

Structural Anisotropy-Driven Atomic Mechanisms of Phase Transformations in the Pt–Sn System

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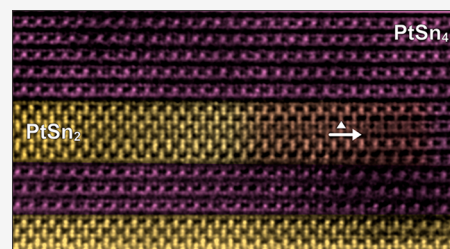
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Supporting Information

ABSTRACT: Using *in situ* atomic-resolution scanning transmission electron microscopy, atomic movements and rearrangements associated with diffusive solid to solid phase transformations in the Pt–Sn system are captured to reveal details of the underlying atomistic mechanisms that drive these transformations. In the PtSn₄ to PtSn₂ phase transformation, a periodic superlattice substructure and a unique intermediate structure precede the nucleation and growth of the PtSn₂ phase. At the atomic level, all stages of the transformation are templated by the anisotropic crystal structure of the parent PtSn₄ phase. In the case of the PtSn₂ to Pt₂Sn₃ transformation, the anisotropy in the structure of product Pt₂Sn₃ dictates the path of transformation. Analysis of atomic configurations at the transformation front elucidates the diffusion pathways and lattice distortions required for these phase transformations. Comparison of multiple Pt–Sn phase transformations reveals the structural parameters governing solid to solid phase transformations in this technologically interesting intermetallic system.

KEYWORDS: PtSn₄, PtSn₂, Pt₂Sn₃, *in situ* STEM, atomic resolution, phase transformation



Over the past two decades, interest in better understanding solid to solid phase transformations, particularly in nanomaterials, has grown considerably due to the potential for tuning their mechanical, catalytic, and electronic properties.^{1–4} In these solid to solid phase transformations—driven by temperature, pressure, or other stimuli^{5–7}—the transitions are either purely structural (martensitic) or combined structural and compositional (diffusive).^{8–10} The latter type, involving long-range movements of the atoms in the system that can proceed at an adjustable speed, is both more desirable and more difficult to study. In nanoscale materials, these transformations are sometimes accompanied by the formation of unexpected polymorphs and intermediate structures^{11–17} that further complicate the pathway. To effectively control these diffusive solid to solid phase transformations, understanding the underlying atomistic mechanisms that dictate the transition pathway between the product and parent phases is critical and needs direct measurements. Therefore, developing atomic-level experimental methods for characterizing transitional stages in these diffusive solid to solid phase transformations, as they happen, is necessary to gain insight into the key factors driving these phase transformations. Combining *in situ* and *ex situ* atomic-resolution measurements using analytical aberration-corrected scanning transmission electron microscopy (STEM) satisfies requirements for such measurements and has been shown to be effective for studying martensitic phase transformation.^{18,19} Here, we employed direct STEM imaging^{13,16–18,20} and diffraction²¹ to study local structure and crystallographic orientations and spatially resolved energy dispersive x-ray

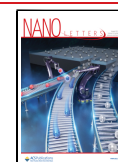
(EDX) spectroscopy^{15,22,23} to determine the composition of intermediate structures and phases.

In this work, we adapted analytical atomic-resolution STEM for *in situ* measurements to shed new light on the atomistic mechanisms of solid to solid phase transformations in the Pt–Sn system,^{24,25} going step by step from PtSn₄ to PtSn₂, then from PtSn₂ to Pt₂Sn₃, and, finally, from PtSn₂ to PtSn. These include transformations from highly anisotropic, layered, orthorhombic PtSn₄ to isotropic cubic PtSn₂ and then to another layered hexagonal Pt₂Sn₃. Atomic rearrangements captured by STEM imaging and spectroscopy showed a direct correlation between the type of anisotropy present in the crystal structure of either parent or product phase and the pathway involved in transforming to the product phase. Because the individual phases of the Pt–Sn system are still challenging to grow, especially in the form of thin films and nanoparticles^{20,26} needed for applications in catalysis and spintronics,^{27–30} successive solid to solid phase transformations discussed here serve as a promising alternative to direct synthesis.^{4,7}

