

# 1 Electrothermoplasmonic Trapping and Dynamic Manipulation of 2 Single Colloidal Nanodiamond

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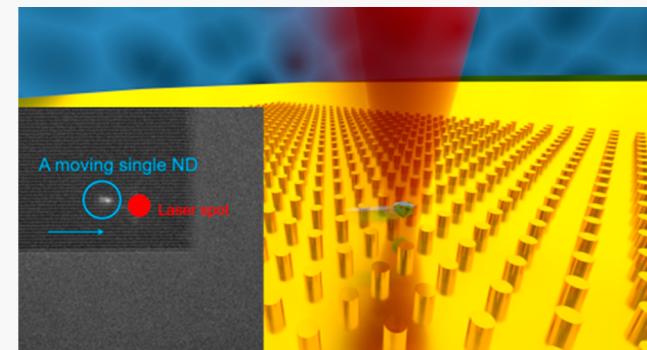
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4 **ABSTRACT:** Low-power trapping of nanoscale objects can be  
5 achieved by using the enhanced fields near plasmonic nano-  
6 antennas. Unfortunately, in this approach the trap site is limited to  
7 the position of the plasmonic hotspots and continuous dynamic  
8 manipulation is not feasible. Here, we report a low-frequency  
9 electrothermoplasmonic tweezer (LFET) that provides low-power,  
10 high-stability and continuous dynamic manipulation of a single  
11 nanodiamond. LFET harnesses the combined action of the laser  
12 illumination of a plasmonic nanopillar antenna array and low-  
13 frequency alternating current (ac) electric field to establish an  
14 electrohydrodynamic potential capable of the stable trapping and  
15 dynamic manipulation of single nanodiamonds. We experimentally  
16 demonstrate the fast transport, trapping, and dynamic manipulation  
17 of a single nanodiamond using a low-frequency ac field below 5 kHz and low-laser power of 1 mW. This nanotweezer platform for  
18 nanodiamond manipulation holds promise for the scalable assembly of single photon sources for quantum information processing  
19 and low noise quantum sensors.



20 **KEYWORDS:** nanotweezers, photothermal effect, electrohydrodynamics, nanodiamond

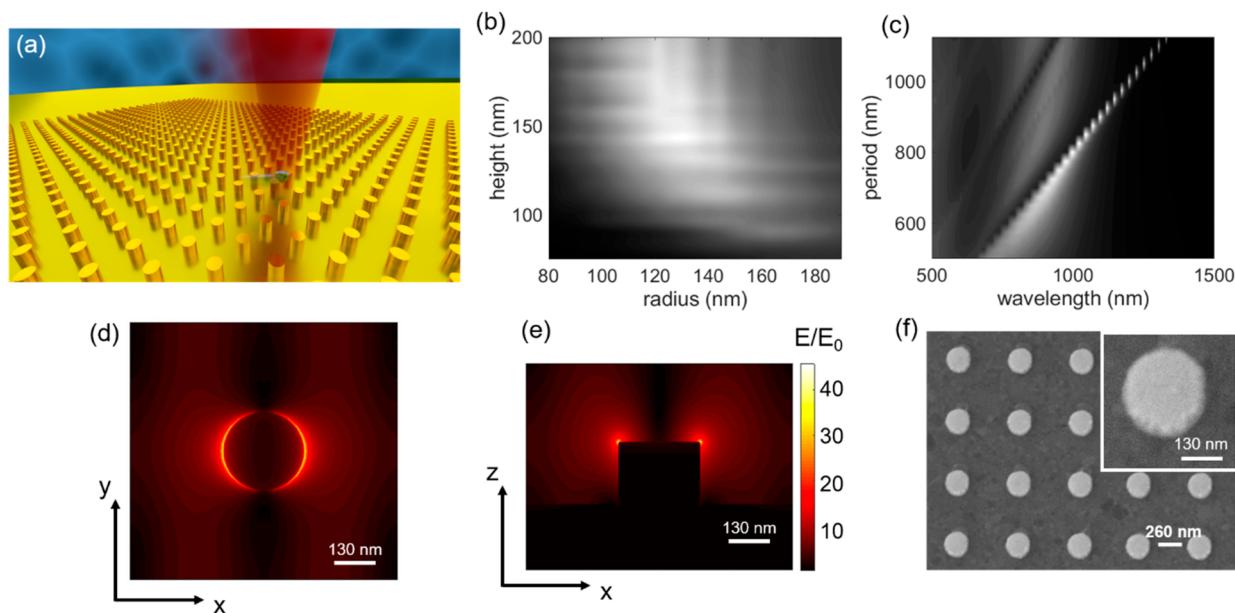
21 **P**lasmonic nanoantennas, which can confine and enhance  
22 local electromagnetic fields, are a powerful platform for  
23 the development of optical nanotweezers.<sup>1–5</sup> Upon resonant  
24 illumination, plasmonic nanoantennas create highly localized  
25 and enhanced electromagnetic fields within volumes well-  
26 beyond the diffraction-limit<sup>6,7</sup> and thus produce narrower and  
27 deeper trapping potential wells than the conventional optical  
28 tweezers. This capability facilitates trapping nanoscale particles  
29 using relatively low laser power with high stability.<sup>8–10</sup> The use  
30 of optical forces from plasmonic nanoantennas is particularly  
31 suited for the trapping and positioning of colloidal nano-  
32 diamonds (NDs) near plasmonic cavities to enhance the light–  
33 matter interaction. In particular, for quantum photonics  
34 applications NDs have been identified as stable quantum  
35 emitters<sup>11–13</sup> capable of providing single photon emission  
36 under room-temperature conditions.<sup>14</sup> To enhance their  
37 emission properties, it is crucial to develop methods that can  
38 rapidly trap and couple them to respective nanophotonic  
39 cavities. Because of their small sizes, NDs are more accessible  
40 for trapping using plasmonic nanotweezers in comparison to  
41 the conventional optical tweezers. However, on-chip plasmonic  
42 nanotweezers are only able to trap the nanoscale objects at the  
43 specific position of the hotspot defined by the plasmonic  
44 nanoantenna and do not possess dynamic manipulation  
45 capability along the plasmonic substrate. Prior work by Lin  
46 et al.<sup>15</sup> reported the continuous dynamic manipulation of  
47 single nanoscale charged gold nanoparticles using a thermo-

