

1 Fundamental Investigation of Reactive-Convective Transport: Implications for

2 Long-Term Carbon dioxide (CO₂) Sequestration

3 Md Fahim Shahriar¹, Aaditya Khanal^{1*}

4 ¹The Jasper Department of Chemical Engineering, The University of Texas at Tyler

5 Corresponding author E-mail address: aadityakhanal@uttyler.edu

6 Abstract

7 The density-driven convection coupled with chemical reaction is the preferred mechanism for permanently storing
8 CO₂ in saline aquifers. This study uses a 2D visual Hele-Shaw cell to evaluate and visualize the density-driven
9 convection formed due to gravitational instabilities. The primary goal of the experiments is to understand the various
10 mechanisms for the mass transfer of gaseous CO₂ into brine with different initial ionic concentrations and flow
11 permeability. Moreover, the impact of CO₂ injection locations, reservoir dipping angle, and permeability heterogeneity
12 is also investigated. We observed that the presence of salts resulted in earlier onset of convection and a larger
13 convective finger wavelength than the case with no dissolved salts. Additionally, a higher lateral mixing between CO₂
14 fingers is observed when dipping is involved. The CO₂ dissolution, indicated by the area of the pH-depressed region,
15 depends on the type and concentration of the ions present in the brine and is observed to be 0.38-0.77 times compared
16 to when no salt is present. Although convective flow is slowed in the presence of salts, the diffusive flux is enhanced,
17 as observed from both qualitative and quantitative results. Moreover, the reduced formation permeability, introduced
18 by using a flow barrier, resulted in numerous regions not being swept by the dissolved CO₂, indicating an inefficient
19 dissolution. We also investigated the effect of discrete high-conductivity fractures within the flow barriers, which
20 showed an uneven vertical sweep and enhanced flow channeling. Lastly, the parameters regarding CO₂ leakage risk
21 during storage are identified and discussed.

22 **Keywords:** CO₂ Sequestration; Convective dissolution; Reactive Dissolution; Rayleigh-Taylor instability;
23 Heterogeneity; Permeability Contrast

24 1. Introduction

25 The atmospheric concentration of carbon dioxide (CO₂) in May 2021 was recorded as 419.5 parts per million (ppm),
26 which is approximately 50% higher than at the beginning of the industrial revolution [1]. The growing concerns about
27 CO₂ emissions have led to investigations of possible carbon capture and storage (CCS) methods. CO₂ geological
28 storage in depleted oil reservoirs or saline aquifers is a preferred CCS
29 method to capture emissions from large point sources [2]. CO₂
30 geological storage has four primary CO₂ trapping mechanisms:
31 structural, residual, dissolution, and mineral trapping, as shown in
32 Fig. 1 [3]. Capillary or residual trapping rate is the highest at the
33 beginning of the CO₂ storage period; however, dissolution trapping
34 becomes more dominant throughout time (Fig. 1). Dissolution
35 trapping captures almost two-thirds of CO₂ injected in the storage
36 volume [4–6]. Even though molecular diffusion of CO₂ in brine is
37 slow, the rate of dissolution trapping is accelerated by other mechanisms, including density-driven convection or
38 Rayleigh-Taylor instability, dispersion, and advection [7,8]. Several factors can affect the convective-dissolution
39 phenomenon for different geological sites and, therefore, need to be considered for accurate CO₂ storage prediction.
40 For example, density-driven convection is reduced for geological storages with lower vertical permeability, resulting
41 in advection becoming the dominant force with increased transverse mixing [4,9]. Nevertheless, the effect of density-
42 driven convection in CO₂ geological storage is significant, as it prevents CO₂ leakage by sinking CO₂ rather than rising
43 to shallower formations [4]. Various studies have identified and extensively studied different aspects of the transport
44 mechanisms during the density-driven convection or Rayleigh-Taylor instability, including diffusion period,
45 convective period, and constant flux [10–15].

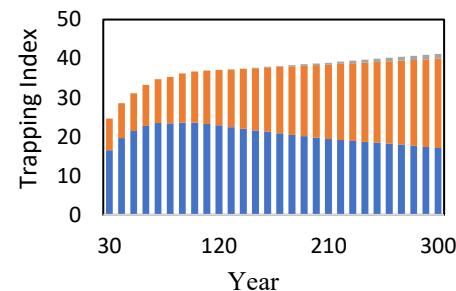


Figure 1. Fraction of CO₂ sequestered by residual trapping (blue), dissolution (orange), and mineralization (gray) (Modified from Khanal and Shahriar [3]). Dissolution is the primary storage mechanism especially with the passage of time.

46 Hele-Shaw cell, a simple structure usually formed by a narrow gap between two transparent flat plates, has
47 been used in multiple experimental setups to visualize and mimic the Rayleigh-Taylor instability or Rayleigh
48 convection (RC) formed during CO₂ storage in different geological structures [16–19]. Kneafsey and Pruess [20]
49 conducted laboratory visualization studies and quantitative CO₂ absorption tests in transparent Hele-Shaw cells to
50 investigate the dissolution-induced density-driven convection phenomenon. The quantitative measurements showed
51 that the density-driven convection initiated faster than predicted. Backhaus et al. [21] studied the density-driven

52 convection for a lighter fluid (water) placed over a heavier fluid (propylene glycol). The initial instability and quasi-
53 steady-state were explained by analyzing the convective time and velocity scales, finger width, and mass transport
54 rate. The test was conducted at standard atmospheric conditions with a Rayleigh number (Ra) of 6000-90,000 [21].
55 Meanwhile, a smaller Ra range ($100 < Ra < 1700$) was adopted in the works of Slim et al. [22]. Potassium
56 permanganate (KMnO₄) in water was used as an analog for CO₂ in brine at atmospheric conditions, describing the
57 dissolution-driven convective behavior from the first contact up to 65% average saturation.

58 Developing scaling relationships, correlations, and models can provide important insight into the convection
59 dissolution properties of CO₂ for different geological storage [2,23–26]. Motjaba et al. [23] developed two scaling
60 relationships, one between Rayleigh and Sherwood numbers and the other between Rayleigh numbers and CO₂
61 convective flux for under 3.45 MPa and $182 < Ra < 20860$. Robust scaling relations between compensated flux and
62 transition times between successive regimes in the system for different salt types (NaCl and CaCl₂) were examined by
63 Mahmoodpour et al. [27]. The results showed that different salt types affect both the short and long-term dynamics of
64 convective dissolution. Faisal et al. [24] obtained correlations between the Rayleigh number and the mass of total
65 dissolved CO₂. Other useful scaling laws, including the onset of the convection, and wavelength of the initial
66 convective instabilities, were also identified and discussed [5,28,29].

67 Different visualization techniques, including the Schlieren method, particle image velocimetry (PIV), laser-
68 induced fluorescence (LIF), and interferometry method, have been adopted to observe the formation and growth of
69 convective finger structures in Rayleigh convection [2,4,30,31]. Zhang et al. [2] presented a vortex model of CO₂
70 adsorption into the water to characterize the interfacial mass transfer coefficient for the continuous convective period.
71 The study adopted particle image velocimetry (PIV) and laser-induced fluorescence (LIF) to calculate the Hele-Shaw
72 cell's solute concentration distribution and instantaneous liquid velocity. Moreover, another recent work by Zhang et
73 al. [32] used a different experimental setup using the UV-induced fluorescence method to investigate gas-liquid
74 interphase mass transfer in a Hele-Shaw cell.

75 Mahmoodpour et al. [27] provided critical insight into visualizing the dissolution-driven convection at high-
76 pressure conditions (up to 535.3 psi). They devised a novel Hele-Shaw apparatus withstanding high pressure and
77 presented CO₂ dissolution-driven convective behavior in a confined brine-saturated porous medium. Tang et al. [4]
78 designed a Hele-Shaw cell rated to 70 MPa and Ra of 346 to investigate the convection parameters, including critical

79 onset time of convection, dissolution rate, and gravitational instabilities. This study used the micro-schlieren technique
80 to conduct the visual inspection. Pressure-volume-temperature (PVT) testing was conducted at 293.15 to 423.15 K
81 and pressure ranging from 14 to 24 MPa.

82 The chemical composition of the CO₂ storage site also significantly affects the reactive transport of dissolved
83 CO₂, as observed in the works of Thomas et al. [33]. Their work investigated the dissolution of CO₂ into an aqueous
84 solution of bases MOH, where M⁺ is an alkali metal cation. For bases MOH, the convection is enhanced for the
85 counter-ion M⁺ sequence of Li⁺ < Na⁺ < K⁺ < Cs⁺. The experimental investigation revealed that the base concentration
86 in solution strongly impacts the nonlinear finger instability, where higher concentration leads to faster instability and
87 shorter time for onset of convection. Furthermore, despite M⁺ ions not actively participating in the geochemical
88 reactions during the dissolution process, the nature of different M⁺ ions varies in the instability development. Loodts
89 et al. [34] observed the effect of pressure, temperature, and NaCl concentration on CO₂ dissolution properties. Their
90 study suggested that increasing CO₂ pressure or reducing temperature or salt concentration leads to higher convective
91 instability. However, temperature has a minimal effect on CO₂ dissolution properties, so controlling the temperature
92 is not essential for the reproducibility of experimental studies [34]. Thomas et al. investigated the effect of salinity by
93 the dissolution of gaseous CO₂ in pure water, Antarctic water, and 0.5-5 M NaCl dissolved in water [35]. The results
94 showed that higher salt concentration delays the formation of instabilities, resulting in delayed onset of convection.
95 Moreover, increased convection pattern wavelength and decreased fingers' velocity and the growth rate increased the
96 salt concentration. Kim and Kim [36] derived and solved linear stability equations for the effect of chemical reactions
97 in an initially quiescent vertical Hele-Shaw cell. Their nonlinear numerical simulation showed that chemical reactions
98 enhance the diffusive flux; however, by retarding the onset of buoyancy-driven convective motion, convective flux is
99 weakened.

100 Formation dip angle is another key factor of consideration for safely storing CO₂ on subsurface geological
101 sites, as it significantly impacts spatial migration distribution during CO₂ dissolution [37–39]. For larger dip angles,
102 the supercritical CO₂ phase could change to a gas phase during upward migration in the reservoir up-dip direction,
103 where the reservoir formation temperature and hydrostatic pressure are lower [38]. As a result, reservoirs with higher
104 dip angles have more chance of CO₂ leakage during geological storage. Jang et al. [37] simulated the effect of dip
105 angle and salinity of CO₂ storage. For formation dip angles of 0°, 5°, and 10°, the migrated CO₂ distances were 60%,

106 73.3%, and 86.7%, respectively, compared to a 15° dip angle in the 200th year of CO₂ migration. Therefore, with a
107 larger formation dip angle, there is a higher possibility of spatial CO₂ migration. They concluded that reservoirs with
108 higher dip angles and salinity have low CO₂ geological storage safety. Wang et al. [38] observed similar effects of
109 formation dip, where the total CO₂ storage amount is inversely proportional to the formation dip angle. The impact of
110 dip angle is more prominent in storage reservoirs with higher porosity and permeability [38]. As Jing et al. [35]
111 observed, higher salinity and high dip angle are not conducive to CO₂ geological storage. However, the effect of
112 salinity is more significant than that of dip angle on the CO₂ liquid phase mass fraction.

113 During the injection of CO₂ in deep saline aquifers, the natural fractures present in the formation may
114 propagate, or new fractures may be induced in the reservoir. The fracture networks in a hydrocarbon reservoir play a
115 vital role in fluid transport from the pores to the wellbore as they are significantly more conductive than the matrix
116 [40–43]. The same principle is applicable during the CO₂ sequestration operation, which makes it difficult to predict
117 the movement of plumes during the injection of CO₂ in fractured porous media. The highly permeable conductive
118 fracture networks can act as a pathway for fluid movement, potentially allowing CO₂ migration to neighboring aquifers
119 or the surface through the cap rock [44,45]. Bond et al. [44] demonstrated improved CO₂ migration prediction by
120 incorporating structural geological fractures in the model. Knowing the spatial distribution of fractures, their
121 orientation, conductivity, and overall contribution to the effective permeability are desirable for geological sites whose
122 permeability is controlled mainly by faults and fractures [45].

123 Despite the considerable investigative analyses on CO₂ dissolution-driven convection using the Hele-Shaw
124 cell, some areas need further investigation. To the best of our knowledge, very few studies have considered the
125 combined effect of dipping angle and varying salinity on CO₂ dissolution flow behavior. Since both the dipping angle
126 and use of salts inhibit CO₂ dissolution, their combined effect needs to be further investigated. Moreover, the effect of
127 varying CO₂ injection points needs further insight into quantitative analysis to make informed decisions about
128 choosing optimal injection points. Additionally, despite the considerable theoretical model and experimental
129 investigation on the effect of salinity in CO₂ dissolution, most of the work considers NaCl and ignores the presence
130 of other salts. And to the best of our knowledge, the effect of fractures in heterogeneous media is yet to be presented.
131 This experimental study critically investigates the effect of these factors to address this knowledge gap. Furthermore,
132 this study discusses how the findings of this work can be translated into improving storage efficiency by providing

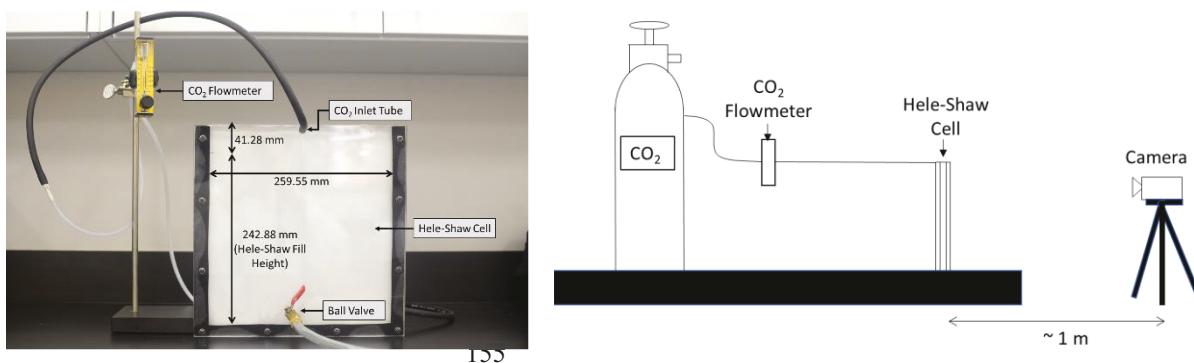
133 key insight into CO₂ convection dynamics. This study primarily investigates the effect of varying several parameters
134 on the CO₂ mass transfer mechanisms in atmospheric conditions. However, since the actual reservoir conditions are
135 at higher pressure and temperature, we provide a brief preliminary evaluation of the effect of high pressure and
136 temperature in the discussions section and leave a more detailed investigation for future studies. Regardless, the
137 fundamental qualitative investigation considered in this study still translates to actual reservoir conditions with similar
138 non-dimensional numbers and thus can be accurately scaled without any loss of generality.

139 The remainder of the paper is organized as follows. **Section 2** provides the experimental setup, and **Section**
140 **3** presents the methods, image processing sequence, and a brief overview of the experiments considered. **Sections 4**
141 **and 5** present the qualitative visualization of CO₂ dissolution in homogeneous media without and with salt,
142 respectively. **Section 6** considers the qualitative visualization of CO₂ dissolution in heterogeneous media.
143 Additionally, the effect of fractures on CO₂ dissolution is also discussed. **Section 7** provides our results in terms of
144 quantitative data. Lastly, **Section 8** is devoted to the discussion, and **Section 9** presents the main conclusions of this
145 study.

146 **2. Experimental Set-up**

147 Experiments on atmospheric conditions (1.01 bar and 22° C) are performed in a Hele-Shaw cell composed of two
148 transparent 0.5 in (12.9 mm) thick plexiglass separated by precision silicone shims with a thickness of 1 mm along
149 the sidewalls. Although glass can have better thickness tolerance than plexiglass, this study considers plexiglass due
150 to higher transparency, allowing enhanced clarity. The front plexiglass panel was drilled at the bottom of the cell and
151 was fitted with a ball valve with a diameter of 0.25 in (6.35 mm). The internal cell dimensions were a length and
152 height of 259 mm and 284 mm, respectively. The height (H) of the water column is 243 mm, as shown in **Fig. 2a**.
153 This port was used to fill and drain the reactor of the experimental fluid.