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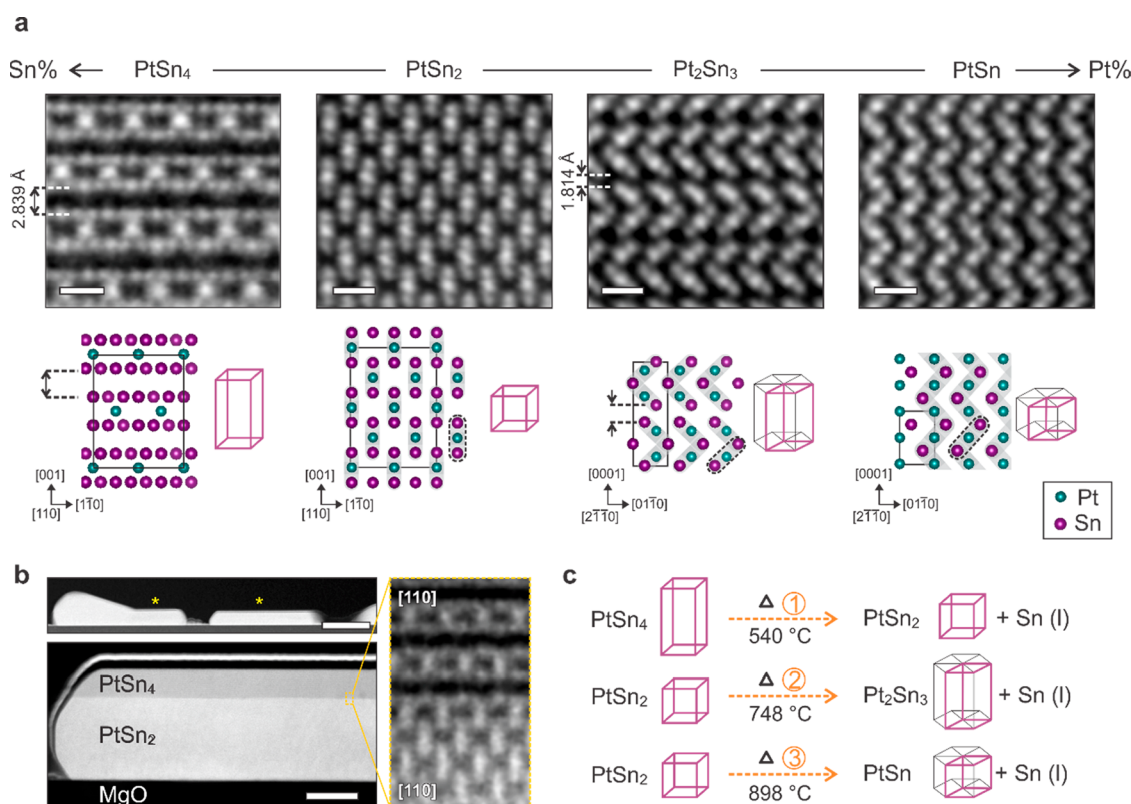


Figure 1. (a) High-resolution HAADF-STEM images of the four Pt_xSn_y phases discussed in this study. Scale bars are 5 Å. Atomic models and crystal lattice structures are illustrated below the HAADF-STEM images. Structural gaps in two-dimensional-like PtSn₄ and Pt₂Sn₃ are marked with dashed lines. Color codes for the atomic models: cyan for Pt and magenta for Sn. (b) Cross-sectional low- and medium-magnification HAADF-STEM images of an initial PtSn₄/PtSn₂ lamella grown on a MgO substrate. Scale bars are 500 nm (top) and 50 nm (bottom). A high-resolution image of the PtSn₂/PtSn₄ interface is presented on the right, showing the atomic structure of the interface. See Figure S1 for the EDX maps of the interface. (c) Simple schematic of three phase transformations (①, ②, and ③) among these four Pt_xSn_y phases, including their known bulk transition temperatures.²⁵

RESULTS

In Figure 1a, the atomic-resolution HAADF-STEM images are presented in order of a decreasing Sn/Pt atomic ratio. While the phases all exhibit clearly distinct crystal structures, they share an important resemblance that influences diffusive transformations among them. PtSn₄ has an orthorhombic crystal structure with space group *Ccee* (No. 68). It consists of layers separated with large 2.839 Å³¹ gaps along the [001] direction, which are clearly visible when viewed perpendicular to that direction. Each unit cell of PtSn₄ includes two such layers, with each layer composed of a Pt{0001} plane sandwiched between two Sn{0001} planes. PtSn₂ has a cubic structure with space group *Fm $\bar{3}$ m* (No. 225). While higher-symmetry PtSn₂ seems drastically dissimilar to lower-symmetry PtSn₄, note that introducing additional Pt atoms into the gaps in PtSn₄ results in a structure similar to that of PtSn₂; this is an important structural relationship that is discussed further below. Pt₂Sn₃ and PtSn both have hexagonal structures with space group *P6₃/mmc* (No. 194). While these two phases exhibit seemingly very similar atomic configurations, seen in the atomic-resolution images and crystal structure diagrams in Figure 1a, lower-symmetry Pt₂Sn₃ has 1.814 Å³² interlayer gaps perpendicular to the [0001] direction that are absent in PtSn. Note the interesting similarity among the atomic structures of PtSn₂, Pt₂Sn₃, and PtSn. A common Sn–Pt–Sn (SPS) elemental building unit is present in each of them, corresponding to nearly identical bonding configurations in

the SPS units across all three phases (Figure S1). Each of the four phases studied possesses unique convergent beam electron diffraction (CBED) patterns and electron energy-loss spectroscopy (EELS) spectra (Figures S2 and S3), which were employed to identify phases during *in situ* STEM heating experiments whenever atomic-resolution imaging was not possible.