electric field induced by the photothermal heating of a porous<sup>48</sup> gold plasmonic substrate submerged in an ionic surfactant. In<sup>49</sup> this approach, ionic surfactants (cetyltrimethylammonium<sup>50</sup> chloride) are introduced in the colloidal solution, which in<sup>51</sup> the presence of a thermal gradient establishes the attractive<sup>52</sup> thermophoretic force to enable the trapping of the gold<sup>53</sup> colloids with single particle resolution. This approach has yet<sup>54</sup> to be translated to the continuous dynamic manipulation of<sup>55</sup> single nanoscale dielectric objects.<sup>15–18</sup> Another recently<sup>56</sup> reported technique by Gosh et al. employs dielectric microrods<sup>57</sup> that are coated with plasmonic nanodisks.<sup>19</sup> The dielectric<sup>58</sup> microrods are large enough to be optically trapped with<sup>59</sup> conventional optical traps, while the light coupled to the<sup>60</sup> plasmonic nanodisks serves as plasmonic tweezers that can trap<sup>61</sup> the nanoscale objects in solution. Dynamic manipulation of the<sup>62</sup> trapped nanoscale objects is achieved in an indirect fashion by<sup>63</sup> manipulating the optically trapped dielectric microrods.<sup>64</sup>

Here, we capitalize on the latest advances in electro-<sup>65</sup> thermoplasmonic tweezers<sup>20–22</sup> and present a method termed<sup>66</sup>

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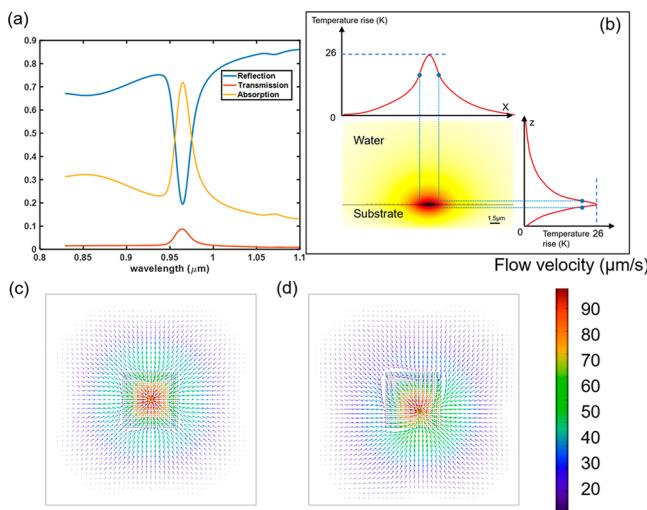


**Figure 1.** (a) Three-dimensional illustration of the ac-assisted nanotweezer enabled by an array of plasmonic nanoantenna capable of dynamic manipulation of single nanodiamonds. A single nanodiamond is trapped and repositioned by the laser motion relative to the nanoantenna array. (b) The absorption cross section efficiency map obtained from sweeping the radius and height of the gold nanopillars in an array. (c) The absorption cross section efficiency spectra when we parametrically sweep the lattice constant of an infinite array. The resonance wavelength is tuned to 973 nm. (d) The simulated electric field distribution in one unit-cell on the top of the nanopillar. (e) The simulated vertical electric field distribution in one unit-cell across the center of the nanopillar. In (d,e), the plane-wave excitation is along the z-direction with a 973 nm wavelength. The color bar is normalized by the incident light electric field amplitude. (f) The SEM image of the fabricated gold-nanopillar array.

