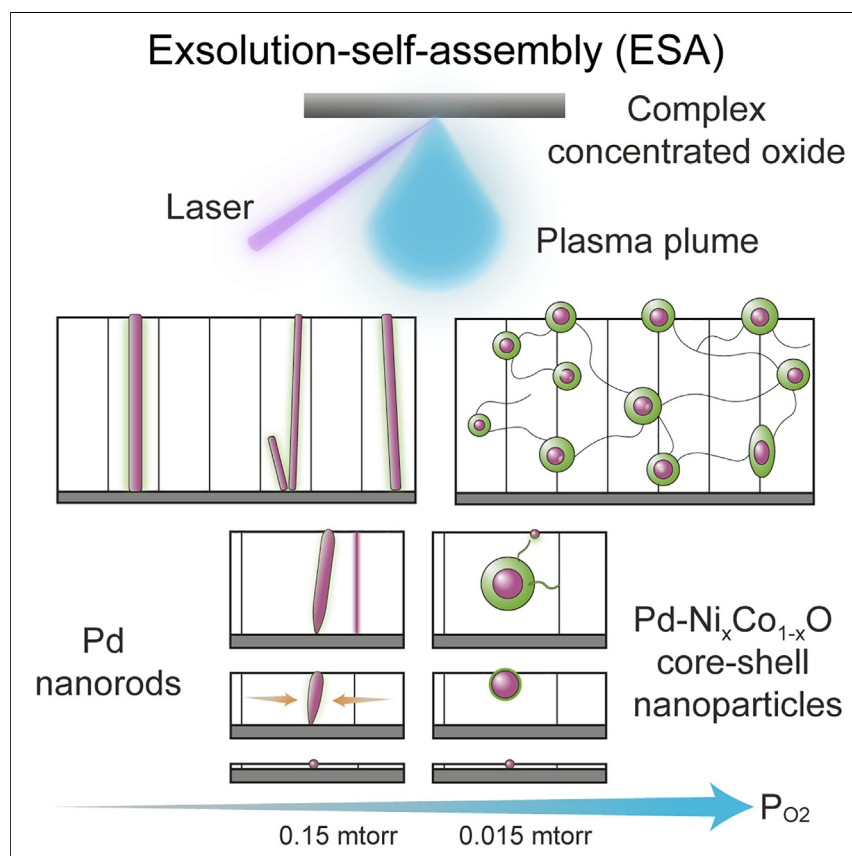


## Article

## Designing nanostructure exsolution-self-assembly in a complex concentrated oxide



A new exsolution-self-assembly (ESA) approach was demonstrated during thin film growth, combining aspects of immiscibility-driven one-step self-assembly with defect-interaction-driven two-step exsolution to yield novel vertically aligned nanostructures and particle-in-matrix nanocomposites. The one-step ESA synthesis is applied to a model complex concentrated oxide, an emerging class of materials including high-entropy and entropy-stabilized oxides, whose ability to stabilize five or more cations in a single crystal structure provides a combinatorial space for new materials with broad electrochemical and information storage applications.

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**Highlights**

Tunable nanostructures in a complex concentrated oxide by exsolution-self-assembly

An approach to direct the evolution of nanoparticles and nanorods via oxygen pressure

Cation reducibility predicts nanostructure chemical composition, tailoring growth

Mass transport through extended defects governs local morphology

**Benchmark**

First qualification/assessment of material properties and/or performance

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Article

# Designing nanostructure exsolution-self-assembly in a complex concentrated oxide

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## SUMMARY

Complex concentrated oxides (CCOs) are an emerging material class that includes high-entropy and entropy-stabilized oxides whose properties stem from disorder-induced electronic structure and chemistry caused by stabilizing solid solutions of >5 cations. Integrating CCOs into composites will expand material design beyond the single-phase paradigm. We demonstrate tunable CCO-derived nanostructures by a simple method: exsolution-self-assembly (ESA), a one-step approach to direct the evolution of nanoparticles and nanorods in nanocomposites. We have developed a fundamental understanding of driving forces and formation mechanisms in CCOs using atomic-scale probes, which reveal that ESA can be directed using Ellingham's model of cation reducibility in a model CCO perovskite  $\text{LaFe}_{0.7}\text{Ni}_{0.1}\text{Co}_{0.1}\text{Pd}_{0.05}\text{Ru}_{0.05}\text{O}_{3-\delta}$ . This approach enables tailored growth of multielement nanorod and nanoparticle composite structures whose formation is correlated with electronic conductivity that exceeds 0.1 S/cm at room temperature. Given the vast combinatorial space of CCOs, ESA is expected to be highly extensible via the integration of various compositions and crystal structures.

## INTRODUCTION

Recently, unprecedented materials have been realized by stabilizing a relatively large number (>5) of cations in solid solution, wherein phase stability is enhanced or provided by high configurational entropy associated with compositional complexity.<sup>1</sup> This emerging class of complex concentrated oxides (CCOs)—including high-entropy oxides (HEOs) with elevated configurational entropy and entropy-stabilized oxides (ESOs), which are unequivocally stabilized by entropy<sup>2</sup>—provides a novel tuning route beyond conventional chemical composition to improve material properties by controlling the distributions of lattice disorder, bond energy, defect formation, and cation multivalency.<sup>3,4</sup> Although research on CCOs/HEOs/ESOs has focused mainly on single-phase materials, our very recent work shows that ESO phase separation into a composite provides novel control over electrical transport mechanisms by directing ionic and electronic current through the different composite phases, highlighting the potential of composite micro- and nanostructures<sup>5,6</sup> to finely tailor the properties of these emerging materials and broaden their potential applications.

Exsolution<sup>7,8</sup> is a multistep synthesis understood by considering atomistic interactions between charged point defects. A reducible oxide is exposed to a reducing

## PROGRESS AND POTENTIAL

Complex concentrated oxides (CCOs) include high-entropy and entropy-stabilized oxides, whose ability to stabilize five or more cations in a single-crystal structure provides a vast design space for new materials. In a model CCO, we demonstrate exsolution-self-assembly (ESA), a highly extensible synthesis for diverse and exotic nanocomposites. ESA combines the benefits of reducibility-controlled metal cation exsolution with the spontaneity of immiscibility-driven self-assembly to produce intricate tailorable nanostructures in a single synthetic step. Application of ESA to CCOs and advanced entropy-designed chemical compositions and crystal structures will yield multidimensional and multielement nanostructures for next-generation electrochemical, electronic, and information storage devices. The synthesis has the potential to create highly tunable nanocomposites that combine the properties and functions of emerging high-entropy materials with precisely designed embedded nanostructures.

atmosphere, forming oxygen vacancies that electrostatically promote the accumulation of metal cations into nanoparticles anchored (i.e., partially submerged) in the surface of a material<sup>9</sup> and/or embedded (submerged) within the material.<sup>10</sup> Exsolution offers tunable and stable nanostructures, flexible material selection, and structure modulation of the oxide matrix,<sup>11–14</sup> but it requires relatively long treatment time at high temperature under reducing gases. Unlike exsolution, nanocomposite self-assembly occurs spontaneously in one step during thin film deposition by physical vapor deposition following one of two mechanisms: diffusion, nucleation, and growth for materials systems with limited miscibility, and (pseudo)spinodal decomposition for miscible materials.<sup>15</sup> Self-assembled nanostructures have final shape, size, spatial distribution, and crystallographic orientation relative to the matrix governed by free energy minimization.<sup>5</sup>

