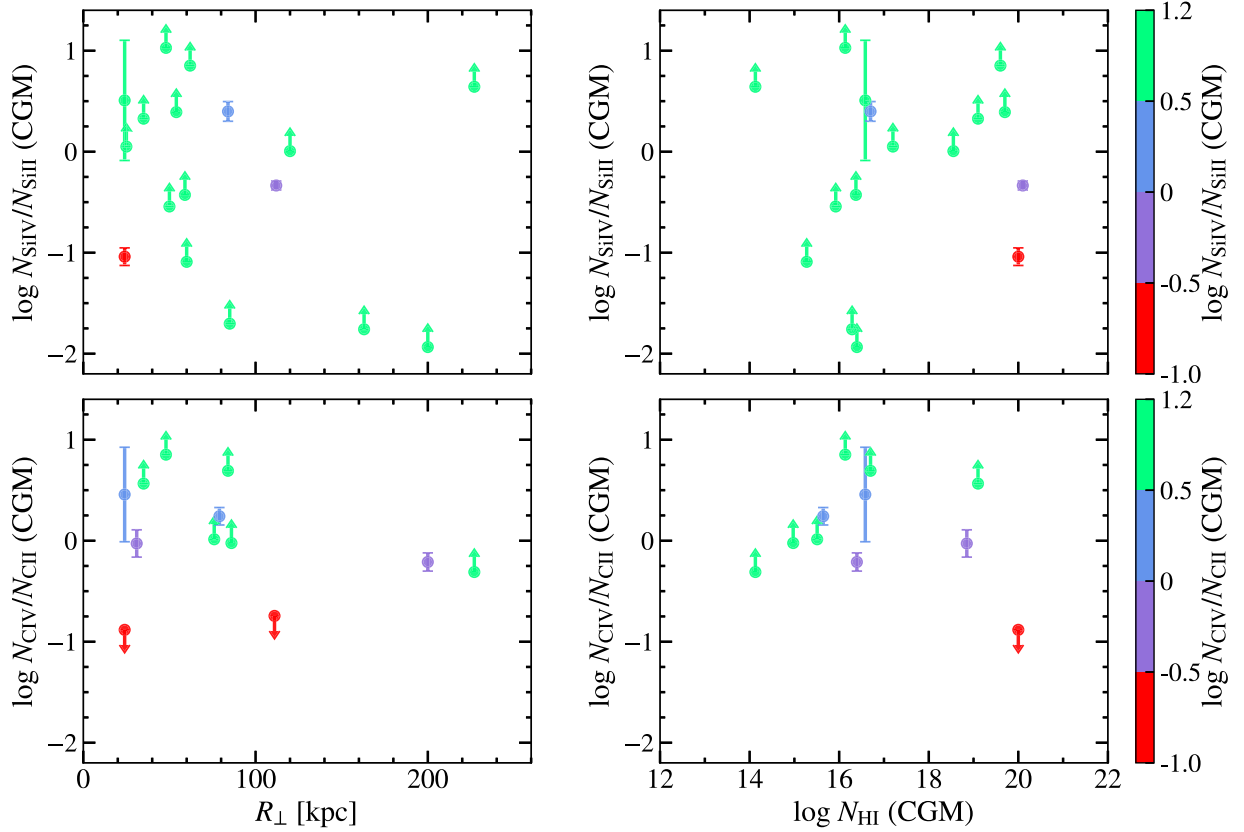


Figure 18. Ionic ratios for the CGM sightlines in this work. $\log N_{\text{SiIV}}/N_{\text{SiII}}$ and $\log N_{\text{CIV}}/N_{\text{CII}}$ vs. R_{\perp} constraints are shown at top left-hand and bottom left-hand, respectively. The ionic ratios $\log N_{\text{SiIV}}/N_{\text{SiII}}$ and $\log N_{\text{CIV}}/N_{\text{CII}}$ vs. $N_{\text{HI,CGM}}$ are shown at top right-hand and bottom right-hand, respectively. Lower limits are shown in green because the value of the corresponding ionic ratio is ambiguous. As in our DLA sample, our $\log N_{\text{SiIV}}/N_{\text{SiII}}$ measurements yield a larger number of direct constraints than other ionic ratios, and we therefore rely on these measurements as our indicator of ionization state for the CGM.

