

## Ionic Conductivity, Salt Partitioning, and Phase Separation in High-Dielectric Contrast Polyether Blends and Block Polymer Electrolytes

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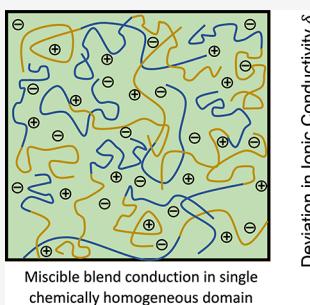
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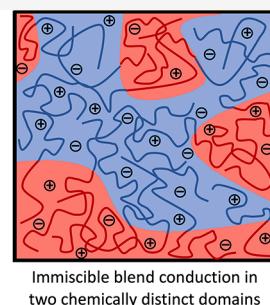
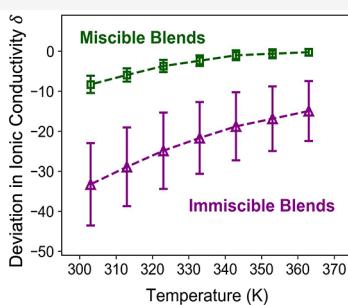
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Miscible blend conduction in single chemically homogeneous domain



Immiscible blend conduction in two chemically distinct domains

**ABSTRACT:** Small-molecule battery electrolytes are composed of mixtures of high-polarity and low-viscosity solvents at compositions that optimize ionic conductivity. In this work, we examined analogous polymer blends composed of one component with rapid segmental dynamics to provide low viscosity and another with a high dielectric constant ( $\epsilon$ ) to enhance ion dissociation. We investigated the inherent tradeoff between polymer polarity and segmental dynamics limiting ionic conductivity through the analysis of ionic conductivity of electrolytes containing lithium bis(trifluoromethanesulfonyl)imide (LiTFSI) with different polarity hosts: poly(allyl glycidyl ether) (PAGE) ( $\epsilon \approx 9$ ), poly[(cyanoethyl glycidyl ether)-*co*-(*n*-butyl glycidyl ether)] [P(CEGE-*co*-nBGE)] ( $\epsilon \approx 24$ ), and poly(cyanoethyl glycidyl ether) (PCEGE) ( $\epsilon \approx 36$ ). Two high-polarity-contrast polymer blends, PAGE/P(CEGE-*co*-nBGE)/LiTFSI and PAGE/PCEGE/LiTFSI, were prepared. While PAGE/PCEGE/LiTFSI blends were immiscible at all compositions, PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends were miscible at LiTFSI concentrations above  $r = 0.065$ . The immiscibility of PAGE/PCEGE/LiTFSI blends imposed a negative deviation in ionic conductivity from a calculated linear average of the two single-polymer electrolytes. This negative deviation was decreased in magnitude to less than 10% in miscible PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends between 30 and 90 °C. To understand the changes in the effective interaction parameter in the presence of LiTFSI, we investigated the disordered-state small-angle X-ray scattering (SAXS) of a diblock polymer, PAGE-*b*-PCEGE, across a range of LiTFSI concentrations. By fitting SAXS profiles of this copolymer using the random phase approximation and an adjustable contrast model, we found that the effective interaction parameter decreased monotonically as the LiTFSI concentration increased. At low concentrations, LiTFSI was primarily solvated in the high-polarity PCEGE-rich phase.

### INTRODUCTION

Due to an increased deployment of intermittent renewable energy sources combined with the electrification of the transportation and chemical industries, the need has emerged for advances in materials for energy storage.<sup>1,2</sup> Electrolytes represent a major component in batteries, facilitating ion transport between the cathode and anode. Widely used small-molecule electrolytes (SMEs), composed of a cyclic carbonate (high-polarity component) and a linear ether or carbonate (low-viscosity component), have optimum conductivities because they balance both rapid molecular dynamics and high ion dissociation.<sup>3,4</sup> There is a continuing interest in replacing SMEs with nonvolatile polymer electrolytes that offer advantages in enabling new battery chemistries, mitigating dendrite growth, increasing energy density, and maximizing the transference number of the electrolyte.<sup>5–7</sup> However, most

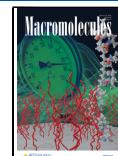
polymer electrolytes contain a single polymer or block copolymer with a non-conducting structural component and ion-conducting component in combination with lithium salt, in contrast with current SME blends used in lithium-ion batteries today.<sup>8–14</sup>

Recent work has focused on understanding the mechanism of ion transport in polymer electrolytes<sup>8,15–18</sup> and the structure–property relationships governing ionic conductiv-

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ity.<sup>19–23</sup> We contributed to this research area through molecular dynamics simulations and experimental investigations of the influence of segmental dynamics and polarity of the host polymer on ion transport in polymer electrolytes.<sup>24–29</sup> We and others have concluded that poly(ethylene oxide) (PEO), a benchmark for polymer electrolytes, combined rapid segmental dynamics with optimized ionic solvation to enhance ionic conductivity at high temperatures (*ca.* 10<sup>−4</sup> S/cm at 70 °C).<sup>30–35</sup> Our simulations established two polarity regimes where ion dissociation and segmental dynamics were, respectively, limiting the ionic conductivity.<sup>25</sup> We corroborated these simulation predictions through experimental investigations on two series of polymer electrolytes based on low-dielectric poly(vinyl ether)s and higher-dielectric poly(glycidyl ether)s. The poly(vinyl ether)-based electrolytes had glass transition temperatures ( $T_g$ s) significantly below room temperature and low dielectric constants ( $\epsilon$ )—a measurement of polarity—in the range of 1.3 to 9.0.<sup>28</sup> In this low-polarity regime, ionic conductivity increased monotonically as the dielectric constant increased due to the increasing dissociation of ions in the polymer hosts. However, in our latest work on poly(glycidyl ether)s with dielectric constants in between 9.0 and 35, we revealed that a further increase in the dielectric constant of host polymers inevitably leads to stronger interchain interactions, slower segmental dynamics, and therefore decreased ionic conductivity.<sup>29</sup> Such an inherent tradeoff between polymer polarity and segmental dynamics limits ionic conductivity for single-polymer–host electrolytes and presents PEO as the optimal neutral polymer host for high temperature applications.

Inspired by similar tradeoffs in SMEs, we hypothesized that blending two polymers, one with rapid segmental dynamics (low viscosity) and another with high polarity (strong ionic solvation), may be a viable strategy to optimize ionic conductivity. In earlier work, we approached this problem by modeling ternary polymer–polymer–salt blend electrolytes based on coarse-grained molecular dynamics simulations.<sup>27</sup> Our results revealed that at moderate- to high-polarity contrast, improving the miscibility of the polymer hosts enables ionic transport in these ternary electrolyte blends.