Here, we present the exsolution-self-assembly (ESA) method, whereby nanostructures form spontaneously from a miscible single-phase solid solution during film growth due to both immiscibility of the new nanophase in the matrix and atomistic interactions between point defects. We demonstrate tunable nanostructure growth in a CCO by ESA, which should provide improved synthetic control of nanoparticles and nanorods across a broader phase space than is currently possible.<sup>16</sup> We present a case study demonstrating simple access of different nanostructures via control of oxygen partial pressure ( $P_{O_2}$ ) during pulsed laser deposition (PLD, Figure 1A). ESA using a CCO precursor enables novel multielement composite morphologies with finely tuned subphase chemical composition that are not possible by conventional exsolution or self-assembly independently. Moreover, the demonstrated fabrication method is shown to work on a variety of substrate materials with application potential in industrial semiconductors and electrochemical devices. By mapping the distribution of distinct subphase structure and chemistry with nanometer spatial resolution, we show that Ellingham's model of cation reducibility predicts nanostructure chemical composition, whereas mass transport via extended defects governs local morphology formation.<sup>17,18</sup> This is demonstrated in  $\text{LaFe}_{0.7}\text{Co}_{0.1}\text{Ni}_{0.1}\text{Ru}_{0.05}\text{Pd}_{0.05}\text{O}_{3-\delta}$ , a model CCO inspired by the  $\text{LaFeO}_3$  perovskite, which has promising applications in solid oxide fuel/electrolysis cell electrodes, hydride batteries, supercapacitors,<sup>19–21</sup> gas sensors, electro-mechanical devices, (photo)electrocatalysis, and photovoltaics due to its abundance, nontoxicity, chemical stability, and direct band gap.<sup>22,23</sup> However, a key drawback of some  $\text{LaFeO}_3$ -based materials is sluggish electronic conductivity ( $<0.1$  S/cm), weak sensing response, low catalytic selectivity, and poor charge separation that seriously restricts performance.<sup>22,24</sup> More important, the electronic conductivity of our nanocomposites exceeds 0.1 S/cm at room temperature, with transport correlated to nanostructure formation primarily through B-site cation (i.e., Ni and Co) mixed valency and site occupancy in the CCO matrix. Given the vast combinatorial space of CCOs, ESA is expected to be highly extensible through application to novel compositions and crystal structures, with the understanding presented here enabling one to take advantage of chemical complexity in a rational way.

## RESULTS

To demonstrate the potential of ESA to access diverse nanocomposite microstructure with simple process variations, we varied the  $P_{O_2}$  during PLD to precisely tune the  $V_O^{\bullet}$  concentration, which in turn controls cation exsolution and matrix decomposition (Figure 1M). At low  $P_{O_2}$ ,  $V_O^{\bullet}$  formation in the perovskite oxide destabilizes B-site cations, inducing exsolution and matrix decomposition. The orthorhombic/pseudocubic (121)/(110)<sub>c</sub> X-ray diffraction (XRD) peak of the CCO matrix shifts

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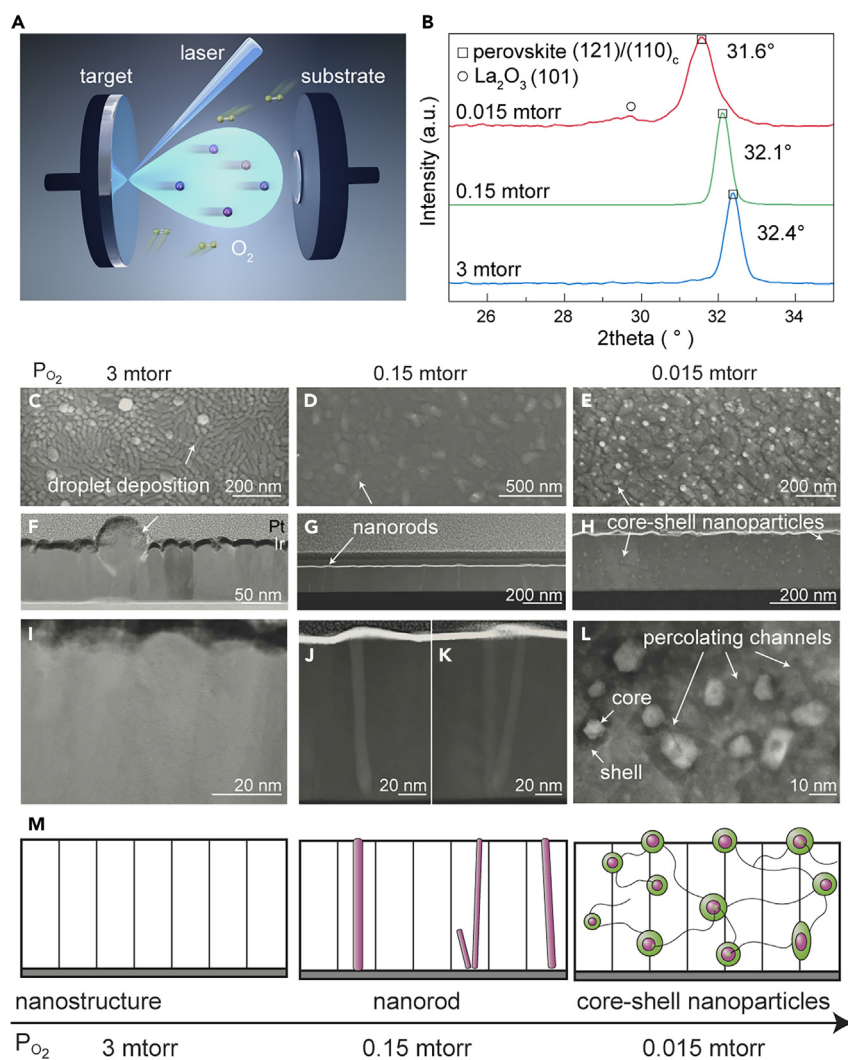
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**Figure 1. Tuning the morphology of exsolution-self-assembled nanostructures via  $P_{O_2}$**

(A) Depositing  $\text{LaFe}_{0.7}\text{Co}_{0.1}\text{Ni}_{0.1}\text{Ru}_{0.05}\text{Pd}_{0.05}\text{O}_{3-\delta}$  thin films with PLD.

(B) XRD patterns of the perovskite thin films deposited at 3, 0.15, and 0.015 mTorr. Matrix diffraction peaks shift left with decreasing  $P_{O_2}$ , indicating that higher  $V_O^{\bullet}$  content expands the unit cell.

(C–E) SEM secondary electron (SE) images of the films deposited at 3 (C), 0.15 (D), and 0.015 mTorr (E), showing top view of self-assembled nanophases.

(F and I) Low- (F) and high-magnification (I) STEM bright-field images of the film deposited at 3 mTorr, with no self-assembled phase observed.

(G, J, and K) Low- (G) and high-magnification (J and K) STEM HAADF images of the film deposited at 0.15 mTorr showing the cross-section of self-assembled nanophases. Self-assembled nanorods grow from the bottom of the thin film to the top surface.

(H and L) Low- (H) and high-magnification (L) STEM HAADF images of the film deposited at 0.015 mTorr. Self-assembled core-shell nanoparticles joined with percolating channels both embedded inside the perovskite matrix and exposed on thin film surface.

from 32.4° to 31.6° with decreasing  $P_{O_2}$  from 3 to 0.015 mTorr (Figures 1B, S1A, and S1B; Note S1), indicating unit cell expansion due to lower  $V_O^{\bullet}$  formation energy in oxygen-deficient environments.<sup>25,26</sup>  $\text{La}_2\text{O}_3$  (101) secondary phase peaks appear after deposition at 0.015 mTorr, illustrating that the perovskite structure collapses locally. Elevated  $V_O^{\bullet}$  concentration with declining  $P_{O_2}$  was corroborated by tracking oxygen species by O1s X-ray photoelectron spectroscopy (XPS) in

