

Deep Learning Model of Hiss Waves in the Plasmasphere and Plumes and Their Effects on Radiation Belt Electrons

Sheng Huang^{1*}, Wen Li¹, Qianli Ma^{1,2}, Xiao-Chen Shen¹, Luisa Capannolo¹, Miroslav Hanzelka^{1,5}, Xiangning Chu³, Donglai Ma², Jacob Bortnik², and Simon Wing⁴

¹Center for Space Physics, Boston University, Boston, MA, USA.

²Department of Atmospheric and Oceanic Sciences, University of California, Los Angeles, CA, USA.

³Laboratory for Atmospheric and Space Physics, University of Colorado Boulder, Boulder, CO, USA.

⁴Applied Physics Laboratory, The Johns Hopkins University, Laurel, MD, USA.

⁵Department of Space Physics, Institute of Atmospheric Physics of the Czech Academy of Sciences, Prague, Czechia.

* Correspondence:

Sheng Huang (hs2015@bu.edu); Wen Li (wenli77@bu.edu)

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Abstract

Hiss waves play an important role in removing energetic electrons from Earth's radiation belts by precipitating them into the upper atmosphere. Compared to plasmaspheric hiss that has been studied extensively, the evolution and effects of plume hiss are less understood due to the challenge of obtaining their global observations at high cadence. In this study, we use a neural network approach to model the global evolution of both the total electron density and the hiss wave amplitudes in the plasmasphere and plume. After describing the model development, we apply the model to a storm event that occurred on 14 May 2019 and find that the hiss wave amplitude first increased at dawn and then shifted towards dusk, where it was further excited within a narrow region of high density, namely a plasmaspheric plume. During the recovery phase of the storm, the plume rotated and wrapped around Earth, while the hiss wave amplitude decayed quickly over the nightside. Moreover, we simulated the overall energetic electron evolution during this storm event, and the simulated flux decay rate agrees well with the observations. By separating the modeled plasmaspheric and plume hiss waves, we quantified the effect of plume hiss on energetic electron dynamics. Our simulation demonstrates that, under relatively quiet geomagnetic conditions, the region with plume hiss can vary from L=4 to 6 and can account for up to an 80% decrease in electron fluxes at hundreds of keV at L>4 over three days. This study highlights the importance of including the dynamic hiss distribution in future simulations of radiation belt electron dynamics.

1 Introduction

Hiss waves are a type of whistler mode, broadband emission that typically exists in the Earth's high density plasmasphere and plume regions (Chan & Holzer, 1976; Hayakawa et al., 1986; Larkina &

Likhter, 1982; Meredith, 2004; Ripoll et al., 2020; Thorne et al., 1973). Since their early discovery (Dunckel & Helliwell, 1969; Russell et al., 1969), hiss waves have been extensively studied, and many of their properties have been revealed (Hayakawa & Sazhin, 1992; Li et al., 2015a; Tsurutani et al., 2015).

Through cyclotron resonant interactions, hiss can pitch-angle scatter electrons with energies ranging from tens of keV up to several MeV (Horne & Thorne, 1998; Li et al., 2007; Ma et al., 2016; Ni et al., 2014). They are responsible for creating the slot region between the inner and outer radiation belts and are believed to be the main driver of the outer belt electron decay during quiet times (Lam et al., 2007; Ma et al., 2015), thus playing an important role in controlling the structure and dynamics of the radiation belts.

Hiss waves are believed to have multiple generation mechanisms, which are still under active research (e.g., Bortnik et al., 2009; Green, 2005; Liu et al., 2020). Lightning-generated whistlers from low altitudes can propagate and evolve into hiss (Bortnik et al., 2003; Sonwalkar & Inan, 1989), but they account for only a portion of the wave power at frequencies > 2 kHz at $L < 3.5$ (Meredith et al., 2006). In recent years, more and more observations and ray-tracing simulations have linked hiss waves with chorus waves propagating into the plasmasphere (Bortnik et al., 2008; Chen et al., 2012a, 2012b; Church & Thorne, 1983; Santolík et al., 2006). This correlation is supported by statistical analyses of wave distribution (Agapitov et al., 2018; Meredith et al., 2013) as well as direct observations through event analyses (Bortnik et al., 2009; Li et al., 2015b). In addition to lightning-generated whistlers and chorus waves propagating into the plasmasphere, electron cyclotron instability can also be a possible energy source for hiss by locally amplifying it to observable levels (Kennel & Petschek, 1966; Thorne et al., 1979). Although the wave growth rate is generally weak (Church & Thorne, 1983; C. Y. Huang et al., 1983), recent studies have shown that the high-frequency hiss waves may be locally generated (Fu et al., 2021; Meredith et al., 2021). In addition, the sharp density gradient near the plasmapause and a fresh injection of anisotropic hot electrons drifting from the nightside plasma sheet can aid in generating intense low-frequency hiss, particularly favored when plasmaspheric plumes are present (Chen et al., 2014; Li et al., 2013; Su et al., 2018; Wu et al., 2022). Plume hiss is thus gaining more and more attention due to its potential role in controlling radiation belt dynamics (Summers et al., 2008). In the era of Van Allen Probes, hiss is found to be prevalent inside plumes (Shi et al., 2019; W. Zhang et al., 2019), and both observations and simulations recognize its importance in precipitating electrons in the outer radiation belt (Li et al., 2019; Ma et al., 2021; Millan et al., 2021; Qin et al., 2021). However, the observation of plume hiss is highly limited during individual events due to a lack of global coverage, and simulations are usually performed based on the statistical properties of plume hiss. Therefore, the spatiotemporal evolution of plume hiss and its effects on energetic electron dynamics remain elusive, though they are believed to critically affect the loss rate of energetic electrons in radiation belts.

In this study, we propose a deep learning approach to model the global evolution of hiss and total electron density, inspired by Bortnik et al. (2016; 2018). Deep learning techniques have shown promising results in space weather modeling by analyzing information from large datasets (Chu et al., 2017a, b; 2021; Ma et al., 2022, Wing et al., 2005, 2022). We present the methodology for our model in section 2. In section 3, we analyze the model performance and apply it to a geomagnetic storm event where the complete evolution of plume hiss is predicted. Then, we simulate the energetic electron evolution based on the modeled hiss and total electron density, and quantify the effects of plume hiss. In section 4, we discuss our findings, followed by our conclusions in section 5.

