

Optical spectroscopic detection of Schottky barrier height at a two-dimensional transition-metal
dichalcogenide/metal interface

*Du Chen^{1,2}, Surendra Babu Anantharaman³, Jinyuan Wu^{2,4}, Diana Y. Qiu^{2,4}, Deep Jariwala³,
Peijun Guo^{1,2,*}*

¹Department of Chemical and Environmental Engineering, Yale University, New Haven, CT
06520

²Energy Sciences Institute, Yale University, West Haven, CT 06516

³Department of Electrical and Systems Engineering, University of Pennsylvania, Philadelphia, PA
19104

⁴Department of Mechanical Engineering and Materials Science, Yale University, New Haven, CT
06520

Abstract

Atomically thin two-dimensional transition-metal dichalcogenides (2D-TMDs) have emerged as semiconductors for next-generation nanoelectronics. As 2D-TMD-based devices typically utilize metals as the contacts, it is crucial to understand the properties of the 2D-TMD/metal interface, including the characteristics of the Schottky barriers formed at the semiconductor-metal junction. Conventional methods for investigating the Schottky barrier height (SBH) at these interfaces predominantly rely on contact-based electrical measurements with complex gating structures. In this study, we introduce an all-optical approach for non-contact measurement of the SBH, utilizing high-quality WS₂/Au heterostructures as a model system. Our approach employs a below-bandgap pump to excite hot carriers from the gold into WS₂ with varying thicknesses. By monitoring the resultant carrier density changes within the WS₂ layers with a broadband probe, we traced the dynamics and magnitude of charge transfer across the interface. A systematic sweep of the pump wavelength enables us to determine the SBH values and unveil an inverse relationship between the SBH and the thickness of the WS₂ layers. First-principles calculations reveal the correlation between the probability of injection and the density of states near the conduction band minimum of WS₂. The versatile optical methodology for probing TMD/metal interfaces can shed light on the intricate charge transfer characteristics within various 2D heterostructures, facilitating the development of more efficient and scalable nano-electronic and optoelectronic technologies.

Keywords

Transition-metal dichalcogenides, Schottky junction, two-dimensional materials, pump-probe spectroscopy, transient reflection

Introduction

Transition metal dichalcogenides (TMDs) are a family of emerging two-dimensional (2D) materials featuring many attractive properties. Compared with previously well-established 2D materials like graphene and hexagonal boron nitride (*h*-BN), the presence of appropriate band gaps in TMDs make them promising candidates for various electronic and optoelectronic devices, such as field-effect transistors,^{1,2} valleytronics,^{3,4} photodetectors,⁵⁻⁷ photovoltaics,^{8,9} and light-emitting diodes.¹⁰⁻¹³ One common issue limiting the performance of these applications is the contact resistance at the TMD-metal interface,^{14,15} which mainly originates from the Schottky barrier formed between the TMDs and metals.^{16,17}

The formation of the Schottky barrier is primarily attributed to the energy band alignments between the TMDs and metals,¹⁴ and metal-induced gap states.¹⁸ A minimal or even zero Schottky barrier height (SBH) is favorable in most applications. Although many theoretical studies and calculations are available,¹⁹⁻²² experimental measurement of the SBH often requires a complex gated-transistor structure,^{18,23-25} or specialized atomic force microscopies.^{14,15,26} The extraction of SBH by the above-mentioned methods is based on the electrical I-V curve or contact potential measurement. Here, we demonstrate an optical, pump-probe based approach to achieve non-contact measurement of the SBH, using WS₂-Au interface as a model system.

Gold is one of the most widely used electrode materials and is known to form a Schottky contact with sulfide-based TMDs, *e.g.*, MoS₂ and WS₂,^{14,24,27,28} since gold has a work function of approximately 5.3 eV,²⁹ exceeding the electron affinity of WS₂ (~4.5 eV).³⁰ This energy difference leads to an SBH of approximately 0.8 eV. In a heterostructure formed by WS₂ and gold, electrons in the gold experience a potential barrier and are hindered from flowing into the WS₂. It has been reported that light can excite electrons with energies large enough to overcome the barrier and be injected into the WS₂, which then form excitons or recombine with the intrinsic holes within the WS₂.^{28,31} In this communication, we used near-infrared (NIR) pump that lies below the bandgap of WS₂ to induce hot-electron injection from gold to WS₂. Owing to the strong absorption bleaching by photoexcited free carriers and excitons,³² a broadband probe is applied to monitor the carrier populations within the WS₂. By sweeping the energy of the pump photons, we observed certain thresholds, below which the energy of the pump photon is insufficient to induce significant amount of electron injection into WS₂, permitting the extraction of the SBH. Our work

demonstrates an alternative, all-optical, non-contact method to measure the SBH between TMDs and metals, and sheds light into the electron transport channels at the TMD-metal interface.