The Pt_xSn_y film used for this study consists of islands that are composed of PtSn₄ and PtSn₂ phases, and PtSn₄/PtSn₂ lamella-structured islands were chosen for phase transformation experiments. See section S1 of the Supporting Information for details. Starting from a PtSn₄/PtSn₂ lamella, three types of phase transformations were studied: ① PtSn₄ to PtSn₂, ② PtSn₂ to Pt₂Sn₃, and ③ PtSn₂ to PtSn (Figure 1c). For each of the three diffusive phase transformations, we determine the atomistic mechanism and crystallographic orientation relationship (OR) characteristic of that transformation. Quantitative characterization of transformation kinetics (e.g., reaction rates, diffusion rates of atoms, and exact transition temperatures) is outside the scope of this study.

PtSn₄ to PtSn₂ (layered orthorhombic to cubic) Phase Transformation. To examine the transformation from PtSn₄ to PtSn₂, PtSn₄/PtSn₂ lamellae were gradually heated inside the transmission electron microscope (Figure 2a and Video S1). As the temperature reaches 250 °C, a dark lateral stripe contrast arises in the PtSn₄ region and expands because of transformation to the PtSn₂ phase. Propagation of product PtSn₂ is faster along the horizontal ⟨110⟩ directions than along

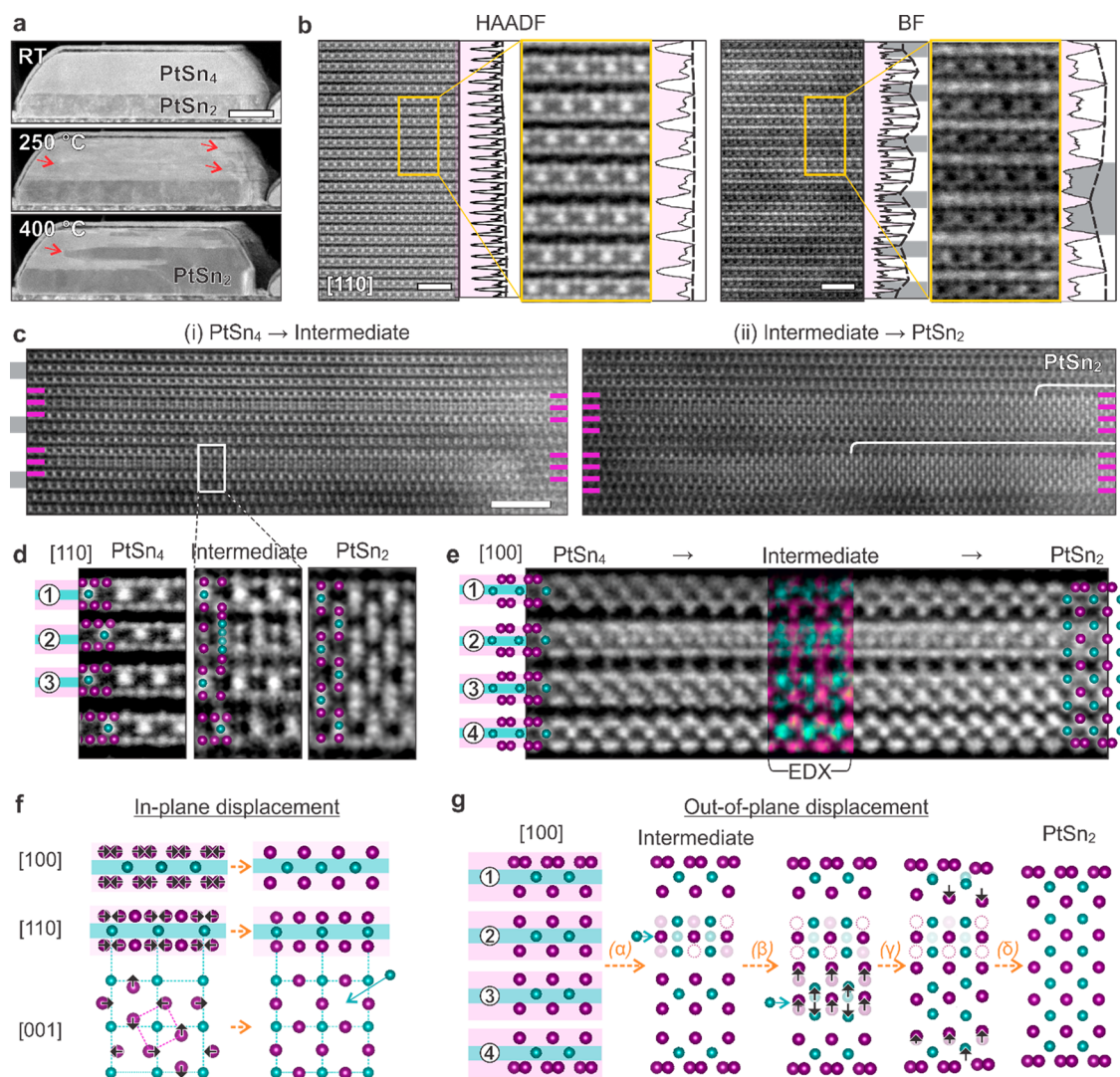


Figure 2. Nucleation stage in the PtSn_4 to PtSn_2 phase transformation. (a) *In situ* PtSn_4 to PtSn_2 phase transformation captured in a series of STEM images. LAADF-STEM is used here to enhance the contrast between the two phases. The growth of PtSn_2 in the PtSn_4 region is marked with red arrows. The scale bar is 100 nm. (b) HAADF- and BF-STEM images of PtSn_4 at 150 °C viewed along the $[110]$ direction. The insets show magnified images from the region inside the yellow boxes. Intensity line profiles are plotted to the right of each image, highlighting the emergence of lateral stripe contrast in the BF-STEM images. Scale bars are 2 nm. (c) Pair of HAADF-STEM images showing the initial stages of the PtSn_4 to PtSn_2 phase transition at 210 °C. The scale bar is 3 nm. (i) Emergence of three-layer-thick (highlighted with magenta lines) intermediate phases in PtSn_4 . The position of the layers with a dark contrast in BF-STEM is indicated with shaded gray lines on the left. (ii) Evolution of the intermediate phase into PtSn_4 . Four-layer PtSn_4 regions that transform into PtSn_2 are marked