2 Data and Deep Learning Model

2.1 Van Allen Probes Data

We train the model using observations from the twin Van Allen Probes (also known as RBSP; Mauk et al., 2013) throughout the majority of their operational time (2013–2019). The Electric and Magnetic Field Instrument Suite and Integrated Science (EMFISIS; Kletzing et al., 2013) suite onboard RBSP provides in-situ measurements of the field and waves with a time resolution of ~ 6 s for the survey mode. Total electron density (N_e) is inferred from the upper hybrid resonance frequency (Kurth et al., 2015) based on the measurements from the High Frequency Receiver (HFR). The WaveForm Receiver (WFR) measures wave activity, which we use to calculate the amplitude of hiss waves following Li et al. (2015a) summarized as follows:

- 1) wave ellipticity > 0.7 ;
- 2) wave planarity > 0.2 ;
- 3) spectral frequency range over 20 – 4000 Hz.

When the satellites are outside the plasmasphere or plume (according to the wave power of electron cyclotron harmonic waves; Shen et al., 2019), the wave amplitude is set to 0.2 pT to indicate no hiss wave. The whole hiss wave dataset has a similar trend to the statistics by Li et al. (2015a) that hiss wave tends to occur on the dayside during enhanced levels of substorm activity (not shown here). The satellite location is also used for training purposes, including L shell, magnetic local time (MLT), and magnetic latitude (MLAT). The MLT is converted into $\sin(\text{MLT}/12 \cdot \pi)$ and $\cos(\text{MLT}/12 \cdot \pi)$ to account for the discontinuity at $\text{MLT}=24$. Additionally, the spin-averaged electron fluxes measured by the Magnetic Electron Ion Spectrometer (MagEIS) instrument (Blake et al., 2013) in the Energetic Particle Composition and Thermal Plasma (ECT) suite (Spence et al., 2013) are used to compare with the results of radiation belt simulations using our density and wave models.

2.2 Geomagnetic Indices

To model both the electron density and wave amplitude at a specific location observed by satellites, we use the geomagnetic indices SML, SMU, Hp30, and SYM-H, which measure the level of geomagnetic disturbance at different latitudes. The SML and SMU indices (Gjerloev, 2012; Newell & Gjerloev, 2011; from SuperMAG Web Service) provide better time coverage (to include recent year data) compared to the more commonly used AL and AU indices. The Hp30 index (Matzka et al., 2021; from GFZ German Research Centre for Geosciences) is designed to improve the temporal resolution of Kp index from 3 h to 30 min. To capture the most variation in the data without introducing many artifacts from interpolation, all satellite observations and geomagnetic indices are interpolated to a time resolution of 1 minute.

2.3 Deep Learning Model

We adopt a similar model structure to that of Huang et al. (2022), as illustrated in Figure 1. In this framework, geomagnetic indices are used as the inputs to a neural network module, known as Long Short-Term Memory (LSTM; Hochreiter & Schmidhuber, 1997). LSTM is well-suited for modeling data sequences in time-series format and can effectively capture the temporal evolution within the data (Karim et al., 2018; Siami-Namini et al., 2019). The extracted output feature H at time t_n can be viewed as a representation of the inner magnetospheric state at time t_n , described solely based on the geomagnetic indices. Subsequently, H is used to fit the satellite observations (both total electron density and hiss wave amplitude), with corresponding satellite location as an input (see Section 2.4).

124 By employing LSTM to process geomagnetic indices alone (without any RBSP data), the temporal
125 evolution is decoupled from the location information, which enables our model to simultaneously
126 learn the complex spatial dependence and the smooth transition along the satellite orbital
127 observations over time.

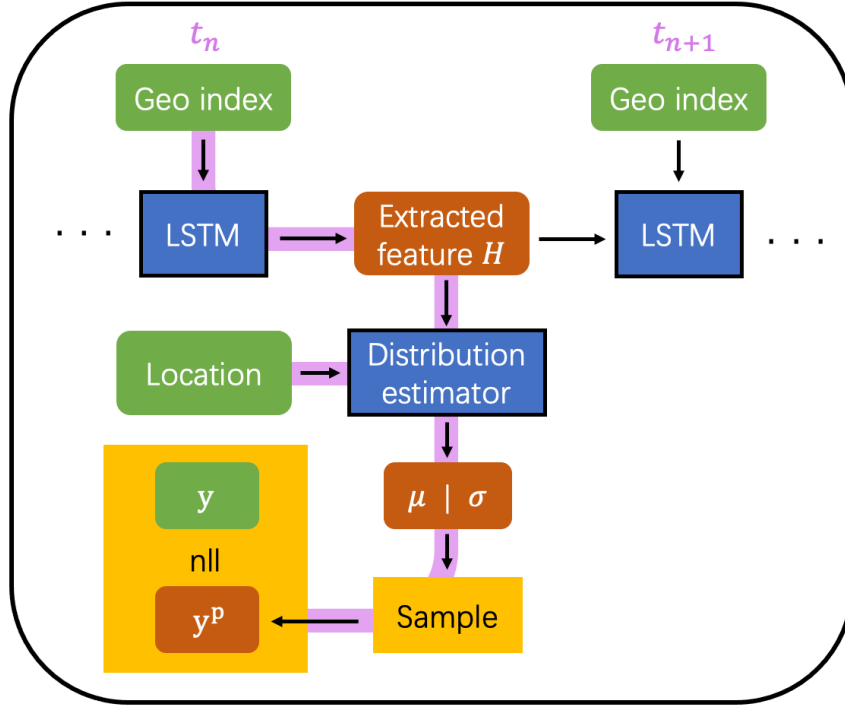
128 As hiss wave amplitude varies significantly in different regions, models tend to estimate the average
129 activity while treating the variation as noise, thus underestimating the wave activity (trained with the
130 same model structure as Huang et al. (2022) on hiss wave; not shown here). To better capture the
131 dynamic nature of hiss wave activity, instead of directly predicting a quantity in a deterministic
132 approach, we use a neural network module that estimates the wave probability distribution (modeling
133 both the mean μ and standard deviation σ) at a specific location and time. This approach essentially
134 introduces an estimation of the uncertainty (Blundell et al., 2015) in the data and is critical to model
135 quantities with large variations (Tasistro-Hart et al., 2021). We avoid applying significant smoothing
136 to the RBSP data to retain the full information carried in the variation. We sample a prediction y^p
137 from the modeled mean and standard deviation and calculate the negative log likelihood

138
$$nll = \sum_i (\log \sqrt{2\pi\sigma^2} + \frac{(y_i - \mu)^2}{2\sigma^2})$$

139 between the observation and model prediction. This process essentially maximizes the possibility of
140 measuring the observed quantity given the estimated distribution. The calculated loss is then used to
141 update the model parameters through the standard backpropagation procedure.

