

BAL outflow in quasar B0254-3327B: analysis and comparison with other extreme UV outflows

Doyee Byun¹,¹★ Nahum Arav,¹ Maryam Dehghanian,¹ Gwen Walker¹ and Gerard A. Kriss²

¹Department of Physics, Virginia Tech, Blacksburg, VA 24061, USA

²Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA

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ABSTRACT

We have identified a broad absorption line (BAL) outflow in the *HST*/STIS spectrum of the quasar QSO B0254-3327B at velocity $v = -3200 \text{ km s}^{-1}$. The outflow has absorption troughs from ions such as Ne VIII, Na IX, Si XII, and Ne V. We also report the first detection of S XIV absorption troughs, implying very high ionization. Via measurement of the ionic column densities, photoionization analysis, and determination of the electron number density of the outflow, we found the kinetic luminosity of the outflow system to be up to ~ 1 per cent of the quasar’s Eddington luminosity, or ~ 5 per cent of the bolometric luminosity, making it a potential contributor to AGN feedback. A solution with two ionization phases was needed, as a single phase was not sufficient to satisfy the constraints from the measured ionic column densities. We find that the ionization parameter of the very high-ionization phase of the outflow is within the expected range of an X-ray warm absorber. We also examined the physical properties of the outflow of Q0254-334 along with previously studied extreme UV outflows, with a total sample of 24 outflow systems, finding a weak negative correlation between outflow velocity and distance from the central source, with larger distances corresponding to slower velocities. The very high-ionization phase of the Q0254-334 outflow has one of the highest ionization parameters of UV absorption outflows to date, which we attribute to the presence of S XIV.

Key words: galaxies: active – quasars: absorption lines – quasars: individual: QSO B0254-3327B.

1 INTRODUCTION

Quasar absorption outflows are often invoked as likely contributors to active galactic nucleus (AGN) feedback (e.g. Silk & Rees 1998; Scannapieco & Oh 2004; Yuan et al. 2018; Vayner et al. 2021; He et al. 2022). They are detected via blueshifted absorption troughs in the rest frame of $\lesssim 40$ per cent of quasars (e.g. Hewett & Foltz 2003; Trump et al. 2006; Dai, Shankar & Sivakoff 2008; Knigge et al. 2008; Allen et al. 2011). In order to contribute to AGN feedback, outflow systems theoretically require a kinetic luminosity (\dot{E}_k) of at least ~ 0.5 per cent (Hopkins & Elvis 2010) or ~ 5 per cent (Scannapieco & Oh 2004) of the quasar’s luminosity, which we interpret to be the Eddington luminosity (L_{Edd}) following the reasoning described by Miller et al. (2020d), as opposed to the bolometric luminosity (L_{Bol}). Past studies have found outflows that fit one or both of these criteria (e.g. Moe et al. 2009; Borguet et al. 2013; Chamberlain, Arav & Benn 2015; Leighly et al. 2018; Choi et al. 2020; Miller et al. 2020a; Walker, Arav & Byun 2022; Byun, Arav & Hall 2022b; Choi et al. 2022b; Byun, Arav & Hall 2022c).

In order to find the value of \dot{E}_k , it is important to find the mass flow rate (\dot{M}), a method for which involves finding the electron number density (n_e) and ionization parameter (U_H) to measure the distance (R) of the outflow from the central source (Borguet et al. 2012b). Multiple quasar outflows have been analysed via this method

(e.g. Hamann et al. 2001; de Kool et al. 2001; Byun, Arav & Walker 2022a; Walker, Arav & Byun 2022). For ionized outflows, the ionization parameter can be determined by measuring the column densities of ions, and comparing them with simulated values based on a range of U_H and hydrogen column density (N_H). Multiple outflow analysis studies have been conducted using the spectral synthesis code CLOUDY (Ferland et al. 2017) for this method (e.g. Xu et al. 2018; Miller et al. 2020a; Walker, Arav & Byun 2022; Byun, Arav & Walker 2022a).

This paper presents the analysis of the absorption outflow of the quasar QSO B0254-3327B (hereafter Q0254-334), using the method described above, based on *HST*/STIS observational data, ultimately finding the ratio between \dot{E}_k and L_{Edd} .

The paper is structured as follows: Section 2 describes the observation and data acquisition of Q0254-334; Section 3 discusses the process of finding the ionic column densities, N_H , U_H , and n_e of the outflow; Section 4 shows our analysis results of whether the outflow’s kinetic luminosity is sufficient to contribute to AGN feedback, and compares our results with studies of other outflows; and Section 5 concludes and summarizes the paper. In our analysis, we adopted a cosmology of $h = 0.696$, $\Omega_m = 0.286$, and $\Omega_\Lambda = 0.714$ (Bennett et al. 2014). We used the PYTHON astronomy package ASTROPY (Astropy Collaboration et al. 2013, 2018) for our cosmological calculations. We also used SCIPY (Virtanen et al. 2020), NUMPY (Harris et al. 2020), and PANDAS (v1.2.4, McKinney 2010; Reback et al. 2021) for the majority of our numerical computations, as well as MATPLOTLIB (Hunter 2007) for plotting our figures.

★ E-mail: dbyun@vt.edu

2 OBSERVATION AND DATA ACQUISITION

Q0254-334 (J2000; RA = 02:56:47.84, DEC = −33:15:26.16, $z = 1.863$) was observed with *HST*/STIS on 2001 February 17 as part of the program SNAP 8681, and on 2001 March, 4 as part of the program GO 8569, with the G230L and G140L gratings, respectively. Prior to this, it was also observed with *HST*/FOS in 1994. Due to the limited wavelength range of the FOS data relative to that of STIS, we have focused on the STIS data for the purpose of this analysis. After retrieving the data from the Mikulski Archive for Space Telescopes, we have co-added the two STIS spectra, and corrected the combined spectrum for galactic reddening and extinction with $E(B - V) = 0.0205$ (Schlafly & Finkbeiner 2011), and the extinction model by Fitzpatrick (1999). The co-added and dereddened spectrum of the two observations, covering observed wavelengths 1138.6–3156.6 Å, along with the FOS spectrum, is shown in Fig. 1.

We have identified a broad absorption line (BAL) outflow system at $v = -3200 \text{ km s}^{-1}$, with its ionic absorption troughs marked by red vertical lines in Fig. 1. Troughs exist of species such as Ne VIII, Na IX, and Si XII, as well as excited state transitions such as O IV* and Ne V*. Arav et al. (2020) define a BAL in the extreme UV range as a continuous absorption feature with normalized flux $I \leq 0.9$ over a width of $\Delta v \gtrsim 1500 \text{ km s}^{-1}$, at least -3000 km s^{-1} blueward of the centre of emission. We have verified that the outflow is a BAL outflow by confirming the width of the Si XII $\lambda 499.406$ and Ne VIII $\lambda 780.324$ troughs (see Fig. 2). The normalized flux was found by modelling the quasar’s continuum via a spline model that gave the minimum possible continuum above the absorption, and the most prominent emission features (e.g. O VI) with Gaussians. The presence of the Ne V* feature allowed us to find the value of n_e , as shown in Section 3.2.

3 ANALYSIS

3.1 Ionic column densities

As the ionic column densities (N_{ion}) of the outflow are crucial in finding the physical properties of the outflow, we used two different methods to find them based on the absorption troughs: the apparent optical depth (AOD) method in which we assume uniform and homogeneous covering (Savage & Sembach 1991); and the partial covering (PC) method in which we include a covering factor $C < 1$ (Barlow, Hamann & Sargent 1997; Arav et al. 1999a, 1999b).

The AOD method allows us to find upper limits and lower limits of ionic column densities with its relative simplicity, while the PC method lets us find more accurate measurements of ions with doublet features (e.g. de Kool, Korista & Arav 2002; Arav et al. 2005; Borguet et al. 2012a; Byun, Arav & Hall 2022b). As done by Byun, Arav & Hall (2022c) for the quasar J024221.87+004912.6, we selected the appropriate method for computing the column density of each ion.

The AOD method involves the following relation between intensity and optical depth (Spitzer 1978; Savage & Sembach 1991):

$$I(\lambda) = I_0(\lambda)e^{-\tau(\lambda)}, \quad (1)$$

where $I(\lambda)$ is the intensity as a function of wavelength, $I_0(\lambda)$ is the intensity without absorption, and τ is the optical depth. Finding the optical depth enables computation of the column density, as they have the following relation:

$$\tau(v) = \frac{\pi e^2}{m_e c} f \lambda N(v), \quad (2)$$

where $\tau(v)$ is the optical depth as a function of velocity, e is the elementary charge, m_e is the mass of an electron, and $N(v)$ is the column density per unit velocity. Integrating $N(v)$ over the velocity range of an ion’s absorption trough results in the ion’s column density.

