



# An Independent Determination of the Distance to Supernova SN 1987A by Means of the Light Echo AT 2019xis

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## Abstract

Accurate distance determination to astrophysical objects is essential for the understanding of their intrinsic brightness and size. The distance to SN 1987A has been previously measured by the expanding photosphere method and by using the angular size of the circumstellar rings with absolute sizes derived from light curves of narrow UV emission lines, with reported distances ranging from 46.77 to 55 kpc. In this study, we independently determined the distance to SN 1987A using photometry and imaging polarimetry observations of AT 2019xis, a light echo of SN 1987A, by adopting a radiative transfer model of the light echo developed in Ding et al. We obtained distances to SN 1987A in the range from  $49.09 \pm 2.16$  kpc to  $59.39 \pm 3.27$  kpc, depending on the interstellar polarization and extinction corrections, which are consistent with the literature values. This study demonstrates the potential of using light echoes as a tool for distance determination to astrophysical objects in the Milky Way, up to kiloparsec level scales.

*Unified Astronomy Thesaurus concepts:* Polarimetry (1278); Supernovae (1668); Large Magellanic Cloud (903); Distance measure (395)

## 1. Introduction

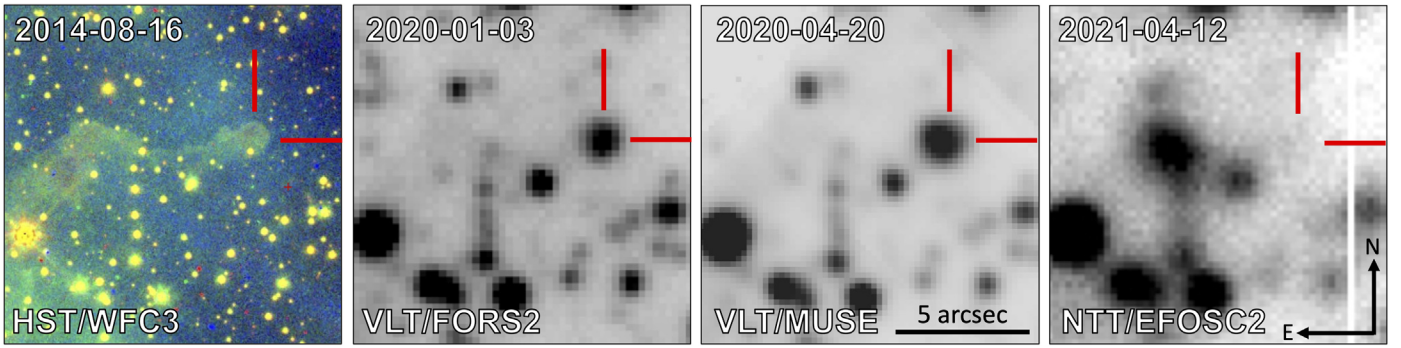
Distance determination to astrophysical objects is fundamental for improving our knowledge about these objects. Without knowing the distance to objects, their intrinsic brightness and size cannot be determined. Furthermore, accurate distances to extragalactic objects are crucial for the determination of the Hubble constant,  $H_0$ , and measuring the expansion rate of the universe (Riess et al. 1998; Perlmutter et al. 1999). The distance to the Large Magellanic Cloud (LMC) is important because the Cepheids in the LMC are used as an anchor for the cosmic distance ladder (e.g., Riess et al. 2019). Probably the most accurate distance to the LMC within 1% was reported by Pietrzyński et al. (2019), who determined the distance,  $d = 49.59 \pm 0.09$  (statistical)  $\pm 0.54$  (systematic) kpc. They applied the interferometry-based surface brightness–color relation derived from single late-type stars to fully detached eclipsing binaries with similar spectral types.

SN 1987A, which exploded in the LMC, has also a relatively well-determined distance. Branch (1987) determined the distance to SN 1987A of  $55 \pm 5$  kpc by matching a distance-independent photometric angular radius of the photosphere to a distance-dependent spectroscopic angular radius, which is known as the Baade method (Baade 1926), or the expanding photosphere method (Kirshner & Kwan 1974). Using the same method, Höflich (1988) obtained a distance of  $48 \pm 2$  kpc based on full non-LTE models by fitting a spectral sequence of

eight spectra obtained during the first 7 months. Eastman & Kirshner (1989) developed a model atmosphere for SN 1987A and determined the distance to SN 1987A of  $49 \pm 6$  kpc using the expanding photosphere method. Hanuschik & Schmidt-Kaler (1991) also applied the Baade method for distance determination and derived a distance of  $52.3 \pm 1.5$  kpc. Panagia et al. (1991) have compared the angular size of the circumstellar rings of SN 1987A in the Hubble Space Telescope (HST) images with absolute sizes determined from light curves of narrow UV emission lines and derived a distance of  $51.2 \pm 3.1$  kpc to SN 1987A. They also determined a distance of  $50.1 \pm 3.1$  kpc to the center of the LMC using radial velocities of the SN relative to the center of the LMC. Schmidt et al. (1992) used IR photometry of SN 1987A and determined a distance of  $49 \pm 3$  kpc using the expanding photosphere method. Gould (1994) revisited the supernovae-ringing method and the measurements by Panagia et al. (1991) and determined a distance to the SN of  $52.7 \pm 2.6$  kpc and to the LMC of  $53.2 \pm 2.6$  kpc. Based on new measurements of the ionized-emission light curves and the angular size of the ring around SN 1987A, Gould (1995) derived an upper limit for the distance to the SN of  $D < 46.77 \pm 0.76$  kpc. Crotts et al. (1995) investigated the light echoes from the inner region of SN 1987A to define the 3D geometry of the circumstellar nebula and used the Panagia et al. (1991) method to determine the distances to the SN and LMC of  $d_{\text{SN}} = 51.7 \pm 3.1$  kpc and  $d_{\text{LMC}} = 51.9 \pm 3.1$  kpc, respectively. Mitchell et al. (2002) conducted a detailed spectroscopic analysis of SN 1987A and modeled the observed spectra from day 1 to 81 using hydrodynamical models. Furthermore, they used the spectral-