with magenta lines. (d) Comparison of the PtSn_4 , intermediate, and PtSn_2 phases from HAADF-STEM images in the $[110]$ direction with overlaid atomic models (cyan for Pt and magenta for Sn). (e) HAADF-STEM image capturing the phase evolution from PtSn_4 to PtSn_2 viewed from the $[100]$ projection. EDX elemental maps are overlaid in the region of the intermediate phase. The image is compared with HAADF-STEM images from bulk PtSn_4 and PtSn_2 in Figure S5. Schematics describe (f) in-plane and (g) out-of-plane atomic displacements during the PtSn_4 to PtSn_2 phase transition. Orange arrows indicate the change in the structure due to phase transformation, and black arrows on the atoms represent the associated atomic displacements. The cyan arrows with Pt atomic models indicate the suggested Pt diffusion pathway. Sn vacancies at layer ② are depicted with white circles with a magenta line color. Out-of-plane displacements are described by a combination of four unique displacement behaviors.

the $[001]$ direction, as evidenced by the laterally stretched contrast in the STEM images. At ~ 540 °C, in addition to the previously formed PtSn_2 , a liquid Sn region develops near the edge of the island (Figure S4). Eventually, all of the PtSn_4 transforms into PtSn_2 and liquid Sn. The liquid Sn solidifies into β -Sn, if the sample is cooled to room temperature (RT), but it vaporizes when heated above 540 °C. The crystallographic OR of the PtSn_4 to PtSn_2 phase transformation is $\text{PtSn}_4(001) \parallel \text{PtSn}_2(001)$ and $\text{PtSn}_4(110) \parallel \text{PtSn}_2(110)$, according to CBED data (Figure S4). The product PtSn_2 phase grows faster along the in-plane direction of the coherent planes.

To investigate the atomic mechanism of the phase transformation, a phase transition was initiated by slowly heating a $\text{PtSn}_4/\text{PtSn}_2$ lamella from RT to 300 °C (Video S2). First, at 150 °C, structural modulations develop in parent PtSn_4 . These modulations are subtle, resulting in no contrast change in HAADF-STEM images. However, they are clearly visible in bright-field (BF) STEM images as lateral dark contrasts appearing once every four to five PtSn_4 layers (Figure 2b and Figure S5). This periodic contrast is also visible in high-magnification BF-STEM images obtained from the other main in-plane $[100]$ direction (Figure S5). This periodic feature

