

In our experiment, the weak coherent input state is p 1285 nm to the single-photon level. To generate photon silicon-nitride ring microresonator (200-GHz free spec from the SFWM pump at the two weak coherent-state freq mode 1, we strongly filter it through a chain of fiber pair intensity to avoid contaminating the input state. with three strong pulsed C-band BS-FWM pumps before en

in the fiber, the transformed O-band fields are separated by a series of beam splitters and detected by superconducting nanowire single-photon detectors (SNSPDs) at each of the three levels. Detection events are then recorded. The nonlinear phase (hence the interaction strength) is scanned by varying the FWM pumps (hence the interaction strength) is scanned. A subset of the resulting measured normalized singles rates for the weak coherent states and is plotted in Fig. 1b). In the absence of multiphoton emission, the measured delay histogram over the inverse microresonator frequency for the coherent state in Fig. 1b, it is evident that the input state is fundamentally nonclassical in nature. To estimate the multiphoton parameter via a separate heralded experiment, the losses are based on measurements concurrent with the multiphoton measurements in the absence of multiphoton emission. This is the first time, to the best of our knowledge, that a multiphoton parameter has been observed in a three-mode microresonator with > 2 modes has been observed in a three-mode microresonator with > 2 modes, which would allow arbitrary control over each mode.

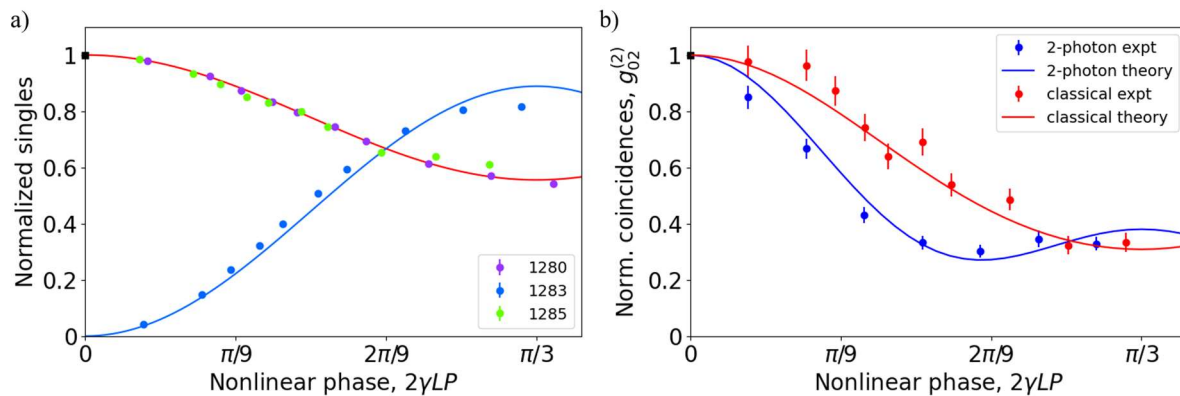


Figure 1. a) Singles rates for the two-mode weak coherent-state input, computed assuming Poisson statistics. b) Comparison between experimental and theoretical results. In both figures, the point at zero nonlinear phase represents the

Our experiment demonstrates the possibility of achieving a large number of modes in the quantum regime. Through use of a microresonator, our system has the potential to enhance the scalability of quantum information processing by relying on beamsplitter transformations, including boson sampling, which show promise for the scalable quantum information platform.

References

- [1] J. Preskill, "Quantum Computing in the NISQ era and beyond," *Quantum*, vol. 2, p. 79, 2018.
- [2] E. T. Kjaer, "A Scheme for Efficient Quantum Computation with Linear Optics," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [3] D. E. Browne, "Resource-Efficient Linear Optics Computation," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [4] N. M. M. Ueda, "Universal Quantum Computation with Continuous-Variable Modes," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [5] S. C. Lee, "Ramsey Interference with Two-Mode Fock States," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [6] C. J. H. van der Wal, "Translation of Quantum States by Frequency Mixing in Fock States," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [7] C. J. H. van der Wal, "Frequency-Domain Quantum Interference with Correlated Photons," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [8] B. E. A. Saleh, "Spectral Photonic Lattices with Long-Range Coupling," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [9] H. E. T. Kjaer, "Electro-Optic Frequency Beam Splitters and Tritters for High-Speed Quantum Computing," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [10] H. E. T. Kjaer, "High-Dimensional Discrete Fourier Transform in Quantum Optics," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [11] H. E. T. Kjaer, "Bayesian Tomography of High-Dimensional On-Chip Biphoton States," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [12] C. J. H. van der Wal, "Frequency-Domain Quantum Interference with Correlated Photons," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [13] R. O. T. Kjaer, "Three-Way Frequency Beam Splitters and Tritters for High-Speed Quantum Computing," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.
- [14] S. E. T. Kjaer, "A large-scale continuous-variable quantum information platform," *Phys. Rev. Lett.*, vol. 111, p. 110501, 2013.