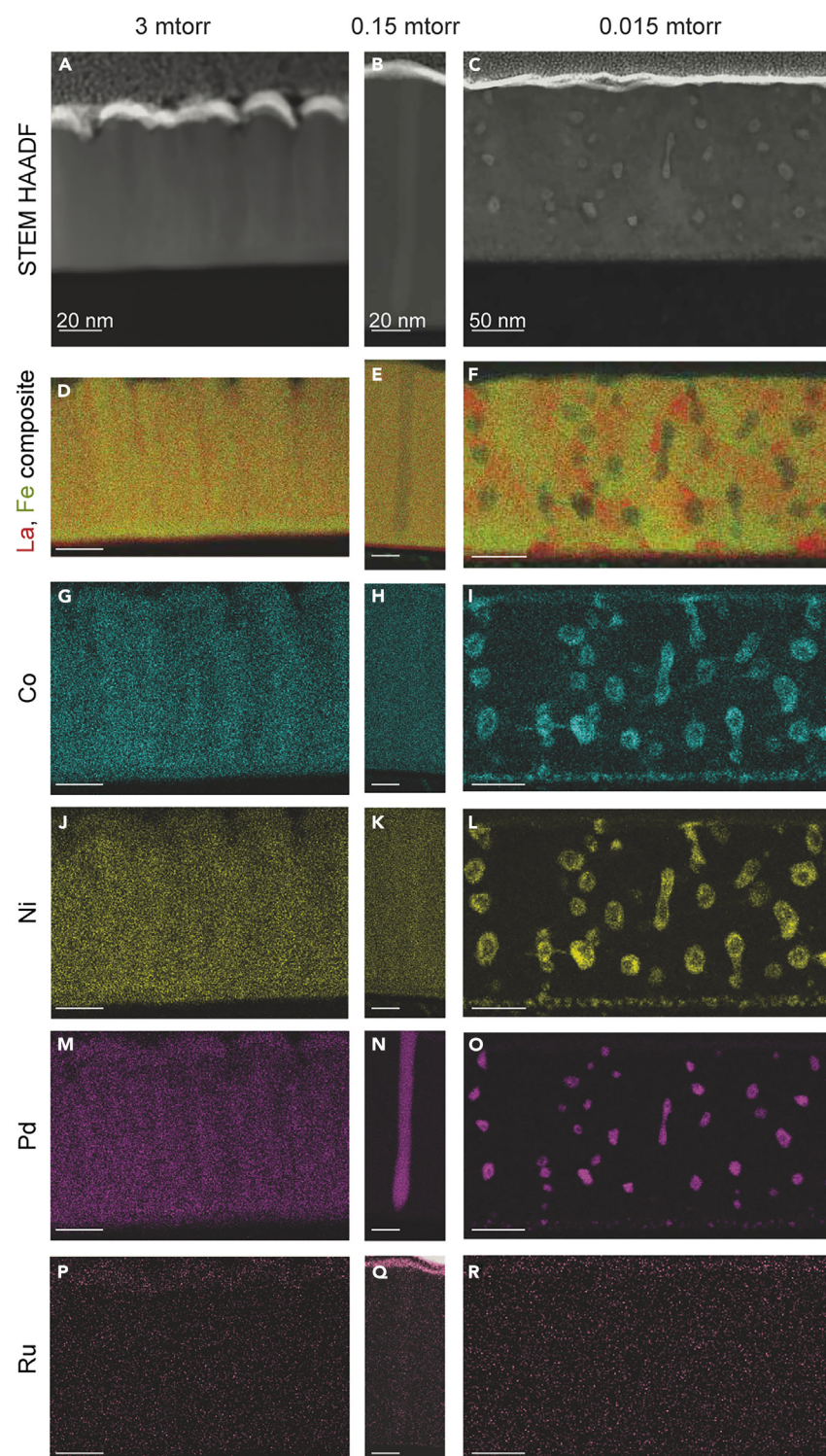


the film's top-most  $\sim 10$  nm (Figures 6A and S2; Table S1; Note S2). When  $P_{O_2} = 3$  mTorr, no self-assembled phases are observed (Figures 1C, 1F, and 1I). When  $P_{O_2} = 0.15$  mTorr, nanorods form in the matrix (Figures 1G, 1J, and 1K), with an average width of 10 nm. Most grow from thin film bottom to the top surface continuously, with an average length of 120 nm. White dots in the scanning electron microscopy (SEM) top-view image (Figure 1D) are nanorod tips. When  $P_{O_2} = 0.015$  mTorr, abundant and well-dispersed core-shell nanoparticles form, with sizes of 7–26 nm (average 19 nm; Figures S27 and S28; Note S16), high population density ( $332 \mu\text{m}^{-2}$ ), and interconnected by percolating channels (Figures 1E, 1H, and 1L) similar to those reported in studies of two-step Fe exsolution.<sup>11,27</sup> In addition, similar phase composition and crystallographic textures are observed in films grown on amorphous and crystalline oxides and metals (Figures S1A–S1C; Note S1), suggesting the general applicability of this method to different device architectures.

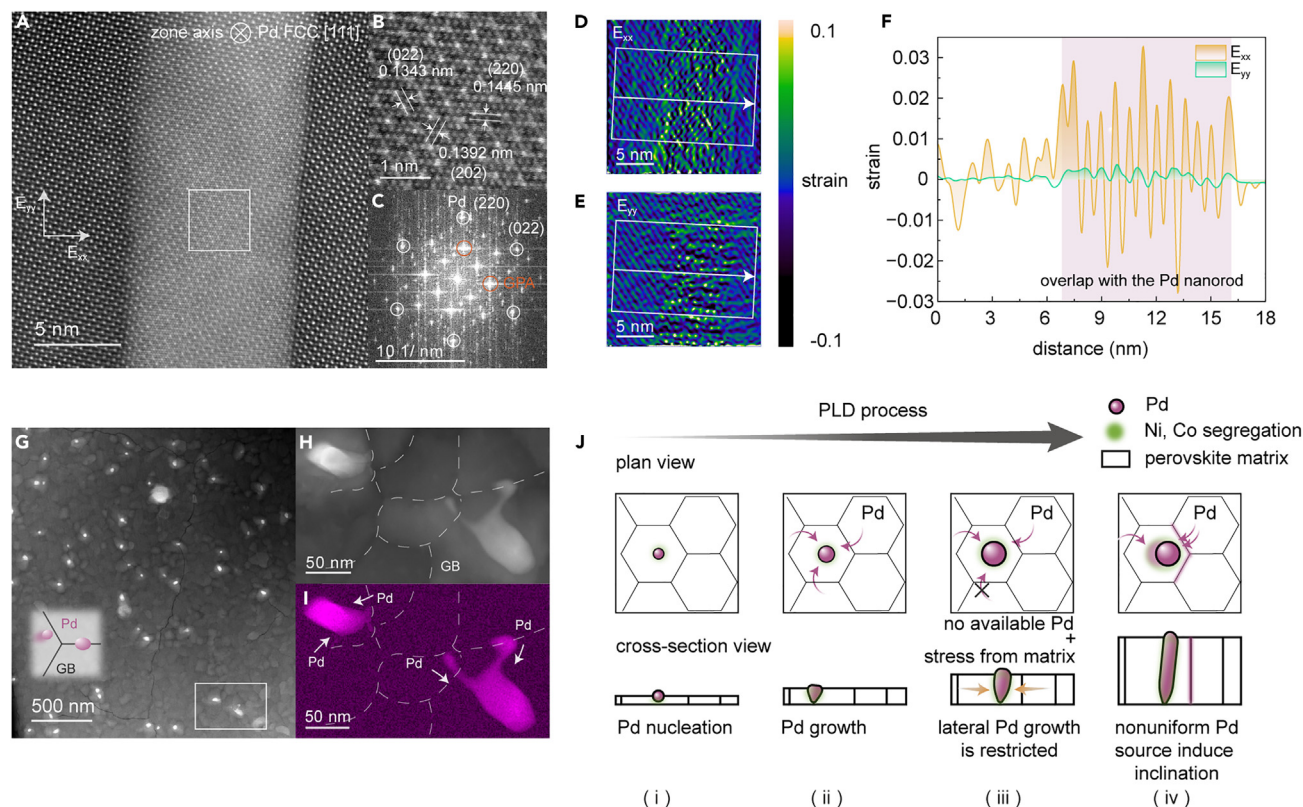
The chemical heterogeneity associated with the composite microstructures was characterized using scanning transmission electron microscopy energy-dispersive X-ray spectroscopy (STEM EDS) and electron energy loss spectroscopy (EELS) (Figures 2 and S3; Note S3). At  $P_{O_2} = 3$  mTorr, uniform elemental distribution (Figures 2A, 2D, 2G, 2J, 2M, 2P, S3A, and S3B) demonstrates that B-site cations remain stable in matrix, except Ru. At  $P_{O_2} = 0.15$  mTorr, nanorods are Pd with subtle Ni and Co segregation at matrix–rod interfaces (Figures 2B, 2E, 2H, 2K, 2N, 2Q, S3C–S3F, and S7). At  $P_{O_2} = 0.015$  mTorr, core-shell nanoparticles comprise Pd cores and  $\text{Ni}_x\text{Co}_{1-x}\text{O}$  shells (Figures 2C, 2F, 2I, 2L, 2O, 2R, S3G, and S3H); La and Fe show nonuniform distribution in the surrounding matrix (Figure 2F), consistent with XRD detection of  $\text{La}_2\text{O}_3$  (Figure 1B). All of the films show Ru surface segregation (Figures 2P and 2R), which atom probe tomography (APT) confirms (Figure S4). Surface accumulation of Ru at 3 mTorr reflects its greater tendency toward reducibility compared to Pd, in line with its weak metal–support interaction compared with Pd in  $\text{LaFeO}_3$ <sup>28,29</sup> (Note S4).