154



155

(a)

(b)

156 **Figure 2.** (a) The Hele-Shaw cell with dimensions and (b) Schematic of the experimental setup

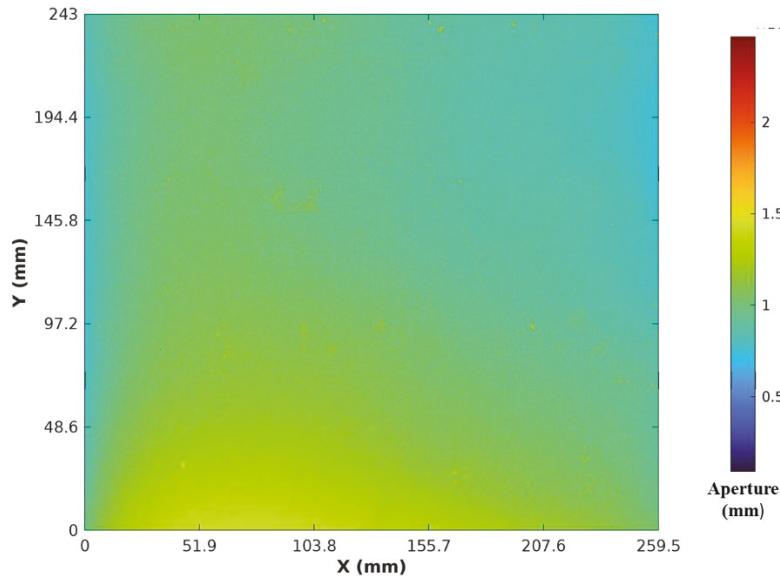
157 At the top, three holes at the center of the front plexiglass were drilled so that the CO₂ could be securely injected into
 158 the cell using an 18-gauge dispensing needle, as shown in **Fig. 2a**. A digital camera (Nikon D7000 with 50 mm lens)
 159 was focused on the cell to take pictures at 20 seconds intervals for 2.5 hours. We maintain a 4 cm distance from the
 160 injection point to the top of the water to avoid disturbing and introducing shear stress at the interface. The experiments
 161 used the same protocol to ensure a similar controlled environment. The schematic of the experimental setup is
 162 presented in **Fig. 2b**.
 163

164 The flow of incompressible fluid in a Hele-Shaw is governed by the Navier-Stokes equation, coupled with
 165 the species conservation equation for cases with chemical reactions [24,46]. The conservation equations for mass,
 166 momentum, and species also implicitly consider the Boussinesq approximation, which assumes that the solute
 167 concentration affects the local density without causing an expansion or contraction of the fluid [47]. The permeability
 168 of the Hele-Shaw cell, k , with a cell aperture of b (for this study, 1 mm), can be obtained from the fundamental
 169 conservation laws as [48]:

$$170 \quad k = \frac{b^2}{12} = \frac{0.001^2}{12} = 8.33 \times 10^{-8} \text{ m}^2 \quad (1)$$

171 As shown in **Eq. 1**, permeability is proportional to the square of the cell square; thus, slight variation in aperture can
 172 result in localized variation of permeability in the cell. As a result, the aperture variation must be checked to ensure
 173 acceptable thickness tolerance. We use the method proposed by Detwiler et al. [49] and Faisal et al. [24] to map the
 174 aperture variation of the Hele-Shaw used in this study. Using ImageJ, the light intensity transmitted through the cell
 175 is analyzed and normalized based on the average value. Then, the average cell aperture, calculated to be around 1.04

176 mm, is found by filling the cell with a known fluid volume. The fluctuations in cell aperture vary proportionally with
177 the normalized light intensity; therefore, by combining the normalized light intensity with average cell aperture, the
178 aperture variation in the cell up to the height of the water column is measured, shown in **Fig. 3**.



179

180 **Figure 3.** Aperture map resulting from the transmitted light analysis on the experimental 1 mm Hele-Shaw cell configuration

181 **Fig. 3** shows that the device used in this study is reasonably homogenous with negligible aperture variations, mainly
182 towards the bottom of the cell. The corresponding Rayleigh number (Ra), a dimensionless number characterizing the
183 system by expressing the ratio of free convection to diffusion, is calculated using the permeability (k) of the porous
184 medium and is represented as shown in **Eq. 2**:

$$191 \quad \text{Ra} = \frac{\Delta \rho g k H}{\mu \varphi D_{CO_2}} \quad (2)$$

185 Where g is the acceleration of gravity [m/s²], H is the height of the water column [m], μ is the dynamic viscosity of
186 water [kg/m.s], D_{CO_2} is the molecular diffusion coefficient of CO₂ in water [m²/s], $\Delta \rho$ is the increase in density due
187 to CO₂ dissolution [kg/m³], k is the permeability of the medium [m²], and φ is porosity. For our experimental cases
188 without any heterogeneity, $\varphi = 1$. The aperture (1 mm) considered in this study is significantly small compared to the
189 height of the water column (243 mm), so the flow environment is mathematically analogous to the Darcy flow in
190 isotropic porous media. Therefore, the results presented in this study follow the same flow properties compared to

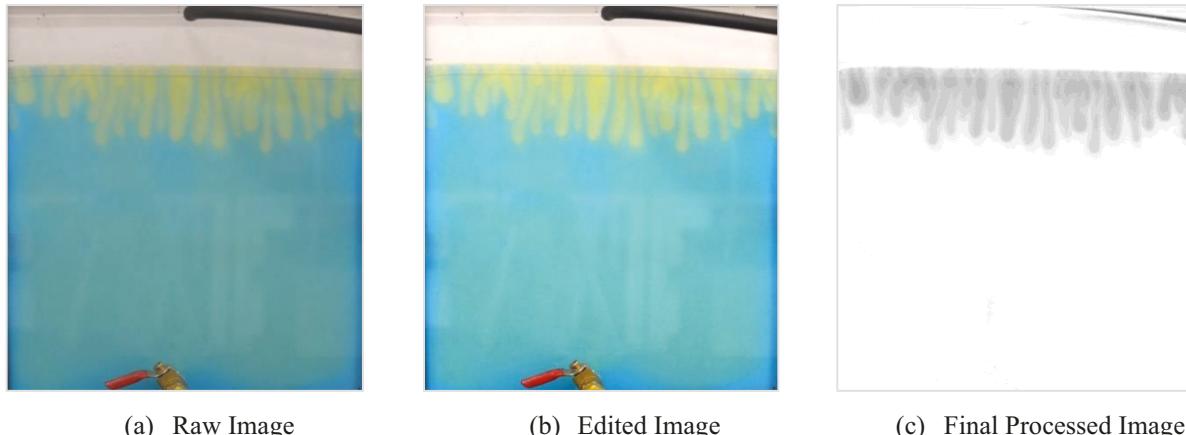
192 porous media [24,46]. However, additional porous material will introduce mechanical dispersion, in addition to
193 convection, due to the flow path tortuosity and will require more in-depth analysis.[50].

194 For the heterogeneous cases, porosity, permeability, and corresponding Rayleigh number are calculated and
195 presented in Section 3.2. We obtain a value of $\text{Ra} = 41782$ for the homogeneous base cases conducted in this study. It
196 should be noted that the Rayleigh number will decrease slightly with an increase in salinity, as observed by Thomas
197 et al. [35]. In their experiment, for salinities of 1M and 2M NaCl, the Rayleigh number decreased by 0.77% and 1.5%,
198 respectively, compared to the Rayleigh number of the water system. We expect salinity to cause a similar decreasing
199 pattern in Rayleigh number for cases with dissolved salts. The values for calculating the Rayleigh number were
200 adopted from Faisal et al. [24], who used a bromocresol green pH indicator at a room temperature of 22 °C, similar to
201 our experiments. As other theoretical studies and experimental investigations observed, for $\text{Ra} \geq 4\pi^2$, natural
202 convection is the predominant flow [24,51]. All the experiments in this study meet that criterion for prevailing
203 convective flow.

204 3. Experimental Procedure and Image Processing

205 The parameters considered for image analysis are the number of fingers formed, average finger length (mm), average
206 wavelength (mm), area of pH depressed region (mm²), and standard deviation of finger length (mm). Finger length is
207 calculated as the vertical distance from the interface to the tip of the finger. In contrast, the average wavelength is the
208 interface length divided by the number of fingers.

209 We followed the same image processing sequence adopted in the work of Kneafsey et al. [52]. Using ImageJ
210 FIJI, we subtracted the initial condition (no CO₂) applied from all the images, resulting in images with less background
211 noise in RGB format. Next, the RGB images were split into three color channels: red, green, and blue. The blue color
212 channel containing mostly noise was discarded, while the two other channels were added, as shown in **Fig. 4**. Lastly,
213 we inverted the grayscale lookup table changing dark to light, making the CO₂-dissolved pH-depressed regions darker.
214 Additional image processing and noise removal were performed on some images for better clarity.



215 **Figure 4.** Image Processing sequence showing (a) unedited image from the camera, (b) edited image with enhanced contrast, (c) grayscale
216 image, which shows the final image used for analysis.

217 Before performing the experiments (**Table 1**), we conducted a repeatability test by comparing the qualitative

218 and quantitative results of three runs for a test case where CO₂ was injected at the middle port of the Hele-Shaw cell
219 filled with a solution prepared by mixing deionized, de-aired water with 0.0114% w/v Bromocresol Green (BCG).

220 This solution is referred to as the *Control Fluid* (CF) in this study. BCG is a pH indicator that changes color from blue

221 at a pH above 5.4 to yellow at a pH of around 3.9 [52]. The application of pH indicators is a widely accepted method
222 for studying convective mixing [16,24,52,53]. Thomas et al. [54] indicated that the color indicator does not

223 significantly affect the development of convective dynamics. Taheri et al. [55] reported no change in water's properties
224 for a small amount of BCG. The solution in their experiment was prepared by mixing distilled deionized water with

225 0.025 wt. % BCG. Therefore, color-based pH indicators can be used when trying to quantify convective-dissolution
226 properties [54,55]. The Hele-Shaw cell was filled at approximately 85% with 63 mL of CF. It is crucial to use a small

227 flow rate so that the water interface at the top of the cell is not disturbed, which can add external perturbation to the
228 interface and influence the viscous finger formation strictly due to gravitational instability. Therefore, we selected a

229 flow rate of 0.59 L/min which showed no impact on the fluid interface, based on multiple runs with different flow
230 rates. The selected flow rates resulted in a slow stream of bubbles when the CO₂ injection needle was placed in a

231 beaker of water [24].

232 The qualitative results showing the convective pattern for the same time scale are presented in **Fig. 5**, which

233 shows good repeatability for the finger size, number, and overall patterns when identical experimental conditions are

234 used for the runs. We also calculated the 95% confidence interval for various parameters that characterize the fluid
235 flow presented later in the quantitative measurement section (**Section 7**).



Run 1

Run 2

Run 3

236 **Figure 5.** Three runs for CO₂ introduced at 0.59 L/min to the left side of the cell at 85 min (CF, 1 mm aperture, vertically oriented flat glass)

237 The experimental investigations considered in our study can be divided into four cases, as presented in **Table**

238 **1.** In **Case A**, we consider the effect of the CO₂ injection point by varying the placement of the CO₂ injection point in
239 the Hele-Shaw cell. Moreover, we also observe the effect of salinity by dissolving 1M CaCl₂ in the Control Fluid.

240 To observe the effect of initial cation concentration, we use MgCl₂ and CaCl₂ with varying concentrations in **Case B**,
241 as shown in **Table 1**. The concentration of commonly dissolved solids in aquifer fluids can vary significantly based
242 on the reservoir [56–60]. For example, the reported concentration of chloride ranges from 1.5 M (in Weyburn
243 reservoir, Canada) to as high as 5.4 M (in Rose Run aquifer, Ohio) [57,59]. In comparison, several studies calculated
244 Calcium concentration ranging from 0.05 to 0.9 M (2004 to 36072 ppm) and Magnesium concentration to be 0.02 M
245 to 0.24 M (486 to 5834 ppm) [56–60]. The Mg⁺⁺ and Ca⁺⁺ cations are used in this study as they have received limited
246 attention compared to numerous studies for Na⁺ cations. Furthermore, Mg⁺⁺ and Ca⁺⁺ cations form a significant
247 fraction of the ions in worldwide saline aquifers [61]. This study aims to provide a comparative view of the impact of
248 different salts; therefore, a uniform concentration of 1M and 2M is considered, which falls in the range of chloride
249 concentration reported in different saline aquifers [57,59], as well as comparable to the concentration used in the work
250 of Mahmoodpour et al. [62].

251 Additionally, we also introduce a slight dipping ($\leq 1^\circ$) to mimic a dipping aquifer by raising the right side of
 252 the Hele-Shaw cell slightly higher than the left. When dipping is involved, the modified Ra is represented as shown
 253 in **Eq. 3** [52]:

$$258 \quad Ra = \frac{\Delta \rho g \sin \theta k H}{\mu \varphi D_{CO_2}} \quad (3)$$

254 *Where, sinθ* is the parameter dictating the effect of gravity due to the dipping introduced. Thus, for a dipping angle
 255 of 1° , our corresponding Ra is 41775. It should be noted that, as we increase the dipping angle, the corresponding Ra
 256 can reduce significantly. For example, Kneafsey et al. [52] reduced their Ra value by a factor of 2 by orienting the cell
 257 30° from the horizontal level.

259 Although **Cases A** and **B** consider homogeneity throughout the cell, CO₂ geological storage sites are
 260 heterogeneous. Barriers like calcite and shale layer are common forms of heterogeneity that can significantly dictate
 261 the density-driven natural convection mechanism [10]. Moreover, different configurations and geometries of the
 262 heterogeneities result in various convective flow directions and movements. Although several studies have introduced
 263 different heterogeneous patterns in the system [10,63,64] to observe and quantify the CO₂ dissolution-driven
 264 convection pattern, the effect of different salts in these heterogeneous systems is yet to be investigated. Therefore, in
 265 **Case C**, we add heterogeneities in the Hele-Shaw system by introducing regions with lower permeability values and
 266 observing the convection pattern change for solutions containing different salts. Moreover, despite the recognized
 267 importance of heterogeneous layers, the impact of fractures present in these heterogeneous formations has not been
 268 studied extensively [40,41]. Therefore, we created narrow spacing inside heterogeneous formations for Case D to
 269 mimic the convective flow through fractured regions. The CO₂ flow rate in all the experiments in **Case B, C, and D**
 270 is kept constant at a minimal flow rate of 0.59 L/min to prevent any disturbance at the water interface on top of the
 271 cell.

272 **Table 1.** The experiments considered in this study and the investigated parameters

Effect of Parameters Observed	Solution Used	CO ₂ Injection Point	Case Number	Dipping Applied	Heteroge neous
CO ₂ injection point and <i>Control Fluid (CF)^a</i>		Side	A1	No	No

Effect of salinity	CF	Middle	A2	No	No
	1 mole CaCl ₂ dissolved in CF	Middle	A3	No	No
Dipping effect and presence of salt types with varying concentration	1 mole CaCl ₂ dissolved in CF	Side	B1	Yes	No
	2 mole CaCl ₂ dissolved in CF	Side	B2	Yes	No
	1 mole MgCl ₂ dissolved in CF	Side	B3	Yes	No
	2 mole MgCl ₂ dissolved in CF	Side	B4	Yes	No
Effect of Heterogeneity	CF	Side	C1	No	Yes
	1 mole CaCl ₂ dissolved in CF	Side	C2	No	Yes
	1 mole MgCl ₂ dissolved in CF	Side	C3	No	Yes
	1 mole NaCl dissolved in CF	Side	C4	No	Yes
Effect of Fracture	CF	Side	D1	No	Yes

^aControl Fluid (CF) is prepared by mixing deionized 1 L of de-aired water with 0.0114% w/v BCG.

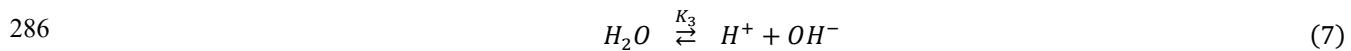
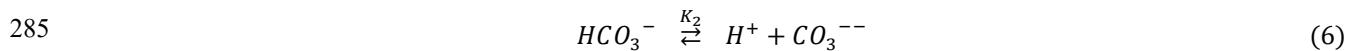
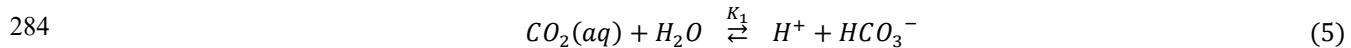
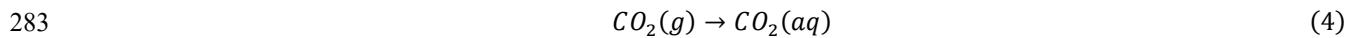
273

274 4. Qualitative Visualization for Homogeneous Case with No Salinity

275 4.1. CO₂ introduced to the side of the cell (Case – A1)

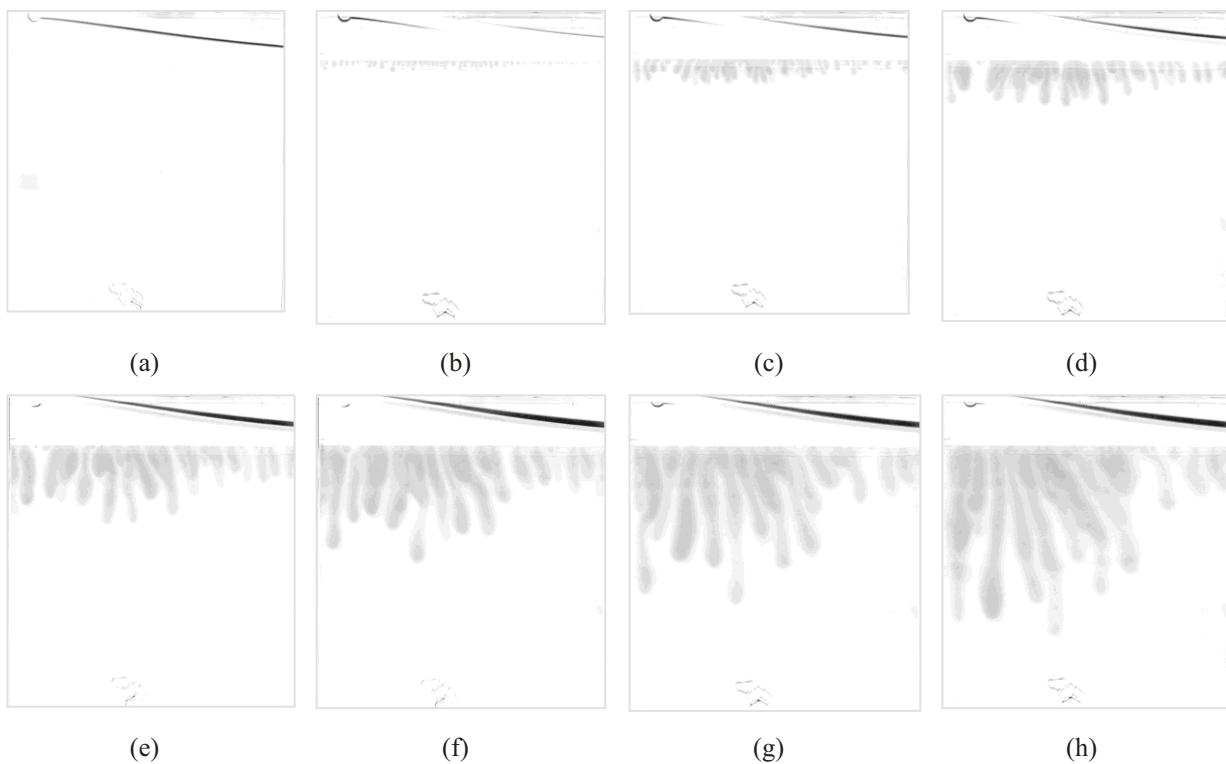
276

277 We introduce the gaseous CO₂ from the leftmost injection port for this case. The development of convection-
 278 driven fingering patterns is similar to what has been observed in previous studies [20,22,35,65]. CO₂ dissolution via
 279 diffusion was the only applicable mechanism at the initial stage of the visualization test, which we call the CO₂
 280 induction phase. During the induction phase, there is a uniform change in pH near the gas-water interface, indicating
 281 the diffusion of CO₂, where the gaseous CO₂ dissolves in water to form aqueous CO₂ and forms equilibrium with
 282 carbonic acid, undergoing the following reactions [61]:



287 During this period, the diffusive layer at the interface gradually expanded without any deformation or
288 instability. We observed the formation of fingers at around 135 ± 15 seconds, demonstrating the end of the induction
289 phase and onset of the convection process causing instability in the form of fingers, as shown in **Table 2**. We
290 considered the time frame for which the fingers are clearly visible; therefore, the actual onset of convection should be
291 slightly earlier since the earliest initiation of finger forming marks the onset of convective flow.