To experimentally realize polymer blend electrolytes, it is important to develop an in-depth understanding of the lithium salt solvation in polymer blends and the influence of adding lithium salts on polymer–polymer interactions. Current experimental studies on polymer blend electrolytes and block polymers typically combine PEO with a variety of polymers as the second component.<sup>36–39</sup> For instance, Nakamura et al. established a thermodynamic theory to describe the solvation of ions in PEO and polystyrene (PS) blends and block copolymers.<sup>40</sup> Ions were solvated in PEO-rich domains due to the lack of binding moieties in PS. Gao et al. reported nontrivial miscibility changes of PEO and poly(1,3,6-trioxocane) blends with lithium bis(trifluoromethanesulfonyl)-imide (LiTFSI).<sup>41</sup> Recently, Shah et al. demonstrated the effect of LiTFSI addition on poly(ethylene oxide)-*b*-poly(methyl methacrylate) (PEO-*b*-PMMA) copolymers.<sup>36</sup> Their results provided experimental validation for the ionic self-consistent field theory developed by Olvera de la Cruz et al.<sup>42,43</sup> At a theoretical level, Wang et al.<sup>40,44</sup> and Hall and co-workers<sup>8,15,16</sup> calculated the thermodynamics and ion correlations in polymer blends and diblock copolymers, demonstrating preferential solvation of ions in high-polarity polymers. Brown et al. applied classical fluid density functional theory to elucidate how the

microphase morphology of diblock copolymers is affected by salt.<sup>8</sup>

Many of the above studies are based on blends of a low dielectric constant polymer such as PS or PMMA ( $\epsilon_{\text{PMMA}} \approx 3$ ) with PEO ( $\epsilon_{\text{PEO}} \approx 9$ ), resulting in systems with low dielectric contrast ( $\Delta\epsilon = |\epsilon_1 - \epsilon_2| \approx 5–6$ ). In contrast, typical SMEs based on ethylene carbonate (EC,  $\epsilon_{\text{EC}} = 89.8$ ) and dimethyl carbonate (DMC,  $\epsilon_{\text{DMC}} = 3.1$ ) have a  $\Delta\epsilon$  of *ca.* 87. Therefore, in this study, we build upon the previous studies of polymer blends and block polymers and increase the dielectric contrast ( $\Delta\epsilon$ ) approaching values encountered in SME blends. Toward this objective, two blends with polarity contrasts of  $\Delta\epsilon \approx 17$  and 27 were prepared. Polymer hosts in these blends were both poly(glycidyl ether)s with a polyether backbone that can complex Li<sup>+</sup> ions. Although these polymers are still in the moderate polarity contrast regime when compared to the SMEs, this study probes the miscibility and ionic conductivity in the highest polarity contrast blend to date, providing system design insights that may help further optimize blend electrolytes for this application. The miscibility and the composition-dependent conductivities of these polymer blends with LiTFSI were studied. We further elucidated the change of the effective Flory–Huggins interaction parameter ( $\chi_{\text{eff}}$ ) in response to increased salt loadings in a high-polarity-contrast diblock polymer melt.

## EXPERIMENTAL SECTION

**Materials.** *N,N*-Dibenzyl-2-aminoethanol ( $\geq 98.0\%$ ), triethylaluminum (1.0 M in hexanes), triisobutylaluminum (1.0 M in hexanes), allyl glycidyl ether ( $\geq 99\%$ ), *n*-butyl glycidyl ether ( $\geq 99\%$ ), *m*-chloroperoxybenzoic acid (70–75% stabilized with water), allyl alcohol ( $\geq 99\%$ ), acrylonitrile ( $\geq 99\%$ ), sodium hydroxide ( $\geq 97.0\%$ ), lithium bis(trifluoromethanesulfonyl)imide, and tetrahydrofuran (THF) were all purchased from commercial sources and used as received. A mono( $\mu$ -alkoxo)bis(alkylaluminum) (MOB) polymerization initiator  $[(\text{Bn})_2\text{NCH}_2\text{CH}_2(\mu_2\text{-O})\text{Al}(\text{iBu})_2\cdot\text{Al}(\text{iBu})_3]$ <sup>45</sup> and 2-cyanoethyl glycidyl ether (CEGE)<sup>46</sup> were prepared as previously described and stored under a nitrogen atmosphere.

**Polymer Purification.** After polymerization, the reaction vial was taken out from the N<sub>2</sub>-filled glovebox and *ca.* 20 mL of CH<sub>2</sub>Cl<sub>2</sub> was added. The vial was left on a stirrer at room temperature. After 2–5 h, the polymer was dissolved in CH<sub>2</sub>Cl<sub>2</sub> to yield a homogenous solution. The polymer solution was transferred to a 50 mL Falcon tube. 1 mL of 0.1 M HCl in methanol was added, followed by the addition of 20 mL of DI water. The solution was mixed by shaking vigorously, which gave a cloudy mixture. The Falcon tube was centrifuged for 15 min at 11,000 rpm. After centrifuging, the mixture exhibited three layers. The top layer, mainly an aqueous solution, and the middle layer, a white cloudy mixture, were decanted. The above procedure was repeated for another two times with a mixture of 1 mL of 0.1 M HCl in methanol and 20 mL of DI water, followed by two times with 20 mL of DI water. The resulting clear CH<sub>2</sub>Cl<sub>2</sub> solution of the polymer was transferred to a clean 20 mL scintillation vial. The solvent was removed by a rotary evaporator. The resulting viscous polymer was further dried under vacuum overnight.

**Polymer Characterization.** <sup>1</sup>H NMR spectroscopy was performed on a 400 MHz Agilent NMR spectrometer at room temperature and referenced to the residual solvent signal of CDCl<sub>3</sub> (7.26 ppm) or CD<sub>2</sub>Cl<sub>2</sub> (5.33 ppm). Differential scanning calorimetry (DSC) was performed on a DSC250 (TA Instruments) with a RCS90 electric chiller attachment. 2.0–5.0 mg of each polymer was placed in an aluminum Tzero DSC pan with a Tzero hermetic lid. The DSC runs were conducted under a nitrogen atmosphere. The samples were first cooled from room temperature to −90 °C, then heated from −90 to 40 °C at 10 °C/min, held at 40 °C for 1 min, cooled to −90 °C at 10 °C/min, and held at −90 °C for one min. This protocol was repeated two more times but to 20 °C instead of 40 °C. The third

heating and cooling ramps were at 5 °C/min. The glass transition temperature was measured using the third heating scan. Size exclusion chromatography (SEC) was carried out on an Agilent system with a 1260 Infinity isocratic pump, degasser, and thermostatted column chamber held at 30 °C containing an Agilent PLgel 10 μm MIXED-B column with an operating range of 500–10,000,000 g/mol relative to PS standards. Chloroform was used as the mobile phase. This system was equipped with an Agilent 1260 refractometer and Infinity Bio-inert multi-detector suite featuring dual-angle static and dynamic light scattering detection. The optical microscope images were recorded under a Zeiss Axio Scope.A1 with a 20× N-Achroplan objective lens. The microscope was coupled with a Linkam LTS420 temperature-controlled stage. The samples were prepared in a nitrogen glovebox and covered with a glass coverslip.

**Broadband Dielectric Spectroscopy.** Dielectric constant measurements were performed with a broadband dielectric spectrometer (BDS) from Novocontrol. The instrument has a high-resolution dielectric Alpha analyzer (frequency range of 10<sup>-5</sup>–10<sup>7</sup> Hz) and can achieve a temperature range from −160 to 400 °C via nitrogen gas and activated heaters. The temperatures within the cell and the Alpha analyzer are controlled by WinDETA software. A parallel plate-oriented sample cell for viscous samples was used for the dielectric measurements. The viscous polymer samples were loaded on a gold-plated lower electrode with a diameter of 40 mm and covered with a smaller gold-plated electrode with a diameter of 20 mm. Teflon spacers of a known surface area and thickness were placed between the electrodes to control the thickness of the polymer sample. The measurement cell was kept closed by tightening the cell closing plate with a spring. Samples were measured from 10<sup>-2</sup> to 10<sup>7</sup> Hz and in a temperature range of 0 to 90 °C.

**General Procedure for the Preparation of Polymer/Salt Blends.** The polymers were brought into a dry nitrogen glove box and dried in vacuo below 10 mTorr at room temperature for one day. Approximately 0.2 g of polymer was added to a tared scintillation vial with a stir bar. The vial containing polymer was then dried again in vacuo overnight and reweighed. The dry weight was used to calculate how much LiTFSI would be needed. After the corresponding amount of LiTFSI was weighed and added to the vial, ca. 0.2 mL of anhydrous THF was added to the vial, which was then allowed to stir at room temperature for 24 h. The sample was then dried in vacuo at room temperature until it reached the weight before the solvent was added.