### Nanophase identification, composition, and formation mechanism

To confirm that the Pd-containing nanorods were Pd metal and to understand the nanorod growth mechanism, atomic-resolution STEM high-angle annular dark field (HAADF) imaging and geometric phase analysis (GPA) were used to identify the nanorod crystal structure and map local strain in the matrix (Figures 3, S5, and S6). When viewing the perovskite matrix along the perovskite  $[100]_c$  zone axis, the Pd metal nanorod lattice was directly visible in the  $[111]$  zone-axis orientation (Figures 3A–3C). To decipher the formation mechanism of the Pd nanorods, GPA was applied to Figure 3A to map the perovskite strain near the Pd nanorod (Figures 3D and 3E). The tensile in-plane lateral strain of the matrix is 10 times larger than the tensile out-of-plane vertical strain at the interface with the Pd nanorod (Figure 3F). The in-plane Pd lattice constant is compressed by 3.4% compared with standard Pd (Figure S6; Table S2), indicating that nanorod lateral growth is restricted by the matrix. The compressive stress contributes to the balance of forces/energies that defines the diameter of the Pd nanorod and thus the amount of laterally diffusing Pd that can be accommodated in each growth layer.<sup>15</sup> In addition, Pd nanorods grow inclined at/near grain boundaries (GBs) with Pd segregation (Figure 3G). It is hypothesized that Pd sourced from GBs creates nanorods inclined toward GBs (Figures 3H and 3I). Therefore, the growth of Pd is restricted by both the elastic energy and the availability of Pd in surrounding GBs (Figure 3J). Ni and Co segregation at the matrix–Pd nanorod interface, observed by STEM EDS (Figures 2H, 2K, and S7) may result from local tensile strain of the perovskite, which lowers the  $V_O^\bullet$  formation



**Figure 2. Elemental mapping of exsolution-self-assembled nanostructures tuned by  $P_{O_2}$**   
(A–C) STEM HAADF images of the CCO-derived nanocomposite thin films.  
(D–N) STEM EDS mapping of (D–F) La and Fe, (G–I) Co, (J–L) Ni, and (M and N) Pd.  
(P–R) STEM EELS mapping of Ru.



**Figure 3. Phase identification and formation mechanism of exsolution-self-assembled Pd nanorods**

(A–C) Cross-section STEM HAADF image of a Pd nanorod embedded in the perovskite matrix (A). Local magnified image of area marked in (A), showing Pd atom arrangement and Pd {220} plane spacing (B), and FFT pattern of (A) showing the Pd nanorod along the [111] zone axis and the matrix along the [100]<sub>c</sub> zone axis (C).

(D and E) Strain mapping of perovskite matrix in (A) along lateral in-plane direction ( $E_{xx}$ ) (D) and vertical out-of-plane direction ( $E_{yy}$ ) (E) via GPA. Directions of  $E_{xx}$  and  $E_{yy}$  are labeled in (A).

(F) Strain profile of the interface matrix with Pd nanorod along the direction marked in (D and E).  $E_{xx}$  shows tensile strain approximately 10 times higher than  $E_{yy}$  at the interface, indicating that compressive stress from perovskite matrix restricts growth of Pd nanorod in lateral direction.

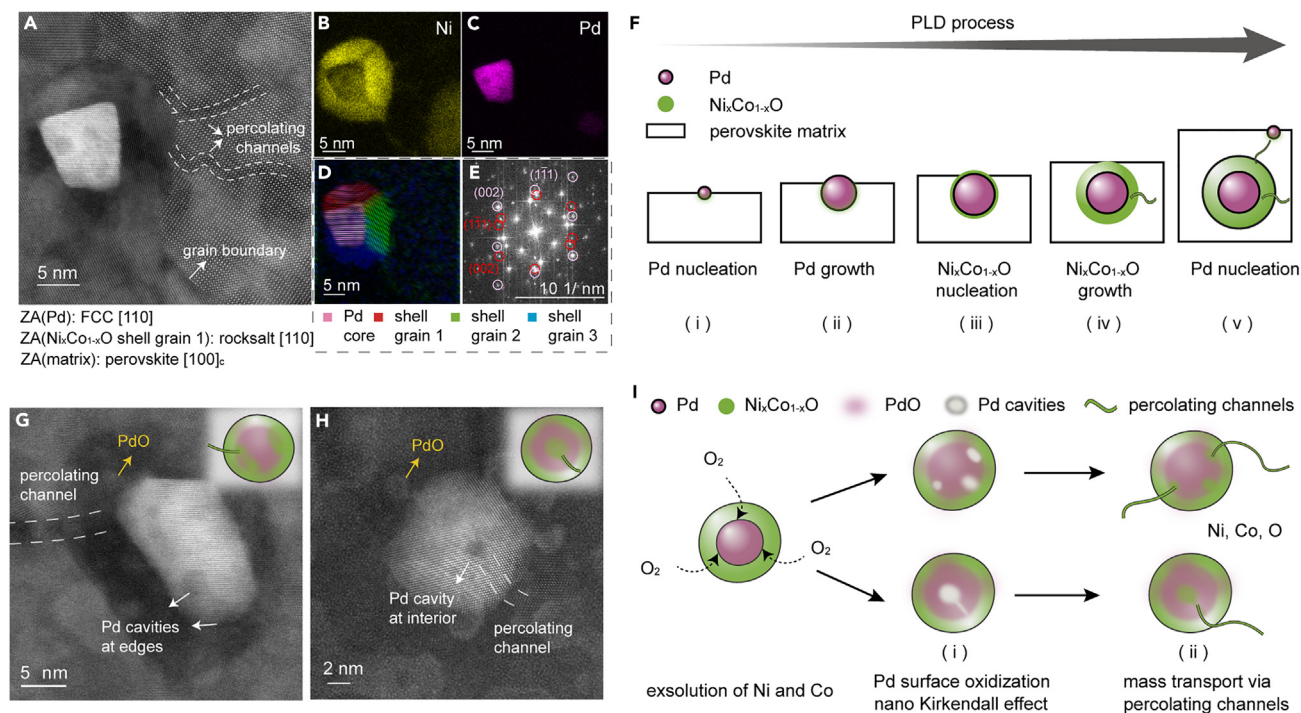
(G–I) Plan-view STEM HAADF image of Pd nanorods embedded in the perovskite matrix (G), (H) is the local magnified image marked in (G), and corresponding EDS mapping of Pd (I). Plan-view observation shows the Pd nanorods are located at or close to GBs, and Pd nanorods show inclined growth as shown as the schematic inset in (G).