292 The fingers increased in width and depth as time progressed, traveling downwards vertically. The fingers
293 started merging as they grew in width, as evidenced by the decrease in the number of fingers and an increase in the
294 wavelength (please refer to **Section 7.1** for the quantitative measurement). Moreover, as time progressed, we observed
295 cell-scale convection, a phenomenon bringing fresh solution (solution not laden with CO₂) to the upper surface from
296 the center, as shown by small new fingers forming at the gas-water interface. The formation of distinctive new fingers
297 (also referred to as "nascent fingers") between established fingers can be seen as early as 27 minutes into the test. This
298 phase is identified as the reignition phase, which lasts until the end of the experiment, also observed in other studies
299 [22,35].



300

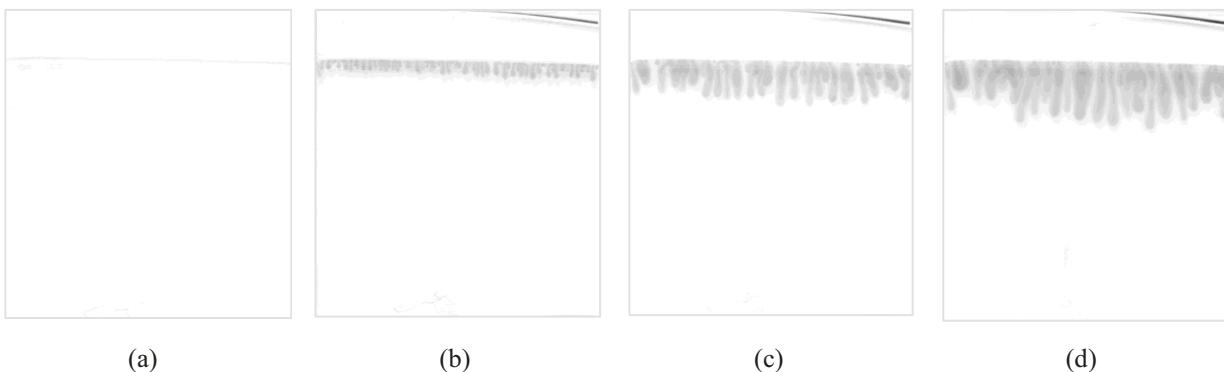
301 **Figure 6.** CO₂ introduced to the left side of the cell (CF, 1 mm aperture, vertically oriented flat glass) (a) Initial (no CO₂) (b) 9 minutes
302 (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h) 154 minutes

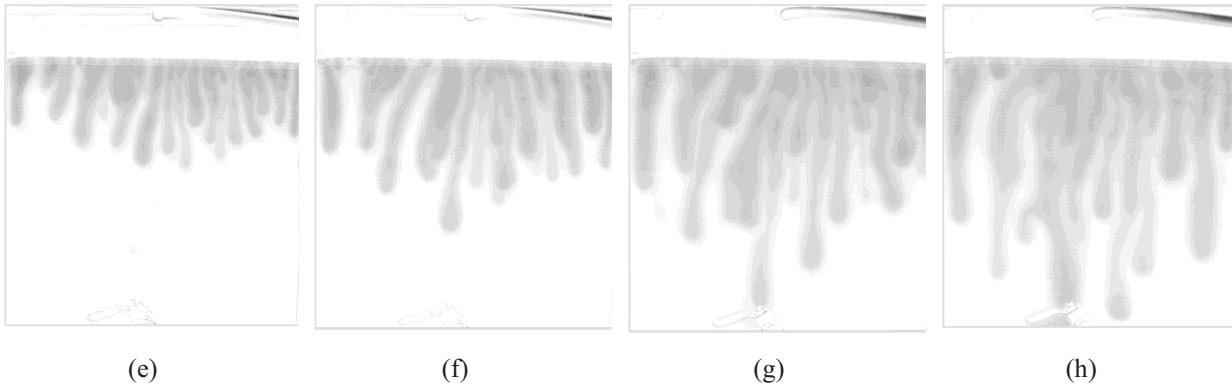
303 The cell-scale convection also leads to the non-vertical movement (slight curve formed) made by the
304 established fingers, which would otherwise move vertically due to density change and gravitational pull. The curved
305 movement of the fingers can be observed as early as 9 minutes into the visualization test, shown in **Fig. 6 (b)**, which
306 becomes more prominent as time progresses.

307 The whole area of the Hele-Shaw can be classified into two regions: active and inactive zones. The active
308 zone can be referred to as the entire vertical region under the CO₂ injection point. In contrast, the region farther from
309 the injection point can be identified as the inactive zone. As seen from **Fig. 6**, when CO₂ is injected into the left side
310 of the cell, the area farther away from it, i.e., the inactive zone, has partially slowed finger formation. This difference
311 can be attributed to the upward fluid flow due to cell-scale convection, which slows down the speed of the fingers.
312 However, since the active zone has continuous CO₂ injection from the top, the effect of cell-scale convection in
313 slowing down the finger flow is not as significant. Furthermore, due to the increasing difference in finger length
314 between the active and inactive zones, we see an increase in the standard deviation of the finger length (refer to **Section**
315 **7.1** for a quantitative overview). Thus, injection near the boundary is clearly shown to have a different convective
316 finger pattern from the injection at the middle of the cell, conducted in ours and several other studies [10,35,52].

317 **4.2. CO₂ introduced to the middle of the cell (Case - A2)**

318 To observe the effect of different injection points, as a slight variation of the first base **Case A1** we changed the
319 injection point for **Case A2** (**Fig. 7**). We observed the CO₂ induction phase followed by a convection period, like the
320 first base case. Furthermore, we also observed similar cell-scale convection during this visualization test.





321
322
323 **Figure 7.** CO₂ introduced to the center of the cell (CF, 1 mm aperture, vertically oriented flat glass) (a) Initial (no CO₂) (b) 9 minutes (c)
324 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h) 154 minutes

325 Changing the CO₂ injection point also affected the layout of the active and inactive zones. For the first base case
326 with CO₂ injection at the extreme left side of the cell, the inactive zone was much farther away at the right side of the
327 cell. However, by injecting CO₂ in the middle of the Hele-Shaw cell, the inactive zones (the left and right sides of the
328 cell) were relatively close to the active zone. The dissolution was symmetrical around the hypothetical vertical central
329 line. At similar experimental conditions, Taheri et al. [10] observed the convection fingers at the side of the boundary
330 cell moving faster than the middle of the cell. They named this phenomenon the 'side boundary effect,' which is not
331 observable in our study. This can be due to the smaller gap width (0.25 mm) or larger Hele-Shaw cell dimensions (500
332 cm × 500 cm) adopted in their study.

333 As seen in **Fig. 6**, the fingers in the inactive zone had slower finger formation than those in the active zone,
334 leading to a high standard deviation with time. However, **Fig. 7** shows that despite the inactive zone having slowed
335 finger formation, it is not as significant as the previous base case **A1**, also evidenced by the smaller standard deviation
336 in finger length compared to the previous test (see **Section 7.1**). This visualization proves that injection points can
337 impact finger evolution. In the field scale, injection nearby vertical no-flow boundaries may impact the convective
338 finger formations, so they should be considered during CO₂ geological storage.

339 Additionally, we observe the finger wavelength to be higher in **Case A3**, as the number of fingers formed is
340 less in the presence of a salt. Furthermore, when comparing **Case A1** and **A2**, the wavelength at the onset of convection
341 is higher when the injection point is at the side (**Case A1**). This difference can be attributed to the fact that during CO₂
342 injection at the side, the fingers at the inactive zone are not formed at the onset time of convection, thus reducing the

343 number of fingers formed, resulting in a higher wavelength. The slower finger growth in the inactive zone for **Case**
 344 **A1** is also observable in **Fig. 4 (b)**. We also considered the wavelength at 9 minutes, as shown in **Table 2**. **By that**
 345 **time**, all the fingers are visible and fully formed throughout the interface. At that period, the wavelength is almost
 346 similar for **Case A1**(4.55 mm) and **A2** (5.41 mm).

347

348 **Table 2.** Comparison and visual representation of the fingers at onset time of convection for Case A

Experiment Case	Solution Used	Picture at Onset Time of Convection ^a	Time (sec) ^b	Wavelength (mm) ^c At Onset of Convection	9 minutes
Case A1	CF		135 ± 15	5.3	4.55
Case A2	CF		135 ± 15	3.56	5.41
Case A3	1 mole CaCl ₂ dissolved in CF		105 ± 15	6.33	6.83

349 ^aThe region under the injection point is cropped and edited for better visibility

350 ^bFor the onset time of convection, we consider the time for which the fingers are clearly visible

351 ^cWavelength is the ratio of the length of interface and the number of fingers for that particular interval

352

353 Onset time of convection, or the time at which the host phase experiences instability in the dissolution-driven
 354 convection flow, is an important aspect of safe CO₂ storage and risk assessment. A shorter onset time of convection is
 355 desired, as convection enhances the mixing of CO₂ in brine, allowing faster CO₂ movement away from the caprock
 356 downwards in the brine [35]. It should be noted that the onset time of convection is defined differently by several
 357 studies [24,28]. This study considers the qualitative onset time of convection, referring to the time point when finger
 358 formation is clearly observable, as provided in **Table 2**, with the corresponding finger wavelength at the onset of
 359 convection. **Table 2** shows that the injection point does not impact the onset time of convection, which is observed to

360 be around 135 seconds (**Case A1 and A2**). However, the convection initiates earlier in the presence of a salt (CaCl₂).
361 This result agrees with Jiang et al. [66], who observed the onset time to be 130 s for pure water, 100 s for 0.25 wt.%
362 saline water, and 20 s for 1.00 wt.% saline water. Several other studies reported the onset of convection as the time
363 when the mass flux deviates from pure diffusion [11,20,22]. Using stability analysis theory, the minimum theoretical
364 onset time of convection, provided by Lindeberg and Wessel-Berg [67], is represented in **Eq. 8**.

$$365 \quad t_c = c1 \left(\frac{\mu \varphi \sqrt{D_{CO_2}}}{\Delta \rho g k} \right)^2 \quad (8)$$

366 For Ra>>40, the horizontal wavelength which most easily causes instability is presented in **Eq. 9**.

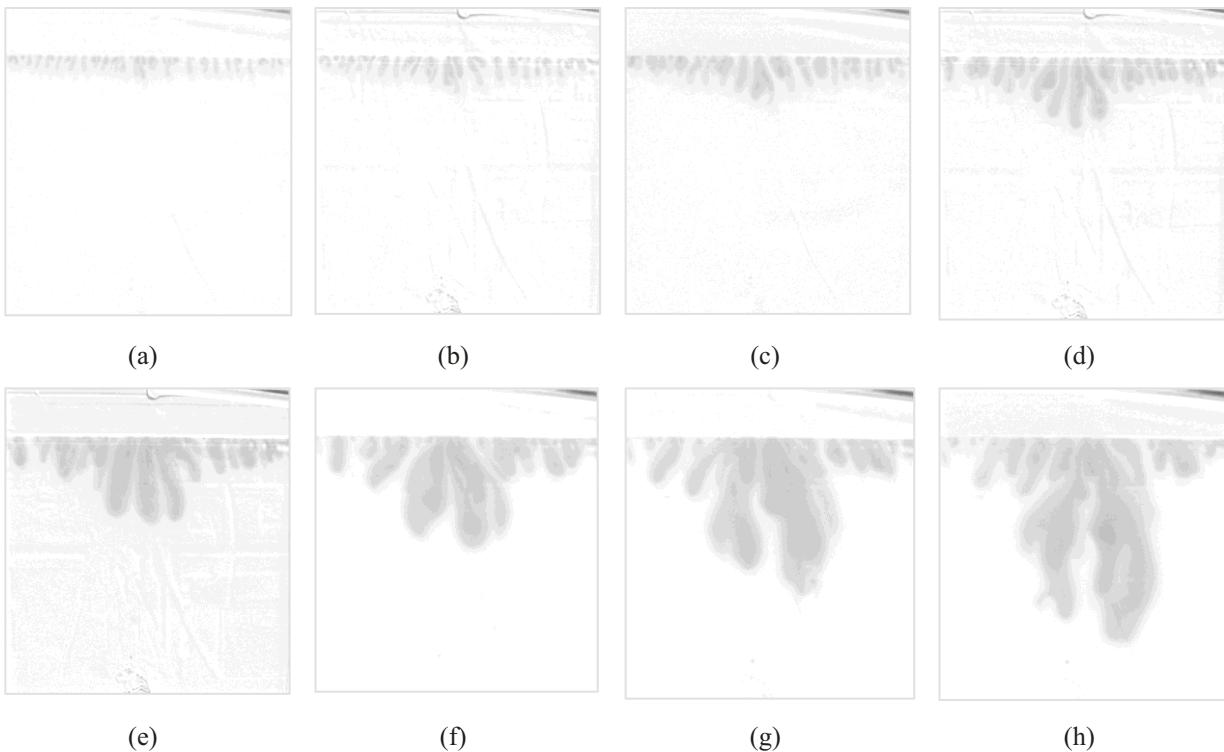
$$367 \quad \lambda_c = c2 \left(\frac{\mu \varphi \pi D_{CO_2}}{\Delta \rho g k} \right) \quad (9)$$

368 Different studies have reported the values of c1 and c2 differently. Based on numerical stability analysis, the minimum
369 theoretical value is reported as c1 = 48.7 and c2 = 96.23 [67]. Using Ra values ranging from 40 to 370 and at 4 MPa,
370 Kneafsey and Pruess experimentally evaluated c1 of 48 to 177 [20]. On the other hand, Pau et al. [28] reported a much
371 higher range of c1 from 1796 to 3670. Similarly, Pruess and Zhang et al. [68] reported a c1 value of 1155.6 and
372 1411.5 in their study. In our study, the c1 value was obtained to be 5770. Furthermore, the value of c2 ranged from
373 195 to 290. Notably, the critical wavelength and onset of convection predicted by linear stability analysis are usually
374 smaller than the experimental value. The discrepancy can be attributed to the fact that the very onset of fingers is
375 differently defined, and the fingers may not be detectable before they are visually observable in experimental studies
376 [11]

377 **5. Qualitative Visualization for Homogeneous Case with Varying Salinity**

378 **5.1. 1 mole CaCl₂ dissolved in Control Fluid – CO₂ introduced to the middle of the cell (Case A3)**

379 The presence of salt stabilizes the convective dissolution of CO₂, as predicted in the theoretical work by Loodts et al.
380 [34] and also shown experimentally in Thomas et al. [35]. By adding salt or increasing salt concentration, the solubility
381 of CO₂ decreases, consequently reducing the density gradient between the fresh solution and the CO₂ mixed solution.
382 Moreover, the system is also stabilized by the increasing viscosity of the solution caused by higher concentrations of
383 salt [34,69]. Since the formation and development of convection-driven flow depends on density-driven gravitational
384 instabilities, the convection process is slowed down in a stable system [34,35,52]. This is also evident from our
385 visualization tests, as seen in **Fig. 8**, and our quantitative measurements (**Section 7**).



