

Disc settling and dynamical heating: histories of Milky Way-mass stellar discs across cosmic time in the FIRE simulations

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Accepted 2023 October 27. Received 2023 September 26; in original form 2023 March 28

ABSTRACT

We study the kinematics of stars both at their formation and today within 14 Milky Way (MW)-mass galaxies from the FIRE-2 cosmological zoom-in simulations. We quantify the relative importance of cosmological disc settling and post-formation dynamical heating. We identify three eras: a Pre-Disc Era (typically $\gtrsim 8$ Gyr ago), when stars formed on dispersion-dominated orbits; an Early-Disc Era (≈ 8 –4 Gyr ago), when stars started to form on rotation-dominated orbits but with high velocity dispersion, σ_{form} ; and a Late-Disc Era ($\lesssim 4$ Gyr ago), when stars formed with low σ_{form} . σ_{form} increased with time during the Pre-Disc Era, peaking ≈ 8 Gyr ago, then decreased throughout the Early-Disc Era as the disc settled and remained low throughout the Late-Disc Era. By contrast, the dispersion measured today, σ_{now} , increases monotonically with age because of stronger post-formation heating for Pre-Disc stars. Importantly, most of σ_{now} was in place at formation, not added post-formation, for stars younger than ≈ 10 Gyr. We compare the evolution of the three velocity components: at all times, $\sigma_{\text{R, form}} > \sigma_{\phi, \text{form}} > \sigma_{\text{Z, form}}$. Post-formation heating primarily increased σ_{R} at ages $\lesssim 4$ Gyr but acted nearly isotropically for older stars. The kinematics of young stars in FIRE-2 broadly agree with the range observed across the MW, M31, M33, and PHANGS-MUSE galaxies. The lookback time that the disc began to settle correlates with its dynamical state today: earlier-settling galaxies currently form colder discs. Including stellar cosmic-ray feedback does not significantly change disc rotational support at fixed stellar mass.

Key words: methods: numerical – Galaxy: evolution – Galaxy: formation – galaxies: disc – galaxies: evolution – galaxies: kinematics and dynamics.

1 INTRODUCTION

The present-day kinematics of stars encode a galaxy’s dynamical history. A star’s orbit is set by both its initial dynamics, as inherited from the star-forming interstellar medium (ISM) at its formation, and all post-formation dynamical perturbations that the star experienced over its lifetime. Thus, stellar kinematics provide dual insight into the past state of the ISM and the dynamical processes that shaped the galaxy. However, the extent to which a star’s current motion results from its initial state or from post-formation perturbations remains poorly understood. The inability to reliably disambiguate these two remains a significant obstacle in connecting the kinematics and morphology of a galaxy today to its formation history (for example Somerville & Davé 2015; Naab & Ostriker 2017).

Overcoming this ambiguity requires understanding the kinematics with which stars formed throughout the entire lifetime of a galaxy. In the solar neighbourhood, the youngest stellar populations share the kinematics of the gas disc from which they formed: young

stars are on nearly circular orbits with total velocity dispersions, $\sigma_{\text{tot}} \approx 25 – 30 \text{ km s}^{-1}$ (Edvardsson et al. 1993; Casagrande et al. 2011). However, the kinematics of stars that formed at earlier times are likely different from those forming in the Milky Way (MW) today, based on observations of high-redshift galaxies, which provide statistics on how the initial state of stars evolved over cosmic time. Such observations reveal that MW-like thin discs today are the endpoint of a metamorphosis: galactic discs at higher redshifts ($z \gtrsim 1$ –2) were thicker, clumpier, and more turbulent than their thin counterparts at lower redshifts (for example Förster Schreiber et al. 2009; Stott et al. 2016; Birkin et al. 2023; though see Rizzo et al. 2020). The high turbulence of MW-mass progenitors probably resulted from the combination of stellar feedback, given their higher star formation rates (SFRs; for example Madau & Dickinson 2014; Förster Schreiber & Wuyts 2020), gravitational instabilities (for example Elmegreen et al. 2007; Dekel et al. 2009; Lehnert et al. 2009), and higher merger and accretion rates, leading to higher gas fractions and surface densities (for example Genzel et al. 2015; Tacconi, Genzel & Sternberg 2020).

Thus, high-redshift observations show that disc galaxies ‘settled’ over cosmic time (Kassin et al. 2012), becoming less turbulent and

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more rotationally supported since $z \sim 1-2$ (for example Wisnioski et al. 2015; Johnson et al. 2018). Specifically, the velocity dispersion, σ , of gas discs decreased while their rotational velocity, v_ϕ , increased. Although most works used emission lines from ionized gas to measure high-redshift kinematics, recent works indicate that the dispersions of atomic and molecular gas, while $\approx 15 \text{ km s}^{-1}$ lower, evolved similarly (for example Übler et al. 2019; Girard et al. 2021). Thus, stars, which are born from cold dense gas, formed with larger dispersions at earlier times when turbulence was higher, and later formed with progressively lower dispersions as turbulence decreased and the galaxy's disc settled.

Alongside understanding how stars kinematically formed, it is equally important to understand how they evolved post-formation. Over their lifetimes, stars undergo repeated scattering interactions with perturbations within the disc, such as giant molecular clouds (GMCs; Spitzer & Schwarzschild 1951; Wielen 1977; Lacey 1984), spiral arms (Barbanis & Woltjer 1967; Carlberg & Sellwood 1985; Minchev & Quillen 2006), and bars (Saha, Tseng & Taam 2010; Grand et al. 2016). Stars are effectively collisionless and dissipationless, and thus these interactions result in enduring increases to their random orbital energy, or 'dynamical heating'. In turn, a stellar population's σ increases over its lifetime. Additional heating may arise from feedback-driven outflows and potential fluctuations (El-Badry et al. 2016), the galaxy's global asymmetries – including torquing from the misalignment of the stellar and gaseous disc (Roškar et al. 2010) – cosmic accretion, mergers, and interactions with satellites (for example Quinn, Hernquist & Fullagar 1993; Brook et al. 2004; Villalobos & Helmi 2008; Belokurov et al. 2018). However, the relative impact of each mechanism remains uncertain.

Furthermore, the scattering environment of the disc evolves across cosmic time. At earlier times, a galaxy's gas fraction and gas surface density were higher (for example Shapley 2011; Swinbank et al. 2012; Tacconi, Genzel & Sternberg 2020), such that scattering by GMCs was likely more frequent. Prior to the formation of a disc, galaxies generally lacked strong spiral structure and thus spiral-driven heating (for example Margalef-Bentabol et al. 2022). On the other hand, certain processes that can drive non-adiabatic rapid heating were more prevalent earlier, including mergers (for example Conselice, Blackburne & Papovich 2005) and feedback-driven outflows (for example Genzel et al. 2015).

One can understand the relative strength of these heating mechanisms by comparing the observed relationship between velocity dispersion and stellar age with predictions from various heating models. The magnitude of each component of σ (along R , ϕ , or Z) provides insight into the relative strength of different mechanisms: spiral structure increases the in-plane dispersion, but is generally inefficient at increasing the vertical, whereas GMCs heat stars more isotropically and can redirect in-plane heating (for example Carlberg & Innanen 1987; Jenkins & Binney 1990; Sellwood 2013).

The MW is an ideal test-bed for such comparisons because it offers unparalleled access to the 3D motions, positions, and ages of individual stars. In particular, the observed monotonic increase in velocity dispersion with age for stars in the solar neighbourhood provides an important benchmark (for example Strömborg 1946; Wielen 1977; Casagrande et al. 2011). Nordström et al. (2004) and Holmberg, Nordström & Andersen (2009) argued that this observed trend is consistent with MW stars forming on 'discy' orbits that then underwent dynamical heating by the combination of GMCs and spiral arms. Mackereth et al. (2019b) and Ting & Rix (2019) extended this analysis to include stars located outside the solar neighbourhood ($R = 3-14 \text{ kpc}$) and found similar results: the observed relations are consistent with stars younger than $\approx 8 \text{ Gyr}$ having formed with a low,

constant dispersion and subsequently being heated by both GMCs and spiral arms. However, as before, the observed dispersions of older stars (age $\gtrsim 8 \text{ Gyr}$) are higher than those predicted by heating models (Mackereth et al. 2019b). This discrepancy could arise if old stars were born kinematically hotter, or if their dynamical heating differed from what current models predict. Likely, both factors contribute, but disentangling this requires insight into the MW's early history.

Fortunately, recent MW observations now allow us to characterize this bygone era. Astrometric data from Gaia (Gaia Collaboration 2016) along with stellar properties from large spectroscopic surveys, such as APOGEE (Majewski et al. 2017), GALAH (De Silva et al. 2015; Buder et al. 2018), and LAMOST (Cui et al. 2012), now provide 6D kinematics and elemental abundances for stars throughout the Galaxy. Most notably, this data revealed that the MW underwent a massive merger – the Gaia-Sausage-Enceladus (GSE) – $\approx 8-11 \text{ Gyr}$ ago (for example Belokurov et al. 2018; Helmi et al. 2018). The GSE merger, probably the most major merger of the MW, contributed most of the stars in the MW's stellar halo, and it likely dynamically heated MW stars to more eccentric halo-like orbits, and potentially it triggered intense star-formation and even the formation of the MW's low- α disc (for example Haywood et al. 2018; Gallart et al. 2019; Mackereth et al. 2019a; Bonaca et al. 2020; Das, Hawkins & Jofré 2020). Recent works suggest that the MW developed this thick – but continuously-settling – disc early in its history, with estimates for the onset of coherent rotation ranging from ≈ 10 to 13 Gyr ago (Belokurov & Kravtsov 2022; Conroy et al. 2022; Xiang & Rix 2022), and that the MW's thin-disc began to form later, $\approx 8-9 \text{ Gyr}$ ago. However, post-formation dynamical heating leads to ambiguity about the nature and timing of these transitions, highlighting the difficulties inherent to separating birth kinematics from post-formation dynamical heating, especially for old metal-poor stars.

Cosmological simulations of the formation of MW-mass galaxies provide important tools for understanding the relative impacts of cosmological disc settling and post-formation dynamical heating (see Naab & Ostriker 2017; Vogelsberger et al. 2020, for recent reviews). Cosmological zoom-in simulations provide the high resolution needed to model a galaxy's internal structure, including the multiphase ISM that sets the birth kinematics of stars and thus the dynamical structure of the disc, within an accurate cosmological context, allowing for non-equilibrium dynamical perturbations, mergers, mass growth, and so forth. Many works used such cosmological zoom-in simulations to show that (1) MW-like galaxies form radially 'inside-out' and vertically 'upside-down', transitioning from a thick, turbulent galaxy at early times to a thin rotationally dominated disc at later times, and (2) the present-day kinematics of stars result from their kinematics at birth and subsequent dynamical heating, both of which change over cosmic time (for example Bird, Kazantzidis & Weinberg 2012; Brook et al. 2012; Martig, Minchev & Flynn 2014; Grand et al. 2016; Bonaca et al. 2017; Ceverino et al. 2017; Ma et al. 2017; Buck et al. 2020; Agertz et al. 2021; Bird et al. 2021).

We build upon a series of papers that used FIRE-2 simulations to study the formation of galactic discs. Stern et al. (2021) found that virialization of the inner circumgalactic medium (CGM), that is, the onset of a stable hot halo, coincided with the formation of a thin disc, a change from 'bursty' to 'steady' star formation, and the suppression of gaseous outflows; the presence of a hot halo can allow infalling gas to align coherently with a galaxy's disc before it accretes onto the disc (Hafen et al. 2022). Additional works found that the virialization of the inner CGM and the associated transition from bursty to steady star formation coincided with the transition from 'thick-disc' to 'thin-disc' formation (Yu et al. 2021, 2022)

and the emergence of a rotationally supported thin gaseous disc from a previously quasi-isotropic ISM (Gurvich et al. 2023). More recently, Hopkins et al. (2023) investigated the physical causes of such transitions. While the initial formation of a disc depends on the development of a sufficiently centrally concentrated mass profile, the transition from bursty to smooth star formation generally arises later, when the absolute depth of the gravitational potential (within the radius of star formation) crosses a critical threshold. This mass threshold is also similar to the threshold for inner CGM virialization.

In this work, we use the FIRE-2 cosmological zoom-in simulations to study cosmological disc settling and dynamical heating across the formation histories of MW-mass galaxies. We study the kinematics of stars both at birth and today, emphasizing different phases of disc-galaxy formation. We examine how stellar formation kinematics manifest themselves today, that is, how post-formation dynamical heating alters the initial orbits of stars. By doing so, we aim to understand how the current kinematics of the MW and nearby galaxies relate to their formation histories.

2 METHODS

2.1 FIRE-2 simulations

We analyse 14 cosmological zoom-in simulations of MW/M31-mass galaxies from the Feedback in Realistic Environments (FIRE) project,¹ which include both dark matter and baryons (gas and stars). The simulations use the GIZMO gravity plus hydrodynamics code (Hopkins 2015) with the meshless finite-mass (MFM) Godunov method, which provides adaptive spatial resolution while maintaining excellent conservation of energy and (angular) momentum.

We ran these simulations using the FIRE-2 physics model (Hopkins et al. 2018). FIRE-2 models the dense, multiphase ISM and incorporates metallicity-dependent radiative cooling and heating processes for gas across temperatures $10\text{--}10^{10}$ K, including free-free, photoionization and recombination, Compton, photoelectric and dust collisional, cosmic ray (CR), molecular, metal-line, and fine structure processes. The simulations self-consistently generate and track 11 elements (H, He, C, N, O, Ne, Mg, Si, S, Ca, Fe), including a model for sub-grid mixing/diffusion via turbulence (Hopkins 2016; Su et al. 2017; Escala et al. 2018). FIRE-2 also includes photoionization and heating from a spatially uniform redshift-dependent ultraviolet (UV) background (Faucher-Giguère et al. 2009) that reionizes the universe at $z \approx 10$.

Crucially for this work, FIRE-2 simulations resolve the phase structure in the ISM, allowing gas to collapse into GMCs and form stars (Benincasa et al. 2020; Guszejnov et al. 2020). Star formation occurs in locally self-gravitating, Jeans-unstable, dense ($n_{\text{SF}} > 1000 \text{ cm}^{-3}$), molecular (following Krumholz & Gnedin 2011) gas. After reaching these criteria, a gas cell probabilistically converts to a star particle in a local free-fall time. Our simulations, like all simulations, have a finite resolution. As such, our star-formation model is an extrapolation of the collapse below the resolution scale. That said, numerous resolution tests carried out in Hopkins et al. (2018) indicate that uncertainties induced by our resolution limit are minor; for example, the GMC mass function and the relation between internal gas velocity dispersion and GMC size for these simulations are both independent of resolution.

Each star particle represents a single stellar population sampled from a Kroupa (2001) initial mass function, with mass and metallicity

inherited from its progenitor gas cell. Once formed, star particles follow stellar evolution models that tabulate feedback event rates, luminosities, energies, and mass-loss rates from STARBURST99 v7.0 (Leitherer et al. 1999). FIRE-2 implements the major channels for stellar feedback, including core-collapse and white-dwarf (Ia) supernovae, stellar winds, radiation pressure, photoionization, and photoelectric heating.

