

Generation of Large Vortex-Free Superfluid Helium Nanodroplets

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Superfluid helium nanodroplets are an ideal environment for the formation of metastable, self-organized dopant nanostructures. However, the presence of vortices often hinders their formation. Here, we demonstrate the generation of vortex-free helium nanodroplets and explore the size range in which they can be produced. From x-ray diffraction images of xenon-doped droplets, we identify that single compact structures, assigned to vortex-free aggregation, prevail up to 10^8 atoms per droplet. This finding builds the basis for exploring the assembly of far-from-equilibrium nanostructures at low temperatures.

The visualization of dopant nanostructures formed inside a helium droplet offers an opportunity to parse the factors driving their self-assembly in a cold, superfluid environment, including the effect of intermolecular forces among the dopant materials and their attraction to quantized vortices, if present. While helium nanodroplets have long been used as a matrix for cooling and preparing dopants for their spectroscopic studies [1,2], it is only through x-ray coherent diffractive imaging that the visualization of both the droplets and dopant nanostructures has become possible [3,4]. One testament to this imaging technology is the captured in situ configurations of xenon-traced vortex filaments in submicron-sized superfluid helium droplets [5–10].

Vortices with quantized circulation are a manifestation of superfluidity and play an essential role in the rotational dynamics of both Bose–Einstein condensates and superfluid helium [11–14]. Although these vortices are fascinating and intriguing by themselves, their presence

dominates nanostructure formation since many dopants are easily attracted to them [9,10,15–18]. On the other hand, in their absence, the dopants may be randomly distributed inside the confined droplet space because of the droplet’s superfluid state and the weak interaction between helium and the dopant [1,2,19,20]. Additionally, since any heat associated with the doping process and dopant aggregation are rapidly dissipated through the evaporation of helium atoms, the dopants could form amorphous particles with a fractal-like substructure influenced by intermolecular forces [19,21–24]. In small helium nanodroplets containing a few tens of thousands of atoms, examples of spectroscopically-identified nanostructures include: silver forming compact clusters at one or several sites in the droplet [20], hydrogen cyanide assembling into linear chains [25], water forming the smallest observed ice nanostructure with only six molecules [21], weakly-bound magnesium aggregating in a foam structure [26,27], and a core-shell structure of a multi-component doped droplet [28–30].

To explore and image self-organized structures, especially those that can only be formed in superfluid droplets, the presence of vortices needs to be suppressed. One means of controlling their presence is to generate

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the droplets by expanding cold helium gas. Here, nanodroplets may be stochastically formed from the initial condensation of helium into small clusters that further grow through collision with other helium clusters at some distance away from the nozzle [31]. This way of droplet generation may preclude the acquisition of angular momentum from shear forces prompted by the co-flowing helium fluid. In this paper, we demonstrate that droplets produced from gas condensation using a conical nozzle are larger and have smaller rotational distortion as compared to those produced using a pinhole nozzle in previous experiments [5,8,32,33]. Smaller shape deformations indicate a smaller angular momentum of the droplets, which likely contain few if any vortices. Because of this effect, we could identify two major types of xenon nanostructures, filaments and compact, that are respectively assigned to vortex-induced [6–10,34–36], and vortex-free aggregation. Finally, we map the occurrence of these two structures based on the droplet size and find that droplets smaller than ~ 200 nm in diameter are conducive to imaging far-from-equilibrium nanostructures; opening routes for studying self-assembly in a self-bound superfluid droplet.