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**Figure 1.** Light echo AT 2019xis observed on 2020 January 3 with FORS2 in the V band (this work), 2020 April 20 with the Multi Unit Spectroscopic Explorer (MUSE; Prog. ID 2104.D-5041(A), PI Kuncarayakti), and 2021 April 12 with the European Southern Observatory (ESO) Faint Object Spectrograph and Camera (EFOSC2) in the V band (as part of PESSTO; Smartt et al. 2015), compared to a Hubble Space Telescope (HST) archival image (Prog. ID. 13401, PI Fransson) observed in the F814W, F606W, and F502N bands with the Wide Field Camera 3 (WFC3). The red lines mark the position of AT 2019xis. The light echo propagated eastward from SN 1987, which is located  $4'.05$  west from AT 2019xis.

fitting expanding atmosphere method to derive the distance to SN 1987A of  $50 \pm 5$  kpc.

Distances to variable sources can also be determined by means of light echoes (e.g., Trümper & Schönfelder 1973; Sparks 1994; Predehl & Schmitt 1995; Predehl et al. 2000; Heinz et al. 2016). When light from a bright source, such as a supernova, reaches a cloud of dust, some of the light is scattered in different directions and eventually reaches an observer on Earth. Sparks (1994) proposed to measure distances to historic supernovae by observing highly polarized circles of scattered light around these objects. Because the maximum polarization occurs for a scattering angle of  $90^\circ$ , the radius of the light echoes can be used as a ruler and corresponds to the length  $ct$ , where  $t$  is the time since the supernova exploded and  $c$  is the speed of light. Using that technique, Sparks et al. (2008) found the geometric distance to V838 Monocerotis, an unusual, variable star with a light echo that appeared after an outburst in 2002. Even if the scattering angle of maximum polarization is not  $90^\circ$ , the method may still be used if the scattering angle is known. Kervella et al. (2014) obtained imaging polarimetry of the reflection nebula around the long-period Galactic Cepheid RS Puppis and utilized two polarization models (one based on Milky Way dust mixture and the other assuming Rayleigh scattering) to retrieve the scattering angle from the degree of polarization. By taking into account the dust distribution in the nebula, they adjusted a model of the phase lag of photometric variations over specific nebular features to calculate the distance to RS Puppis.

We determined the distance to SN 1987A by means of a light echo of SN 1987A, AT 2019xis. We obtained imaging polarimetry observations of AT 2019xis and in combination with the publicly available photometry of AT 2019xis, applied the radiative transfer model of the SN 1987A light echo developed in Ding et al. (2021) to determine the distance to SN 1987A.

The light echo, AT 2019xis ( $\alpha = 05:36:13.700$ ,  $\delta = -69:16:24.70$ ), was discovered by the Optical Gravitational Lensing Experiment (OGLE) on 2019 October 15 (Gromadzki & Wyrzykowski 2019). It was initially classified on 2019 December 28 by ePESSTO+ as an SN 1987A-like Type II supernova (Antilen et al. 2019) at  $\sim 70$  days after the shock breakout. However, AT 2019xis was later reclassified as a light echo of SN 1987A from the tip of a dust cloud (Taubenberger et al. 2019), which is also visible in archival HST images, located at an angular distance of  $4'.05$  from SN

1987A. Figure 1 shows images of AT 2019xis observed at three different epochs compared to the HST image.

In Section 2 we describe the observations and data reduction, in Section 3 we describe the model and method for the distance determination, and in Section 4 we report the results and discuss the present findings in comparison with the literature.

## 2. Data and Observations

### 2.1. Imaging Polarimetry

We obtained imaging linear polarimetry of the light echo AT 2019xis with the FOCal Reducer/low dispersion Spectrograph 2 (FORS2; Appenzeller et al. 1998) mounted on the primary focus of European Southern Observatory’s (ESO) Very Large Telescope (VLT) Antu (Prog. ID 2104.C-5031(A), PI Cikota). The data were taken on 2020 January 3, during good weather conditions. The object was at an airmass between 1.4 and 1.5 during the observations. The DIMM seeing was between  $\sim 0''.8$  and  $1''.0$ , and the Moon was set.

The observations were taken in three different passbands, the FORS2 V\_HIGH ( $\lambda_0 = 555$  nm, FWHM = 123.2 nm), R\_SPECIAL ( $\lambda_0 = 655$  nm, FWHM = 165 nm), and I\_BESS ( $\lambda_0 = 768$  nm, FWHM = 138 nm) filters, at four half-wave plate (HWP) angles of  $0^\circ$ ,  $22.5^\circ$ ,  $45^\circ$ , and  $67.5^\circ$ . The redundancy of four HWP angles reduces the flat-fielding issue and cancels out other instrumental effects (see Patat & Romaniello 2006). The sequence was repeated twice in the V and I bands to increase the signal-to-noise ratio and taken only once in the R band.