As mentioned above, the PC method involves a covering factor $C < 1$, which follows the relation shown in the equations below (Arav et al. 2005):

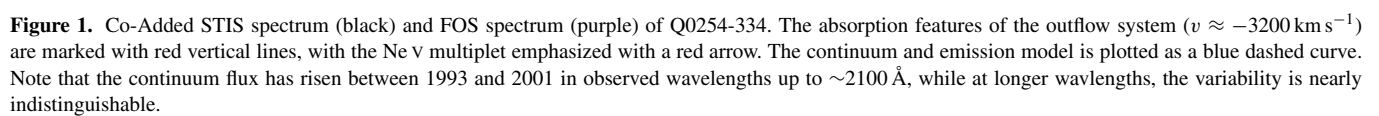
$$I_R(v) - [1 - C(v)] = C(v)e^{-\tau(v)} \quad (3)$$

$$I_B(v) - [1 - C(v)] = C(v)e^{-2\tau(v)}, \quad (4)$$

where $I_R(v)$ and $I_B(v)$ are the intensities at the red (longer wavelength) and blue (shorter wavelength) troughs of a doublet transition, $C(v)$ is the covering factor as a function of velocity, and τ is the optical depth.

For each ion, we converted the spectrum from wavelength space to velocity space, using the redshift of the quasar and the wavelengths of the ionic transition lines (see Figs 3,4). We then chose integration ranges for each ion that covered visible absorption features and minimized blending effects with other lines. For instance, the O VI doublet had heavy blending between the red and blue troughs (see Fig. 3f). We thus chose a range where the overlap between the red and blue troughs would be minimized and computed a lower limit to the column density with the AOD method. As there were no discernible absorption troughs of Ly γ , C III, and S IV*, we measured their AOD column density with integration range $v \approx -4500$ to -2000 km s^{-1} to match the Ne VIII width, and treated them as upper limits. Due to the severe blending in the multiplet of S IV $\lambda\lambda 744.904$, 748.393, and S IV* $\lambda 750.221$ (see panel 4 of Fig. 1), we were unable to pinpoint the column density of the resonance state S IV 0 from this trough. However, as there was no discernible absorption trough of S IV $\lambda 809.656$, we were able to find an upper limit of its column density. Similarly, the trough of O IV $\lambda 787.711$ blended with O IV* $\lambda 790.190$ and potentially with the neighbouring S V $\lambda 786.468$. As such, we were unable to find the column density of the resonance state O IV 0, and could only find a lower limit of the O IV* column density. We were also limited to finding a lower limit of the Ne VIII column density based on an AOD measurement due to the saturation of the doublet troughs. In the doublet of S XIV, we determined the red trough of S XIV $\lambda 446$ to be contaminated due to the visible blueward absorption compared to the blue trough $\lambda 418$ (see Fig. 4 plot c). Due to this limitation, we measured the AOD column density of S XIV based on the blue trough. We determined that it was safe to treat this column density as a measurement due to its shallower depth relative to similarly ionized troughs with comparable oscillator strengths (e.g. Si XII).

The integrated column densities are shown in Table 1. The rightmost column shows the values adopted for the photoionization solution described in Section 3.3. The errors have been propagated from the error in the flux and a conservative 20 per cent error has been added in quadrature to the adopted column density errors to account for the uncertainty in the continuum level due to the subjectivity of the model (Xu et al. 2018). This uncertainty is demonstrated in the column density calculation of O IV* based on the different continuum models shown in Fig. 5. The maximum, minimum, and intermediate continuum fits in the region are shown as blue, red, and purple dashed lines, respectively. The O IV* absorption is marked with a grey vertical line. The AOD measurement of the O IV* column density is $28.0^{+2.8}_{-2.0} \times 10^{14} \text{ cm}^{-2}$ for the intermediate continuum, while the higher and lower continuum levels yield results of $32.6^{+2.8}_{-2.0} \times 10^{14}$ and $24.4^{+2.8}_{-2.0} \times 10^{14} \text{ cm}^{-2}$. This indicates a ± 15 per cent difference in column density depending on continuum level, and a 20 per cent



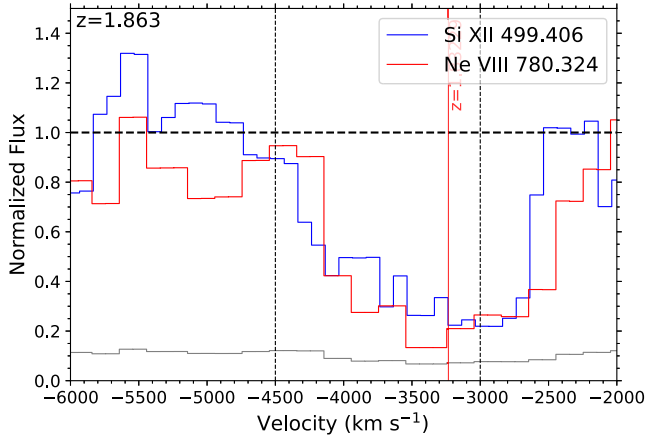


Figure 2. Normalized flux in the region of the Si XII $\lambda 499.406$ and Ne VIII $\lambda 780.324$ absorption. The dashed horizontal line shows the continuum level, and the dotted vertical lines show the region between $v = -4500$ and -3000 km s^{-1} , which represents the range for the EUV BAL criteria, as defined by Arav et al. (2020). The normalized flux for Si XII falls below 0.9 in this region. While the Ne VIII normalized flux lies above 0.9 at -4500 km s^{-1} , it is within uncertainty range to fall below the threshold. To demonstrate, the error in the Ne VIII flux is shown in grey.

difference when including the individual errors. Note that the column density of Ne V was based on a Gaussian fit of the troughs of its different energy states, which we further describe in Section 3.2.

3.2 Ne v Gaussian fitting

As seen in Fig. 6 (top panel), the Ne v multiplet of the outflow is blended into a singular trough. The involved transitions are Ne v 0 ($\lambda = 480.415 \text{ \AA}$), Ne v* 411 ($\lambda = 481.227, 481.366, 481, 371 \text{ \AA}$), and Ne v* 1109 ($\lambda = 482.990, 482.994 \text{ \AA}$). To remedy the blending, we modelled the individual energy states of Ne v by fitting Gaussian profiles for each of the expected absorption features, and running a best-fit algorithm to best match the data. The free parameters used were the optical depth of the ground state Ne v trough, the width of the trough, and $\log n_e$. We assumed the AOD scenario and adjusted the depths of the excited state troughs to match the oscillator strengths of the transition lines, as well as the abundance ratios $N(\text{Ne v}^*)/N(\text{Ne v } 0)$ from the CHIANTI 9.0.1 atomic data base (Dere et al. 1997, 2019). We assumed a temperature of 10 000 K in our CHIANTI computations. A similar process of finding n_e via the ratios between the different energy states of Ne v is demonstrated by Miller et al. (2020a) and is especially illustrated in their Fig. 3.

We have found that the optical depth $\tau = 0.69 \pm 0.11$, FWHM = $2360 \pm 170 \text{ km s}^{-1}$, and $\log n_e = 3.6 \pm 0.1 [\text{cm}^{-3}]$. Using the modelled troughs, we have calculated the column densities of each energy state of Ne v, as shown in Table 1. Since the value of n_e is crucial in finding the distance of the outflow from the central source (as described in Section 4.1), we later ran a simulation with the spectral synthesis code CLOUDY (version c17.00, Ferland et al. 2017) in order to verify the temperature of the outflow. With the two-phase high-ionization solution later described in Section 3.3 as our input parameters, the simulation yielded a temperature of $T \approx 27\,000 \text{ K}$. Calculating the electron number density with this temperature yielded $\log n_e = 4.0^{+0.1}_{-0.1}$. As such, we adopted this value of $\log n_e$ for the purpose of our analysis. The total column density of Ne v based on this computation is in agreement with the value based on the $T = 10\,000 \text{ K}$ assumption.

As an alternate method of modelling the blended trough of the Ne v multiplet, we used the trough of Si XII $\lambda 499$ as a template to create a profile of two blended Gaussians (see Fig. 7). We then ran a best-fit algorithm to model the absorption of each energy state, leaving the width of the profile as a fixed parameter (see Fig. 6 bottom panel). This resulted in an electron number density of $\log n_e = 4.3 \pm 0.1 [\text{cm}^{-3}]$, which is only ~ 0.3 dex higher than the simple Gaussian fitting.

While the difference in the electron number density shifts the kinetic luminosity to lower values relative to those shown in Table 3, the kinetic luminosities remain in agreement within error. We thus focus on the results based on the Gaussian model throughout the paper. The parameters are described in further detail in Sections 4 and 4.1.

3.3 Photoionization solution

We used the measured ionic column densities to constrain the values of the hydrogen column density (N_H) and ionization parameter (U_H), as done in previous works (e.g. Xu et al. 2019; Walker, Arav & Byun 2022; Byun, Arav & Walker 2022a; Byun, Arav & Hall 2022b, c). For this purpose, we used a grid of simulated models produced with CLOUDY (Ferland et al. 2017) with a range of N_H and U_H values as input parameters, modelling the ionic abundances at different N_H and U_H . We used the ionic column densities as shown in Table 1 to set upper and lower limits to these parameters assuming solar metallicity. We adopted a spectral energy distribution (SED) that would match the V-band flux of Q0254-334 found on NED, the UV continuum flux measured at three separate points, as well as the X-Ray fluxes observed with Chandra at energy ranges from 0.5–7 keV (see Fig. 8). The resulting photoionization solution can be seen in Fig. 9. Note that there are limitations to this SED due to potential variability between the different observations that were referenced for its construction. Rao, Turnshek & Nestor (2006) reports a V-band magnitude of 16 and cite Wright et al. (1982), who in turn discuss observations made with the 3.9-m Anglo-Australian telescope on 1978 November 28 and on 1978 December 5. Chandra observations of Q0254-334 were made on 2000 January 2 and 2000 February 15. We have also calculated the α_{ox} spectral index based on our SED, using the following equation (Tananbaum et al. 1979; Sobolewska, Gierliński & Siemiginowska 2009):

$$\alpha_{\text{ox}} = 0.3838 \log \frac{L(2 \text{ keV})}{L(2500 \text{ \AA})}, \quad (5)$$

which yielded a result of $\alpha_{\text{ox}} = -1.58$. This is somewhat higher than the range of α_{ox} values of LBQS broad absorption line quasars, which were reported by Gallagher et al. (2006) (-2.58 to -1.65).

