

The need for complementary techniques for reliable characterization of MoS₂-like layers

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ABSTRACT

The observation of characteristic A_{1g} and E_{2g}^1 peaks, at around 382 cm⁻¹ and 408 cm⁻¹ respectively, in Raman spectroscopy is considered evidence of 2H-structured MoS₂, probably the most extensively-studied transition-metal dichalcogenide. Here, using a combination of X-ray diffraction, X-ray photoelectron spectroscopy, and resonant Raman spectroscopy, we show that detection of A_{1g} and E_{2g}^1 modes in Raman spectra alone may not necessarily imply the presence of MoS₂. A series of Mo-S films, $\approx$ 20-nm-thick, are grown on single-crystalline Al₂O₃(0001) substrates at 1073 K as a function of H₂S partial pressure, p_{H_2S} ($=$ 0, 0.01%, 0.1%, and 1% of total pressure) *via* ultra-high vacuum dc magnetron sputtering of a Mo target in 20 mTorr (2.67 Pa) Ar/H₂S gas mixtures. In pure Ar discharges and with p_{H_2S} up to 0.1%, i.e. $p_{H_2S} \leq 2.67 \times 10^{-3}$ Pa, we obtain body centered cubic (bcc), 110-textured films with lattice parameter a increasing from 0.3148 nm (in pure Ar), to 0.3151 nm (at $p_{H_2S} = 2.67 \times 10^{-4}$ Pa), and 0.3170 nm (at $p_{H_2S} = 2.67 \times 10^{-3}$ Pa), which we attribute to increased incorporation of S in the Mo lattice. With 1% H₂S, i.e. $p_{H_2S} = 2.67 \times 10^{-2}$ Pa, we obtain 000 l oriented 2H-structured MoS_{2.0 $\pm$ 0.1} layers. Raman spectra of the thin films grown using 0.1% (and 1%) H₂S, show peaks at around 380 cm⁻¹ (382 cm⁻¹) and 412 cm⁻¹ (408 cm⁻¹), which could be interpreted as A_{1g} and E_{2g}^1 Raman modes for 2H-MoS₂. By comparing the Raman spectra of MoS_{2.0 $\pm$ 0.1} and Mo:S thin films, we identify differences in A_{1g} and E_{2g}^1 peak positions and intensities of defect-sensitive peaks relative to the A_{1g} peaks that can help distinguish pure MoS₂ from non-stoichiometric MoS_{2-x} and multiphase Mo:S materials.

I. INTRODUCTION

Transition metal dichalcogenides (TMDCs) of the form MX_2 , where M is a transition metal from groups 4, 5, and 6 and X is a chalcogen (S, Se, Te), are van der Waals (vdW) bonded layers made of covalently bonded X-M-X sheets. Since the discovery of free-standing graphene, TMDCs have generated considerable interest for applications in flexible, low power consumption optoelectronics, electrocatalytic hydrogen generation, electron spin and valley based computing technologies, and DNA sequencing.^{1–3} Transition-metal disulfides, such as MoS_2 and WS_2 , are a subset of the TMDCs, historically well-known for their chemical properties with a variety of catalytic applications, more prominently in the processing and purification of liquid transportation fuels.^{4–11} MoS_2 and WS_2 are relatively inexpensive and earth-abundant compared to more precious Ru-based compounds. MoS_2 , in particular, has been extensively studied by the catalysis community^{5–7,9–11} due to its attractive properties, and by the tribology community owing to its layered structure.^{12–14} There are two naturally available crystalline forms of MoS_2 : hexagonal structured 2H- MoS_2 ($\text{P6}_3/\text{mmc}$, with in-plane and out-of-plane lattice parameters, $a = 0.316$ nm and $c = 1.229$ nm, respectively) and rhombohedral 3R- MoS_2 (R3m , $a = 0.316$ nm, $c = 1.837$ nm). Bulk 2H-structured MoS_2 is an indirect bandgap (1.23 eV) semiconductor; in contrast, one-molecule-thick 2H- MoS_2 layer is a direct bandgap (1.9 eV) semiconductor.^{15,16} It is predicted that the bandgaps of bilayer MoS_2 and other TMDCs decrease with increasing applied electric field and eventually they exhibit metallic behavior.¹⁷ Luminescence quantum efficiency of monolayer MoS_2 is $\times 10^4$ higher than its bulk counterpart,^{16,18,19} high carrier mobility,^{20,21} high Seebeck coefficient,²² high photoconductivity,^{23,24} strong exciton bonding,²⁵ and environment sensitivity.^{26,27} MoS_2 and other TMDC layers exhibit quantum spin Hall effect;²⁸ valley Hall effect²⁹ and valley polarization¹ phenomena, which helped create valleytronics, an exciting new

field of physics.^{30–32} With their ability to form atomically-thin diodes,^{33–36} exhibit strong size-dependent morphology and electronic structure,³⁷ and potentially tunable optoelectronic properties,^{38,39} TMDCs are attractive as electrically switchable light sources,⁴⁰ high-performance logic devices,⁴¹ atomically-thin electronic^{21,24,42–46} and optoelectronic devices (e.g., solar cells, photodiodes),^{22–24,33,35,44,47–53} flexible transistors,^{54–56} chemical sensors,⁵⁷ and more.