first assess the ionization state along each sightline. We approach this by calculating the ratios of the column densities of high-ionization to low-ionization transitions of the same species. Logarithmic ratios of $N_{\text{SiIV}}/N_{\text{SiII}}$, $N_{\text{CIV}}/N_{\text{CII}}$, and $N_{\text{AlIII}}/N_{\text{AlII}}$ versus N_{HI} for our DLA sample are shown in Figure 17. We also include $N_{\text{SiII}}/N_{\text{OI}}$ because this ratio should yield -1.2 dex in neutral gas (assuming solar abundance ratios). If this ratio is above -1.2 dex, then this suggests the material is highly ionized (Prochaska et al. 2015).

The measured ionic ratio constraints range from -2.5 dex to $+0.5$ dex; however, the vast majority of our constraints on $N_{\text{CIV}}/N_{\text{CII}}$, $N_{\text{AlIII}}/N_{\text{AlII}}$, and $N_{\text{SiII}}/N_{\text{OI}}$ are upper limits. Our $N_{\text{SiIV}}/N_{\text{SiII}}$ estimates included the greatest number of direct measurements, and we therefore rely on this ionic ratio as our primary indicator of ionization fraction. Several of our DLAs exhibit $\log N_{\text{SiIV}}/N_{\text{SiII}}$ values $\gtrsim -0.5$ dex (shown in purple in Figure 17), which are indicative of partially ionized conditions (Prochaska et al. 2015). However, the average $\log N_{\text{SiIV}}/N_{\text{SiII}}$ value for these systems is ≈ -1.1 dex, which is consistent with the predominantly neutral conditions of most DLAs (e.g., Prochaska & Wolfe 2000).

We also calculate $N_{\text{SiIV}}/N_{\text{SiII}}$ and $N_{\text{CIV}}/N_{\text{CII}}$ ionic ratios for our CGM sightlines, shown in Figure 18. Here, we color any lower limits as green. As for our DLA sightlines, our $N_{\text{SiIV}}/N_{\text{SiII}}$ estimates include the greatest number of direct

measurements. All except two sightlines exhibit ionic ratios $\gtrsim -0.5$ dex, which is indicative of predominantly ionized conditions.

A.2. Metallicity Constraints for the DLA Sample

As explained in Section 5.4, to estimate the metallicity of our absorption systems, we make use of the quantity ($\{X_i/H_j\}$) introduced by (Prochaska et al. 2015, see Equation (5)). In cases in which the ionization correction is small, such as for DLAs, we assume that $\{X_i/H_j\} = [X/H]$. We adopt solar elemental abundances from Asplund et al. (2009).

In Figure 19, we show $\{\text{SiII}/\text{HI}\}$, $\{\text{FeII}/\text{HI}\}$, $\{\text{CII}/\text{HI}\}$, and $\{\text{OI}/\text{HI}\}$ versus N_{HI} for our DLA sample. Points are color-coded by the value of $\log N_{\text{SiIV}}/N_{\text{SiII}}$ for the system (as in Figure 17). This latter ionic ratio is low in the vast majority of these sightlines (< -0.5 dex), implying that ionization corrections are small and that the former ratios are indeed representative of the metallicity. Our values of $\{\text{SiII}/\text{HI}\}$ range from ~ -2.7 to -0.4 dex and have an average of -1.35 dex. The range of $\{\text{FeII}/\text{HI}\}$ values is similar (~ -2.8 to -0.7 dex, with an average of -1.76 dex). These metallicities are typical of those exhibited by unbiased DLA samples at this epoch (see Figure 16; Neeleman et al. 2013).