67 low frequency electrothermoplasmonic tweezer (LFET) that  
 68 uses low-frequency alternating current (ac) electric field  
 69 combined with a laser illumination for the trapping and  
 70 continuous dynamic manipulation of a single ND (70–100 nm  
 71 in diameter) suspended in deionized water environment. Using  
 72 LFET, we can trap a single ND on the top of an array of a  
 73 plasmonic nanoantenna and dynamically manipulate the single  
 74 ND by simply moving the laser spot or microscope stage, as  
 75 shown in Figure 1a. Our nanotweezers platform is composed  
 76 of an array of gold nanopillars illuminated with a near-IR laser  
 77 (973 nm wavelength) and perpendicularly applied low-  
 78 frequency ac electric field to optically induce thermal gradients  
 79 and distort the ac electric field, respectively. To obtain a  
 80 maximized photothermal conversion efficiency, while keeping  
 81 the laser power low, we parametrically optimize the absorption  
 82 cross section of one unit-cell by sweeping the radius, height,  
 83 and lattice constant of the gold nanopillar array, as shown in  
 84 Figure 1b,c. The coupling of the gold nanopillar array and the  
 85 induced localized plasmon resonance of the single antennas  
 86 results in an enhanced and spatially confined electric field  
 87 distribution, as shown in Figure 1d,e, so that the photothermal  
 88 conversion efficiency is enhanced. The optimized gold  
 89 nanopillars are of 130 nm in radius and 150 nm in height  
 90 and are placed on a 50 nm thick gold film on a glass substrate.  
 91 The lattice constant of the array is 720 nm, which ensures that  
 92 the highest absorption cross section is centered around 973 nm  
 93 wavelength. We also notice that tuning the size and period of  
 94 the gold nanopillar does not have a significant effect on the  
 95 distribution of the applied ac electric field as shown in Figures  
 96 S1 and S2. We attribute this to the fact that the wavelength of  
 97 the ac field is much larger than the tuning range of the size and  
 98 period of the gold nanopillar. As a result, the optically  
 99 optimized design with a radius of 130 nm and period of 720  
 100 nm was chosen for this work. The scanning electron

101 microscopy (SEM) image of the fabricated sample is depicted 102  
 103 in Figure 1f. The fabrication was performed using the standard 104  
 105 nanofabrication approach composed of electron-beam lithog- 106  
 107 raphy and lift-off (see Methods). 108