As one may expect, properties of the 2D layers are sensitive to the layer composition, crystallinity, vacancies, substitutional dopants, and domain boundaries, number of layers (i.e. thickness), stacking sequence, and relative orientations of the layers. For example, band gaps of pseudobinary alloys (e.g., $\text{MoS}_{2-y}\text{Se}_y$, $\text{Mo}_{1-x}\text{W}_x\text{S}_2$) made by substitution of either cations (e.g., Mo, W) or anions (e.g., S, Se) vary continuously with the alloy composition.^{58–61} Tunable band gaps and photocatalytic activities as well as unique electronic properties (e.g., MoS_2/WS_2) can be achieved by vertical stacking of TMDCs.⁶¹ In polydomain (analogous to polycrystallinity in bulk crystals) TMDC layers, density and structure of domain boundaries (e.g., zigzag, twin, tilt) strongly influence the electronic and optoelectronic properties.^{62–65} Doping of monolayer MoS_2 affects its structural, electronic and magnetic properties;⁶⁶ S-vacancies in TMDCs (e.g., MoS_2 and WS_2) increase conductance,⁶⁷ lead to asymmetry in electron and hole conduction characteristics,⁶⁸ and modify optical properties of the layers;⁶⁹ incorporation of O-atoms in MoS_2 affects its electronic and optoelectronic characteristics;^{70–73} H-atom adsorption and point defects in MoS_2 influence its ferromagnetic properties;⁷⁴ impurities on the substrate supporting MoS_2 layers can affect the MoS_2 surface work function and its electronic conductivity;^{75,76} moisture and ambient light affect MoS_2 field-effect transistor performance;²⁷ surface defects (e.g., step edges and step structure) dictate thermochemical stability of MoS_2 layers.⁷³ Placing graphene on MoS_2 can affect the MoS_2 band gap,⁷⁷ whose nature and magnitude depends on the interlayer orientation.⁷⁸ The

interlayer interaction strength in multilayer heterostructures can vary with the layer composition.⁷⁹ Therefore, accurate knowledge of, and control over the TMDC layer composition and crystallinity are critical for the realization of TMDC-based devices with optimal performance.

Several methods exist for the synthesis of TMDCs, especially MoS₂ and WS₂.⁸⁰ Here, we list a select few that have been used to grow mono/few-layer TMDCs: (i) Probably the simplest method of preparing individual sheets or hetero-layered stacks of TMDCs is *via* mechanical exfoliation followed by layer transfer onto a substrate of choice. A variant to this approach is liquid exfoliation process, using which a variety of TMDCs (e.g., MoS₂, MoTe₂, TaSe₂, etc.) are dispersed in solution and deposited as individual layers.⁸¹ (ii) A more commonly used method for the synthesis of TMDC layers is *via* chemical vapor transport of metal-oxide and elemental S powder or H₂S gas as precursors.^{63,82–86} A variant of this method uses metal-halides (e.g., MoCl₅, WCl₆, NbCl₅) instead of metal-oxides, and S precursors to grow TMDCs (MoS₂, WS₂, and NbS₂),^{87–90} e.g., $\text{MoCl}_5 + \text{S} \rightarrow \text{MoS}_2$. (iii) Another method involves sulfidation of deposited metal (Mo or W) thin films using elemental S or H₂S gas to form MoS₂ or WS₂,^{91–96} e.g., $\text{Mo} + \text{S}$ or $\text{H}_2\text{S}_{(\text{g})} \rightarrow \text{MoS}_2$. (iv) TMDC layers have also been deposited *via* decomposition of single-source precursors,^{97–100} e.g., thermolysis of ammonium tetrathiomolybdate ($[\text{NH}_4]_2\text{MoS}_4 \rightarrow \text{MoS}_2$); physical vapor deposition (PVD), e.g., sputtering;^{13,101–103} and evaporation/sublimation of TMDCs.^{104,105} Some of these methods have been successfully used for the large-scale synthesis of TMDC layers. However, there are several aspects of the growth processes that are either unknown or not well understood. In particular, the extent of O-incorporation in TMDCs (e.g., MoS_{2-y}O_y) either during growth or during subsequent processing are largely unknown. This knowledge is essential because TMDCs can oxidize at fairly low temperatures in presence of moisture.^{12,14} A related poorly understood aspect is the effect of deposition approach on metal-chalcogen content

in TMDC layers (e.g., thermolysis of thiomolybdates can result in excess S and sulfurization of Mo using H₂S can lead to S-vacancies).