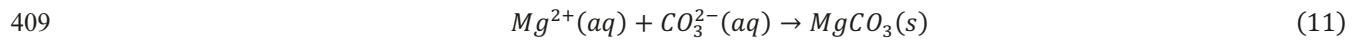
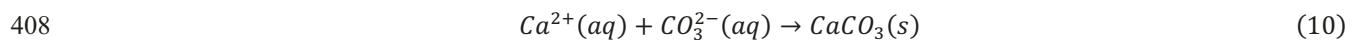
386
387 **Figure 8.** CO₂ introduced to the center of the cell (1 mole CaCl₂ dissolved in CF- 1 mm aperture, vertically oriented flat glass) (a) 9 minutes
388 (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h) 154 minutes

389 Although the symmetry around the hypothetical vertical line is maintained, the presence of CaCl₂ severely
390 hindered the finger growth rate in the inactive zones, resulting in a large standard deviation throughout the test, as
391 seen in **Fig. 8**. Moreover, we also observed a reduction in the number of fingers formed, the maximum being 38 at
392 540 seconds. Unlike the previous base cases where fingers were formed close to one another, the fingers formed in
393 this case were placed at a minimal distance from each other, as observed in **Fig. 8 (b)**. The visualization tests show
394 the importance of knowing the constituents and their composition level for a saline aquifer beforehand to predict the
395 expected CO₂ storage behavior.

396 The pore fluid in saline aquifers contains different compositions of ionic solutions, with common components
397 including Calcium, Magnesium, Potassium, Iron, Chloride, Sulphate, etc., which leads us to investigate the CO₂
398 dissolution-driven convection phenomenon in CaCl₂ and MgCl₂. [61,70]. The presence of different minerals in
399 sandstone formations has a direct effect on CO₂ mineral trapping and also plays a role in altering reservoir transport
400 properties [61,71]. For example, Anorthite (CaAl₂Si₂O₈), a common Ca-bearing feldspar, in sedimentary rock aids in

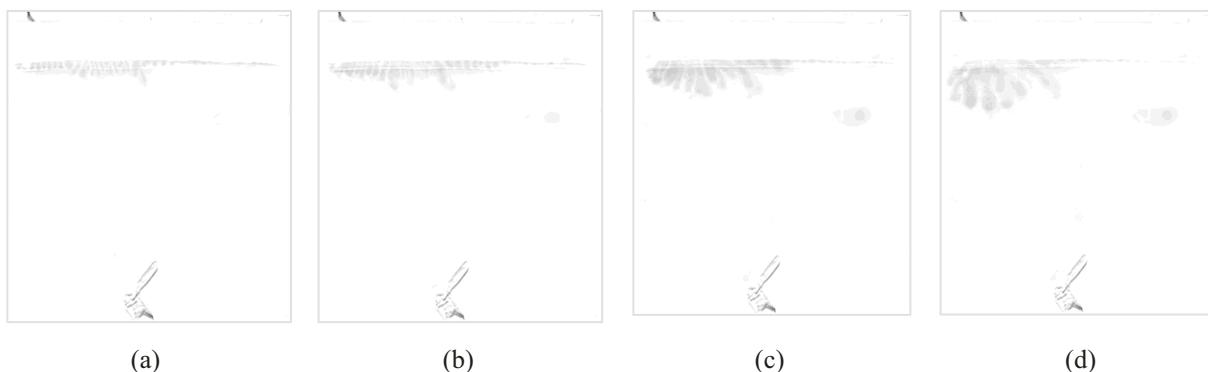
401 faster reaction kinetics and higher CO₂ reactivity, thus enhancing the reservoir rock porosity as well [71]. We refer
402 our readers to the review of Silva et al. [61] for the detailed geochemical aspects of CO₂ sequestration in deep saline
403 aquifers.

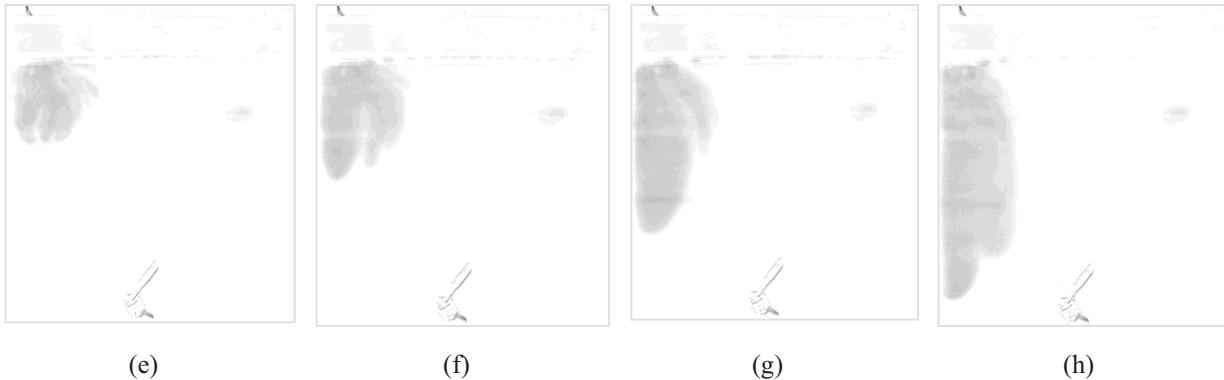
404 As observed by Liu et al. [70], in the presence of salt ions like NaCl, CaCl₂, or MgCl₂ in pore fluid, water
405 molecules bind with the "solvates," leaving less water for CO₂ dissolution. The precipitation reactions for Ca²⁺ and
406 Mg²⁺ ions result in the formation of Calcite (CaCO₃) and Magnesite (MgCO₃), respectively, as shown in **Equations**
407 **10 and 11** [61].



410 Even though most theoretical benchmarks and simulation studies consider ideally horizontal geological
411 storage sites, it is exceedingly rare in real-life scenarios. Most of the storage sites have a formation dip angle, thus
412 making it another key factor of consideration for safely storing CO₂ on subsurface geological sites [37–39]. To
413 visualize the effect of different salts with varying concentrations with a minimal dipping angle, we introduced CO₂ to
414 the left side of the cell. The right side of the Hele-Shaw cell was placed about 1° higher than the left side for the
415 experiments in **Case B**. This allows us to inspect the CO₂ convective dissolution in regions farther from the injection
416 point and the effect of dipping. Furthermore, the qualitative findings from the experiments allow investigation of the
417 changes due to different salts in saline aquifers. We prepare the solution by dissolving salt with different
418 concentrations in the CF for the visualization tests.

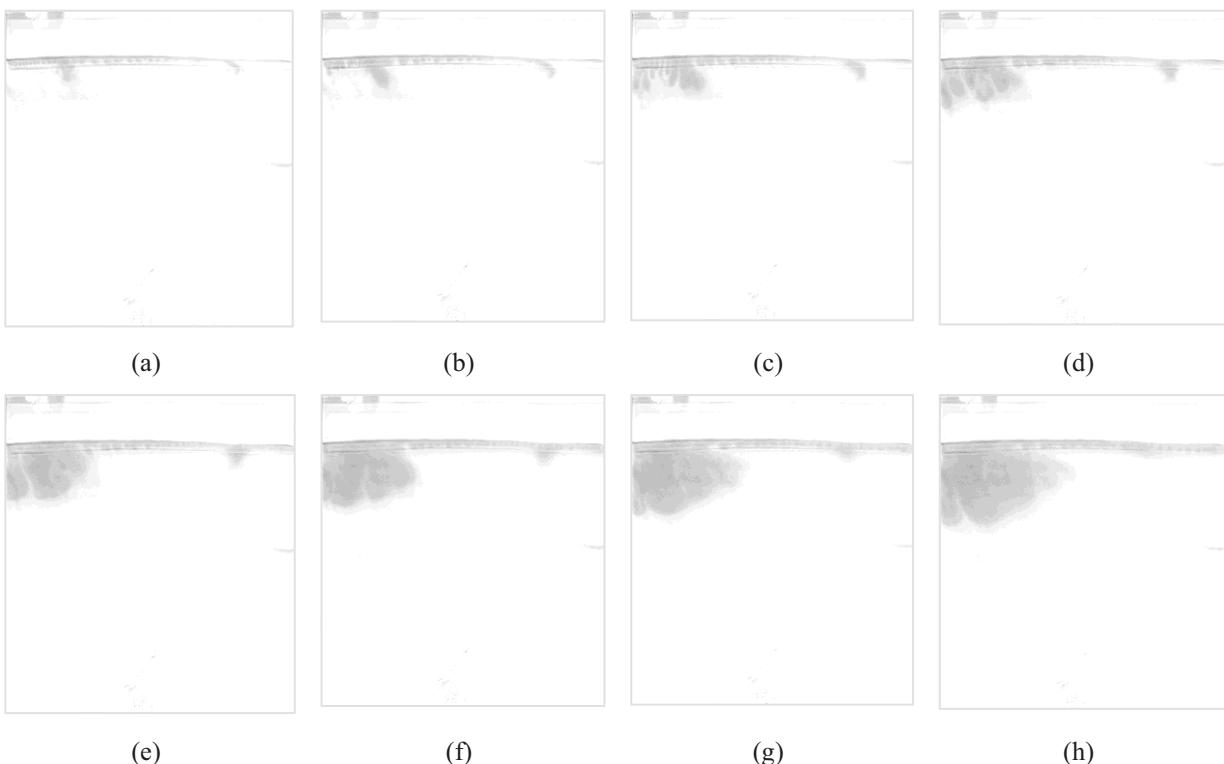
419 **5.2. CaCl₂ dissolved in Base Case Fluid – CO₂ introduced to the side of the cell (Case B1 and Case B2)**





420

421 **Figure 9.** CO₂ introduced to the side of the cell at 0.59 L/min (1 mole CaCl₂ dissolved in CF- 1 mm aperture, vertically oriented flat glass, **Case**
422 **B1)** (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h) 154
423 minutes



424

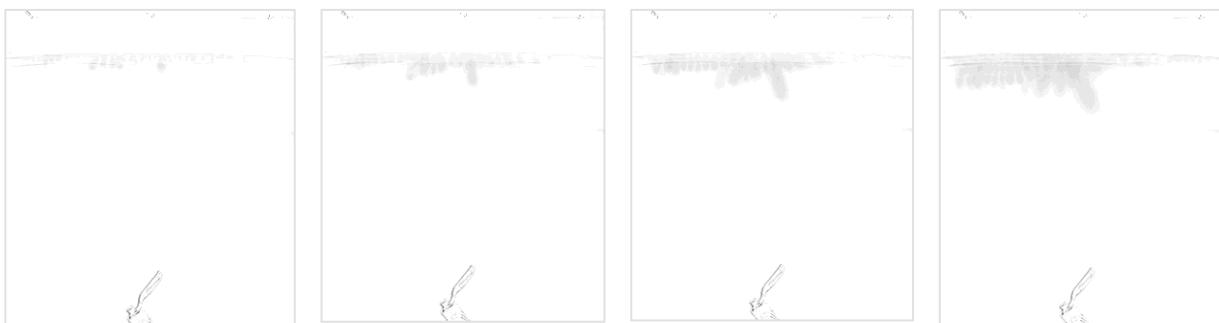
425 **Figure 10.** CO₂ introduced to the side of the cell at 0.59 L/min (2 mole CaCl₂ dissolved in CF- 1 mm aperture, vertically oriented flat glass,
426 **Case B2)** (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h)
427 154 minutes

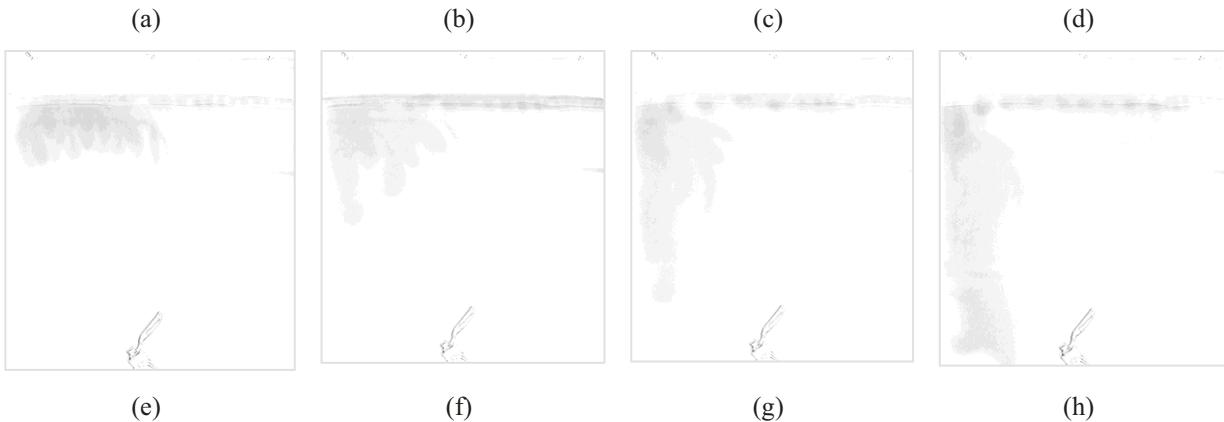
428 The effect of the CO₂ injection point near the boundary and the dipping angle is observable by comparing
429 **Case A3, Case B1, and B2.** For **Case B1 and B2**, the fingers formed near the boundary can't move freely due to
430 restricted space and merge with the nearby fingers, resulting in bigger oval-shaped finger formations, as seen in **Fig.**
431 **9 and 10.** Moreover, due to the dipping angle, with the right side of the cell being placed higher, the finger formation
432 on the right side is severely impacted. The finger formed on the right side moves laterally and mixes with the left side
433 fingers, aiding in the oval-shaped finger formation. The formation of nascent-fingers is observed, and the cycle of
434 lateral finger movement and finger-merging continues throughout the experiment. Whereas, without any boundary
435 effect, as represented by CO₂ injection in the middle (**Case A3**), the coalesced fingers in the active zones protrude in
436 different directions rather than forming a large oval-shaped finger formation.

437 The effect of increased salt concentration is visualized in **Case B2**, which has a higher CaCl₂ concentration,
438 resulting in a lower CO₂ solubility. Moreover, the formation and development of convective-driven flow are slowed
439 in this case, as shown in **Fig. 8.** The vertically downward finger movement is slowed significantly, alluding to the
440 reduced impact of gravity-driven convectional instabilities.

441 **5.3. Effect of MgCl₂ with different concentration (Case B3 and Case B4)**

442 **Fig. 11 and 12** show that the solution with MgCl₂ has higher CO₂ convective flow than CaCl₂, which is
443 evident from a higher pH depressed area (indicating dissolved CO₂) in the pictures. The CO₂ dissolved area indicated
444 by the pH-depressed region follows the order of CaCl₂ 2M solution < MgCl₂ 2M solution < CaCl₂ 1M solution <
445 MgCl₂ 1M solution. Detailed quantitative analysis of the pH depressed zone for the salts are provided in **Section 8.**
446 Besides that, we observe the same effect of dipping angle on CO₂ dissolution in the presence of MgCl₂. Lateral mixing
447 is observed in both lower and higher concentrations of MgCl₂. However, the vertical movement of the finger is
448 considerably slowed down at 2M MgCl₂ solution, as shown in **Fig. 12.**

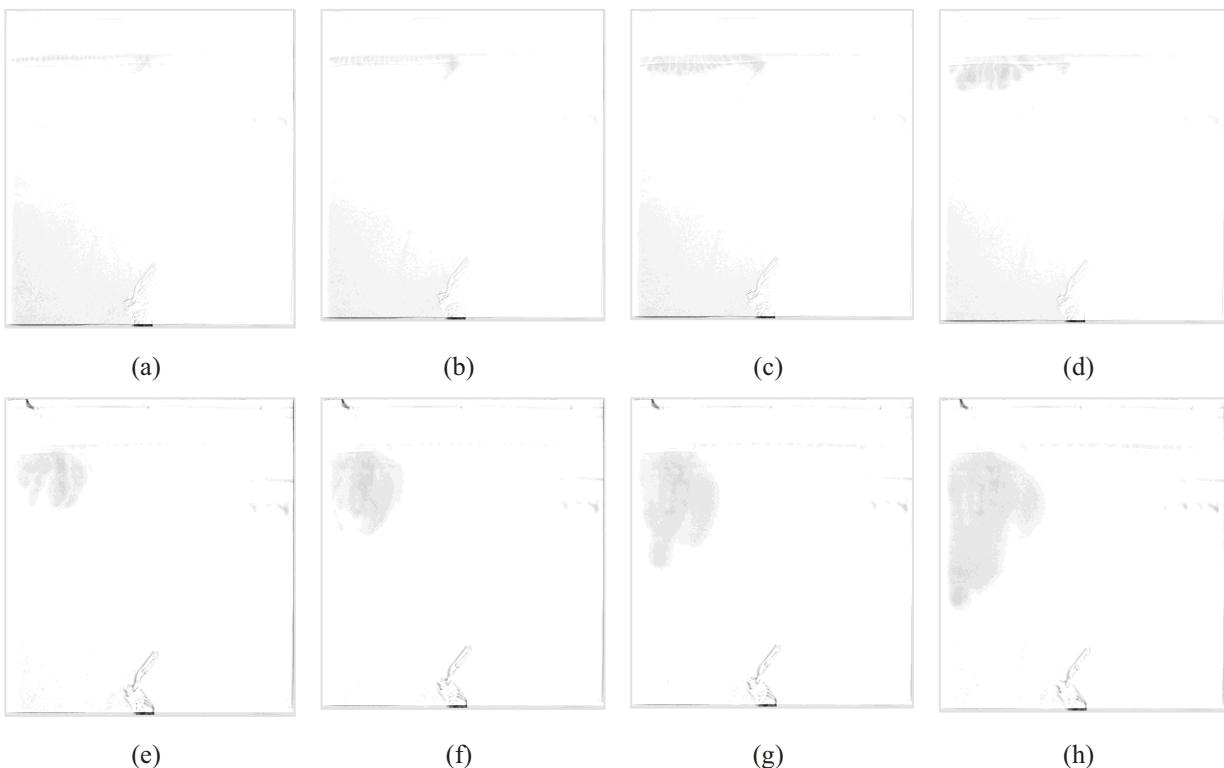




449

450 **Figure 11.** CO₂ introduced to the side of the cell at 0.59 L/min (1 mole MgCl₂ dissolved in CF- 1 mm aperture, vertically oriented flat glass,
451 **Case B3)** (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h)
452 154 minutes