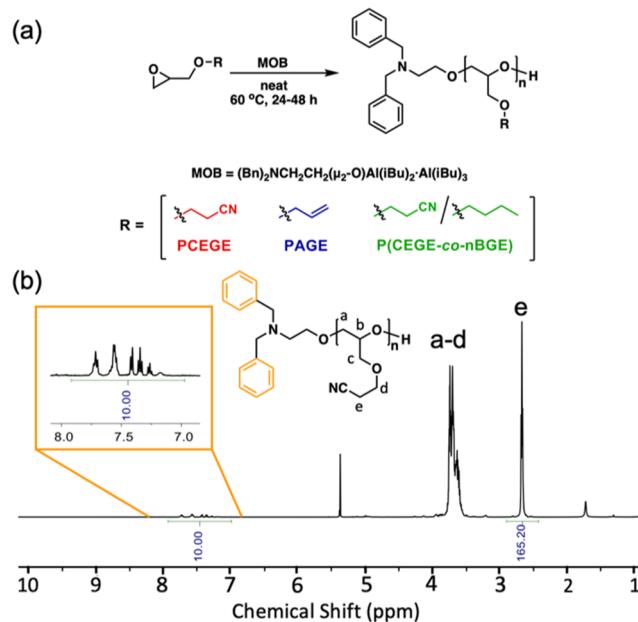
**Electrochemical Impedance Spectroscopy.** All data were collected on a BioLogic MTZ-35 Impedance Analyzer with an intermediate temperature system utilizing a through-plane conductivity cell. The electrode plates as well as the EIS chamber were brought into the glove box with a Teflon washer (surface area ca. 50 mm<sup>2</sup>) with the same diameter as the electrode. The Teflon washer served to contain the polymer electrolyte and fix the distance between the electrodes. The washer was placed on one electrode, the dry polymer electrolyte was added to fill the washer, and the other electrode was placed on top and screwed in place. The electrode assembly was placed into the sample holder and sealed before being brought out of the glove box. The EIS probes and temperature probe were then inserted into the sample holder. The sample chamber was placed into an insulated temperature stage. Each experiment was run in a frequency range that encompassed the full data set at all temperatures (30 to 90 °C with a step increase of 10 °C). Impedance data were collected during both a heating cycle and a cooling cycle. The two measurements were performed for each polymer electrolyte sample. The sample was allowed to equilibrate for an hour at each temperature before data collection. The washer inner diameter and thickness was used for the sample dimensions when calculating ionic conductivity. Ionic conductivities and standard deviations were reported from the average of four data sets.

**Small-Angle X-ray Scattering.** All data were collected on a SAXSLabs Ganesha 300K with an X-ray wavelength of 1.5418 Å. Kapton tape and a stainless-steel washer to contain the polymer samples were brought into a nitrogen-filled glovebox. The stainless-steel washer was placed on the Kapton tape, and a polymer sample was then added to fill the washer with another piece of Kapton tape

placed on top to cover the sample. The sample was sealed before being brought out of the glove box. Samples were measured in a temperature range of 30 to 80 °C with a step increase of 10 °C. Samples were allowed to equilibrate for 30 min at each temperature before data collection on heating.

## RESULTS AND DISCUSSION

A high-polarity polyether with a dielectric constant higher than 20 is required in order to explore polymer–polymer miscibility and ion transport in intermediate dielectric contrast polymer blend electrolytes.<sup>44</sup> Poly(cyanoethyl glycidyl ether) (PCEGE) possesses polar cyano functional groups pendant to the polyether backbone that increase the overall polymer polarity. The synthesis of PCEGE was previously reported by Cantor et al. using the Vandenberg catalyst to polymerize cyanoethyl glycidyl ether (CEGE);<sup>46</sup> however, the resulting PCEGE was insoluble in most organic solvents and was therefore not well characterized. We hypothesized that the limited solubility of PCEGE was a result of its high molecular weight, which is typical of polymers synthesized with the Vandenberg catalyst.<sup>47,48</sup> Therefore, we synthesized our PCEGE using a previously reported mono( $\mu$ -alkoxo)bis(alkylaluminum) (MOB) initiator/catalyst composition that has demonstrated control over epoxide ring-opening polymerizations up to intermediate molecular weights (ca. 10–60 kg/mol) (Figure 1a).<sup>45,49</sup> The polymerization was performed to target a



**Figure 1.** (a) Synthesis of PCEGE, PAGE, and P(CEGE-co-nBGE). (b) <sup>1</sup>H NMR spectroscopy (500 MHz) of PCEGE in CD<sub>2</sub>Cl<sub>2</sub>.

number-average molecular weight ( $M_n$ ) of 10 kg/mol. The polymer was obtained as a viscous yellow liquid and exhibited good solubility in organic solvents such as dichloromethane, THF, and chloroform, allowing for comprehensive solution-phase characterization of its chemical structure and physical properties.  $M_n$  of PCEGE was determined via end-group analysis based on <sup>1</sup>H NMR spectroscopy (Figure 1b, Table 1) and SEC analysis relative to PS standards (Table 1). End group analysis yielded a  $M_n$  of 10.2 kg/mol and SEC analysis calculated a  $M_n$  of 10.2 kg/mol, which were close to the target of 10 kg/mol. The molecular weight dispersity ( $D$ ) was

Table 1. Polymer Characterization Data

polymer	$M_n^a$ (kg/mol)	$M_n^b$ (kg/mol)	$D$
PAGE	11.5	10.3	1.27
P(CEGE- <i>co</i> -nBGE)	12.8	12.4	1.22
PCEGE	10.2	10.2	1.24

<sup>a</sup>Calculated based on end-group analysis of <sup>1</sup>H NMR spectroscopy.

<sup>b</sup>Calculated based on GPC analysis relative to PS standards.

determined to be 1.24 by SEC relative to PS standards (Table 1). We then copolymerized CEGE with a low-polarity monomer, *n*-butyl glycidyl ether (nBGE), in a 50:50 weight ratio to synthesize a material with intermediate polarity. The resulting copolymer P(CEGE-*co*-nBGE) had  $M_n$  of 12.8 kg/mol according to end-group analysis and  $D = 1.22$  by SEC. For the low-polarity, low-viscosity component, we chose poly(ally glycidyl ether) (PAGE), which exhibits a polarity similar to PEO, but without semi-crystallinity.<sup>9</sup> PAGE/LiTFSI electrolytes have been previously investigated and exhibit conductivities of *ca.*  $>10^{-5}$  S/cm at room temperature.<sup>9</sup>

The dielectric permittivities of PCEGE, P(CEGE-*co*-nBGE), and PAGE were determined by broadband dielectric spectroscopy (BDS) and taken as a measurement of their polarities. The average of the plateau region of the measured real permittivity as a function of frequency was taken as the dielectric constant ( $\epsilon$ ) (Figure 2a). The dielectric permittivity

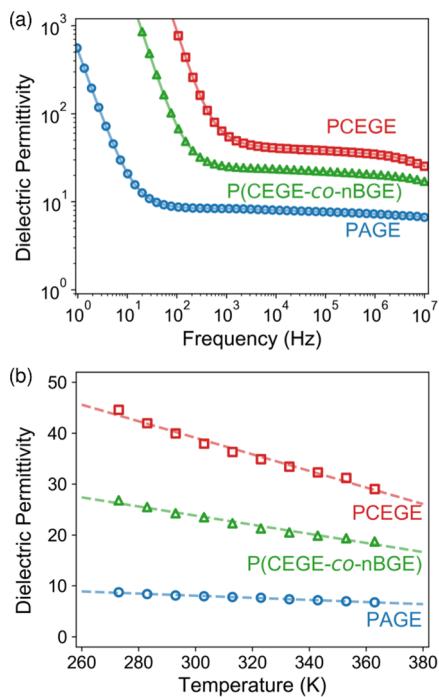


Figure 2. (a) BDS of PAGE, P(CEGE-*co*-nBGE), and PCEGE at 30 °C. (b) Dielectric permittivities of PAGE, P(CEGE-*co*-nBGE), and PCEGE in the temperature range 0 to 90 °C.

of PCEGE increased from 32 to 45 as the temperature decreased from 90 to 0 °C, which was three to four times higher than that of PAGE at all temperatures (Figure 2b), indicating the high polarity contrast between these two polymers. The dielectric permittivity of P(CEGE-*co*-nBGE) increased from 19 to 25 with decreasing temperature from 90 to 0 °C, still maintaining a high contrast with PAGE.

