

# Chalcogen isotopes reveal limited volatile contribution from late veneer to Earth

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**Abstract:** The origins of Earth's volatiles are highly debated. Comparing the chalcogen isotope ratios in the bulk silicate Earth (BSE) to those of its possible building blocks, chondrites, can be used to constrain the origins of Earth's volatiles; however, this comparison may be complicated by the isotope effects during protoplanetary differentiation, which largely remain poorly understood. Using first-principles calculations, we find that core-mantle differentiation does not significantly fractionate selenium and tellurium isotopes, while equilibrium evaporation from early planetesimals would enrich selenium and tellurium in heavy isotopes in the BSE. The sulfur, selenium, and tellurium isotopic signatures of the BSE reveal that protoplanetary differentiation plays a key role in establishing most of Earth's volatiles and a late veneer does not significantly contribute to the BSE's volatile inventory.

**One-Sentence Summary:** The Earth obtained its volatiles during its main accretion stage, but not through a patchy veneer arrived late.

## Main Text:

Understanding how Earth accreted its volatile elements, particularly its life-essential elements, is key to understanding the planetary evolution and the habitability of terrestrial planets. One group of hypotheses suggests that proto-Earth accreted mainly from nearly volatile-poor materials from the inner Solar system, with subsequent addition of volatile elements supplied by volatile-rich materials from the outer Solar system to the bulk silicate Earth (BSE) after core formation had ceased (1–4).

The late-addition model, often referred as a “late veneer”, was originally proposed to explain the abundances of highly siderophile elements (HSEs) in the BSE (5, 6), as metal/silicate partitioning indicated their near complete segregation into the metal phase during Earth’s core formation (7). This model was subsequently adopted to account for chalcogen (sulfur (S), selenium (Se), and tellurium (Te)) abundances in the BSE (3) because their concentrations in the Earth’s mantle are much higher than expected from their metal/silicate partition coefficients (8, 9). Further, mantle peridotites were found to have chalcogen/HSE abundance ratios similar to those of carbonaceous chondrites, providing evidence for the late delivery of carbonaceous-chondrite-like material to Earth’s mantle (3), although it remains debated whether the chalcogen/HSE ratios measured in peridotites represent BSE’s signatures because of the complicated history of these peridotites (10). However, recent experiments suggest that S metal/silicate partition coefficients at core-forming conditions (~20-60 GPa and > 3000 K) (11) are much smaller than that at low pressures, and the modeled S abundance in the BSE after core formation could even be higher than the measured value for the Earth’s mantle. Thus, a late veneer is not needed to explain the S abundance, and possibly volatile elements more generally, in the BSE.

An alternative possibility is that the proto-Earth accreted from volatile-rich materials, and the present-day volatile element abundances in the BSE was set primarily by protoplanetary differentiation processes including planetesimal evaporation and core formation (12–14). Chalcogen isotopes can provide independent constraints on the origin and evolution of Earth’s volatile elements. Earth’s mantle has a lighter S isotope composition ( $\delta^{34}\text{S}$   $= [({}^{34}\text{S}/{}^{32}\text{S})_{\text{sample}}/({}^{34}\text{S}/{}^{32}\text{S})_{\text{std}} - 1] \times 1000$  ‰, where std refers to the standard sample) than any type of chondrites (15, 16) (Fig. S1), whereas its Se and Te isotopic compositions ( $\delta^{82/76}\text{Se}$   $= [({}^{82}\text{Se}/{}^{76}\text{Se})_{\text{sample}}/({}^{82}\text{Se}/{}^{76}\text{Se})_{\text{std}} - 1] \times 1000$  ‰;  $\delta^{128/126}\text{Te}$   $= [({}^{128}\text{Te}/{}^{126}\text{Te})_{\text{sample}}/({}^{128}\text{Te}/{}^{126}\text{Te})_{\text{std}} - 1] \times 1000$  ‰) are similar to those of carbonaceous chondrites but significantly heavier than those of enstatite chondrites (17–19), a postulated dominant isotopic component of Earth’s building materials (20). The sub-chondritic  $\delta^{34}\text{S}$  of the BSE, which cannot be a result of a late veneer, is plausibly explained by evaporation from molten planetesimals, with insignificant S isotope fractionation during core formation (21). This suggests that protoplanetary differentiation may have played a key role in establishing Earth’s volatile element inventory (12–14), rather than a late veneer. In contrast, the similarity in  $\delta^{82/76}\text{Se}$  and  $\delta^{128/126}\text{Te}$  between the BSE and carbonaceous chondrites was argued to be the result of the late addition of carbonaceous-chondrite-like material to Earth (17, 19). Such a conclusion can be drawn only if protoplanetary differentiation processes did not fractionate Se and Te isotopes. However, Se and Te have cosmochemical and geochemical characteristics similar to S (22, 23), and it is important to assess their isotope fractionation during evaporation and core formation.