142 To address the issue of an unbalanced dataset in training the hiss wave model, we have implemented
143 a weighted sampler. While we dedicate considerable attention to geomagnetically active times, it is
144 important to note that quiet times are more common and generally exhibit low wave activities. When
145 the model is trained on the entire dataset, it tends to learn more efficiently from weaker waves,
146 resulting in an underestimation of wave activity. To mitigate this imbalance, we use a weighted
147 sampler that selects training samples based on a probability proportional to the largest wave
148 amplitude within the subsequent 1-hour period. Consequently, periods with stronger wave activity
149 are more likely to be included in the training process than those with weak wave activity, leading to a
150 model with improved performance during geomagnetically active times.

151 We include more details of the model structure and optimization procedure in Appendix A.



152

153 **Figure 1.** Model structure and workflow. Purple line: data flow at time t_n ; Green box: model input;
 154 Blue box: neural network model modules; Red box: (intermediate) model output; Yellow box: data
 155 operation. After the hidden state H is encoded by LSTM from the geomagnetic indices, a probability
 156 distribution is estimated at the satellite location, and a prediction y^p is sampled from this distribution.
 157 The negative log likelihood (nll) is calculated between the prediction y^p and satellite observation y ,
 158 and is further used to update the model parameters through backpropagation. y denotes either total
 159 electron density or hiss wave amplitude.

160 2.4 Data Processing

161 The data from 2013 to 2019 is divided into 7-day blocks, with 70% randomly assigned as the training
 162 set, 20% as the validation set, and 10% as the test set. The period 13-19 May 2019 is also kept in the
 163 test set for further simulation (see Section 3.2). This division into 7-day blocks is chosen to avoid
 164 data leakage that is common in time-series modeling, and is short enough to allow for a large number
 165 of blocks, and long enough to prevent information leakage, while also considering long-term
 166 seasonal and solar cycle variations. After the training time range is settled (7-day blocks that belong
 167 to the training set combined), during each runtime we generate training samples with a weighted
 168 sampler using the following procedure. 1) Before the training starts, both Van Allen Probe A and B
 169 observations that fall within these 7-day blocks are assigned with a sequence of weights. Each weight
 170 that corresponds to a certain timestamp is calculated to be proportional to the largest wave amplitude
 171 within the subsequent 1-hour period. The resulting weight sequence has the same length as satellite
 172 observations. 2) During the training, starting times of the satellite observations are randomly picked
 173 given the weight sequence, and for each selected time, a period of 10-hour that follows the selected
 174 time is used in the training process. Each 10-hour period of observation is then paired with the
 175 corresponding 10 hours of geomagnetic indices and the preceding 24 hours of historical
 176 geomagnetic indices at 1-min resolution to provide information on the state of the inner
 177 magnetosphere. In summary, for each 10-hour period, the model takes 24 hours of historical
 178 geomagnetic indices as inputs, followed by another 10 hours of geomagnetic indices and satellite

location (L , $\sin(\text{MLT}/12*\pi)$, $\cos(\text{MLT}/12*\pi)$, MLAT) at the same time. The model predicts the total electron density and hiss wave amplitude within the 10-hour period. The negative log likelihood is calculated between observation and model prediction over each 10-hour period. Loss is accumulated over a number of sequences trained at the same time, until it backpropagates to update the neural network parameters.

3 Results

3.1 Model Performance

The overall model performance is shown in Figure 2 for total electron density (A-C) and hiss wave amplitude (D-F) for different datasets, respectively. The x-axis represents the observed quantity y (density or hiss wave amplitude), while the y-axis represents the corresponding modeled quantity y^p . The color represents how many $y-y^p$ pairs are located in that region. The red dashed diagonal line indicates a perfect model prediction ($y=y^p$). The darker areas, concentrated near the red line, indicate good model performance for the majority of the data. This is also quantified by the correlation coefficient between $\log_{10} y$ and $\log_{10} y^p$ denoted by “ r ” in each panel. The model performance for electron density is similar to that of Huang et al. (2022), where the Pearson correlation coefficient for the test dataset (Figure 2C) is about 0.9, close to that of the validation set (Figure 2B) indicating that the model generalization ability is good. The mean square error (mse) is 0.16, indicating that the model generalizes and performs very well in modeling electron density. For the hiss wave amplitude, there is more spread of the darker areas (Figure 2F) with $r=0.74$, $\text{mse}=0.53$ for the test dataset, which suggests that the model performance is worse for the hiss amplitude than electron density. This is partly because the wave activity is highly dynamic, exhibiting fluctuations on short timescales, and thus is less predictable compared to the cold plasma density. Nevertheless, by adopting a probability-based approach, our model reproduces the general global wave evolution fairly well, as presented in the following section.