X-ray coherent diffractive imaging was performed at the Nano-sized Quantum Systems end station of the European XFEL’s Small Quantum Systems scientific instrument [37,38]. This imaging technique takes snapshots of the droplet size and shape and, if doped, the structure of the dopant aggregates [3,4]. The helium droplets were produced using a conical nozzle with a throat diameter of $150\text{ }\mu\text{m}$, a half-opening angle of 3° , and a channel length of 9.3 mm . The nozzle was attached to a Parker valve and operated at a constant stagnation pressure of 20 bar, while the nozzle temperature T_0 was varied from 5 K to 14 K. Variable amounts of xenon were introduced in a gas doping cell installed along the droplet’s flight path. At ~ 1.05 meters from the nozzle exit, the droplets reached the interaction volume where they intersected the x-ray beam that had a photon energy of 1 keV, pulse energies in the range of 3 mJ to 5 mJ, and a focus diameter of $\sim 1.5\text{ }\mu\text{m}$ (FWHM) [38]. The scattered light from a diffraction event was collected up to 6° by an area detector located ~ 370 mm from the interaction volume [37]. A detailed description of the experiment and other relevant information are given in the Supplementary Material (SM) [39].

Figures 1(a) to 1(c) show examples of diffraction images for pure (i.e., undoped) droplets produced at $T_0 = 5.0\text{ K}$, 7.3 K and 10.0 K , respectively. The size distributions are depicted in Fig. 1(d) where the droplet size N_{He} is customarily given as the number of helium atoms and is related to the droplet’s radius R_D through $N_{\text{He}} = (R_D/0.222\text{ nm})^3$ [1,2]. The nozzle temperatures explored here encompass two droplet production regimes: the fragmentation of liquid helium close to the nozzle throat at $T_0 = 5\text{ K}$, and the condensation of helium gas after the

nozzle throat at $T_0 \geq 10\text{ K}$ [1,2]. The size distribution at $T_0 = 5\text{ K}$ is bimodal with a broad size range, whereas, that for $T_0 \geq 10\text{ K}$ is more narrow with an almost log-normal shape. Under similar stagnation conditions, the droplets produced at the gas condensation regime using the conical nozzle are about two to three orders of magnitude larger than those produced using the $5\text{ }\mu\text{m}$ pinhole nozzle [1,2], also see Fig. S7 in SM [39].

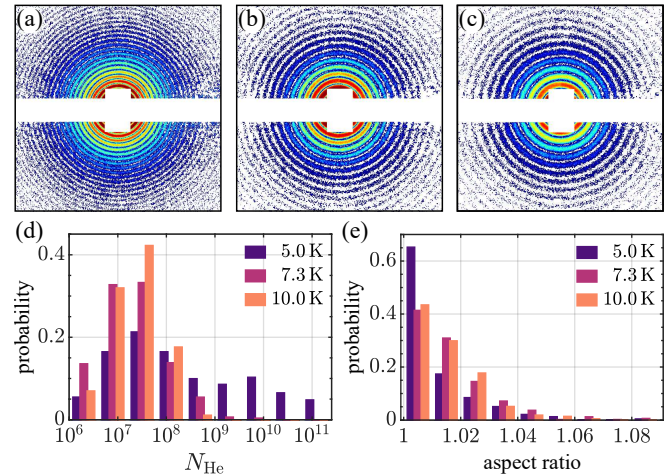


FIG. 1. Selected diffraction images of pure droplets obtained at a nozzle temperature of (a) 5.0 K, (b) 7.3 K, and (c) 10.0 K. The size distributions are given in panel (d), and their aspect ratio distributions in panel (e). The numbers of diffraction images used for the size (aspect ratio) distributions are 291 (236) for 5.0 K, 766 (488) for 7.3 K, and 272 (214) for 10.0 K.