Results and Discussion

We mechanically exfoliated monolayer and few-layer WS_2 onto smooth gold substrates (see Fig. S1a). The thickness of the flakes was determined using atomic force microscopy (AFM), revealing layer numbers of 1, 5, and 10 (see Fig. S1b), assuming a monolayer thickness of 0.6 nm.³³ Pump-probe transient reflection (TR) experiments were carried out on these samples under ambient conditions. We utilized an optical objective to focus probe pulses onto the sample with a sub-5- μm spot size. The pump source, ranging from 1100 nm (1.12 eV) to 2700 nm (0.46 eV), was generated using an optical parametric amplifier (OPA) powered by a Pharos amplifier. The Pharos amplifier produces an output wavelength of 1030 nm and an output pulse duration of 170 fs at a repetition rate of 2 kHz. Supercontinuum probe pulses (530~720 nm) were generated by focusing a small portion of the 1030-nm fundamental output onto a 4-mm thick YAG crystal (see Fig. S2). Steady-state optical reflectance was measured using the same pump-probe microscopy setup, with the pump light turned off. All the optical measurements were performed at ambient conditions.

We observed two distinct resonant dips in the reflectance spectra (Fig. 1a), corresponding to the A and B excitons in WS_2 at the K point. These resonant dips display a redshift with increasing layer number, consistent with findings from a previous band-structure study on WS_2 .³⁴ The TR experimental results for a 1L- WS_2 /Au sample under 400-nm excitation are presented in Fig. 1b, expressed in terms of $\Delta R/R$. Here, $\Delta R/R$ is defined as $[R(t) - R(0)]/R(0)$, where $R(0)$ represents the reflectance without pump excitation, and $R(t)$ is the reflectance at delay time t after the pump excitation. The kinetics of $\Delta R/R$ extracted at the resonant dip wavelength, as shown in Fig. 1b, exhibits a typical photobleaching response. This suggests the rapid formation of free electron-hole pairs within the first 500 fs, followed by the creation of excitons from thermally equilibrated hot carriers.³⁵ The TR results distill a high sensitivity of reflectance to an excess density of charge carriers and excitons, indicating the potential for detecting additional charge carriers injected from gold to WS_2 when overcoming the Schottky barrier.

We then performed TR experiments on both an Au/glass sample and the 1L- WS_2 /Au/glass sample using a pump wavelength of 1100 nm, and the resulting transient spectral maps of $\Delta R/R$ are shown in Fig. 2a and 2b, respectively. The pump photon energy of approximately 1.1 eV is

notably below the bandgap of WS₂.³⁴ In Fig. 2a, a robust and broad transient $\Delta R/R$ signal is observed on the bare gold film. Fig. 2b presents the $\Delta R/R$ transient spectral map collected on the 1L-WS₂/Au sample. Notably, the $\Delta R/R$ comprises a broadband bleaching response consistent with the TR result obtained from gold, along with a prominent positive lobe of $\Delta R/R$ centered at approximately 642 nm, corresponding to the WS₂ A-exciton resonance.

To isolate the true transient response originating from 1L-WS₂, we subtracted the $\Delta R/R$ response (Fig. 2a) of the bare gold film from that of 1L-WS₂/Au (Fig. 2b). The resultant $\Delta R/R$ spectral map, shown in Fig. 2c, distills the distinct feature arising from the monolayer WS₂ alone. We extracted the temporal kinetics of $\Delta R/R$ using the data in Fig. 2c at the wavelength of 640 nm and present the results in Fig. 2d. Interestingly, our analysis shows that the ‘background-corrected’ $\Delta R/R$ kinetics of 1L-WS₂ exhibits a biexponential decay, resembling the response with above-bandgap pump excitation (see Fig. 1b). Similar transient reflectance spectral are also obtained for both 5L-WS₂/Au and 10L-WS₂/Au samples as illustrated in Fig. S3. All three samples with different thickness of WS₂ feature a band renormalization effect at the first 1 ps that can be attributed to a high concentration of excess free electrons present. This finding suggests that the TR response under 1100-nm pump excitation can be ascribed to pump-induced charge carriers in WS₂.^{28,36-38} To explore this further, we conducted pump power-dependent measurements on both the 1L-WS₂/gold sample and a sample with few-layer WS₂ on gold. As summarized in Fig. S4, we observed a linear relationship between the peak value of $\Delta R/R$ and the pump fluence. This linear relationship rules out the possibility of multi-photon absorption of the 1100-nm pump photons by WS₂, thereby confirming the charge injection-induced $\Delta R/R$ response from gold to WS₂ with 1100-nm pump excitation.