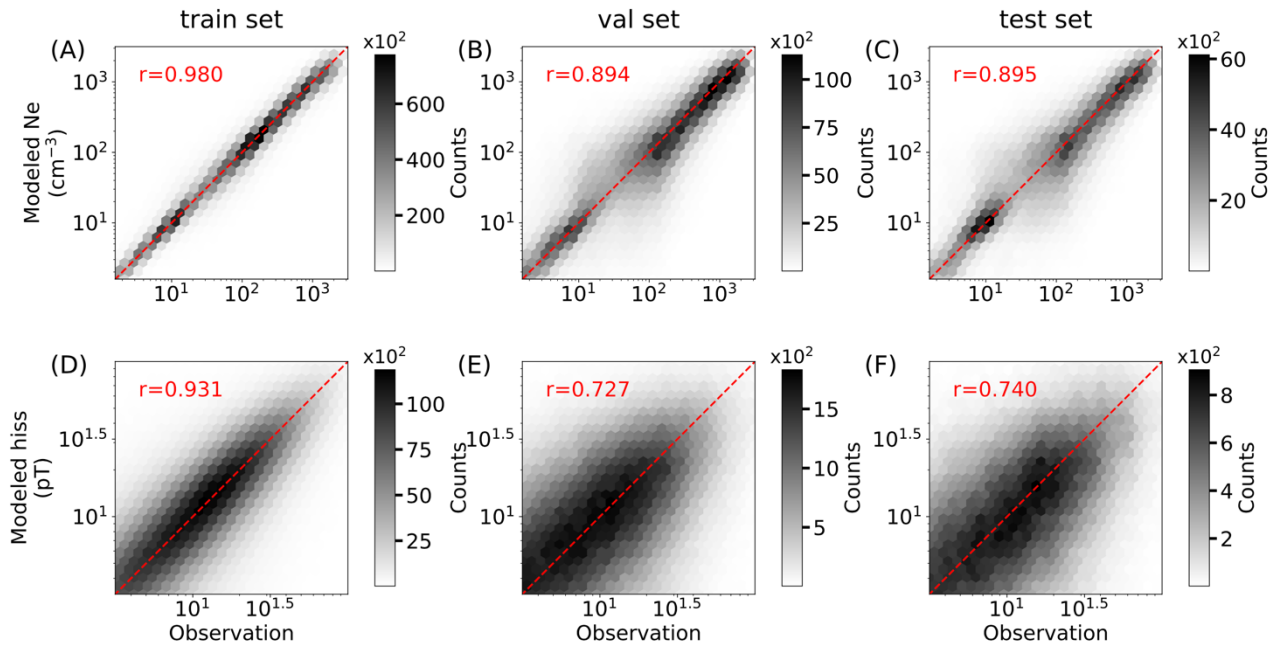


Figure 2. Overall model performance of electron density (A-C) and hiss wave amplitude (D-F) for different datasets. X-axis: observed quantity; Y-axis: modeled quantity. The correlation coefficient is calculated and labeled as “r”. The Red dashed line denotes the data pair where the model prediction matches the observed value perfectly.

3.2 Event Study

We present a case study focusing on the global evolution of hiss waves and evaluate their effects on the energetic electron dynamics during a storm event on 14 May 2019, which is intentionally excluded from the training set. RBSP observations reveal the formation of a plasmaspheric plume and intensification of hiss waves over 13-19 May 2019, as shown in Figure 3. The SYM-H and SML indices (Figure 3A) peaked on May 14 when RBSP was on the dayside and observed a clear signature of the plasmaspheric plume (first by RBSP-A and later by RBSP-B, marked with black arrows in Figure 3B and 3C, respectively). Hiss wave amplitude intensified during the event (Figures 3D and 3E). Panels (F)–(H) show binned satellite observations of energetic electron fluxes at energies of 132 keV, 235 keV, and 470 keV, respectively. The electron flux increased by an order of magnitude from $L \sim 5$ to $L \sim 3$ within several hours during the main phase of the storm, which occurred at 7 UT on May 14. After the storm main phase, the electron flux decayed gradually over the subsequent days due to radial diffusion and pitch-angle scattering by waves that we will model later. We plot the modeled electron density and hiss wave amplitude in panels (B)–(E) and show a line-by-line comparison between the model (orange) and the observation (blue) during the event. Overall, the model accurately captures the evolution of the plasmopause location, especially during the latter half of the event when SYM-H and SML were very quiet while K_p varied. There were instances when RBSP measured very low density ($<10 \text{ cm}^{-3}$), but the model predicted slightly higher density ($\sim 30 \text{ cm}^{-3}$). Although the relative error is significant, the absolute error remains relatively low. The modeled hiss wave amplitude generally follows the observations, successfully capturing most of the peak values.