The glass transition temperature ( $T_g$ ) is another important property that significantly affects the ionic conductivities of polymer electrolytes.<sup>19,28,29,50–55</sup> The strong dipole–dipole attractions between the cyano functional groups in PCEGE and P(CEGE-*co*-nBGE) results in strong polymer–polymer interactions, similar to conventional polymers like poly(acrylonitrile-*co*-butadiene).<sup>56</sup> Consequently, the  $T_g$ s of  $-37$  and  $-56$  °C, respectively, for PCEGE and P(CEGE-*co*-nBGE) were much higher than that of PAGE ( $-78$  °C). Upon dissolving LiTFSI into the polymers, the  $T_g$  of the polymer electrolytes increased relative to the pure-polymer  $T_g$ s.

Having characterized the polarity and thermal properties of the polymers, we then investigated ion transport in electrolytes composed of binary polymer/LiTFSI blends at a concentration of  $r = 0.125$  ( $r$  denotes the molar ratio between LiTFSI and polymer repeat units), which had previously been observed to be the optimal salt loading for PAGE homopolymers, beyond which LiTFSI solubility becomes an issue.<sup>9</sup> The ionic conductivities ( $\sigma$ ) of the single-polymer electrolytes were measured with electrochemical impedance spectroscopy (EIS) under a nitrogen atmosphere. The conductivities are plotted as  $\sigma(T)$  versus  $1000/T$  (K<sup>-1</sup>) in Figure 3a, which shows the ionic

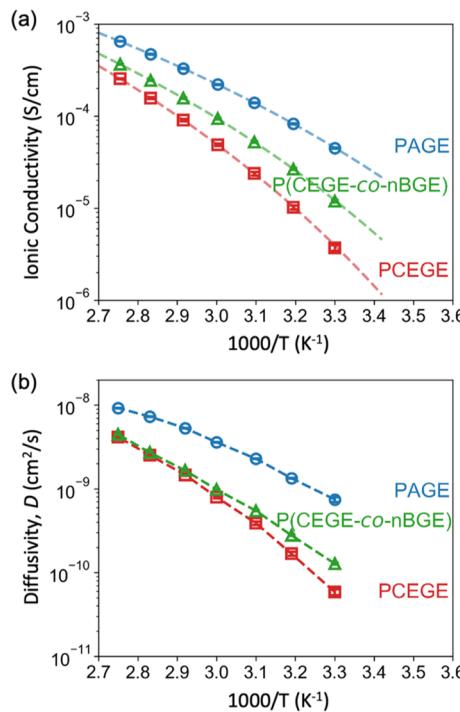


Figure 3. (a) Ionic conductivities and (b) diffusivities of PAGE-LiTFSI, P(CEGE-*co*-nBGE)-LiTFSI, and PCEGE-LiTFSI electrolytes ( $r = 0.125$ ) in a temperature range from 30 to 90 °C. The dashed lines in (a) represent VTF fits to the ionic conductivity data.

conductivities measured for PCEGE/LiTFSI, P(CEGE-*co*-nBGE)/LiTFSI, and PAGE/LiTFSI blends as a function of temperature (30–90 °C). The conductivities of the PCEGE/LiTFSI electrolytes were more than 1 order of magnitude lower than those of PAGE/LiTFSI at low temperatures ( $\leq 40$  °C). As the temperature increased, the conductivity of the PCEGE/LiTFSI blend increased from  $10^{-6}$  to  $10^{-4}$  S/cm. Such a strong temperature dependence suggested that the ionic conductivity of PCEGE was limited by segmental dynamics, which is in good agreement with our previous simulations and

experimental results in the high-dielectric permittivity regime.<sup>29</sup> The conductivity data were fit to the Vogel–Tamman–Fulcher (VTF) equation (eq 1), which is commonly used to model temperature-dependent conductivity in polymer electrolytes.

$$\sigma(T) = \frac{A}{T^{1/2}} \exp\left[-\frac{E_a}{R(T - T_0)}\right] \quad (1)$$

The VTF fit of experimental ionic conductivity data allows for analyzing  $\sigma(T)$  as a function of temperature by qualitatively separating the effects of charge carrier concentration, often related to the prefactor  $A$  and the activation energy,  $E_a$ . In our analysis, the Vogel temperature ( $T_0$ ) was included as a fitting parameter instead of being set to the typical  $T_g$ –50 value because of the possible misleading results caused by improper choice of  $T_0$ .<sup>57</sup> Our fitting gave  $T_0$  close to the value of  $T_g$ –50 for the single-polymer electrolytes (see Table 2), suggesting that, in the range of temperatures investigated, the mechanism of ion conduction follows this VTF form rather than a direct Arrhenius dependence.<sup>57</sup>

**Table 2.** VTF Fitting Parameters with Variance for Single-Polymer Electrolytes with LiTFSI ( $r = 0.125$ )

	$A$ (K <sup>1/2</sup> ·S/cm)	$E_a$ (kJ/mol)	$T_0$ (K)
PAGE	5.95 ± 0.84	10.0 ± 0.4	168 ± 4
P(CEGE- <i>co</i> -nBGE)	9.28 ± 1.00	10.9 ± 0.3	179 ± 2
PCEGE	21.3 ± 3.9	12.4 ± 0.5	184 ± 3

According to the VTF analysis, the pre-factor  $A$  increased from 5.95 K<sup>1/2</sup>·S/cm for PAGE/LiTFSI to 9.28 K<sup>1/2</sup>·S/cm for P(CEGE-*co*-nBGE)/LiTFSI and further to 21.3 K<sup>1/2</sup>·S/cm for PCEGE/LiTFSI. This increase in the pre-factor was consistent with the increasing polarity of the host polymer, which is expected to increase the free charge carrier concentration by increasing the solubility of LiTFSI in the polymer matrix.<sup>25,26</sup>

Despite the increase in charge carrier solubility in the more polar polymers, the ionic conductivities decreased with the polarity of the single-polymer electrolytes in agreement with coarse-grained molecular dynamics simulations by Wheatle et al. in the high-dielectric regime<sup>24,25</sup> and the corroborating experimental results.<sup>29</sup> Our analysis revealed that the  $E_a$  of single-polymer electrolytes increased with the polymer polarity. The higher  $E_a$  in PCEGE/LiTFSI and P(CEGE-*co*-nBGE)/LiTFSI electrolytes was attributed to stronger polymer–polymer and ion–polymer interactions, ultimately resulting in lower ion diffusivities ( $D$ ). In this higher dielectric constant regime, the ionic conductivity was governed by  $T_g$ , which gives insights into relative chain mobility and not ion solubility. The observed trend in ionic conductivity was therefore expected as the  $T_g$  of the polymers increased from PAGE to P(CEGE-*co*-nBGE) to PCEGE. We calculated the diffusivities of the three single-polymer electrolytes using a method reported by Sokolov et al.<sup>58,59</sup> The real conductivities ( $\sigma'$ ) of PCEGE and PAGE were first plotted against frequency ( $\nu$ ) and then fit to the power law presented in eq 2 (Figure S33). The fitting yielded a critical frequency ( $\nu^*$ ) that characterizes the boundary between the DC plateau and the hopping regime and a power-law exponent ( $a$ ). The diffusivity was calculated using eq 3, where  $n_p$  is the number density of charge pairs in the electrolyte.