We conducted first-principles calculations (see details in Supplementary Materials) to obtain the equilibrium Se and Te isotope fractionation factors between silicate and metal and between

vapor and silicate. Selenium and Te are redox-sensitive elements, but there is no literature study investigating the Se and Te species in silicate melts under different redox conditions. By analogy to S, which predominantly occurs as  $S^{2-}$  at  $\log fO_2 < \text{FMQ-1}$  (1 log unit lower than the fayalite-magnetite-quartz (FMQ) buffer) in silicate melts (24, 25), the Se and Te species are likely dominated by  $Se^{2-}$  and  $Te^{2-}$  in silicate melts at the  $fO_2$  conditions of core formation for Earth ( $\log fO_2 < \text{FMQ-4}$ ) (26, 27), respectively. We model the silicate melts using  $Mg_{32}Si_{32}O_{95}Se$ ,  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Se$  (model "pyrolite"), and  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Te$  compositions. The metallic melts are modeled by two multicomponent alloys,  $Fe_{87}Ni_4Si_{10}C_2O_2H_2S_2Se$  and  $Fe_{87}Ni_4Si_{10}C_2O_2H_2S_2Te$ .

The structures of Se-bearing melts at 5-100 GPa and 3000 K were derived from first-principles molecular dynamics (FPMD) simulations. The FPMD results show that Se is mainly bonded to Mg and Si atoms in  $Mg_{32}Si_{32}O_{95}Se$  melt with a Se-Si distance of 2.25-2.29 Å and a Se-Mg distance of 2.47-2.57 Å (Fig. S2), while it is primarily bonded to Fe atoms with a Se-Fe distance of  $\sim 2.2$  Å in  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Se$  melt (Fig. S3). In  $Fe_{87}Ni_4Si_{10}C_2O_2H_2S_2Se$  metallic melt, Se is dominantly bonded to Fe atoms with a Se-Fe distance of 2.20-2.30 Å (Fig. S4). Similarly, Te is also mainly bonded to Fe atoms in  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Te$  and  $Fe_{87}Ni_4Si_{10}C_2O_2H_2S_2Te$  melts (Fig. S5). The bonding configurations of Se and Te in silicate and metallic melts are similar to those of S under relatively reducing conditions (21).

The average force constants,  $\langle F \rangle$ , of Se and Te in melts are controlled by their bonding structures and were calculated using the snapshots from the FPMD simulations based on the harmonic approximation (see Methods). The  $\langle F \rangle$  of Se in  $Mg_{32}Si_{32}O_{95}Se$  melt increases from  $\sim 133$  N/m at 5 GPa to  $\sim 365$  N/m at 82 GPa (Fig. S6) because the coordination numbers of Se-Si and Se-Mg bonds significantly increase with pressure (Fig. S2). The  $\langle F \rangle$  of Se in  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Se$  melt falls on the  $\langle F \rangle$  vs pressure trend of  $Mg_{32}Si_{32}O_{95}Se$  melt, suggesting a negligible compositional effect on  $\langle F \rangle$  of Se in silicate melts (Fig. S6). Similarly, the  $\langle F \rangle$  of Se in the metallic melt increases by  $\sim 1.6$  times from 11 to 100 GPa due to the considerable increase in the coordination number of the Se-Fe bond with compression (Fig. S4). However, the  $\langle F \rangle$  difference between silicate and metallic melts does not significantly change with pressure, which is only 5-25 N/m at 10-80 GPa (Fig. S6). The  $\langle F \rangle$  difference of Te between  $Mg_{30}NaCa_2Fe_4Si_{24}Al_3H_2O_{89}Te$  and  $Fe_{87}Ni_4Si_{10}C_2O_2H_2S_2Te$  melts is  $< 10$  N/m at  $\sim 44$  GPa. For comparison, the  $\langle F \rangle$  difference of S between the reducing silicate melt and the metallic melt is  $< 30$  N/m at  $< 80$  GPa (21).