We generated cosmological zoom-in initial conditions at $z \approx 99$ using MUSIC (Hahn & Abel 2011). Each zoom-in region is embedded within a cosmological box of side length 70.4–172 Mpc. The simulations use a flat ΛCDM cosmology with parameters broadly consistent with Planck Collaboration (2020): $h = 0.68\text{--}0.71$, $\Omega_{\Lambda} = 0.69\text{--}0.734$, $\Omega_{\text{m}} = 0.266\text{--}0.31$, $\Omega_{\text{b}} = 0.0455\text{--}0.048$, $\sigma_8 = 0.801\text{--}0.82$, and $n_s = 0.961\text{--}0.97$.

Our simulation set consists of 14 MW/M31-mass galaxies: eight isolated galaxies from the Latte suite, and six galaxies from the ELVIS on FIRE suite of Local Group (LG)-like pairs of galaxies. The Latte suite, with the exception of m12z, has an initial baryon particle mass of 7070 M_{\odot} , although stellar mass loss leads to star particles having typical masses $\approx 5000 \text{ M}_{\odot}$, and dark matter particle mass of $3.5 \times 10^5 \text{ M}_{\odot}$. m12z has initial baryon and dark matter particle masses of 4200 M_{\odot} and $2.1 \times 10^5 \text{ M}_{\odot}$, respectively. The ELVIS on FIRE suite has a mass resolution $\sim 2 \times$ better than the Latte suite: Romeo and Juliet have initial baryon particle masses of 3500 M_{\odot} , while Romulus and Remus and Thelma and Louise have 4000 M_{\odot} .

Both star and dark-matter particles have fixed gravitational force softenings (comoving at $z > 9$ and physical thereafter) with a Plummer equivalent of $\epsilon_{\text{star}} = 4 \text{ pc}$ and $\epsilon_{\text{dm}} = 40 \text{ pc}$. Gas cells have fully adaptive gravitational softenings that match the hydrodynamic kernel smoothing: softening lengths are $\approx 20\text{--}40$ pc at typical ISM densities ($n \sim 1 \text{ cm}^{-3}$) and $< 1 \text{ pc}$ in the densest regions.

Certain aspects of the FIRE-2 ISM model have been subject to scrutiny within the literature. In particular, Kim et al. (2023) and Wibking & Krumholz (2023) argued that the equilibrium gas cooling in FIRE-2 transitions between the warm and cold phase in the neutral ISM at unphysically low pressures and densities due to our order of too-low free-electron fraction at a given gas density. This implies that the diffuse ISM in FIRE-2 is entirely in the cold neutral phase. However, their idealized tests did not encompass how FIRE-2 simulations actually implement gas cooling; instead, they only tested one aspect of gas cooling in FIRE-2 and only in the presence of a metagalactic UV background, with no local sources. The free-electron fraction and ISM phase structure in actual FIRE-2 simulations agree much better with the observations presented in those works.² Furthermore, the specifics of the ISM cooling model are largely inconsequential for the properties that we examine here, because the thermal pressure essentially never dominates over the turbulent pressure in the cold, dense ISM resolved in these simulations. Indeed, as Hopkins et al. (2018) showed, dramatic changes to the assumed ISM cooling rates lead to only minor effects on most galaxy dynamical properties. Indeed, we tested re-simulating one of our simulations (m12i) over the last 500 Myr using the newer FIRE-3 model (Hopkins et al. 2023), which includes numerous updates to gas cooling, and this results in nearly identical velocity dispersions for cold gas and stars.

²The FIRE collaboration has made numerous improvements to the gas cooling model in recent years: the default cooling in FIRE-2 since 2020, and in FIRE-3, includes all terms claimed to be missing from earlier FIRE models. See the GIZMO user guide: http://www.tapir.caltech.edu/~phopkins/Site/GIZMO_files/gizmo_documentation.html

¹FIRE project web site: <http://fire.northwestern.edu>

For our main results (Sections 3.1, 3.3, 3.4, and 3.5), we present properties averaged across our galaxy sample, not including the lower-mass galaxies m12r and m12z or Juliet from the ELVIS on FIRE suite. However, in Sections 3.6, 3.7, and 3.8, where we explore dependence on galaxy mass, we include all 14 galaxies. m12r and m12z have stellar masses of 1.8 and $1.5 \times 10^{10} M_{\odot}$, lower than the MW's and M31's stellar masses of $\approx 5 \times 10^{10}$ and $\approx 10^{11} M_{\odot}$, respectively. Furthermore, m12r and m12z have stellar kinematics that are weakly or never dominated by disc-like motion. Juliet has a stellar mass closer to the MW's, at $3.3 \times 10^{10} M_{\odot}$, and disc-dominated kinematics at $z = 0$, but we exclude it, because a rapid gas accretion event late in its history leads to anomalous properties that skew the sample-averaged properties towards trends unique to Juliet. We will examine these cases in more detail in future work.

2.2 Measuring stellar dynamics today and at formation

Motivated primarily by measurements of the MW, we select stars that are in each galaxy's 'solar annulus': $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc today. In future work, we will examine trends with radius, to understand the drivers of post-formation dynamical heating and their application to MW observations in greater detail; briefly, we find that the dispersion today generally decreases with R today, but that neither the dispersion at formation nor the dispersion today depend on R at $R \gtrsim 1\text{--}2$ kpc. The effect of the vertical selection on our results is minimal, as we show in Appendix A.

For all of our results besides those in section 8, we show 'disc-wide' (or more aptly, 'annulus-wide') kinematics. That is, we compute median velocity and velocity dispersions using all the stars (in a given age bin) within the 'solar annulus'. However, in section 8, we show 'local' kinematics, meaning we compute the median velocity and dispersion within smaller patches (of diameter 150 or 500 pc) spanning an annulus centred at 8 kpc, and then average across all patches.

A key advantage of simulations is their ability to study properties both today and when stars formed. In both cases, we select stars on the basis of their current coordinates. However, when determining properties at formation, we only include stars that formed *in situ*, which we define following Bellardini et al. (2022) as forming within a total distance of 30 kpc comoving from the galactic centre. We do not apply this *in-situ* selection when measuring properties today, to more directly compare with observations. Across our simulations, *ex-situ* stars comprise only ~ 4 per cent of all stars currently within the disc ($R < 20$ kpc and $|Z| < 3$ kpc), but they can make up a non-trivial percentage within our oldest age bins: we show the effect of our *in-situ* selection on velocity dispersion versus age in Appendix B.

We compute stellar positions and velocities relative to the galaxy's centre. We determine the orientation of the galaxy at each snapshot (independently) by computing the rotation tensor and axis ratios of the principal axes of the stellar disc, which we define via the moment of inertia tensor of the 25 per cent youngest star particles inside a radius enclosing 90 per cent of the total stellar mass within 10 kpc physical. We compute properties at formation using the galaxy's orientation at that time, not at $z = 0$.

Our simulations store 600 snapshots from $z = 99\text{--}0$, with a time between snapshots of 20–25 Myr, and the past 20 Myr (immediately before today) includes 10 snapshots spaced by 2 Myr. We measure a star's 'formation' properties at the first snapshot after it formed, so a star particle on average experienced ≈ 10 Myr of evolution when we record its properties. This limits to some extent our ability to measure the earliest dynamical evolution, especially because, in our simulations, stars are born clustered in GMCs (for example

Benincasa et al. 2020; Guszejnov et al. 2020), where they can be subject to various gravitational interactions on short time-scales. Thus, 'formation' properties in our analysis necessarily refer to the combination of birth properties and early dynamical evolution.

When computing the velocities of a stellar population, we select star particles using 250 Myr bins of age, which corresponds to the dynamical/orbital time at the solar annulus today. When plotting velocities, we use the mass-weighted median velocity of all star particles within a given age range; we find nearly identical results using the mass-weighted mean instead. To ensure that outlier stars do not skew our results, we compute the velocity dispersion by finding half of the difference between the mass-weighted 16th- and 84th-percentile velocities of all stars in a given age range, which is equivalent to the standard deviation for normal distribution. We compute these quantities (in each age bin) individually for each galaxy and then show the mean and standard deviation across our 11 galaxies.

3 RESULTS

3.1 Stellar kinematics versus age

3.1.1 Defining the eras of disc evolution

Motivated by the kinematic trends that we explore and by previous analyses of these galaxies (for example Ma et al. 2017; Faucher-Giguère 2018; Garrison-Kimmel et al. 2018; Yu et al. 2021; Gurvich et al. 2023), we divide our galaxies' histories into three eras, which we term the 'Pre-Disc', 'Early-Disc', and 'Late-Disc' Eras. Each of our galaxies began in the Pre-Disc Era, when the gas lacked coherent rotation, such that stars formed on nearly isotropic orbits. Then, each galaxy formed a stable, coherent disc and transitioned into the Early-Disc Era. Specifically, the start of the Early-Disc Era marks the time when all subsequent generations of stars that formed with rotationally dominated kinematics, with $(v_{\phi} > \sigma_{\text{tot}})_{\text{form}}$. Table 1 lists the lookback time, t_{lb} , that each galaxy transitioned into the Early-Disc Era, $t_{\text{lb}}[(v_{\phi}/\sigma_{\text{tot}})_{\text{form}} > 1]$.

These discs were kinematically hot, with high σ , at the start of the Early-Disc Era, but they became increasingly settled, with increasingly cold and rotationally dominated kinematics, throughout it. As a result, $(v_{\phi}/\sigma_{\text{tot}})_{\text{form}}$ grew most rapidly throughout this era. Eventually, the growth of $(v_{\phi}/\sigma_{\text{tot}})_{\text{form}}$ slowed and the galaxies transitioned into the Late-Disc Era, when the kinematics at formation did not evolve much as stars formed on highly rotational orbits in a dynamically cold disc. On average, galaxies in our sample transitioned from the Pre-Disc to the Early-Disc Era ≈ 8 Gyr ago and from the Early-Disc to the Late-Disc Era ≈ 4 Gyr ago, though with significant galaxy-to-galaxy scatter in these ages.

Hopkins et al. (2023) argued that two aspects of the galaxy's gravitational potential drive the transitions between these eras. The shape of the potential, specifically the formation of a sufficiently centrally concentrated mass profile, drives the initial formation of a disc, that is, the Pre-Disc to Early-Disc transition. Separately, the absolute depth of the potential (within the radius of star formation) drives the transition from bursty to smooth star formation, which roughly coincides with the transition from the Early-Disc to Late-Disc era, although this connection may be less direct. Similarly, our Pre-Disc and Early-Disc Eras generally correspond to the 'bursty star formation' phase, in Stern et al. (2021), Gurvich et al. (2023) and Yu et al. (2021), while our Late-Disc Era corresponds to their 'steady star formation' phase.

Table 1. Stellar properties at $z = 0$ of the FIRE-2 simulations of MW/M31-mass galaxies that we analyse, in order of decreasing stellar mass.

Name	M_{star}^{90} [$10^{10} \text{ M}_{\odot}$]	v_{ϕ}^{disc} [km s^{-1}]	σ_z^{disc} [km s^{-1}]	$\sigma_{\text{tot}}^{\text{disc}}$ [km s^{-1}]	σ_z^{local} [km s^{-1}]	$\sigma_{\text{tot}}^{\text{local}}$ [km s^{-1}]	$t_{\text{lb}} \left[\left(\frac{v_{\phi}}{\sigma_{\text{tot}}} \right)_{\text{form}} > 1 \right]$ [Gyr]	$t_{\text{lb}} \left[\left(\frac{v_{\phi}}{\sigma_{\text{tot}}} \right)_{\text{now}} > 1 \right]$ [Gyr]	$t_{\text{lb}} \left[\left(\frac{v_{\phi}}{\sigma_z} \right)_{\text{now}} > \sqrt{3} \right]$ [Gyr]	$t_{\text{lb}}^{\text{burst}}$ [Gyr]
m12m	10.0	288	16.7	50.1	13.3	33.1	9.21	6.90	7.68	3.81
Romulus	8.0	265	20.4	55.9	16.6	41.8	7.42	7.16	7.16	4.90
m12b	7.3	272	17.9	56.6	13.1	37.5	7.42	7.42	7.42	6.32
m12f	6.9	256	19.9	74.9	13.5	31.6	7.42	6.39	6.39	5.01
Thelma	6.3	230	21.2	62.8	19.4	43.9	4.35	4.35	4.25	2.57
Romeo	5.9	249	14.1	40.3	12.2	31.3	11.0	10.2	10.2	6.52
m12i	5.3	238	20.1	52.4	12.8	27.5	6.65	6.39	6.39	3.14
m12c	5.1	240	23.1	62.8	19.5	38.4	6.49	6.14	6.39	3.70
m12w	4.8	163	16.6	83.9	5.5	29.8	4.09	4.6	5.89	0
Remus	4.0	216	13.9	40.3	10.7	25.0	7.93	7.93	7.93	5.88
Juliet	3.3	211	16.6	45.6	15.0	32.6	4.35	4.61	4.61	4.40
Louise	2.3	177	14.6	37.5	11.0	26.1	7.16	6.9	7.16	5.56
m12z	1.8	92.6	21.4	67.6	19.0	39.0	0.51	–	–	–
m12r	1.5	127	17.4	47.1	7.3	15.4	5.9	0.26	2.3	–
Sample mean	6.0	236	18.0	56.1	13.5	32.4	7.19	6.76	6.98	4.31

The first column lists the galaxy name: ‘m12’ indicates an isolated galaxy; otherwise the galaxy is in a LG-like pair. M_{star}^{90} is the stellar mass within R_{star}^{90} , the radius enclosing 90 per cent of the stellar mass within 20 kpc. v_{ϕ}^{disc} , σ_z^{disc} , and $\sigma_{\text{tot}}^{\text{disc}}$ are the disc-wide median azimuthal velocity, vertical velocity dispersion, and total velocity dispersion, respectively, of stars younger than 250 Myr in a cylindrical radius $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc from the midplane. σ_z^{local} and $\sigma_{\text{tot}}^{\text{local}}$ are the local vertical and total velocity dispersion of stars younger than 100 Myr measured in 500 pc patches centred at $R = 8$ kpc. The first three lookback times are the onset of disc formation, defined as the oldest age/lookback time that stars permanently passed above a disc-wide threshold: $v_{\phi}/\sigma_{\text{tot}} > 1$ at formation, $v_{\phi}/\sigma_{\text{tot}} > 1$ measured at present, and $v_{\phi}/\sigma_z > 1.8$ measured at present. The last column lists the lookback time when the galaxy transitioned from bursty to steady star formation from Yu et al. (2021). The last row shows the average quantities for our sample (excepting m12r, m12z, and Juliet, see text).

3.1.2 Median velocity versus age

We start by analysing the general dynamic evolution of our simulated galaxies. Fig. 1 (top) shows the sample-averaged azimuthal, radial, vertical, and total median velocities at formation (dashed lines) and at present (solid lines) versus stellar age. We label the Late-, Early-, and Pre-Disc Eras on the top of each column, and we mark the start of the Early and Late-Disc Eras with shaded bars. The time that one era ended and another began is not exact, in part because Fig. 1 shows sample-averaged quantities, so the shaded bars represent the approximate transition time and not an instantaneous boundary.