The brightness of 2019xis at the time of the polarimetry observations was  $18.93 \pm 0.03$  mag in the I band, as measured by the OGLE project (Udalski et al. 2015). The exposure times used were 400 s, 250 s, and 200 s, which led to a signal-to-noise ratio of  $\sim 80$ ,  $\sim 85$ , and  $\sim 100$  per exposure in the V, R, and I bands, respectively.

The linear polarization was calculated following the standard approach (ESO 2015; Cikota et al. 2017a). We measured the flux of the light echo in the ordinary and extraordinary beams at all HWP angles using the aperture photometry function from PythonPhot.<sup>8</sup> We used an aperture radius of  $1''.3$ , with the inner and outer sky radii of  $1''.5$  and  $2''.5$ , respectively. Tests with slight variations of the aperture and sky radii produced consistent results within the  $1\sigma$  uncertainties.

<sup>8</sup> <https://github.com/djones1040/PythonPhot>

**Table 1**  
The Imaging Polarimetry Measurements of the Light Echo AT 2019xis

| Start UT Date       | Passband | $q$<br>(%)      | $u$<br>(%)      | $\theta$<br>(deg) | $p$<br>(%)      | $q_{\text{ISP}}$<br>(%) | $u_{\text{ISP}}$<br>(%) | $p_{\text{ISPcorr}}$<br>(%) |
|---------------------|----------|-----------------|-----------------|-------------------|-----------------|-------------------------|-------------------------|-----------------------------|
| 2020-01-03 03:52:16 | <i>V</i> | $3.30 \pm 0.63$ | $0.05 \pm 0.67$ | $0.4 \pm 5.5$     | $3.18 \pm 0.63$ | $-0.20 \pm 0.07$        | $0.51 \pm 0.08$         | $3.42 \pm 0.64$             |
| 2020-01-03 05:02:54 | <i>R</i> | $3.67 \pm 0.82$ | $0.36 \pm 0.95$ | $2.8 \pm 6.4$     | $3.51 \pm 0.82$ | $-0.48 \pm 0.09$        | $0.50 \pm 0.12$         | $3.99 \pm 0.83$             |
| 2020-01-03 05:32:44 | <i>I</i> | $3.97 \pm 0.48$ | $0.28 \pm 0.38$ | $2.0 \pm 3.5$     | $3.92 \pm 0.48$ | $-0.59 \pm 0.09$        | $0.30 \pm 0.05$         | $4.51 \pm 0.49$             |

**Note.** The columns labeled  $q_{\text{ISP}}$  and  $u_{\text{ISP}}$  denote the Stokes  $q$  and  $u$  parameters of the interstellar polarization (ISP), as determined from field stars. The column labeled  $p_{\text{ISPcorr}}$  denotes the ISP corrected polarization of AT 2019xis. Both the observed polarization and the ISP corrected polarization have been corrected for the polarization bias.

The normalized Stokes  $q$  and  $u$  parameters were calculated as described in the FORS2 User Manual (ESO 2015; see also Cikota et al. 2017a and Chu et al. 2022):

$$q = \frac{2}{N} \sum_{i=0}^{N-1} F(\theta_i) \cos(4\theta_i),$$

$$u = \frac{2}{N} \sum_{i=0}^{N-1} F(\theta_i) \sin(4\theta_i), \quad (1)$$

where  $N$  is the number of the half-wave retarder plate angles;  $\theta_i = 22.5^\circ \times i$ , with  $0 \leq i \leq 3$ ; and  $F(\theta_i)$  is the normalized flux difference between the ordinary ( $f^o$ ) and extraordinary ( $f^e$ ) beams:

$$F(\theta_i) = \frac{f^o(\theta_i) - f^e(\theta_i)}{f^o(\theta_i) + f^e(\theta_i)}. \quad (2)$$

The polarization position angles of the raw measurements have been corrected for the HWP zero angle chromatic dependence (Table 4.7 in ESO 2015). Finally, we calculated the degree of linear polarization,  $p$ , and the polarization angle,  $\theta$ :

$$p = \sqrt{q^2 + u^2}, \quad \theta = \frac{1}{2} \arctan(u/q) \quad (3)$$

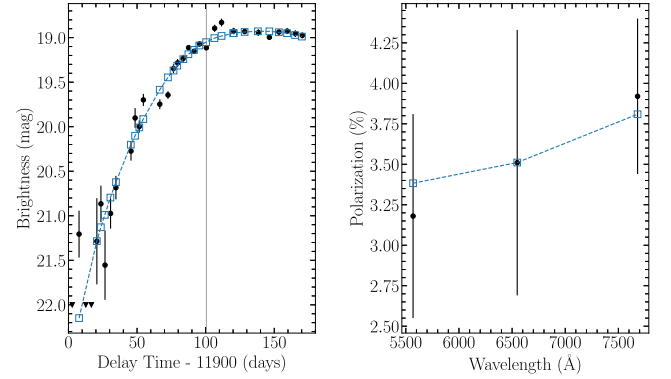
and apply a polarization bias correction following Wang et al. (1997). The imaging polarimetry results are summarized in Table 1 and shown in Figure 2.

### 2.2. Light Curve of AT 2019xis

We used the publicly available light curve<sup>9</sup> of AT 2019xis from the OGLE project database (Udalski et al. 2015). The light echo was observed as part of their OGLE-IV sky survey with the 1.3 m Warsaw University Telescope located at the Las Campanas Observatory in Chile in the *I* band at 29 epochs between 2019 October 2 and 2020 March 13. There are also upper limit (nondetection) measurements available before 2019 October 2. The light curve is shown in Figure 2.