The equilibrium fractionation value ( $10^3 \ln \alpha_{\text{silicate-melt}}$ ) of  $^{82/76}Se$  and of  $^{128/126}Te$  between silicate and metallic melts are derived from the  $\langle F \rangle$  differences using the high-temperature approximation of the Bigeleisen-Mayer equation (28) (see Supplementary materials). The silicate melt is only marginally enriched in heavy Se and Te isotopes relative to the metallic melt: at 3000 K, the  $10^3 \ln \alpha_{\text{silicate-metal}}$  of  $^{82/76}Se$  is  $< +0.012$  ‰ at 10-60 GPa (Fig. 1), and the  $10^3 \ln \alpha_{\text{silicate-metal}}$  of  $^{128/126}Te$  is essentially zero at  $\sim 44$  GPa.

We further model the Se and Te isotope fractionation during Earth's core formation using two endmember models, equilibrium and Rayleigh distillation models. Residual Se and Te abundances in the BSE after core formation are determined by their metal-silicate partition coefficients. Previous experiments (8) report a large Se partition coefficient (up to  $\sim 120$  at 3000 K) between silicate and metal ( $D_{Se}^{\text{metal/silicate}}$ ) at low pressures ( $< 20$  GPa), which is similar to the silicate-metal S partition coefficient ( $D_S^{\text{metal/silicate}}$ ). However, subsequent experiments (29) at  $> 40$  GPa suggest that S becomes much less siderophile at Earth's core-forming conditions than previously estimated from extrapolation of low-pressure data (8). If this applies to Se,

$D_{\text{Se}^{\text{metal/silicate}}}$  would also be smaller than the low-pressure data at Earth's core-forming conditions. Accordingly, we use a large range of 30-120 for  $D_{\text{Se}^{\text{metal/silicate}}}$ . Our results show that core-mantle differentiation can only shift the BSE's  $\delta^{82/76}\text{Se}$  by at most +0.01 ‰ and +0.05 ‰ for equilibrium and Rayleigh distillation models (Fig. 1), respectively, indicating that core formation cannot explain the BSE's  $\delta^{82/76}\text{Se}$  if Earth accreted predominantly from enstatite-chondrite-like material (20). Similarly, core formation can only shift the BSE's  $\delta^{128/126}\text{Te}$  by at most +0.005 ‰ even if  $D_{\text{Te}^{\text{metal/silicate}}}$  is  $\sim 300$  and  $\delta^{34}\text{S}$  by  $< +0.1$  ‰ (21), which cannot explain the  $\delta^{34}\text{S}$  difference between the BSE and enstatite chondrites.