(J) Schematic illustration of mechanism of Pd nanorod growth. Pd nucleates at the beginning of PLD (i), then surrounding Pd diffuse to the nucleus to grow laterally and vertically (ii). The lateral growth of Pd is restricted by compressive stress from matrix and limited availability of Pd from the surroundings, forming Pd nanorods (iii). Nonuniform Pd sourcing from grains and GBs leads to inclined growth of Pd nanorods (iv).

energy<sup>30</sup> to reduce  $\text{Ni}^{3+}$  and  $\text{Co}^{3+}$  and accommodates the larger-radii  $\text{Ni}^{2+}$  and  $\text{Co}^{2+}$  species.

In the core-shell composites, self-assembled nanostructures were confirmed to be metal-oxide  $\text{Pd-Ni}_x\text{Co}_{1-x}\text{O}$  core-shells using crystal structure and composition mapping by fast Fourier transform (FFT) analysis (Figures S9–S11), STEM EDS (Figures 4B, 4C, S9, and S10), and APT (Figures 5 and S26; Note S15). Representative core-shell nanoparticles and percolating channels connecting them are shown in Figure 4A and Note S8. Crystal structures of the perovskite matrix, nanoparticle core and rock salt shell, and  $\text{La}_2\text{O}_3$  were resolved based on inverse FFT (IFFT) patterns (Figure S9F). FFT patterns of the core show FCC {111} and {200} with lattice parameters of Pd metal (Figure S9I). FFT patterns of the rock salt shell (Figure S9G) shows a lattice parameter of 4.192 Å, which is between NiO (4.177 Å) and CoO (4.261 Å).<sup>31,32</sup> Elemental mapping (Figures 4B, 4C, and S8B–S8D) is consistent with the FFT phase



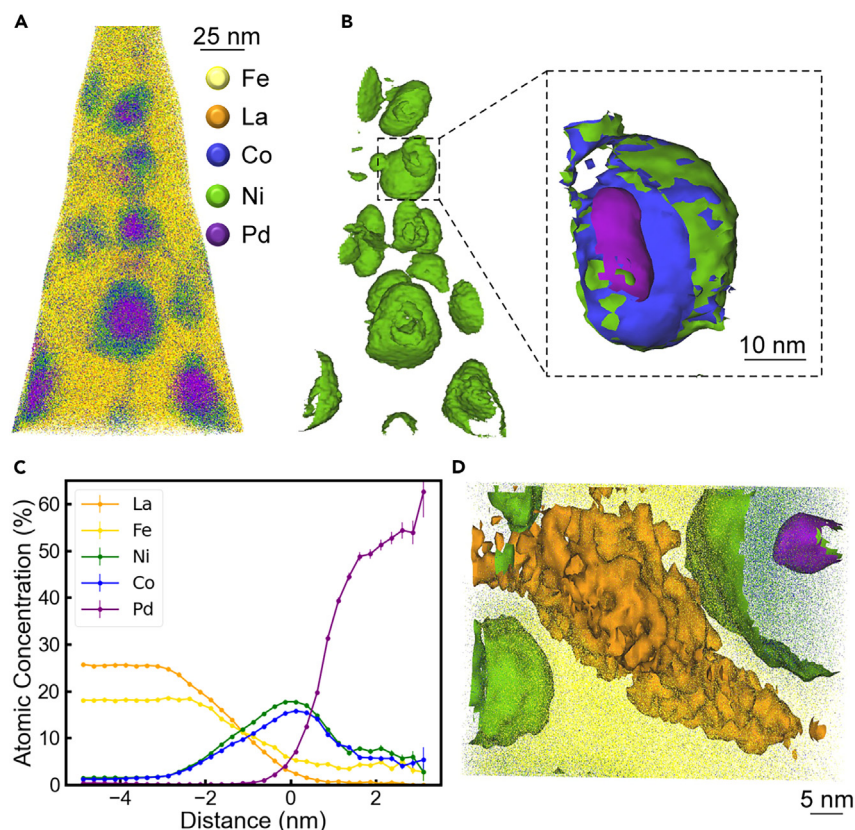


**Figure 4. Phase identification and formation mechanism of core-shell Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O nanoparticles**

(A) STEM HAADF image of a Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticle embedded in a GB of the perovskite matrix, where the matrix is along perovskite [100]<sub>c</sub> zone axis (ZA) and the Pd core is along FCC [110] ZA. (B and C) Corresponding STEM EDS mapping of Ni (B) and Pd (C). (D) IFFT composite of Pd core (pink) and rock salt Ni<sub>x</sub>Co<sub>1-x</sub>O shell grains in (A). The shell is composed of 3 grains marked as shell grain 1 (red), 2 (green), and 3 (blue). The shell grains nucleate and grow coherently at different facets of the Pd core. (E) Composite FFT patterns of shell grain 1 (red) and Pd core (pink), showing Ni<sub>x</sub>Co<sub>1-x</sub>O shell epitaxially grows along Pd surface. (F) Formation mechanism of Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticles: Pd nucleation (i) and growth (ii); Pd seeds exsolution of Ni and Co and nucleation of Ni<sub>x</sub>Co<sub>1-x</sub>O at surface of Pd core (iii) and Ni<sub>x</sub>Co<sub>1-x</sub>O grow to form a shell by mass transport by percolating channels (iv). A new Pd nucleus forms on the surface, initiating another core-shell particle (v). (G and H) STEM HAADF images of self-assembled Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticles showing cavities at the edge (G) and interior of Pd cores (H). These inverse-core-shell structures are filled with Ni<sub>x</sub>Co<sub>1-x</sub>O and connected by percolating channels, as shown in the schematic inset. (I) Mechanism of inverse core-shell structure: Pd cavities arise from nano-Kirkendall effect induced by local surface oxidation of Pd during exsolution of Ni and Co. Ni<sub>x</sub>Co<sub>1-x</sub>O grows inside Pd cavities via diffusion along percolating channels.

analysis, confirming that self-assembled core-shell nanoparticles consist of a Pd metal core with a rock salt Ni<sub>x</sub>Co<sub>1-x</sub>O shell. The shell Ni:Co ratio ranges from 1.5 to 1.9 (1.7 average) in 7 core-shell nanoparticles, indicating that the shell is a Ni-rich rock salt Ni<sub>0.63</sub>Co<sub>0.37</sub>O (Figure S26; Note S15).

The 3D core-shell nanoparticle reconstruction and chemical composition analysis by APT confirm that Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanostructures are interconnected with percolating channels (Figures 5A–5C and S8; Videos S1, S2, S3, S4, S5, S6, and S7). Both APT tomography (Figures 5A–5C and S8) and STEM EDS mapping (Figure S13) illustrate that percolating channels are composed of Co, Ni, and O. Although percolating channels show similar chemical composition with Ni<sub>x</sub>Co<sub>1-x</sub>O shells, the atomic structure was not resolved. Based on prior perovskite reduction studies yielding Fe exsolution,<sup>11,27</sup> Co-rich CoO<sub>4</sub> ordered sublayers<sup>33</sup> and Ruddlesden-Popper domain boundaries,<sup>27</sup> we infer that percolating channels are Ni- and Co-rich transition sublayers<sup>33</sup> and serve as mass transport pathways from matrix to Ni<sub>x</sub>Co<sub>1-x</sub>O shell.<sup>33</sup> APT 3D reconstruction (Figure 5D) shows La<sub>2</sub>O<sub>3</sub> next to the nanoparticle, which is consistent with FFT analysis (Figures S9F, S9J, and



**Figure 5. APT reconstruction of exsolution-self-assembled Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticles interconnected by percolating channels**

(A) Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticles in Fe-rich perovskite matrix. Ni- and Co-rich channels are visible as blue-green filaments connecting the nanoparticles.