In a low-perturbative regime, the alteration in reflectance of a semiconductor near its bandgap can be described as $\Delta R = \frac{\partial R}{\partial E_g} \frac{dE_g}{dT} \Delta T + \frac{\partial R}{\partial N} \Delta N$, where ΔT denotes the temperature rise, E_g is the bandgap, and ΔN represents the change in the charge-carrier concentration of the semiconductor.³⁹ Following the pump excitation, the initial response of $\Delta R/R$ is dominated by the dynamics of charge carrier, succeeded by the slower decay of lattice temperature. In the case of the 1L-WS₂/Au sample, the lattice heating in both gold and WS₂ is evident in the long-lived $\Delta R/R$ signal, persisting for hundreds of picoseconds (ps) after pump excitation, as illustrated in Fig. S5. This extended feature of $\Delta R/R$, emerging relatively late and peaking around 1000 ps, arises from a

redshift of the exciton resonance of WS₂ with an increasing temperature.⁴⁰ This temperature rise in WS₂ is induced by heat transfer from the nearby photothermally excited gold, as well as relaxation of hot carriers injected from gold.

In the initial time window (*i.e.*, the first 20 ps after pump excitation), the ‘background-corrected’ signal of $\Delta R/R$ is primarily influenced by the change in charge carrier concentration in WS₂. Consequently, the rapid rise of $\Delta R/R$ following pump excitation is associated with a growing concentration of injected electrons in the WS₂. A similar response of $\Delta R/R$ was observed for three selected samples (*i.e.*, 1L, 5L and 10L-thick WS₂ on gold). We then extracted ‘background-corrected’ $\Delta R/R$ kinetics at the resonance reflectance dips for the three samples, namely 630 nm, 635 nm, and 640 nm, as illustrated in Fig. 1a. Informed by earlier studies, which demonstrate that the charge transfer between WS₂ and graphene occurs within the first 20 fs,⁴¹ we anticipate that the charge injection dynamics cannot be resolved by our TR experiments. Nevertheless, as summarized in Fig. 3, the $\Delta R/R$ kinetics display a rapid rise upon pump excitation, followed by a bi-exponential decay comprising both fast and slow components with distinct time constants.

We fitted the $\Delta R/R$ kinetics using a bi-exponential function, $\frac{\Delta R}{R} = Ae^{-\frac{t}{\tau_f}} + Be^{-\frac{t}{\tau_s}} + C$, where τ_f and τ_s represent the time constants for the fast and slow components, as summarized in Fig. 3d. The fast decay time constant τ_f , ranging from 70 to 100 fs, can be attributed to the cooling of injected hot electrons from a non-thermalized to a thermalized distribution. This timescale of τ_f aligns with prior literature reports on hot carrier thermalization in few-layer TMDs.^{42,43} Briefly, when the pump photon energy exceeds the SBH (denoted as Φ_B), the injected electrons initially possess a non-thermal Gaussian distribution determined by the difference between pump photon energy and Φ_B . The hot electrons undergo carrier-carrier scattering and rapidly reach a Fermi-Dirac distribution in the low-dimensional WS₂ with reduced dielectric screening.

After relaxing to the conduction band edge, the thermalized electrons in WS₂ traverse the WS₂/Au interface, eventually decaying back to the gold substrate. This process is captured by the slow time constant, τ_s , which ranges from 1.5 to 3.4 ps and decreases with an increasing layer number of WS₂. Similar processes have been observed in WS₂/graphene heterostructures, where a carrier back-flow from WS₂ to graphene takes approximately 1 ps.^{36,37,44-46} It is worth noting that the interlayer charge transfer time τ_s is, unlike τ_f , influenced by the interface of WS₂ and gold

rather than WS₂ alone. The faster charge transfer observed in thicker WS₂ can be attributed to improved contact between WS₂ and gold, as well as a lower Schottky barrier, as discussed later. A more prolonged $\Delta R/R$ signal, persisting for hundreds to thousands of ps, follows the initial two exponential decay components. We attribute to this long-lived $\Delta R/R$ component to the rise in lattice temperature in WS₂ due to hot carrier relaxation, and potentially the extended trapping of electrons at the surface states.³⁷ In the subsequent discussion, our focus is solely on the first 20 ps, which provides insights into the charge injection process of interest. We note that at this early time range, the significant heating effect in WS₂ caused by the heat transfer from gold can be neglected.