We now consider the effect of evaporative loss from molten planetesimals that formed proto-Earth. Previous studies (30, 31) suggest that the net isotope fractionation between vapor and melt during planetesimal evaporation strongly depends on the evaporative conditions (30), e.g., the vapor pressure or the vapor saturation degree. If the vapor saturation degree is far lower than 100%, the kinetic effect dominates the net isotope fractionation (30, 31), and the residual melt is always enriched in heavy isotopes after evaporation. This might plausibly explain the heavy  $\delta^{82/76}\text{Se}$  ratio in the BSE relative to the enstatite-chondrite-like material (17), but it cannot explain the BSE's negative  $\delta^{34}\text{S}$  value relative to chondrites (15, 16). When planetesimals undergo evaporation in the presence of nebular  $\text{H}_2$  under a total pressure of  $\sim 10^{-4}$  bar, the vapor saturation degree approaches 100%, and the net isotope fractionation is equal to the equilibrium isotope fractionation between vapor and melt ( $10^3\ln\alpha_{\text{vapor-silicate}}$ ) (30, 31). Our previous thermodynamic calculations using the GRAINS code show that S mainly occurs as  $\text{H}_2\text{S}$  in the vapor phase (Fig. 2), which is enriched in heavy S isotopes relative to the silicate melt (21). Modeling indicates that the sub-chondritic  $\delta^{34}\text{S}$  signature in the BSE can be explained by the evaporative loss of  $\sim 90\%$  S mainly as  $\text{H}_2\text{S}$  from molten planetesimals due to the positive  $10^3\ln\alpha_{\text{vapor-silicate}}$  of S isotopes (Fig. 2d).

Following the framework for S isotopes, we conducted thermodynamic calculations with solar abundances (22) under  $1\text{e}^{-4}$  bar using GRAINS to determine the Se and Te species in the vapor phase. The results show that Se in the vapor phase mainly occurs as atomic Se (g), whose fraction increases from  $\sim 83\%$  at 1300 K to  $\sim 97\%$  at 1500 K (Fig. 2a). The dominant species for Te in the vapor phase is also atomic Te (g), whose fraction is  $> 99\%$  at 1300-1500 K. In contrast, S in the vapor phase mainly occurs as  $\text{H}_2\text{S}$  (g) and/or HS (g) at 1300-1500 K (Fig. 2b). The different species for Se and S in the vapor phase reflect a higher electronegativity for S, leading to a negative Gibbs free energy of formation of  $\text{H}_2\text{S}$  (g), whereas the Gibbs free energy of formation is positive for  $\text{H}_2\text{Se}$  (g) and  $\text{H}_2\text{Te}$  (g) (Fig. S7). That is, at 1300-1500 K and  $10^{-4}$  bar total pressure,  $\text{H}_2\text{S}$  (g) is a stable phase, but  $\text{H}_2\text{Se}$  (g) and  $\text{H}_2\text{Te}$  (g) are not. We also conducted thermodynamic simulations with solar elemental abundances but the H concentration decreased by one order of magnitude, which corresponds to a more oxidizing condition than the solar nebula. The results show that the fractions of major Se and Te species in the vapor phase do not significantly change compared with the foregoing case.

To determine the  $10^3\ln\alpha_{\text{vapor-silicate}}$ , we conducted first-principles calculations to derive the  $\langle F \rangle$  in all species and estimated the  $\langle F \rangle$  in the vapor phase based on their fractions. The vapor phase is enriched in light Se and Te isotopes relative to the silicate melt with the  $10^3\ln\alpha_{\text{vapor-silicate}}$  of  $\delta^{82/76}\text{Se}$  ranging from -0.1 to -0.2 ‰ and the  $10^3\ln\alpha_{\text{vapor-silicate}}$  of  $\delta^{128/126}\text{Te}$  from -0.031 to -0.023 ‰ at 1300-1500 K (Fig. 2c and Fig. S8). Thus, Se and Te isotopes show an opposite fractionation direction to that of S (Fig. 2d), such that evaporative loss of S, Se, and Te from molten planetesimals will enrich the residual melt with light S isotopes but heavy Se and Te isotopes.