Fig. 1 (top left) shows the median azimuthal (rotational) velocity, v_{ϕ} , versus age. For most of the Pre-Disc Era, stars formed with $v_{\phi, \text{form}} \approx 0\text{--}20 \text{ km s}^{-1}$: in the early galaxy, stars did not coherently rotate. Eventually, $v_{\phi, \text{form}}$ began to increase as the disc began to settle, and the galaxy transitioned into the Early-Disc Era. The slight growth of $v_{\phi, \text{form}}$ towards the end of the Pre-Disc Era arises because we average over multiple galaxies, so earlier-transitioning galaxies bias the sample’s mean: within individual galaxies, the rise of $v_{\phi, \text{form}}$ is generally (but not always) sharp and coincident with the transition to the Early-Disc Era (see Fig. 3). $v_{\phi, \text{form}}$ strongly rose throughout the Early-Disc Era, but its growth weakened once the galaxy transitioned into the Late-Disc Era (≈ 4 Gyr ago), likely because the total mass within our radial range changed little during this period (Garrison-Kimmel et al. 2018). $v_{\phi, \text{now}}$ primarily follows the same trends with age as $v_{\phi, \text{form}}$. However, stars that formed in the Pre-Disc Era have significantly larger v_{ϕ} now than at formation, because they were torqued onto more coherent rotational orbits as the disc subsequently formed and settled. Likely, gas-rich mergers that seeded the disc’s orientation also torqued up these existing stars (for example Santistevan et al. 2021). Bellardini et al. (in preparation) show that this increase in v_{ϕ} corresponds to an increase in specific angular momentum for these pre-existing stars, that is, they were indeed torqued. In contrast, stars that formed in the Late-Disc Era on nearly circular orbits now have slightly smaller v_{ϕ} now than at formation because of dynamical heating/scattering.

Fig. 1 (top row, column 2 and column 3) shows the median radial velocity, v_R , and the median vertical velocity, v_z , versus stellar age. Previous work analysing FIRE simulations found that bursty star formation in the Pre-Disc galaxy drove significant gaseous outflows (for example Muratov et al. 2015; Anglés-Alcázar et al. 2017; Feldmann et al. 2017; Sparre et al. 2017; Pandya et al. 2021; Stern et al. 2021). These outflows sometimes remained star-forming and produced stars that inherited their progenitor’s outflowing dynamics (El-Badry et al. 2016; Yu et al. 2021). We find complementary results: stars that were born in the Pre-Disc Era had positive (outward) $v_{R, \text{form}}$, indicating the presence of star-forming outflows, and $v_{Z, \text{form}}$ that rapidly fluctuated around 0. Once the early disc formed, $v_{R, \text{form}}$ decreased to 0, in part because feedback-driven outflows declined and eventually ceased: throughout the rest of the Early-Disc and Late-Disc Eras, stars had $v_{R, \text{form}} \approx v_{Z, \text{form}} \approx 0$.

At present, both $v_{R, \text{now}}$ and $v_{Z, \text{now}} \approx 0$ for stars that formed in all eras, meaning that old stars’ subsequent dynamical evolution overrode any preference for outward/upward/downward motion. This difference also reflects our selection of stars that remain within the disc today. Our galaxies largely formed radially ‘inside-out’ (for example Garrison-Kimmel et al. 2018), such that old *in-situ* stars typically formed inside of our current radial range of 6–10 kpc, and as El-Badry et al. (2016) and Ma et al. (2017) discussed for the FIRE simulations, old *in-situ* stars typically reside at larger radii today, compared to when they formed, because bursts of feedback-driven outflows generated large fluctuations in the gravitational potential at early times that moved these stars outward. We note that this change in radius for old stars is distinct from ‘radial migration’, in that the vast majority of these stars have moved coherently outward, while radial migration entails a roughly equal number of inward and outward moving stars. As such, old stars in our sample are more likely than young stars to be near their apocentres, and as a result, have $v_{R, \text{now}}$ closer to 0. Bellardini et al. (in preparation) will examine this radial redistribution in detail.

Fig. 1 (top right) shows the median total velocity, v_{tot} , versus age. The growth of the galaxy’s mass drove much of the evolution

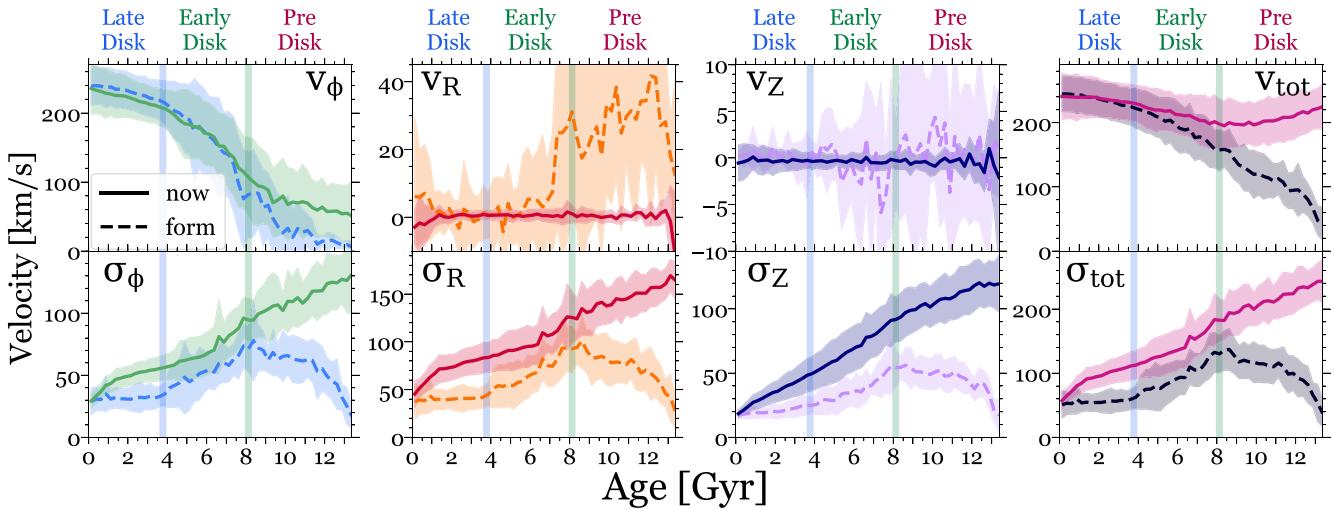


Figure 1. Median velocities and velocity dispersions for stars versus age, within cylindrical $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc at $z = 0$. Age bins have a width of 250 Myr. The lines show the average and the shaded regions show the standard deviation across 11 MW-mass galaxies, with velocities measured at $z = 0$ (solid) and at the time of formation (dashed, including only *in-situ* stars). The shaded vertical bars show when our galaxies transitioned from the Pre-Disk to the Early-Disk Era (≈ 8 Gyr ago) and the Early-Disk to the Late-Disk Era (≈ 4 Gyr ago), on average: we define these eras in Section 3.1.1. *Top row:* Median velocity versus age. For most of the Pre-Disk Era, stars formed with little net rotation, but $v_{\phi, \text{form}}$ rapidly increased in the Early-Disk Era when the disc initially formed, and it increased at a lower rate in the Late-Disk Era when the disc further settled. Stars that formed in the Pre-Disk Era have $v_{\phi, \text{now}} > v_{\phi, \text{form}}$ from dynamical torquing after formation, whereas stars that formed in the Late-Disk Era now have slightly smaller $v_{\phi, \text{now}}$ than at formation from post-formation dynamical heating. v_R and v_Z fluctuated at formation in the Pre- and Early-Disk Eras, because of frequent mergers, rapid accretion, bursty star formation, and gaseous outflows, but today they are 0 across all ages. v_{tot} , $v_{\text{tot, form}}$ steadily increased at all times, but $v_{\text{tot, now}}$ is nearly independent of age, because all stars have been dynamically heated to orbit near the virial velocity. *Bottom row:* Velocity dispersion versus age. All three components and the total show the same trends with age. σ_{form} increased throughout the Pre-Disk Era as the galaxy grew, peaking ≈ 8 Gyr ago, but decreased throughout the Early-Disk Era and remained constant throughout the Late-Disk Era. By contrast, σ_{now} monotonically increases with stellar age because of the dynamical heating of stars after their formation. Thus, the kinematics of stars today do not simply reflect their kinematics at formation.

of v_{tot} , because the underlying gravitational potential determines its virial velocity. In turn, the evolution of $v_{\text{tot, form}}$ indicates that the galaxy's mass increased over time, most rapidly in the Pre-Disk Era and only minimally in the Late-Disk Era. The evolution of $v_{\text{tot, form}}$ largely resembles that of $v_{\phi, \text{form}}$, with the exception of the Pre-Disk Era, when $v_{\text{tot, form}}$ grew steadily, while $v_{\phi, \text{form}}$ remained near 0. This is not surprising: v_{tot} depends only on the magnitude of each component, while v_{ϕ} also depends on the sign. In the Pre-Disk Era, stars rotated, just not coherently, such that the median $v_{\phi, \text{form}} \approx 0$ while $\sqrt{v_{\phi, \text{form}}^2} > 0$. However, once the galaxy transitioned into the Early-Disk Era, stars had larger and larger fractions of their (increasing) $v_{\text{tot, form}}$ in rotation. At present, v_{tot} has almost no dependence on age, with all stars near $v_{\text{tot, now}} \approx v_{\text{virial}} \approx 200\text{--}240 \text{ km s}^{-1}$, as a result of virialization and dynamical heating.

3.1.3 Velocity dispersion versus age

Fig. 1 (bottom) shows the azimuthal, radial, vertical, and total stellar velocity dispersion, σ , versus age, at both the time of the star's formation and at present. All components exhibit similar trends with age, albeit with different normalizations. Here we focus on general trends with age; we compare the three components in Section 3.3.

σ_{form} reflects the turbulence in star-forming gas and shows distinct behaviour in each kinematic era. The oldest stars formed with the lowest σ , reflecting their low $v_{\text{tot, form}}$ in the low-mass galaxy progenitor at early times. As the galaxy grew during the Pre-Disk Era, σ_{form} grew, reflecting the deepening of the gravitational potential in addition to an increase in absolute gas turbulence, from accretion,

mergers, bursty and more vigorous star formation, and outflows (for example Ma et al. 2017; El-Badry et al. 2018a; Hung et al. 2019; Flores Velázquez et al. 2021).

σ_{form} reached a maximum during the transition to the Early-Disk Era. This peak in σ_{form} reflects the competing effects of the galaxy continuing to become more massive (increasing $v_{\text{tot, form}}$) but also starting to settle into a disc, decreasing the relative contribution to σ_{form} . σ_{form} steadily decreased throughout the Early-Disk Era as the disc settled. This steady decrease in turbulence could occur only after the initial onset of the disc, when it could self-regulate to maintain a turbulent Toomre $Q \sim 1$ in star-forming gas (for example Hopkins, Quataert & Murray 2012; Faucher-Giguère, Quataert & Hopkins 2013; Wisnioski et al. 2015; Krumholz et al. 2018; Gurvich et al. 2020; Orr et al. 2020). As such, equilibrium models imply that a gas disc maintains $\sigma/v_{\phi} \approx M_{\text{gas}}(< r)/M_{\text{tot}}(< r) \approx f_{\text{gas}}$, so σ reduced as f_{gas} did.

Finally, upon the transition to the Late-Disk Era, σ_{form} reached a floor and remained nearly constant, reflecting the plateauing gas fraction and correspondingly steady and low turbulence and SFR during this era (Ma et al. 2017; Garrison-Kimmel et al. 2018; Hung et al. 2019; Yu et al. 2021; Gurvich et al. 2023).

While σ_{form} shows distinct trends in each era, σ_{now} shows markedly different behaviour, monotonically increasing with age across all eras. Thus, the additional dispersion from post-formation dynamical heating, in combination with σ_{form} , produces the observed continuous increase in dispersion with stellar age; neither process alone explains it. This also means that σ_{now} does not directly and simply reflect the formation history of the galaxy. This difference is most evident for stars that formed in the Pre-Disk Era, which have the highest values of σ_{now} today but formed with the smallest values of σ_{form} . By

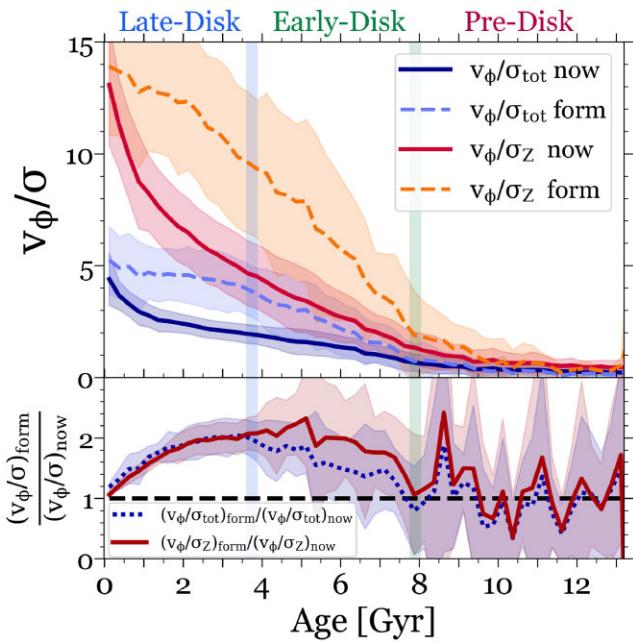


Figure 2. *Top:* The ratio of median v_ϕ to σ , that is, how rotationally supported or ‘discy’ the stars are, versus age. We average across 11 galaxies, and the shaded region shows the galaxy-to-galaxy standard deviation. Solid lines show (v_ϕ/σ) measured at present, while dashed lines show (v_ϕ/σ) measured when the star formed. In the Pre-Disk Era, stars formed with $v_\phi/\sigma \approx 0$, but once the disc began to form, $(v_\phi/\sigma_Z)_{\text{form}}$ rapidly increased until today, so vertical settling continued until today. By contrast, $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ also initially increased rapidly during the Early-Disk Era, but it remained roughly constant in the Late-Disk Era. *Bottom:* $(v_\phi/\sigma)_{\text{form}} / (v_\phi/\sigma)_{\text{now}}$ for σ_{tot} (dotted) and σ_Z (solid), which measures how ‘discy’ stars were when they formed relative to today. Both ratios have similar values and exhibit similar trends with age. Stars that formed in the Early- and Late-Disk Eras have ratios above 1 (and near 2 for most ages), because post-formation dynamical heating increased their dispersions. Pre-Disk stars have $(v_\phi/\sigma)_{\text{now}} \approx (v_\phi/\sigma)_{\text{form}} \approx 0$, with fluctuations from mergers/accretion and feedback.

contrast, stars forming in the Early-Disk Era exhibit an age evolution that follows a similar slope at formation and at present, modulo a normalization offset. Finally, stars that formed in the Late-Disk Era have σ_{now} increasing with age while σ_{form} is constant. We examine the age dependence of post-formation dynamical heating in detail in Section 3.4.

In conclusion, the age dependence of σ_{form} differed in the Pre-Disk, Early-Disk, and Late-Disk Eras, reflecting the distinct kinematics of these eras. At present, σ_{now} increases monotonically with age, with a relation to σ_{form} that differs for each era.

3.1.4 Rotational support versus age

The ratio v_ϕ/σ is a dimensionless measure of the degree of rotational support, that is, how ‘discy’ or ‘dynamically cold’ a disc is. Fig. 2 (top) shows the ratio of median v_ϕ to σ_Z and σ_{tot} , versus stellar age, averaged across our 11 galaxies. We do not show the ratios with σ_R or σ_ϕ , because they exhibit qualitatively similar trends as σ_{tot} , albeit with higher normalizations.