As with rotating viscous liquid drops [40], a rotating superfluid helium droplet can have a series of equilibrium shapes depending on its angular velocity and angular momentum [41,42]. One way to describe a droplet’s shape is through its aspect ratio AR , which is defined as the ratio between the semi-major and semi-minor axes of the projected droplet density onto a two-dimensional plane. Figure 1(e) shows the aspect ratio distribution for $T_0 = 5.0\text{ K}$, 7.3 K and 10.0 K with average values of (1.011 ± 0.012) , (1.016 ± 0.013) , and (1.015 ± 0.012) , respectively. Collectively, 96.4% (99.9%) of the droplets generated using the conical nozzle have an aspect ratio smaller than 1.05 (1.10), making them close to spherical, and the largest observed AR is 1.3. These values are considerably smaller than those obtained from previous investigations using different kinds of nozzles. For instance, values as high as 2.3 have been reported for droplets produced from a $5\text{ }\mu\text{m}$ pinhole nozzle at 20 bar and 5 K with $\sim 65\%$ of the droplets having $AR \leq 1.05$ [43]. Another experiment using a flow-cryostat but with the same $5\text{ }\mu\text{m}$ pinhole nozzle at $T_0 \simeq 5\text{ K}$ reported a mean value $\langle AR \rangle = 1.059 \pm 0.005$, with 1% of the droplets exhibiting $AR \geq 1.4$ [33]. For droplets produced at 80 bar and 5.4 K using a trumpet-shaped nozzle with a throat diameter of $100\text{ }\mu\text{m}$ and a half-opening angle of 20° , aspect ratios as

large as ~ 3 were observed, and 92.9% of the droplets have close-to-spherical shape [44]. Further elaboration about the droplets produced using these three nozzles is included in the SM [39].

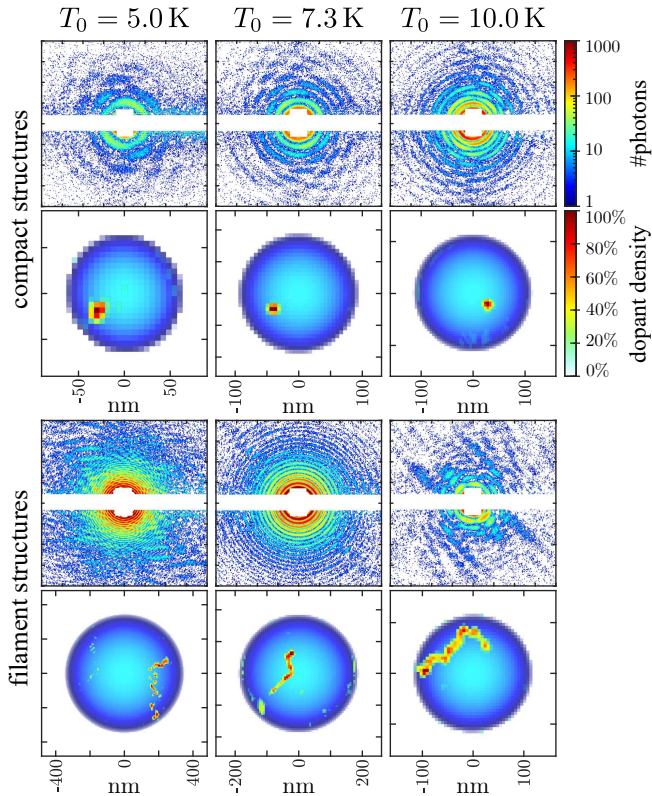


FIG. 2. Major types of xenon nanostructures inside helium droplets generated using a conical nozzle at different nozzle temperatures T_0 . Shown are selected examples of compact structures from vortex-free droplets and filaments from vortex-hosting ones. For each structure type, the first row shows the measured diffraction image, whereas the second shows the numerically reconstructed xenon-doped droplets using a modified DCDI, see the SM [39]. The nm-scale is different in each figure. The droplet is shown in blue colors, while the droplet density scales with the brightness. The color scales for the diffraction and the dopant reconstruction are shown on the upper right portion.

The consequence of generating close-to-spherical and large superfluid nanodroplets is reflected on the types of nanostructures observed inside them. Figure 2 shows diffraction images and numerically reconstructed xenon clusters in droplets generated at different nozzle temperatures. The image reconstruction algorithm is a modified version of the Droplet Coherent Diffractive Imaging (DCDI) [6,39]. In each reconstruction, the droplet is represented by dark blue to blue colors, while the xenon clusters are represented by dark red to light green colors. Two major kinds of structures were observed: compact (upper panel) and filaments (lower panel). The compact structures are assigned to vortex-free aggregation, are located at some distance away from the droplet's cen-

ter, and have roughly circular shapes with radii ranging from 10 nm to 15 nm. Because of the high thermal conductivity of superfluid helium, these compact structures possibly grow with an amorphous or even fractal morphology as similarly observed in bulk superfluid helium [23,24,45]. Compact structures larger than 20 nm are also observed (see Fig. S5 in SM [39]), but instead of having a more circular shape, they seem to be composed of two clusters, where each may have initially been formed at different aggregation sites before colliding and fusing into a cluster-cluster aggregate. However, because the heat of condensation is not enough to melt them, the two clusters remain distinct without reforming into one big clump [23].