A single-phase solution was insufficient to satisfy the constraints from the ionic column densities. To remedy this issue, we formulated a two-phase solution in which a high- and very high-ionization phase exist co-spatially. We deduced that the two phases would be co-spatial based on the kinematic similarity between the high-ionization troughs and the very high-ionization troughs. Specifically, Fig. 6 shows that we get a very good fit for Ne v (a high-ionization line), using the velocity template of Si XII (a very high-ionization line, see Fig. 7). We find that the two-phase solution satisfies more of the constraints set by the measured ionic column densities (reduced $\chi^2 = 5.1$, as opposed to 22.3 for the one-phase solution). To cover the range of possible metallicities, we have also applied models of metallicity $Z \approx 4.68 Z_{\odot}$ (Ballero et al. 2008; Miller et al. 2020b), which are shown in the lower panel of Fig. 9. The results are favourable towards the supersolar metallicity solution, of which the reduced χ^2 values are

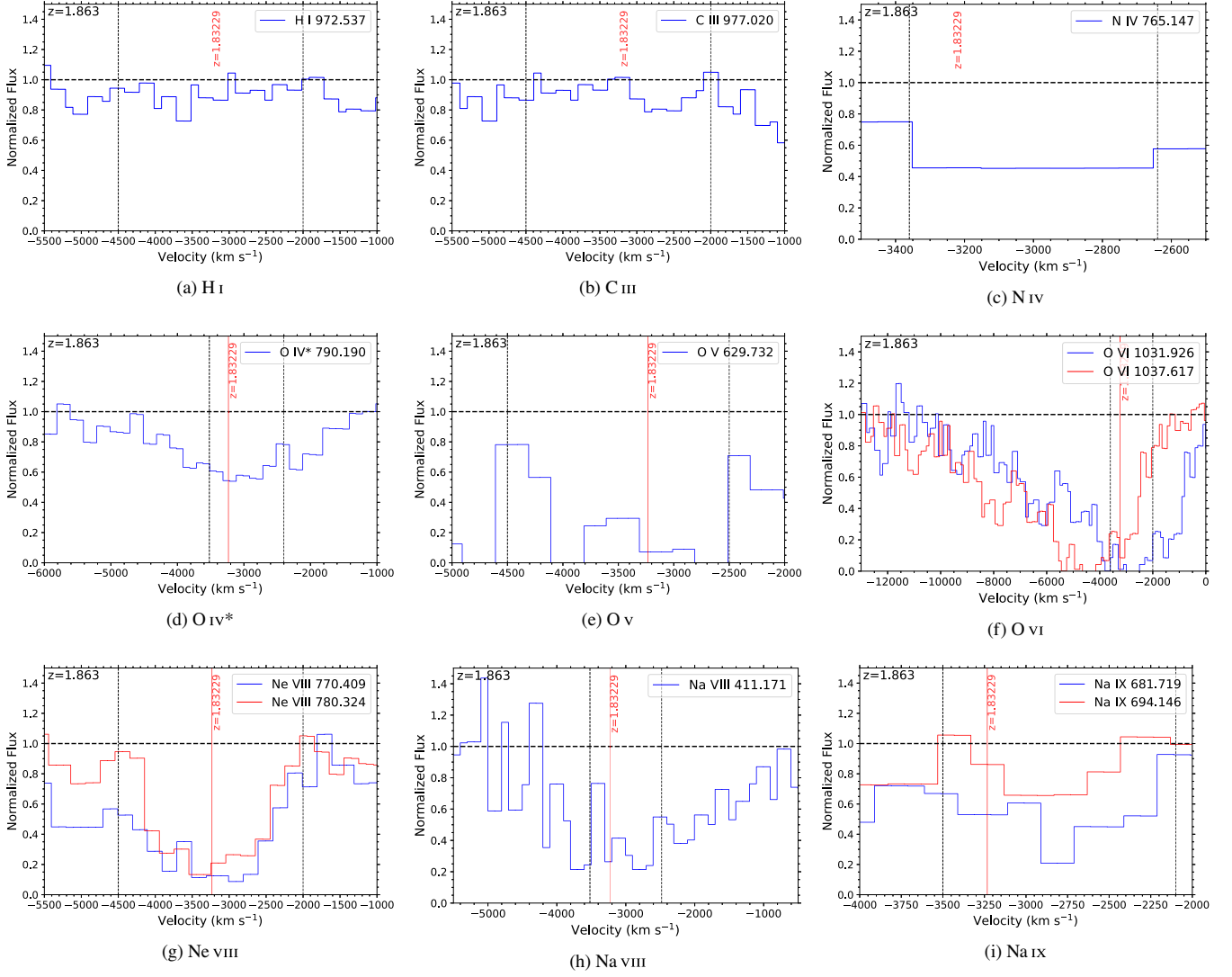


Figure 3. Absorption troughs of the Q0254-334 outflow plotted in velocity space. The velocity of the outflow at $z = 1.89229$ is marked with red vertical lines. The integration range used to calculate the column densities is marked with dotted vertical lines, while the continuum level is indicated by the dashed horizontal line.

16.0 and 0.5, for the one-phase and two-phase solutions, respectively. As discussed by Arav et al. (2013), the inability for a one-phase ionization solution to reasonably fit the measurements and limits of N_{ion} necessitates the adoption of a two-phase solution. This is further demonstrated by comparing the modelled column densities of H I, Na IX, and S XIV from supersolar one-phase and two-phase solutions to the observed ones, as shown in Table 2. Note that the reported measured column density of H I is an upper limit based on Ly γ . Due to the larger discrepancy between modelled and measured column densities in the solar abundance solutions (e.g. model H I column density upwards of ~ 20 times larger than measured), they have been excluded from the table. As can be seen in the table, the two-phase solution yields modelled column densities with a maximum 2σ difference between modelled and measured column densities, while the one-phase solution yields a 4–8 σ difference. As such, the comparison of modelled column densities favours the two-phase solution. The U_{H} and N_{H} values found using χ^2 analysis are shown in Table 3.

We compared the N_{H} and U_{H} values found using the Q0254-334 SED with those found using the SED of the quasar HE0238-1904 (hereafter HE0238, Arav et al. 2013), as the latter SED has been

adopted for quasar outflow analysis in several past papers (e.g. Miller et al. 2020a; Walker, Arav & Byun 2022; Byun, Arav & Walker 2022a; Byun, Arav & Hall 2022b, c). We report the $\log N_{\text{H}}$ and $\log U_{\text{H}}$ values derived from the HE0238 SED in Table 4. Comparing these values with those found in Table 3 shows that while the one-phase solutions are in agreement within error, the two-phase solutions show a discrepancy in the $\log U_{\text{H}}$ values that range up to ~ 0.5 dex. We conducted the same procedure using the Si XII template fit discussed in Section 3.2, of which the results are shown in Table 5.

3.4 Black hole mass calculation

Black hole masses of AGN are often found using the emission features of Mg II (Bahk, Woo & Park 2019) or C IV (Vestergaard & Peterson 2006; Coatman et al. 2017). However, as the STIS spectrum of Q0254-334 lacked both features, we looked to the O VI emission to compute the mass of the central black hole. We referred to the method described by Tilton & Shull (2013), measuring the O VI FWHM to find the mass.