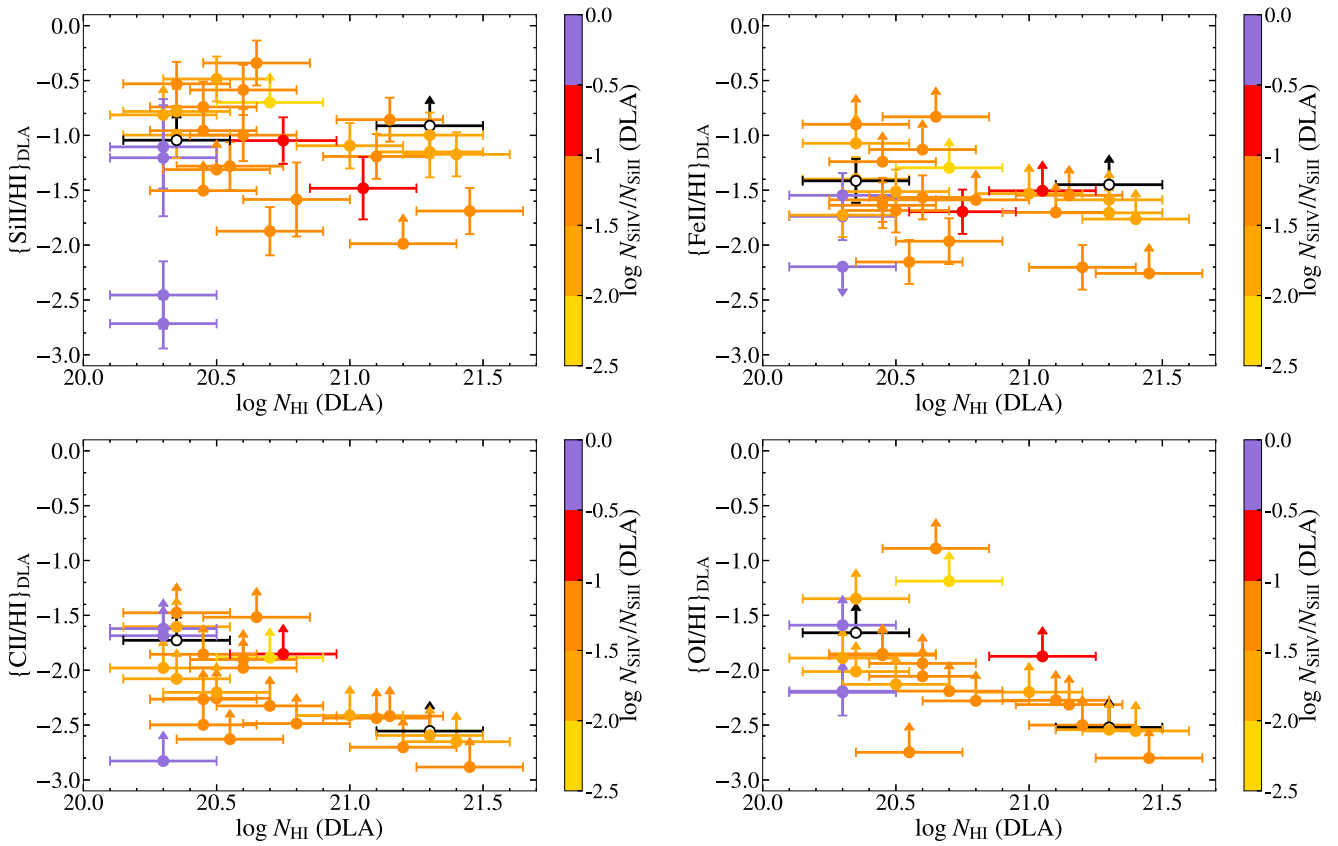


Figure 19. Ionic ratio metallicities estimated from Si II (top left-hand panel), Fe II (top right-hand panel), C II (bottom left-hand panel), and O I (bottom right-hand panel) column densities plotted vs. H I column density for our DLA sample. The color bar shows the value of the ionic ratio $\log N_{\text{SiIV}}/N_{\text{SiII}}$, which is an indicator of the ionization fraction. Purple points mark systems with $\log N_{\text{SiIV}}/N_{\text{SiII}} > -0.5$ dex, i.e., systems which are predominantly ionized and for which our ionic ratio metallicities are likely overestimated. The black open circles mark sightlines with an ambiguous ionization fraction due to heavy blending of the Si IV line profiles. In general, the measured DLA metallicities are consistent with the Δv_{90} -metallicity relation from (Neeleman et al. 2013, see Figure 16).