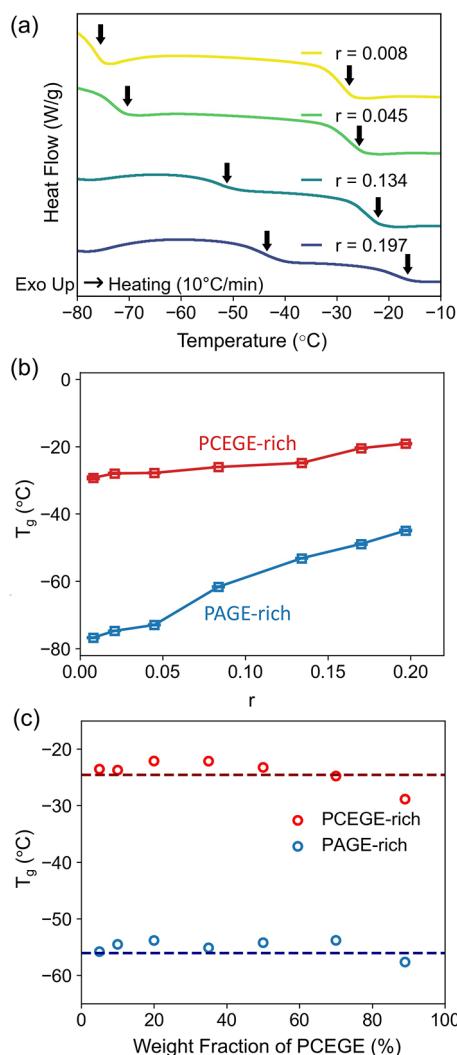
$$\sigma(\nu) = \sigma_0 \left[ 1 + \left( \frac{\nu}{\nu^*} \right)^a \right] \quad (2)$$

$$D = \frac{\pi \nu^*}{3} \left( \frac{3}{4\pi n_p} \right)^{2/3} \quad (3)$$

The ion diffusivities of PAGE/LiTFSI, P(CEGE-*co*-nBGE)/LiTFSI, and PCEGE/LiTFSI electrolytes increased as the temperature increased from 30 to 90 °C (Figure 3b). The diffusivity in high-polarity PCEGE/LiTFSI and P(CEGE-*co*-nBGE)/LiTFSI was significantly lower than that of PAGE/LiTFSI, particularly when the temperature was below 50 °C. Together, our results suggested that slow segmental dynamics, as characterized by a high energy barrier and limited ion diffusivity, decreased the ionic conductivity of high-polarity polymer electrolytes.

The inherent trade-off between high ion dissociation and rapid segmental dynamics makes it difficult to improve the ionic conductivity of polymer electrolytes. In other contexts, blending high-polarity and low-viscosity polymers has been studied as a strategy to overcome the transport limitations that arise in single-component polymer hosts.<sup>41,60–63</sup> Motivated by such studies, we examined the miscibility of 50:50 wt % PAGE/PCEGE blends as a function of LiTFSI concentration. Over the entire range of LiTFSI concentration investigated ( $0 < r < 0.2$ ), two distinct  $T_g$ s were observed as shown in Figure 4a, suggesting the coexistence of two phases. Both  $T_g$ s increased as LiTFSI concentration increased, consistent with increasing ion concentration in both phases. The difference in  $T_g$  between phases decreased as LiTFSI concentration increased as shown in Figure 4b, but the system did not converge to a single  $T_g$  within the LiTFSI concentrations investigated. We examined the changes in  $T_g$  of PAGE/PCEGE/LiTFSI blends at a single concentration of LiTFSI ( $r = 0.125$ ) while varying the composition of PAGE/PCEGE (Figure 4c). DSC traces of the entire series of the high-dielectric-contrast PAGE/PCEGE blends revealed two  $T_g$ s, one at *ca.* –54 °C (PAGE-rich phase) and the other at *ca.* –25 °C (PCEGE-rich phase). The  $T_g$ s of each phase in the blends were close to those of the single-polymer electrolytes. The composition of each phase was likely similar to the single-polymer electrolytes, indicating near-total immiscibility within the investigated composition range of the high dielectric-contrast PAGE/PCEGE system. The immiscibility was corroborated by optical microscopy of the blends revealing a multiphase morphology at room temperature (Figure S14).

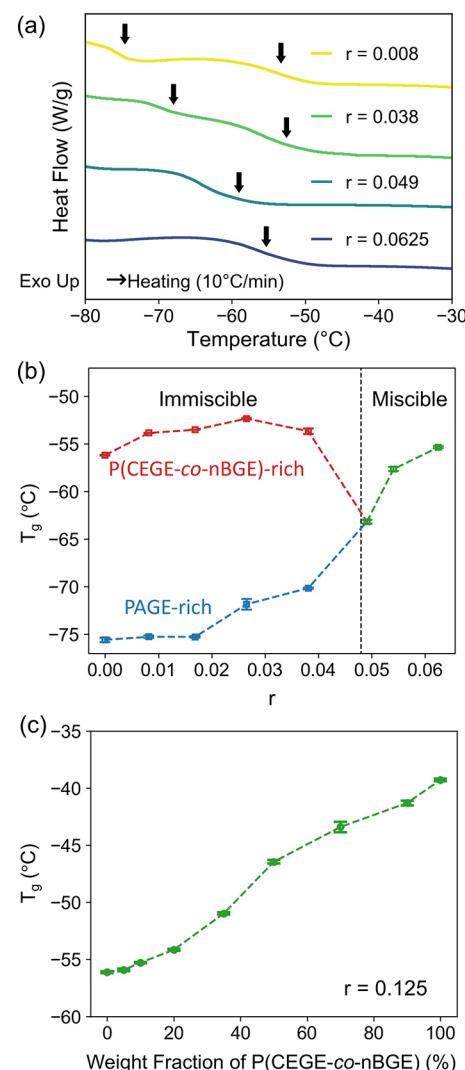
The miscibility of lower dielectric-contrast PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends was investigated in a similar manner. Varying amounts of LiTFSI were added to a symmetric polymer blend consisting of 50:50 wt % PAGE/P(CEGE-*co*-nBGE). At low LiTFSI concentration ( $r < 0.04$ ), two glass transition temperatures were observed by DSC as shown in Figure 5a. However, when the concentration of LiTFSI was increased above  $r = 0.049$ , the blends became miscible as evidenced by a single  $T_g$  observed on the DSC. In the PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends, the higher concentration of LiTFSI resulted in improved miscibility between the two polymers, eventually leading to a compatibilized polymer blend as shown in Figure 5b. At a LiTFSI concentration of  $r = 0.125$ , the PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends showed a single  $T_g$  at all compositions (Figure S13) with the  $T_g$  increasing with the weight fraction of P(CEGE-*co*-nBGE) (Figure 5c). Our



**Figure 4.** (a) DSC traces and (b)  $T_g$  of immiscible 50:50 wt % PAGE/PCEGE blends at different LiTFSI concentrations. (c)  $T_g$  of immiscible PAGE/PCEGE with a LiTFSI concentration of  $r = 0.125$  at different weight fractions of PCEGE. Red/blue dashed lines represent the  $T_g$  of single-polymer electrolytes at the same LiTFSI concentration.