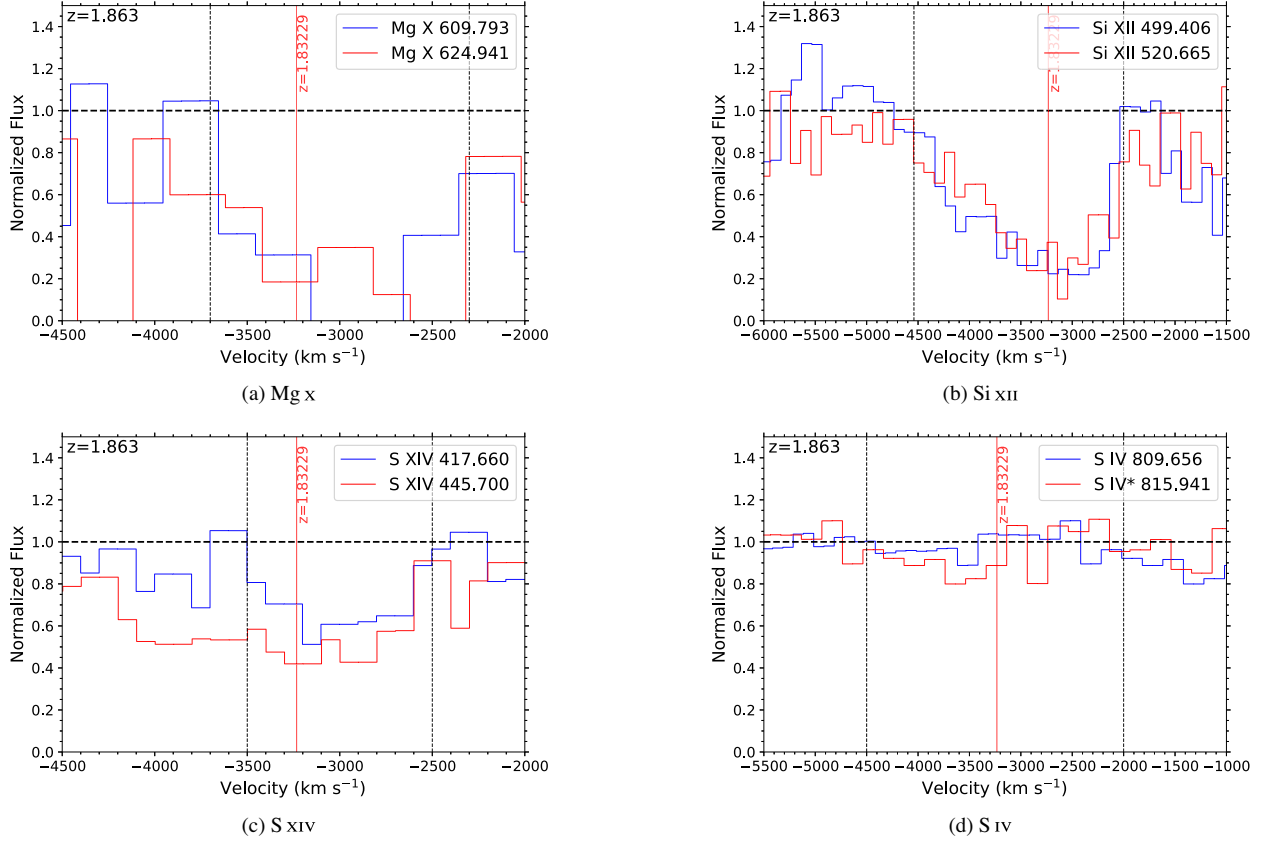


Figure 4. Absorption troughs of Mg X, Si XII, S XIV, and S IV* in the outflow of Q0254-334. Format and notation are identical to those of Fig. 3.

Table 1. Q0254-334 outflow column densities from STIS observations. The numbers next to the Ne V* excited states denote the energies in cm^{-1} . The values are in units of 10^{14} cm^{-2} .

Ion	AOD	PC	Adopted
H I	38^{+7}_{-4}	—	$<38^{+11}_{-4}$
C III	$1.8^{+0.3}_{-0.2}$	—	$<1.8^{+0.4}_{-0.2}$
N IV	$4.6^{+0.6}_{-0.4}$	—	$>4.6_{-1.0}$
O IV*	$28.0^{+2.8}_{-2.0}$	—	$>28.0_{-6.0}$
O V	50^{+8}_{-5}	—	$>50_{-11}$
O VI	114^{+10}_{-4}	—	$>114_{-23}$
Ne V total	196^{+22}_{-22}	—	196^{+45}_{-45}
Ne V 0	106^{+18}_{-18}	—	—
Ne V* 411	72^{+12}_{-12}	—	—
Ne V* 1109	18^{+3}_{-3}	—	—
Ne VIII	229^{+25}_{-10}	—	$>229_{-47}$
Na VIII	$29^{+11.6}_{-2.9}$	—	$>29_{-6.6}$
Na IX	$38.2^{+8.4}_{-7.2}$	$48.2^{+6.2}_{-6.1}$	$48.2^{+11.5}_{-11.3}$
Mg X	366^{+87}_{-46}	—	$>336_{-87}$
Si XII	360^{+57}_{-22}	—	$>360_{-75}$
S IV total	27^{+2}_{-2}	—	$<27^{+5}_{-2}$
S IV 0	$7.2^{+0.9}_{-0.6}$	—	—
S IV*	20^{+3}_{-2}	—	—
S XIV	198^{+20}_{-13}	—	198^{+45}_{-42}

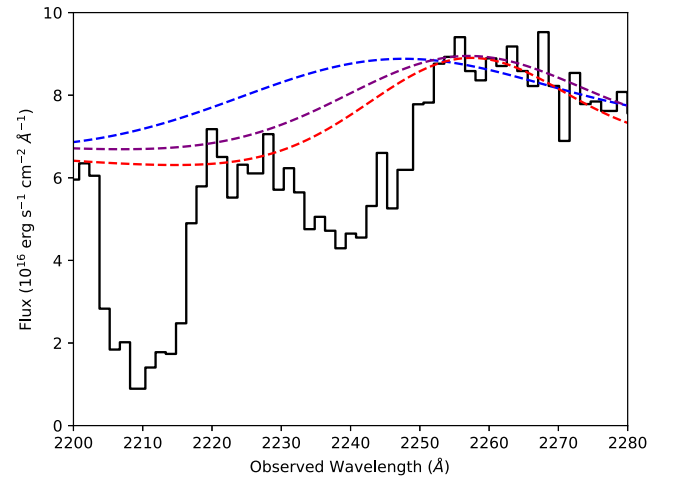


Figure 5. The O IV region of the Q0254-334 spectrum showing the possible continuum fits, which lead us to use 20 per cent error bars on the column densities (See text).

Although Tilton & Shull (2013) specify the use of two Gaussians to fit each line of the emission doublet, we opted to fit one Gaussian per line instead, as the lower signal to noise ratio of the STIS spectrum did not warrant the more detailed modelling method. We employed a best-fit algorithm adjusting the amplitude of the blue emission line, the ratio between the blue and red lines, and the FWHM of the blue line. The ratio between the blue and red line amplitudes was constrained between 1–1.5, and the widths of the two features

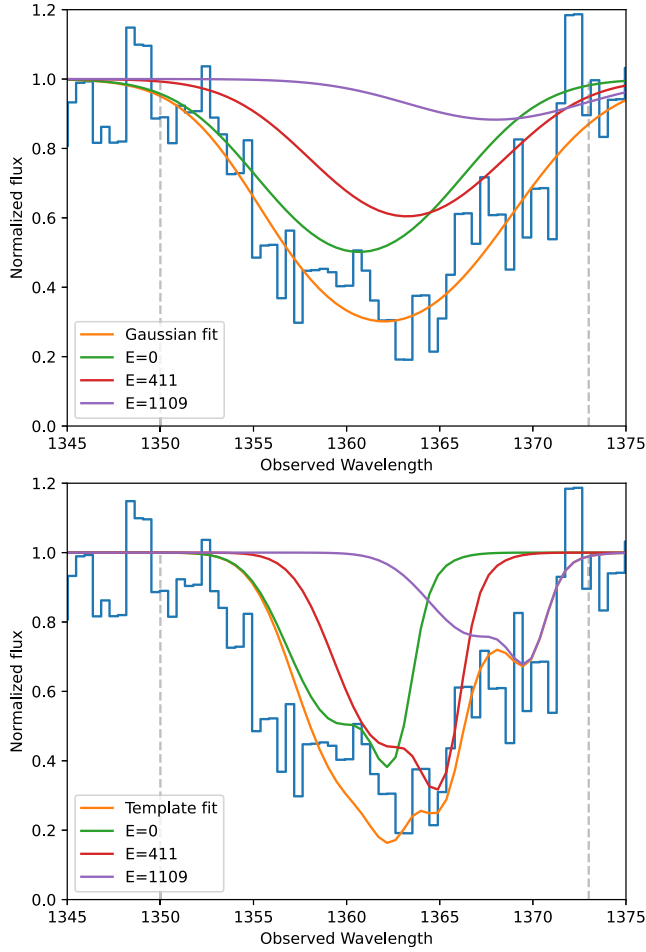


Figure 6. Modelling of the Ne V absorption troughs, created by fitting Gaussians (top) and by using Si XII as a template (bottom). The vertical dashed lines represent the range of data used for our fitting. The green curve is the modelled absorption of the resonance state of Ne V. The red curve shows the combined absorption of the $E = 411 \text{ cm}^{-1}$ level lines, and the purple curve shows the absorption of the $E = 1109 \text{ cm}^{-1}$ level lines. The absorption features from multiple lines of the same excited states have been combined within the figure. The orange curve represents the total combined modelled absorption of the Ne V multiplet.

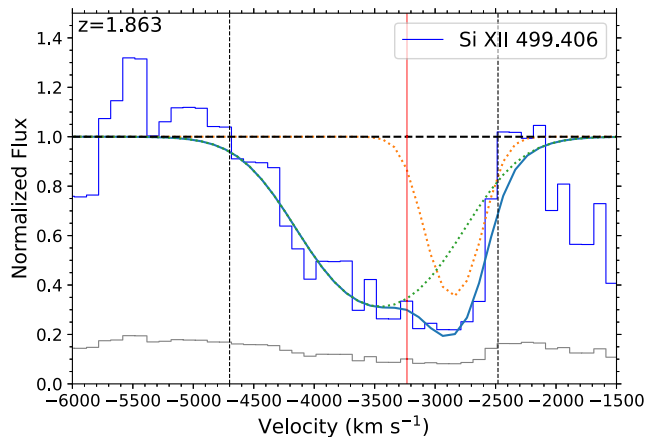


Figure 7. Fitting of a two-Gaussian profile to the absorption trough of Si XII. The dotted curves show the individual Gaussians in the profile, while the blue curve shows the blended profile of both Gaussians. The dotted vertical lines represent the range of data that was used for fitting the Gaussians.