In the Pre-Disk era, $(v_\phi/\sigma_{\text{tot}})_{\text{form}} < 1$ (by definition). When v_ϕ, form permanently exceeded $\sigma_{\text{tot}, \text{form}}$ (≈ 8 Gyr ago on average), a galaxy transitioned into the Early-Disk Era, when successively younger populations were born with increasing $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ and $(v_\phi/\sigma_Z)_{\text{form}}$. After the galaxy transitioned into the Late-Disk Era (≈ 4 Gyr ago

on average), $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ grew only slightly, so stars during this era formed with the same degree of total rotational support. By contrast, $(v_\phi/\sigma_Z)_{\text{form}}$ continued to increase throughout the Late-Disk Era, with only slightly slower growth than in the Early-Disk Era, reflecting the more significant reduction in $\sigma_{Z, \text{form}}$ over time. Thus, we find more significant vertical ‘settling’ during the Late-Disk Era.

Relative to formation, stars have lower v_ϕ/σ as measured today. To better understand how the post-formation change in v_ϕ/σ depends on age, Fig. 2 (bottom) shows the ratio $(v_\phi/\sigma)_{\text{form}} / (v_\phi/\sigma)_{\text{now}}$ versus age. This ratio indicates how ‘discy’ a stellar population formed relative to how discy it is today, from the combined effect of post-formation dynamical heating and coherent torquing.

Younger stars, which formed in the Late-Disk Era, became increasingly less discy with age, as post-formation dynamical heating increased their dispersion. Stars older than ≈ 2 Gyr exhibit ratios near 2, meaning that they are currently half as discy as when they formed. For stars that formed in the Early-Disk Era, this trend reverses, as the post-formation change in disciness lessens with age, and the oldest stars in this era are now equally discy as when they formed. This occurs because $(v_\phi/\sigma)_{\text{form}}$ decreased with age faster than the post-formation change in v_ϕ/σ increased. Now, the ratio with σ_{tot} exhibits behavior nearly identical to that of the ratio with σ_Z , indicating that both ‘3D disciness’ and ‘vertical disciness’ evolved after formation identically relative to their values at formation.

Stars that formed in the Pre-Disk Era have values of $(v_\phi/\sigma)_{\text{form}} / (v_\phi/\sigma)_{\text{now}}$ that fluctuate around 1. Some stars have higher v_ϕ/σ now than when they formed, from post-formation torquing from events like mergers. Also, the division of intrinsically small velocity values amplifies the fluctuations.

3.2 Case studies

While Figs 1 and 2 show trends averaged across 11 galaxies, Fig. 3 shows these trends for three individual galaxies: m12f (blue), m12m (pink), and Romeo (green).

Romeo is one of our closest analogues to the MW today, residing in a LG-like pair and having a stellar mass of $5.9 \times 10^{10} M_\odot$, similar to the MW. Furthermore, recent observational analysis suggests that the MW’s disc began to form ≈ 12 Gyr ago (Belokurov & Kravtsov 2022; Conroy et al. 2022; Xiang & Rix 2022), and as we quantify further below, of all our galaxies, Romeo’s disc began to form the earliest and thus nearest in time to this estimate for the MW, likely because of its forming in LG-like environment (Garrison-Kimmel et al. 2018; Santistevan et al. 2020). Fig. 3 (left) shows that Romeo’s v_ϕ, form rapidly increased and its $(v_\phi/\sigma_Z)_{\text{form}}$ began to rise ≈ 12 Gyr ago and it transitioned into its Early-Disk Era ≈ 11 Gyr ago. Although Romeo transitioned the earliest, its evolution largely follows the same trends as most of our other galaxies.

m12f highlights the effect of major mergers on stellar kinematics. Its v_ϕ, form and $(v_\phi/\sigma_Z)_{\text{form}}$ began to increase ≈ 10 Gyr ago, during its transition to the Early-Disk Era. However, both decreased ≈ 8 Gyr ago, when a galaxy flew by and torqued its disc. After this, v_ϕ, form and $(v_\phi/\sigma_Z)_{\text{form}}$ increased again. m12f then experienced a major merger ≈ 6.9 Gyr ago. Although the dispersion was high during the merger, afterwards $\sigma_{Z, \text{form}}$ sharply decreased while $(v_\phi/\sigma)_{\text{form}}$ increased. Thus, a major merger led to a more rotationally supported disc and may have triggered its transition to the Late-Disk Era. This agrees with previous work that found that gas-rich mergers can trigger the settling of galactic discs, particularly if they occur when the potential is sufficiently deep to prevent outflows after bursts of merger-triggered star formation (for example Hopkins et al. 2009;

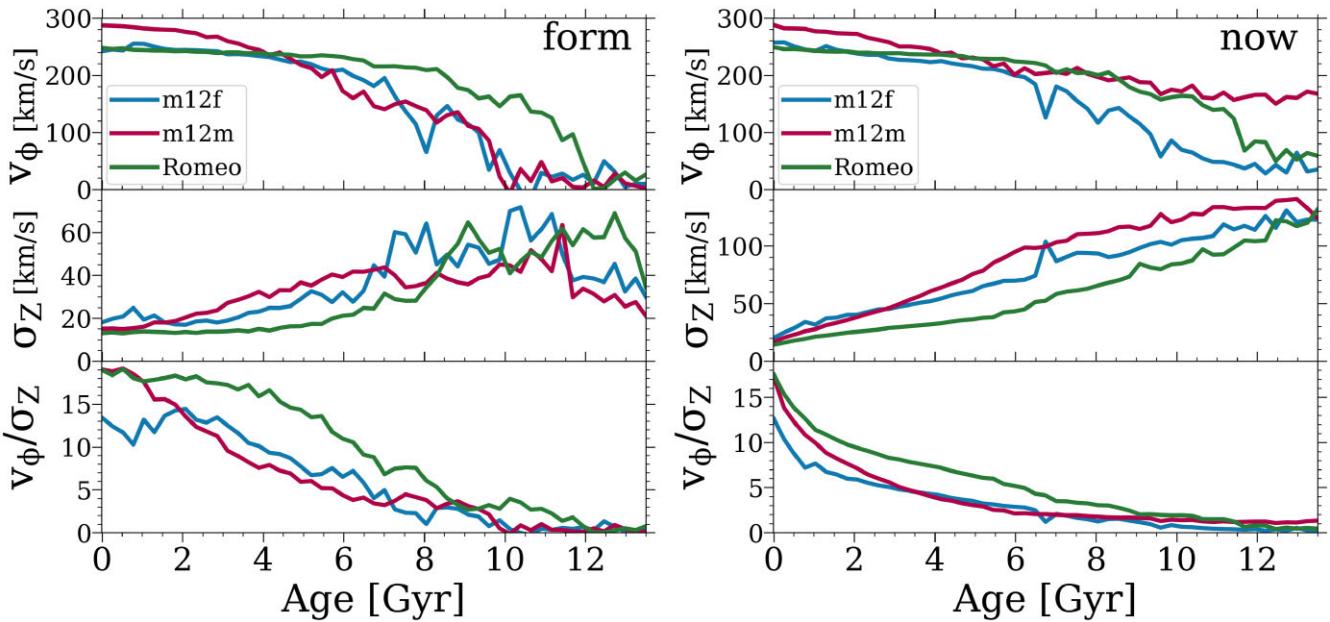


Figure 3. Median azimuthal velocity and velocity dispersion of stars versus age for three different galaxies: m12f, m12m, and Romeo. Left shows velocities at the time of formation, while right shows velocities today. *Top:* Median azimuthal velocity, v_ϕ . Romeo is the earliest settling galaxy in our sample; it exhibited rapid growth in v_ϕ at formation ≈ 2 Gyr before either m12m or m12f. However, by $z = 0$ the later-settling m12m exhibits higher v_ϕ at almost all ages, including its oldest stars. Thus, kinematics at present are not necessarily indicative of kinematics at formation. *Middle:* Vertical velocity dispersion, σ_Z . At formation, all of the galaxies show large fluctuating velocity dispersions for old stars, which peaked at different times, but this peak does not reflect itself in the current velocity dispersion versus age, at right. m12f underwent a major merger ≈ 7 Gyr ago, corresponding to a spike in the velocity dispersion in both panels. *Bottom:* Ratio of median azimuthal velocity to vertical velocity dispersion, as a measure of rotational support. Romeo is the most rotationally supported galaxy at all ages, measured both today and at formation, though m12m catches up for young stars today. Despite differing histories, all three galaxies have similar velocity dispersions for young stars today, although they are more separated in v_ϕ/σ_Z .

Moreno et al. 2021; Santistevan et al. 2021; McElroy et al. 2022; He et al. 2023).

At present, the high dispersion during the merger and the sharp decline afterward remain discernible. Although the merger more dramatically impacts the in-plane components of the dispersion (notably, σ_ϕ at present spikes from 60 km s^{-1} to 150 km s^{-1} back down to 55 km s^{-1} within 1 Gyr during the merger), for consistency, we display only the vertical component. Thus, the present-day kinematics of stars, specifically ‘spikes’ in the velocity dispersion versus age, can reveal past merger events, as we will explore in future work. After its disc began to settle, m12f underwent a smaller, prograde merger ≈ 1.5 Gyr ago that triggered a significant burst of star formation (Yu et al. 2021), increased σ_{form} , and decreased $(v_\phi/\sigma)_{\text{form}}$. Thus, certain mergers occurring in the Pre- or Early-Disc Eras increased the galaxy’s rotational support, while mergers occurring in the Late-Disc Era decreased the disc’s rotational support.

We next highlight m12m, whose evolution deviates from our other galaxies. m12m has the largest stellar mass at $z = 0$ in our sample, $\approx 10^{11} M_\odot$, more similar to M31. Furthermore, m12m underwent numerous early mergers, including two particularly massive mergers ≈ 10 and 9 Gyr ago, but it had a quiescent merger history over the past 8 Gyr.

m12m has our second earliest transition into the Early-Disc Era: its $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ permanently surpassed unity ≈ 9 Gyr ago. Despite this, m12m has yet at $z = 0$ to transition fully into its Late-Disc Era: its $v_{\phi,\text{form}}$ continues to increase (albeit only slightly), all but its vertical component of σ_{form} have yet to become constant, and $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ continues to strongly increase. This may be because m12m began to develop a bar ≈ 4.8 Gyr ago, leading to a bar-buckling event that formed a significant X-shaped bulge composed primarily of younger

stars (age < 4.8 Gyr; Debattista et al. 2019). m12m’s kinematics are also unique at present; most notably, its $v_{\phi,\text{now}}$ exhibits no significant decrease for old stars; instead, all ages have $v_{\phi,\text{now}} \gtrsim 160 \text{ km s}^{-1}$, which is almost $3 \times$ faster than our sample’s average, likely as a result of its multiple major mergers. Of our three case studies, m12m generally had the lowest $\sigma_{Z,\text{form}}$ for stars older than ≈ 8 Gyr, but at present m12m has the highest $\sigma_{Z,\text{now}}$ at almost all ages, because its stars have been more dynamically heated.

Comparing these case studies provides useful context and caution for interpreting population-wide trends. First, the dispersion of young stars does not necessarily reflect the dispersion of older stars: across the three galaxies, the difference in $\sigma_{Z,\text{now}}$ is only a few km s^{-1} for the youngest stars (age < 250 Myr), but jumps to 14 km s^{-1} at 1 Gyr old and 57 km s^{-1} at 6 Gyr old. In fact, the rate at which σ_{now} increases with age varies between galaxies. On a related note, how σ_{form} compares between galaxies may be different from how σ_{now} does. That is, stars of a certain age may now have larger/smaller dispersions than those in another galaxy despite forming with smaller/larger dispersions. For example, Romeo generally had the highest values of $\sigma_{Z,\text{form}}$ for stars older than 6 Gyr, while m12m had the lowest. However, this ordering flips at present: today, σ_Z in m12m is higher than in Romeo. Thus, the dispersion with which stars formed does not set their present dispersion, because the amount and age dependence of post-formation heating can differ significantly between galaxies.

On the contrary, v_ϕ/σ_Z provides a similar comparison between galaxies at both formation and present-day. $(v_\phi/\sigma_Z)_{\text{now}}$ maintains its relative ordering between galaxies from formation and even largely preserves the trends with age exhibited by each individual galaxy: Romeo still has the highest values for stars younger than ≈ 10 Gyr, m12f and m12m still have similar values for intermediate-aged stars,

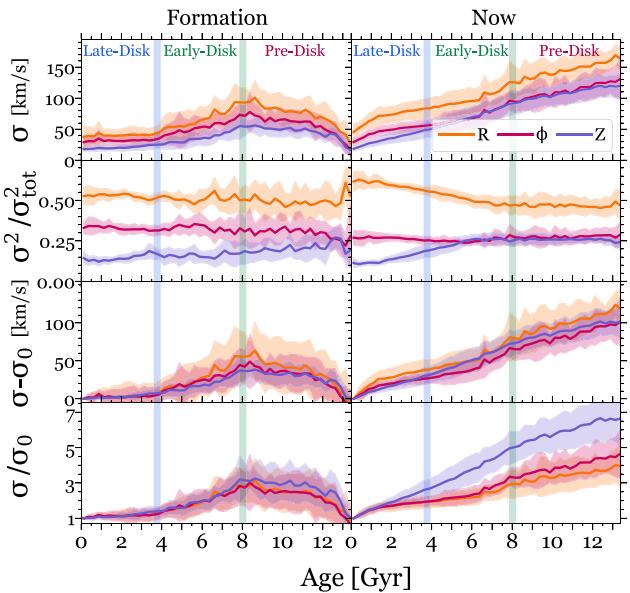


Figure 4. Comparing the evolution of the three components of velocity: lines show the radial, azimuthal, and vertical velocity dispersion at formation (left) and today (right) versus age. We average across 11 galaxies, and the shaded regions show the galaxy-to-galaxy standard deviation. *Top row:* Velocity dispersion, σ . At formation, all components exhibit constant ordering, with $\sigma_R > \sigma_\phi > \sigma_Z$, and similar dependence on age. Today, σ_R dominates at all ages, while σ_ϕ and σ_Z are near identical for older stars. However, σ_Z falls below σ_ϕ for stars that formed in the Late-Disk Era. *Second row:* Ratio of each component's σ^2 relative to the total at that age, as a metric of the relative kinetic energy in each component. At both formation and today, σ_R^2 dominates the kinetic energy. At formation, kinetic energy partition is nearly independent of age, while today, it is nearly independent for ages > 6 Gyr, but because of post-formation heating, the importance of σ_Z^2 decreases while σ_R^2 increases for younger and older stars. *Third row:* Absolute change in the velocity dispersion relative to σ_0 , the dispersion of the youngest stars (age < 250 Myr at $z = 0$) for each component. At formation and today, all components are nearly self-similar, indicating that each component changes with age by about the same absolute amount (although at present, the radial component is slightly higher, because of its stronger post-formation heating). *Bottom row:* Velocity dispersion normalized by its value for the youngest stars, which quantifies the fractional change in σ over time. At formation, all components display similar values and age dependence, such that intermediate-age stars formed with 3 times higher σ than the youngest or oldest stars did. As measured today, σ_Z exhibits the strongest fractional dependence on age, despite it showing the weakest absolute dependence on age in the top panels, as a result of post-formation heating.

and m12f still has the lowest values for young stars. Although the differences between each galaxy's v_ϕ/σ_Z are less pronounced today than at formation, present day observations of $(v_\phi/\sigma)_{\text{now}}$ (versus age) give a better representation of a galaxy's formation history than σ_{now} and v_ϕ, now alone.