Raman spectroscopy is commonly used for the characterization of 2D layered materials including TMDCs.^{106,107} In case of 2H-MoS₂, the characteristic first order Raman signatures of bulk, phase pure crystal are observed at ~32 cm⁻¹, the rigid layer mode E_{2g}^2 associated with the vibration of adjacent S-Mo-S sheets against each other, at 382~383 cm⁻¹, the E_{2g}^1 mode corresponding to in-plane vibration of S and Mo atoms within a sheet, and at 407~408 cm⁻¹, the A_{1g} mode due to out-of-plane vibration of S atoms within a sheet. (We note that in the back-scattering geometry, the characteristic E_{1g} mode at ~285 cm⁻¹ arising from in-plane vibration of S atoms within a sheet is forbidden.^{108–113}) The observation of Raman peaks characteristic of the 2D layer are generally considered as sufficient evidence for the formation/presence of that material. Here, we focus on the validity of such a conclusion.

In this report, we report on the growth and characterization of Mo-S thin films with varying S content sputter-deposited using a Mo target in 20 mTorr (2.67 Pa) Ar/H₂S gas mixtures as a function of H₂S partial pressure, p_{H_2S} at 1073 K on Al₂O₃(0001) substrates. Using X-ray diffraction (XRD), we determine that films sputter-deposited in pure Ar are body centered cubic (bcc) Mo with 110 texture and those grown using Ar + H₂S gas mixtures with $p_{H_2S} = 2.67 \times 10^{-4}$ Pa and at 2.67×10^{-3} Pa are bcc-structured Mo:S solid solutions. Using higher $p_{H_2S} = 2.67 \times 10^{-2}$ Pa, we obtain basally-oriented 2H-structured MoS_x film. Using X-ray photoelectron spectroscopy (XPS), we determine S-contents in the films grown using $p_{H_2S} = 2.67 \times 10^{-3}$ Pa and 2.67×10^{-2} Pa as $\approx$ 64.3 at.% and 66.7 at.%, respectively. Interestingly, resonant Raman spectra obtained using 633 nm laser from both these films show peaks due to the E_{2g}^1 and A_{1g} Raman modes characteristic of

MoS₂ and peaks associated with structural defects due to longitudinal acoustic (LA) phonons at the M point, referred to as LA(M). Based on our results, we suggest that the mere observation of the E_{2g}^1 and A_{1g} peaks in Raman spectra may not be sufficient to conclude the existence of phase-pure, crystalline MoS₂. We find differences in E_{2g}^1 and A_{1g} peak positions and relative intensities of LA(M)/ A_{1g} peaks can help determine the crystalline quality of MoS_x phase. Based on our results, we suggest that additional characterization maybe necessary to confirm the presence of MoS₂ and to accurately determine the composition and phase purity of the MoS₂-like layers.

II. EXPERIMENTAL

The Mo-S thin films are deposited on single-side polished, $10 \times 2 \times 0.5$ mm³, Al₂O₃(0001) substrates (miscut $< 0.5^\circ$, 99.99% purity, from MTI Corp.) at temperature $T_s = 1073$ K in an ultra-high vacuum (UHV, base pressure $< 5.0 \times 10^{-9}$ Torr, 6.67×10^{-7} Pa) system, described in Refs. ¹¹⁴⁻¹¹⁶. All the details of the substrate cutting, cleaning, and mounting procedures are presented in Ref. ¹¹⁴. Prior to deposition, the substrates are degassed at 1273 K in the UHV chamber until the base pressure is below 8×10^{-7} Pa. With T_s set to 1073 K, we measure up to 100 K differences in temperature along the length of the sample.

We use dc magnetron source with a 50.8 mm diameter $\times$ 3.175 mm thick Mo target (99.95% pure from ACI Alloys Inc.) for the sputter-deposition of Mo and Mo-S thin films using 2.67 Pa Ar + H₂S gas mixtures with 0, 0.01%, 0.1%, and 1% H₂S partial pressures, i.e. $p_{H_2S} = 0$, 2.67×10^{-4} Pa, 2.67×10^{-3} Pa, and 2.67×10^{-2} Pa, respectively. The deposition time is 30 min. The Mo target power is set to 50 W. The target voltages during deposition are 230 ~ 232 V, 243 ~ 248 V, 256 ~ 258 V, 292 ~ 295 V at $p_{H_2S} = 0$, 2.67×10^{-4} , 2.67×10^{-3} , and 2.67×10^{-2} Pa, respectively. Details of the Mo-S deposition procedure are presented in Ref. ¹¹⁴.