Combining the isotope fractionation data from our study and the literature metal/silicate partition coefficients (8, 9, 29), we model the abundances and isotopic compositions of S, Se, and Te in the BSE after early planetesimal evaporation followed by late-stage core formation. The Earth's building material, which is required to estimate the initial abundances and isotopic compositions of S, Se, and Te, remains highly debated. Dauphas (2017) (20) modeled the isotopic evolution of multi-elements (including O, Ca, Ti, Cr, Ni, Mo, Ru, and Nd) of the Earth's mantle during Earth's accretion and found that the accreted material is mainly characterized by the enstatite-chondrite-like composition with only minor fractions of carbonaceous chondrite-like material. The best-fitting model of Earth's accreting material that can reproduce most of these isotopic signatures consists of 71% enstatite chondrite, 24% ordinary chondrite, and 5% carbonaceous chondrite (20). That is, the Earth's accreting materials come from both CC (carbonaceous) and NC (non-carbonaceous) reservoirs, which is also supported by Earth's nucleosynthetic K and Zn isotope anomalies (32–35). Furthermore, the mass-independent Mo and Nd isotopic composition of Earth's mantle may reflect a mixture between NC and unsampled *s*-process-enriched CC material compared to known chondrites (36–38), consistent with the Dauphas's model. In contrast, Schiller et al. (2020) found that the mass-independent  $\mu^{54}\text{Fe}$  of Earth's mantle is lower than other chondrites but only overlaps with the value of CI chondrite (39), indicating that most of Fe in Earth's mantle derived from inward-drifting CI-like material. If so, large fractions of Cr and Ni in Earth's mantle would be delivered by the same material as Fe is intermediate in metal affinity between Cr and Ni, which, however, cannot explain Earth's Ni and Cr isotopic signatures (40–43). Overall, we use the Dauphas's model in this study due to the isotopic similarity of Earth and this model for many elements.

The initial  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$ , and  $\delta^{128/126}\text{Te}$  of Earth's building material are estimated to be  $\sim -0.21$  ‰,  $-0.35$  ‰, and  $+0.02$  ‰, respectively. The initial S, Se, and Te abundances in the accreting material are  $\sim 4.5$  wt%, 20 ppm (3), and 1.8 ppm (18, 19), respectively. Our models show that approximately 85 to 95% early evaporative loss can reproduce the present-day BSE's  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$ , and  $\delta^{128/126}\text{Te}$ , as well as the S, Se, and Te abundances after core formation over the modeled  $D_{\text{S}^{\text{metal/silicate}}}$ ,  $D_{\text{Se}^{\text{metal/silicate}}}$ , and  $D_{\text{Te}^{\text{metal/silicate}}}$  ranges, without the need for a late veneer (Fig. 3).

If the evaporative loss of S, Se, and Te is greater than 95%, the BSE would have S, Se, and Te abundances lower than the present-day BSE's values after core formation, and its  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$ , and  $\delta^{128/126}\text{Te}$  would deviate from the current BSE's values. In this case, a late veneer is required to increase the S, Se, and Te abundances to the level of the current BSE, but the final  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$  and  $\delta^{128/126}\text{Te}$  of the BSE would also be affected by the late-veneer material. To evaluate the late-veneer effect on the  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$ , and  $\delta^{128/126}\text{Te}$  of the bulk silicate reservoir, we conduct Monte Carlo simulations with a late veneer characterized by carbonaceous-chondrite-like material. The models show that the  $\delta^{82/76}\text{Se}$  and  $\delta^{128/126}\text{Te}$  of the BSE can always be reproduced, no matter how much Se and Te in the present-day BSE was added by a late veneer (Fig. S9). However, the  $\delta^{34/32}\text{S}$  of the BSE will be close to the late-veneer value if the amount of S added by the late veneer is too high, and no more than  $\sim 30\%$  of the present-day BSE's S budget is allowed to be added by the late veneer to reproduce the BSE's  $\delta^{34/32}\text{S}$  (Fig. S9).

Our modeling suggests that the abundances and isotope compositions of S, Se, and Te in the BSE can be explained by protoplanetary differentiation rather than a late veneer. Even though a late veneer is allowed, the amounts of S, Se, and Te added by the late veneer should not exceed  $\sim 30\%$  of the present-day BSE's budgets. If the late veneer volatile-rich material is sourced from the outer Solar System, e.g., carbonaceous chondrite (1–4), our result constrains that the mass of