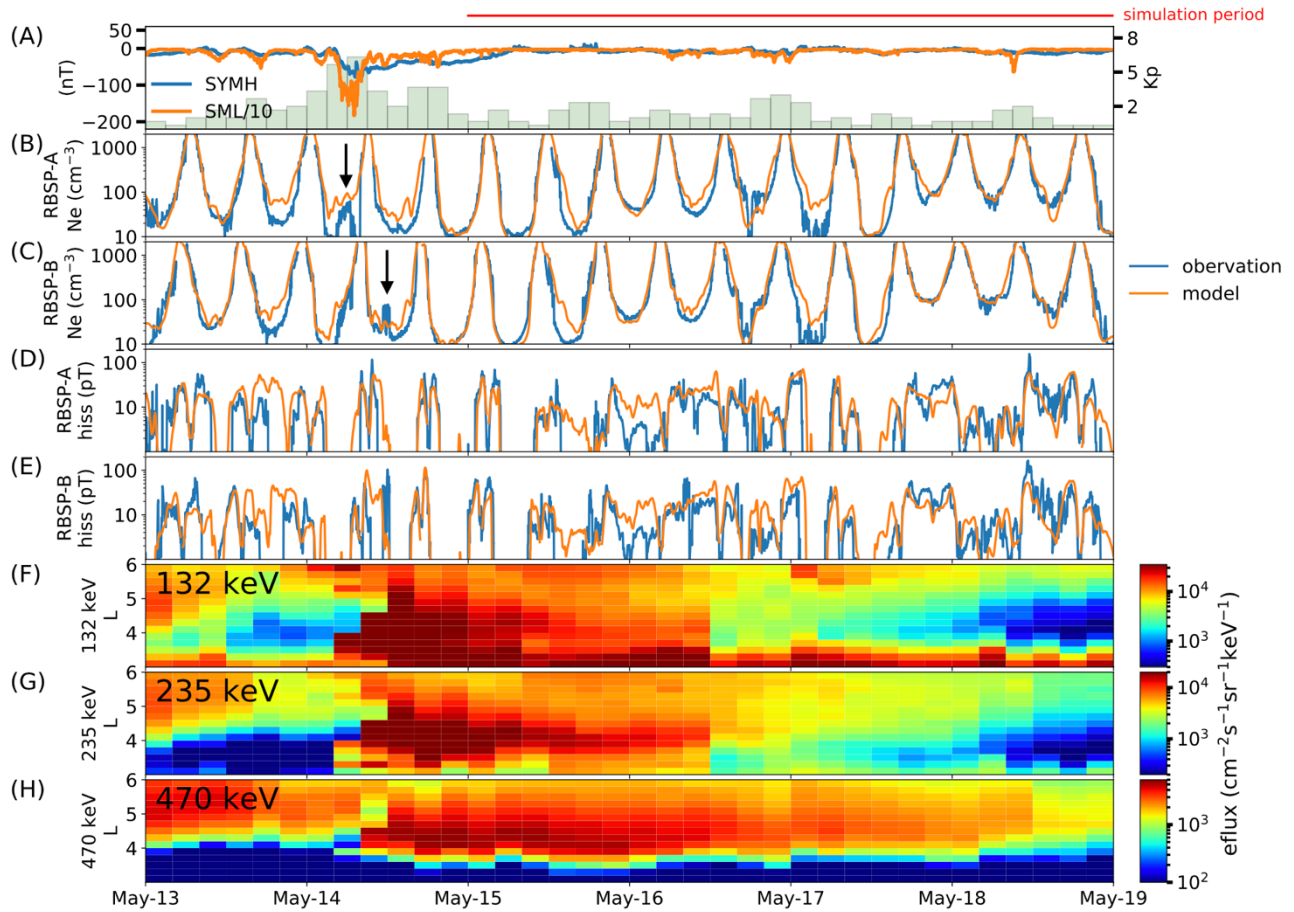
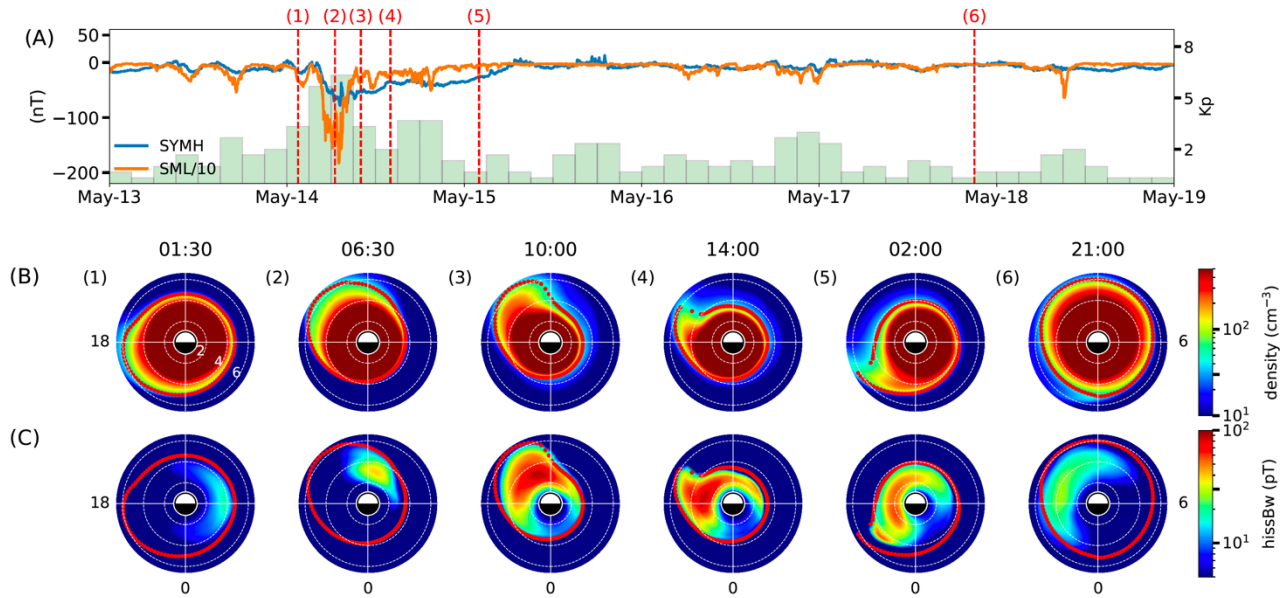


Figure 3. Overview of the geomagnetic storm during 13 – 19 May 2019. (A) SYM-H, SML, and Kp indices during the event. (B-C) Comparison between modeled electron density (orange) and satellite observation (blue) for RBSP-A and -B, respectively. Black arrows indicate plume features observed by the satellites. (D-E) Same as (B-C), but for hiss wave amplitude. (F-H) Measured spin-averaged electron flux at different energy channels.

Figure 4 provides several snapshots illustrating the modeled global evolution of both electron density and hiss waves, allowing for a more comprehensive understanding of their dynamics during and following the storm event. As indicated by SYM-H and SML in panel (A), we select six specific times (1-6) to examine the modeled electron density (B) and hiss wave amplitude (C) before, during, and after the storm. Before the storm onset (1), the plasmasphere was relatively quiet and extended up to $L=6$ on the dusk side. Correspondingly, hiss wave activity was low, which is expected during quiet conditions (Kim et al., 2015; Li et al., 2015). As the storm intensified (2) with higher Kp and decreased SYM-H, the plasmasphere was pushed to the dayside due to the enhanced convection electric field, and hiss waves were intensified in the dawn-to-noon sector, probably related to the enhanced injection from the nightside plasma sheet. As the storm progressed (3), the plasmaspheric plume was formed, and the region with strong hiss waves shifted to the dusk side. The intensified waves predominantly occurred at high L, showing a good spatial correlation with the plume, in agreement with the statistical results of plume hiss (Shi et al., 2019). During the recovery phase from (3) to (5), the model predicted a rotating and narrowing plume, consistent with physical simulation results (De Pascuale et al., 2018), with hiss waves rotating and decaying simultaneously. After the storm, instances of persistent moderate hiss wave activity were observed (6). During the entire

251 period, the majority of the wave power was concentrated near the plasmapause, in agreement with
 252 statistical results (Malaspina et al., 2017).