As discussed earlier, the SBH at the WS₂/Au interface is approximately 0.8 eV. This determination is drawn from literature reports on a threshold behavior in the current-voltage curves. Specifically, substantial drain-source currents were only detected when the bias voltage between the TMD and gold exceeded the Schottky barrier.^{15,18,24,25,47} Here, we hypothesize that the efficiency of pump-induced gold-to-WS₂ charge injection varies with the pump photon energy. By sweeping the pump photon energy from above to below the SBH, we expect the charge injection efficiency to decrease substantially, thereby permitting the determination of the SBH.

We then performed TR experiments on the three WS₂/Au samples (*i.e.*, 1L, 5L and 10L-thick WS₂ on gold) with systematically varied pump wavelengths under a fixed pump fluence. The ‘background-corrected’ $\Delta R/R$ kinetics are summarized in Fig. 4. Notably, the 1100-nm pump excitation induces $\Delta R/R$ signals in all three samples, and the magnitude of $\Delta R/R$ decreases with increasing pump wavelength. Importantly, we observed different threshold behaviors for the three samples. In the case of 1L-WS₂/Au, the charge-transfer-induced $\Delta R/R$ signal is barely recognizable under 1200-nm pump wavelength (Fig. 4a). However, for 5L-WS₂/Au, we observe $\Delta R/R$ even at 2000-nm pump excitation (Fig. 4b). For 10L-WS₂/Au, a pump wavelength of 2400 nm still results in a discernable $\Delta R/R$ signal associated with charge carrier injection. Since the absolute magnitude of the $\Delta R/R$ signal at the initial time (<0.5 ps) is proportional to the concentration of injected electrons, these results imply that the three WS₂ sample of different thicknesses possess different SBH, with thicker (thinner) WS₂ forming a lower (higher) Schottky barrier with gold.

To extract the SBH, we plotted the trend of integrated $\Delta R/R$ with the pump photon energy in Fig. 4d-f for the three samples. In these plots, we integrated $\Delta R/R$ over a short, early time window

(*i.e.*, the first 4 ps) to exclude effects due to interface-trapped electrons and material heating. As expected, a threshold behavior in the pump photon energy is observed for all three samples, and no $\Delta R/R$ signal and the associated charge injection are observed below the respective threshold pump photon energy for each sample. Based on the cut-off pump photon energy, we can extract the SBH for the three samples with 1L, 5L and 10L WS₂ to be approximately 0.85 eV, 0.65 eV and 0.50 eV. Note that the measurements for each sample under a certain pump photon energy are repeated three times and then averaged together. The standard errors of the integrated kinetics (see Fig. S6) indicate a high reproducibility.

In addition to the decreasing trend of SBH with the number of WS₂ layers, a distinct kink is evident in Fig. 4d and 4f for 1L-WS₂ and 5L-WS₂, respectively. In these cases, the rise in integrated $\Delta R/R$ with pump photon energy displays two different slopes. Notably, such a trend is not observed in the thicker, 10L-WS₂. We attribute this observation to the evolution of the WS₂ band structure with increasing layer number, as the SBH at the WS₂/Au interface correlates with the conduction band of WS₂. Considering a simple Schottky barrier formed at this interface, neglecting surface states and Fermi level pinning, the SBH of the junction, Φ_B , depends solely on the difference between the work function of gold (W_{Au}) and the electron affinity of WS₂ (χ_e), given by $\Phi_B = W_{Au} - \chi_e$.^{15,48} WS₂ undergoes a crossover from an indirect to a direct bandgap as the thickness increases from a monolayer to multilayer. Consequently, the bandgap experiences a redshift with an increasing layer number due to the lowering of the local conduction band minimum (CBM) at the midpoint between the Γ and the K points of the Brillouin zone, namely the Q valley.^{15,34,49} This shift of CBM with layer number is responsible for the observed decrease in Φ_B from TR experiments on thicker WS₂ interfaced with gold. We note that similar trend of SBH with the number of TMD layer has been reported on the MoS₂/Au interface, where the Φ_B decreases from 0.8 eV to 0.2 eV when the layer number goes from one to five.¹⁵

We hypothesize that the observed kinks in the 1L-WS₂/Au and 5L-WS₂/Au samples (Fig. 4d and 4e), contrasted with their absence in the 10L-WS₂/Au sample (Fig. 4f), stem from variations in the density of states (DOS) near the CBM of WS₂. Specifically, we anticipate that the injection probability of pump-excited hot electrons from gold to WS₂, surmounting the Schottky barrier, should exhibit a linear dependence on the available states in WS₂. To test this hypothesis, we performed density functional theory (DFT) calculations of the band structures and DOS of 1L- and

5L-WS₂, as presented in Fig. 5a and 5b, respectively. In the DFT results, the CBM was chosen as the zero-energy point, hence the vertical axis indicates the excess energy of injected hot electrons. The maximal excess energy of injected hot electrons is equal to the pump photon energy minus the SBH.