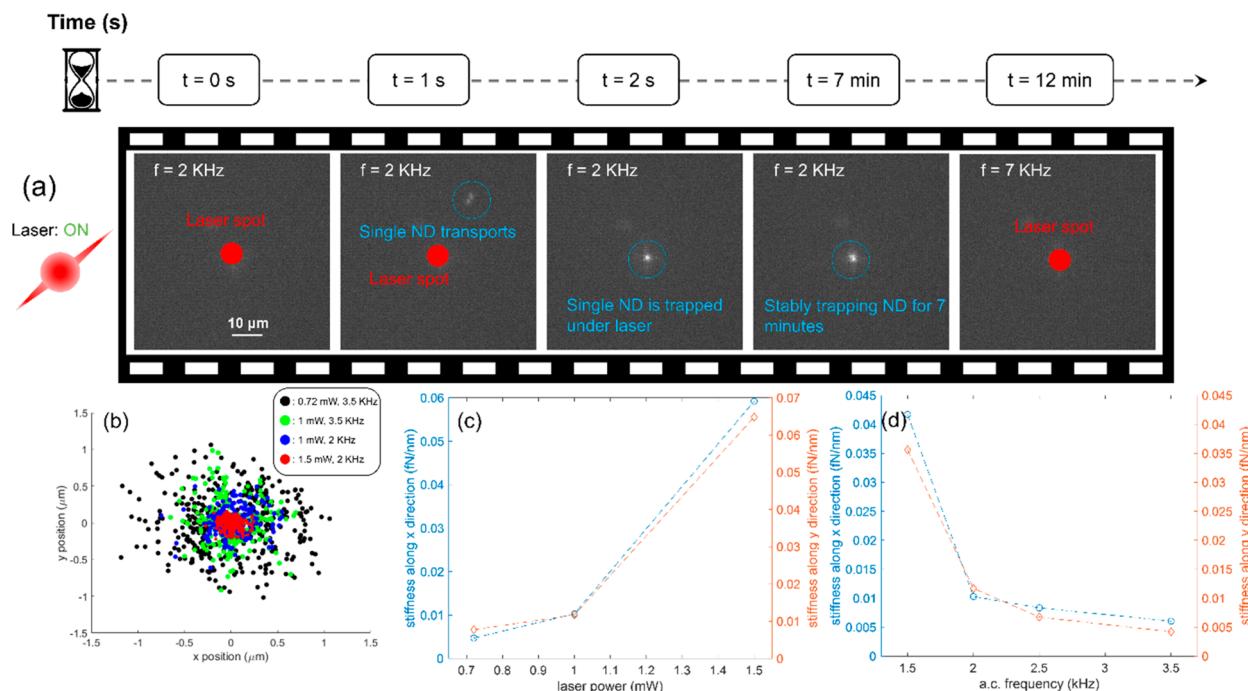
109 Using the optimized design, the absorption, reflection, and 110  
 111 transmission spectra are numerically calculated and plotted in 112  
 113 Figure 2a. The absorption cross-section is obtained by solving 114  
 115 the Maxwell's equation using a finite-difference-time-domain 116  
 117 commercial solver, Lumerical FDTD. The enhanced light 118  
 119 absorption leads to temperature rise and thermal gradient in 120  
 121 the fluid. The simulated temperature that rises for a 1 mW 122  
 123 laser illumination and focused to a spot size of 1  $\mu$ m radius is 124  
 125 depicted in Figure 2b. An ac electric field is applied across the 126  
 127 microfluidic channel to establish the electrothermoplasmonic 128  
 129 (ETP) flow because of the action of the laser-induced thermal 130  
 131 gradient and applied ac field, which enables fast transport and 132  
 133 the in-plane manipulation of trapped NDs. The gold nanopillar 134  
 135 array distorts the local ac electric field to create both normal 136  
 137 and tangential ac electric-field components. These tangential 138  
 139 components exert lateral Coulombic force on the diffuse 140  
 141 charges in the electrical double layer (EDL), which produces 142  
 143 local ac electro-osmotic (ACEO) motion of the fluid. The total 144  
 145 flow fields at a position of 50 nm above the array when the 146  
 147 laser is focused at the center and off the center of the array are 148  
 149 depicted in Figure 2c,d, respectively. They show the ability of 150  
 151 our nanotweezers to reposition the trapping location site by 152  
 153 relocating the laser spot on the gold nanopillar array. 154  
 155 Additionally, the primary force responsible for the localization 156  
 157 of the trapped nanodiamonds along the out-of-plane direction 158  
 159 is due to the particle–surface interaction force, which results 160  
 161 from the interaction between the surface charges of the particle 162  
 163 and their image charges in the conducting plane.<sup>23</sup> This 164  
 165 particle surface interaction force increases at lower ac 166  
 167 frequencies below the charge relaxation frequency of the 168



**Figure 2.** (a) The simulated absorption, reflection, and transmission spectra of an infinite array of gold nanopillar array. The absorption spectrum is centered at a peak wavelength of 973 nm to maximize the photothermal conversion efficiency at this wavelength. (b) The simulated temperature rise in the vicinity of the nanopillar array using 1 mW laser of 973 nm and 1  $\mu\text{m}$  laser spot radius. Line plots indicate the thermal confinement is better along the vertical direction than along the horizontal direction because of thermal spreading into the high thermal conductivity gold film. The total flows including ACEO and ETP flows simulated using the same laser settings as (b) with the laser spot at (c) the center of the array and (d) off from the center of the array. The induced temperature field from (b) in conjunction with an applied ac electric field generates a strong ETP flow centered around the laser spot position. By moving the laser position, the flow fields are spatially translated along the plasmonic substrate.

fluid. The interplay between the microfluidic flows, optical gradient force, and particle–surface interaction force enhanced at low ac electric-field frequencies enables the rapid particle loading process, localization of a single ND near the gold nanopillar array, and dynamic manipulation of a single ND. A detailed explanation of the respective contributions of the forces acting on a single nanodiamond in the system is presented in the Supporting Information along with the numerical analysis of the particle–surface interaction force, repulsive thermophoretic force, and optical gradient force. Furthermore, an illustration of the direction of the forces is depicted in Figure S3. The electrohydrodynamic physics to model the flows is simulated using the finite-element-solver in COMSOL Multiphysics software.

Experimental demonstration was performed using diluted solution of NDs (Sigma-Aldrich fluorescent nanodiamond) with an average particle size of 70–100 nm in deionized water. The NDs are diluted to a concentration of  $5.65 \times 10^5$  particle/mL. A linearly polarized laser beam with 973 nm wavelength was focused to a spot size of 1  $\mu\text{m}$  radius on the gold nanopillar array using a water immersion objective lens with a numerical aperture of 1.2. Figure 3a shows the process of fast transport, trapping, and releasing of a single ND particle under different ac frequencies, as shown in SI Video 1. The laser power was set to 1 mW and the ac frequency was initially set to a low frequency of 2 kHz to guarantee the trapping. First, we illuminated the nanopillar region with the laser but without the ac field applied and no trapping phenomenon was observed. Then an ac electric field at a frequency of 2 kHz was applied perpendicular to the gold film. The applied peak-to-peak voltage was 10 V and the microfluidic chamber height is 120  $\mu\text{m}$ , corresponding to an ac electric field amplitude of 83 333