Similar to previous observations [5–10], the filaments are also assigned to xenon clusters tracing the vortex length. In contrast, however, the number of vortices in our experiment is rather few (< 6), and no Bragg peaks, whose occurrence in a diffraction attests to the presence of a vortex lattice [5,8], were observed. The off-centered single vortices in Fig. 2 convey that the xenon distribution along the vortex line is uneven; rather, each is dotted by distinct nanometer-sized (10 to 20 nm) xenon clusters. Additionally, the filament is not smooth and its curvature is not as expected for an off-centered vortex [15,46]. Instead, it has oscillations along its length that specify underlying dynamics of how dopants approach a vortex [16,17,45,47] and indicate the presence of helical Kelvin waves relevant to quantum turbulence and the cascade of excitations in superfluids [11,48,49]. This observation of single vortices is instrumental to understanding vortex nucleation and decay in a superfluid helium droplet, as similarly observed in Bose-Einstein condensates [50,51].

Crucial to distinguishing these two nanostructures is the droplet's angular momentum, which may be partitioned between quantum vortices and capillary waves. Almost all droplets in earlier imaging experiments were generated in the liquid fragmentation regime and contained multiple vortices [5–10,32,33,44,52]. These droplets probably acquire angular momentum through the shear flow of the expanding helium fluid [32–34]. A smaller nozzle orifice, or higher velocities due to higher stagnation pressures of the expanding fluid, increases the action of shear forces that translates to a larger vorticity of the fluid flow. Within a few nanoseconds during flight, the droplets become superfluid via evaporative cooling and reach a temperature of ~ 0.4 K after ~ 100 ps [18] with the angular momentum conserved. If the acquired angular momentum is sufficient, one or more stable quantized vortices are nucleated. If more than enough is acquired, multiple vortices form a triangular lattice with the same quantum of circulation and sense of rotation [11,14] and bear most of the droplet's angular momentum [8,41,42]. At high angular momentum, the velocity fields from the vortices and the capillary waves induce deformations

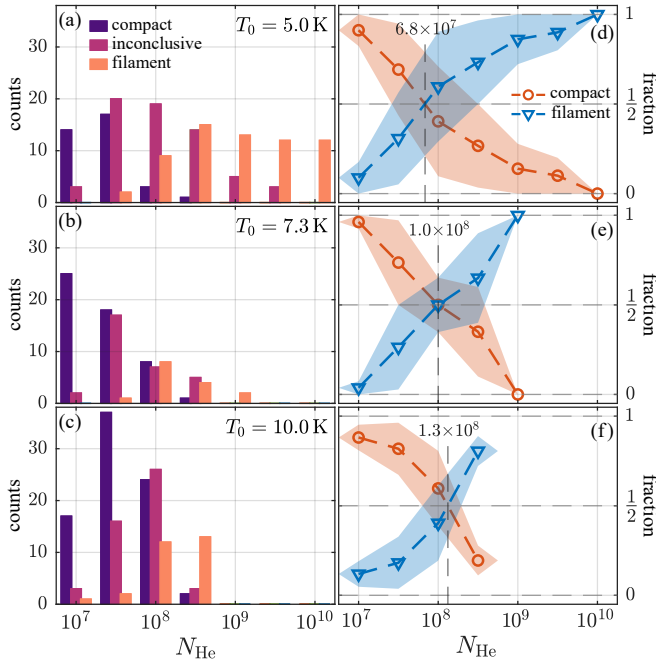


FIG. 3. Mapping structure formation as a function of the droplet size. The counts (left) and fractions (right) of the identified compact and filament xenon nanostructures for $T_0 = 5$ K, 7.3 K, and 10 K (from top to bottom). Structures that could not be assigned to either categories are categorized as inconclusive (see Fig. S5 in SM [39]) and serve as identification error for the fraction diagrams.