Assuming equal probability of injection for hot electrons into any lower-lying energy states, we integrated the DOS with respect to energy below the maximum excess energy level. This integration yields the total number of states available for electron injection. The energy cutoff in the integration corresponds directly to the pump photon energy. The results of this integration are shown in Fig. 5c for 1L-WS₂ and in Fig. 5d for 5L-WS₂. Notably, and in line with the TR experiments, a distinctive kink is evident in the plot of integrated DOS versus energy for 1L-WS₂, stemming from the peak in the DOS corresponding to an M₀-type van Hove singularity occurring at the six-fold degenerate valley along the Λ high-symmetry path from Gamma to K (sometimes referred to as the Q valley). In contrast, such a kink is absent in the plot for 5L-WS₂, attributable to its smoothly varying DOS with energy. Qualitatively, this disparity between 1L-WS₂ and 5L-WS₂ arises from the near continuum of conduction bands in the latter, providing a greater, and more smoothly varying pool of available DOS for hosting the injected electrons.

Conclusion

In summary, we have introduced and validated an all-optical method for probing the SBH at interfaces between 2D-TMD and metals, utilizing the WS₂/Au interface as a representative model system. Through nuanced comparison of transient reflection results obtained from WS₂/Au and a bare gold substrate, we have identified the spectral and temporal signatures associated with interfacial charge transfer. By systematically varying the pump wavelength across the near- to mid-infrared range, traversing the SBH, we have quantified the SBH for WS₂/Au samples with varying WS₂ thicknesses from monolayer to 10 layers, and discussed how the charge injection efficiency correlates with the density of states in WS₂. Importantly, this non-contact approach is versatile and can be generalized for investigating the energy alignment of diverse 2D-TMD/metal interfaces, providing insights into their interfacial quality, and advancing the understanding and control of 2D materials in the context of nanoelectronic and nanophotonic applications.

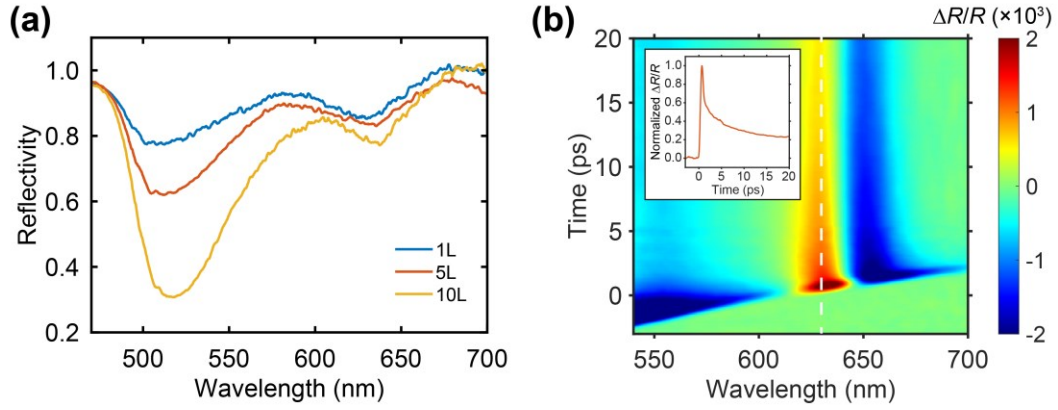


Figure 1. (a) Steady-state reflectance spectra of 1L, 5L and 10L-WS₂ on gold; the reflectance were normalized by the gold substrate. (b) Transient $\Delta R/R$ spectral map of 1L-WS₂/Au excited by above-bandgap pump with a wavelength of 400 nm. Note that the chirp of the TR data hasn't been corrected, which however doesn't change the conclusion. The inset shows the kinetics of $\Delta R/R$ extracted at 640 nm, corresponding to the positive peak of photon bleaching features (indicated by the white-dashed line in the figure).