The as-deposited Mo-S/Al₂O₃(0001) thin films are characterized using XRD in a Bede D1 high-resolution X-ray diffractometer, resonant Raman spectroscopy with a 633 nm (1.96 eV) laser at 0.4 mW in a Renishaw In-Via Raman confocal microscope system, and XPS in a Kratos Analytical AXIS Ultra DLD following the procedures outlined in Ref. ¹¹⁴. The Raman data presented here are acquired with a resolution of 2 cm⁻¹. We use Gaussian (Lorentzian) functions to extract XRD and Raman peak positions and full widths at half maxima (FWHM).^{117–119} In our Raman measurements, we can access the Raman modes E_{2g}^1 and A_{1g} for MoS₂ but not the low frequency E_{2g}^2 mode as it is obscured due to Rayleigh scattered light. The Raman excitation energy (1.96 eV) corresponds to the direct bandgap of MoS₂, ~1.96 eV, at the K point in the Brillouin zone (BZ).^{109,120} This correspondence with an optical transition can be used to enhance Raman signal intensities and probe phonon modes that are at the BZ boundary (resonant Raman scattering).¹¹¹ XPS data are obtained from the air-exposed Mo-S samples deposited with $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa and 2.67×10^{-4} Pa. The same samples are later sputter-etched with 4 keV Ar⁺ ions rastered across 2×2 mm² with 50 μ A extractor current. Casa XPS software package¹²¹ was used to correct for the background based on Shirley algorithm¹²² and to determine using a Gaussian function with a Lorentzian character [GL(30)] the positions and integrated intensities of Mo 3*d* and S 2*p* peaks. Using the relative sensitivity factors of 3.32 and 0.668, respectively, for Mo 3*d* and S 2*p* (from Casa XPS element library), we calculate the S/Mo contents in the films from the ratios of the integrated intensities of these peaks.

III. RESULTS AND DISCUSSION

Figure 1 is a plot of symmetric XRD 2θ - ω scans obtained from the films deposited using different $p_{\text{H}_2\text{S}}$ values with intensities normalized to the intensity of the Al₂O₃ 0006 reflection and

plotted on logarithmic scales. The two highest intensity peaks labeled s are Al₂O₃ 0006 and 00012 reflections at 2 θ values 41.68° and 90.74°, respectively, from the 0001-oriented single-crystalline α -Al₂O₃ (R $\bar{3}$ c) substrate. We also find two relatively lower intensity peaks at 2 θ values 20.5° and 64.5° in all the scans due to the forbidden Al₂O₃ 0003 and 0009 reflections, respectively.¹²³

Table I. X-ray diffraction peak positions (2 θ in degrees) tabulated for data shown in Fig. 1. For comparison, 2 θ_0 values expected for bcc-Mo,¹²⁴ bulk 2H-MoS₂,¹²⁵ and 3R-MoS₂¹³⁷ are also presented. We define $\Delta = 100 \times (2\theta - 2\theta_0) / 2\theta_0$ as the percentage deviation between the expected and observed 2 θ values for a given phase. Data in italics and within the parenthesis correspond to the 3R-MoS₂ phase.

bcc-Mo		Mo [pure Ar]	Mo:S [0.01% H ₂ S]	Mo:S [0.1% H ₂ S]	2H-MoS ₂ 3R-MoS ₂	MoS _x [1% H ₂ S]	
hkl	2 θ_0	2 θ Δ %	2 θ Δ %	2 θ Δ %	hkil	2 θ_0	2 θ Δ %
110	40.50	40.46 -0.1	40.46 -0.1	40.2 -0.7	0002 0003	14.40 <i>14.45</i>	14.36 -0.3 (-0.6)
220	87.63	87.58 -0.06	87.52 -0.1	—	0004 0006	29.03 <i>29.14</i>	28.94 -0.3 (-0.7)
222	115.97	116.02 0.04	115.72 -0.2	—	0006 0009	44.16 <i>44.34</i>	44.04 -0.3 (-0.7)
					0008 00012	60.16 <i>60.42</i>	59.99 -0.3 (-0.7)