that resemble the equilibrium shapes of rotating classical viscous droplets [8,41,42], i.e., from spherical to oblate and to triaxial prolate shapes at higher angular momenta [5,8,32,33,44]. The angular momentum required for a single straight vortex to reside in the droplet's center is equal to $N_{\text{He}}\hbar$ [15,41,42,46,53]. If less angular momentum is present, an off-axis curved vortex may be created that precesses around the droplet's center and whose angular momentum decreases with its curvature [15,16,46]. On the other hand, if the acquired angular momentum is insufficient for vortex nucleation, it may be stored as capillary waves, and, in this case, the droplets can only rotate by adopting prolate-shape configurations [41,42,53]. In the present experiment, the close-to-spherical droplets suggest that they contain small angular momenta, consistent with the observation of the absence or a small number of vortices, see SM for further discussion [39]. However, the exact partition of the small angular momenta between few vortices and capillary waves requires further theoretical and experimental studies, especially for sub-micron to millimeter-sized superfluid drops [54].

The finding that compact structures are mainly observed in smaller droplets and filaments in larger ones can locate the droplet sizes where the formation of either structure is favored. Figure 3 shows the abundance and fractions of these xenon structures for $T_0 = 5$ K (top),

7.3 K (middle), and 10 K (bottom). The vertical dashed lines on the fraction panel indicate that the transition crossover, or where the onset of vortex-induced xenon aggregation is facilitated, occurs at $N_{\text{He}} \approx 10^8$, almost independent of how the droplets are generated. Nevertheless, the means of droplet generation affects the relative abundance of either structures. For instance, filaments are dominant in the liquid fragmentation regime at $T_0 = 5$ K, and conversely, compact structures in the gas condensation regime at $T_0 \geq 10$ K. The observation of filaments in the gas condensation regime is interesting since it suggests that angular momentum was acquired during droplet growth, possibly from droplet collision. A simple kinematic model to estimate the acquired angular momentum shows that for the same collision velocity the angular momentum per atom increases with droplet size, and plausibly why smaller droplets are vortex free [39]. In other words, larger droplets require slower collision velocities to have the same angular momentum per atom as compared to smaller droplets. Moreover, these larger droplets also need less angular momentum per atom to maintain a stable curved vortex [15,39]. These estimates corroborate the observed size-dependent vortex stability in superfluid helium droplets. Finally, at 12 K and 20 bar, only compact structures were found, but both structures with a similar crossover were observed again when the stagnation pressure was increased to 60 bar, see Fig. S8 in SM [39].

In summary, we have demonstrated that large superfluid helium nanodroplets generated using a large-diameter conical nozzle are close to spherical where the angular momentum is likely shared between a few vortices and surface capillary waves. Additionally, we distinguished two main types of dopant aggregation, where their relative abundance is droplet size dependent. The size-dependent transition from vortex-free to vortex-induced structures also emphasizes the size dependence of the rotational state of superfluid helium droplets. For instance, xenon-traced vortex filaments are only readily observed in droplets with more than 10^8 helium atoms. This size not only benchmarks the onset where vortex-induced nanostructure formation starts to be dominant, but it also indicates the droplet size amenable for the preparation of self-organized nanostructures without the influence of quantum vortices. This study could be a beginning for x-ray imaging of far-from-equilibrium nanostructures from various kinds of dopant materials with different intermolecular forces (e.g. hydrogen or metallic bonding). Without the influence of vortices, it also becomes possible to image dynamics (e.g. nanosplasma ignition) occurring between the dopants and the droplet (in contrast to the observation in Ref. [55]), and how charges are distributed on the droplet's surface [10].

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