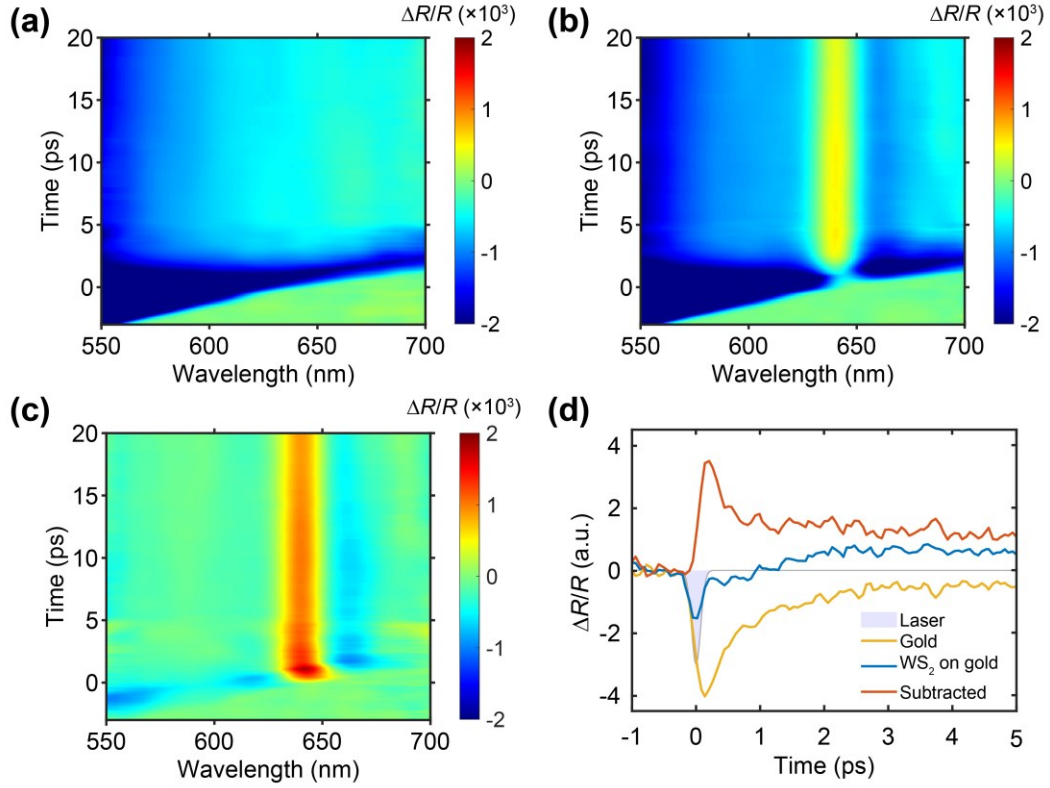


Figure 2. (a) Transient $\Delta R/R$ spectral map of the bare gold substrate obtained with pump excitation at 1100 nm. (b) Transient $\Delta R/R$ spectral map for the 1L-WS₂/Au sample with a pump wavelength of 1100 nm. (c) 'Background-corrected' transient $\Delta R/R$ spectral map for the 1L-WS₂/Au sample obtained by subtraction of the gold substrate response. (d) Kinetics of $\Delta R/R$ extracted at the positive peak of WS₂ for bare gold substrate (yellow line) and 1L-WS₂/Au (blue line), yielding the 'background-corrected' kinetics (orange line). The instrument response is shown by the gray area.