Table I shows 2 θ values of all the film peaks observed in the XRD data for all the Mo-S samples along with 2 θ_0 values expected for bcc bulk Mo ($a_0 = 0.3147$ nm)¹²⁴ and 000*l* (with $l = 2, 4, 6,$ and 8) reflections of 2H-MoS₂ ($c_0 = 1.229$ nm).¹²⁵ The XRD data (black curve in Figure 1) from the sample deposited using pure Ar, i.e., $p_{\text{H}_2\text{S}} = 0$, shows three peaks corresponding to 110, 220, and 222 reflections of bcc-Mo. Relative intensity of 222 peak with respect to 110 peak is 0.03, which is lower than the value expected for randomly oriented polycrystalline Mo,¹²⁴

indicative of 110-texture. We measure lattice parameters a from each of the peak positions and determine the arithmetic average a to be 0.3148 ± 0.0002 nm. This value is 0.03% larger than the a_o of bulk bcc Mo,¹²⁴ suggestive of residual stresses in the film. For the films grown using Ar with 0.01% and 0.1% H₂S, XRD data in Figure 1, green and blue curves respectively, show three, and one peaks at slightly lower 2θ values than those expected for pure bcc-Mo (see Table I). We assign the peak positions to bcc-structured solids and determine a values as 0.3151 ± 0.0001 nm and 0.3170 ± 0.0001 nm, respectively, for the films grown using 0.01% and 0.1% H₂S. We note that these a values are 0.13% and 0.73% higher than a_o of compositionally pure bcc-Mo.¹²⁴ We attribute the larger lattice parameters to incorporation of S in the Mo lattice and refer to these films as Mo:S solid solutions.¹²⁶

The red curve in Figure 1 is a typical XRD scan acquired from the film sputter-deposited using $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-2}$ Pa, i.e. Ar + 1% H₂S gas mixture. (A portion of this data is presented in Ref. ¹¹⁴) We find four peaks at 2θ values comparable to those expected for 2H-MoS₂ (see Table I), which we identify as 2H-MoS₂ 000 l ($l = 2, 4, 6,$ and 8) reflections, indicative of 0001-oriented growth. We rule out the 3R-MoS₂ phase based on the differences in the peak positions observed between the expected and experimentally obtained 2θ values (see Table I). From each of the peak positions, we extract out-of-plane lattice parameters c and calculate the average c value to be 1.233 ± 0.0002 nm, which is 0.3% larger than the c_o expected for stoichiometric bulk 2H-MoS₂.¹²⁵ Full width half maximum of the 0002 peak is 0.56° . We will refer to this film as MoS_x/Al₂O₃(0001).

In the following sections, we compare and contrast the Raman and XPS spectra obtained from two samples, hereto referred to as Mo:S/Al₂O₃(0001) and MoS_x/Al₂O₃(0001), sputter-deposited in Ar with 0.1% H₂S ($p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa) and 1% H₂S ($p_{\text{H}_2\text{S}} = 2.67 \times 10^{-2}$ Pa), respectively. Figure 2 shows representative Raman spectra obtained from the Mo:S/Al₂O₃(0001)

(blue curve) and $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$ (red curve) samples using a 633 nm (1.96 eV) laser excitation. The intensities are as measured and plotted without translation on a logarithmic scale. In the Raman data of $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$, we find peaks at $\sim 382 \text{ cm}^{-1}$ and $\sim 408 \text{ cm}^{-1}$, the characteristic E_{2g}^1 and A_{1g} Raman modes, respectively for MoS_2 . Interestingly, Raman spectra from $\text{Mo:S}/\text{Al}_2\text{O}_3(0001)$ show peaks at $\sim 380 \text{ cm}^{-1}$ and at $\sim 412 \text{ cm}^{-1}$, values comparable to those expected for MoS_2 . However, XRD data (Figure 1, blue curve) from the same sample does not show any peaks due to MoS_2 . We note that the E_{2g}^1 and A_{1g} peak intensities are approximately five-fold higher in the Raman spectra of $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$ than in the Raman spectra of $\text{Mo:S}/\text{Al}_2\text{O}_3(0001)$. For both monolayer and multilayer MoS_2 , the E_{2g}^1 and A_{1g} mode peak positions have been found to be sensitive to strain (both modes redshift with increasing strain) and doping (A_{1g} mode broadens and redshifts).^{127–130} Defects induced by ion bombardment¹²⁸ or by electron irradiation¹²⁷ lead to an increase in the separation $\Delta\omega$ ($= A_{1g} - E_{2g}^1$) between the characteristic modes as the E_{2g}^1 mode redshifts and the A_{1g} mode blueshifts, along with the activation of M-point phonons that are at the center of the BZ boundary. (For reference, $\Delta\omega$ is $\sim 19 \text{ cm}^{-1}$ for defect-free monolayer MoS_2 and $\Delta\omega$ is around 25 cm^{-1} for multilayer MoS_2 .^{120,131}) For rf sputtered MoS_2 films that are sulfur-deficient or oxidized, larger values of $\Delta\omega$ ($> 25 \text{ cm}^{-1}$) have been observed.^{132,133} For our $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$, we measure $\Delta\omega$ as $26 \pm 4 \text{ cm}^{-1}$, comparable to the reference value. For $\text{Mo:S}/\text{Al}_2\text{O}_3(0001)$ layers, we find $\Delta\omega$ to be considerably larger, $32 \pm 4 \text{ cm}^{-1}$, presumably due to defects.