453



454

455 **Figure 12.** CO₂ introduced to the side of the cell at 0.59 L/min (2 mole MgCl₂ dissolved in CF- 1 mm aperture, vertically oriented flat glass,
456 **Case B4)**(a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h)
457 154 minutes

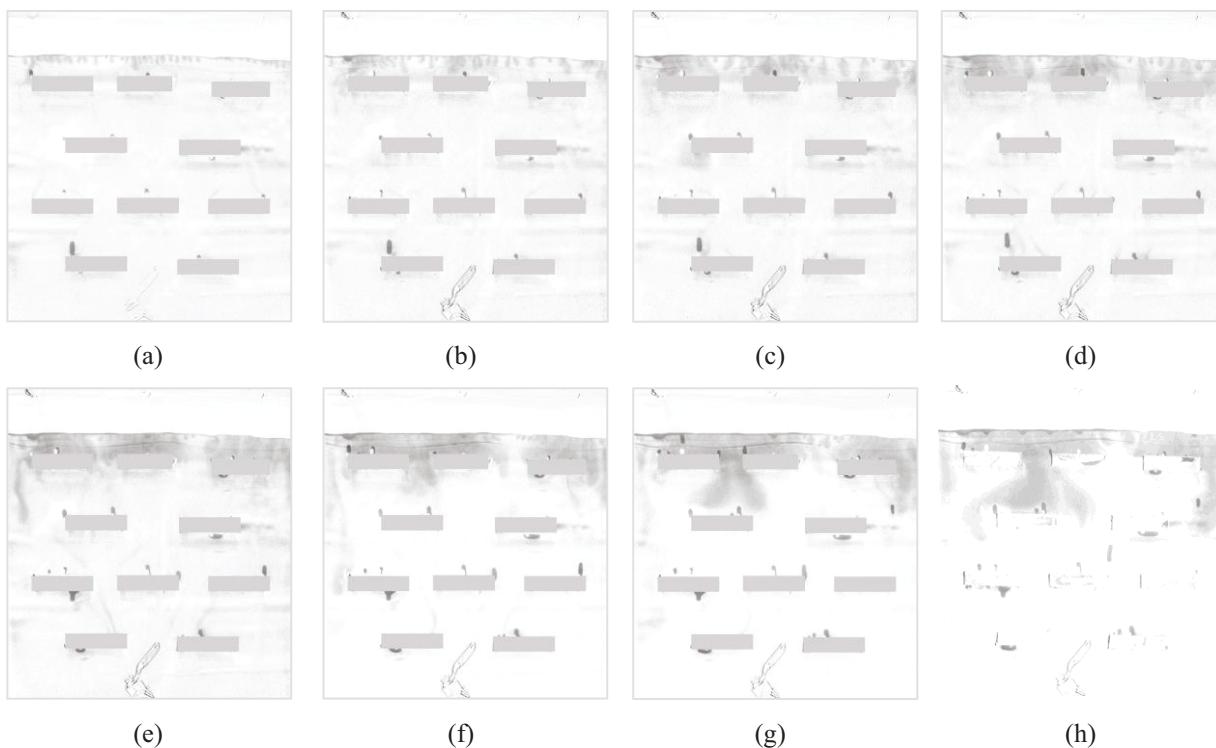
458 **6. Qualitative Visualization for Effect of Permeability Heterogeneity (Case C and D)**

459 All the storage sites have some form of vertical and horizontal permeability heterogeneities introduced by calcite
460 layers or shales, which play a major role in dictating fluid flow [10,64]. Different convective-flow patterns (gravity
461 fingering, channeling, or dispersion) can be observed based on the heterogeneity medium. Depending on the
462 heterogeneity parameters, the average CO₂ mass flux at the top boundary can reach different constant values after
463 decreasing at the initial stage and then increasing, as shown by Ranganathan et al. [72].

464 In this section, we investigate the effect of heterogeneities on convective behavior by setting up multiple 0.05
465 mm thick graphite strips (99.5% graphite flexible sheet) inside the Hele-Shaw cell, thus reducing the aperture to 0.05
466 mm for these regions. The length and height of these graphite strips are 50 mm and 12.7 mm, respectively.
467 Furthermore, to prevent any airborne hydrocarbon contamination, the cell was cleaned thoroughly using distilled
468 water, and the graphite strips were replaced with new ones prior to conducting the experiments. The graphite strips
469 were selected for ease of removal and adhesion to the cell walls. Although silicone shims were used for setting up the
470 Hele-Shaw, they merely act as a boundary and do not affect the flow properties inside the cell. Therefore, the surface
471 properties (e.g., wettability, interfacial tension, etc.) remain the same for all the experiments where the effect of
472 heterogeneity is evaluated. We have highlighted the graphite strips in the figures to visualize them better. The
473 permeability for the region with graphite strip was $2.083 \times 10^{-8} \text{ m}^2$, which was calculated by assuming that the
474 subsection behaves like an independent miniature Hele-Shaw cell with an aperture size of 0.05 mm by using **Eq. 1**.
475 Additionally, to determine the effective permeability of the Hele-Shaw cell, we consider the permeability of both the
476 regions with and without the graphite strips and calculate the weighted average based on the area. The updated porosity
477 is calculated as the ratio of available volume inside the Hele-Shaw to the total volume and is found to be 0.948. The
478 effective permeability for the medium with heterogeneities is calculated as $7.68 \times 10^{-8} \text{ m}^2$, with a corresponding
479 Rayleigh number of 40663, which is lower than the permeability and Ra number obtained for our homogeneous cases
480 by 7.72% and 2.68%, respectively. Although the permeability calculated using this method may not strictly represent
481 the actual permeability due to the associated boundary effects, the overall trend remains valid.

482 **6.1. Effect of Heterogeneity (Case C1, C2, C3, and C4)**

483 **Fig. 13** displays the CO₂ dissolution behavior in the presence of heterogeneous barriers. The downward velocity of
484 the fingers is slowed down along with decreased CO₂ dissolution, as evidenced by the smaller pH-depressed region
485 compared to the homogenous case. Furthermore, as shown in **Fig. 13 (e-h)**, upon passing the first set of heterogeneous
486 layers, rather than traveling downwards vertically, the CO₂ finger moves in a curved manner towards the side of the
487 heterogeneous strips in the second layer. The curved finger travel alludes to a preferential CO₂ movement path, which
488 can vary on the geometry of the heterogeneity. Heterogeneity also promotes flow channeling, resulting in a large
489 region that CO₂ does not encounter.



490 **Figure 13.** CO₂ introduced to the side of the cell at 0.59 L/min (CF, 1 mm aperture, vertically oriented flat glass with heterogeneous layers,
491 **Case C1**) (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes (h)
492 154 minutes

493

494 **Fig. 14, 15, and 16** represent the CO₂ dissolution behavior in the presence of 1M CaCl₂ solution, 1M MgCl₂ solution,
495 and 1M NaCl solution, respectively. For NaCl and CaCl₂, as the convection is slower than CF, the lateral merging of
496 CO₂ fingers is more dominant once it encounters the heterogeneous barrier. However, as shown in **Fig. 15**, for MgCl₂,

497 the vertical movement is more dominant compared to both NaCl and CaCl₂. This prevalent vertical finger travel for
498 1M MgCl₂ solutions can be attributed to faster CO₂ dissolution in the MgCl₂ solution than in NaCl and CaCl₂ solutions.
499 The more rapid dissolution allows the fingers to quickly pass the heterogeneous barriers, leaving less time for lateral
500 mixing.

501

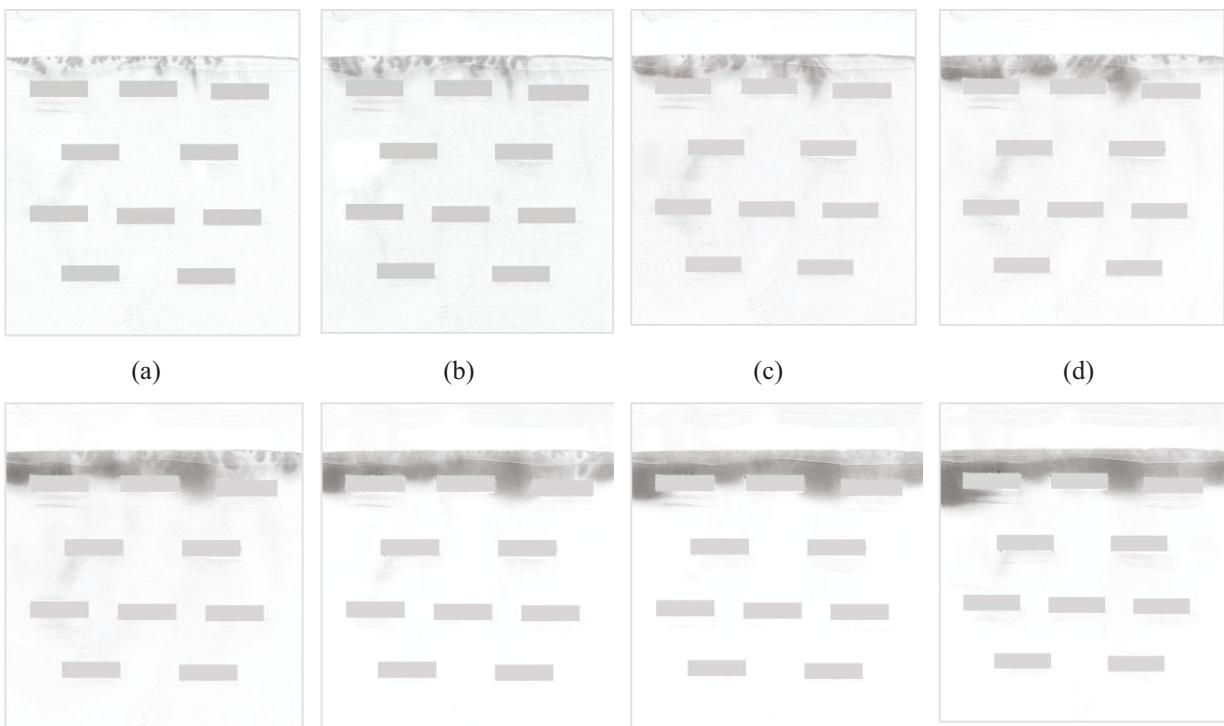


502 **Figure 14.** CO₂ introduced to the side of the cell at 0.59 L/min (1 mole CaCl₂ dissolved CF- 1 mm aperture, vertically oriented flat glass with
503 heterogeneous layers, **Case C2**) (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes
504 (g) 118 minutes (h) 154 minutes

505



506 **Figure 15.** CO₂ introduced to the side of the cell at 0.59 L/min (1 mole MgCl₂ dissolved CF- 1 mm aperture, vertically oriented flat glass with
507 heterogeneous layers, **Case C3**) (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes
508 (g) 118 minutes (h) 154 minutes



(e)

(f)

(g)

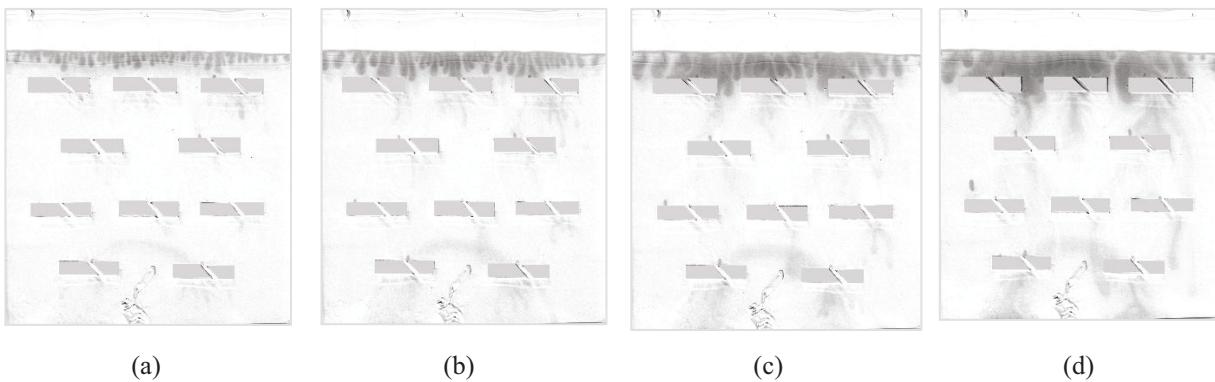
(h)

509 **Figure 16.** CO₂ introduced to the side of the cell at 0.59 L/min (1 mole NaCl dissolved CF- 1 mm aperture, vertically oriented flat glass with
510 heterogeneous layers, **Case C4**) (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes
511 (g) 118 minutes (h) 154 minutes

512 **6.2. Effect of Fractures in Heterogeneous Layers (Case D1)**

513 Information about natural fractures and their interaction with CO₂-acidified brine in the subsurface can play a crucial
514 role in assessing the CO₂ storage parameters. Compared to the surrounding rock matrix, natural fractures in the
515 geothermal reservoirs have a higher permeability region [73]. Fractures also have high conductivity, and the presence
516 of an interconnected fracture system can provide potential CO₂ escape routes during storage by promoting the
517 increased spatial spreading of CO₂ plume, thus allowing CO₂ migration pathway through the cap rocks to neighboring
518 aquifers or surfaces [45,74].

519 Despite the recognized importance of heterogeneous layers, the impact of fractures present in these heterogeneous
520 formations on the convective-flow pattern has not been visually studied extensively [40,41]. Therefore, we study the
521 effect of fractures in heterogeneous regions by creating thin slits (4 mm spacing at 45° angle) between the graphite
522 shims, as shown in **Fig. 17**. **Fig. 17** shows that the CO₂ spatial spreading is much faster due to the fractures compared
523 to cases without fractures in heterogeneity (**Case C1-4**). This can be attributed to an increase in effective permeability
524 in the system. The regions underneath the top fracture layer show the point of the CO₂ dissolution path, as new CO₂
525 fingers are visible underneath them in **Fig. 17 (b)**. Alluding to the channeling effect, a preferential dissolution path is
526 clearly visible around an hour into the experiment shown in **Fig. 17 (e)**.

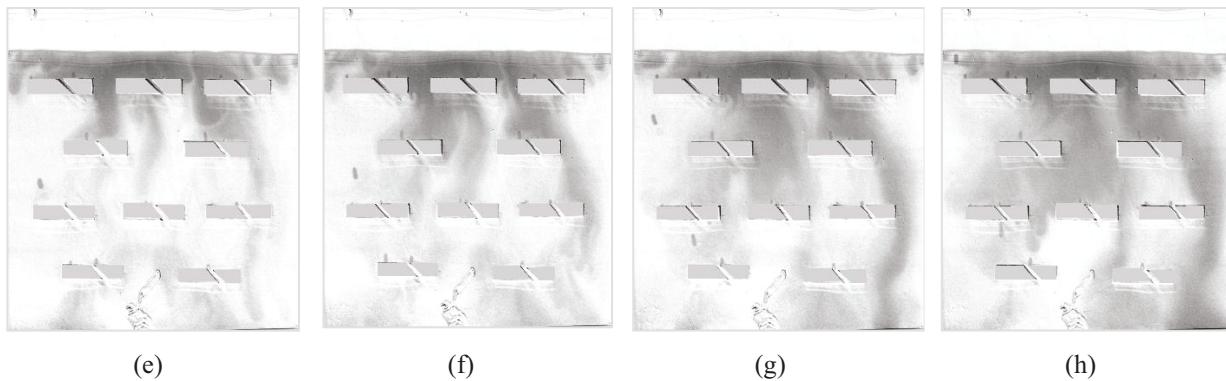


(a)

(b)

(c)

(d)

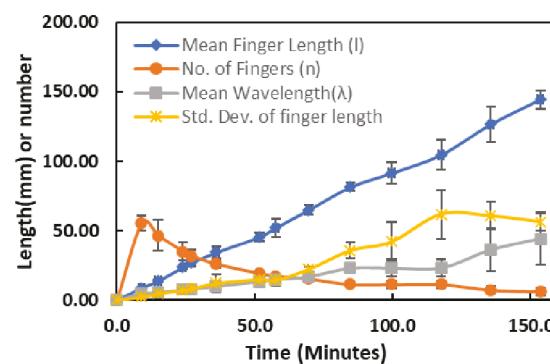


527 **Figure 17.** CO₂ introduced to the side of the cell at 0.59 L/min (CF- 1 mm aperture, vertically oriented flat glass with fractured heterogeneous
528 layers, **Case D1**) (a) 9 minutes (b) 15 minutes (c) 24 minutes (d) 36 minutes (e) 58 minutes (f) 85 minutes (g) 118 minutes
529 (h) 154 minutes

530 7. Quantitative measurements

531 The number of fingers formed, average finger length (mm), average wavelength (mm), and standard deviation of
532 finger length are the parameters that characterize the finger evolution dynamics. Finger length is calculated as the
533 vertical distance from the interface to the tip of the finger. In contrast, the average wavelength is the interface length
534 divided by the number of fingers. It should be noted that fingers with different tips at the end moving in different
535 directions are considered separate fingers in this study.