3.3 Comparing velocity components

All three spatial components of velocity dispersion share similar trends with age, but we now compare them in detail. Fig. 4 shows the radial, azimuthal, and vertical velocity dispersions, averaged across the 11 galaxies, at formation (left) and at present (right).

Fig. 4 (top row) compares the dispersions. The relative order of the components at formation is constant with age: $\sigma_{R, \text{form}} > \sigma_{\phi, \text{form}} > \sigma_{Z, \text{form}}$. At present, $\sigma_{R, \text{now}}$ dominates at all ages. However,

$\sigma_{Z, \text{now}}$ now matches $\sigma_{\phi, \text{now}}$ for stars that formed in the Early- and Pre-Disk Eras. Thus, the similarity between $\sigma_{Z, \text{now}}$ and $\sigma_{\phi, \text{now}}$ for intermediate-aged and old stars arises from vertical post-formation dynamical heating, not from formation. Measurements of σ_{now} in the Solar neighbourhood find $\sigma_R > \sigma_\phi > \sigma_Z$ (for example Wielen 1977; Dehnen & Binney 1998; Nordström et al. 2004). Although our results agree for stars that formed in the Late-Disk Era, we find that $\sigma_{\phi, \text{now}} \approx \sigma_{Z, \text{now}}$ for stars that formed in the Early and Pre-Disk Eras, potentially indicating that post-formation heating of σ_Z is greater for older stars in our galaxies than in the MW.

Fig. 4 (second row) shows the ratio of each component's σ^2 to the total σ^2 at that age, which compares each component's contribution to the overall kinetic energy in dispersion. Because $\sigma_{\text{total}}^2 \equiv \sigma_R^2 + \sigma_\phi^2 + \sigma_Z^2$, the sum of this ratio in all three components is 1. Remarkably, the ordering and relative value of each component at the time of formation are independent of age. That is, σ_{form} was distributed across its three components the same way throughout the galaxy's entire formation history, with $\sigma_{R, \text{form}}^2$, $\sigma_{\phi, \text{form}}^2$, and $\sigma_{Z, \text{form}}^2$ accounting for ≈ 51 per cent, 32 per cent, and 17 per cent of $\sigma_{\text{tot,form}}^2$, on average. This roughly constant partitioning across cosmic time is striking, given that it persists across the eras of disc formation, including the transition from dispersion-dominated to rotation-dominated orbits, from bursty to steady star formation, and the abatement of galactic outflows. Given this consistent partitioning at formation during both early and late times, we posit that the drivers of (star-forming) gas dispersion common in the early galaxy, such as star-burst driven outflows, and those common in the later galaxy, such as spiral arms, have similar impacts on how the energy in dispersion at formation is partitioned between components. However, we relegate testing this hypothesis to future work.

Currently, $\sigma_{R, \text{now}}$ dominates in all ages, as Fig. 4 (top) shows. However, unlike at formation, the relative contributions of $\sigma_{R, \text{now}}^2$ and $\sigma_{Z, \text{now}}^2$ are constant with age only for stars older than ≈ 6 Gyr, which formed in the Pre-Disk Era and the beginning of the Early-Disk Era. $\sigma_{R, \text{now}}^2$ still contributes ≈ 50 per cent of the random kinetic energy of these old stars, while $\sigma_{\phi, \text{now}}^2$ and $\sigma_{Z, \text{now}}^2$ both contribute ≈ 25 per cent.

For stars younger than 6 Gyr, the relative contribution of $\sigma_{R, \text{now}}^2$ increases for younger stars, even surpassing its contribution at formation, because young stars have been heated most strongly along their radial component, as we show below. In contrast, the relative contribution of $\sigma_{Z, \text{now}}^2$ decreases for younger stars, more directly reflecting the low $\sigma_{Z, \text{form}}$ of these young stars. Interestingly, stars of all ages have approximately 25 per cent of their random kinetic energy from $\sigma_{\phi, \text{now}}^2$, slightly less than when they formed.

We define σ_0 as the dispersion of our youngest stars (age < 250 Myr at $z = 0$). To examine the degree of self-similarity in the evolution of the different velocity components, Fig. 4 (third row) shows $\sigma - \sigma_0$, the absolute change in the velocity dispersion with age. All velocity components have more similar evolution in $\sigma - \sigma_0$ than in σ . This is most self-similar at formation, but it also holds to a lesser degree at present. In particular, $\sigma_{R, \text{form}}$ exhibits a greater absolute change only for intermediate-aged stars (≈ 6 –10 Gyr), whereas $\sigma_{R, \text{now}}$ has a greater absolute change at all ages, because $\sigma_{R, \text{now}}$ had the largest amount of dispersion added by post-formation dynamical heating for all but intermediate-aged stars, as we discuss below.

Fig. 4 (bottom row) shows σ/σ_0 , the fractional change in the velocity dispersion with age, relative to today. At formation, this exhibits significant self-similarity, with all components showing near identical evolution for stars that formed in the Early- and Late-Disk Eras. In particular, at the onset of the Early-Disk Era, stars

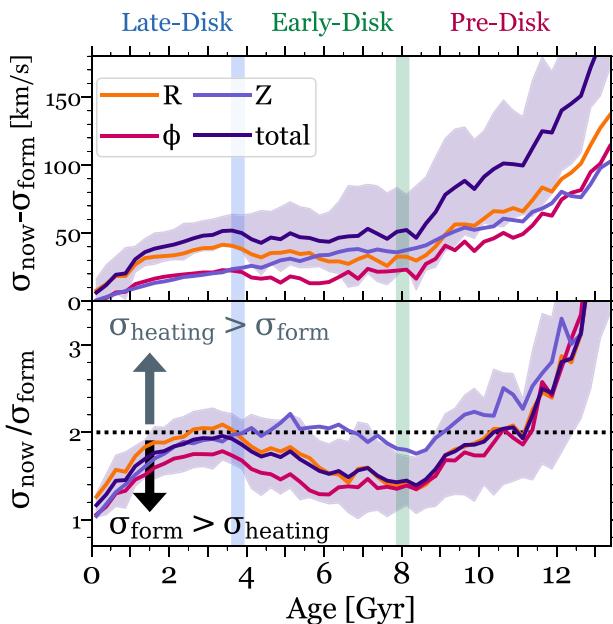


Figure 5. Comparison between the velocity dispersion today versus at formation, which quantifies the amount of post-formation dynamical heating, versus age, for different components of velocity. We average across 11 galaxies, and the shaded region shows the galaxy-to-galaxy standard deviation. *Top:* Absolute change in dispersion from formation to today. The amount of heating increases with age. However, for all components except σ_Z , the heating is constant for ages $\approx 3\text{--}8$ Gyr, and thus, intermediate-age stars were heated by a similar amount. *Bottom:* Ratio of dispersion today relative to that at formation, which shows an ‘S’-shaped dependence on age. The dotted horizontal line at 2 divides stars whose current dispersion is primarily from formation (below) versus post-formation (above). The combination of small dispersions at birth and rapid post-formation heating leads to a steadily increasing ratio with age up to ≈ 2 Gyr. Intermediate-age stars have a present dispersion mostly inherited at formation, because they formed when the intrinsic dispersion within the galaxy was highest. The oldest stars have the largest contribution from post-formation heating. *For stars younger than 10 Gyr, most of their dispersion was in place at formation, though σ_Z shows nearly equal contributions from formation and subsequent heating at ages 4–8 Gyr.*

formed with $\approx 3 \times$ the dispersion as young stars today for all velocity components. However, the striking self-similarity at formation does not hold at present. Although all components evolve relatively similarly for ages 0–3 Gyr, the vertical component increasingly dominates for older stars, with the oldest stars having vertical dispersions $\approx 7 \times \sigma_{Z,0}$. Because σ_Z had similar fractional evolution at formation, its distinctive behaviour at present must arise from a greater fractional amount of post-formation heating for stars that formed in the Early- and Pre-Disk Eras, which we explore next.

3.4 Disc settling versus post-formation heating

As we have shown, both disc settling and dynamical heating shape the relation between σ_{now} and stellar age. We now quantify the relative importance of these two processes versus age.

Fig. 5 shows the absolute and fractional impact of post-formation dynamical heating on the present-day radial, azimuthal, vertical, and total velocity dispersion of stars versus age. Most notably, both absolute and fractional post-formation dynamical heating follow

distinct trends with age for stars that formed in each of our three kinematic eras.

Fig. 5 (top) shows the difference between dispersion at present and formation, $\sigma_{\text{now}} - \sigma_{\text{form}}$, versus age, which directly conveys the amount of dispersion added post-formation. We focus primarily on the total dispersion, noting that the radial and azimuthal components show similar behaviour. However, we discuss the vertical dispersion separately below, because it shows qualitatively different behaviour.

For stars that formed in the Late-Disk Era, the amount of dispersion added post-formation increases with age, but at a declining rate, consistent with the idea that post-formation dynamical heating from scattering processes causes the observed increase in σ_{now} with age (for example Spitzer & Schwarzschild 1951). However, stars that formed in the Early-Disk Era have been dynamically heated by the same total amount, $\approx 50 \text{ km s}^{-1}$, regardless of age. Considering the monotonic increase of σ_{tot} in Fig. 1, one may not intuitively expect this plateau. For example, stars that formed at the start of the Early-Disk Era (≈ 8 Gyr ago) currently have $\sigma_{\text{tot, now}}$ that is almost 80 km s^{-1} higher than stars born at the end of the Early-Disk Era (≈ 4 Gyr ago), but have 5 km s^{-1} less σ_{tot} added post-formation. Thus, the monotonic increase of $\sigma_{\text{tot, now}}$ with age for stars that formed in the Early-Disk Era ($\approx 4\text{--}8$ Gyr ago) arose not from the dependence of post-formation heating on age, but rather, from the amount of disc settling in σ_{form} versus age.

Once a disc started to form, scattering with structures within the now well-defined disc plane likely drove most post-formation heating, though interactions with satellite galaxies likely also contributed. As stars were dynamically heated, they experienced larger radial and vertical oscillations. Larger vertical oscillations decrease the amount of time that stars spend near the midplane and increase their vertical velocities when passing through the midplane. In turn, the ability of stars to be heated along the in-plane direction decreases as vertical oscillations increase. Similarly, stars with large radial eccentricities encounter spiral structure at different orbital phases; the effects of these varied encounters tend to average out, such that spiral-driven heating becomes progressively weaker for increasingly eccentric orbits (Binney & Tremaine 2008). Thus, stars that formed with low dispersions in the Late-Disk Era were initially efficiently heated. However, this efficiency decreased accordingly as their dispersions increased over their lifetimes, explaining the near-exponential behaviour for $\sigma_{\text{tot, now}} - \sigma_{\text{tot, form}}$ shown in the upper row of Fig. 5. Similarly, because stars that formed in the Early-Disk and Pre-Disk Eras had large dispersions at birth, they likely underwent little post-formation heating over the course of the Late-Disk and Early-Disk Era (though old stars likely experienced significant heating within the Pre-Disk Era).

On the other hand, post-formation vertical heating shows different trends than in-plane heating across the Late and Early-Disk Eras. The amount of σ_Z added post-formation never plateaus but instead monotonically increases with age by $\approx 7 \text{ km s}^{-1}$ per Gyr. Therefore, post-formation vertical heating is likely a steady, continuous process, implying a different dynamical origin than in-plane post-formation heating. In future work, we will pursue a detailed analysis of the origin of in-plane and vertical heating.

For stars that formed in the Pre-Disk Era, the amount of post-formation heating increases steadily with age for all components, including vertical. Older stars have more dispersion added post-formation, not only because they have had more time to be heated, but also because the early galaxy experienced the most frequent mergers, starbursts, and potential fluctuations that increased the dispersion of existing stars (for example El-Badry et al. 2018b). Likely, this phase accounted for most of the heating of stars that formed in the Pre-Disk

Era. In later eras, more localized, temporally stable processes (that is, scattering by non-axisymmetric structures within the disc) primarily heated the stars.

Fig. 5 (bottom) shows the ratio of dispersion at present to dispersion at formation, $\sigma_{\text{now}}/\sigma_{\text{form}}$, versus age, which quantifies the fraction of the present dispersion that existed at formation versus arose from post-formation heating. Values of $\sigma_{\text{now}}/\sigma_{\text{form}} > 2$ indicate that post-formation heating contributes more to the current dispersion than the dispersion from formation does, while values < 2 indicate that most of the dispersion was in place at the time of formation.

The evolution of the total dispersion largely matches that of the radial and azimuthal components. $\sigma_{\text{now}}/\sigma_{\text{form}}$ has an ‘S’-shaped dependence on age, somewhat different from that in the top panel. Post-formation heating increases rapidly with age for stars that formed in the Late-Disc Era. The oldest stars of this era have current (radial, vertical, and total) dispersions with approximately equal contributions from heating and formation. Furthermore, stars that formed ≈ 4 Gyr ago, during the transition to the Late-Disc Era, experienced the largest fractional increase from post-formation heating of all stars younger than 10 Gyr (for all but the vertical dispersion). These stars formed in a dynamically settled thin disc with low dispersion and had more time to experience dynamical heating than stars that formed afterward.

However, this dependence on age strikingly reverses in the Early-Disc Era, when $\sigma_{\text{now}}/\sigma_{\text{form}}$ decreases with age. As Fig. 5 (top row) showed, the absolute amount of dispersion added post-formation is constant throughout this era. However, as Fig. 1 showed, σ_{form} increases with age throughout this era. Thus, the fractional contribution of σ_{form} at birth became increasingly important with age in this era. Besides stars younger than ≈ 1 Gyr, the oldest stars from the Early-Disc Era have the lowest dependence on post-formation heating, with only 40 per cent of their current σ_{tot} from post-formation heating, corresponding to their maximum σ_{form} and comparatively low amount of post-formation heating. In essence, stars were not able to be dynamically heated efficiently during this Early-Disc Era, because they already were born with high dispersions.

For stars that formed in the Pre-Disc Era, $\sigma_{\text{now}}/\sigma_{\text{form}}$ increases rapidly with age, because σ_{form} decreases with age (given the galaxy’s lower mass), while the amount of post-formation heating (and thus σ_{now}) increases with age (from mergers, starbursts, and strong outflows). Thus, the current total dispersion of stars is primarily due to post-formation heating for stars older than ≈ 11 Gyr, and is primarily inherited at formation for stars younger than 11 Gyr.

The evolution of the vertical component largely resembles that of the other components for stars that formed in the Late-Disc Era. However, while $\sigma_{\text{now}}/\sigma_{\text{form}}$ of the in-plane components noticeably decrease with age for stars that formed in the Early-Disc Era, the vertical component remains relatively constant near 2. In fact, $\sigma_{Z,\text{now}}/\sigma_{Z,\text{form}} \approx 2$ for stars with ages between 2–10 Gyr, meaning that the current vertical dispersions of stars with ages 2–10 Gyr have equal contributions from post-formation heating and formation. Thus, for most stars, the fractional impacts of vertical post-formation heating and vertical disc settling are about the same. For stars younger than 2 Gyr, $\sigma_{Z,\text{form}}$ is more important, while for stars older than 10 Gyr, vertical post-formation heating is.