253

254 **Figure 4.** Snapshots of a geomagnetic storm event during 13 – 19 May 2019. (A) SYM-H, SML, and
 255 Kp indices during the event. (B) Modeled total electron density on the equatorial plane at different
 256 times, indicated by red dashed lines in panel (A). The contour of electron density of 50 cm^{-3} is
 257 overplotted as a red line to indicate the plasmapause. White dashed circles represent $L=2, 4$, and 6 .
 258 (C) Same as panel (B), but for hiss wave amplitude.

259 3.3 Event Simulation

260 We use the UCLA 3-D diffusion code (Ma et al., 2015, 2018) to simulate the energetic electron
 261 evolution, considering radial diffusion and local resonant interactions with hiss waves. The
 262 simulation starts at 00 UT on May 15, following a period of significant local electron acceleration
 263 period and the extension of the plasmapause beyond $L = 4$. During the following four quiet days, the
 264 electron flux gradually decayed, providing a unique opportunity to model the effects of pitch angle
 265 scattering caused by hiss waves. The observed electron fluxes at 00 UT on May 15 are used as the
 266 initial condition for all L shells, as well as the time-varying boundary conditions at $L=2.6$ and $L=6$.
 267 The energy range in the simulation is set from 374 keV to 4.5 MeV at $L=2.6$ and from 40 keV to 1
 268 MeV at $L=6$, maintaining the conservation of the first adiabatic invariant. The pitch angle gradients
 269 of phase space density at $\alpha=0^\circ$ and $\alpha=90^\circ$ are set to be 0. The modeling results of energetic electron
 270 fluxes are not sensitive to the energy boundary condition assumptions because the energy diffusion
 271 coefficients due to hiss are much smaller than the pitch angle diffusion coefficients (e.g., Ni et al.,
 272 2013; Thorne et al., 2013). Radial diffusion coefficients are calculated using the formulation by Liu
 273 et al. (2016) with pitch angle dependence from Schulz (1991, p229). The pitch angle, momentum,
 274 and their mixed diffusion coefficients are computed based on the total plasma density and hiss wave
 275 amplitude obtained from the deep learning model with a time cadence of 5 min. The wave frequency
 276 spectrum is derived from the Van Allen Probes statistics (Li et al., 2015), and wave normal angles are
 277 assumed to be quasi field-aligned near the magnetic equator, gradually becoming highly oblique at
 278 higher latitudes (Ni et al., 2013). The deep learning model provides the time-varying total electron

density and hiss wave amplitude as functions of L shell and MLT at the equator, which are used as inputs to the 3-D diffusion code.

Figure 5 shows the modeled MLT-averaged hiss wave amplitude (A) and the simulated energetic electron flux evolution (B-D) in the same energy channels as shown in Figure 3. At the start of the simulation on May 15, the energetic electron fluxes were initially high in the outer radiation belt. As a result of both radial diffusion and scattering by hiss waves, the electron flux gradually decayed over the following 1-3 days. Instances of faster decay and slumps in the electron flux were successfully reproduced by the simulation at 0 and 18 UT on May 17, consistent with the RBSP observations. These slumps can be attributed to the enhanced wave activity, which causes stronger pitch angle scattering. To quantify the role of plume hiss in energetic electron dynamics, we divided the modeled global distribution of hiss waves into plume hiss and plasmaspheric hiss based on the modeled total electron density. We defined the plume as the region with a total electron density in the range 20–200 cm^{-3} , as identified from the global maps of modeled electron density, in agreement with typical plume statistics (Darrouzet et al., 2008; Moldwin, 2004). Although this definition may include the outer plasmasphere, as well as attached or detached plumes, it serves our purpose as this region exhibits similar characteristics that allow access for energetic electrons, potentially providing a source of free energy for whistler mode wave intensification (e.g., Li et al., 2013; Shi et al., 2019). Figure 5E displays the modeled plume hiss, characterized by an MLT-averaged wave amplitude of $\sim 10\text{--}20$ pT. The majority of the plume hiss was located at $L \sim 5$, although the coverage was sometimes extended to $L > 6$. Despite its high variability, a clear trend emerged during the first three days, indicating that the inner edge of the plume hiss moved from $L = 4$ to 5 due to the refilling of the plasmasphere after the storm.

To assess the impact of plume hiss on energetic electron flux, we conducted simulations considering only plasmaspheric hiss and compared them with simulations that included the effects of both plasmaspheric and plume hiss (the simulated electron fluxes are denoted as J_1 and J_2 , respectively). The difference in electron fluxes between these simulations, quantified by $(J_1 - J_2)/J_1$, represents the sole effect of plume hiss, as shown in panels (F-H). When the plume hiss effect was included, there was a consistent decrease in electron fluxes over the 100–500 keV energy range. After a few days of simulation, the plume hiss accounted for an $\sim 80\%$ decrease in 132 keV electron flux and a $\sim 40\%$ decrease in 470 keV electron flux at $L \sim 4.5$, near the heart of the outer radiation belt. At higher L , the plume hiss also contributed significantly to electron losses, resulting in a $\sim 30\%$ – 70% decrease in electron flux at $L \sim 5.5$. It is worth noting that the hiss wave activity depicted in Figure 5A is relatively modest, but the peak wave amplitude reached up to ~ 100 pT. The averaged value of hiss wave amplitude during the recovery phase of this event is lower than the averaged statistical wave amplitude (~ 100 pT) on the dayside during strong geomagnetic conditions with $AL^* < -500$ nT (Li et al., 2015). It is interesting to note that there have been instances where hiss wave amplitudes in plumes exceeded 1000 pT (Su et al., 2018). Therefore, we expect that plume hiss waves would have a much stronger impact during periods of higher geomagnetic activity.