Appendix B

Comments on Individual Systems: Physical Origins

To gain insight into the origins of the material detected along our CGM sightlines, we refer to the detailed analysis of the properties and origins of circumgalactic material in the FIRE-2 cosmological zoom simulations described in Hafen et al. (2019). This work used particle tracking to distinguish halo gas that has accreted from the IGM, that originates from an outflow from the host galaxy, or that originates from gas ejected from another galaxy. They performed this analysis for a suite of simulations of halos that evolve to have $M_h \sim 10^{10-12} M_\odot$ at low redshift, assessing the median metallicity as a function of radius for each of these three CGM components at both $z=2$ and $z=0.25$. Overall, these authors find that the material originating in winds on average exhibits only slight changes in metallicity with increasing radial distance, that wind material from the central galaxy on average has higher metallicity than winds from satellites, and that material accreting from the IGM exhibits metallicities $\lesssim 0.01 Z_\odot$ at distances $> 0.5 R_{\text{vir}}$. We caution that these statements apply to the median metallicities as a function of radius, and therefore can only be used to indicate the most statistically likely origins of the gas in our sightlines.

B.1. An Unusually Low-Metallicity CGM System: J0004-0844

One of our CGM systems (in sightline pair J0004-0844 with $R_\perp = 35$ kpc and $z_{\text{abs}} = 2.75$), indicated with the purple square in Figure 15, has a metallicity $Z/Z_\odot < 10^{-2.09}$, i.e., within the range of metallicities measured in the $z \sim 2$ IGM (Schaye et al. 2003;

Simcoe et al. 2004), and consistent with the lowest-metallicity $z \sim 2$ CGM sightlines probing distances $R_\perp > 100$ kpc from their galaxy counterparts (Simcoe et al. 2006; Mackenzie et al. 2019; Lofthouse et al. 2020). This sightline has a large H I column density ($N_{\text{HI}} > 10^{19.1} \text{ cm}^{-2}$) but exhibits no detectable absorption from low or intermediate metal ions. Our upper limit on the column density of O I places a firm limit on the metallicity of this system at approximately two orders of magnitude lower than that measured for the associated DLA ($Z/Z_\odot = 10^{-0.1 \pm 0.24}$). The average velocity of this material as traced by C IV absorption is within $< 10 \text{ km s}^{-1}$ of that of the DLA, which suggests that this material is gravitationally bound, even under the assumption that the system is hosted by a $10^{10} M_\odot$ halo (see the purple square in Figure 14).

Given that almost none of the gas originating in winds in the FIRE-2 simulations at $z=2$ has a metallicity as low as our limit for this system (see Figure 17 in Hafen et al. 2019), as well as the quiescent kinematics of the observed absorption profile, we posit that this system most likely arises from accreting IGM material. Moreover, we emphasize that our detection of such a low-metallicity system is likely to be only possible due to the absence of material originating in galactic winds along this sightline.

B.2. A High-Metallicity, High Velocity-Offset CGM System: J2103+0646

The CGM system detected toward sightline pair J2103+0646 (represented by the gold square in Figure 15) is one of

the most metal-enriched in our sample, with an enrichment level near that of the ISM of the lensed LBG cB58 (Pettini & Cooke 2012). The metallicity of this system is also an order of magnitude above the metallicity of the corresponding DLA despite being at $R_{\perp} = 31$ kpc. Considering that at this redshift ($z_{\text{abs}} = 2.14$) the virial radius is likely $R_{\text{vir}} < 90$ kpc for the host galaxy, this implies that this sightline probes gas at $\gtrsim 0.30 R_{\text{vir}}$. The Hafen et al. (2019) analysis of $M_h \sim 10^{12} M_{\odot}$ halos at such small radial distances predicts that both accreted IGM material and gas originating from winds may plausibly be enriched to this level in this inner CGM environment.