In addition to the characteristic modes, due to resonant Raman scattering, the curves in Figure 2 reveal several peaks at $\sim 179 \text{ cm}^{-1}$, $\sim 229 \text{ cm}^{-1}$, 460 cm^{-1} , $\sim 528 \text{ cm}^{-1}$, $\sim 568 \text{ cm}^{-1}$, $\sim 599 \text{ cm}^{-1}$, $\sim 640 \text{ cm}^{-1}$, and $\sim 820 \text{ cm}^{-1}$ for $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$; we find fewer peaks at nearly the same values, $\sim 183 \text{ cm}^{-1}$, $\sim 227 \text{ cm}^{-1}$, $\sim 459 \text{ cm}^{-1}$, and $\sim 641 \text{ cm}^{-1}$, for $\text{Mo:S}/\text{Al}_2\text{O}_3(0001)$. We note that peaks at

$\sim 227 \text{ cm}^{-1}$ and at $\sim 820 \text{ cm}^{-1}$ are also associated with MoO_3 ,^{109,128} but we do not see the other Raman modes of MoO_3 at $\sim 158 \text{ cm}^{-1}$, $\sim 285 \text{ cm}^{-1}$, $\sim 666 \text{ cm}^{-1}$, and 994 cm^{-1} . We therefore rule out the presence of MoO_3 and assign the peaks to Raman modes of MoS_2 as follows. The intense, asymmetric peak at $\sim 459 \text{ cm}^{-1}$ [$\sim 460 \text{ cm}^{-1}$] visible in the Raman spectra of $\text{Mo:S/Al}_2\text{O}_3(0001)$ [$\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$] is interpreted to be a convolution of two peaks – one from a second order phonon process involving the longitudinal acoustic (LA) phonon at the M point, referred to as 2LA(M) at $\sim 454 \text{ cm}^{-1}$ and another, a first order infrared-active optical phonon A_{2u} at $\sim 465 \text{ cm}^{-1}$. The observation of the Raman inactive A_{2u} mode is attributed to resonant excitation.¹¹¹ The frequency of the LA(M) phonon based on that of the 2LA(M) phonon has been calculated to be $\sim 227 \text{ cm}^{-1}$ and its appearance determined to be due to confinement of phonons due to defects such as S-vacancies¹²⁷ and grain boundaries.¹¹¹ Other studies have shown that the relative intensity of the LA(M) phonon mode with respect to the characteristic peaks of MoS_2 increases with structural disorder.^{128,134} Based on these previous studies,^{109,111,120} we assign the peak at $\sim 459 \text{ cm}^{-1}$ [$\sim 460 \text{ cm}^{-1}$] seen in the Raman spectra of $\text{Mo:S/Al}_2\text{O}_3(0001)$ [$\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$] as due to $2\text{LA(M)} + A_{2u}$ phonons; the peaks at $\sim 183 \text{ cm}^{-1}$ [$\sim 179 \text{ cm}^{-1}$], $\sim 227 \text{ cm}^{-1}$ [$\sim 229 \text{ cm}^{-1}$], and $\sim 641 \text{ cm}^{-1}$ [$\sim 640 \text{ cm}^{-1}$] to $A_{1g} - \text{LA(M)}$, LA(M) , and $A_{1g} + \text{LA(M)}$ phonons, respectively; the peaks at $\sim 528 \text{ cm}^{-1}$, $\sim 568 \text{ cm}^{-1}$, $\sim 599 \text{ cm}^{-1}$, and at $\sim 820 \text{ cm}^{-1}$ only seen in the spectrum of $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$ can be assigned to $E_{1g} + \text{LA(M)}$, $2E_{1g}$, $E_{2g}^1 + \text{LA(M)}$, and $2A_{1g}$ phonons, respectively. The fact that we observe overall higher intensities for all Raman peaks, relatively lower intensity LA(M) phonon peak (measured with respect to A_{1g} peak) and additional peaks in the Raman spectrum of the $\text{MoS}_x/\text{Al}_2\text{O}_3(0001)$ sample strongly suggests the presence of highly crystalline MoS_2 phase. In contrast, the higher intensity of the LA(M) phonon peak (compared to the A_{1g} peak) and larger $\Delta\omega$ for the $\text{Mo:S/Al}_2\text{O}_3(0001)$ sample are indicative of higher degree of structural disorder, S-

vacancies, and/or smaller size MoS₂ grains in the film. This is plausible and consistent with the absence of XRD peaks associated with MoS₂ phase in this sample.