late veneer cannot exceed 0.2% of the mass of Earth's mantle. This allows an estimate of the maximum budgets of other volatile elements delivered along with chalcogens by the late veneer. Using the elemental abundances in carbonaceous chondrite (3, 44–46) and the BSE (47), we find that the late veneer contributes at most 5% of the H and 40% of the C in the current BSE (45, 47–49). **The limited amount of water from a late veneer indicates that Earth may have obtained most of its water during early accretion and differentiation (50) from its major building block – enstatite-chondrite-like material (20) and/or through the interaction between primordial hydrogen atmosphere and proto-Earth. Also, the allowed mass of late veneer would not significantly contribute to the budgets of moderately volatile elements in the BSE, which is supported by the mass-dependent Zn and Cu isotopic signatures of Earth's mantle (51, 52).** In contrast, the maximum mass of late veneer constrained from S, Se, and Te isotopes can supply 100% of the N in the current BSE ( $2.2 \pm 1.2$  ppm, (45, 47, 53)). The pre-late veneer BSE would have a higher C/N ratio than the present-day BSE, which is also supported by metal-silicate partition coefficients (14, 54). In addition, at most 30% of the HSEs in the current BSE were added by a late veneer. This is consistent with recent experiments (55, 56) yielding low metal-silicate partition coefficients for Pd and Pt and hinting for possible higher concentrations of platinum-group elements in the mantle after core formation. Overall, our results show that protoplanetary differentiation may play a dominant role in establishing Earth's most volatile element inventory, and the late veneer would have a limited contribution to the BSE's volatiles, with the exception of N.

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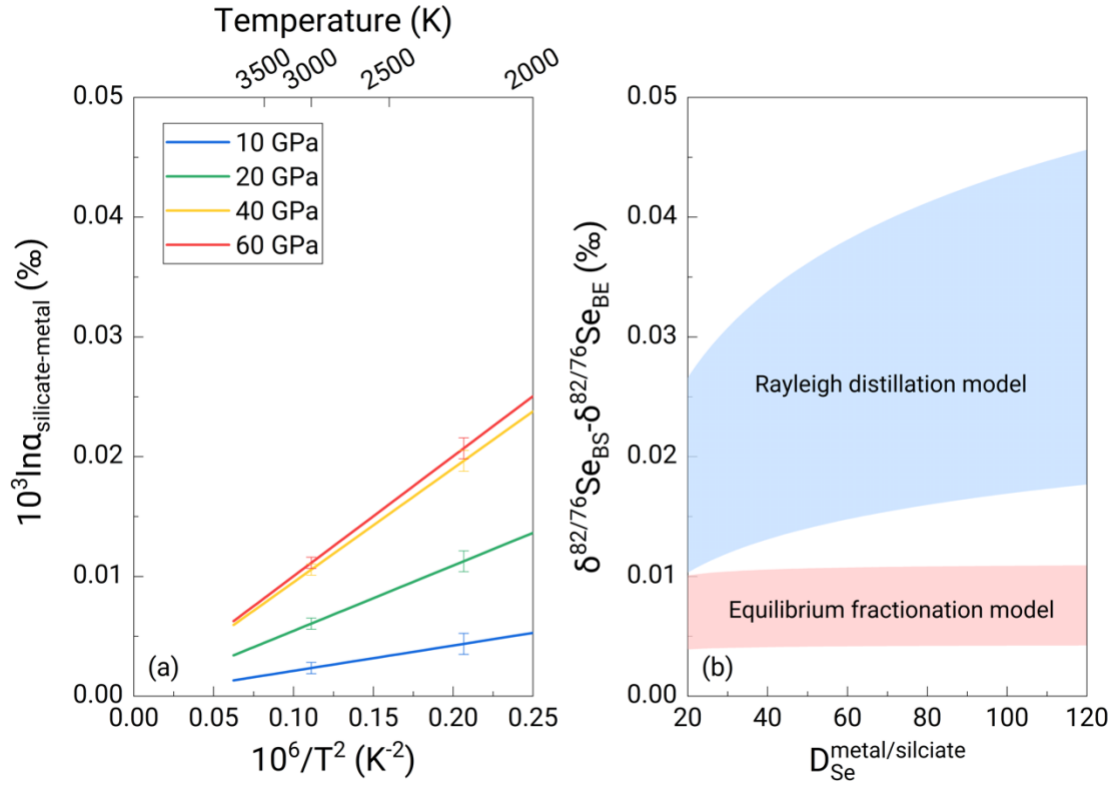
## Supplementary Materials