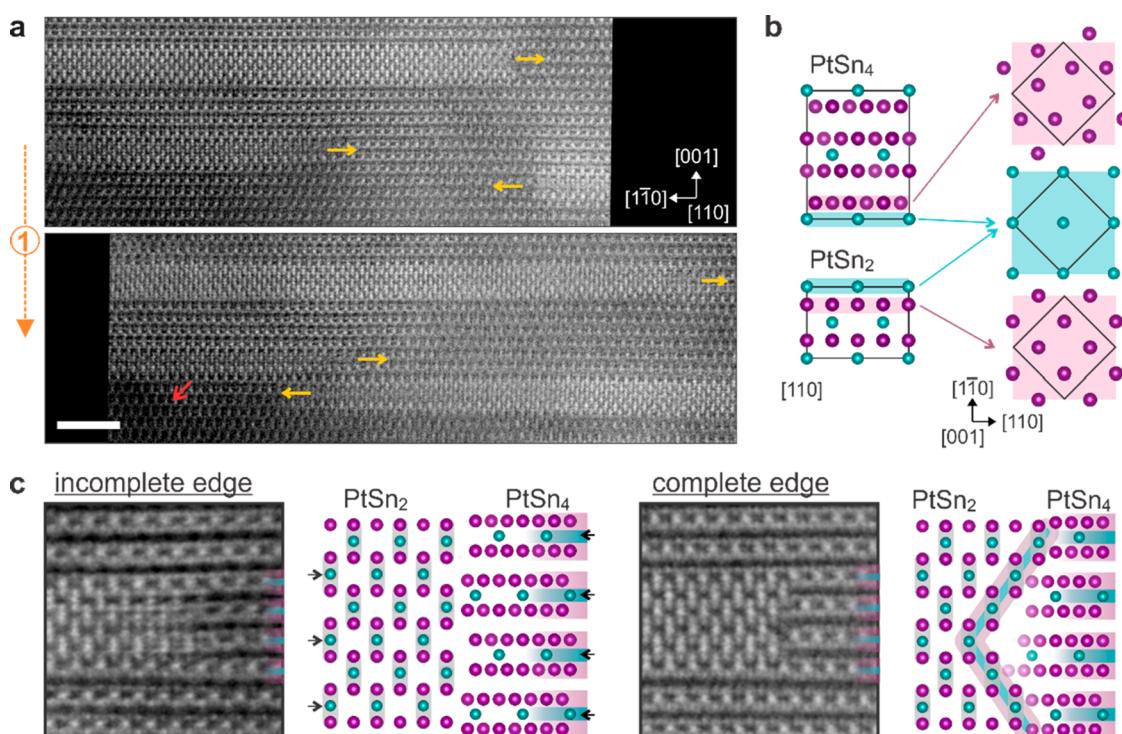


Figure 3. (a) Pair of HAADF-STEM images showing the growth of PtSn_2 slabs in PtSn_4 . The growth direction of each slab is indicated by yellow arrows located on the propagation front. A depleted region with a dark contrast is indicated by the red arrow. The scale bar is 3 nm. (b) Atomic models showing the common $\text{Pt}\{001\}$ and semicoherent $\text{Sn}\{001\}$ planes in PtSn_2 and PtSn_4 . (c) HAADF-STEM images and corresponding atomic models of incomplete and complete edges of the propagation front of a PtSn_2 slab in a PtSn_4 matrix. Pt diffusion paths are indicated by black arrows. The complete edge between two phases corresponds to a set of $\text{PtSn}_2\{112\}$ planes, which are marked with shaded lines.

remains the same when the sample is cooled from 150 °C to RT, indicating a permanent change in the structure. We speculate that this periodic contrast in the BF (not visible in HAADF) images stems from periodic strain in the structure resulting from ordered vacancy formation (mostly lighter Sn vacancies), producing an enlarged superlattice that can be readily accommodated in layered PtSn_4 . Similar two-dimensional (2D)-like structural modulations at increased temperatures have been observed in In_2Se_3 .³³

At ~ 200 °C, a new intermediate phase appears in the layers between the dark bands of the BF-STEM image; the intermediate phase develops into PtSn_2 upon further heating (Figure 2c and Figure S5). The intermediate phase exhibits an atomic structure that is different not only from those of the parent (PtSn_4) and product (PtSn_2) phases but also from those of any existing Pt_xSn_y phases. It is confined within three PtSn_4 layers (①–③ in Figure 2d) and spans laterally. The HAADF-STEM image and EDX elemental maps of this intermediate phase (Figure 2d and Figure S5) reveal major atomic rearrangement in the middle layer ② and minor changes in the two capping layers, ① and ③. The atomic modifications needed for the formation of the intermediate phase from PtSn_4 can be described by a combination of in-plane and out-of-plane displacements, as depicted schematically in panels f and g of Figure 2. Whereas in-plane displacement of Sn atoms (which transforms a $\text{Sn}\{001\}$ plane in PtSn_4 into a $\text{Sn}\{001\}$ plane of PtSn_2) occurs in all three layers, the atomic movements in the out-of-plane direction are distinct in each layer. Out-of-plane displacement of Pt atoms occurs only in the middle layer, ②, where a $\text{Pt}\{001\}$ plane splits into two $\text{Pt}\{001\}$ planes (step α in Figure 2g). Here, Pt is supplemented via in-plane atomic influx, which can be deduced from the unaltered contrast of

the neighboring ① and ③ layers, and Sn atoms seem to diffuse along the in-plane direction generating Sn vacancies. Similar out-of-plane atomic displacements occurring at lower temperatures, albeit far more subtle, may account for the periodic feature captured in BF-STEM images at 150 °C.

Next, the three-layer-thick intermediate phase interacts with one adjacent PtSn_4 layer to transform into PtSn_2 . Such evolution of atomic rearrangements from PtSn_4 to the intermediate phase and to the PtSn_2 phase was captured in a single HAADF-STEM image (Figure 2e). The in-plane displacement of Sn atoms observed in the intermediate phase also occurred in all remaining layers except for the two outermost boundary $\text{Sn}\{001\}$ planes. Along with the out-of-plane Pt splitting in layer ② from the intermediate phase formation (step α), the $\text{Pt}\{001\}$ plane in layer ③ also splits in the out-of-plane direction, generating two $\text{Pt}\{001\}$ planes (Figure 2g). Simultaneously, out-of-plane shifts of $\text{Sn}\{001\}$ planes toward the center of layer ② occur in all of the layers (β – δ step). Eventually, four layers of PtSn_4 turn into a 1.5-unit cell (6- $\text{Pt}\{001\}$ -plane-thick) PtSn_2 slab. Note that such a transformation requires diffusion of additional Pt into the matrix along with the removal of Sn. Pt diffuses into layers ② and ③ primarily within the $\text{Pt}\{001\}$ layer along the in-plane direction. As one can see in the atomic models from the $[001]$ direction, the $\text{Pt}\{001\}$ plane is less densely packed than the $\text{Sn}\{001\}$ plane and has some vacant positions, which can be used as the diffusion pathway. Darker patches with a modified atomic configuration are observed near the edge of the new PtSn_2 slabs, indicating formation of Pt-depleted regions after the diffusion of Pt into the PtSn_2 phase (Figure 3a). Sn atoms diffuse out to the lateral surface primarily along the in-plane