### 2.3. Interstellar Polarization and Dust Extinction

The LMC itself is very dust rich, and there is also some Galactic foreground dust along the line of sight toward the LMC. Fitzpatrick & Walborn (1990) found a total reddening along the SN 1987A sight line of  $E(B - V) \simeq 0.16$  mag, with the LMC component of  $E(B - V) \simeq 0.10$  mag and the Galactic component of  $E(B - V) \simeq 0.06$  mag. This is consistent with the result by Walker & Suntzeff (1990) who determined an  $E$



**Figure 2.** Light curve in the *I* band (left panel) and polarization measurements (right panel) of the light echo AT 2019xis. Also shown is the simultaneous best fit with the simulated model (blue squares and dashed lines in both panels). The vertical line in the left panel marks the time of the polarization observations, and the triangles are nondetections. The observations shown in this plot were not corrected due to interstellar polarization (ISP) or extinction (see case (i) in Section 3).

$(B - V) = 0.17 \pm 0.02$  mag from photometry of stars near SN 1987A. Furthermore, Fitzpatrick & Walborn (1990) measured the reddening of the star Sk-69° 203,  $E(B - V) = 0.19$  mag, located in the 30 Doradus region 2'2 north of the SN 1987A position.

There are also polarization measurements along the line of sight to SN 1987A available in the literature. Schmidt (1976) reported the Galactic foreground ISP  $p = 0.40\% \pm 0.13\%$  and a position angle of  $20^\circ$  for the region where SN 1987A is located. Jeffery (1991) reported broadband polarimetry of SN 1987A taken at the La Plata Observatory by M. Mendez (2023, private communication) from late 1988 to early 1989. Mendez determined ISP along the sight line to SN 1987A that can be described with the following Serkowski parameters (Serkowski et al. 1975):  $p_{\text{max}} \simeq 1.05\%$ ,  $\lambda_{\text{max}} \simeq 5400 \text{ Å}$  and  $K \simeq 1.1$  and a polarization angle of  $\theta \simeq 34^\circ$ .

The dust cloud from which the light echo propagates is, however, located 4'05 (about  $\sim 200$  pc; Taubenberger et al. 2019) east from SN 1987A. In this region the dust properties along the sight line may differ compared to the sight line toward SN 1987A.

To investigate the dust extinction toward AT 2019xis, we examined the publicly available MUSE data of AT 2019xis (Prog. ID 2104.D-5041(A), PI Kuncarayakti) and looked for the Na I D  $\lambda 5890 \text{ Å}$  and  $\lambda 5896 \text{ Å}$  absorption lines which are well-known tracers of dust and gas in the Milky Way (Poznanski et al. 2012). The spectra of AT 2019xis show no obvious Na I D lines; however, a bright (14.4 G mag) star, Gaia EDR3 4657665228220710912, located 3'5 west from AT 2019xis, displays weak Na I D lines. The median of the

<sup>9</sup> <http://ogle.astrouw.edu.pl/ogle4/transients/transients.html>



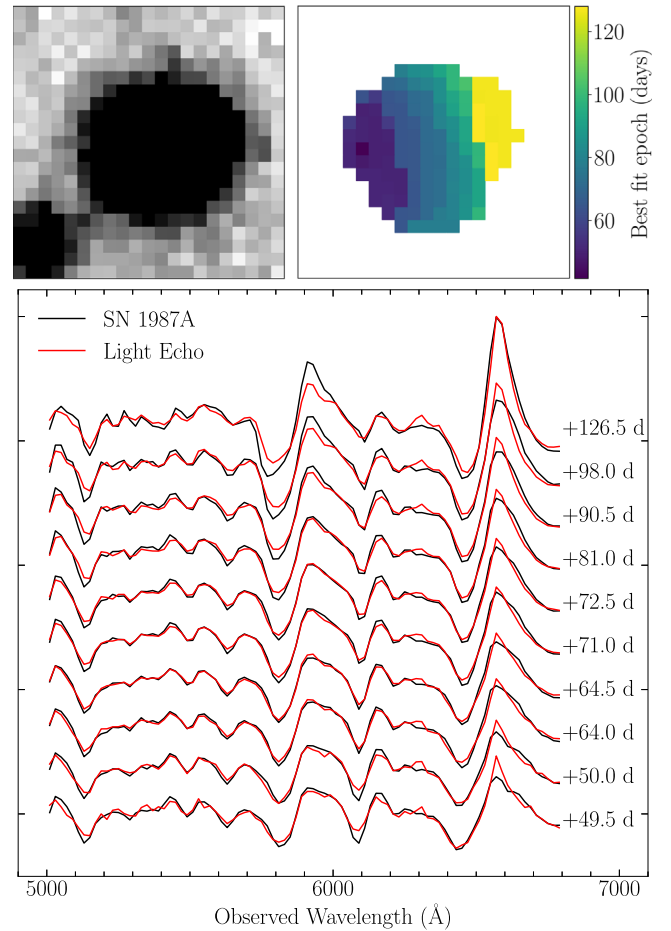
geometric distance posterior from the Earth to the star is  $r_{\text{med}} \sim 14.7$  kpc (Bailer-Jones et al. 2021). Therefore, this line of sight is expected to include a large fraction of Galactic dust. The equivalent widths (EWs) of the Na I D absorption lines are  $\text{EW}(\text{D2}) = 0.23 \pm 0.01 \text{ \AA}$  and  $\text{EW}(\text{D1}) = 0.22 \pm 0.01 \text{ \AA}$ . Following the empirical relation (valid for Milky Way dust) between the sodium absorption and dust extinction by Poznanski et al. (2012), this corresponds to a reddening of  $E(B - V) = 0.05 \pm 0.01 \text{ mag}$  (see their Equation (9)). Note that this value is comparable to the Galactic reddening component determined by Fitzpatrick & Walborn (1990).