Materials and Methods

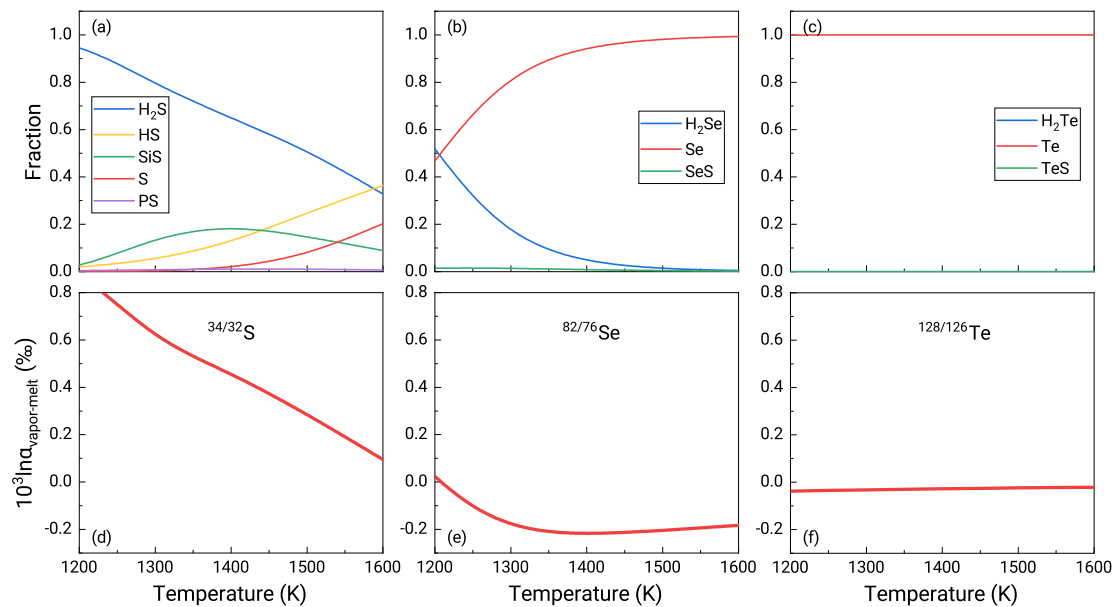
Supplementary Text

Figs. S1 to S9

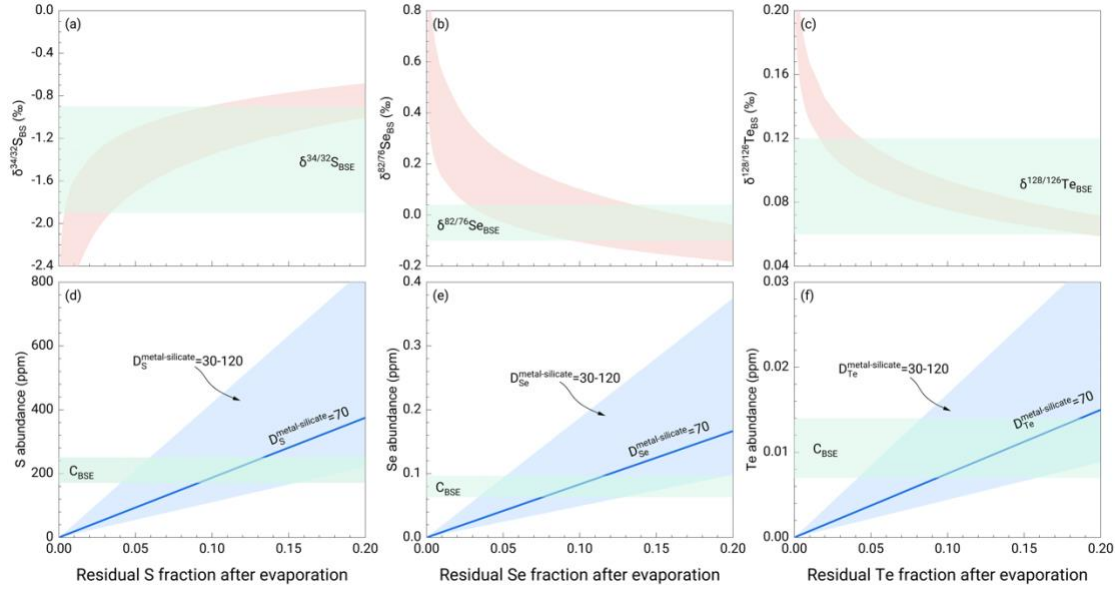
References (41–53)