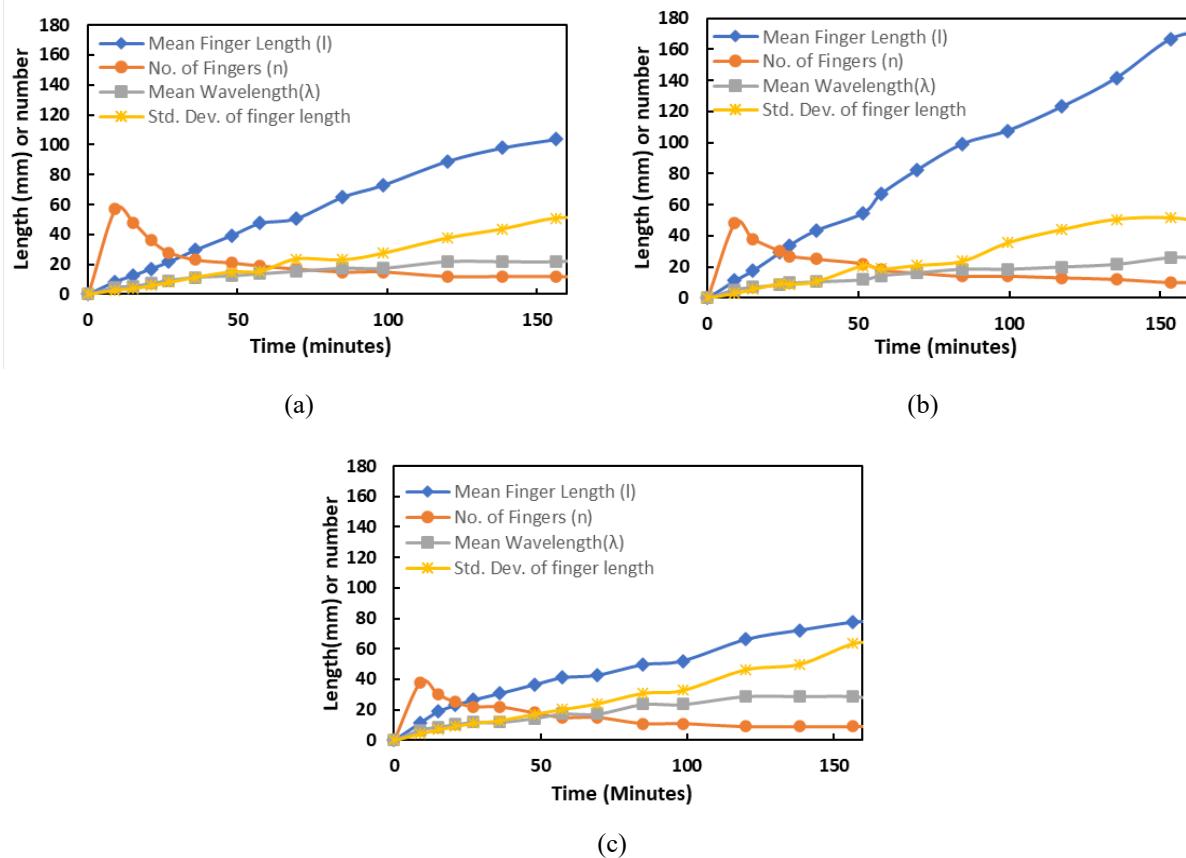
536 Before presenting the results of our experimental investigations, we perform a repeatability test for the
537 quantitative data (refer to **Fig. 4**). As shown in **Fig. 18**, for the three runs, we calculate the average of the finger
538 parameters and consider the deviation with a 95% confidence interval. Although we observed a slight difference in
539 the number of fingers formed initially, the deviation
540 becomes negligible in the middle and later parts of the
541 experiment. We also observed a deviation from the mean
542 finger length value in the latter part of the experiment. This
543 deviation in the early part of the experiment can be due to
544 the uncertain nature of convective finger merging and flow
545 patterns as the finger grows. The acceptable level of
546 deviation observed in **Fig. 18** shows the repeatability of our
547 experiments.



548 **Figure 18.** Quantitative data of the finger length parameters for
549 the repeatability experiments.

548 **7.1. Effect of injection point and presence of salt**

549 The image analysis for **Case A1-A3 (Fig. 19 a and b)** reveals average finger length having a similar trend, which
550 denotes that the average finger length formation is independent of the injection point. Moreover, the number of fingers
551 formed remains similar, regardless of the CO₂ injection point. However, the standard deviation of finger length is
552 higher when CO₂ is injected into the left side of the cell. This is because the fingers underneath the CO₂ injection point
553 quickly grew in size while the area on the right side observed less CO₂ dissolution, causing a higher imbalance in
554 finger size.



555

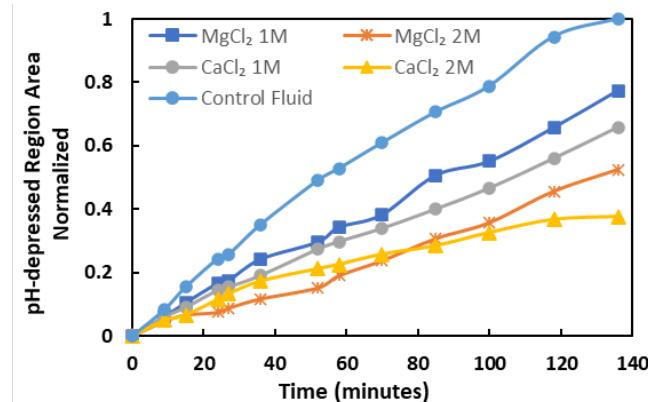
556 **Figure 19.** Number of fingers formed, wavelength, and finger length parameters for (a) Case A1 - CO₂ introduced to the side of the cell and
557 (b) Case A2 - CO₂ introduced to the center of the cell (c) Case A3- 1 mole CaCl₂ dissolved in CF - CO₂ introduced to the center of the cell

558 The effect of salinity on finger evolution can be observed in **Fig. 19 (c)**. The effect of convection is attenuated
559 due to the presence of CaCl₂, resulting in less finger interaction than in the control fluid. This is evident by the reduced
560 number of fingers formed. Similar effects of salt were observed in other studies [34–36]. A higher standard deviation

561 of finger length is also observed due to the significant difference in finger size between the active zone (area under
562 the CO₂ injection point) and the inactive zones (area at the side of the Hele-Shaw cells).

563 **7.2. Effect of different salts with varying concentrations**

564 To illustrate the effect of the presence of different salts with varying concentrations, the pH-depressed region for **Case**
565 **B1-B4** is calculated and compared with the CO₂ dissolution in the CF. **Fig. 20** shows the results normalized with
566 respect to the pH-depressed area for the CF. CO₂ is injected at the side of the Hele-Shaw cell for all the experiments,
567 to control the effect of different injection points on the convective-dissolution pattern. It should be noted that a higher
568 pH-depressed area refers to more CO₂ dissolution. As seen in **Fig. 20**, higher concentration leads to slower CO₂
569 dissolution for the same salt type, which is in perfect agreement with other studies [34,35,52]. The CO₂ dissolved area
570 indicated by the pH-depressed region follows the order of CaCl₂ 2M solution < MgCl₂ 2M solution < CaCl₂ 1M
571 solution < MgCl₂ 1M solution. At the beginning part of the experiment, all the salt types have almost similar
572 dissolution areas. It can be due to the impact of minimal convection-driven flow at the beginning of the experiment.
573 Therefore, it can be safe to assume that the effect of salinity is not significant throughout the induction phase and for
574 a small period after the onset of convection.



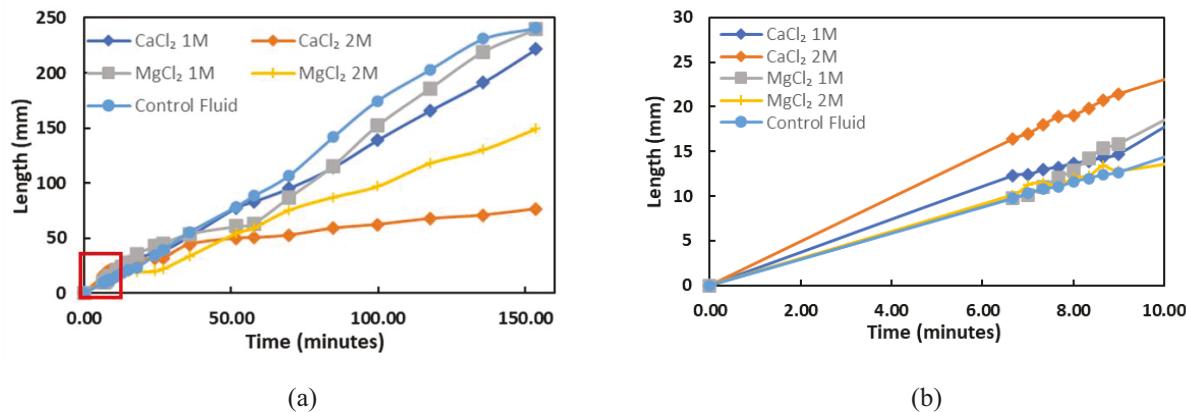
575

576 **Figure 20.** Area of the pH-depressed region (normalized) for different salts with varying concentration

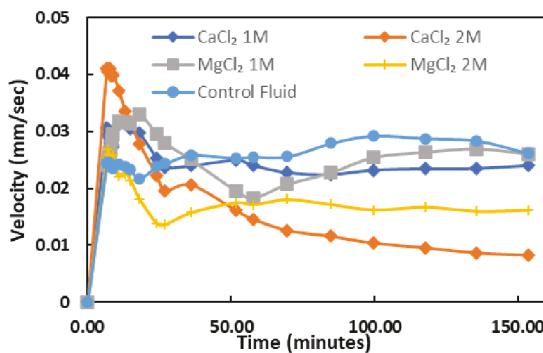
577 The downward progression of fingers (also referred to as 'mixing length' by Thomas et al. [35]) is another factor
578 worthy of investigation as it relates to the depth of investigation during CO₂ storage. To calculate the downward
579 progression of the fingers, we consider the distance from the gas-water interface to the tip of the longest finger for any
580 time, presented in **Fig. 21 (a)**. The vertical depth of investigation is quite similar for the CF, CaCl₂ 1M, and MgCl₂ 1M
581 solutions. However, at higher concentrations (CaCl₂ 2M solution and MgCl₂ 2M solution), we observe a significant

582 decrease in the vertical progression of fingers. Using NaCl with varying concentrations, Thomas et al. observed a
 583 similar effect [33]. Increasing salinity slows the convective flow, causing less CO₂ to dissolve, consequently reducing
 584 the vertical depth of investigation. It should also be noted that dipping can reduce the vertical depth of investigation,
 585 as it promotes lateral mixing over vertical travel.

586 Another interesting phenomenon can be observed while considering the earlier part of the experiments,
 587 indicated by the red squared box in **Fig. 21 (a)**, also magnified as presented in **Fig. 21 (b)**. As shown in **Fig. 21 (b)**,
 588 although the higher salt concentration in solution leads to a less vertical progression of fingers, the temporal finger
 589 evolution is higher at the earlier stage, also clearly observable from the velocity of finger evolution shown in **Fig. 22**.
 590 This indicates that the presence of salt enhances the diffusive flux which is the dominant mass-transfer method at an
 591 earlier stage. However, as time progresses, convection flows become the predominant flow causing the finger
 592 progression rate to slow. This experimental investigation also perfectly agrees with Kim and Kim's observation from
 593 the numerical simulation [36].



594 **Figure 21.** (a) Temporal evolution of the vertical progression of fingers for different cases (b) Magnified view of the red squared region

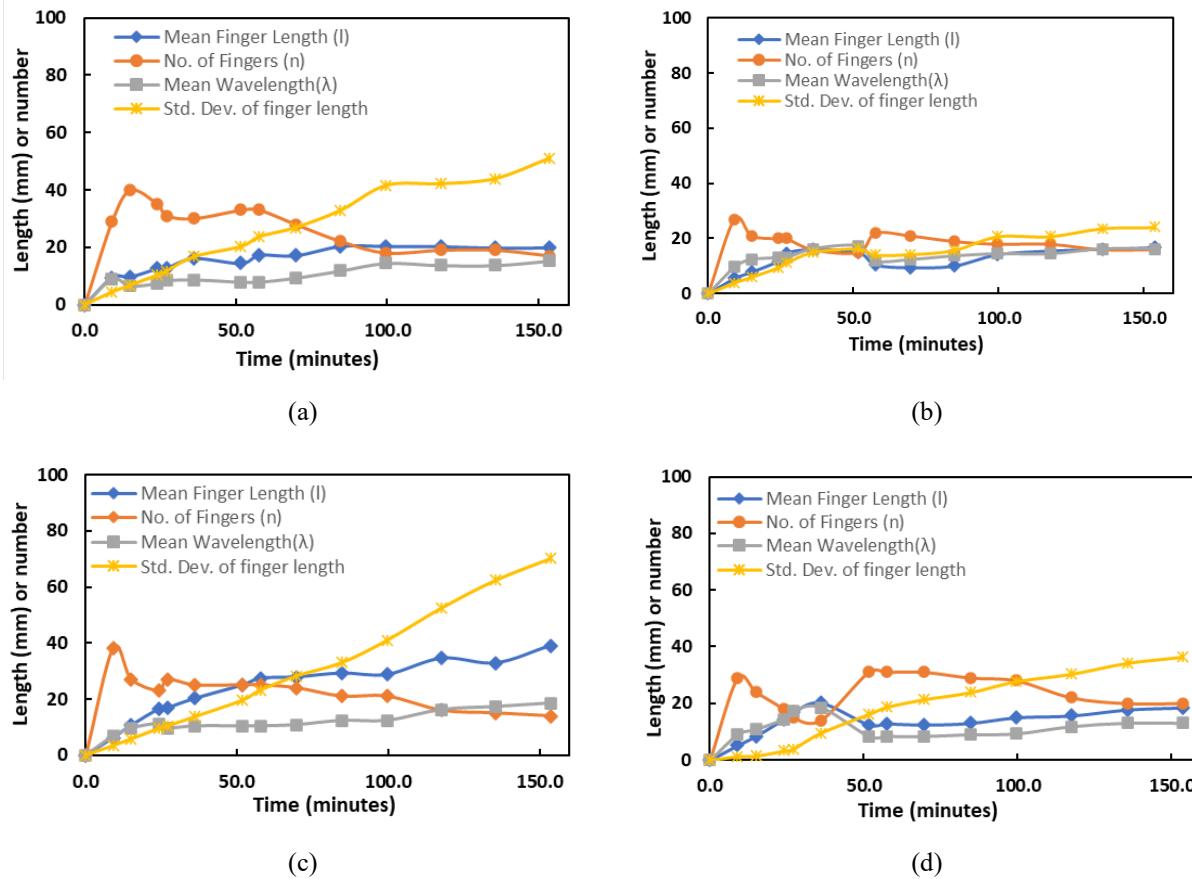


595

596 **Figure 22.** The velocity of vertical finger evolution for different cases
 597 **Fig. 23** presents the image analysis result for the salts with different concentrations. We observe a slow increase in
 598 the mean finger length at the beginning of the experiment, and the number of fingers decreases as the neighboring
 599 fingers merge. However, after around 36 minutes, the nascent fingers with short finger lengths start to appear at the
 600 gas-water interface, which reduces the mean finger length. Moreover, the mature fingers start to merge, resulting in
 601 fewer dominant fingers, further decreasing the mean finger length, which is the weighted average of all the fingers.
 602 However, it should be noted that despite our observation of a reduction in the mean finger length, there is still vertical
 603 progression of fingers throughout the experiment, as evident in **Fig. 21**. This trend continues throughout the latter part
 604 of the experiment, as the larger finger merges and forms the oval shape formation, there is a noticeable decrease in
 605 the mean finger length, also causing a sharp increase in the standard deviation. Moreover, as the contributing factor
 606 in determining mean finger length is the vertical finger travel, the lateral mixing of fingers due to the dipping angle
 607 also causes a reduction in the mean finger length.

608 The convective finger formation and dissolution rate depend significantly on the type and concentration of
609 salt in the CF. **Fig 23 (b)** shows that the number of fingers formed is considerably slowed in case of a higher
610 concentration of CaCl_2 . There is a sharp increase in the number of fingers formed between 52 and 58 minutes since
611 the fingers on the right side became apparent in that period. This shows that the CO_2 dissolution rate on the right side
612 is significantly lower than on the left. Both the injection point and the presence of a dipping angle can cause this
613 variance in dissolution rate.