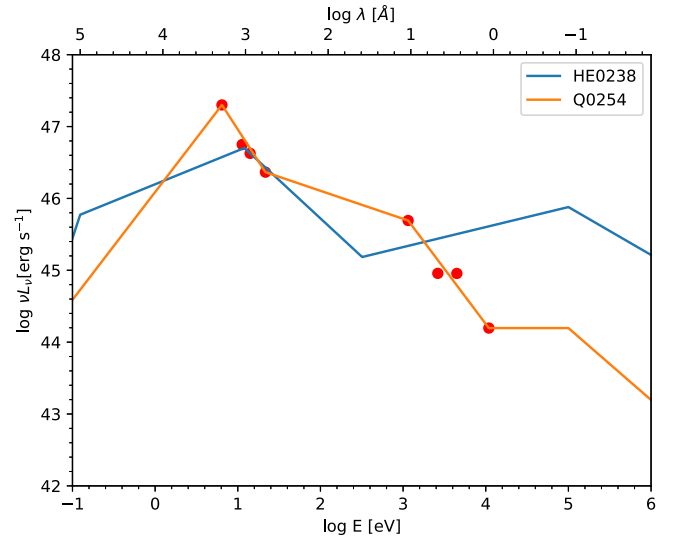


Figure 8. Comparison of SED shapes of HE0238 (Arav et al. 2013) and Q0254-334. The Q0254-334 SED was formed by using the V-band magnitude (the first red dot, Rao, Turnshek & Nestor 2006), UV continuum flux measured at three different wavelengths (rest wavelengths $\lambda = 574, 880, 1097 \text{ Å}$; second, third, and fourth red dots), and the X-ray fluxes reported by Chandra (5th–8th dots). The HE0238 SED was scaled to match the UV continuum flux for the sake of this comparison. We use the Q0254-334 SED for the analysis in this paper.

were fixed to be equal to each other. For the resulting fit, we found a ratio of 1, normalized amplitude $A = 0.23 \pm 0.11$, and $\text{FWHM} = 4800 \pm 900 \text{ km s}^{-1}$. This, along with the measured flux of $F_\lambda = 7.7^{+1.0}_{-1.0} \times 10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1}$ at rest wavelength 1050 Å , resulted in a black hole mass of $M_{\text{BH}} = 5.3^{+5.5}_{-2.7} \times 10^9 M_\odot$, and Eddington luminosity $L_{\text{Edd}} = 6.6^{+6.9}_{-3.4} \times 10^{47} \text{ erg s}^{-1}$. Note that we have limited our Gaussian fit to the red wing of the emission feature, as the blue wing has been contaminated by the absorption outflow (see Fig. 10). While this contamination has contributed significantly to the uncertainty, we were unable to find alternative emission features with which to estimate the black hole mass.

4 DISCUSSION

4.1 Distance of the outflow from the central source

With the parameters we found as described in Section 3, we could calculate the distance of the outflow from the central source, as well as the kinetic luminosity of the outflow. The distance can be found based on the definition of the ionization parameter U_{H} :

$$U_{\text{H}} \equiv \frac{Q_{\text{H}}}{4\pi R^2 n_{\text{H}} c}, \quad (6)$$

where Q_{H} is the emission rate of ionizing photons, R is the outflow distance from the source, n_{H} is the hydrogen number density, and c is the speed of light. Solving the equation for R gives us

$$R = \sqrt{\frac{Q_{\text{H}}}{4\pi U_{\text{H}} n_{\text{H}} c}}. \quad (7)$$

For highly ionized plasma, $n_{\text{e}} \approx 1.2 n_{\text{H}}$ (Osterbrock & Ferland 2006) and the values of U_{H} and n_{e} were found in Section 3.

We followed the method of other works (e.g. Miller et al. 2020a; Byun, Arav & Walker 2022a; Byun, Arav & Hall 2022c) to find Q_{H} and integrated over the SED mentioned in Section 3.3, limiting our

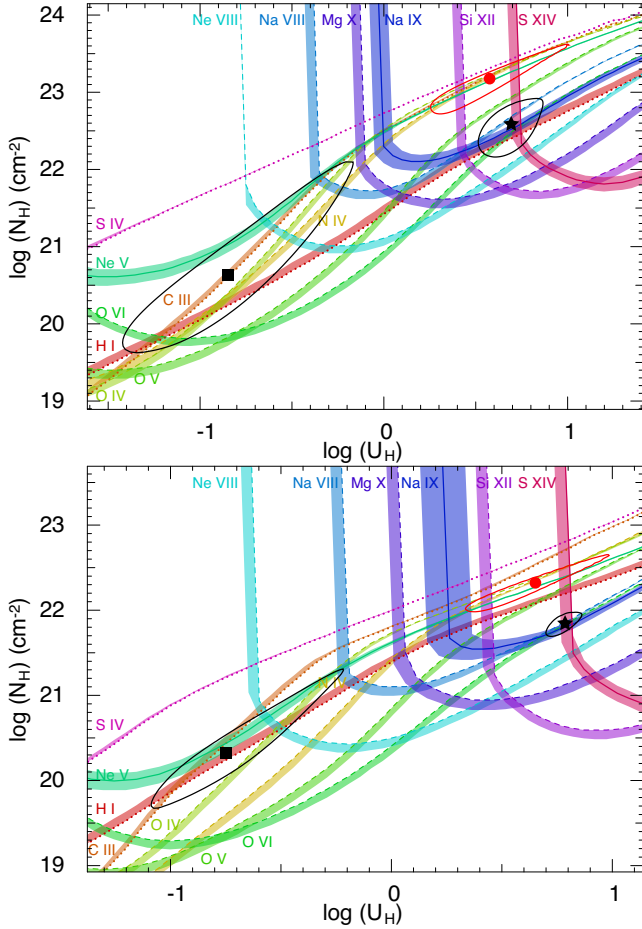


Figure 9. Plots of the Hydrogen column density (N_{H}) versus ionization parameter (U_{H}), assuming solar (top) and supersolar (bottom) metallicities, with the Q0254-334 SED shown in Fig. 8. Constraints on the parameters are based on measured column densities as shown in Table 1. Measurements are shown as solid curves, while upper and lower limits are represented with dotted and dashed curves, respectively. The coloured bands represent the uncertainties in the constraints. The red circle shows the one-phase solution of N_{H} and U_{H} , while the black square and star show the high-ionization and very high-ionization phases of the two-phase solution, respectively. The 1σ uncertainties of the solutions are shown as black/red ellipses.

Table 2. The measured and modelled column densities of H I, Na IX, and S XIV of the Q0254-334 outflow. The second and third columns denote the supersolar one-phase solution and the σ difference between modelled and measured values; the fourth and fifth columns show the same for the two-phase solution. The values are in units of 10^{14} cm^{-2} .

Ion	Measured	$4.68Z_{\odot}$ 1-p	$\Delta\sigma$	$4.68Z_{\odot}$ 2-p	$\Delta\sigma$
H I	$<38^{+11}_{-11}$	123	8	59	2
Na IX	48^{+12}_{-11}	96	4	55	0.6
S XIV	200^{+40}_{-40}	40	-4	190	-0.3

range to energies over 1 Ryd. This yielded the bolometric luminosity $L_{\text{Bol}} = 2.40^{+0.24}_{-0.24} \times 10^{47} \text{ erg s}^{-1}$ and $Q_{\text{H}} = 9.33^{+0.94}_{-0.94} \times 10^{56} \text{ s}^{-1}$. The distance estimates of the outflow calculated with this value are shown in Tables 3 and 5.

4.2 Contribution of the outflow to AGN feedback

For an outflow to contribute to AGN feedback, its kinetic luminosity must be at least ~ 0.5 (Hopkins & Elvis 2010) or ~ 5 per cent (Scannapieco & Oh 2004) of the quasar's Eddington luminosity. Assuming an incomplete spherical shell, the mass flow rate can be calculated as follows:

$$\dot{M} \simeq 4\pi\Omega R N_{\text{H}} \mu m_{\text{p}} v, \quad (8)$$

followed by the kinetic luminosity:

$$\dot{E}_{\text{k}} \simeq \frac{1}{2} \dot{M} v^2, \quad (9)$$

where Ω is the global covering factor, $\mu = 1.4$ is the mean atomic mass per proton, v is outflow velocity, and m_{p} is the mass of a proton (Borguet et al. 2012a). We assumed $\Omega = 0.2$, as C IV BALs are found in ~ 20 per cent of quasars (Hewett & Foltz 2003). We use the Ω associated with C IV BALs, since our high-ionization phase has troughs from ions of very similar ionization potential. For example, in our spectrum, we detect O IV $\lambda 787$. O IV has an ionization potential of 64 eV, which is quite similar to the C IV ionization potential of 77 eV. Assuming supersolar metallicity, this calculation yielded a kinetic luminosity of $\log \dot{E}_{\text{k}} = 44.76^{+0.20}_{-0.18} [\text{erg s}^{-1}]$ for the one-phase solution, and $\log \dot{E}_{\text{k}} = 44.98^{+0.15}_{-0.07} [\text{erg s}^{-1}]$ for the two-phase solution, leaving a ~ 0.2 dex difference between the solutions.