In summary, the monotonic rise of the current velocity dispersion with stellar age arises from the combined effects of both cosmological disc settling and the dynamical processes that heat stars after they form. In understanding the origin of the current shape of the relation between the total velocity dispersion and age, we conclude that: (1) post-formation heating dominates the dispersion for stars that formed in the Late-Disc Era (age $\lesssim 3.5$ Gyr), when the dispersion at

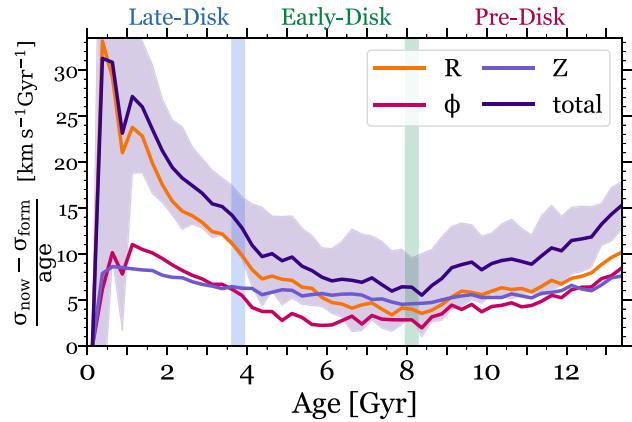


Figure 6. The effective rate of dynamical heating since formation, $(\sigma_{\text{now}} - \sigma_{\text{form}})/\text{age}$, versus stellar age, for each component of velocity. We average across 11 galaxies, and the shaded region shows the galaxy-to-galaxy standard deviation. The youngest stars have the highest effective rate of heating, implying that the rate of dynamical heating was most rapid immediately after stars formed and became more gradual over time. Increases in σ_R strongly dominate for stars that formed in the Late-Disc Era, but all components exhibit broadly similar effective rates for stars that formed in the Early- and Pre-Disc Eras.

formation was flat with age, (2) disc settling dominates for stars that formed in the Early-Disc Era (age ≈ 4 –8 Gyr), when post-formation heating was flat with age, and (3) post-formation heating strongly dominates for stars that formed in the Pre-Disc Era (age $\gtrsim 10$ Gyr). Overall, post-formation dynamical heating only dominates the current dispersion for stars that formed ≈ 3 Gyr before the disc started to form, or, equivalently, the stellar velocity dispersion at present is primarily inherited at formation and not from post-formation dynamical heating for stars younger than ≈ 10 Gyr.

Our results agree with Yu et al. (2022), who also analysed these FIRE-2 galaxies and found that most stars currently on ‘thin-disc-like’, ‘thick-disc-like’, or ‘spheroid-like’ orbits were born that way, that is, the type of orbit a star is currently on primarily reflects its formation orbit, not its post-formation evolution.

3.5 Effective rate of post-formation heating

Stars that formed in the Pre-Disc Era have the most dispersion added post-formation, as Fig. 5 (top) shows. One might assume (naively) that the amount of heating simply scales with age because older stars have had more time to be heated. However, Fig. 5 (top) showed that the amount of post-formation heating does not necessarily increase with age in all eras. For example, stars 8 Gyr old have experienced the same absolute amount of post-formation heating as stars 4 Gyr old. To provide more insight, we examine the ‘effective’ rate that stars today experienced post-formation heating, on average, via $(\sigma_{\text{now}} - \sigma_{\text{form}})/\text{age}$. This quantity is not the actual instantaneous rate of increase in dispersion. Rather, it represents the overall, time-averaged rate at which the dispersion increased since formation. Thus, one can think of this effective rate as the amount of post-formation heating normalized to the stellar populations’ age.

Fig. 6 shows this effective rate versus stellar age, for radial, vertical, azimuthal, and total dispersions. First, we focus on the total dispersion: $(\sigma_{\text{tot,now}} - \sigma_{\text{tot,form}})/\text{age}$. The youngest stars experienced the largest effective rate of increase, which rapidly decreases with age for stars that formed during the Late-Disc Era. For stars that formed in the Early-Disc Era, the effective rate was low, reaching

a minimum near the transition to the Pre-Disc Era. This strong age dependence implies that, for stars that formed in a disc, most dynamical heating occurred shortly after birth, and their heating rate slowed as they became progressively dynamically hotter (excepting the vertical component). By contrast, stars that formed in the Pre-Disc Era had effective rates of heating that increase with age, likely because more dramatic, galaxy-wide, non-equilibrium perturbations like mergers and accretion, starbursts, and galactic winds were more prominent. The number of these perturbation events that stars in the Pre-Disc era experienced increases with stellar age, such that the oldest stars have experienced the most cumulative heating.

Furthermore, for stars that formed in the Late-Disc Era, the increase in σ_R strongly dominates over the other components, so radial heating dominates the total heating. For example, the youngest stars have effective rates near $10 \text{ km s}^{-1} \text{ Gyr}^{-1}$ for σ_ϕ and σ_Z , but above $30 \text{ km s}^{-1} \text{ Gyr}^{-1}$ for σ_R , which may indicate that heating by spiral arms and/or bars dominates most early heating for stars forming in the Late-Disc Era. However, the effective radial rate decreases with age, while the other components remain more constant. Thus, stars that formed in the Early- and Pre-Disc Eras have nearly identical rates for all components of σ , indicating that heating was more isotropic in the early galaxy. Thus, post-formation heating operated differently over cosmic time. Any model of post-formation heating thus should treat early cosmic times (prior to the initial formation of the disc) and later cosmic times (after the formation of the disc) separately.

3.6 Disc kinematics versus galaxy mass

So far, we showed trends averaged across our 11 galaxies that are similar to the MW in stellar mass. Now, we analyse all 14 galaxies in our sample individually, including the lower-mass m12r, m12z, and Juliet to expand the mass range (see Table 1 for details).

Fig. 7 shows how stellar kinematics depend on the current stellar mass (within R_{90}) of a galaxy. We focus on σ and v_ϕ/σ_Z , as before. Although we do not show it, v_ϕ increases strongly with stellar mass, which follows from the deeper gravitational potential of more massive galaxies (see also Hopkins et al. 2023). We show the isolated galaxies from the Latte suite as circles and the galaxies in LG-like pairs from the ELVIS on FIRE suite as stars.

Fig. 7 (top panels) shows σ_Z (top row) and v_ϕ/σ_Z (second row), for stars of all ages. If we exclude the two lowest-mass galaxies, m12r and m12z, whose discs have yet to or have just begun to settle, σ_Z increases with stellar mass. This trend is remarkably strong among LG-like galaxies, with the exception of Thelma, which underwent the most major mergers of our sample and has a higher σ_Z than galaxies of similar or higher mass. However, unlike σ_Z , v_ϕ/σ_Z of all stars does not depend on mass at $> 2 \times 10^{10} \text{ M}_\odot$. Interestingly, galaxies in LG-like pairs have lower dispersions and more rotational support than isolated galaxies on average. These colder kinematics likely relate to the fact that on average our LG-like galaxies formed discs earlier than our isolated galaxies did, as we show below.

Fig. 7 (bottom panels) shows the same, but for stars younger than 250 Myr, which more closely reflects the current dynamical state of the ISM and is less affected by a galaxy's integrated merger and heating histories. For young stars, σ_Z does not depend on stellar mass, while v_ϕ/σ_Z increases modestly with mass, implying that more massive galaxies have more vertically settled gas discs at present.

Although we do not show it, we performed an identical analysis for σ_R , σ_ϕ , and σ_{tot} . The kinematics of all stars show nearly identical trends with stellar mass for all components; however, the in-plane and total kinematics of young stars depend more weakly on stellar mass

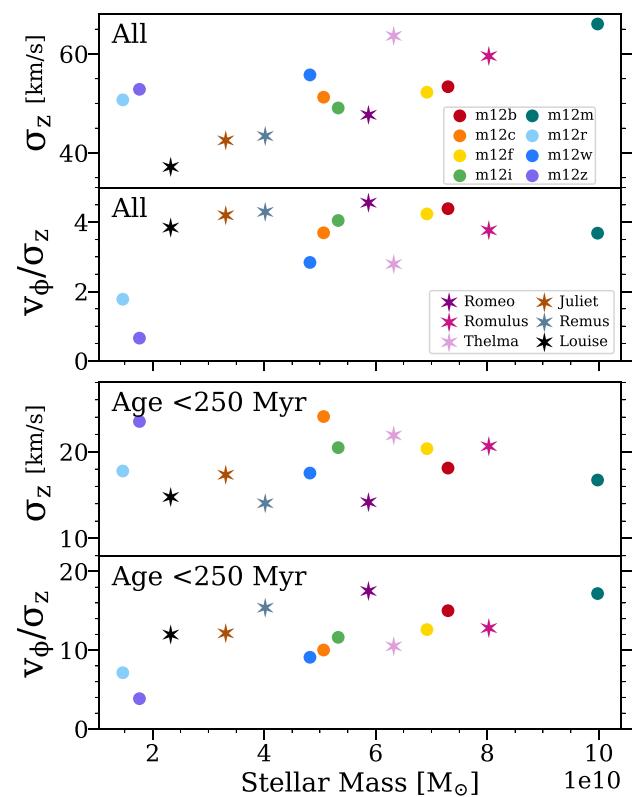


Figure 7. Vertical velocity dispersion, σ_Z , and v_ϕ/σ_Z today versus stellar mass for 14 galaxies. We show the isolated (m12) galaxies via circles and LG-like paired galaxies via stars. *Top panels:* Stars of all ages in the galaxy. The dispersion for our disc-dominated galaxies increases with mass. All galaxies have approximately constant v_ϕ/σ_Z regardless of mass, excepting m12r and m12z. *Bottom panels:* Stars with ages < 250 Myr. Here, σ_Z does not correlate with mass, although v_ϕ/σ_Z is slightly higher at higher mass. Overall, galaxies in LG-like pairs tend to have lower dispersions and more rotational support at a given mass; however, this difference is more evident for all stars than for young stars.

than the vertical kinematics (weakly) do. Ultimately, we conclude that the stellar mass of a galaxy has little effect on the current dispersion of young stars in the solar annulus and at most a modest effect on their degree of rotational support in the mass range we explore.

3.7 Comparisons with observations

Thus far, we have focused on our galaxies' disc-wide kinematics, that is, the median velocities and dispersions for all stars (in a given age bin) within our simulated 'solar annulus' of $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc at $z = 0$. However, observational measurements of nearby external galaxies typically measure kinematics in smaller-scale 'patches' and then average over these spatial patches to discuss trends with age within the galaxy. Similarly, most published dispersion values for the MW are for stars within one 'patch' – the solar neighbourhood. We now examine how the type of measurement – 'disc-wide' or 'local' – impacts our measured velocity dispersions.

We measure a local velocity dispersion within annular bins of side length 150 pc, motivated by the typical ALMA-resolution scale for nearby galaxies, and then compute the average across the annulus at a given radius. For more context of our results on velocity dispersions of young stars, we also determine the local dispersions of cold and

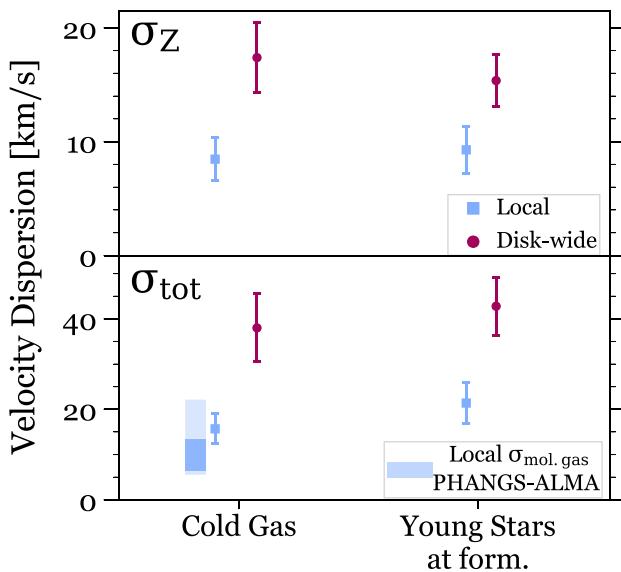


Figure 8. Comparison of velocity dispersion, for a disc-wide versus local measurement, for cold dense gas and young stars: Specifically, σ_z (top) and σ_{tot} (bottom) of cold ($T < 300$ K) and dense ($n > 10 \text{ cm}^{-3}$) gas at present (left) and young (age < 250 Myr) stars at formation (right). The points show the mean across all 14 galaxies, while the error bars show the standard deviation. Circles show the dispersion of gas or stars across $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc, identical to our previous analyses. The squares show the local velocity dispersion within 150 pc patches, with cold gas patches centred at ≈ 4 kpc and young star patches centred at 8 kpc (see the text for further explanation). The shaded region shows the total dispersion of molecular gas in 19 PHANGS-ALMA galaxies: the light shaded region shows the sample’s full range while the darker shaded region shows the sample’s 16th–84th percentile range.

dense gas, $T < 300$ K and $n > 10 \text{ cm}^{-3}$, as a proxy for molecular gas, following Orr et al. (2018).

Fig. 8 shows the average vertical (top) and total (bottom) dispersion of cold dense gas at present (left) and young (age < 250 Myr) stars at formation (right), averaged over our 14 galaxies, for both local dispersions (light blue) and disc-wide dispersions (magenta). Our disc-wide dispersions use the same geometric selection as all of our previous analyses, while our local dispersions are the annulus-averaged dispersion in patches of side length 150 pc centred at $R = 8$ kpc. However, for the local dispersion of cold gas, we instead average between the dispersion within annuli centred at $R \approx 3.5$ and $R \approx 4.9$ kpc (corresponding to 1/3 and 2/3 of our observational radial range), to match the radial range of the PHANGS observations we compare against (see below). The cold gas velocity dispersions measured at $R = 8$ kpc are ≈ 72 per cent of those measured at $R \approx 4.2$ kpc in our simulations.

As the results for σ_z (top) show, cold gas in FIRE-2 galaxies have 1D local velocity dispersions of $\sigma_{\text{1D}} \approx 8.4 \text{ km s}^{-1}$ at $R \approx 4.2$ kpc (as shown, and $\approx 6 \text{ km s}^{-1}$ at $R \approx 8$ kpc not shown) – much lower than the typical disc-wide dispersions in our previous results. We also compare against ALMA-measured dispersions of molecular gas CO(2–1) on ≈ 100 pc scales in PHANGS galaxies from Sun et al. (2020), where we restrict the observed sample to the 19 galaxies that also have resolved stellar kinematics in Pessa et al. (2023). We include only gas observations within $R = 4.2 \pm 2$ kpc, and we measure the local dispersion of cold gas in our FIRE-2 galaxies at concordant radii (≈ 3.5 and 4.9 kpc). We measure the mean σ_{LOS} for each PHANGS galaxy and assume isotropy to convert to a 3D

σ_{tot} . Fig. 8 (bottom left) shows the PHANGS-ALMA sample’s full range of mean σ_{tot} as a shaded light-blue region and 1σ scatter in darker blue. Although our FIRE-2 galaxies have local velocity dispersions in cold gas somewhat higher than typical in PHANGS, all of our FIRE-2 galaxies fall within the observed range. Although a rigorous comparison would require synthetic CO observations, our comparison here indicates that velocity dispersions in cold gas are broadly consistent with those of molecular gas in nearby galaxies.