The CGM sightline has two velocity components: one close to the DLA H I velocity centroid (the “primary”) and one close to the CGM H I velocity centroid (the “secondary”). The primary component has a low N_{CIV} compared to the secondary component, and therefore does not contribute significantly to the total metallicity of the sightline. The C IV profile of the secondary component is offset from the DLA redshift by 263 km s^{-1} , and thus traces gas that would not be gravitationally bound to a halo with $M_h \leq 10^{11} M_{\odot}$. However, the Δv_{90} width of the C IV absorption is low (45 km s^{-1}), suggesting this sightline does not probe a region of the CGM that is being actively enriched by outflows. The primary velocity component has a velocity centroid that is offset by only 32 km s^{-1} from the DLA redshift, and thus likely traces gas that remains gravitationally bound to its host halo.

B.3. Systems that Must Be Wind-Enriched, But Which Have Quiescent Kinematics

The CGM system observed toward sightline pair J0955-0123 (represented by the dark blue square in Figure 15) has a redshift $z_{\text{abs}} = 2.73$ and a projected separation $R_{\perp} = 89$ kpc. At this distance and redshift, this sightline likely probes gas near or beyond the virial radius of a galaxy in a $\geq 10^{12} M_{\odot}$ halo. This system has a high H I column density ($N_{\text{HI}} = 10^{19.2} \text{ cm}^{-2}$) and a metallicity $Z/Z_{\odot} = 10^{-1.43}$. This level of enrichment is for the most part only achieved in the FIRE-2 simulations by wind-enriched gas (Hafen et al. 2019). The median metallicity of material accreted from the IGM and located at the virial radius is

less than $3 \times 10^{-3} Z_{\odot}$ for all simulated halos, even with metal diffusion enriching the IGM particles, i.e., it is more than a factor of ten lower than that measured in the CGM of this system. We therefore suggest this gas likely probes wind-enriched material. The velocity centroid of the low-ionization gas along this sightline is close to that of the DLA ($\delta v_{\text{weight}} = -43 \text{ km s}^{-1}$), and its Δv_{90} width is narrow. The high-ionization gas traced by C IV is similarly near the DLA redshift ($\delta v_{\text{weight}} = -36 \text{ km s}^{-1}$); however, it has a large Δv_{90} width ($\Delta v_{90} = 178 \text{ km s}^{-1}$). This gas is likely bound within halos with masses $\geq 10^{10} M_{\odot}$.

The CGM system in our pair target J2146-0752, with a metallicity $Z/Z_{\odot} = 10^{-0.96}$, is represented by the red square in Figure 15. This sightline probes the CGM of a DLA at $z_{\text{abs}} = 1.85$ at an impact parameter $R_{\perp} = 120$ kpc. Therefore, this particular sightline is likely to be near or beyond the virial radius of typical DLA galaxies at this redshift. The median metallicity of IGM material at this redshift and proper distance as predicted by Hafen et al. (2019, $\leq 3 \times 10^{-3} Z_{\odot}$) is 1.6 dex lower than what we measure along the sightline, suggesting that this gas likely originated from a wind of some type. The CGM absorption of this system is furthermore unusual in that the strengths of several metal transitions are comparable, and in some cases stronger, than those observed in the associated DLA sightline, while at the same time it has a substantially lower H I column density by approximately two orders of magnitude. However, its metallicity is only 0.2 dex below that of the DLA sightline. We observe velocity centroids of $\delta v_{\text{weight}} = 90 \text{ km s}^{-1}$ for low-ionization gas and $\delta v_{\text{weight}} = 59 \text{ km s}^{-1}$ for C IV along this sightline, along with Δv_{90} widths of 87 km s^{-1} and 131 km s^{-1} for low- and high-ionization material, respectively. As with the J2103+0646 system discussed earlier, these quiescent kinematics suggest that although the metallicity of this system implies that it has been enriched by winds, our sightline is not probing the energetics associated with this enrichment.