The ratio between the kinetic luminosity and Eddington luminosity yields $\dot{E}_{\text{k}}/L_{\text{Edd}} = 0.08^{+0.10}_{-0.05}$ per cent for the one-phase solution and $\dot{E}_{\text{k}}/L_{\text{Edd}} = 0.14^{+0.16}_{-0.07}$ per cent for the two-phase solution, which is below the 0.5 per cent threshold. For the sake of completeness, we have also found the ratio between \dot{E}_{k} and the bolometric luminosity L_{Bol} , resulting in $\dot{E}_{\text{k}}/L_{\text{Bol}} = 0.4^{+0.3}_{-0.2}$ and $0.7^{+0.3}_{-0.1}$ per cent (see Table 3). Based on the ratio between \dot{E}_{k} and L_{Edd} , the outflow would be unable to contribute to AGN feedback. It is important to note that the different assumed metallicity values have significant effects on the physical parameters of the outflow, such as a near order of magnitude difference in kinetic luminosity, leading to values that may be sufficient for AGN feedback contribution (see Table 3).

4.3 The two-phase outflow

As mentioned earlier in Section 3.3, the two-phase photoionization solution provides a better fit to the constraints from the measured ionic column densities. While the values of \dot{E}_{k} for the one-phase and two-phase solutions agree with each other within error (see Table 3), there are significant differences to be found in the other parameters, such as distance, N_{H} , and U_{H} .

Note that the difference in N_{H} between the high- and very high-ionization phases is ~ 1.5 orders of magnitude, as well as the difference in U_{H} . Assuming the two phases are co-spatial, the volume filling factor of the high-ionization phase is as follows (Arav et al. 2013; Miller et al. 2020a):

$$f_{\text{V}} = \frac{U_{\text{H,HP}}}{U_{\text{H,VHP}}} \times \frac{N_{\text{H,HP}}}{N_{\text{H,VHP}}} \quad (10)$$

resulting in $\log f_{\text{V}} = -3.1^{+1.1}_{-0.9}$, which follows our expectations from the high-ionization phase's larger n_{e} and smaller N_{H} values compared to those of the very high-ionization phase.

4.4 Connection to X-ray warm absorbers

The two-phase solution for the outflow of Q0254-334 is comparable to the parameters measured in X-ray warm absorbers. For instance, in their analysis of the Seyfert galaxy NGC 3783, Netzer et al. (2003)

Table 3. Physical properties of the Q0254-334 Outflow. The high- and very high-ionization phases for the two-phase solution are assumed to be co-spatial.

Metallicity Solution Phase	One-phase	Z_{\odot}		One-phase	$4.68Z_{\odot}$	
		High	Two-phase Very high		High	Two-phase Very high
$\log(N_{\text{H}})$ [cm ⁻²]	$23.18^{+0.45}_{-0.46}$	$20.63^{+1.47}_{-1.01}$	$22.59^{+0.33}_{-0.43}$	$22.32^{+0.33}_{-0.34}$	$20.32^{+0.99}_{-0.65}$	$21.84^{+0.13}_{-0.15}$
$\log(U_{\text{H}})$ [dex]	$0.58^{+0.43}_{-0.32}$	$-0.85^{+0.68}_{-0.57}$	$0.69^{+0.17}_{-0.18}$	$0.65^{+0.33}_{-0.32}$	$-0.75^{+0.53}_{-0.34}$	$0.79^{+0.08}_{-0.09}$
$\log(n_{\text{e}})^*$ [cm ⁻³]	$4.0^{+0.1}_{-0.1}$	$4.0^{+0.1}_{-0.1}$	$2.5^{+0.7}_{-0.6}$	$4.0^{+0.2}_{-0.2}$	$4.0^{+0.1}_{-0.1}$	$2.5^{+0.5}_{-0.4}$
Distance [pc]	100^{+50}_{-40}		500^{+490}_{-280}	90^{+40}_{-30}		450^{+240}_{-210}
\dot{M} [M _⊙ yr ⁻¹]	1400^{+1000}_{-700}		1800^{+800}_{-300}	180^{+90}_{-60}		290^{+50}_{-90}
$\dot{M}v$ [10 ³⁶ ergs cm ⁻¹]	28^{+21}_{-14}		37^{+16}_{-6}	4^{+2}_{-1}		6^{+1}_{-2}
$\log(\dot{E}_{\text{k}})$ [ergs ⁻¹]	$45.65^{+0.30}_{-0.24}$		$45.78^{+0.07}_{-0.16}$	$44.76^{+0.20}_{-0.18}$		$44.98^{+0.15}_{-0.07}$
$\dot{E}_{\text{k}}/L_{\text{Edd}}$ [per cent]	$0.7^{+1.1}_{-0.4}$		$0.9^{+0.9}_{-0.5}$	$0.08^{+0.10}_{-0.05}$		$0.14^{+0.16}_{-0.07}$
$\dot{E}_{\text{k}}/L_{\text{Bol}}$ [per cent]	$3.5^{+3.6}_{-1.5}$		$4.6^{+1.0}_{-1.4}$	$0.4^{+0.3}_{-0.2}$		$0.7^{+0.3}_{-0.1}$

Note. *The n_{e} of the very high-ionization phase is the n_{e} of the high-phase times the ratio of the high/very high-ionization parameters.

Table 4. Photoionization solution for the Q0254-334 outflow assuming the HE0238 SED.

Metallicity Solution Phase	One-phase	Z_{\odot}		One-phase	$4.68Z_{\odot}$	
		High	Two-phase Very high		High	Two-phase Very high
$\log(N_{\text{H}})$ [cm ⁻²]	$23.08^{+0.38}_{-0.38}$	$20.83^{+0.85}_{-0.59}$	$22.58^{+0.16}_{-0.19}$	$22.29^{+0.36}_{-0.34}$	$20.61^{+0.98}_{-0.94}$	$21.84^{+0.48}_{-0.19}$
$\log(U_{\text{H}})$ [dex]	$0.88^{+0.30}_{-0.24}$	$-0.39^{+0.47}_{-0.40}$	$1.00^{+0.06}_{-0.08}$	$0.87^{+0.29}_{-0.25}$	$-0.24^{+0.54}_{-0.59}$	$1.01^{+0.07}_{-0.09}$

Table 5. Physical Properties of the Q0254-334 Outflow, based on the Si XII template-based fitting of the Ne V absorption. The high- and very high-ionization phases for the two-phase solution are assumed to be co-spatial.

Metallicity Solution Phase	One-phase	Z_{\odot}		One-phase	$4.68Z_{\odot}$	
		High	Two-phase Very high		High	Two-phase Very high
$\log(n_{\text{e}})$ [cm ⁻³]	$4.3^{+0.1}_{-0.1}$	$4.3^{+0.1}_{-0.1}$	$2.8^{+0.7}_{-0.6}$	$4.3^{+0.1}_{-0.1}$	$4.3^{+0.1}_{-0.1}$	$2.8^{+0.5}_{-0.3}$
Distance [pc]	70^{+30}_{-20}		370^{+350}_{-200}	70^{+30}_{-20}		330^{+170}_{-150}
\dot{M} [M _⊙ yr ⁻¹]	1000^{+700}_{-500}		1300^{+500}_{-200}	130^{+60}_{-50}		210^{+30}_{-60}
$\dot{M}v$ [10 ³⁶ ergs cm ⁻¹]	20^{+15}_{-10}		27^{+10}_{-3}	$2.6^{+1.3}_{-0.9}$		$4.4^{+0.6}_{-1.2}$
$\log(\dot{E}_{\text{k}})$ [ergs ⁻¹]	$45.52^{+0.30}_{-0.24}$		$45.64^{+0.06}_{-0.15}$	$44.63^{+0.19}_{-0.17}$		$44.85^{+0.14}_{-0.06}$
$\dot{E}_{\text{k}}/L_{\text{Edd}}$ [per cent]	$0.5^{+0.8}_{-0.3}$		$0.6^{+0.6}_{-0.3}$	$0.06^{+0.08}_{-0.03}$		$0.10^{+0.11}_{-0.05}$
$\dot{E}_{\text{k}}/L_{\text{Bol}}$ [per cent]	$2.6^{+2.6}_{-1.1}$		$3.4^{+0.6}_{-1.0}$	$0.33^{+0.19}_{-0.11}$		$0.55^{+0.23}_{-0.08}$

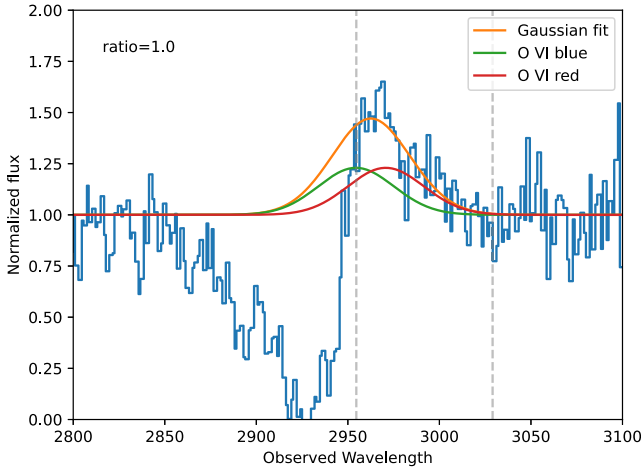


Figure 10. Gaussian fitting of the O VI emission feature. The dashed vertical lines denote the range of data used for the Gaussian fit. The green and red curves are the modelled blue and red emission features, respectively, and the orange curve represents the combined modelled emission.