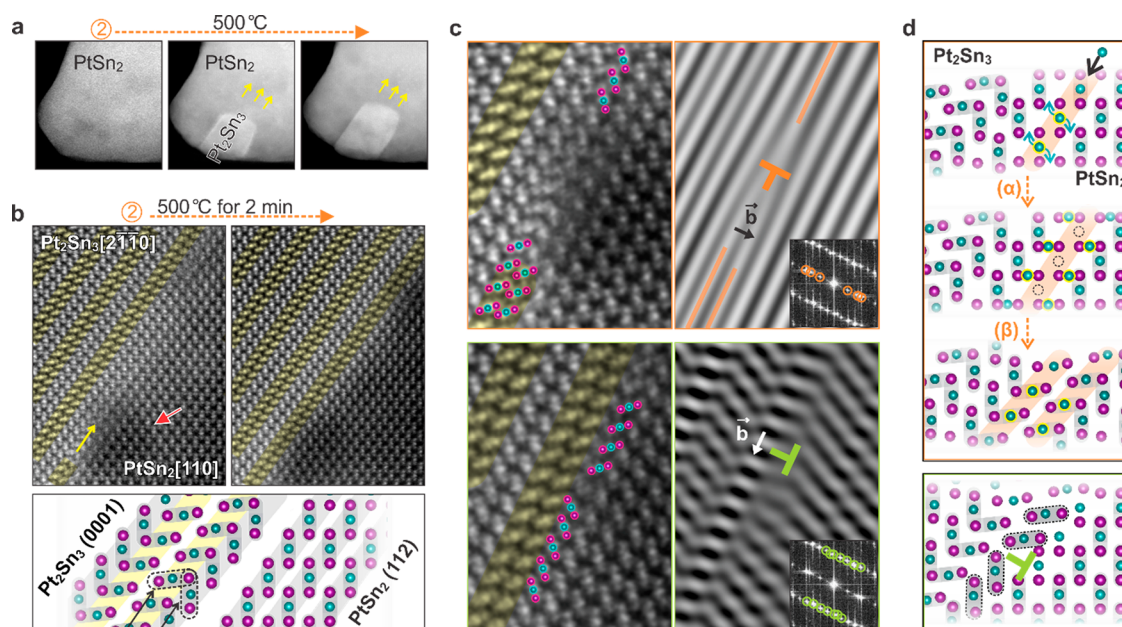


Figure 4. PtSn₂ to Pt₂Sn₃ phase transformation. (a) Set of *in situ* HAADF-STEM images showing the formation and growth of Pt₂Sn₃ in PtSn₂. (b) Atomic-resolution HAADF-STEM images showing the growth of a layer of Pt₂Sn₃ at the PtSn₂/Pt₂Sn₃ boundary. Atomic models of the Pt₂Sn₃/PtSn₂ interface are included at the bottom, where two types of SPS configurations in Pt₂Sn₃ are outlined with black dashed lines and labeled as p (for parallel) and t (for tilted). The Pt₂Sn₃ layer growth direction is indicated by the yellow arrow, and a depleted region with dark contrast is indicated by the red arrow. Pt₂Sn₃{0001} planes composed of t-type SPS are highlighted in yellow. (c) HAADF-STEM images of type 1 (top) and type 2 (bottom) edge dislocations at the PtSn₂–Pt₂Sn₃ interface with the Burgers vector perpendicular and parallel to the interface, respectively. Fourier-filtered images on the right show the location of dislocations (insets, used Fourier spots). (d) Schematics describing the role of the two dislocations in transforming the p-type SPS units into t-type units at the PtSn₂–Pt₂Sn₃ interface.

direction, as liquid Sn regions form only at the side face of the original PtSn₄ region, not at the top or bottom (Video S1).

As the phase transformation progresses, the 1.5-unit cell PtSn₂ slab expands along the in-plane direction in the PtSn₄ matrix while sandwiched between the PtSn₄{001}–PtSn₂{001} interfaces (Figure 3a and Video S3). Such anisotropic growth is favorable, owing to the presence of a coherent PtSn₄{001}||PtSn₂{001} interface as well as the layered structure of PtSn₄. As illustrated in Figure 3b, these two phases possess identical Pt{001} atomic planes and similar Sn{001} planes with a negligible in-plane lattice mismatch (<0.4%), which contributes to the stability of the interface. The large gaps in the layered structure of PtSn₄ provide a buffer in interaction with the neighboring layers and flexibility in the atomic displacements for PtSn₂ formation. Furthermore, the thickness of the 1.5-unit cell PtSn₂ slab (1.934 nm)³⁴ is comparable to the thickness of the four layers of PtSn₄ (1.987 nm),³¹ which further stabilizes the configuration. The layered structure of PtSn₄ dictates every aspect of this phase transformation, including the emergence of the intermediate phase, stabilization of homogeneous nuclei, and directional growth of the product PtSn₂ phase.