We also determined the ISP toward AT 2019xis following the same method as in Chu et al. (2022; see their Figure 2) by calculating the polarization of field stars. The field stars' polarization measurements have been corrected for instrumental polarization, which is negligible in the center of the field but increases up to  $\sim 1\%$  at the borders of the field (Patat & Romaniello 2006; González-Gaitán et al. 2020). The ISP of the field was estimated by taking the weighted mean (with inverse-variance weights) of the field stars' Stokes  $q$  and  $u$ , using 44, 59, and 40 field stars in the V, R and I bands, respectively. The Stokes  $q_{\text{ISP}}$  and  $u_{\text{ISP}}$  are listed in Table 1. Despite a wide range of the field stars' polarization values in all bands, ranging from  $\sim 0.2\%$  to  $\sim 2\%$ , the weighted means in different bands are relatively consistent. The mean polarization in the V, R, and I bands are  $p = 0.53\% \pm 0.08\%$  with a polarization angle of  $\theta = 55.9 \pm 4.0^\circ$ ;  $p = 0.68\% \pm 0.10\%$  with  $\theta = 66.8 \pm 4.2^\circ$ ; and  $p = 0.65\% \pm 0.08\%$  with  $\theta = 76.6 \pm 3.6^\circ$ , respectively. We applied the mean ISP values to correct the polarization measurements of AT 2019xis (see Table 1) and run various simulations, with and without the ISP correction, as explained in Section 3. We also acknowledge that the limitation of this method is that most of the stars used to estimate the ISP are nearby compared to the large distance toward the LMC. This means that there may be more dust along the line of sight toward the LMC producing polarization, which is not being taken into account.

#### 2.4. Geometry of the Dust Cloud

SN 1987's light echoes in particular seem to come from extended dust sheets, of course with some substructure in it, but are in general very large and roughly perpendicular to the line of sight. Many of these light echoes are relatively unchanged for years, i.e., the arcs do not change significantly in shape and apparent motion (e.g., Suntzeff et al. 1988; Xu et al. 1995; Rest et al. 2005). However, AT 2019xis is special because it is one of the brighter SN 1987A light echoes, and it appears as a point source. Figure 1 shows that the light echo coincides with a tip of a dust pillar visible in the HST image, which is rising from a large dust structure in the east (not shown in the cutout). The comparison of the observations taken on 2020 January 3 and 2020 April 20 also shows that AT 2019xis does not change significantly in the apparent position and shape throughout the period of the available photometry from OGLE. Additional observations taken on 2021 April 12 with the European Southern Observatory (ESO) Faint Object Spectrograph and Camera (EFOSC2) mounted on the New Technology Telescope (NTT) show that the AT 2019xis faded away as the light echo propagated further and is scattering from another overdensity  $\sim 5''$  east from AT 2019xis.

AT 2019xis was observed with MUSE at two epochs, on 2020 February 19 and 2020 April 20 (Prog. ID 2104.D-5041



**Figure 3.** Top left panel: light echo AT 2019xis observed with MUSE on 2020 April 20. The size of the cutout is  $4'' \times 4''$  and is oriented so that north is up and east is left. We measured the image quality FWHM of  $\sim 0.7''$  ( $\sim 3.6$  pixel) using neighboring stars. Top right panel: epoch map of the light echo determined by matching template spectra of SN 1987A to each spaxel. The epochs are given relative to the SN explosion. Bottom panel: an example sequence of the light-echo spectra (red lines) compared to the best-matched SN 1987A templates (black lines). The best-fit epochs are indicated.

(A), PI Kuncarayakti). A detailed analysis of the MUSE data is out of the scope of this paper, but we used the extracted light-echo spectra to investigate the shape and morphology of the dust cloud producing AT 2019xis. We cleaned the light-echo spectra using wavelet decomposition (Holschneider et al. 1989; see also Wagers et al. 2010; Cikota et al. 2019) and matched them to SN 1987A spectral templates. The templates have a cadence of 0.5 days and were generated by interpolating original SN 1987A spectra observed at more than 80 epochs by Phillips et al. (1988, 1990). To match the spectra we followed the methods used in the Supernova Identification tool (Tonry & Davis 1979; Blondin & Tonry 2007).

Figure 3 shows the epoch map of the light echo, derived by matching the light-echo spectra from each spaxel to SN 1987A template spectra. The light echo propagates from right to the left in the image (eastwards) and at the diameter of the dust cloud ( $\sim 1''$ ) is much smaller than the angular distance to the supernova ( $\sim 4.05''$ ), so the light waves are propagating nearly parallel. However, the points of equal epochs display some curvature. The curvature can be explained either due to the spherical shape of the cloud in the case of an optically thick cloud or due to distance differences of the scattering plane in

the case of optically thin dust structures. We also note that the light-echo spectra match relatively well with the SN 1987A spectra (see bottom panel in Figure 3) and that the best-fit epochs are continuously increasing from east to west, which implies that there is no significant overlapping of the spectra of different epochs, suggesting that the cloud is optically thick. Additionally, the epoch map is consistent with the dust pillar’s spherical appearance visible in the HST image. Thus, assuming a spherical dust cloud seems to be a reasonable simplification for modeling this light echo in the Section 3.