found the parameters of the absorbing gas composed of three different components, with the oxygen ionization parameter ranging from $\log U_{\text{ox}} = -2.4$ to -0.6 . To effectively compare the U_{H} of Q0254-334 to the U_{ox} values of NGC 3783, we calculated the oxygen ionizing emission rate Q_{ox} as defined below:

$$Q_{\text{ox}} = \int_{\nu(0.54 \text{ keV})}^{\nu(10 \text{ keV})} \frac{L_{\nu}}{h\nu} d\nu \quad (11)$$

such that the ratio $\frac{Q_{\text{ox}}}{Q_{\text{H}}} = \frac{U_{\text{ox}}}{U_{\text{H}}}$. The resulting value of the emission rate was $Q_{\text{ox}} = 3.9^{+0.4}_{-0.4} \times 10^{54} \text{ s}^{-1}$, which is 2.4 ± 0.1 orders of magnitude smaller than Q_{H} . Subtracting 2.4 ± 0.1 from the $\log U_{\text{H}}$ values of the high- and very high-ionization phases leads to $\log U_{\text{ox}} = -3.2^{+0.5}_{-0.5}$ and $-1.6^{+0.3}_{-0.3}$, respectively. The very high-ionization phase has a U_{ox} within the range of U_{ox} values of the NGC 3783 absorbing gas. We note that NGC 3783 is a much lower luminosity AGN than Q0254-334, and that its SED may be different. However, lacking high-quality X-ray spectra of $z \sim 1$ quasars, it is still illuminating to compare the NGC 3783 X-ray wind with the EUV wind seen in Q0254-334.

4.5 Comparison to other extreme UV objects

As the spectrum of Q0254-334 covers observed wavelengths as short as 400 \AA , we found it appropriate to compare it with other quasars observed in the extreme UV range (hereafter EUV500, Arav et al. 2020). We compiled a list of the physical parameters of 28 EUV500 quasar outflow systems analysed in previous works (Arav et al. 2020; Miller et al. 2020a; Xu et al. 2020a; Miller et al. 2020b; Xu et al. 2020b; Miller et al. 2020c; Xu et al. 2020c), and added the parameters of Q0254-334 for comparison, with a total of 29 EUV500 outflow systems. Out of the 29 outflow systems, 24, including the outflow discussed in this paper, have measurements of kinetic luminosity and distance from the source. We compared the parameters of the Q0254-334 outflow such as \dot{E}_{k} , N_{H} , R , and U_{H} , with the other 23 outflow systems.

As seen in Fig. 11, no strong correlation has been found between $\log \dot{E}_{\text{k}}$ and $\log L_{\text{Edd}}$ or between $\log \dot{E}_{\text{k}}$ and $\log L_{\text{Bol}}$. Four of the 24 outflows (~ 16 per cent) are above the threshold of $\dot{E}_{\text{k}}/L_{\text{Edd}} \sim 5$ per cent, while 7 (~ 29 per cent) are between the 0.5 and 5 per cent thresholds. With regards to L_{Bol} , five of the outflow systems (~ 20 per cent)

are above the 5 per cent threshold, while seven (~ 29 per cent) are between the 0.5 and 5 per cent thresholds. Note that while the values of $\log \dot{E}_{\text{k}}$ range between 41–47, $\log L_{\text{Edd}}$ and $\log L_{\text{Bol}}$ range between 47.0–47.9 and 46.6–47.6, respectively, which is much narrower than the range of $\log \dot{E}_{\text{k}}$. It is also indicative of the ability of line-of-sight analysis to identify outflow systems at large ranges of kinetic luminosity as well as velocity. Fig. 12 shows the $\log N_{\text{H}}$ and $\log U_{\text{H}}$ values of the high- and very high-ionization phases of each of the outflow systems. With the exception of the very high-ionization phase of the outflow system of UM425 traveling at -9420 km s^{-1} , the high-ionization phases tend to have values of $\log U_{\text{H}} < 0$, while the very high-ionization phases have $\log U_{\text{H}} > 0$. Note that the very high-ionization phase of the Q0254-334 outflow has a higher $\log U_{\text{H}}$ value relative to the average of the other outflows. We largely attribute this to the detection of S XIV. As can be seen in Fig. 9, the very high-ionization phase solution is at the intersection between the Na IX and S XIV constraints. The $\log U_{\text{H}}$ value is ~ 0.3 dex higher than what it would have been if S XIV were not detected, and the parameters were constrained by other ions such as Si XII. It is also notable that as shown in Fig. 12, the $\log U_{\text{H}}$ of the very high-ionization phase of the Q0254-334 outflow is higher than the average $\log U_{\text{H}}$ of the other outflows. We suspect that future observed outflows with S XIV would yield comparably high U_{H} values.

Note that there is an apparent edge in the range of $\log N_{\text{H}}$ and $\log U_{\text{H}}$ values of the outflows. In Fig. 12, we have indicated the approximate locations of the hydrogen ionization front ($N_{\text{H}} \approx 10^{23} U_{\text{H}}$), as well as the He II ionization front ($N_{\text{H}} \approx 10^{22.2} U_{\text{H}}$). The $N_{\text{H}}/U_{\text{H}}$ ratio for the He II ionization front was calculated based on the average $\log N_{\text{H}}$ value at which the He II to He III ratio is 1:1 in a series of CLOUDY models created with a range of $-2.0 < \log U_{\text{H}} < 1.0$. We used the aforementioned SED of HE0238 for the models, as this SED was used for the analysis of the majority of the EUV500 outflows in question. Interestingly, the $\log N_{\text{H}}$ versus $\log U_{\text{H}}$ values of all of the EUV500 outflows fall under the He II ionization front, which would suggest that they are high ionized BALs (HiBALs) as opposed to low ionized BALs (LoBALs). This is supported by the lack of BALs from low-ionization species. For instance, there is a noticeable lack of absorption where the trough of C II $\lambda 687$ would be, despite a large oscillator strength of $f = 0.336$ (see Fig. 13).

We also examined the ranges of R and v of the outflow systems (see Fig. 14). To examine the correlation between distance and velocity, we conducted a weighted least squares linear fit between $\log R$ and $\log |v|$, taking into account the asymmetry of the reported errors in R . We adopted the weight determination method described by Barlow (2003). The weight of each data point $w_i = 1/V_i$ was determined by the value of V_i :

$$V_i = \sigma_i^2 + \left(1 - \frac{2}{\pi}\right) \alpha_i^2 \quad (12)$$

in which $\sigma_i = \frac{\sigma_i^+ + \sigma_i^-}{2}$ is the mean of the upper and lower errors of $\log R$, while $\alpha_i = \frac{\sigma_i^+ - \sigma_i^-}{2}$.

The weighted linear fit yielded a slope of -1.08 and an intercept of 6.44 , suggesting a negative correlation. To determine the strength of the correlation, we calculated a modified value of the coefficient of determination r^2 that would take into account the weight of each data point. The residual sum of squares SS_{res} was modified so that:

$$SS_{\text{res}} = \sum_i w_i (\log R_i - f_i)^2, \quad (13)$$

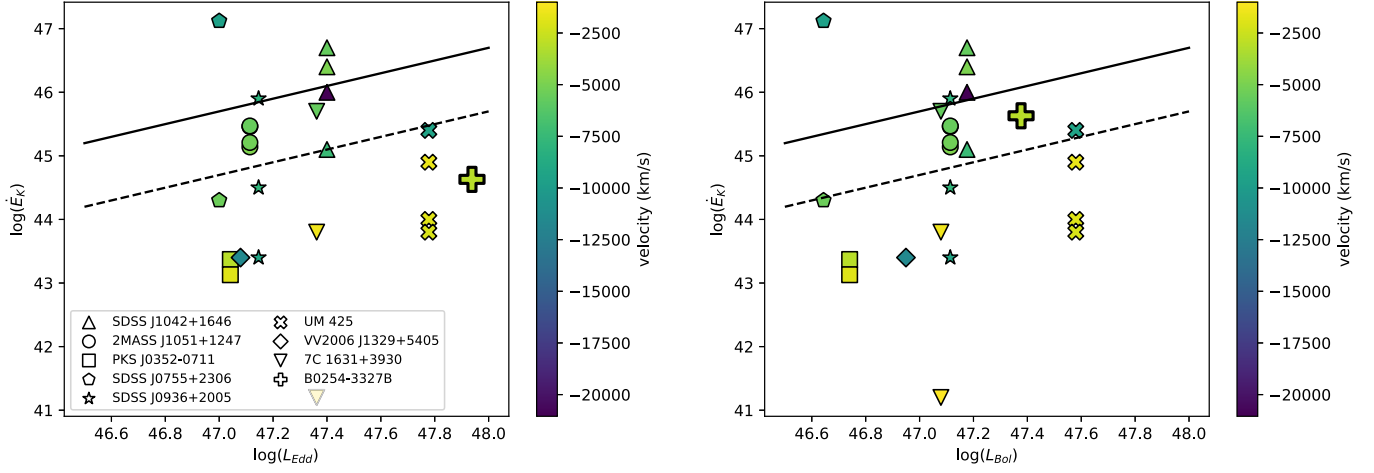


Figure 11. Distribution of $\log \dot{E}_k$ versus $\log L_{\text{Edd}}$ (left) and $\log L_{\text{Bol}}$ (right) of EUV500 outflows. The dashed and solid lines on the left (right) indicate the \dot{E}_k/L_{Edd} (\dot{E}_k/L_{Bol}) thresholds of 0.5 and 5 per cent, respectively. The plus-sign symbol denotes the outflow of Q0254-334 as reported in this paper, while the other symbols denote the parameters of other EUV outflows reported by Arav et al. (2020) and Miller et al. (2020c). The colour map corresponds to the velocities of the outflow systems.