(B) Ni isosurface of Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shell nanoparticles in (A) with the matrix removed to better reveal their 3D structure. An isolated core-shell particle with Ni (green), Co (blue), and Pd (purple) isosurfaces is shown in the expanded view.

(C) Atomic concentrations of relevant elements along the surface normal of a Pd isosurface. The Pd isosurface at zero distance is shown inset in (B).

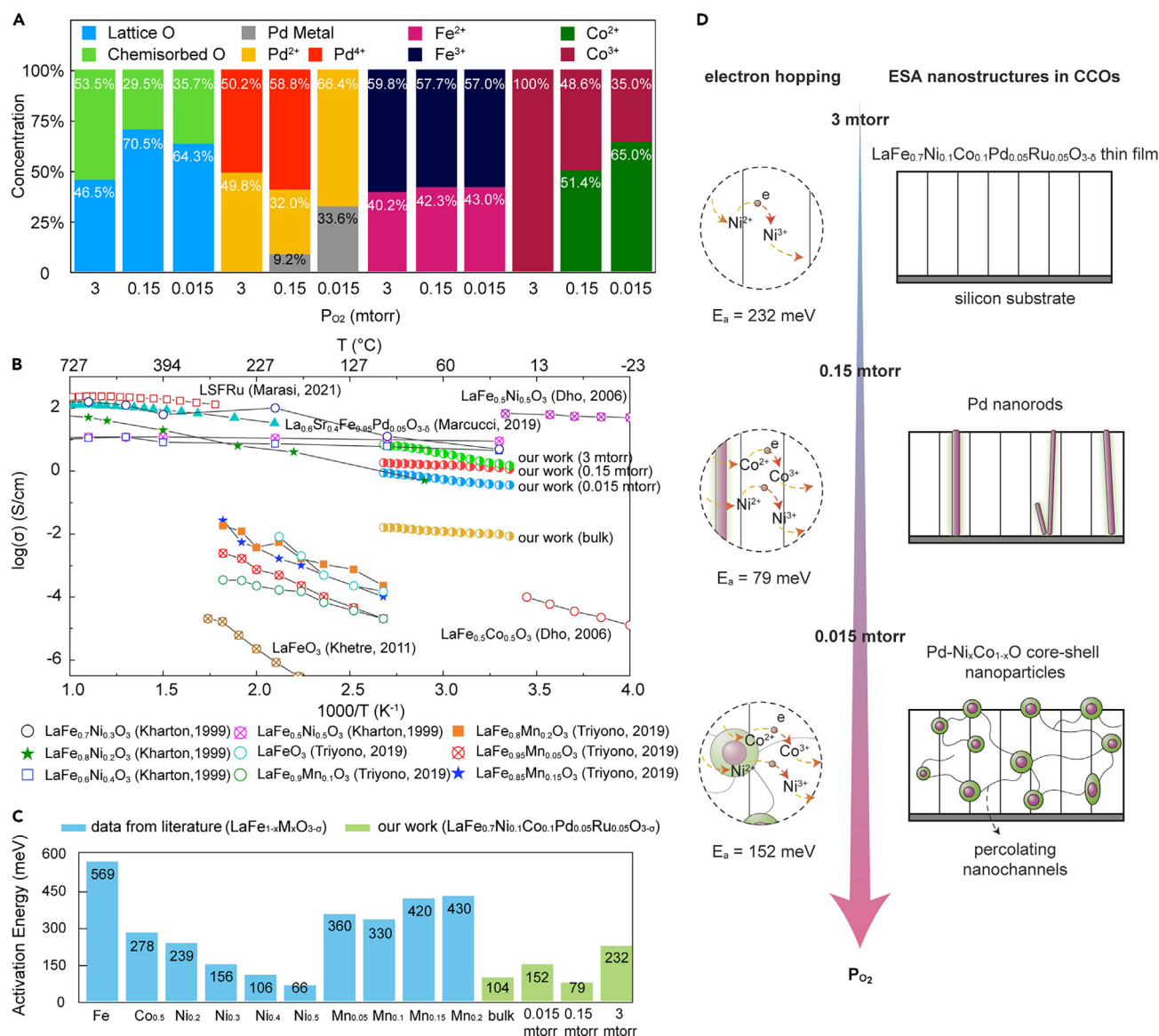
(D) APT 3D reconstruction of secondary phase La<sub>2</sub>O<sub>3</sub> showing its association with the core-shell nanoparticle.

S12A–S12I; Note S7) and STEM EDS mapping (Figures 2F, S9B, S12J, and S12K). Compared with the film deposited at 0.15 mTorr, in which no La<sub>2</sub>O<sub>3</sub> phase is observed, La<sub>2</sub>O<sub>3</sub> in the thin film deposited at 0.015 mTorr illustrates that the exsolution of Ni and Co promote local degradation of the perovskite phase associated with A-site separation.<sup>34</sup> In addition, the association of the secondary phase La<sub>2</sub>O<sub>3</sub> with an adjacent core-shell nanostructure indicates that exsolution of Ni<sub>x</sub>Co<sub>1-x</sub>O takes place around self-assembled Pd.

### Tunable mixed-valence structure and its role in electrical transport

To understand the effect of nanostructure formation on material properties, electrical conductivity measurements identified the conduction mechanism as n-type electronic transport by small polaron hopping (Figure 6B). Importantly for potential applications, electronic conductivity values are tunable via  $P_{O_2}$  from 0.1 to 10 S/cm, and in line with the most-conductive LaFeO<sub>3</sub>-type materials reported. Temperature-dependent cyclic voltammetry (CV) measurements enabled the calculation of conductivity activation energies ( $E_a$ ) ranging from 79 to 232 meV (Figures S23 and





**Figure 6. Tunable mixed-valence structure and its role in polaron hopping electronic transport**

(A) XPS quantitative analysis of valence states of B-site cations and oxygen species in thin films deposited at different  $P_{O_2}$ .

(B and C) Electronic conductivity as function of temperature (B) and activation energy,  $E_a$  (C), of bulk PLD precursor and thin films, with comparison to typical LaFeO<sub>3</sub>-based perovskites in the literature (Table S3).

(D) Influence of  $P_{O_2}$  on mixed-valence structure of thin films to modulate  $E_a$  of electronic conductivity.

S24; Tables S4 and S5; Notes S12 and S13), similar to expected values for Co- and Ni-facilitated polaron hopping in B-site substituted LaFeO<sub>3</sub>-type materials (Figure 6C; Table S4).<sup>35–40</sup> The importance of the mixed-valence structure of cations (i.e., their content and distribution)<sup>41</sup> is apparent by the fact that the conductivity of the CCO films decreases with decreasing  $P_{O_2}$ .

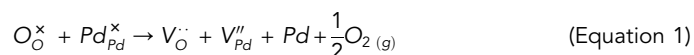
Thus, XPS was used to quantify the cation valence state in the films, which is tunable via  $P_{O_2}$  due to the binding of  $V_O^\bullet$  to reduced cations in the matrix lattice and nanostructure formation (Figures 6A and S17–S19; Table S3; Note S10). Given the extensive Pd nanostructure formation, contributions of Pd<sup>2+</sup> and Pd<sup>0</sup> increase with

decreasing  $P_{O_2}$  (increasing  $V_O^{\bullet\bullet}$ ). Fe valence is nearly insensitive to  $V_O^{\bullet\bullet}$  concentration given its lower reducibility than Co and Ni.<sup>42</sup> Co valence is highly sensitive to  $P_{O_2}$ , because  $Co^{2+}$  usually accompanies  $V_O^{\bullet\bullet}$  formation for electroneutrality,<sup>43</sup> with the  $Co^{2+}$  concentration here increasing from 0% to 65% as  $P_{O_2}$  decreased from 3 to 0.015 mTorr. Although it is difficult to quantify Ni valence due to peak overlap of Ni 2p<sup>3/2</sup> with La 3d<sup>3/2</sup> in XPS (Note S14), qualitative XPS analysis shows mixed Ni<sup>2+</sup>/Ni<sup>3+</sup> in the 3 mTorr film, corresponding well with the fact that Ni is more reducible than Co<sup>42</sup> (Figure S25).