Most importantly, Fig. 8 shows that the dispersion of cold gas in our simulations is significantly larger (≈ 3 times) if we measure it disc-wide rather than locally. Similarly, stars at formation have disc-wide σ_z and σ_{tot} that are ≈ 2 times larger than their local values at $R = 8$ kpc. These values broadly agree with those from Orr et al. (2020; $\approx 15\text{--}30 \text{ km s}^{-1}$), who examined the spatially resolved dispersion of neutral gas in seven of our isolated (m12) simulations. This motivates the need to use local velocity dispersions to compare against spatially resolved observations of nearby galaxies.

Finally, the local σ_z and σ_{tot} for young stars at formation are 1.5 and 1.9 times larger, respectively, than for cold gas (at $R \approx 8$ kpc). On the other hand, young stars at formation have a disc-wide σ_{tot} that is about the same as cold dense gas, while the disc-wide σ_z is lower for stars at formation than for cold dense gas. The higher local dispersion for stars than for cold gas is expected, given that we typically measure them ≈ 12 Myr after formation. However, the weaker (and even negative) change from cold gas to stars at formation for the disc-wide dispersion suggests that rapid dynamical heating has less of an effect on the galaxy-wide dispersion, which may be more influenced by spiral arms, warps, etc. We defer a more detailed investigation to future work.

Fig. 7 showed that the disc-wide vertical velocity dispersion of young stars in FIRE-2 galaxies is $\approx 15\text{--}25 \text{ km s}^{-1}$ higher than the MW’s $\approx 10 \text{ km s}^{-1}$ (Nordström et al. 2004; Casagrande et al. 2011; Anders et al. 2023). However, there are key differences in the spatial scales over which we measured our dispersions in Fig. 7 versus the scales used for MW observations. Given the non-trivial differences between the disc-wide and local velocity dispersions shown in Fig. 8, we now revisit our previous comparison between galaxy kinematics and stellar mass from Fig. 7, and now include observed values from the MW and nearby galaxies. This more detailed comparison between stellar kinematics in our simulations, in the MW, and in nearby galaxies is especially important because detailed MW observations lay the bedrock for much of our understanding of disc formation and dynamical heating. However, the MW is only one galaxy, and it remains unclear how representative the MW is of similarly massive galaxies.

For these comparisons, we compute the average local velocity dispersion of our galaxies, as opposed to the disc-wide dispersion of our previous results. Specifically, motivated by the typical measurement scales of the observations that we compare against here, we measure the ‘local’ velocity dispersion within annular bins centred at 8 kpc and then compute the average across the annulus.

Fig. 9 shows the local σ_{tot} (top) and $v_{\phi}/\sigma_{\text{tot}}$ (bottom) for young (age < 100 Myr) stars versus stellar mass. We also show the observed velocity dispersion of young stars from Casagrande et al. (2011) for the MW, Dorman et al. (2015) for M31 (Andromeda), Quirk et al. (2022) for M33 (Triangulum), and Pessa et al. (2023) for 19 galaxies from PHANGS-MUSE. We show velocity dispersions for only young stars; in McCluskey et al. (in preparation), we will pursue a more rigorous comparison of the full relation between velocity dispersion and age using mock observations.

An important caveat is that each of these observational works measured stellar populations and velocity dispersions somewhat

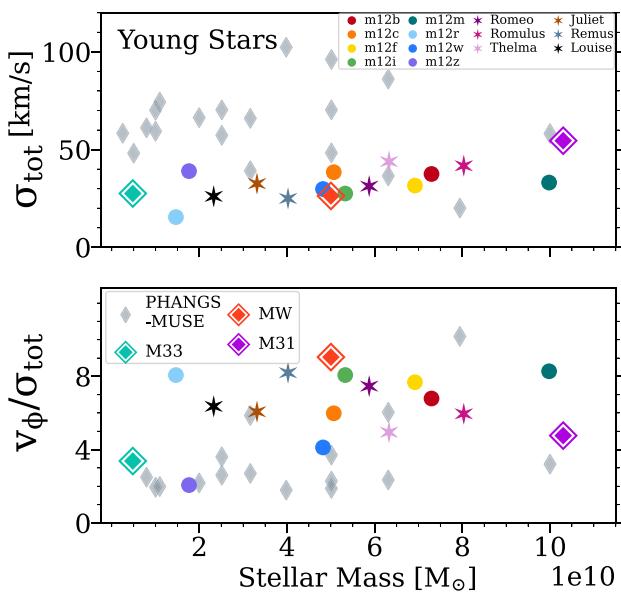


Figure 9. Total velocity dispersion, σ_{tot} , and median $v_{\phi}/\sigma_{\text{tot}}$, today for young (age < 100 Myr) stars versus stellar mass for 14 galaxies. As before, we show the isolated (m12) galaxies via circles and LG-like paired galaxies via stars. We also show observations of young stars in the LG: the MW (Casagrande et al. 2011), M31 (Dorman et al. 2015), and M33 (Quirk et al. 2022), and in 19 nearby star-forming disc galaxies from the PHANGS-MUSE survey (Pessa et al. 2023). Young stars in FIRE-2 galaxies have kinematics that broadly agree with most observed galaxies; their $v_{\phi}/\sigma_{\text{tot}}$ are slightly lower than the MW, but two galaxies are within 10 per cent of the MW. Young stars in the MW are kinematically colder than those in M31, M33, and all but one of the 19 galaxies in PHANGS-MUSE.

differently. In particular, the M31, M33, and MW measurements used individual stars, but the PHANGS-MUSE measurements used integrated spectra from regions of size ≈ 100 pc, which may include some contribution from older stars, leading to larger velocity dispersions. Nevertheless, the larger sample from PHANGS-MUSE provides an important statistical context for both our simulations and the LG. M31 and M33 observations measure resolved stars in multiple apertures of radius ≈ 760 and 350 pc, respectively, while the MW measurement includes stars within 60 pc of the Sun, so this represents only a single region of the MW. As a compromise to matching all of these works, and to limit numerical noise in our simulations, we use an intermediate aperture of diameter 250 pc (unlike in Fig. 8, where we used smaller 150 pc patches to better match ALMA’s resolution).

Moreover, we show the youngest stellar population from each work, but their age binning varies: the average age is 100 Myr for the MW, 30 Myr for M31, 80 Myr for M33, and 100 Myr for PHANGS-MUSE. Therefore, for this comparison only, we measure stars in our simulations today younger than 100 Myr. This reasonably matches the MW and PHANGS-MUSE, but these younger stars in M31 and M33 may have undergone somewhat less post-formation heating than in the MW, PHANGS-MUSE, or our simulations.

Finally, full 3D kinematics are available only for the MW: measurements of M31, M33, and the PHANGS-MUSE sample are limited to line-of-sight dispersions, σ_{LOS} . To estimate a total 3D dispersion, we assume isotropy, such that $\sigma_{\text{tot}}^2 = 3\sigma_{\text{LOS}}^2$, given the wide variety of observed inclination angles across these observations. This is a simplification, given that our simulated galaxies show σ that is not isotropic, with $\sigma_R > \sigma_{\phi} > \sigma_Z$. As such, our assumption may

underestimate the total dispersion of galaxies observed with smaller inclinations (more face-on), including the PHANGS-MUSE sample, which generally has inclinations $\approx 25\text{--}55^\circ$, while our assumption likely overestimates the total dispersion of highly inclined galaxies, such as M31 at 77° .

Our 14 FIRE-2 galaxies have an average local $\sigma_{\text{tot},\text{now}} \approx 32 \text{ km s}^{-1}$ for young stars, with isolated and LG-like galaxies exhibiting similar values, on average. Although our average $\sigma_{\text{tot},\text{now}}$ is slightly higher than the MW and M33’s values of ≈ 26.3 and 27.5 km s^{-1} , it is considerably lower than M31’s value of $\approx 55 \text{ km s}^{-1}$. However, the local total velocity dispersion of young stars shows a slight increase with galaxy mass, unlike the disc-wide vertical dispersion in Fig. 7. In turn, the high dispersion of M31 is likely (at least in part) due to its higher stellar mass. We emphasize that MW’s value of $\sigma_{\text{tot},\text{now}}$ is lower than that of all but 1 of the 21 galaxies in this PHANGS-MUSE sample. That said, five of our galaxies exhibit dispersions within 5 km s^{-1} of the MW’s value, while nine lie within 10 km s^{-1} , and three, the fairly low-mass m12r, Remus, and Louise, have even lower dispersions than the MW at ≈ 15.5 , 25.0, and 26.1 km s^{-1} , respectively.

Sanderson et al. (2020, fig. 2) also compared three FIRE-2 galaxies – m12i, m12f, m12m – against observations of the MW and M31; but for M31, they compared the disc-wide total (3D) dispersion in FIRE-2 against just the local line-of-sight (1D) dispersion of M31. Therefore, our comparison is more accurate and shows better agreement between FIRE-2 and the MW and M31.

Fig. 9 (bottom) compares $v_{\phi}/\sigma_{\text{tot}}$, which is a better, dimensionless metric of ‘disciness’ and the degree of rotational support. For each of the PHANGS-MUSE galaxies, we used their maximum v_{ϕ} in Lang et al. (2020), which they determined using PHANGS-ALMA measurements of CO (2–1) emission. We note that Lang et al. (2020) do not include v_{ϕ} values for three of the PHANGS galaxies; as such, we do not show these three galaxies in the bottom panel. As above, we emphasize that this observational compilation shows that young stars in the MW are kinematically colder, in terms of both σ_{tot} and $v_{\phi}/\sigma_{\text{tot}}$, than those in all but 1 of 21 nearby galaxies, including M31 and M33. Specifically, young stars in the MW have $v_{\phi}/\sigma_{\text{tot}} \approx 9$, while M31, M33, and the average PHANGS galaxy have values of ≈ 4.8 , 3.4, and 3.4. In turn, this initial comparison indicates that the MW is a kinematic outlier compared to similar-mass nearby galaxies; however, we will test this hypothesis more rigorously in McCluskey et al. (in preparation).

Secondly, young stars in FIRE-2 galaxies show a range of $v_{\phi}/\sigma_{\text{tot}}$ that is consistent with the range of the observed sample: all FIRE-2 galaxies reside within the observed range. The average $v_{\phi}/\sigma_{\text{tot}}$ of FIRE-2 is ≈ 6.4 , with LG-like and isolated galaxies having similar values. If anything, the typical FIRE-2 galaxy is kinematically colder than M31, M33, and the typical galaxy in PHANGS-MUSE. Although none of the FIRE-2 galaxies have $v_{\phi}/\sigma_{\text{tot}}$ quite as high as the MW’s value (≈ 9), two galaxies, Remus and m12m, are within 10 per cent of the MW’s value, at 8.3 and 8.2, respectively. Thus, this tentative comparison indicates that young stars in FIRE-2 have statistically realistic/representative values of σ_{tot} and $v_{\phi}/\sigma_{\text{tot}}$, provided that we measure both on spatial scales comparable to those of observations.

3.8 Time of disc onset

Galaxies transitioned from their Pre-Disc to their Early-Disc Era at their ‘time of disc onset’, which marks the start of the continuous process of disc settling. Most of our galaxies transitioned 4–8 Gyr ago, although Romeo transitioned particularly early, ≈ 11 Gyr ago.

We now examine the time of disc onset across our sample, including how it relates to a galaxy's stellar kinematics today. We also reiterate that, unlike in the previous subsection, all values are now disc-wide quantities, just as in the first 6 subsections of our results.

We defined the transition to the Early-Disc Era as when all stellar populations that formed after this time had $(v_\phi/\sigma_{\text{tot}})_{\text{form}} > 1$. Of course, v_ϕ/σ measured today may differ from its formation value. To examine how transition times derived using present-day kinematics compare to our fiducial formation-based times, we also measure the age when $(v_\phi/\sigma_{\text{tot}})_{\text{now}}$ permanently exceeds 1. Furthermore, we show that σ_Z evolved somewhat differently post-formation than the other components, so we additionally test this by also finding the oldest age after which $(v_\phi/\sigma_Z)_{\text{now}}$ permanently exceeds $\sqrt{3}$, under the simplifying (and in detail incorrect) approximation of isotropy. For all of these metrics, we tested a range of threshold values, and while this shifts the exact lookback time that a galaxy transitioned, it has little effect on the relative ordering within our sample, which is our primary metric of interest. We also tested transition times determined using only v_ϕ (both at formation and today) and found that they broadly agree with our v_ϕ/σ derived times, differing by $\lesssim 1$ Gyr for thresholds above $\approx 100 \text{ km s}^{-1}$. However, we limit our analysis to v_ϕ/σ -based metrics, because, unlike v_ϕ alone, v_ϕ/σ is dimensionless, largely independent of galaxy stellar mass in this mass range, and has a logical threshold ($v_\phi > \sigma$).

In summary, we use three thresholds to measure the time of disc onset that defines the transition from the Pre-Disc to Early-Disc Era:

- (i) $v_\phi/\sigma_{\text{tot}} > 1$ at formation
- (ii) $v_\phi/\sigma_{\text{tot}} > 1$ as measured today
- (iii) $v_\phi/\sigma_Z > \sqrt{3} \approx 1.73$ as measured today

Table 1 lists the lookback times/ages of disc onset for each galaxy. These three metrics give times within ≈ 0.5 –1 Gyr of each other for most galaxies. Times based on $(v_\phi/\sigma_Z)_{\text{now}}$ and $(v_\phi/\sigma_{\text{tot}})_{\text{now}}$ agree best with each other. Thus, the present vertical and total kinematics of stars provide similar insights into the disc's formation. However, lookback times from the kinematics at formation are generally earlier, because dynamical heating tends to decrease v_ϕ/σ of stars post-formation. A few galaxies yield substantially different times for different metrics. Notably, formation kinematics indicate that m12r transitioned 5.9 Gyr ago, while current kinematics indicate that it settled only 0.26 Gyr ago, because m12r experienced a merger with an LMC-mass galaxy ≈ 0.5 Gyr ago, which significantly heated existing stars. This lends caution to using present-day kinematics to infer the exact timing of events in a galaxy's formation history.

As Fig. 2 shows, once our galaxies transitioned to the Early-Disc Era, their stars formed on progressively cooler orbits as the disc continuously settled. If, after the time of disc onset, $(v_\phi/\sigma)_{\text{form}}$ grew at a broadly similar rate in each galaxy, then galaxies that had earlier times of disc onset would have higher $(v_\phi/\sigma)_{\text{form}}$ today, simply because they had more time to settle.

Fig. 10 tests this idea, showing the time of disc onset versus v_ϕ/σ_Z of stars younger than 250 Myr for our 14 galaxies. The three rows show our three thresholds in v_ϕ/σ for the onset of disc settling, which all show similar trends. We show isolated galaxies as circles and galaxies in LG-like pairs as stars, as in Fig. 7. Fig. 10 shows that the time of disc onset was ≈ 4.5 –8 Gyr ago ($z \approx 0.4$ –1) for most of our galaxies, although Romeo (see Section 3.2) has the earliest time of disc onset, ≈ 10.2 –11 Gyr ago ($z \approx 2$).