### 3. Model and Distance Determination

Ding et al. (2021) used a Monte Carlo radiative transfer model (MCRTM) to simulate light-echo observations (light curve and polarization) of SN 1987A at the position of AT 2019xis. They used *UBVRI* light curves of SN 1987A obtained at the Sutherland field station of the South African Astronomical Observatory as the input (Catchpole et al. 1987, 1988, 1989; Menzies et al. 1987; Whitelock et al. 1988) and assumed a simplified dust cloud with similar geometry to AT 2019xis to study how the simulated observations vary depending on the dust model properties, dust cloud size, and shape. They adopted optical properties of the dust-grain models developed by Weingartner & Draine (2001) and investigated the simulated observational differences for different dust-grain size distributions for the Milky Way, LMC, and Small Magellanic Cloud (SMC) dust models.

We adopt the MCRTM by Ding et al. (2021) to determine the distance of the SN 1987A. The apparent brightness and shape of the light curve of the echo are dependent on the intrinsic SN light curve, distance, dust cloud size, shape, optical thickness, and the optical properties of the dust grains. Because we know the path-length difference between the light of SN 1987A traveling directly toward Earth and the light echo, the distance of the SN 1987A determines the scattering angle, with which the light from the supernova is scattered. The scattering geometry was illustrated in Ding et al. (2021), where the distance to the SN 1987A was assumed to be 51.4 kpc (J. Ding 2023, private communication). However, assuming that we do not know the distance to SN 1987A, given the dependence of the light echo’s linear polarization on the scattering angle (see also e.g., Sparks 1994), the scattering angle (and thus, the distance) can be estimated by fitting the MCRTM simulations to the polarization and light-curve measurements.

The linear polarization of the light echo is also strongly dependent on the optical thickness (Ding et al. 2021). Optically thick clouds display large spatial variations of the polarization. For example, in an optically thick spherical cloud, the linear polarization increases in the regions closer to the edge of the cloud.

The dust cloud (AT 2019xis) distance to the Earth, optical thickness, and size of the dust cloud are simultaneously constrained by the light-curve and polarization measurements. The distance between the dust cloud and SN 1987A is also constrained by the time delay and the angular distance between AT 2019xis and SN 1987A (see Ding et al. 2021 for details). The dust cloud is assumed to be a homogeneous sphere, whose size and optical thickness are characterized along with the diameter. The dust particles in the cloud are assumed to be prolate spheroids with an aspect ratio of 4. The optical properties of these particles are computed based on the rigorous invariant imbedding T-matrix method (Johnson 1988;

Yang et al. 2019). A size distribution defined by Weingartner & Draine (2001) as “LMC” is used to obtain size-distribution-averaged dust optical properties.

We conduct the MCRTM simulations and fit the observations after applying different extinction and ISP corrections as follows:

(i) Without applying any ISP and extinction correction, i.e., using the original OGLE light-curve and the ISP-uncorrected polarization measurements. The best fit is shown in Figure 2.

(ii) Same as (i) but assuming halved polarization measurements errors. This test is conducted to investigate the effect of the polarization measurements’ precision on the distance errors.

(iii) After applying the ISP correction to the polarization measurements determined from field stars (determined in Section 2.3, see also Table 1).

(iv) After applying an ISP correction, which follows a Serkowski curve with  $p_{\max} = 1.05\%$ ,  $\lambda_{\max} = 5400 \text{ \AA}$ , and  $K = 1.1$ , as determined in Jeffery (1991).

(v) After applying the ISP correction to the polarization measurements determined from field stars (determined in Section 2.3, see Table 1) and a reddening correction to the light curve of  $E(B - V) = 0.06 \text{ mag}$  assuming Milky Way dust with  $R_V = 3.1$  (see Section 2.3). Following Cardelli et al. (1989), this corresponds to the extinction in the  $I$  band of  $A_I = 0.035 \text{ mag}$ . For comparison, please note that in the case of LMC dust with  $R_V = 2.6$ , the extinction is very similar,  $A_I = 0.033 \text{ mag}$ .

(vi) After applying the ISP correction to the polarization measurements determined from field stars (see Section 2.3 and Table 1) and a reddening correction to the light curve of  $E(B - V) = 0.16 \text{ mag}$  assuming Milky Way dust with  $R_V = 3.1$  (see Section 2.3). Following Cardelli et al. (1989), this corresponds to  $A_I = 0.094 \text{ mag}$ . For comparison, assuming LMC dust with  $R_V = 2.6$ , the extinction would be  $A_I = 0.088$ .

(vii)–(xi). Same as case (vi), but instead of using the LMC dust model, we run the fitting for different dust models from Weingartner & Draine (2001): LMC2, Milky Way dust with  $R_V = 3.1$  (MW3.1), MW4.0, MW5.5, and SMC.

(xii) After applying the ISP correction, which follows a Serkowski curve with  $p_{\max} = 1.05\%$ ,  $\lambda_{\max} = 5400 \text{ \AA}$ , and  $K = 1.1$ , as determined in Jeffery (1991), and a reddening correction to the light curve of  $E(B - V) = 0.06 \text{ mag}$  assuming Milky Way dust with  $R_V = 3.1$ . Following Cardelli et al. (1989), this corresponds to  $A_I = 0.035 \text{ mag}$ .