In addition to XPS, the distribution of Co and Ni mixed valency was probed by spatially resolved STEM EELS of a Pd nanorod and core-shell nanostructure, revealing increased amounts of Co<sup>2+</sup> and Ni<sup>2+</sup> surrounding the nanostructures (Figures S20 and S22; Note S11). This indicates that self-assembled Pd triggers Co<sup>3+</sup> and Ni<sup>3+</sup> reduction as part of the formation mechanism by a seed effect that is discussed below.

## DISCUSSION

As demonstrated, ESA yields tunable nanostructures spontaneously in one self-assembly step, while allowing the cation exsolution sequence to be predicted by relative reducibility,<sup>17</sup> highlighting the power of the chemical complexity of CCOs to create diverse and unique nanocomposite materials. In this case, lowering  $P_{O_2}$  causes strongly reducible Pd to take electrons from lattice oxygen during reduction to Pd metal, releasing oxygen (Equation 1). Co and Ni with relatively weak reducibility may also take electrons from lattice oxygen to reach a reduced valence state, but they form a thermodynamically stable suboxide phase instead of further reduction to metal (e.g., Equation 2).<sup>12</sup>



As mentioned above, Ni/Co-coated Pd nanorods are self-assembled by a variant of a previously identified mechanism<sup>15</sup>; however, we observed an intricate core-shell nanostructure growth mechanism that relies on sequential Pd metal phase formation, which seeds the mixed-oxide shell growth and whose constituents use percolating mass transport channels. Recalling that the Ni<sub>x</sub>Co<sub>1-x</sub>O shell is polycrystalline (Figures 4D, S9K–S9N, and S10) and each Ni<sub>x</sub>Co<sub>1-x</sub>O grain shows an epitaxial relation with its corresponding Pd core surface (Figures 4E and S10), it is hypothesized that Ni and Co exsolution and nucleation and growth of Ni<sub>x</sub>Co<sub>1-x</sub>O are triggered by Pd cores via a seed effect<sup>45</sup> (Figure 6F; Note S6). Following Figure 4F, during PLD, (i) exsolved Pd crystals nucleate on the surface of the film; (ii) additional Pd diffuses to and coarsens the nuclei; (iii) Pd crystals seed the self-assembly of reduced Co and Ni due to the Pd–matrix interfacial stress effect (see GPA; Note S5); (iv) Ni<sub>x</sub>Co<sub>1-x</sub>O phase thickens around the Pd to form a shell via diffusion of Ni and Co along percolating channels; and (v) Pd core growth is restricted by Ni<sub>x</sub>Co<sub>1-x</sub>O shell and nanostructure growth is confined by the CCO matrix, causing new Pd nuclei on the surface and generating new core-shell nanoparticles via repeating steps (ii)–(v). Interestingly, cavities appear at both edges and interiors of Pd metal cores (Figures 4G, 4H, S11, and S14; Note S9), supporting the hypothesis that cavities in the Pd metal arise from the nano-Kirkendall effect<sup>46</sup> induced by local surface oxidation of Pd by O<sub>2</sub> released from lattice oxygen during exsolution<sup>47</sup> of Ni and Co (Figure 4I; Note S9). During surface oxidation, Pd vacancies form in the metal due to their faster

outward diffusivity compared to the inward diffusivity of O, generating Pd-depleted cavities that merge at the surfaces and interior of Pd cores (Figure 4I(i)). This allows self-assembling  $\text{Ni}_x\text{Co}_{1-x}\text{O}$  growth within cavities, again via percolating channels (Figures 4I(ii), S15, and S16).

Lastly, interpretation of the correlation of electronic transport with nanostructure formation considers their effects on the CCO perovskite matrix, whose mixed-valence cation network needed for polaron hopping is strongly affected by ESA (Figure 6D). LaFeO<sub>3</sub>-based perovskites facilitate small polaron hopping (Note S12), but given the recent emergence of CCOs, HEOs, and ESOs, their electronic (and generally electrical) transport mechanisms are not well understood because they depend on interrelated factors such as the concentration and type of elements present, amount and type of disorder, and spin states.<sup>2</sup> Hence, the interpretation of electronic conductivity and activation energy focuses on three aspects: cation mixed valence, cation site occupancy/vacancy, and secondary phases. Mixed-valence cation pairs facilitate polaron hopping, and decreasing their separation decreases the activation energy for polaron conduction by increasing electron wavefunction and orbital overlap.<sup>48</sup> Cation vacancies lower the concentration of hopping sites, and thus absolute conductivity value, without affecting the relative amounts of mixed-valence pairs (i.e., without affecting the activation energy). Secondary phases, concomitant with cation vacancies, decrease the total volume of the polaron conducting phase. As detailed in Note S12, the low conductivity ( $\sim 10$  mS/cm) of the ceramic PLD precursor is attributed to insulating La<sub>2</sub>O<sub>3</sub> secondary phases, and the sparse network of B-site cation mixed valency they induce. We believe it is feasible to mitigate the formation of La<sub>2</sub>O<sub>3</sub> and propose several design strategies regarding CCO targets and PLD growth conditions (Note S17).

All thin films have higher conductivity than the precursor, although their conductivities decrease with decreasing  $P_{\text{O}_2}$  (i.e., increasing  $V_{\text{O}}^{\bullet\bullet}$  concentration), suggesting a complex electronic transport mechanism that is sensitive to nanostructure formation and cation mixed valency (Figure 6D). Ru surface accumulation is expected to produce  $\sim 5\%$  B-site vacancies ( $V_{\text{Ru}}^{\text{||}}$ ) compensated by  $V_{\text{O}}^{\bullet\bullet}$  in the matrix. The 3-mTorr film is the only single-phase sample (Figures 1 and S1) and is the most conductive despite also having the highest  $E_a$  of 232 meV (Figures 6B and 6C), which, given no XPS or EELS evidence for  $\text{Co}^{2+}/\text{Co}^{3+}$  mixed valency, is attributed to  $\text{Ni}^{2+}/\text{Ni}^{3+}$ -facilitated polaron hopping (Figures 6A, S20, and S21; Note S11). The 0.15-mTorr film has Pd nanorods with Ni- and Co-enriched nanorod-matrix interfaces (Figures 2H, 2K, and S7) and lower conductivity and  $E_a$  (79 meV) than the 3-mTorr film (Figures 6B and 6C); lower  $E_a$  is attributed to increased  $\text{Co}^{2+}$  concentration, which decreases the Co, and presumably Ni, small polaron pair separation,<sup>48</sup> and approaches that of 50% Ni-substituted LaFeO<sub>3</sub> (66 meV).<sup>35</sup> The 0.015-mTorr film contains connected Pd-Ni<sub>x</sub>Co<sub>1-x</sub>O core-shells and La<sub>2</sub>O<sub>3</sub> volumes; compared to the 0.15-mTorr film, insulating La<sub>2</sub>O<sub>3</sub> lowers conductivity, whereas higher  $E_a$  (152 meV) is attributed to the further increased  $\text{Co}^{2+}$  concentration that widens Co and Ni small polaron pair separation and thus  $E_a$ ,<sup>48</sup> which is comparable to 30% Ni-doped LaFeO<sub>3</sub> ( $E_a = 156$  meV).<sup>35</sup> It is also hypothesized that  $V_{\text{O}}^{\bullet\bullet}$  causes strong electron localization at a subset of cations (e.g., Fe<sup>49</sup>), which halts charge transport near these cation-vacancy associates.<sup>49,50</sup> This reduces the total percolating volume for current carriers (via other mixed-valence cations, e.g., Co and Ni), leading to a more tortuous conduction pathway and lower conductivity.