Figure 3 shows X-ray photoelectron spectra obtained from air-exposed (a) Mo:S/Al₂O₃(0001) and (b) MoS_x/Al₂O₃(0001) samples. In the plots, we find two peaks around 162 eV associated with S 2*p* and three peaks at energies between 229 and 236 eV due to Mo 3*d* peaks. The S 2*p* peaks in the Mo:S/Al₂O₃(0001) [MoS_x/Al₂O₃(0001)] films at 162.1 eV [161.9 eV] and at 163.3 eV [163.1 eV] correspond to S 2*p*_{3/2} and 2*p*_{1/2} peaks, respectively, as expected for S 2– species.¹³⁵ The observed Mo peaks can be attributed to the presence of Mo +4 and Mo +6 species with two sets of Mo 3*d*_{5/2} and Mo 3*d*_{3/2} doublets: for Mo:S/ Al₂O₃(0001) [MoS_x/Al₂O₃(0001)], a Mo 3*d*_{5/2} peak at 229.3 eV [229.1 eV] and Mo 3*d*_{3/2} peak at 232.4 eV [232.2 eV] colored in dark green and another doublet with Mo 3*d*_{5/2} peak at 232.7 eV [232.5 eV], Mo 3*d*_{3/2} peak at 235.9 eV [235.6 eV] colored in brown. The Mo 3*d* doublets with peaks at lower (higher) binding energies correspond to Mo +4 (Mo +6) species.¹³⁵ The presence of Mo +4 species can be attributed to MoS₂ and MoO₂ while Mo +6 species are associated with MoO₃. From earlier studies,¹³⁶ we know that the as-deposited layers in our experiments are free of oxygen. Moreover, we do not detect MoO₃ or MoO₂ in both XRD and Raman spectra obtained from these samples. XPS data (not shown) obtained from sputter-etched samples do not show any peaks associated with Mo oxides. We therefore suggest that any detection of Mo +6 peaks in the samples is likely a result of surface oxidation upon air-exposure.

We note that the XPS of Mo:S/Al₂O₃(0001) film shows a higher integrated intensity for Mo +6 species than the XPS data from MoS_x/Al₂O₃(0001) film suggestive of a higher degree of oxidation in the Mo:S/Al₂O₃(0001) sample. The XPS of the sputter etched Mo:S/Al₂O₃(0001) [MoS_x/Al₂O₃(0001)] film near the Mo 3*d* binding energy shows an asymmetric Mo 3*d* doublet

with peaks at ~ 228.2 eV [~ 227.9 eV] (Mo $3d_{5/2}$), ~ 231.4 eV [~ 231.1 eV] (Mo $3d_{3/2}$) that we identify as due to metallic Mo (0) species. For Mo:S/Al₂O₃(0001), we are unable to resolve the fairly broad and noisy signal at energies associated with S $2p$ peaks; for MoS_x/Al₂O₃(0001), a S $2p_{3/2}$ (S $2p_{1/2}$) peak was resolved at ~ 162 eV (~ 163.2 eV). Based on these results, we suggest that the entire Mo +4 peak intensity is due to MoS₂ and attribute the Mo +6 peaks to surface oxidation of metallic Mo and/or Mo:S solid solution. From the peak areas of S and Mo +4, we determine the effective *sulfide* composition of Mo:S/Al₂O₃(0001) as MoS_{1.8 $\pm$ 0.1} and that of MoS_x/Al₂O₃(0001) as MoS_{2.0 $\pm$ 0.1}. We note that this quantification approach does not account for metallic Mo present in the films and therefore the overall composition of Mo:S/Al₂O₃(0001) film is likely a combination of Mo + MoS_{1.8 $\pm$ 0.1}. The fact that we do not see in Fig. 1 any XRD peaks associated with MoS₂ but only a weak bcc-Mo reflection in the Mo:S/Al₂O₃(0001) films is consistent with our conclusion.