observation of the miscibility change of PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends provides experimental support for a previous theoretical study that reported a decreased effective interaction parameter ( $\chi_{\text{eff}}$ ) in high-polarity polymer blends with added salt.<sup>27,44</sup>

The different behaviors of PAGE/P(CEGE-*co*-nBGE)/LiTFSI (miscible) and PAGE/PCEGE/LiTFSI (immiscible) create an opportunity to investigate how miscibility affects the ionic conductivity of polymer blends with high polarity contrast. As the composition of the high-polarity polymer in the blend increased, the ionic conductivity of the blends decreased. Wide variability was observed in the measurement of PAGE/PCEGE/LiTFSI (Figure 5a) due to its heterogeneity. To better evaluate the conductive properties of these polymer blends, the ionic conductivity of a polymer blend at a given composition and temperature ( $\sigma_{\text{blend},i}$ ) was compared to a conductivity model ( $\sigma_{\text{linear},i}$ ) calculated by assuming a linear mass-averaging of ionic conductivities of the two single-polymer electrolytes. We define the deviation ( $\delta$ ) as

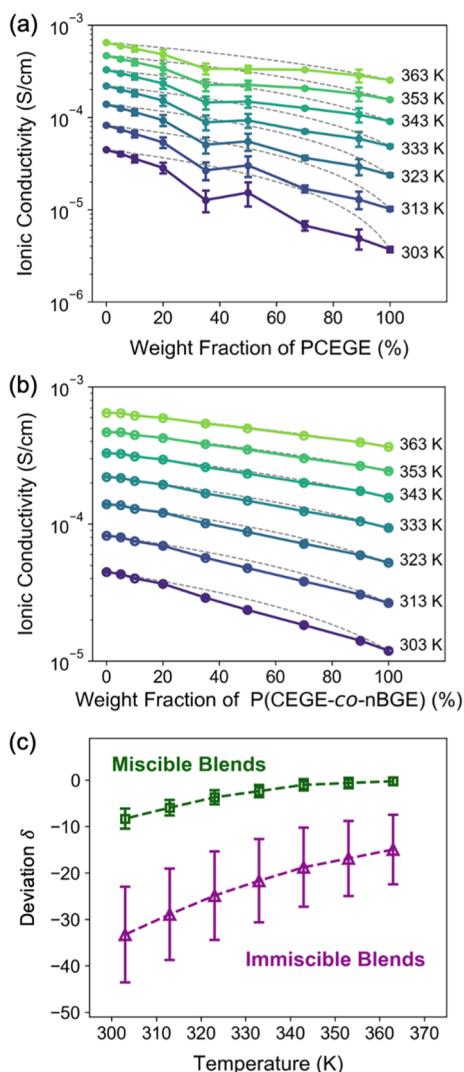


**Figure 5.** (a) DSC traces and (b)  $T_g$  of 50 wt % PAGE/P(CEGE-*co*-nBGE) at different LiTFSI loadings. (c)  $T_g$  of PAGE-P(CEGE-*co*-nBGE) with a LiTFSI loading of  $r = 0.125$  at different weight fractions of P(CEGE-*co*-nBGE).

$$\delta = \frac{1}{n} \sum_{i=1}^n \left( \frac{\sigma_{\text{blend},i} - \sigma_{\text{linear},i}}{\sigma_{\text{linear},i}} \right) \times 100 \quad (4)$$

where  $n$  is the number of conductivity data points collected across all blend compositions for each temperature and/or salt concentration.

In previous simulations, we predicted that the deviation in ionic conductivity of high-polarity-contrast polymer blends was heavily dependent on the compatibility of the two polymers.<sup>27</sup> A negative  $\delta$  was predicted when the two polymers did not interact with one another favorably, while positive  $\delta$  was predicted when the polymer–polymer interactions were more favorable. In addition, the coarse-grained molecular dynamics simulations suggested that in compatible ternary high-polarity-contrast polymer blend electrolytes, the low-polarity, low-viscosity component increases the segmental dynamics of the high-polarity component where the ions are preferentially solvated. As such, the less negative  $\delta$  of PAGE/P(CEGE-*co*-nBGE)/LiTFSI likely results from better compatibility and more favorable polymer–polymer interactions than in immiscible PAGE/PCEGE/LiTFSI (Figure 6). Above 353 K,



**Figure 6.** Ionic conductivities of (a) ternary PAGE/PCEGE/LiTFSI ( $r = 0.125$ ) blends and (b) PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends over a temperature range from 303 to 363 K. Dashed lines (in gray) represent the linear mass-average between the ionic conductivities of the two single-polymer electrolytes. (c) Deviation of ionic conductivities of miscible PAGE/P(CEGE-*co*-nBGE)/LiTFSI blends and immiscible PAGE/PCEGE/LiTFSI blends as a function of temperature.

the deviation is close to 0% for PAGE/P(CEGE-*co*-nBGE)/LiTFSI and the blend is miscible. This behavior contrasts with that of the PAGE/PCEGE/LiTFSI blends, which exhibit a negative  $\delta$  of  $-15$  to  $-33\%$  over the entire temperature range 303–363 K due to immiscibility. However, both miscible and

immiscible blends have a  $\delta$  that increases with temperature, meaning that heating enhances polymer–polymer compatibility.

To investigate the relationship between polymer–polymer interactions and ion solvation in these high-polarity-contrast polymer electrolytes, we investigated the disordered state small-angle X-ray scattering (SAXS) of a model block polymer of PAGE-*b*-PCEGE, which was prepared using the same  $[(Bn)_2NCH_2CH_2(\mu_2\text{-O})Al(iBu)_2\cdot Al(iBu)_3]$  initiator/catalyst and sequential ring-opening polymerization of AGE followed by CEGE (Scheme 1). The resultant PAGE-*b*-PCEGE was characterized with  $^1\text{H}$  NMR spectroscopy and SEC.  $^1\text{H}$  DOSY NMR spectroscopy (Figure S6) demonstrated that the chain end and both blocks had the same diffusion coefficient of *ca.*  $3.0 \times 10^{-6} \text{ cm}^2/\text{s}$ , consistent with full covalent attachment of PAGE and PCEGE blocks. Based on  $^1\text{H}$  NMR, the final  $M_n$  for each block was 7.1 kDa PAGE and 13.1 kDa PCEGE.

Disordered-state SAXS profiles of PAGE-*b*-PCEGE at various LiTFSI concentrations were collected and are shown in Figure 7a at 303 K. The scattering profiles revealed a broad single peak for PAGE-*b*-PCEGE at  $q = 0.3\text{--}0.6 \text{ nm}^{-1}$  at LiTFSI concentrations of  $r = 0.000\text{--}0.065$ , consistent with the disordered state of the block copolymer. The intensity of the scattering peak was sensitive to the LiTFSI concentration. The intensity first increased when a small amount ( $r = 0.011$ ) of LiTFSI was added and then decreased with further addition of LiTFSI until disappearing at  $r = 0.113$ . In addition, the primary scattering peak location ( $q^*$ ) for the pure diblock copolymer ( $r = 0.000$ ) was at  $q^* = 0.451 \text{ nm}^{-1}$  and the addition of LiTFSI ( $r = 0.011$ ) resulted in a small shift to  $q^* = 0.458 \text{ nm}^{-1}$  as shown in Figure 7a.