In summary, we demonstrated tunable nanostructure growth in a CCO by a simple method: ESA. ESA is a one-step approach to direct the evolution of nanoparticles

and nanorods in CCO-derived composites in a perovskite  $\text{LaFeO}_3$ -inspired model CCO, and we present a case study demonstrating how to simply access different nanostructures via the control of oxygen partial pressure. This synthetic approach is expected to be generally applicable to various substrate materials and device architectures as suggested through successful growth on multiple crystalline and amorphous oxide and metal substrates here. A detailed examination of  $\text{LaFe}_{0.7}\text{Ni}_{0.1}\text{Co}_{0.1}\text{Pd}_{0.05}\text{Ru}_{0.05}\text{O}_{3-\delta}$  CCO thin films with atomic-scale electron- and atom-probe imaging and spectroscopies revealed that ESA can be directed using Ellingham's model of cation reducibility, and mass transport via extended defects governs local morphology formation. Furthermore, ESA enables the tailored growth of novel multielement nanorod and nanoparticle composite structures whose formation is correlated with the electronic conductivity of films, which exceeds 0.1 S/cm at room temperature. Given the vast combinatorial space of CCOs, which includes HEOs and ESOs, ESA is expected to be highly extensible via the integration of novel compositions and crystal structures.

## EXPERIMENTAL PROCEDURES

### Resource availability

#### Lead contact

Further information and requests for resources and reagents should be directed to and will be fulfilled by the lead contact, William J. Bowman ([will.bowman@uci.edu](mailto:will.bowman@uci.edu)).

#### Materials availability

Materials generated in this study will be made available on request, but we may require a payment and/or a completed materials transfer agreement if there is potential for commercial application.

#### Data and code availability

The data generated and analyzed in this study are available from the corresponding author on reasonable request. This paper does not report original code. Any additional information required to reanalyze the data reported in this paper is available from the [lead contact](#) upon request.

### Sample preparation

The CCO powders were prepared with a method derived from the Pechini method. Nitrate salts of the constituent metals were mixed in stoichiometric ratios and dissolved in water. Citric acid (3 mol/mol perovskite) and ethylene glycol butyl ether (1 mol/mol perovskite) were subsequently added while the solution was being stirred and heated at 60°C. The solution was dried overnight at 110°C, which resulted in xerogel formation. The xerogel was ground finely in a mortar and pestle and then calcined in a muffle furnace at 800°C for 4 h.

The prepared CCO powders were uniaxially pressed (Strongway 20-Ton Hydraulic Shop Press) at 100 MPa and then followed by cold-isostatic compress (EQ-CIP20-DIE, MTI Corporation) at 200 MPa to make a PLD precursor target pellet. The pellet was sintered at 1,200°C for 12 h in air with a ramp rate of 10°C/min in a tube furnace (Carbolite Gero CC-T1).

Thin films were fabricated using PLD with an Nd:YAG solid laser (Continuum SL III-10, 266-nm wavelength, 10-ns pulse duration) in a high-vacuum deposition system (Neocera Pioneer 180 GLAD PLD System). Silicon (100) single crystal substrates (MTI Corporation) are ultrasonically cleaned in acetone, isopropanol, and methanol

before PLD. The cleaned substrates are fixed at the center of the substrate holder with the substrate-target distance of 57 mm. During PLD, 30,000 laser pulses with an energy fluence of 5.3 J/cm<sup>2</sup> and a laser frequency of 5 Hz were applied to the PLD target to produce a plasma plume deposited on the silicon substrate with a surface temperature of 650°C. We chose this temperature based on prior work showing high crystallinity and density.<sup>51</sup> Mass flow of high-purity O<sub>2</sub> (99.994%) and turbomolecular pump speed were used to tune  $P_{O_2}$  during PLD.

### Characterization

The crystal structure and lattice distortion of the prepared thin films and the CCO PLD target were examined by XRD (Rigaku SmartLab set to 40 kV and 44 mA) in a symmetric  $\theta$ -2 $\theta$  configuration using Cu K $\alpha_1$  radiation (1.54 Å wavelength). The XRD of thin films was operated in general parallel beam/parallel slit analyzer mode with 0D detector, whereas the XRD of the target pellet was conducted in general Bragg-Brentano mode with a 1D detector. Thin film surface morphology was observed with SEM (TESCAN GAIA-3 XHM field-emission SEM [FESEM]) operated at 7 kV with a work distance of 5 mm. The sample surface was sputtered with 5 nm Ir to avoid charging (EMS 150T sputter coater).

STEM specimens were prepared by focused ion beam (FIB) operated at 30 kV for trenching and at 15 or 3 kV for thinning (TESCAN GAIA-3 XHM FESEM). Thin films morphology, phase identification, chemical composition, and local electronic structure were characterized by STEM EDS and EELS using JEOL Grand ARM300CF, equipped with dual 100-mm<sup>2</sup> silicon drift detectors and Gatan GIF Quantum K2. Data acquisition and processing, including GPA and FFT analysis, have been conducted in Gatan Microscopy Suite (GMS 3).

Cation valence states and oxygen species of the prepared thin films were derived and analyzed by XPS carried out using a Kratos AXIS-Supra instrument with a monochromatic Al K $\alpha$  X-ray source. The high-resolution XPS spectra for C1s, O1s, La 3d, Fe 2p, Co 2p, Ni 2p, Pd 3d, and Ru 3d were recorded. The binding energies of all of the spectra were calibrated by standard C1s peak at 284.8 eV for charging effect correction. Peak components of the acquired XPS spectra were analyzed quantitatively using CasaXPS software.

To prepare samples for APT analysis, portions of the TEM specimens were welded onto Si microposts and sharpened by annular milling in an FEI Helios dual-beam FIB microscope. APT was performed with a local-electrode atom-probe (LEAP) 5000 XS (Cameca) at a sample temperature of 30 K and a background pressure of  $5 \times 10^{-11}$  Torr. A 355-nm ultraviolet laser was used to evaporate the sample at a pulse rate of 250 kHz and a pulse energy of 22 pJ. IVAS 3.8.6 was used to reconstruct the data and provide a 3D composition profile.

### Electrical transport measurements

Several 30-nm Ir contacts were deposited onto the surface of the thin films as electrodes via sputtering (EMS 150T sputter coater). In-plane CV tests were carried out using a Biologic sp-200 electrochemical workstation associated with Linkam probe stage (HFS600E-PB4) and Linkam temperature controller (T96 System Controller) for measuring electronic conductivity and deriving activation energy for polaron hopping (Figure S19). Out-of-plane CV tests were conducted for measuring electronic conductivity of the PLD target in which silver ink (Fuel Cell Materials, Nexceris) was coated on both sides as electrodes.



## SUPPLEMENTAL INFORMATION

Supplemental information can be found online at <https://doi.org/10.1016/j.matt.2023.12.012>.

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## AUTHOR CONTRIBUTIONS

H.G. and W.J.B. conceived and designed this research. S.S. prepared the CCO powder under the supervision of K.L.G.-A. H.G. fabricated the CCO PLD target, prepared the thin films, and performed XRD and SEM characterization. H.G. and X.W. prepared TEM samples with FIB. H.G., X.W., M.X., and T.A. performed the STEM-EDS and EELS characterization. H.G. and M.X. conducted the STEM GPA. C.M. performed APT analysis under the supervision of L.J.L. I.T. and H.G. performed the XPS experiments. M.B. conducted the CV tests. J.D.S. performed statistics on self-assembled nanophases. H.G., C.M., M.B., and W.J.B. interpreted the characterization and property data. H.G. wrote the manuscript. All coauthors commented on the manuscript and contributed to the final revision. W.J.B. supervised the study.

## DECLARATION OF INTERESTS

The authors declare no competing interests.

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