**Figure 3.** (a) The sequence of frames shows the fast process of a single ND transported into the hotspot and stably trapped on the gold nanopillar pattern for 7 min at 2 kHz. By tuning up the ac frequency to 7 kHz, the single ND is released rapidly. (b) The scattering plot of the trapped single ND under various ac field and laser power conditions. (c) The stiffness calculated at 2 kHz ac frequency with varying laser power applied. Higher laser power enables higher trapping stiffness. Blue circles represent the stiffness along the x-direction, and orange diamonds represent the stiffness along y-direction.

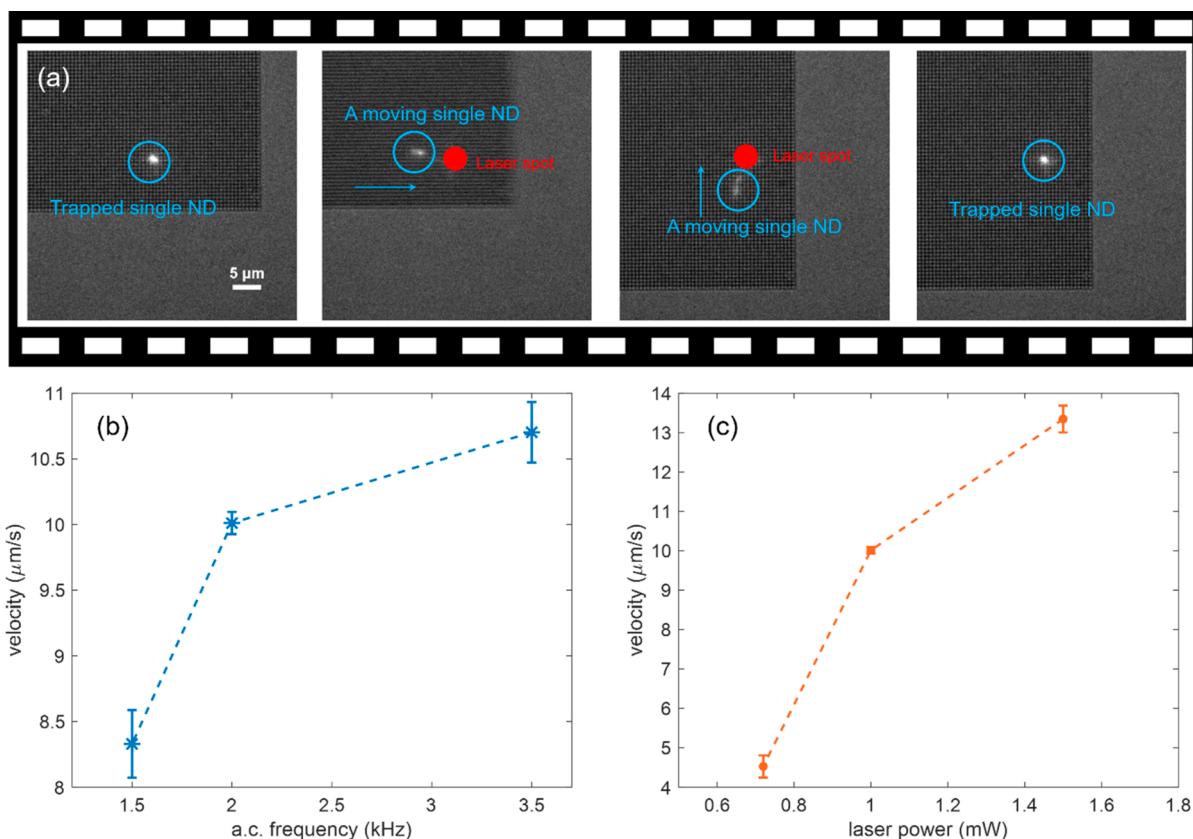
167 V/m. Immediately, an ETP flow was generated inside the  
168 microchannel resulting in the rapid transport of a single ND to  
169 the illuminated gold nanopillar array region. These results  
170 contrast with our previous findings describing opto-thermo-  
171 electrohydrodynamic tweezers (OTET)<sup>20</sup> using a finite array  
172 of gold nanoholes. The difference in the particle trapping  
173 behavior for the nanopillar array versus the nanohole array  
174 arises because the distribution of the induced ac electro-  
175 osmotic flow field vectors act in a manner to induce a strong  
176 centripetal drag force (i.e., radially inward) on the nanoscale  
177 objects toward the nanopillar array. Furthermore, the nano-  
178 pillar array possesses a stronger radially inward ETP flow that  
179 acts to transport target objects toward the illuminated laser  
180 spot. Thus, unlike in OTET<sup>20</sup> the nanoscale objects are not  
181 repelled from the nanopillar array to establish a stagnation  
182 zone as depicted in Figure S5 (see Supporting Information).  
183 Upon arrival of the ND to the laser spot, it is stably trapped at  
184 the position of the laser illumination as long as the laser  
185 illumination and the low-frequency ac field are simultaneously  
186 present. After 7 min, the ac electric field frequency was  
187 increased from 2 to 7 kHz and the ND was rapidly released  
188 from the trap by the axial component of the ETP flow as  
189 depicted in Figure 3a. This is because the particle–surface  
190 interaction force reduces as the ac frequency is increased.  
191 Thus, the drag force from the axial ETP flow and the repulsive  
192 thermophoretic force overcomes the particle–surface inter-  
193 action force to convect the particle away from the surface of  
194 the nanopillar array.