**Fig. 1. Selenium isotope fractionation during core formation.** (a) Equilibrium Se isotope fractionation factors ( $10^3 \ln \alpha$  of  $^{82/76}\text{Se}$ ) between silicate and metallic melts. (b) Modelled  $\delta^{82/76}\text{Se}$  difference between the bulk silicate part ( $\delta^{82/76}\text{Se}_{\text{BS}}$ ) and the bulk Earth ( $\delta^{82/76}\text{Se}_{\text{BE}}$ ) caused by mantle-core differentiation. The Se partition coefficient ( $D_{\text{Se}}^{\text{metal/silicate}}$ ) determines the remaining Se fraction in the silicate part after core formation. Equilibrium and Rayleigh distillation models are considered as two endmember models.



**Fig. 2. Sulfur, selenium, and tellurium species in the vapor phase and the isotope fractionation between vapor and melt.** (a) (b) (c) the fractions of S, Se, and Te species in the vapor evaporated from molten planetesimals with the presence of nebular  $\text{H}_2$ . At 1300-1500 K, Se and Te predominantly occur as  $\text{Se}(\text{g})$  and  $\text{Te}(\text{g})$ , respectively, while S mainly occurs as  $\text{H}_2\text{S}(\text{g})$ . (d) (e) (f) the equilibrium S, Se, and Te isotope fractionation factors between vapor and melt. The residual melt is enriched in light S isotopes but heavy Se and Te isotopes relative to the vapor phase due to the different species of S, Se, and Te in the vapor phase.



**Fig. 3. Abundances and isotopic compositions of S, Se, and Te in the bulk silicate part established by planetesimal evaporation followed by core formation.** The green shades refer to the present-day BSE values (3, 15–17, 19). **The Se and Te isotopic compositions of the BSE were inferred from mantle peridotites (17, 19). The previous estimate of BSE's  $\delta^{82/76}\text{Se}$  from Mid-Ocean-Ridge Basalts (10) is consistent with the value from peridotites.** (a)  $\delta^{34/32}\text{S}_{\text{BS}}$  (b)  $\delta^{82/76}\text{Se}_{\text{BS}}$  and (c)  $\delta^{128/126}\text{Te}_{\text{BS}}$  as a function of the residual fraction after evaporation, as core formation does not induce resolvable S, Se, and Te isotope fractionation. The initial  $\delta^{34/32}\text{S}$ ,  $\delta^{82/76}\text{Se}$ , and  $\delta^{128/126}\text{Te}$  of the building material (71% enstatite chondrite + 24% ordinary chondrite + 5% carbonaceous chondrite) (20) are estimated to be  $\sim -0.21$  ‰,  $-0.35$  ‰, and  $+0.02$  ‰ (Fig. S1), respectively. (d) S (e) Se (f) Te abundances in the bulk silicate part after evaporation followed by core-mantle differentiation. The initial S, Se, and Te abundances in the build material are set as  $\sim 4.5$  wt%, 20 ppm (3), and 1.8 ppm (18, 19), respectively. The blue shades represent the modeled results with  $D_{\text{S}}^{\text{metal/silicate}}$ ,  $D_{\text{Se}}^{\text{metal/silicate}}$ , and  $D_{\text{Te}}^{\text{metal/silicate}}$  ranging from 30 to 120 (8, 9, 29), in which the results at  $D_{\text{S}}^{\text{metal/silicate}}$ ,  $D_{\text{Se}}^{\text{metal/silicate}}$ , and  $D_{\text{Te}}^{\text{metal/silicate}} \sim 70$  are highlighted by the blue lines.