We analyzed the SAXS data within the random phase approximation (RPA) to extract the effective interaction parameter ( $\chi_{\text{eff}}$ ) and investigate ion-partitioning. Within this framework, the total scattering density function  $I_{\text{tot}}(q)$  can be written as

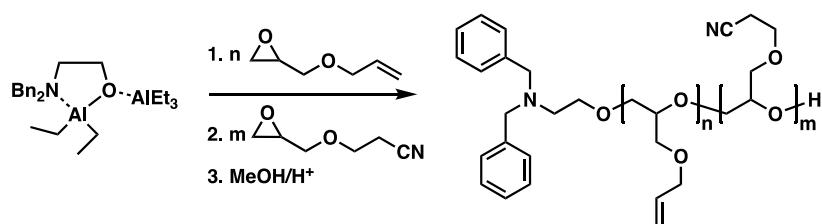
$$I_{\text{tot}}(q) = I_{\text{dis}}(q) + I_{\text{bkg}}(q) \quad (5)$$

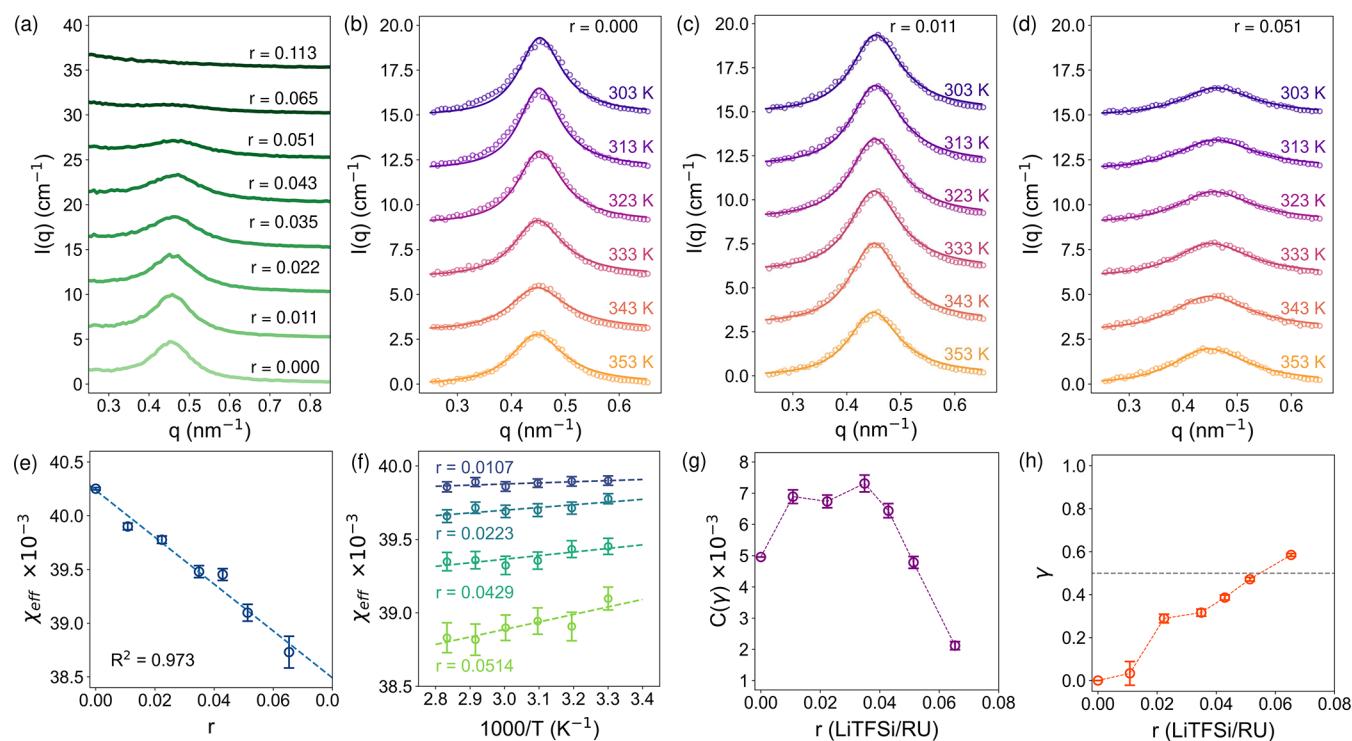
where  $I_{\text{bkg}}(q)$  is the background scattering signal, which is an exponentially decaying function of  $q$  with two adjustable parameters and  $I_{\text{dis}}(q)$  is the scattering profile of a disordered block polymer melt as given by<sup>64</sup>

$$I_{\text{dis}}(q) = C \left[ \frac{S(q)}{W(q)} - 2\chi_{\text{eff}} \right] \quad (6)$$

In this equation,  $\chi_{\text{eff}}$  is the effective Flory–Huggins interaction parameter, which will be different for the diblock system than the previous system containing a blend of the two homopolymers;<sup>65</sup>  $C$  is the X-ray scattering contrast as determined by the electron densities of the polymer phases:

**Scheme 1.** Synthesis of PAGE-*b*-PCEGE Diblock Copolymer Conducted Neat at 60 °C; PAGE Block Was Given 24 h to React Followed by 3 Days for the PCEGE Block





**Figure 7.** (a) SAXS profiles of PAGE-*b*-PCEGE with LiTFSI at various  $r$  at 303 K. Scattering profiles of PAGE-*b*-PCEGE electrolytes at different temperatures at LiTFSI concentrations of (b)  $r = 0.000$ , (c) 0.011, and (d) 0.051. The absolute scattering intensity (background subtracted) for  $T = 353$  K is presented. Data from 343, 333, 323, 313, and 303 K are shifted vertically for clarity by 3, 6, 9, 12, and 15  $\text{cm}^{-1}$ , respectively. The curves represent the best fit based on the adjustable contrast RPA model. (e)  $\chi_{\text{eff}}$  as a function of  $r$  at 303 K. The dashed lines are linear fits of  $\chi_{\text{eff}}$  and  $r$ . (f)  $\chi_{\text{eff}}$  as a function of temperature at different LiTFSI concentrations. The solid lines are fits according to eq 16. (g)  $C(\gamma)$  and (h)  $\gamma$  obtained from RPA fits with respect to increased loading of LiTFSI at 303 K.

and  $S(q)$  and  $W(q)$  are the determinant and the sum of the structure factor matrix as given by

$$S(q) = S_{\text{AA}}(q) + S_{\text{BB}}(q) + 2S_{\text{AB}}(q) \quad (7)$$

$$W(q) = S_{\text{AA}}(q)S_{\text{BB}}(q) - S_{\text{AB}}^2(q) \quad (8)$$

$$S_{\text{AA}}(q) = Ng(f) \quad (9)$$

$$S_{\text{BB}}(q) = Ng(1 - f) \quad (10)$$

$$S_{\text{AB}}(q) = (N/2)[g(q) - g(f) - g(1 - f)] \quad (11)$$

$$g(f) = (2/x^2)[fx + \exp(-fx) - 1] \quad (12)$$

$$x = q^2 R_g^2 \quad (13)$$

Here,  $g(f)$  is the form factor for a Gaussian chain (i.e., Debye function), and  $f$  is the volume fraction of PAGE as calculated using the density of PAGE and PCEGE and the relative composition of PAGE and PCEGE in the diblock copolymer according to  $^1\text{H}$  NMR spectroscopy.  $N$  is a degree of polymerization relative to a reference volume, and  $R_g$  is the radius of gyration of the diblock copolymer.