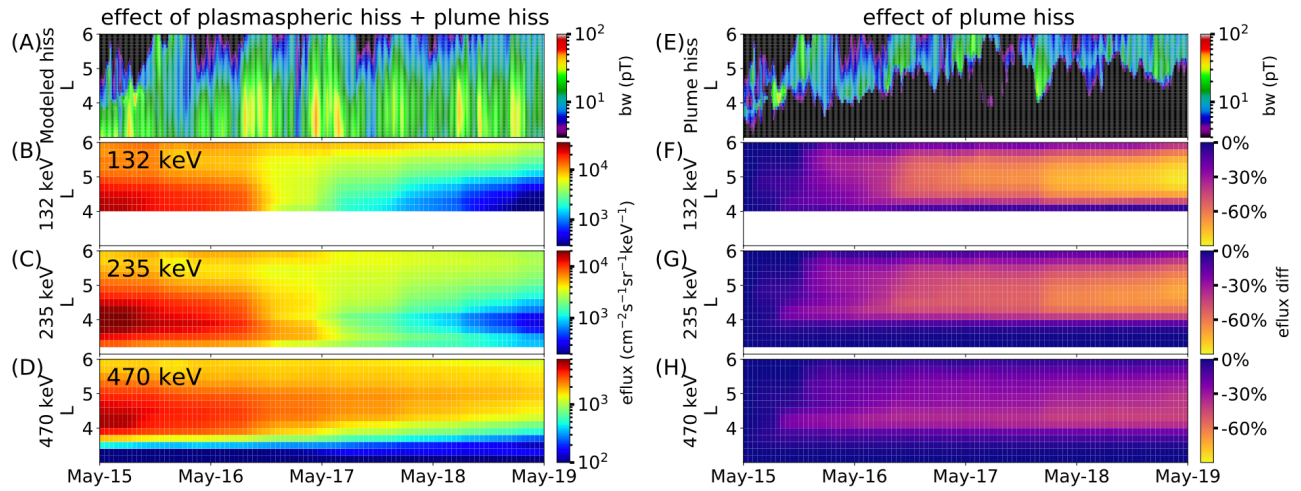


Figure 5. Simulated energetic flux evolution during a quiet period. (A) MLT-averaged hiss wave amplitude as a function of L and time from the deep learning model. (B) Simulated electron flux evolution for 132 keV electrons as a function of L and time, starting at L>4. (C) Same as panel (B) but for 235 keV electrons. (D) Same as panel (B) but for 470 keV electrons. (E) Modeled MLT-averaged plume hiss wave amplitude. (F) Difference in simulated electron flux with and without plume hiss for 132 keV electrons. (G) Same as panel (F) but for 235 keV electrons. (H) Same as panel (F) but for 470 keV electrons.

Figure 6 presents a comparison between the simulated (dashed line) and the observed electron flux evolution (solid line) at L=4.4. This L shell is located in the heart of the outer radiation belt, where the electron flux decay is most prominent. Moreover, choosing L=4.4 ensures that it is sufficiently distant from the simulation boundary, thus the change at this distance is mostly from the simulation itself, minimizing the potential impact of using observations as boundary conditions. In all three energy channels, the simulation exhibits a gradual flux decay from May 15 to 16, followed by a faster decay from 16 to 17. The simulation accurately captures the electron flux decay rate until the end of May 18, when the observation reveals a faster decay of higher-energy electrons. This faster decay could be attributed to the influence of waves other than hiss waves alone, as discussed below.

compare simulated eflux with observation at L=4.4

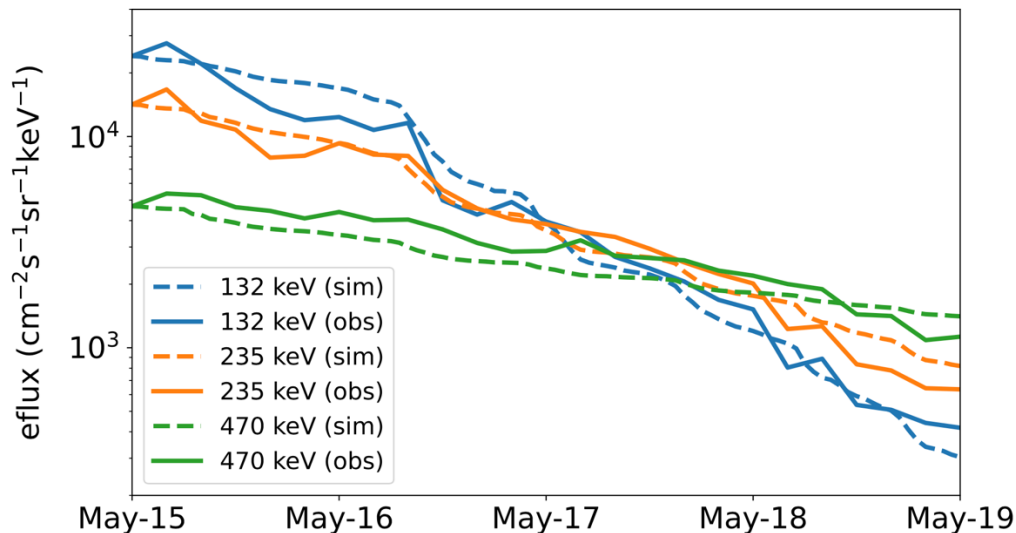


Figure 6. Comparison between the simulated (dashed line) and the observed electron flux evolution (solid line) at $L=4.4$. Each color represents a different energy channel.

4 Discussion

Although the simulated electron flux reproduced the observed flux for most of the period, there was a slightly faster decay rate in the observed flux on the last day of the simulation. Several potential factors could contribute to this discrepancy, which are discussed below.