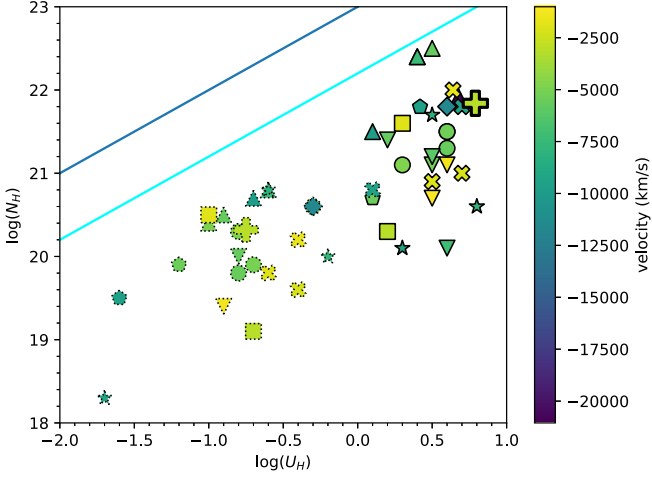


Figure 12. Distribution of $\log N_{\text{H}}$ versus $\log U_{\text{H}}$ of EUV500 outflows. Symbols are coded as they are in Fig. 11. Symbols with dotted outlines denote high-ionization phases, while symbols with solid outlines denote very high-ionization phases. The colour map corresponds to the velocities of the outflow systems. The dark blue and cyan curves show the H I and He II ionization fronts, respectively.

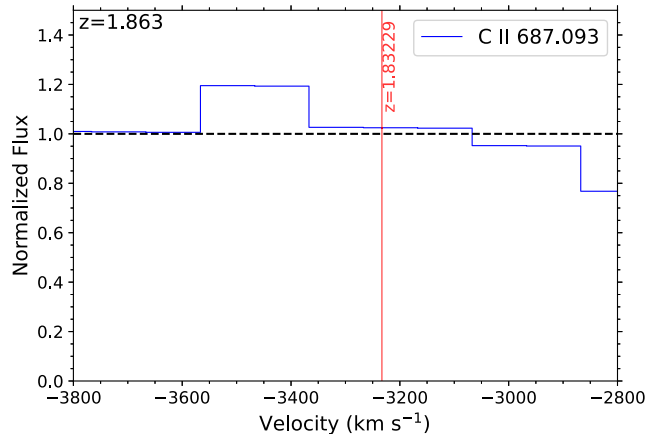


Figure 13. Region in which C II $\lambda 687$ absorption is expected to be found.

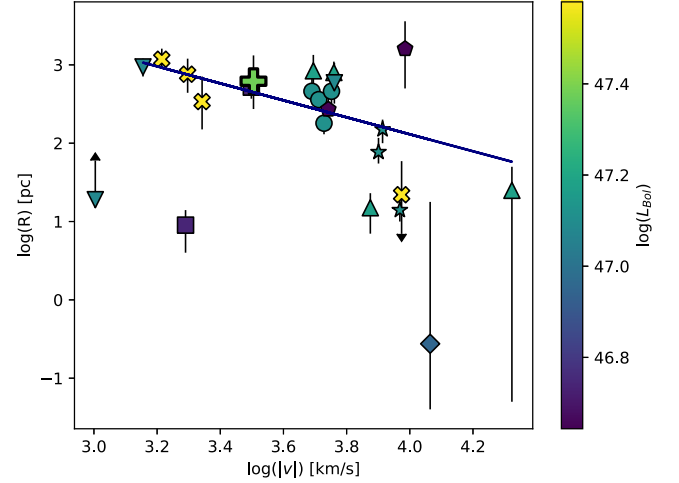


Figure 14. Distribution of $\log R$ versus $\log |v|$ of EUV500 outflows. Symbols are coded as they are in Fig. 11, with errors in $\log R$ indicated with error bars. The dark blue line shows the weighted least squares linear fit of $\log R = -1.08 \times \log |v| + 6.44$, with adjusted $r^2 = 0.28$. The colour map corresponds to the $\log L_{\text{Bol}}$ values of the outflows.

where f_i is the value of $\log R$ according to the linear fit. The total sum of squares SS_{tot} was adjusted so that:

$$SS_{\text{tot}} = \sum_i w_i (\log R_i - \langle \log R \rangle)^2, \quad (14)$$

where $\langle \log R \rangle$ is the weighted mean of $\log R$. The resulting value of $r^2 = 1 - \frac{SS_{\text{res}}}{SS_{\text{tot}}}$ is 0.28, suggesting a weak negative correlation between $\log R$ and $\log |v|$. This correlation is further supported by a Spearman correlation of -0.43 and associated p -value of 0.05. We find it worth mentioning that Choi et al. (2022a) have conducted a similar analysis of iron low ionized BALs (FeLoBALs), in which they determined that the correlation (or anticorrelation) between distance and velocity was dependent on the E1 parameter, which is formulated based on the relationship between the equivalent width of [O III] emission and the ratio between Fe II and H β fluxes (Leighly et al. 2022).

5 SUMMARY AND CONCLUSION

We have identified a BAL outflow in the *HST*/STIS spectrum of the quasar QSO B0254-3327B, of which we have found the ionic column densities (see Table 1). Based on the column densities, we conducted photoionization analysis to find the values of hydrogen column density N_{H} and ionization parameter U_{H} . The results of our analysis are as follows:

(i) *The two-phase solution.* The constraints from the measured ionic column densities required a solution with two ionization phases: the high-ionization phase and the very high-ionization phase. The two-phase solution showed a significant improvement in the χ^2 value compared to that of the one-phase solution (reduced $\chi^2 = 5.1$ versus 22.3).

(ii) *The energetics of the outflow.* We were able to determine the composition of the Ne V trough via Gaussian fitting of the blended features (see Fig. 6), thanks to which we were able to narrow down the electron number density n_e . Through the use of equations (7), (8), and (9), we were able to determine the distance, mass flow rate, and kinetic luminosity of the outflow (see Table 3). There were notable differences in the energetics parameters based on different values of assumed metallicity ($Z = Z_{\odot}$ versus $Z = 4.68Z_{\odot}$, see Tables 3 and 5).

(iii) *Potential contribution to AGN feedback.* As the ratio between the kinetic luminosity and the quasar's Eddington luminosity $\dot{E}_k/L_{\text{Edd}} = 0.9^{+0.9}_{-0.5}$ per cent assuming solar abundance and $0.14^{+0.16}_{-0.07}$ assuming supersolar abundance, its contribution to AGN feedback is model dependent, as the theoretical thresholds for the ratio are ~ 0.5 (Hopkins & Elvis 2010) and ~ 5 per cent (Scannapieco & Oh 2004).

(iv) *Comparison to X-ray warm absorbers.* We have compared the ionization parameter values of the high-ionization and very high-ionization phase to that of the X-ray warm absorber of NGC 3783 analysed by Netzer et al. (2003). Converting U_{H} to the oxygen ionization parameter U_{ox} as defined by Netzer et al. (2003) showed that the U_{ox} of the very high-ionization phase agreed with that of the NGC 3783 absorber within the error.

(v) *Comparison with other EUV500 outflows.* We have also examined the physical parameters of previously studied EUV500 outflows (Arav et al. 2020; Miller et al. 2020a; Xu et al. 2020a; Miller et al. 2020b; Xu et al. 2020b; Miller et al. 2020c; Xu et al. 2020c) and determined that out of the sample of 24 outflow systems with measured kinetic luminosity, up to ~ 50 per cent may contribute to AGN feedback, depending on the theoretical model. The trend between $\log R$ and $\log |v|$ was also analysed via a weighted least squares linear fit, showing a weak negative correlation, indicated by a Spearman rank of -0.43 , and p -value of 0.05 . We have also found that the very high-ionization phase of the Q0254-334 outflow had one of the highest U_{H} values of all UV absorption outflows to date.

The process of finding $\log n_e$ was limited to the examination of a blended trough of Ne V. Further observations and analyses of the quasar may reveal more excited state troughs, which could help to improve the uncertainty in $\log n_e$. Studying additional EUV500 outflows will be essential to a more thorough statistical analysis of their parameters as well.

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DATA AVAILABILITY

The STIS data of Q0254-334 described in this paper may be obtained from the MAST archive at <https://dx.doi.org/10.17909/63xd-4s51>. The FOS data shown for comparison can be obtained at <https://dx.doi.org/10.17909/z7ta-9048>.

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