(xiii) After applying the ISP correction, which follows a Serkowski curve with  $p_{\max} = 1.05\%$ ,  $\lambda_{\max} = 5400 \text{ \AA}$ , and  $K = 1.1$ , as determined in Jeffery (1991), and a reddening correction to the light curve of  $E(B - V) = 0.16 \text{ mag}$  assuming Milky Way dust with  $R_V = 3.1$  (see Section 2.3). Following Cardelli et al. (1989), this corresponds to  $A_I = 0.094 \text{ mag}$ .

### 4. Results and Discussion

As shown in the right panel in Figure 2, the observed trend in the degree of polarization of scattered light from the dust cloud shows a possible increase toward the red wavelengths. However, we note that the measurement errors are relatively large and that within the errors, the degree of polarization may also be consistent with the rise toward the blue wavelengths. A significant polarization, however, is detected in the  $I$  band.

One of the misconceptions of polarization by dust scattering is that polarization is highest in the blue, as quoted in recent papers on the continuum polarization observed along the lines

**Table 2**  
Distance Determination Results with Different Models

| Case   | Dust Model | ISP Correction                               | Extinction Correction | Distance (kpc) | Optical Thickness <sup>a</sup> | Diameter (lt-yr) |
|--------|------------|--|-----------------------|----------------|--------------------------------|------------------|
| (i)    | LMC        | No correction                                | No corrections        | 53.44 ± 2.51   | 3.38 ± 0.23                    | 1.36 ± 0.05      |
| (ii)   | LMC        | No correction (assuming half errors of $p$ ) | No corrections        | 53.45 ± 1.37   | 3.37 ± 0.23                    | 1.36 ± 0.05      |
| (iii)  | LMC        | Field stars (this work)                      | No correction         | 49.09 ± 2.16   | 4.20 ± 0.44                    | 1.29 ± 0.09      |
| (iv)   | LMC        | Jeffery (1991)                               | No correction         | 58.72 ± 3.31   | 2.63 ± 0.23                    | 1.50 ± 0.04      |
| (v)    | LMC        | Field stars (this work)                      | $A_I = 0.035$ mag     | 49.30 ± 2.16   | 3.87 ± 0.07                    | 1.26 ± 0.01      |
| (vi)   | LMC        | Field stars (this work)                      | $A_I = 0.094$ mag     | 49.67 ± 2.17   | 3.50 ± 0.08                    | 1.26 ± 0.02      |
| (vii)  | LMC2       | Field stars (this work)                      | $A_I = 0.094$ mag     | 49.77 ± 2.14   | 3.68 ± 0.32                    | 1.34 ± 0.07      |
| (viii) | MW3.1      | Field stars (this work)                      | $A_I = 0.094$ mag     | 54.87 ± 2.35   | 1.22 ± 0.12                    | 1.68 ± 0.02      |
| (ix)   | MW4.0      | Field stars (this work)                      | $A_I = 0.094$ mag     | 52.65 ± 2.29   | 2.03 ± 0.34                    | 1.39 ± 0.04      |
| (x)    | MW5.5      | Field stars (this work)                      | $A_I = 0.094$ mag     | 49.23 ± 2.12   | 3.72 ± 0.12                    | 1.30 ± 0.03      |
| (xi)   | SMC        | Field stars (this work)                      | $A_I = 0.094$ mag     | 53.19 ± 2.39   | 1.89 ± 0.31                    | 1.42 ± 0.03      |
| (xii)  | LMC        | Jeffery (1991)                               | $A_I = 0.035$ mag     | 58.92 ± 3.34   | 2.46 ± 0.27                    | 1.52 ± 0.04      |
| (xiii) | LMC        | Jeffery (1991)                               | $A_I = 0.094$ mag     | 59.39 ± 3.27   | 2.20 ± 0.28                    | 1.57 ± 0.04      |

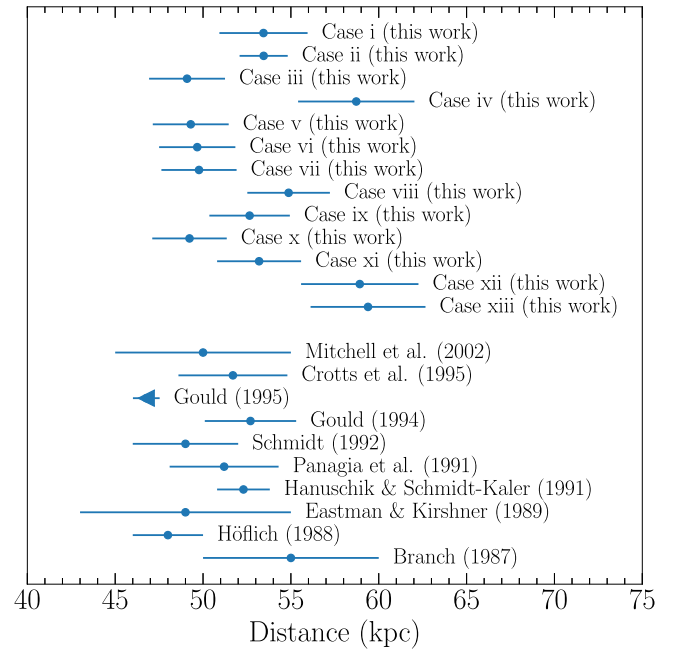
**Note.**

<sup>a</sup> The optical thickness is given at  $0.8 \mu\text{m}$ .

of sight toward Type Ia supernovae (see, e.g., Cikota et al. 2017b, 2019; Chu et al. 2022). However, this is correct only when the cloud is optically thin at the red wavelengths. In contrast, in the case of scattering from optically thick dust clouds, the linear polarization is expected to show a steady increase from  $0.3$  to  $1.0 \mu\text{m}$  (see White 1979; Voshchinnikov & Karjukin 1994; Kartje 1995; Zubko & Laor 2000). Such behavior is, for instance, observed in reflection nebulae (Zellner 1974).