Appendix C Additional Figures

We include a set of figures showing our spectroscopic coverage of Ly α and several metal-line transitions for all DLA

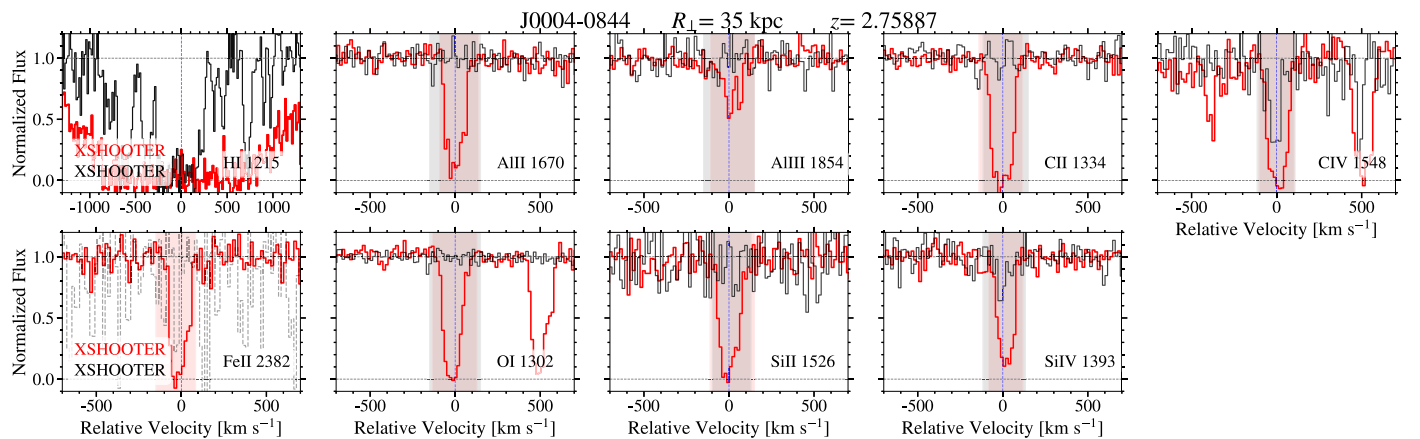


Figure 20. H I and metal-line absorption profiles for all 32 DLA-CGM sightline pairs in our sample. Each set of nine panels shows our coverage of Ly α , Al II, Al III, C II, C IV, Fe II, O I, Si II, and Si IV transitions associated with an individual DLA (red histogram) and the corresponding CGM system (black histogram). The QSO pair ID is shown at the top of each set of nine panels, along with the projected distance between the sightlines at the DLA redshift. The blue-dotted line indicates the systemic velocity of the DLA. The shaded regions indicate the placement of the velocity windows used to measure metal-line absorption strength. In the case of profiles with multiple components, the shaded regions are marked “2” or “3” to indicate a second or third component. The instrument that was used for each sightline is labeled in the bottom left-hand corner of each H I panel. Transitions which are not used in this analysis due to extreme blending are shown with dotted histograms. The complete figure set (32 images) is available in the online journal.

(The complete figure set (32 images) is available.)