Fig. 23 shows that the number of fingers formed for $MgCl_2$ is almost similar to that of $CaCl_2$. In the latter half of the experiment, we observed the fingers in $MgCl_2$ having a higher mean finger length, particularly due to the faster vertical travel of the finger on the left side of the cell. Moreover, since the finger formation in the right side of the cell is significantly slow in all the cases, we observe a higher standard deviation in the later part of the experiment for the $MgCl_2$ experiments. Furthermore, the lower standard deviation of fingers in both the 2M solution cases can be attributed to the higher salinity and dipping angle preventing vertical travel and promoting lateral mixing of the fingers.



620 **Figure 23.** Number of fingers formed, wavelength and finger length parameters for (a) 1 mole CaCl₂ dissolved in Control Fluid – CO₂
 621 introduced to the side of the cell and (b) 2 mole CaCl₂ dissolved in Control Fluid – CO₂ introduced to the side of the cell (c) 1 mole MgCl₂
 622 dissolved in Control Fluid – CO₂ introduced to the side of the cell (d) 2 mole MgCl₂ dissolved in Control Fluid – CO₂ introduced to the side of
 623 the cell

624 8. Discussion

625 The structural morphology (formation dipping angle, pore-networks), geological properties (rock porosity,
 626 permeability, presence of fractures or barriers, etc.), and the accurate geochemical composition of the storage site play
 627 a significant role in transporting and storing CO₂ in the brine in dissolved form. Therefore, a fundamental
 628 understanding of these parameters and their role in convective CO₂ transport is essential to optimize the existing CO₂
 629 storage operations and select appropriate new CO₂ storage sites. In addition, understanding the horizontal migration
 630 of supercritical CO₂ plumes due to structural dipping and point of injection is essential to reduce the chance of CO₂
 631 leakage from fractured caprocks. The empirical observations of the parameters made in our study are summarized and

632 presented in **Table 3**. This study critically evaluates these parameters to provide some important insights, which are
633 discussed below:

634 **8.1. Effect of salt concentration and dipping angle on convective dissolution**

635 Our study showed that the presence of salts could significantly slow the convective flow, and the convective flow rate
636 can depend on the salt type and its concentration. The CO₂ dissolution pattern was also noticeably different in the
637 presence of salt, moving in a piston-like-displacement manner. Although the convective flow is slowed down, we
638 observed enhanced diffusive flux, the dominant mass-transfer method at the earlier stage of CO₂ storage. Comparing
639 the convective CO₂ flow pattern without and with salts with different concentrations, we can conclude that lower
640 salinity storage sites are more favorable for safe CO₂ storage, despite benefits like faster dissolution and earlier onset
641 time of convection.

642 **Table 3.** Empirical observations on the parameters observed

Parameters Observed	Empirical Observations
CO ₂ Injection Point	<ul style="list-style-type: none">Injection point near boundary has different convective flow pattern than injection at the middle
Salinity	<ul style="list-style-type: none">Enhances the diffusive fluxSlows down the convective flowReduces CO₂ dissolution rate with increasing salinity
Dipping	<ul style="list-style-type: none">Promotes lateral mixing compare to vertical finger travelReduces CO₂ storage with increasing dipping angle (however, the effect is not as significant as salinity)
Fractures	<ul style="list-style-type: none">Promotes high spatial CO₂ plume spreadingHighly fractured regions have a high risk of CO₂ leakage during storage
Heterogeneity	<ul style="list-style-type: none">Different convection flow patterns are possible based on heterogeneous patterns

	• Uneven CO ₂ sweep with large unaffected regions
--	--

643

644 Even though most of the experimental studies and simulations generalize geological storage sites to be horizontal,
645 the majority of the sites have gradients as an effect of diagenesis and other geological phenomena. Although the effect
646 of dip angle is less significant than that of salinity, the presence of the formation dip can significantly impact the
647 space migration of CO₂ mixed fluid flow, as observed in our study. Due to the dominance of lateral mixing of CO₂
648 fingers over vertical travel, where dipping is involved, the depth of investigation is reduced for storage sites with a
649 formation angle. A larger formation dip angle is not conducive to CO₂ geological storage as it promotes spatial CO₂
650 migration, risking long-term CO₂ storage. Further experimental studies on the effect of dipping angles on the depth of
651 investigation can be done at a field-scale level for better theoretical benchmarking.

652 **8.2. Effect of flow barriers (boundary, heterogeneities, and fractures) on convective dissolution**

653 In our study, the injection point near the boundary had a considerably different convective flow pattern than the
654 injection in the middle. From this visual investigation, we can infer that the vertical faults present at the storage site
655 can cause a change in the convective flow pattern. Additionally, our characterization of convective flow for
656 heterogeneities with fractures reveals how the CO₂ plume can migrate/spread through the fractures. Since the fractured
657 regions can be highly conductive and increase CO₂ spatial spreading, fractures are not conducive to CO₂ storage. In
658 addition, the formation of preferential channels for CO₂-acidified fluid flow is also observed. However, it should be
659 noted that the channel growth and preferential fluid flow are more complex in storage sites and thus needs more
660 attention.

661 Furthermore, the convective-flow pattern and the effective permeability of the storage site can vary
662 significantly based on the fracture patterns. Therefore, the fraction patterns used to derive correlations will only apply
663 to reservoirs with a similar pattern, showing the importance of considering all possible fracture network characteristics
664 to critically estimate and model the complex fracture geometry for a particular storage site.

665

666

667

668 **8.3. Effect of Higher Pressure and Temperature**

669 CO₂ storage sites have higher pressure and temperature than atmospheric conditions, leading to several
670 studies focusing on the effects of high pressure and high temperature (HPHT) on CO₂ dissolution. Outeda et al. [16]
671 investigated the impact of varied pressure (1.5-5 bar) on CO₂ dissolution at a constant temperature of 25 °C. Increasing
672 pressure resulted in high CO₂ mass influx to the aqueous phase, increasing mixing length. Furthermore, instabilities
673 are noticed earlier at higher pressure. At the lowest pressure of 1.5 bar, the instabilities were observable at around 60
674 s, whereas, for 5 bar, the instabilities could be seen as early as 10 s. However, factors like the wavelength of fingers
675 or shape were independent of pressure change. Another CO₂ dissolution study at HPHT condition (140-240 bar and
676 20-150 °C) by Tang et al. [4] showed that increasing pressure resulted in a higher Rayleigh number. Moreover, the
677 study concluded that pressure has a larger influence on CO₂ solubility than the diffusion coefficient at reservoir HPHT
678 condition.

679 Despite considerable CO₂ dissolution visualization studies being performed, most do not consider a porous
680 medium. Therefore Amarasinghe et al. [53] investigated CO₂ dissolution in a porous media using 400-600 μm glass
681 beads at 105 bar and 50 °C. The results were compared to Vosper et al. [75], who used similar glass bead sizes but
682 different atmospheric conditions (1 bar and room temperature). A comparison of the results revealed that the CO₂
683 transport rate is considerably higher at higher pressure and temperature (HPHT) conditions (reaching a depth of 40
684 mm in 20 minutes) compared to atmospheric conditions (reaching 40 mm in 100 minutes). This higher transport rate
685 is also attributed to CO₂ being supercritical at the HPHT condition. Amarasinghe et al. [13] also investigated CO₂
686 transport in heterogeneous porous media by packing the Hele-Shaw cell with glass beads of different permeability at
687 100 bar and 50 °C.

688 The effect of brine composition on CO₂ dissolution at 34.5-36.9 bar and 50 °C was investigated in the works
689 of Mahmoodpour et al. [62]. They observed a late onset time of convection and lower finger growth rate in solutions
690 containing NaCl and CaCl₂ at HPHT conditions, similar to our experimental observations at atmospheric conditions.
691 Furthermore, a study by Jiang et al. [14] at 38 °C and 90-110 bar reached similar findings, showing that CO₂ solubility
692 is higher at high pressure, leading to faster CO₂ solution accumulation at the top interface. Furthermore, under
693 supercritical pressure, the density difference of the CO₂ solution is considerably higher, thus promoting finger
694 formation and development. In contrast, salinity seemed to inhibit finger formation and migration.

695 The findings from the CO₂ dissolution studies under HPHT conditions show that higher pressure and
696 temperature are favorable for faster CO₂ transport. However, since there is no significant change in finger morphology
697 due to variation in pressure and temperature, it is safe to assume that the fundamental aspects of reactive-convective
698 transport in atmospheric conditions found in this study will also hold in high-pressure and temperature conditions.
699 Future studies will consider a more complicated interpretation involving quantitative analyses with the effect of
700 different salt types and concentrations at high pressure and supercritical conditions.

701 **9. Conclusion**

702 This study provided a visual investigation into the effects of different parameters that dictates the CO₂ geologic
703 storage: injection point and pressure, presence of salts with varying concentration, and presence of heterogeneities
704 and dipping angle using the Hele-Shaw cell. Moreover, quantitative results were provided, which can be used for
705 theoretical modeling. Our results show favorable conditions for CO₂ geological storage, including none or low
706 reservoir dipping. Furthermore, the absence of salinity is preferable; however, since all geological storage aquifer
707 fluids will have some dissolved salts, selections should be made based on reservoirs with the lowest salinity. We
708 observed the qualitative onset time of convection at around 135 s with a corresponding finger wavelength of 3.56-5.3
709 mm. In the presence of salt, the convection initiated at about 105 s with a corresponding finger wavelength of 6.33
710 mm. This difference shows that the initial diffusion is higher in brine with high salt concentrations but does not
711 translate to enhanced convection and dissolution. The CO₂ dissolved area for different salts, indicated by the pH-
712 depressed region, are considerably smaller than when no salt is present, and it follows the order of CaCl₂ 2M solution
713 < MgCl₂ 2M solution < CaCl₂ 1M solution < MgCl₂ 1M solution. A factor not considered in this study but worthy of
714 further investigation is the synergistic effect of multiple salts in saline aquifers which can be studied by creating
715 solutions with a chemical composition similar to the brine by mixing several salts. By identifying the critical factors
716 controlling the convective mixing, this study also provides new insights into the possible line of future work on the
717 transport mechanisms during dissolution-driven CO₂ convective flow.

718 **Acknowledgements:**

719 We would like to thank Mr. Edward Farina from the UT Tyler Machine Shop for his significant help with the material
720 selection, fabrication, and troubleshooting of the apparatus used in this study. We thank the anonymous reviewers for
721 their edits and suggestions, which improved our manuscript. This material is based upon work partly supported by the

722 National Science Foundation Award under CBET-2245484. Any opinions, findings, and conclusions or
723 recommendations expressed in this material are those of the author(s) and do not necessarily reflect the views of the
724 National Science Foundation.

725 **Table of Abbreviations**

2D	Two Dimensional
3D	Three Dimensional
BCG	Bromocresol Green
CaCl ₂	Calcium Chloride
CCS	Carbon Capture and Storage
CF	Control Fluid
CO ₂	Carbon Dioxide
HPHT	High Pressure and High Temperature
KMnO ₄	Potassium Permanganate
LIF	Laser-Induced Fluorescence
NaCl	Sodium Chloride
PIV	Particle Image Velocimetry
PVT	Pressure-Volume-Temperature
Ra	Rayleigh number
RC	Rayleigh convection

RGB Red-Green-Blue

sCO₂ Supercritical CO₂

UV Ultraviolet

726

727 References

728 [1] M.F. Shahriar, A. Khanal, The current techno-economic, environmental, policy status and perspectives of
729 sustainable aviation fuel (SAF), *Fuel.* 325 (2022) 124905. <https://doi.org/10.1016/j.fuel.2022.124905>.

730 [2] Z. Zhang, Q. Fu, H. Zhang, X. Yuan, K.T. Yu, Experimental and Numerical Investigation on Interfacial Mass
731 Transfer Mechanism for Rayleigh Convection in Hele-Shaw Cell, *Ind. Eng. Chem. Res.* 59 (2020) 10195–
732 10209. <https://doi.org/10.1021/acs.iecr.0c01345>.

733 [3] A. Khanal, M.F. Shahriar, Physics-Based Proxy Modeling of CO₂ Sequestration in Deep Saline Aquifers,
734 *Energies.* 15 (2022) 4350. <https://doi.org/10.3390/en15124350>.

735 [4] Y. Tang, Z. Li, R. Wang, M. Cui, X. Wang, Z. Lun, Y. Lu, Experimental Study on the Density-Driven Carbon
736 Dioxide Convective Diffusion in Formation Water at Reservoir Conditions, *ACS Omega.* 4 (2019) 11082–
737 11092. <https://doi.org/10.1021/acsomega.9b00627>.

738 [5] R. Khosrokhavar, G. Elsinga, R. Farajzadeh, H. Bruining, Visualization and investigation of natural
739 convection flow of CO₂ in aqueous and oleic systems, *J. Pet. Sci. Eng.* 122 (2014) 230–239.
740 <https://doi.org/10.1016/j.petrol.2014.07.016>.

741 [6] R. Khosrokhavar, A. Eftekhari, R. Farajzdadeh, K.H. Wolf, H. Bruining, Effect of Salinity and Pressure on
742 the Rate of Mass Transfer in Aquifer Storage of CO₂, *IOR 2015 - 18th Eur. Symp. Improv. Oil Recover.*
743 (2015). <https://doi.org/10.3997/2214-4609.201412114>.

744 [7] H. Emami-Meybodi, Stability analysis of dissolution-driven convection in porous media, *Phys. Fluids.* 29
745 (2017). <https://doi.org/10.1063/1.4974275>.

746 [8] C.A. Dorao, Dispersion phenomena in gas liquid systems, *J. Nat. Gas Sci. Eng.* 5 (2012) 25–30.
747 <https://doi.org/10.1016/j.jngse.2011.08.004>.

748 [9] M.R. Soltanian, M.A. Amooie, Z. Dai, D. Cole, J. Moortgat, Critical Dynamics of Gravito-Convective Mixing
749 in Geological Carbon Sequestration, *Sci. Rep.* 6 (2016) 35921. <https://doi.org/10.1038/srep35921>.

750 [10] A. Taheri, O. Torsæter, E. Lindeberg, N.J. Hadia, D. Wessel-Berg, Qualitative and quantitative experimental
751 study of convective mixing process during storage of CO₂ in heterogeneous saline aquifers, *Int. J. Greenh.*

752 Gas Control. 71 (2018) 212–226. <https://doi.org/10.1016/j.ijgge.2018.02.003>.

753 [11] A. Taheri, O. Torsæter, D. Wessel-Berg, M. Soroush, Experimental and Simulation Studies of Density-
754 Driven-Convection Mixing in a Hele-Shaw Geometry with Application for CO₂ Sequestration in Brine
755 Aquifers, in: All Days, SPE, 2012. <https://doi.org/10.2118/154908-MS>.

756 [12] W. Amarasinghe, I. Fjelde, Y. Guo, CO₂ dissolution and convection in oil at realistic reservoir conditions: A
757 visualization study, J. Nat. Gas Sci. Eng. 95 (2021). <https://doi.org/10.1016/j.jngse.2021.104113>.

758 [13] W.S. Amarasinghe, I. Fjelde, A.M.N. Flaata, Visual investigation of CO₂ dissolution and convection in
759 heterogeneous porous media at reservoir temperature and pressure conditions, Greenh. Gases Sci. Technol.
760 11 (2021) 342–359. <https://doi.org/10.1002/ghg.2055>.

761 [14] L. Jiang, Y. Wang, G. Lu, J. Yang, Y. Song, Experimental Study on the Density-Driven Convective Mixing
762 of CO₂ and Brine at Reservoir Temperature and Pressure Conditions, Energy and Fuels. (2022).
763 <https://doi.org/10.1021/acs.energyfuels.2c00054>.

764 [15] B. An, D. Solorzano, Q. Yuan, Viscous Fingering Dynamics and Flow Regimes of Miscible Displacements
765 in a Sealed Hele-Shaw Cell, Energies. 15 (2022) 5798. <https://doi.org/10.3390/en15165798>.

766 [16] R. Outeda, C. El Hasi, A. D’Onofrio, A. Zalts, Experimental study of linear and nonlinear regimes of density-
767 driven instabilities induced by CO₂ dissolution in water, Chaos. 24 (2014) 013135.
768 <https://doi.org/10.1063/1.4868040>.

769 [17] A. Singh, Y. Singh, K.M. Pandey, Viscous fingering instabilities in radial Hele-Shaw cell: A review, Mater.
770 Today Proc. 26 (2019) 760–762. <https://doi.org/10.1016/j.matpr.2020.01.022>.

771 [18] M. De Paoli, M. Alipour, A. Soldati, How non-Darcy effects influence scaling laws in Hele-Shaw convection
772 experiments, J. Fluid Mech. (2020). <https://doi.org/10.1017/jfm.2020.229>.

773 [19] M.C. Kim, Linear and nonlinear analyses on the onset of gravitational instabilities in a fluid saturated within
774 a vertical Hele-Shaw cell, Chem. Eng. Sci. 126 (2015) 349–360. <https://doi.org/10.1016/j.ces.2014.12.032>.

775 [20] T.J. Kneafsey, K. Pruess, Laboratory experiments and numerical simulation studies of convectively enhanced
776 carbon dioxide dissolution, Energy Procedia. 4 (2011) 5114–5121.

777 https://doi.org/10.1016/j.egypro.2011.02.487.

778 [21] S. Backhaus, K. Turitsyn, R.E. Ecke, Convective instability and mass transport of diffusion layers in a Hele-
779 Shaw geometry, *Phys. Rev. Lett.* 106 (2011) 1–4. <https://doi.org/10.1103/PhysRevLett.106.104501>.