Based on eqs 5–13, the intensity and shape of the scattering peak is determined by  $\chi_{\text{eff}}$  and  $C$ , which are affected by the addition of LiTFSI to PAGE-*b*-PCEGE because of changes in microphase separation and the X-ray scattering contrast. Taking this into consideration, we adopted the adjustable contrast model developed by Balsara and co-workers to capture changes in both  $\chi_{\text{eff}}$  and  $C$ .<sup>36,66,67</sup> In this model,  $C$  is expressed as a function of  $\gamma$ , a parameter that describes the

preference of LiTFSI to be dissolved in either PAGE-rich ( $\gamma = 1$ ) or PCEGE-rich domain ( $\gamma = 0$ ).

$$C(\gamma) = v_{\text{ref}} \left( \frac{b_{\text{AGE}} + p(\gamma) \left( \frac{n_{\text{LiTFSI}}}{n_{\text{AGE}}} \right) b_{\text{LiTFSI}}}{v_{\text{AGE}} + p(\gamma) \left( \frac{n_{\text{LiTFSI}}}{n_{\text{AGE}}} \right) v_{\text{LiTFSI}}} - \frac{b_{\text{CEGE}} + (1 - p(\gamma)) \left( \frac{n_{\text{LiTFSI}}}{n_{\text{CEGE}}} \right) b_{\text{LiTFSI}}}{v_{\text{CEGE}} + (1 - p(\gamma)) \left( \frac{n_{\text{LiTFSI}}}{n_{\text{CEGE}}} \right) v_{\text{LiTFSI}}} \right)^2 \quad (14)$$

$$p(\gamma) = \frac{n_{\text{AGE}}(\gamma)}{n_{\text{AGE}}(\gamma) + n_{\text{CEGE}}(1 - \gamma)} \quad (15)$$

In this model,  $b_{\text{AGE}}$  and  $b_{\text{CEGE}}$  are the respective X-ray scattering length densities of the polymer blocks,  $v_{\text{AGE}}$  and  $v_{\text{CEGE}}$  are the molar volumes of the repeat units, and  $v_{\text{ref}}$  is the reference volume and is set to  $0.1 \text{ nm}^3$ . The important feature of this method is that the disordered scattering peak can be fitted with three adjustable parameters  $\chi_{\text{eff}}$ ,  $R_g$ , and  $\gamma$  to gain insights into the thermodynamics of ion partitioning in polymer electrolytes.

The first key observation was that the primary peak position remained relatively constant while the peak intensity decreased upon heating (Figure 7b). Similar trends were observed at higher LiTFSI concentrations (Figure 7c,d). The fit parameters varied with temperature and LiTFSI concentration. At 303 K,  $\chi_{\text{eff}}$  monotonically decreases with increasing  $r$ , suggesting that LiTFSI facilitates more favorable polymer–polymer interactions between PAGE and PCEGE. The temperature depend-

ence of  $\chi_{\text{eff}}$  is also shown in Figure 7f at LiTFSI concentrations of  $r = 0.0107, 0.0223, 0.0429$ , and  $0.0514$ . The solid lines are the fits for each concentration using the empirical equation.

$$\chi_{\text{eff}} = \frac{A}{T} + B \quad (16)$$

For all LiTFSI concentrations,  $\chi_{\text{eff}}$  decreased with increasing temperature, indicating better polymer miscibility at higher temperatures. This observation corroborated our earlier temperature-dependent ionic conductivity deviations from linearity ( $\delta$ ): For PAGE/PCEGE/LiTFSI and PAGE/P-(CEGE-*co*-nBGE)/LiTFSI blends, the deviation ( $\delta$ ) approached zero as temperature increased. The decrease of  $\chi_{\text{eff}}$  did not explain the increased scattering intensity as  $r$  was increased from 0.0 to 0.011. The plot of the X-ray scattering contrast,  $C(\gamma)$ , versus  $r$  (Figure 7g) shows an increase in  $C(\gamma)$  from  $5.1 \times 10^{-3} \text{ cm}^{-1}$  for the pure PAGE-*b*-PCEGE ( $r = 0.0$ ) to  $7.4 \times 10^{-3} \text{ cm}^{-1}$  at a LiTFSI concentration of  $r = 0.035$ . Further increases in LiTFSI concentration led to decreases in  $C(\gamma)$ . Therefore, the initial SAXS intensity increase at a low LiTFSI concentration may be attributed to the increase in contrast,  $C(\gamma)$ , after the addition of a small amount of LiTFSI that partitioned primarily to PCEGE-rich domains. Significantly, the adjustable contrast model was able to distinguish between simultaneous changes in  $\chi_{\text{eff}}$  and  $C(\gamma)$ . The eventual disappearance of the primary scattering peak at  $r = 0.113$  is attributed to both an increase in polymer compatibility as well as the decrease in scattering contrast between the two blocks.

For SMEs, a molecular-level picture of the ion solvation consists of the high-polarity component being present in the immediate solvation shell surrounding the ions.<sup>6</sup> Our previous molecular dynamics study on high-polarity-contrast polymer blend electrolytes with  $r = 0.065$  also predicted that the high-polarity component is predominantly in the vicinity of solvated ions.<sup>27</sup> Here, when  $r$  was less than 0.05,  $\gamma$  was below 0.5, revealing a stronger preference for LiTFSI to be partitioned in the high-polarity PCEGE-rich domain (Figure 7h). At  $r = 0.011$ , almost all LiTFSI was solvated in the PCEGE-rich domain. As the LiTFSI concentration increased further, little selective partitioning was evident, leading to a much smaller contrast,  $C(\gamma)$  (Figure 7g), and a more uniform distribution of ions in the system while  $\chi_{\text{eff}}$  simultaneously decreased. The disappearance of apparent scattering peaks at high LiTFSI concentrations (c.f., Figure 7a) was attributed to the decrease in both  $\chi_{\text{eff}}$  and  $C(\gamma)$ . To put these results in context, as with the SME blend case, the LiTFSI preferentially partitioned in the more polar PCEGE-rich phase at low concentration. Because the LiTFSI was preferentially solvated in one phase, the scattering contrast and therefore the SAXS signal intensity increased. As the amount of LiTFSI was increased, the polymer blocks became more miscible. The  $\text{Li}^+$  ions interact with both polyethers and dilute the unfavorable segmental interactions between incompatible polymers, which explains why  $\chi_{\text{eff}}$  of the system decreases. The LiTFSI compatibilized the two blocks through shared binding of  $\text{Li}^+$  ions, decreasing the scattering contrast of the system as the LiTFSI was more uniformly distributed throughout the polymer matrix.

## CONCLUSIONS

We investigated the relationship between polarity-contrast, phase behavior, and ionic conductivity in polyether blend electrolytes. PAGE/P(CEGE-*co*-nBGE)/LiTFSI was miscible at high LiTFSI concentrations, while the immiscibility of

PAGE/PCEGE/LiTFSI blends gave rise to a strong negative deviation from the linear average of the ionic conductivity of PAGE/LiTFSI and PCEGE/LiTFSI. This negative deviation became less significant at higher temperature due to more favorable polymer–polymer interactions and improved miscibility. For the miscible PAGE/P(CEGE-*co*-nBGE)/LiTFSI blend electrolyte, the negative deviation in ionic conductivity was less than what was found in the immiscible PAGE/PCEGE/LiTFSI, becoming near zero when the temperature was above 353 K. This observation stresses the importance of polymer–polymer host compatibility in designing high-polarity-contrast polymer blend electrolytes. Effective polymer–polymer interactions mediated by the presence of LiTFSI were investigated by SAXS of disordered PAGE-*b*-PCEGE. We found that the effective interaction parameter decreased monotonically as the salt loading increased using the adjustable contrast model. At low concentrations, LiTFSI was primarily solvated in the high-polarity PCEGE domain as demonstrated by the low values for the partition function,  $\gamma$ . Future investigation into polymer blend electrolytes should focus on optimizing polarity contrast and miscibility to balance high ionic dissociation with rapid segmental dynamics and ultimately exploit a positive deviation in ionic conductivity as predicted by Wheate et al.<sup>27</sup>

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.macromol.2c02023>.

<sup>1</sup>H NMR spectroscopy, size exclusion chromatography, differential scanning calorimetry, broadband dielectric spectroscopy, and electrochemical impedance spectroscopy (PDF)

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**Notes**

The authors declare no competing financial interest.

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