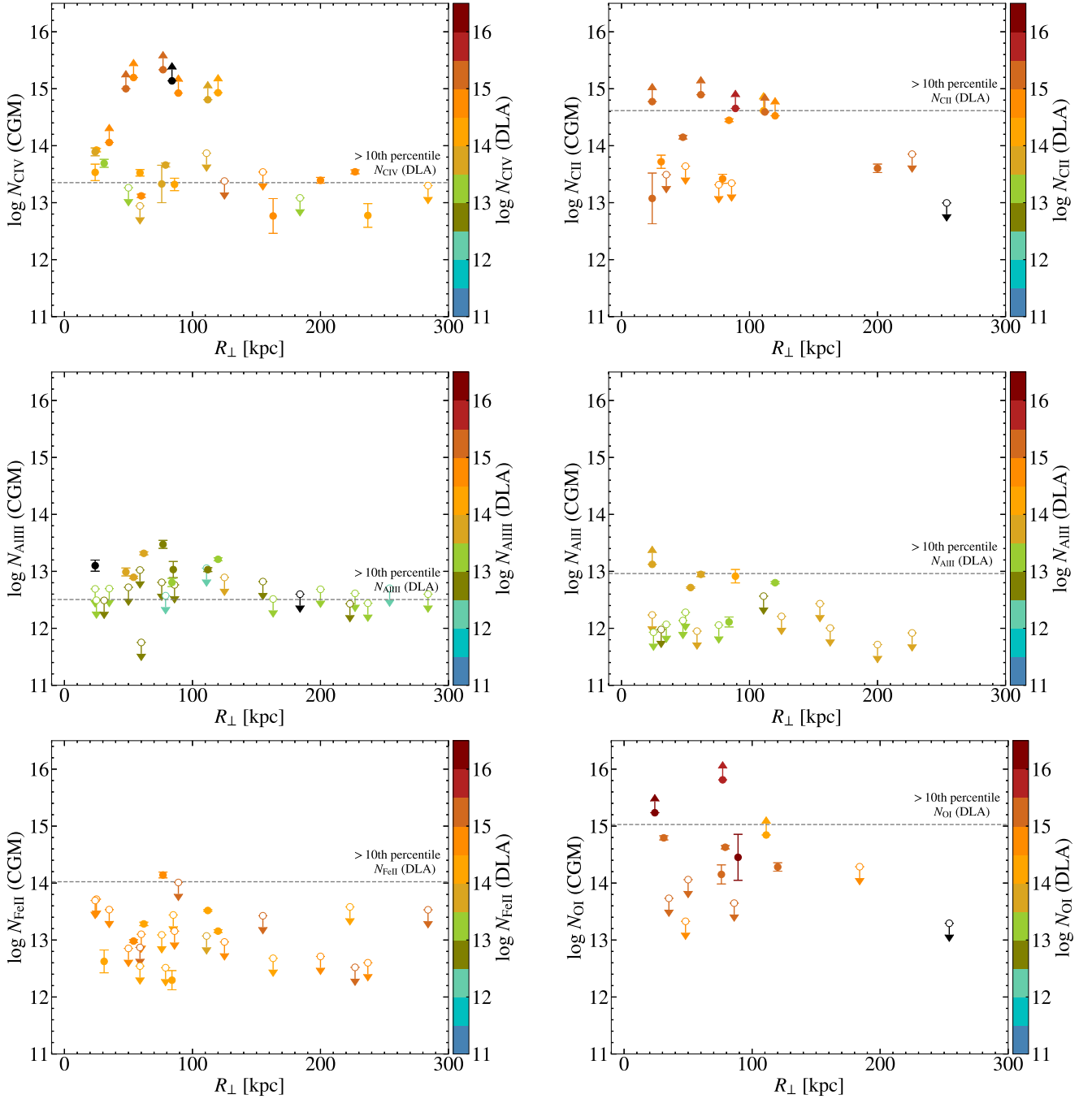


Figure 21. Column densities of C IV (top left-hand panel), C II (top right-hand panel), Al III (middle left-hand panel), Al II (middle right-hand panel), Fe II (bottom left-hand panel), and O I (bottom right-hand panel) in our CGM sightlines vs. proper distance (R_{\perp}). Colors represent the corresponding DLA column density for that ion. Black points indicate ambiguous DLA column density values. Open symbols indicate that our constraint on the CGM sightline column is an upper limit. The horizontal-dashed lines represent the threshold above which 90% of metal-line column densities for the DLA sightlines fall.

and CGM sightlines. An example is shown in Figure 20. Figure 21 shows the column densities of several metal ions measured along our CGM sightlines versus proper distance (R_{\perp}).

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