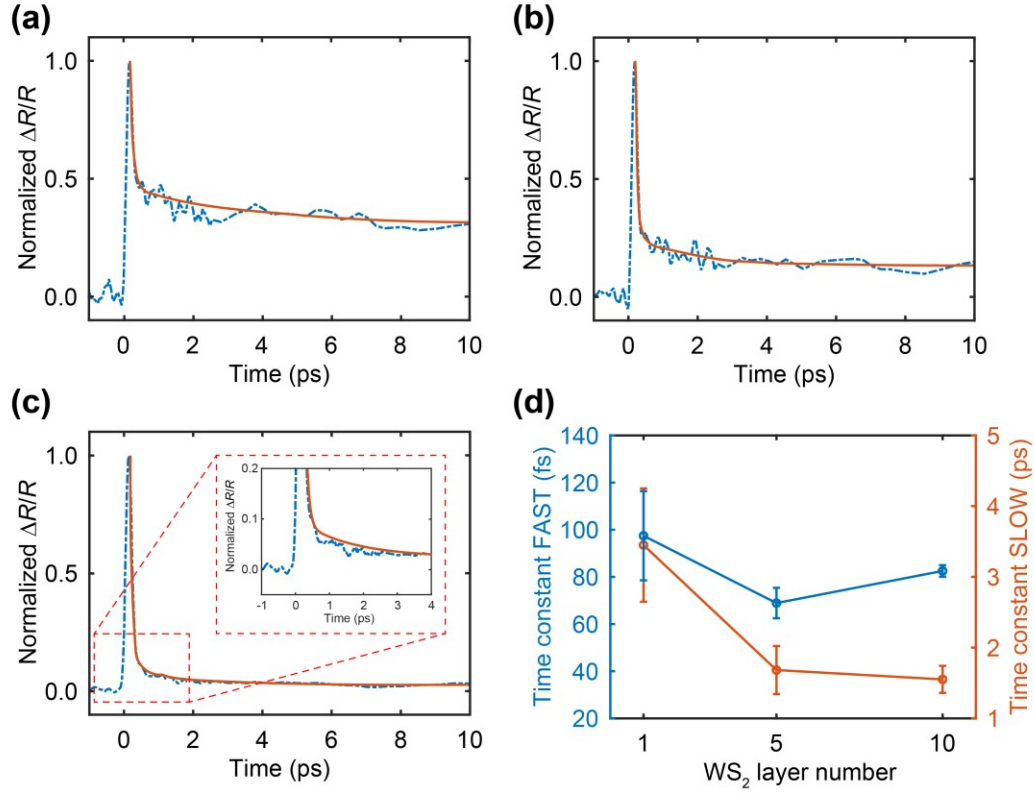


Figure 3. Normalized $\Delta R/R$ kinetics and its bi-exponential fitting for (a) 1L-WS₂/Au, (b) 5L-WS₂/Au, and (c) 10L-WS₂/Au extracted at the corresponding B exciton wavelengths. The inset in (c) shows the zoom-in view of the $\Delta R/R$ signals after the initial 1 ps. (d) The fitted exponential time constants for the fast (left axis) and slow (right) components of the dynamics for the three samples.