195 The trapping stability under various laser power and ac  
196 frequency conditions are described in the scatter plots in  
197 Figure 3b. The trapping stability for the data presented in  
198 Figure 3a is depicted in blue. The trap stiffness was determined  
199 by using the equipartition theorem approach<sup>24</sup> as described in  
200 Section 4 of Supporting Information. Figure 3c depicts the  
201 trapping stiffness along the *x*- and *y*-direction under varying  
202 laser power conditions for a constant ac field frequency of 2  
203 kHz. The results show that the trapping stiffness is increasing  
204 with increasing laser power. We attribute this enhancement in  
205 the in-plane trapping stiffness to the fact that the in-plane  
206 optical gradient force and the radial ETP flow drag force  
207 increases with higher laser power. The variation of the trapping  
208 stiffness with the applied ac frequency under a fixed laser  
209 power of 1 mW is depicted in Figure 3d. The result shows an  
210 inverse relation whereby the trapping stiffness is increasing  
211 with decreasing ac field frequency. This is because for lower ac  
212 frequency, the particle–surface interaction force along the  
213 vertical direction and local ACEO flow are both strength-  
214 ened.<sup>25,26</sup> In general, as depicted in Figure 3b, the trapping  
215 stability increases with increasing laser power and decreasing  
216 ac field frequency. The laser power still needs to be kept within  
217 the limits that prevents excessive temperature rise to prevent  
218 laser-induced bubble formation. Thus, LFET provides the  
219 means to harness the synergistic interaction of optical gradient  
220 force and electrostatic particle–surface interaction force to  
221 achieve the stable trapping and dynamic manipulation of a  
222 single nanodiamond. It should be noted here that the gold  
223 nanopillar array provides strong electric field enhancement,  
224 which can enhance the optical gradient force and the  
225 photothermal heating of the gold nanopillar array. The  
226 photothermal heating of the gold nanopillar array gives rise  
227 to a positive (i.e., repulsive) thermophoretic force. This  
228 thermophoretic force is repulsive in the absence of ionic  
229 surfactants in the fluid medium and thus repels the suspended

230 nanodiamond from the nanoantenna.<sup>17,27</sup> A quantitative  
231 analysis is presented in Section 7 of Supporting Information  
232 to show the contribution from the thermophoretic force that  
233 repels the nanodiamond from the surface of the gold  
234 nanopillars. We also performed experiments whereby the  
235 nanodiamond was initially trapped using the ac field and laser  
236 illumination. Subsequently, the ac field was turned off with the  
237 laser still on and the nanodiamond was observed to escape  
238 from the trap as shown in SI Video 3. This result suggests that  
239 the optical gradient force alone was not sufficient to stably trap  
240 the nanodiamond in the presence of the repulsive thermo-  
241 phoretic force.