IV. CONCLUSIONS

In summary, using a combination of materials characterization techniques, we show that the detection of A_{1g} and E_{2g}^1 Raman modes does not necessarily imply that the material is 2H-MoS₂. To this purpose, we deposited a series of Mo-S thin films with varying S-contents by reactive sputtering of a Mo target in 2.67 Pa Ar/H₂S gas mixtures on Al₂O₃(0001) substrates at 1073 K. XRD and XPS measurements indicate the formation of bcc-structured Mo films at $p_{\text{H}_2\text{S}} = 0$ Pa, 2.67×10^{-4} Pa, a bcc Mo:S solid solution with 64 at.% S at $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa, and a 0001-oriented 2H-MoS₂ film at $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-2}$ Pa. Raman spectroscopy results of the films deposited at $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa and 2.67×10^{-2} Pa show characteristic A_{1g} and E_{2g}^1 peaks, commonly attributed to 2H-MoS₂, however, with noticeable differences in the peak separations ($\Delta\omega = A_{1g} - E_{2g}^1$) and peak intensities of the A_{1g} , E_{2g}^1 , and the defect sensitive LA(M) phonon modes. Based on

these results, we suggest that careful interpretation of Raman spectra along with complementary characterization data are highly desirable not only for identification of MoS₂ and other TMDCs but also for accurate determination of the composition and crystallinity of non-stoichiometric compounds.

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Data Availability Statement: The data that support the findings of this study are available from the corresponding author upon reasonable request.

Author Declarations

Conflict of interest: The authors have no conflicts to disclose.

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FIGURES

Figure 1. Symmetric $2\theta:\omega$ X-ray diffraction (XRD) data plotted on a logarithmic scale, obtained from samples sputter deposited in Ar + H₂S atmosphere on Al₂O₃(0001) substrates, with $p_{\text{H}_2\text{S}}$ (partial pressure of H₂S during sputter deposition) = 0%, 0.01%, 0.1%, and 1% (black, green, blue, and red curves) of total gas pressure (2.67 Pa). Film reflections from Mo and MoS_x are labeled as shown while substrate reflections are denoted by s. A portion of the red curve is adapted with permission from A. Deshpande, K. Hojo, K. Tanaka, P. Arias, H. Zaid, M. Liao, M. Goorsky, and S.K. Kodambaka, ACS Appl Nano Mater, (2023). Copyright (2023) American Chemical Society.

Figure 2. Raman spectra obtained from MoS_x samples sputter deposited with $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa (blue) and 2.67×10^{-2} Pa (red), acquired with a 633 nm wavelength laser and plotted on a logarithmic scale. The first order modes – E_{2g}^1 (~ 382 cm⁻¹, red curve; ~ 380 cm⁻¹ blue curve) and A_{1g} (~ 408 cm⁻¹, red curve; ~ 412 cm⁻¹, blue curve) are visible in both curves along with second order modes excited due to resonant Raman scattering – $A_{1g} \pm \text{LA(M)}$, $2\text{LA(M)} + A_{2u}$. The LA(M) mode at (~ 227 cm⁻¹, blue curve; ~ 229 cm⁻¹, red curve) associated with defects, is observed in both curves but with a reduced relative intensity (with respect to the characteristic A_{1g} peak) for the red curve. The difference in the peak positions $\Delta\omega = E_{2g}^1 - A_{1g}$ for both samples is also labeled. Red curve is reprinted with permission from A. Deshpande, K. Hojo, K. Tanaka, P. Arias, H. Zaid, M. Liao, M. Goorsky, and S.K. Kodambaka, ACS Appl Nano Mater, (2023). Copyright (2023) American Chemical Society.

Figure 3. X-ray photoelectron spectra (XPS) obtained from air-exposed Mo-S films deposited with $p_{\text{H}_2\text{S}} =$ (a) 2.67×10^{-3} Pa (blue) and (b) 2.67×10^{-2} Pa (red). Open symbols are raw data while Gaussian-Lorentzian (GL (30)) fits to the data, carried out after a Shirley background (green dashes) subtraction are shown in solid lines – dark green for Mo +4 (MoS_x) and brown for Mo +6 (MoO₃). For $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-3}$ Pa, we determine the film S-content as ~ 64 at.% and it is more oxidized in comparison to the sample sputter-deposited using $p_{\text{H}_2\text{S}} = 2.67 \times 10^{-2}$ Pa, whose composition is MoS_{2.0 $\pm$ 0.1}. Red curve is reprinted with permission from A. Deshpande, K. Hojo, K. Tanaka, P. Arias, H. Zaid, M. Liao, M. Goorsky, and S.K. Kodambaka, ACS Appl Nano Mater, (2023). Copyright (2023) American Chemical Society.