Once the nucleation of PtSn₂ is complete, it grows primarily by expanding laterally without forming an intermediate phase. Two distinct morphologies of the propagation front are observed: diffuse (incomplete) and sharp (complete) edges (Figure 3c). While the incomplete edge shows gradual contrast variation across the boundary, the complete edge exhibits clear PtSn₂{112} planes that terminate in SPS units at the boundary. The incomplete edge forms as a result of the in-plane atomic diffusion (Pt influx and Sn efflux), and those edges with an appropriate stoichiometry can be stabilized after formation of

the PtSn₂{112} planes. This distinct edge at the propagating front indicates the high stability of the PtSn₂{112} planes, a characteristic that will be revisited in the following sections.

PtSn₂ to Pt₂Sn₃ (cubic to layered hexagonal) Phase Transformation. After further gradual heating to 500 °C (or slightly higher), a subsequent PtSn₂ to Pt₂Sn₃ phase transformation was observed. As discussed above (Figure 1a), the product Pt₂Sn₃ phase also has a layered structure but with narrower (1.814 Å) gaps between the {0001}Pt₂Sn₃ planes. In *in situ* experiments that examined this transition, pulse heating was applied to ensure the better thermal stability of the TEM holder and more consistent atomic-resolution STEM imaging (Figure S6 and Videos S4 and S5). This phase transformation is initiated via nucleation of Pt₂Sn₃ at the surfaces of PtSn₂ particles, after which the new Pt₂Sn₃ phases grow anisotropically through the parent material (Figure 4a and Video S4). In Figure 4b, two sequential atomic-resolution HAADF-STEM images show the growth of one new layer of Pt₂Sn₃ at the boundary between PtSn₂ and Pt₂Sn₃. The growth direction of the new phase is Pt₂Sn₃{110}, which is one of the in-plane directions of Pt₂Sn₃.

The crystallographic OR of the PtSn₂ to Pt₂Sn₃ transformation was analyzed using CBED patterns and atomic-resolution STEM images as PtSn₂{112}||Pt₂Sn₃{0001} and PtSn₂{110}||Pt₂Sn₃{2110} (Figure S6). This crystallographic OR can be better understood in terms of the structural similarity between the two phases. The layer of Pt₂Sn₃ (Pt₂Sn₃{0001}) is comprised of an array of V-shaped structural units made of two tilted SPS units (labeled p and t in Figure 4b). One side of the layer (p side) forms a structure almost identical (except for a minor lattice mismatch of $\Delta d_{\text{in-plane}} = 4.9\%$) to the PtSn₂{112} plane, which is a stable facet of PtSn₂

as discussed above (Figure 3c) [$\Delta d_{\text{in-plane}} = (d_{\text{PtSn}_2\{110\}} - d_{\text{Pt}_2\text{Sn}_3\{2\overline{1}10\}})/d_{\text{PtSn}_2\{110\}}$]. Consequently, a stable semicoherent interface of $\text{Pt}_2\text{Sn}_3\{0001\}/\text{PtSn}_2\{112\}$ with misfit dislocations is produced. The phase transformation progresses in the direction parallel to the interface, preserving the semicoherent state (Video S5).

Two types of edge dislocations are frequently observed at this interface: one located at the propagation front of Pt_2Sn_3 with a Burgers vector perpendicular to the interface (type 1) and the second with a Burgers vector parallel to the interface (type 2) (Figure 4c). In type 1 dislocations, a single interfacial $\text{PtSn}_2\{112\}$ layer from the parent phase splits into two SPS layers of $\text{Pt}_2\text{Sn}_3\{0001\}$. The splitting is accompanied by reorientation of the p-type SPS in PtSn_2 to the t-type configuration, resulting in Pt_2Sn_3 . The type 2 dislocations, which originate from the in-plane lattice mismatch between the $\text{PtSn}_2\{110\}$ and $\text{Pt}_2\text{Sn}_3\{2110\}$ planes at the interface, are misfit dislocations. Both dislocations play critical roles in the transformation of PtSn_2 into Pt_2Sn_3 .