780 [22] A.C. Slim, M.M. Bandi, J.C. Miller, L. Mahadevan, Dissolution-driven convection in a Hele-Shaw cell, *Phys.*
781 *Fluids.* 25 (2013). <https://doi.org/10.1063/1.4790511>.

782 [23] S. Mojtaba, R. Behzad, N.M. Rasoul, R. Mohammad, Experimental study of density-driven convection effects
783 on CO₂ dissolution rate in formation water for geological storage, *J. Nat. Gas Sci. Eng.* 21 (2014) 600–607.
784 <https://doi.org/10.1016/j.jngse.2014.09.020>.

785 [24] T.F. Faisal, S. Chevalier, Y. Bernabe, R. Juanes, M. Sassi, Quantitative and qualitative study of density driven
786 CO₂ mass transfer in a vertical Hele-Shaw cell, *Int. J. Heat Mass Transf.* 81 (2015) 901–914.
787 <https://doi.org/10.1016/j.ijheatmasstransfer.2014.11.017>.

788 [25] H. Wang, F. Torabi, F. Zeng, H. Xiao, Experimental and numerical study of non-equilibrium dissolution and
789 exsolution behavior of CO₂ in a heavy oil system utilizing Hele-Shaw-like visual cell, *Fuel.* 270 (2020).
790 <https://doi.org/10.1016/j.fuel.2020.117501>.

791 [26] M.C. Kim, C. Wylock, Linear and nonlinear analyses of the effect of chemical reaction on the onset of
792 buoyancy-driven instability in a CO₂ absorption process in a porous medium or Hele-Shaw cell, *Can. J.*
793 *Chem. Eng.* 95 (2017) 589–604. <https://doi.org/10.1002/cjce.22694>.

794 [27] S. Mahmoodpour, B. Rostami, M.R. Soltanian, M.A. Amooie, Convective Dissolution of Carbon Dioxide in
795 Deep Saline Aquifers: Insights from Engineering a High-Pressure Porous Visual Cell, *Phys. Rev. Appl.* 12
796 (2019). <https://doi.org/10.1103/PhysRevApplied.12.034016>.

797 [28] G.S.H. Pau, J.B. Bell, K. Pruess, A.S. Almgren, M.J. Lijewski, K. Zhang, High-resolution simulation and
798 characterization of density-driven flow in CO₂ storage in saline aquifers, *Adv. Water Resour.* 33 (2010) 443–
799 455. <https://doi.org/10.1016/j.advwatres.2010.01.009>.

800 [29] A.C. Slim, Solutal-convection regimes in a two-dimensional porous medium, *J. Fluid Mech.* 741 (2014) 461–
801 491. <https://doi.org/10.1017/jfm.2013.673>.

802 [30] C. Wylock, A. Rednikov, B. Haut, P. Colinet, Nonmonotonic Rayleigh-Taylor Instabilities Driven by Gas–
803 Liquid CO₂ Chemisorption, *J. Phys. Chem. B.* 118 (2014) 11323–11329. <https://doi.org/10.1021/jp5070038>.

804 [31] X. Cao, M. Zhang, A.S. Mujumdar, Q. Zhong, Z. Wang, Effects of ultrasonic pretreatments on quality, energy
805 consumption and sterilization of barley grass in freeze drying, 2018.
806 <https://doi.org/10.1016/j.ulstsonch.2017.06.014>.

807 [32] Z. Zhang, H. Zhang, X. Yuan, K.T. Yu, Effective UV-Induced Fluorescence Method for Investigating
808 Interphase Mass Transfer of Single Bubble Rising in the Hele-Shaw Cell, *Ind. Eng. Chem. Res.* 59 (2020)
809 6729–6740. <https://doi.org/10.1021/acs.iecr.9b07106>.

810 [33] C. Thomas, V. Loodts, L. Rongy, A. De Wit, Convective dissolution of CO₂ in reactive alkaline solutions:
811 Active role of spectator ions, *Int. J. Greenh. Gas Control.* 53 (2016) 230–242.
812 <https://doi.org/10.1016/j.ijggc.2016.07.034>.

813 [34] V. Loodts, L. Rongy, A. De Wit, Impact of pressure, salt concentration, and temperature on the convective
814 dissolution of carbon dioxide in aqueous solutions, *Chaos.* 24 (2014). <https://doi.org/10.1063/1.4896974>.

815 [35] C. Thomas, S. Dehaeck, A. De Wit, Convective dissolution of CO₂ in water and salt solutions, *Int. J. Greenh.*
816 *Gas Control.* 72 (2018) 105–116. <https://doi.org/10.1016/j.ijggc.2018.01.019>.

817 [36] M. Chan Kim, Y. Hwan Kim, The effect of chemical reaction on the onset of gravitational instabilities in a
818 fluid saturated within a vertical Hele-Shaw cell: Theoretical and numerical studies, *Chem. Eng. Sci.* 134
819 (2015) 632–647. <https://doi.org/10.1016/j.ces.2015.05.066>.

820 [37] J. Jing, Y.L. Yang, Z.H. Tang, Effects of formation dip angle and salinity on the safety of CO₂ geological
821 storage - A case study of Shiqianfeng strata with low porosity and low permeability in the Ordos Basin, China,
822 *J. Clean. Prod.* 226 (2019) 874–891. <https://doi.org/10.1016/j.jclepro.2019.04.038>.

823 [38] F. Wang, J. Jing, T. Xu, Y. Yang, G. Jin, Impacts of stratum dip angle on CO₂ geological storage amount and
824 security, *Greenh. Gases Sci. Technol.* 6 (2016) 682–694. <https://doi.org/10.1002/ghg.1594>.

825 [39] A. Kumar, M. Noh, G.A. Pope, K. Sepehrnoori, S. Bryant, L.W. Lake, Reservoir simulation of CO₂ storage
826 in deep saline aquifers, *SPE - DOE Improv. Oil Recover. Symp. Proc.* (2004). <https://doi.org/10.2118/89343>.

827 ms.

828 [40] A. Khanal, R. Weijermars, Comparison of Flow Solutions for Naturally Fractured Reservoirs Using Complex
829 Analysis Methods (CAM) and Embedded Discrete Fracture Models (EDFM): Fundamental Design
830 Differences and Improved Scaling Method, *Geofluids*. 2020 (2020). <https://doi.org/10.1155/2020/8838540>.

831 [41] A. Khanal, R. Weijermars, Distinguishing Fracture Conductivity and Fracture Flux: A Systematic
832 Investigation of Individual Fracture Contribution to Well Productivity, in: *Proc. 8th Unconv. Resour. Technol.*
833 Conf., American Association of Petroleum Geologists, Tulsa, OK, USA, 2020. <https://doi.org/10.15530/urtec-2020-2176>.

835 [42] R. Weijermars, A. Khanal, Flow in fractured porous media modeled in closed-form: Augmentation of prior
836 solution and side-stepping inconvenient branch cut locations, *Fluids*. 5 (2020).
837 <https://doi.org/10.3390/fluids5020051>.

838 [43] R. Weijermars, A. Khanal, High-resolution streamline models of flow in fractured porous media using discrete
839 fractures: Implications for upscaling of permeability anisotropy, *Earth-Science Rev.* 194 (2019) 399–448.
840 <https://doi.org/10.1016/j.earscirev.2019.03.011>.

841 [44] C.E. Bond, R. Wightman, P.S. Ringrose, The influence of fracture anisotropy on CO₂ flow, *Geophys. Res.*
842 Lett. 40 (2013) 1284–1289. <https://doi.org/10.1002/grl.50313>.

843 [45] M. Iding, P. Ringrose, Evaluating the impact of fractures on the performance of the In Salah CO₂ storage site,
844 *Int. J. Greenh. Gas Control.* 4 (2010) 242–248. <https://doi.org/10.1016/j.ijggc.2009.10.016>.

845 [46] A. V. Gorin, On Analogy Between Convective Heat and Mass Transfer Processes in a Porous Medium and a
846 Hele-Shaw Cell, in: 2012: pp. 139–159. https://doi.org/10.1007/978-3-642-21966-5_6.

847 [47] F. Graf, E. Meiburg, C. Härtel, Density-driven instabilities of miscible fluids in a Hele-Shaw cell: Linear
848 stability analysis of the three-dimensional Stokes equations, *J. Fluid Mech.* 451 (2002) 261–282.
849 <https://doi.org/10.1017/s0022112001006516>.

850 [48] L. Paterson, Fingering with miscible fluids in a Hele Shaw cell, *Phys. Fluids.* 28 (1985) 26–30.
851 <https://doi.org/10.1063/1.865195>.

852 [49] R.L. Detwiler, S.E. Pringle, R.J. Glass, Measurement of fracture aperture fields using transmitted light: An
853 evaluation of measurement errors and their influence on simulations of flow and transport through a single
854 fracture, *Water Resour. Res.* 35 (1999) 2605–2617. <https://doi.org/10.1029/1999WR900164>.

855 [50] A. Khanal, M.F. Shahriar, Optimization of CO₂ Huff-n-Puff in Unconventional Reservoirs with a Focus on
856 Pore Confinement Effects, Fluid Types, and Completion Parameters, *Energies.* 16 (2023) 2311.
857 <https://doi.org/10.3390/en16052311>.

858 [51] G. Lu, Y. Liu, L. Jiang, T. Ying, Y. Song, B. Wu, Study of Density Driven Convection in a Hele-Shaw Cell
859 with Application to the Carbon Sequestration in Aquifers, *Energy Procedia.* 114 (2017) 4303–4312.
860 <https://doi.org/10.1016/j.egypro.2017.03.1576>.

861 [52] T.J. Kneafsey, K. Pruess, Laboratory flow experiments for visualizing carbon dioxide-induced, density-driven
862 brine convection, *Transp. Porous Media.* 82 (2010) 123–139. <https://doi.org/10.1007/s11242-009-9482-2>.

863 [53] W. Amarasinghe, I. Fjelde, J.Å. Rydland, Y. Guo, Effects of permeability on CO₂ dissolution and convection
864 at reservoir temperature and pressure conditions: A visualization study, *Int. J. Greenh. Gas Control.* 99 (2020).
865 <https://doi.org/10.1016/j.ijggc.2020.103082>.

866 [54] C. Thomas, L. Lemaigre, A. Zalts, A. D’Onofrio, A. De Wit, Experimental study of CO₂ convective
867 dissolution: The effect of color indicators, *Int. J. Greenh. Gas Control.* 42 (2015) 525–533.
868 <https://doi.org/10.1016/j.ijggc.2015.09.002>.

869 [55] A. Taheri, E. Lindeberg, O. Torsæter, D. Wessel-Berg, Qualitative and quantitative experimental study of
870 convective mixing process during storage of CO₂ in homogeneous saline aquifers, *Int. J. Greenh. Gas Control.*
871 66 (2017) 159–176. <https://doi.org/10.1016/j.ijggc.2017.08.023>.

872 [56] F. Liu, P. Lu, C. Zhu, Y. Xiao, Coupled reactive flow and transport modeling of CO₂ sequestration in the Mt.
873 Simon sandstone formation, Midwest U.S.A., *Int. J. Greenh. Gas Control.* 5 (2011) 294–307.
874 <https://doi.org/10.1016/j.ijggc.2010.08.008>.

875 [57] Z. Li, M. Dong, S. Li, L. Dai, Densities and Solubilities for Binary Systems of Carbon Dioxide + Water and
876 Carbon Dioxide + Brine at 59 °C and Pressures to 29 MPa, *J. Chem. Eng. Data.* 49 (2004) 1026–1031.

877 https://doi.org/10.1021/je049945c.

878 [58] R.J. Rosenbauer, T. Koksalan, J.L. Palandri, Experimental investigation of CO₂-brine-rock interactions at
879 elevated temperature and pressure: Implications for CO₂ sequestration in deep-saline aquifers, Fuel Process.
880 Technol. 86 (2005) 1581–1597. <https://doi.org/10.1016/j.fuproc.2005.01.011>.

881 [59] B. Zerai, B.Z. Saylor, G. Matisoff, Computer simulation of CO₂ trapped through mineral precipitation in the
882 Rose Run Sandstone, Ohio, Appl. Geochemistry. 21 (2006) 223–240.
883 <https://doi.org/10.1016/j.apgeochem.2005.11.002>.

884 [60] P. Ranganathan, P. Van Hemert, E.S.J. Rudolph, P.Z.J. Zitha, Numerical modeling of CO₂ mineralisation
885 during storage in deep saline aquifers, Energy Procedia. 4 (2011) 4538–4545.
886 <https://doi.org/10.1016/j.egypro.2011.02.411>.

887 [61] G.P.D. De Silva, P.G. Ranjith, M.S.A. Perera, Geochemical aspects of CO₂ sequestration in deep saline
888 aquifers: A review, Fuel. 155 (2015) 128–143. <https://doi.org/10.1016/j.fuel.2015.03.045>.

889 [62] S. Mahmoodpour, B. Rostami, M.R. Soltanian, M.A. Amooie, Effect of brine composition on the onset of
890 convection during CO₂ dissolution in brine, Comput. Geosci. 124 (2019) 1–13.
891 <https://doi.org/10.1016/j.cageo.2018.12.002>.

892 [63] T. Xu, Y.-C.C. Jin, Simulation the convective mixing of CO₂ in geological formations with a meshless model,
893 Chem. Eng. Sci. 192 (2018) 187–198. <https://doi.org/10.1016/j.ces.2018.07.016>.

894 [64] A. Taheri, O. Torsæter, E. Lindeberg, N.J. Hadia, D. Wessel-Berg, Effect of Convective Mixing Process on
895 Storage of CO₂ in Saline Aquifers with Layered Permeability, Adv. Chem. Res. 3 (2021).
896 <https://doi.org/10.21926/acr.2101012>.

897 [65] V. Loodts, C. Thomas, L. Rongy, A. De Wit, Control of Convective Dissolution by Chemical Reactions:
898 General Classification and Application to CO_{2} Dissolution in Reactive Aqueous
899 Solutions, Phys. Rev. Lett. 113 (2014) 114501. <https://doi.org/10.1103/PhysRevLett.113.114501>.

901 [66] L. Jiang, S. Wang, D. Liu, W. Zhang, G. Lu, Y. Liu, J. Zhao, Change in convection mixing properties with

902 salinity and temperature: CO₂ storage application, *Polymers* (Basel). 12 (2020).
903 <https://doi.org/10.3390/POLYM12092084>.

904 [67] E. Lindeberg, D. Wessel-Berg, Upscaling studies of diffusion induced convection in homogeneous and
905 heterogeneous aquifers, *Energy Procedia*. 4 (2011) 3927–3934. <https://doi.org/10.1016/j.egypro.2011.02.331>.

906 [68] K. Pruess, K. Zhang, Numerical Modeling Studies of The Dissolution-Diffusion-Convection Process During
907 CO₂ Storage in Saline Aquifers, Lawrence Berkeley Natl. Lab. (2008) 1–33.
908 <https://publications.lbl.gov/10.2172/7D178882-C5E3-4A78-8B79-633C227DDED7>.

909 [69] H. Qiblawey, M. Arshad, A. Easa, M. Atilhan, Viscosity and density of ternary solution of calcium chloride
910 + sodium chloride + water from T = (293.15 to 323.15) K, *J. Chem. Eng. Data*. 59 (2014) 2133–2143.
911 <https://doi.org/10.1021/je500070k>.

912 [70] B. Liu, B.S. Mahmood, E. Mohammadian, A.K. Manshad, N.R. Rosli, M. Ostadhassan, Measurement of
913 solubility of CO₂ in NaCl, CaCl₂, MgCl₂ and MgCl₂ + CaCl₂ brines at temperatures from 298 to 373 K and
914 pressures up to 20 mpa using the potentiometric titration method, *Energies*. 14 (2021).
915 <https://doi.org/10.3390/en14217222>.

916 [71] M. Sorai, T. Ohsumi, M. Ishikawa, K. Tsukamoto, Feldspar dissolution rates measured using phase-shift
917 interferometry: Implications to CO₂ underground sequestration, *Appl. Geochemistry*. 22 (2007) 2795–2809.
918 <https://doi.org/10.1016/j.apgeochem.2007.06.019>.

919 [72] P. Ranganathan, R. Farajzadeh, H. Bruining, P.L.J.J. Zitha, Numerical Simulation of Natural Convection in
920 Heterogeneous Porous media for CO₂ Geological Storage, *Transp. Porous Media*. 95 (2012) 25–54.
921 <https://doi.org/10.1007/s11242-012-0031-z>.

922 [73] A. Jafari, T. Babadagli, Effective fracture network permeability of geothermal reservoirs, *Geothermics*. 40
923 (2011) 25–38. <https://doi.org/10.1016/j.geothermics.2010.10.003>.

924 [74] R. March, F. Doster, S. Geiger, Assessment of CO₂ Storage Potential in Naturally Fractured Reservoirs With
925 Dual-Porosity Models, *Water Resour. Res.* 54 (2018) 1650–1668. <https://doi.org/10.1002/2017WR022159>.

926 [75] H. Vosper, K. Kirk, C. Rochelle, D. Noy, A. Chadwick, Does numerical modelling of the onset of dissolution-

927 convection reliably reproduce this key stabilization process in CO₂ Storage?, *Energy Procedia*. 63 (2014)
928 5341–5348. <https://doi.org/10.1016/j.egypro.2014.11.566>.

929