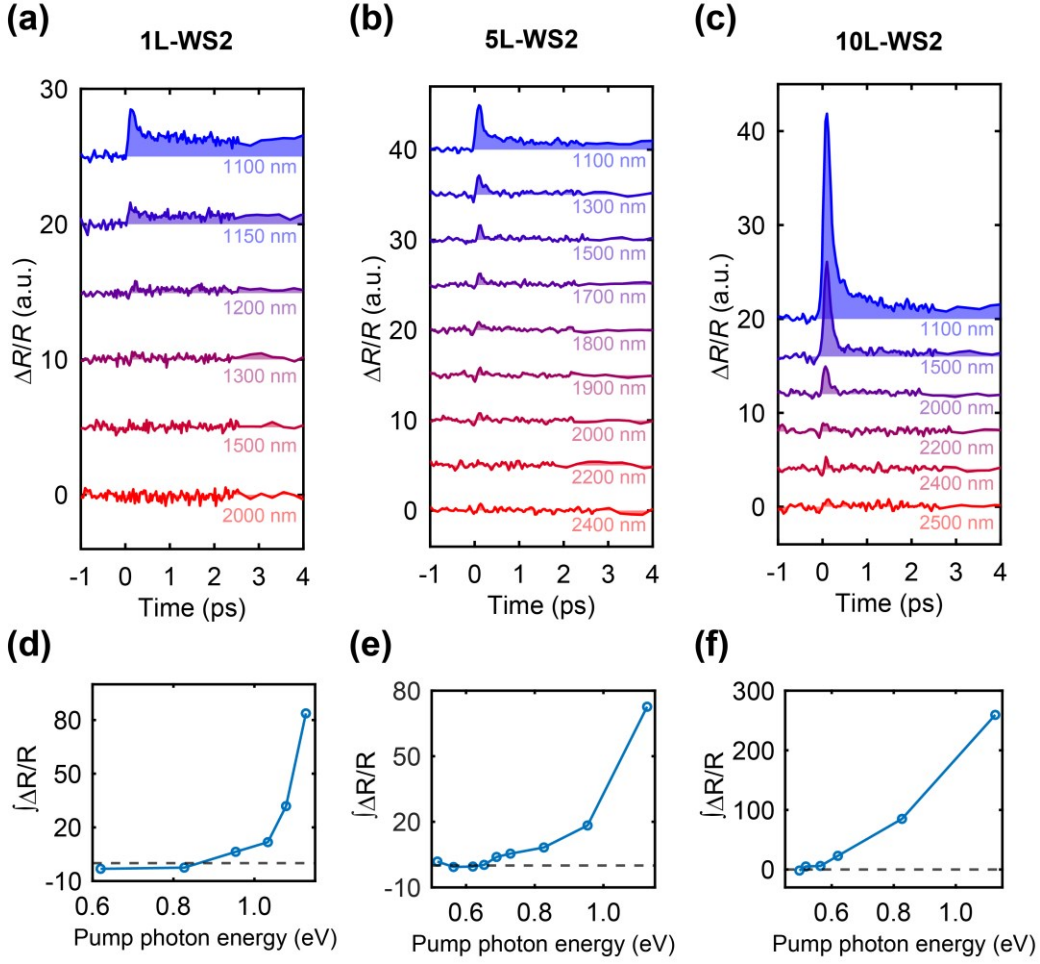


Figure 4. The pump-wavelength dependent kinetics of (a) 1L-WS₂/Au (b) 5L-WS₂/Au (c) 10L-WS₂/Au extracted at the corresponding B exciton wavelengths. The kinetics are offset for clarify of presentation. The $\Delta R/R$ kinetics are integrated over the first 4 ps for (d) 1L-WS₂/Au, (e) 5L-WS₂/Au, and (f) 10L-WS₂/Au, with pump wavelength converted to pump photon energy in eV.

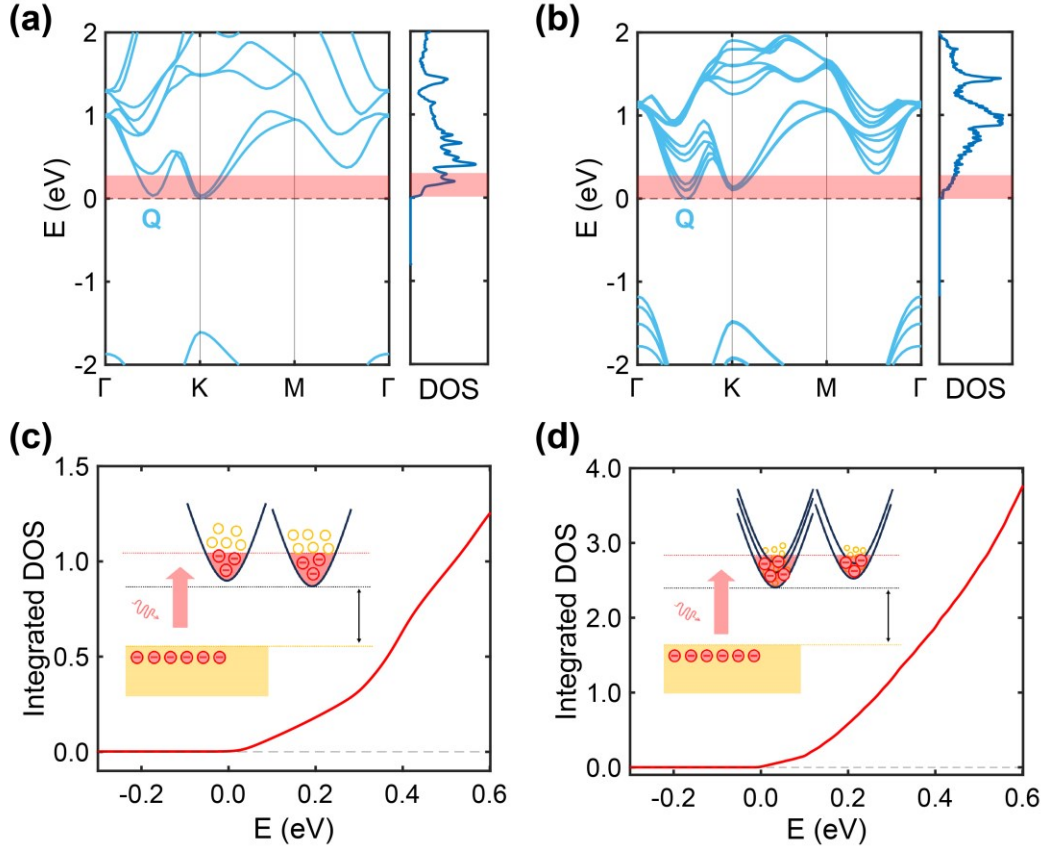


Figure 5. Calculated DFT band structures and conduction-band density of states (DOS) of (a) 1L-WS₂ and (b) 5L-WS₂. The conduction band minimum at each diagram is set as the zero-energy point. The shadow indicates the available states for injected electrons with energy larger than the Schottky barrier. The DOS integrated over the shadowed energy window are shown in (c) for 1L-WS₂ and (d) for 5L-WS₂.

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Author contributions

P.G. and D.J. conceived the project. D.C. performed optical measurements. S.B.A. and D.J. fabricated the samples. J.W. and D.Y.Q. performed DFT calculations. D.C. wrote the manuscript, with inputs from co-authors.

Conflicts of Interest

There are no conflicts of interest to declare.

Correspondence

peijun.guo@yale.edu

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