At the type 1 dislocations, the Pt atom at the middle of each SPS in the $\text{PtSn}_2\{112\}$ layer moves to the nearby sites between two Sn atoms, splitting one $\text{Pt}\{112\}_{\text{PtSn}_2}$ atomic plane into two planes and producing new t-type SPS units (step α in Figure 4d). Next, the two layers of the t-type SPS units relax and form the $\text{Pt}_2\text{Sn}_3\{0001\}$ layers. The additional Pt needed to fill the new Pt sites and to form Pt_2Sn_3 diffuses in parallel to the $\text{Pt}_2\text{Sn}_3\{0001\}$ layers, leaving Pt-depleted regions that can be readily observed as dark patches near the interface (Figure S6). Similar to the PtSn_4 to PtSn_2 phase transformation, the propagation and atomic rearrangements are highly anisotropic following the layered structure. The type 2 misfit dislocation drives flipping of the SPS unit between p and t (Figure 4c). The tensile strain in the PtSn_2 side produces an open space for easier atomic movement, resulting in flipping of the SPS unit and influx of Pt. Dislocation-guided transformations between two solid phases has been observed in other systems.^{8–10} These observations, highlighted by the alignments of the SPS layers, again point to a dominant role of the structural characteristics of the parent and product phases in facilitating diffusive transformation.

Different from the cases discussed above, phase transformations from cubic PtSn_2 to hexagonal PtSn , where neither the parent phase nor the product phase has well-defined anisotropic structural characteristics, occur isotropically (section S2), further supporting the critical role of the lattice structure in the atomistic mechanisms of phase transformations. Our results demonstrate that solid to solid phase transformations follow a pathway that lies structurally between the parent and product phases, and the anisotropy in their lattice structures plays the most significant role in determining the atomic pathway. Atomic planes parallel to layers in a structure provide natural pathways for a facile elemental diffusion, and the gaps between the layers act as a buffer for local atomic modulations and rearrangements.

CONCLUSION

Our analysis of diffusive solid to solid phase transformations in the Pt–Sn system, based on direct atomic-resolution *in situ* STEM characterization, details how the relationship between the crystal structures of the parent and product phases dictates the atomistic mechanisms of the phase transformation. In the PtSn_4 to PtSn_2 phase transformation, in which highly

anisotropic PtSn_4 with a layered structure transforms into isotropic PtSn_2 with a cubic structure, the transformation is initiated at 150 °C by forming a 2D periodic prestructure guided by the symmetry of the parent PtSn_4 . Upon further heating to ~200 °C, nucleation of an intermediate structure takes place following the prestructure. Next, the intermediate structure transforms into PtSn_2 and propagates in PtSn_4 . The resulting slab of PtSn_2 with a thickness of 1.5 unit cells is coherently surrounded by the parent PtSn_4 phase. These atomic movements and rearrangements are driven by two structural factors: the presence of coherent planes at the interface of PtSn_4 and PtSn_2 and the high stability and flexibility of layered structures in PtSn_4 . In the PtSn_2 to Pt_2Sn_3 phase transformation, where higher-symmetry PtSn_2 with a cubic structure transforms into lower-symmetry Pt_2Sn_3 with a layered structure at ≥ 500 °C, the transformation initiated on the surfaces propagates anisotropically along the in-plane direction of the layers of Pt_2Sn_3 mediated by the dislocations. Here again, the presence of a semicoherent interface between the phases and the stability and flexibility of the layered structure of the product Pt_2Sn_3 phase drive the atomic movements and rearrangements. Future computational studies investigating the energetics of the atomic planes and their strain dependency will help to improve our understanding of these observations. In the PtSn_2 to PtSn phase transformation, in which both parent and product phases are isotropic, the atomic movements and rearrangements are also isotropic without the formation of clearly defined interfacial structures.

The observed atomistic mechanisms of these transformations exhibit a clear correlation with the structural characteristics of the parent and product phases. Structural anisotropy in one of the involved phases can dictate a directional phase transformation due to the easy formation of stable habit planes, facile diffusion pathways, and structural flexibility. These results suggest that more research into examining the atomistic mechanisms of solid to solid phase transformations should be conducted to fully elucidate the principles of solid-state reactions and guide the efficient engineering of nanomaterials through phase transformation.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.nanolett.3c02162>.

Methods; structure of the as-grown Pt_xSn_y films (section S1); crystal structures of the PtSn_2 , Pt_2Sn_3 , and PtSn phases (Figure S1); CBED map and atomic models at the major zone axes of all four Pt_xSn_y compounds (Figure S2); low-loss EELS obtained from the four Pt_xSn_y phases (Figure S3); supporting data for the PtSn_4 to PtSn_2 phase transformation (Figure S4); STEM images of the periodic feature and the intermediate phase in the PtSn_4 to PtSn_2 transition (Figure S5); supporting data for the Pt_2Sn_3 to PtSn_2 phase transformation (Figure S6); supporting data for the PtSn_2 to PtSn phase transformation (Figures S7 and S8); and details of the PtSn_2 to PtSn phase transformation (section S2) (PDF)

Video S1 (MP4)

Video S2 (MP4)

Video S3 (MP4)

Video S4 (MP4)

Video S5 (MP4)

Video S6 (MP4)

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Author Contributions

H.Y. and K.A.M. conceived the project. H.Y. performed STEM experiments with input from K.A.M. D.Z. grew Pt–Sn thin films with input from J.-P.W. T.B. contributed to analysis. H.Y. and K.A.M. prepared the manuscript with contributions from all authors. K.A.M. directed all aspects of the project.

Notes

The authors declare no competing financial interest.

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