The observed possible increase in polarization toward the red wavelengths in the case of AT 2019xis may suggest that the scattering cloud is optically thick in all three bands ( $V$ ,  $R$ , and  $I$ ). In optically thick clouds, effective scattering only occurs in the outermost layer of the cloud, and the detailed models presented in Ding et al. (2021) demonstrate that polarization can actually peak in the red wavelengths. Additionally, Ding et al. (2021) have demonstrated that if a scattering cloud with an optical depth close to unity can be resolved spatially, the degree of polarization will exhibit color dependence at deeper layers of the cloud. Specifically, the degree of polarization in redder filters is higher at deeper layers. Thus, multiband polarimetry of spatially resolved dust clouds can be utilized as a simple diagnostic tool to determine the optical depth of clouds in general.

The distances to SN 1987A were constrained by fitting the MCRTM simulations (Ding et al. 2021) to the AT 2019xis light curve and polarization observations. Examples of the simulated light-echo images are shown in Section 3.2 in Ding et al. (2021). The results are shown in Figure 4 in comparison to literature distances and summarized in Table 2. Figure 4 shows that all the distances are in general comparable to the literature values. Also, the distance uncertainties of  $\sim 5\%$  are comparable to the uncertainties achieved with other methods, the expanding photosphere method, and distance determinations using the circumstellar rings. However, our uncertainties are highly dependent on the polarization uncertainty. When comparing the results of cases (i) and (ii), it becomes evident that by reducing the assumed polarization uncertainty by half, the precision of the distance estimate approximately doubles (see Table 2). Furthermore, the degree of polarization also has a high impact on the distance results. After correcting our initial polarization measurements (see case (i) with no corrections applied) for the



**Figure 4.** Distance measurements to SN 1987A determined in this work (cases (i)–(xiii)), compared to the values from the literature. The distances determined by Branch (1987), Höflich (1988), Eastman & Kirshner (1989), Hanuschik & Schmidt-Kaler (1991), Schmidt et al. (1992), and Mitchell et al. (2002) are based on the expanding photosphere method, while Panagia et al. (1991), Gould (1994), Gould (1995), and Crotts et al. (1995) derived the distance based on the circumstellar rings. The distance determined by Gould (1995), marked with a triangle, is an upper limit. For the description of cases (i)–(xiii), see Section 3 and Table 2.

ISP as determined by Jeffery (1991), we find that the estimates of distances to SN 1987A increase to  $\sim 59$  kpc (see cases (iv), (xii), and (xiii)). In contrast, in cases (iii), (v), and (vi), after correcting our measurements due to the ISP based on the field stars, which effectively leads to an increase of the polarization values (see Table 1), the distance estimates to SN 1987A decreased to  $\sim 49$  kpc.

In contrast, the distance is only minimally affected by the extinction correction of the light curve. This can be seen by comparing case (iii), where no extinction correction is applied,



to case (vi), where an extinction correction of  $A_I = 0.094$  mag is applied. The resulting distances are  $49.09 \pm 2.16$  kpc and  $49.67 \pm 2.17$  kpc without and with the extinction correction, respectively (Table 2).

The analysis of cases (vi)–(xi) highlights how different dust models impact distance measurements. When fitting the models with LMC dust, the distance to SN 1987A is estimated to be smaller (around 49 kpc) than when using MW and SMC dust models, which produce distances ranging from 49 to 54 kpc. Although there is some uncertainty introduced by variations in the optical properties of the dust, we consider the use of LMC dust to be a reasonable assumption. This is because the properties of the LMC dust model were derived from independent observations of the LMC, as detailed in Weingartner & Draine (2001). The assumption of the geometry of the scattering dust cloud may also introduce an uncertainty in the distance estimate, but the uncertainty is less than that resulting from the optical properties of the dust particles.



Therefore, assuming the LMC dust model is correct, the ISP correction is the main source of uncertainty in distance estimation. Consequently, the distance to SN 1987A can be constrained to a range between  $49.30 \pm 2.16$  kpc (case v) and  $59.39 \pm 3.27$  kpc (case (xiii)). The weighted mean distance obtained from cases (v), (vi), (xii), and (xiii) is  $52.4 \pm 1.3$  (stat)  $\pm 4.8$  (sys) kpc.

In summary, we have demonstrated that the use of polarization measurements of light echoes, combined with photometric measurements, can be utilized to estimate distances to supernovae or other transients in our Milky Way and beyond, with a level of accuracy that is comparable to other techniques. However, this method does have some limitations as it relies on several assumptions regarding dust model properties, dust-grain size distributions, dust cloud size, and the geometric shape of the dust cloud. Despite these limitations, one of the most significant sources of uncertainty is the polarization uncertainty. By reducing this uncertainty, for example, by obtaining more polarimetry observations, it may be possible to achieve greater precision than what can be obtained with other methods.

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