242 To investigate this further, we have calculated the repulsive  
243 thermophoretic force and optical gradient force along the *z*-  
244 direction (axial) as depicted in Figure S8. The result shows a  
245 peak thermophoretic force with a maximum of about 0.33 pN  
246 at the illumination power of 1 mW. We have compared the  
247 optical gradient force and repulsive thermophoretic force in  
248 the axial direction under varying laser powers ranging from 0.5  
249 to 1.5 mW. The result presented in Figure S8b shows that the  
250 axial repulsive thermophoretic force is larger than the optical  
251 gradient force for all the laser powers considered. Thus, in the  
252 presence of laser illumination of the nanopillar array alone, the  
253 repulsive thermophoretic force would prevent the trapping of  
254 the nanodiamond with optical force alone. This repulsive  
255 thermophoretic force has been reported to prevent stable  
256 plasmonic optical trapping using plasmonic nanoantenna in  
257 earlier works unless the temperature rise is minimized such as  
258 by using very low optical illumination intensity,<sup>28</sup> off-resonant  
259 excitation,<sup>29</sup> or integrating a heat sink.<sup>5</sup> In LFET, with the low-  
260 frequency ac electric field applied the induced electrostatic  
261 particle–surface interaction force has been estimated to be  
262 about 10.5 pN at a location of 20 nm from the gold nanopillar  
263 as depicted in Figure S4, which is 32 times larger than the peak  
264 repulsive thermophoretic force at the same location. The  
265 particle surface interaction force increases as the nanodiamond  
266 approaches the gold nanopillar surface, but it is eventually  
267 balanced by the electrical double-layer repulsion force to  
268 enable the nanodiamond to reach an equilibrium height from  
269 the gold nanopillar surface and prevent the nanodiamond from  
270 spontaneously sticking to the gold nanopillar surface. This  
271 equilibrium height is estimated to be approximately 20 nm  
272 from the surface of the gold nanopillar as depicted in Figure  
273 S4a. Thus, in LFET the electrostatic particle–surface  
274 interaction force induced by applying a low-frequency ac  
275 electric field conveniently overcomes the repulsive thermo-  
276 retic force to still enable the trapping and dynamic  
277 manipulation of a single nanodiamond despite the repulsive  
278 thermophoretic force. Furthermore, the LFET approach allows  
279 us to achieve high stability trapping at low laser power (1  
280 mW), which is smaller than the previously reported powers of  
281 50–100 mW using a laser beam optical trap for the same size  
282 of single nanodiamond (~100 nm).<sup>30,31</sup>

283 The trapping stability remains approximately the same when  
284 the nanodiamond is trapped at the center of the pattern or near  
285 the edge of the nanopillar array as shown in Figure S6 of the  
286 Supporting Information. To verify single nanodiamond  
287 trapping, the trapped nanodiamond is patterned onto the  
288 gold nanopillar array and observed under scanning electron  
289 microscope (SEM). The diameter of the trapped nanodiamond  
290 is about 100 nm, as shown in the SEM image depicted in  
291 Figure S7 of the Supporting Information.



**Figure 4.** (a) The frame-by-frame sequence showing a single ND is stably trapped and manipulated along the gold nanopillar array using 1 mW laser power and a low-frequency ac field of 2 kHz. After both laser and ac field are turned on, a single ND is transported toward the illuminated spot within a few seconds. Subsequently, the laser spot is translated toward the right and the single ND moves rapidly toward the new laser spot and is stably trapped at the new location. (b,c) Transport velocity of the trapped nanodiamond under various ac frequency and laser power conditions. Blue asterisks represent the velocities under laser power from 0.72 to 1.5 mW with 2 kHz ac frequency. Orange dots represent the velocities under ac frequency for 1.5, 2, and 3.5 kHz with 1 mW laser power. For the same experimental condition, we repeated the experiment several times to calculate the average velocity. Error bars indicate the standard deviation of all the acquired velocities under the same experimental condition.

292 Besides trapping and releasing the NDs, the dynamic  
 293 manipulation of the trapped NDs can be achieved using  
 294 LFET, a capability not feasible in conventional chip-based  
 295 plasmonic traps<sup>3,5</sup> based on near-field optical force alone or  
 296 geometry-induced electrostatic traps.<sup>32</sup> The dynamic manipu-  
 297 lation of the trapped NDs is achieved by translating the  
 298 microscope stage or the laser beam focus along the in-plane  
 299 direction. The translation of the laser focus from one position  
 300 to another along the nanopillar array creates new thermal  
 301 hotspots and hence new ETP flow that rapidly convects the  
 302 trapped NDs to the newly illuminated spot along the surface of  
 303 the nanopillar array as demonstrated in Figure 4a and SI Video  
 304 2. The manipulation can be achieved over the whole patterned  
 305 region as long as the laser spot is kept on the nanopillar array.  
 306 The velocity of the transport of the nanodiamond depends on  
 307 the laser power and ac frequency. To clarify, as shown in the  
 308 second and third frames of Figure 4 the single nanodiamond is  
 309 not necessarily always within the laser spot during the motion  
 310 of the laser spot illumination across the nanopillar array. When  
 311 the laser spot is suddenly translated to another position, due to  
 312 the contribution from ETP flow the single nanodiamond is  
 313 delivered to the new laser spot site rapidly. We have also  
 314 presented data on the average velocity of the nanodiamond  
 315 transport given by the displacement between the initial laser  
 316 spot and final laser spot position divided by the time span that  
 317 the single nanodiamond takes to arrive at the new location of

318 the laser spot. Higher laser power generates more photo-  
 319 thermal-heating, so that the ETP flow is accelerated and the  
 320 transport velocity is enhanced as depicted in Figure 4. For  
 321 example, the transport velocity increases from 4.5 to 13.2 μm/s  
 322 when the laser power was increased from 0.72 to 1.5 mW.  
 323 With respect to variation of ac frequency, the nanodiamond  
 324 transport velocity reduces with decreasing ac field frequency.  
 325 We attribute this observation to the stronger electrostatic  
 326 particle–surface force that increases with decreasing ac field  
 327 frequency, resulting in more frictional drag on the nano-  
 328 diamond as the ac frequency is reduced as depicted in Figure  
 329 4b,c. For all of the experimental conditions, the achieved  
 330 transport velocity ranges from 4 to 14 μm/s. This ability to  
 331 quickly deliver the nanodiamond to the position of the laser  
 332 spot enables facile dynamic manipulation across the surface of  
 333 the fold nanopillar array. In SI Video 4, we show a typical video  
 334 used for calculating the transport velocity.