1. The presence of waves other than hiss waves can affect energetic electron dynamics. For example, chorus waves can also scatter electrons in the energy range of hundreds of keV, especially on the nightside where the plasmopause is often located at $L < \sim 5$. When performing simulations that include both chorus and hiss waves, the effects of these waves will be taken into consideration. However, this is beyond the scope of the present study, as we focus solely on modeling the hiss wave distribution in the plasmasphere and plume and their quantitative scattering effects on electrons.
2. The presence of other waves may not scatter particles directly, but instead enhance the efficiency of hiss waves in scattering energetic electrons into the loss cone. Previous studies have shown that when electromagnetic ion cyclotron (EMIC) waves and hiss waves coexist at the same L shell, MeV electrons can be first scattered by hiss waves and subsequently scattered and precipitated by EMIC waves (Drozdov et al., 2020; Ma et al., 2015), resulting in a significant reduction in their lifetimes (Li et al., 2007; Zhang et al., 2017). Fast magnetosonic waves can induce additional scattering at intermediate pitch angles, leading to increased electron losses compared to scattering by hiss alone (Hua et al., 2018). Non-linear phase trapping by chorus waves can accelerate 300-500 keV electrons, which may then resonate with EMIC waves, resulting in their rapid scattering into the loss cone (Bashir et al., 2022). The combined effects of different wave modes on the radiation belt dynamics are beyond the scope of the present study and are left for future investigations.

There are different ways to define plumes used in simulations. In our study, we define the plume region as an area with a total electron density ranging from 20 to 200 cm^{-3} at $L < 6$. This definition typically encompasses the outer plasmasphere or the plume, where energetic electrons ($> \sim 10$ s keV) can access, thus leading to highly variable wave activity over time and space. We have found a considerable amount of hiss wave power at $L > 4$, and the outermost extension of hiss waves has been observed to vary from $L = 4$ to 6, even during relatively quiet periods indicated by the geomagnetic indices. The commonly used density and wave statistical models, which are often expressed as simple functions of Kp and/or AE (Golden et al., 2012; O'Brien & Moldwin, 2003; Saikin et al., 2022; Spasojevic et al., 2015), do not capture such variability since the underlying geomagnetic indices might not exhibit strong variations during the period. These statistical models predict a constant wave power at a given location for a range of geomagnetic indices. Our findings demonstrate that even under relatively quiet conditions, hiss wave activity could exhibit dynamic evolution, and such spatial variation plays a crucial role in the evolution of energetic electron fluxes over time at different L shells, as shown in Figure 5.

5 Conclusions

We have developed a neural network model to simultaneously reconstruct the global evolution of both electron density and hiss wave amplitude in the Earth's plasmasphere and plume. Unlike

377 traditional deterministic models, our approach estimates the distribution of these quantities, allowing
378 for a better representation of variations in the data on both large and small scales.

379 To quantify the evolution and effects of plume hiss, we focused on the storm event that occurred over
380 13 – 19 May 2019, during which RBSP observed the formation of a plasmaspheric plume, followed
381 by a gradual decay in the electron fluxes at a few hundred keV. Our model successfully captured the
382 global evolution of the plume, as well as the plume hiss within it during the entire event. As
383 geomagnetic activity increased, hiss wave power intensified and shifted from dawn to dusk, where
384 the plume was formed later. The plume and plume hiss exhibited a strong spatial correlation and
385 rotated together as the geomagnetic activity became weaker. The plume wrapped around the Earth
386 and became thinner over the nightside, where hiss wave power diminished rapidly. During the
387 recovery phase, the plasmasphere was gradually refilled, and hiss wave activity remained relatively
388 low in general. Our model provided valuable insights into the relationship between the plume
389 structure (as seen in the plasma density) and plume hiss on a global scale.

390 To quantify the impact of plume hiss, we separated the modeled total hiss wave population into
391 plasmaspheric hiss and plume hiss, and simulated the energetic electron flux evolution with and
392 without plume hiss. By including both plasmaspheric and plume hiss, together with radial diffusion,
393 the simulated electron flux decay reproduces the observation very well. The remaining differences in
394 the electron flux decay may be attributed to scattering effects from other waves. Although the MLT-
395 averaged wave amplitude was $\sim 10\text{--}20$ pT, plume hiss alone was responsible for an additional $\sim 80\%$
396 decrease in 132 keV electron flux at $L\sim 4.5$ within 3 days, and $\sim 30\%$ decrease in 470 keV electron
397 flux at $L\sim 5.5$. These results highlight the dynamic nature of hiss wave evolution even during
398 geomagnetically quiet conditions, and emphasize the significant role played by plume hiss in shaping
399 the energetic electron dynamics, especially in the outer radiation belt, which should be considered in
400 future simulations of radiation belt dynamics.

401

6 Data Availability Statement

The Van Allen Probes data from the EMFISIS instrument were obtained from <http://emfisis.physics.uiowa.edu/Flight/>. Data from the ECT instrument were obtained from https://rbsp-ect.newmexicoconsortium.org/data_pub/. The geomagnetic indices used in the model training are available at https://omniweb.gsfc.nasa.gov/form/omni_min.html (SYM-H); <https://supermag.jhuapl.edu> (SML and SMU); <https://www.gfz-potsdam.de/en/hpo-index/> (Hp30). All data used to produce figures, as well as the Python script defining the model structure, are publicly available at <https://doi.org/10.6084/m9.figshare.22817531>

7 Author Contributions

SH and WL developed the study concept and lead the project. SH designed, implemented, and trained the neural network model, and performed the event analysis. QM performed the Fokker-Planck simulation and produced Figure 5. XS processed the wave spectrum data from EMFISIS and generated the dataset of hiss. LC contributed to the model design and event analysis. XNC, DM, JB, MH, and SW contributed to the model design. SH wrote the first draft of the manuscript. All authors contributed to the discussion of the project and edited the manuscript.

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702 **11 Appendix A: Model structure and optimization procedure**

703 We optimize the hyperparameters in our model following the steps described by Huang et al. (2022).

704 After careful tuning, we used the following set of optimal hyperparameters for our model:

- 705 a) 2 LSTM layers, each with a size of 256.
- 706 b) To output with an estimation of mean, 5 fully connected layers with each of size (260, 128, 128,
707 128, 128, 1) and SELU as activation function are applied.
- 708 c) To output with an estimation of standard deviation, 5 fully connected layers with sizes (260, 128,
709 128, 128, 128, 1) and SELU as activation function are applied, with an additional soft-plus
710 operation that converts the output to be positive.
- 711 d) The encoder length is 24 hours.
- 712 e) The decoder length is 10 hours.

713 The detailed script that defines the model structure and weighted sampler can be found in the file
714 uploaded in the figshare archive.