335 In conclusion, we have demonstrated a low-frequency  
 336 electrothermoplasmonic tweezer that harnesses low-frequency  
 337 ac electric field below 5 kHz combined with low laser power  
 338 for the dynamic manipulation of single NDs along a gold  
 339 nanopillar array substrate. The ability of our nanotweezer to  
 340 trap a single ND within a few seconds and rapidly relocate it  
 341 on a plasmonic substrate holds promise for the implementation  
 342 of on-chip single photon sources, plasmonic nanolasers, and  
 343 low-noise quantum sensors.

## 344 ■ METHODS

345 **Gold Nanopillar Array Fabrication.** After a float glass  
346 substrate was cleaned up, we deposited 50 nm of gold film with  
347 5 nm of chromium as the adhesion layer. Then, photoresist  
348 (PMMA A4) was spin-coated on for electron beam lithography  
349 to define the pattern region. After development, another layer  
350 of 150 nm gold was deposited. Finally, we finished the process  
351 using acetone for lift-off, and a square gold nanopillar array  
352 with 80  $\mu\text{m}$  on the side was left on the chip.

353 **Sample Preparation.** After the nanopillar was fabricated,  
354 we sandwiched the gold film by covering it with a glass  
355 coverslip that has a thin coating of indium tin oxide spaced by  
356 a 120  $\mu\text{m}$  thick dielectric spacer to create a microfluidic  
357 channel around the patterns. Two copper wires are connected  
358 to the ITO side on the coverslip and on the gold film to apply  
359 the ac electric field.

360 The fluorescent nanodiamond solution originally had a  
361 concentration of 1 mg/mL in deionized water with nitrogen  
362 vacancy of  $\sim$ 3 ppm and purchased at Sigma-Aldrich. The  
363 nanodiamond was diluted by 100 000 times using deionized  
364 water to generate a sparse enough solution suitable for single  
365 nanodiamond manipulation. The final concentration used in  
366 the experiment was  $5.65 \times 10^5$  particle/mL.

367 **Fluorescence Imaging.** The trapping and imaging were  
368 performed using a custom fluorescent imaging and optical  
369 trapping microscope based on a Nikon Ti2-E inverted  
370 microscope. The suspended particle solution was injected  
371 into the microfluidic channel. A high quantum efficiency  
372 sCMOS camera (Photometrics PRIME 95B) was used to  
373 acquire images at a frame rate of 3.3333 frames per second.  
374 The trapped nanodiamonds were excited under a green light  
375 from a filtered broadband fluorescent illumination lamp  
376 (Nikon INTENSILIGHT C-HGFI). The emitted red light  
377 was collected through the same objective lens and imaged on  
378 the camera. The nanopillar array was illuminated with a 973  
379 nm semiconductor diode laser (Thorlabs CLD101S). The laser  
380 beam was focused with a Nikon 60 $\times$  water-immersion  
381 objective lens (NA, 1.2). The ac electric field was supplied  
382 by a dual-channel function generator (BK Precision 4047B).

## 383 ■ ASSOCIATED CONTENT

## 384 ■ Supporting Information

385 The Supporting Information is available free of charge at  
386 <https://pubs.acs.org/doi/10.1021/acs.nanolett.1c00357>.

387 Video showing the ability of a low-frequency electro-  
388 thermoplasmonic trap to rapidly trap a single nano-  
389 diamond and to release the trapped single ND by  
390 increasing the ac frequency (AVI)

391 Video showing that after the single ND is trapped a  
392 dynamic manipulation is able to be achieved by  
393 translating the laser spot across the nanopillar array  
394 (AVI)

395 Video showing the recorded process of a trapped  
396 nanodiamond escaping from the trap when the ac field  
397 is turned off (AVI)

398 Video showing typical video frames used for measuring  
399 the average manipulation velocity of the nanodiamond  
400 (AVI)

401 Video showing a trapped nanodiamond near the edge of  
402 the pattern, corresponding to the data in Figure S6  
403 (AVI)

404 The ac field with various sizes, force illustration and  
405 particle–surface interaction, ac osmosis simulation,  
406 stiffness analysis, SEM of deposited nanodiamond,  
407 thermal simulation from FEM software, balance between  
408 thermophoretic force and optical gradient force, full  
409 explanations of supporting videos (PDF)

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## 433 Notes

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