Fig. 10 shows that galaxies that experienced earlier disc onset have young stars (age < 250 Myr) with larger v_ϕ/σ_Z today, that is, they currently form stars on dynamically cooler orbits. While we do not show it, we find a similarly strong correlation between the time of

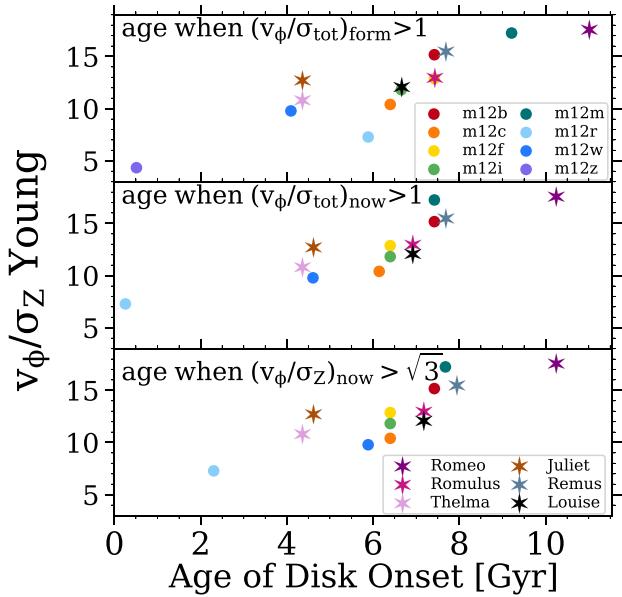


Figure 10. Degree of rotational support, v_ϕ/σ_Z , for young stars (age < 250 Myr) versus the lookback time of disc onset. Circles show isolated (m12) galaxies, and stars show galaxies in LG-like pairs. We examine three metrics to determine the time that each galaxy's disc started to settle, which yield similar but occasionally different times and ordering between galaxies. *Top:* Disc onset time defined as age when $v_\phi/\sigma_{\text{tot}}$ at formation permanently exceeded 1. *Middle:* Disc onset time defined age as when $v_\phi/\sigma_{\text{tot}}$ as measured today permanently exceeded 1. *Bottom:* Disc onset time defined age as when v_ϕ/σ_Z as measured today permanently exceeded 1.8. Overall, using kinematics at formation versus today or measuring different velocity components all lead to similar times of disc onset. All three metrics show that discs that started to settle earlier currently form stars on more rotationally supported orbits, which implies that the MW, which is unusually dynamically cold today, started to settle unusually early.

disc onset and the average v_ϕ/σ_Z and $v_\phi/\sigma_{\text{tot}}$ for both all and young stars. Table 2 lists the Pearson correlation coefficients, which are close to 0.9, and their corresponding p-values, which are $\sim 10^{-5}$ – 10^{-6} , indicating a strong correlation. We find similar results using instead the Spearman rank coefficients.

We also examined correlations with other galaxy properties. Surprisingly, the time of disc onset does not correlate much with the galaxy's current stellar mass within our sample. It also does not correlate with any component of σ , for all or young stars. However, the time of disc onset does correlate with the median v_ϕ of all or young stars, with Pearson coefficients ranging between 0.75 and 0.85 for all stars and between 0.73 and 0.80 for young stars. Despite the lack of correlation with σ , each threshold has its highest correlations with v_ϕ/σ , not with v_ϕ alone, further supporting that $(v_\phi/\sigma)_{\text{now}}$ is a more fundamental (and dimensionless) metric than σ_{now} to compare different galaxies' formation histories.

The strong correlation with the v_ϕ/σ of young stars today is particularly striking, considering that it only moderately correlates with v_ϕ of young stars. This correlation implies that, after disc onset, these discs settled at relatively similar rates, in order to preserve rank ordering. In particular, after disc onset, mergers and internal processes tend not to significantly alter the overall Hubble-time-averaged rate of disc settling.

In general, galaxies in LG-like pairs experienced earlier disc onset than isolated galaxies. This follows from Section 3.6, where LG-like galaxies had lower disc-wide σ and higher disc-wide v_ϕ/σ than

Table 2. Pearson correlations and p-values between the degree of rotational support of stars and different lookback times of galaxy transitions.

Metric	v_ϕ/σ_Z young corr. coeff. (p-value)	$v_\phi/\sigma_{\text{tot}}$ young corr. coeff. (p-value)	v_ϕ/σ_Z all corr. coeff. (p-value)	$v_\phi/\sigma_{\text{tot}}$ all corr. coeff. (p-value)
$v_\phi/\sigma_Z, \text{now} > 1.8$	0.90 (1.3×10^{-6})	0.91 (8.1×10^{-6})	0.86 (6.7×10^{-5})	0.92 (3.3×10^{-6})
$v_\phi/\sigma_{\text{tot}}, \text{now} > 1$	0.91 (6.1×10^{-6})	0.86 (8.8×10^{-5})	0.87 (5.1×10^{-5})	0.93 (2.4×10^{-6})
$v_\phi/\sigma_{\text{tot}, \text{form}} > 1$	0.89 (1.7×10^{-5})	0.89 (2.4×10^{-5})	0.81 (4.0×10^{-4})	0.88 (3.00×10^{-5})
t_{burst}	0.78 (3.0×10^{-3})	0.70 (0.011)	0.55 (0.067)	0.63 (0.028)

The first three rows are the lookback times of the onset of disc settling, using the three different metrics discussed in Section 3.8, while the fourth row is the lookback time corresponding to the galaxy's transition from bursty to steady star formation from Yu et al. (2021). All three times of disc onset show strong correlations with the degree of rotational support: earlier-forming discs are kinematically colder, both for young stars and for all stars, than later-forming discs.

isolated galaxies of similar (current) stellar mass, consistent with their larger radii (Bellardini et al. 2022). The earlier settling times for galaxies in LG-like environments follows from the fact that such galaxies assemble (form stars) earlier than those in isolated environments, which in turn follows from the fact that their dark-matter halos formed earlier (Garrison-Kimmel et al. 2019; Santistevan et al. 2020). Thus, galaxies in LG-like environments, despite ending up at similar stellar mass today as galaxies in isolated environments, formed stars earlier, their discs started to settle earlier, and their discs are radially larger and dynamically cooler today.

This dependence on the LG-like environment provides important insight into the assembly history of the MW (and M31) compared to other galaxies of similar mass today. As discussed in Section 3.7, young stars in the MW are dynamically colder than those in other nearby galaxies. If we naively extrapolate and place the MW's (local) $v_\phi/\sigma_{\text{tot}} \approx 8$ along the relation in Fig. 10, the trends within our simulated sample indicate that the MW disc began to settle earlier than our galaxies' discs ($\gtrsim 11$ Gyr ago). Remarkably, this simple extrapolation agrees with recent estimates of the age of the MW's thick disc: recent work has argued that its coherent rotation may have begun ≈ 12 –13 Gyr ago (Conroy et al. 2022; Xiang & Rix 2022). Thus, the early disc onsets exhibited by our simulated LG-like galaxies provides critical context for this seemingly early assembly of the MW's disc.

4 SUMMARY AND DISCUSSION

4.1 Summary of results

We summarize the main results from our analysis of stellar discs across the formation histories of 14 MW-mass galaxies from the FIRE-2 simulations. Most importantly, the kinematics of stars today do not simply reflect their kinematics at formation, but neither do they simply reflect post-formation dynamical heating; both processes are important. In particular, although the present-day dispersion of stars, σ_{now} , increases monotonically with age, σ_{form} does not.

We identified Three Eras of Disc Formation, when stars formed with different kinematics and experienced different degrees of post-formation dynamical heating:

(1) **Pre-Disc Era** (typically $\gtrsim 8$ Gyr ago): Stars formed on dispersion-dominated orbits, with $(v_\phi/\sigma_{\text{tot}})_{\text{form}} < 1$ and $v_\phi \lesssim 20 \text{ km s}^{-1}$. σ_{form} increased over time, reflecting the deepening of the gravitational potential. Thus, the oldest stars formed with the lowest dispersion. However, the present-day dynamics versus stellar age do not reflect this trend: σ_{now} increases monotonically with age, meaning that the oldest stars now have the highest dispersion. Thus, post-formation dynamical heating is most important during this era. Furthermore, the typical $v_{\phi, \text{now}} \approx 60 \text{ km s}^{-1}$, so the oldest stars often

have been 'spun up' since formation. As a result of both of these trends, $(v_\phi/\sigma_{\text{tot}})_{\text{now}}$ remains below 1, similar to formation.

(2) **Early-Disc Era** (typically ≈ 4 –8 Gyr ago): Began at the time of disc onset, when $v_\phi/\sigma_{\text{tot}} > 1$ and grew rapidly, and ended when the growth of v_ϕ/σ slowed significantly. Similarly, σ_{form} peaked at the beginning of this era, and decreased steadily throughout, reaching a minimum by the end of the era. The amount of dispersion added by post-formation heating was constant with age throughout this era, so the slope of σ_{now} with age is unchanged from that at formation. The majority of σ_{now} for these stars was in place at formation, not added via post-formation dynamical heating.

(3) **Late-Disc Era** (typically $\lesssim 4$ Gyr ago): Began when the growth of $(v_\phi/\sigma_{\text{tot}})_{\text{form}}$ slowed. Stars formed with a nearly constant $\sigma_{\text{tot}, \text{form}}$. $\sigma_{\text{tot}, \text{now}}$ increases monotonically with age, because post-formation dynamical heating increases with age, and the oldest stars in this era have σ_{now} with an equal contribution from post-formation heating and formation.

We examined the evolution of the three different components of the velocity. At formation, $\sigma_R, \text{form} > \sigma_\phi, \text{form} > \sigma_Z, \text{form}$ for all ages. Remarkably, across all eras and ages, the relative amount in each component is relatively unchanged. At present, σ_R, now continues to dominate at all ages, but $\sigma_Z, \text{now} \approx \sigma_\phi, \text{now}$ for stars that formed in the Early- and Pre-Disc Eras (ages $\gtrsim 4$ Gyr). Post-formation dynamical heating primarily increased σ_R at ages $\lesssim 4$ Gyr but acted nearly isotropically for older stars.

We examined the dependence of stellar kinematics on galaxy stellar mass. σ_{now} of all stars increases with galaxy stellar mass, but $(v_\phi/\sigma)_{\text{now}}$ of all stars is nearly constant. Young stars (age < 250 Myr) exhibit opposite trends: σ_{now} is independent of mass, while $(v_\phi/\sigma)_{\text{now}}$ increases with mass.

We compared disc-wide and local ($\lesssim 250$ pc) velocity dispersions for cold gas and young stars at formation: local dispersions are ≈ 3 times lower than disc-wide dispersions for cold gas and ≈ 2 times lower for stars at formation. Thus, the velocity dispersion strongly depends on the spatial scale measured over, which is critical for observational comparisons.

We compared locally measured kinematics in our 14 simulated galaxies to those observed in three LG galaxies and nearby star-forming galaxies, compiling measurements of σ and v_ϕ/σ for young stars (age $\lesssim 100$ Myr) in the MW, M31, M33, and 19 galaxies from PHANGS-MUSE. The MW is significantly kinematically colder than both M31 and M33, and all but one of the PHANGS galaxies. For young stars, all of our simulations show broad agreement with observations, with three of our galaxies even exhibiting lower locally measured σ_{tot} than the MW.

We quantified the time of disc onset, the lookback time that each galaxy's disc began to settle, using thresholds in v_ϕ/σ , measured both today and at formation. The degree of rotational support of young stars today correlates strongly with the time of disc onset: earlier-

settling galaxies currently form colder discs. Most of our galaxies began to settle $\approx 4\text{--}8$ Gyr ago, with galaxies in LG-like environments generally starting to settle earlier than isolated galaxies. Romeo, a galaxy in an LG-like pair, began to settle the earliest, $\approx 10\text{--}11$ Gyr ago, likely the closest to the MW.

In Appendix C, we examine the impact of including self-consistent CR injection/transport from stars, assuming a constant diffusion coefficient, which in the FIRE-2 simulations leads to galaxies with lower stellar masses, lower rotational velocities, and lower dispersions. Overall, the inclusion of CRs does not change significantly v_ϕ/σ at a given stellar mass.

4.2 Discussion

4.2.1 Caveats

We first discuss key limitations and caveats to our analysis of 14 MW-mass galaxies from the FIRE-2 simulations.

First, we reiterate the selection function of our simulated sample. We selected isolated or LG-like paired dark-matter halos from dark-matter-only simulations at $z = 0$ with $M_{200m} \approx 1\text{--}2 \times 10^{12} M_\odot$, without additional selection based on galaxy properties or merger/formation history, etc. Thus, we did not choose our galaxies specifically to be analogues of the MW (or M31 or M33). As such, our galaxy-averaged results represent randomly sampled formation histories of MW-mass galaxies, which should fairly sample the cosmological scatter in formation histories.

Secondly, these FIRE-2 simulations do not model all of the physics that may be relevant for disc formation, in particular, active galactic nucleus (AGN) feedback from black holes. The degree to which AGN may impact stellar dynamics in the solar annuli of MW-mass galaxies (which in this work we define as $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc) is unclear. While AGN feedback is probably most crucial for more massive galaxies and most significantly alters stellar dynamics in a galaxy's innermost regions (Dubois et al. 2013; Weinberger et al. 2017; Mercedes-Feliz et al. 2023; Wellons et al. 2023), growing evidence from simulations indicates that AGN can both directly and indirectly affect the kinematic and structural properties of MW-mass and low-mass galaxies (Dashyan et al. 2018; Koudmani, Henden & Sijacki 2021; Irodou et al. 2022). For example, Grand et al. (2017) found that the primary effect of AGN in MW-mass galaxies was to suppress central star formation, which in turn prevents the formation of overly massive bulges. Indeed, Grand et al. (2017) found that a simulation of a MW-mass galaxy reran from $z = 1$ without AGN feedback had a smaller fraction of stars on disc-like orbits (defined as having orbital circularities > 0.7), more dominant bulges, and enhanced star formation in the outer disc leading to a more extended disc component compared to versions run with AGN feedback.

Furthermore, our analysis of birth kinematics has limitations. We determined a star particle's position and velocity relative to the galaxy's major axes, which are stable and well-defined at present, but not necessarily at earlier times, when accretion and mergers were more significant and even the notion of a single main galaxy progenitor can be ill-defined (Wechsler et al. 2002; Giocoli et al. 2007; Santistevan et al. 2020). Thus, cleanly separating the R, ϕ , and Z components is difficult at early times. However, all components of the dispersion at formation display similar trends as σ_{tot} , which is independent of the determination of the axes, implying that the impact of this limitation on our overall qualitative results is likely minor.

Furthermore, as we discussed in Section 2.2, our simulations store snapshots every 20–25 Myr, so our star 'formation' properties include

this additional ≈ 12 Myr of post-formation dynamical evolution, on average. This early heating can be significant: using an idealized simulation of an isolated MW-like galaxy with sub-parsec resolution (0.05 pc), Renaud et al. (2013) found that stars that formed with $\sigma_{\text{tot,form}} \approx 10 \text{ km s}^{-1}$ were heated to 15 km s^{-1} in just 10 Myr. Similarly, our Fig. 6 shows that σ_{tot} increased rapidly by $\approx 30 \text{ km s}^{-1}$ in the first ≈ 125 Myr after formation. We tested whether this early evolution could change our conclusions about the relative importance of 'formation' dispersion versus post-formation heating. For stars younger than 10 Gyr, $\sigma_{\text{tot,now}}/\sigma_{\text{tot,form}} < 1.95$ for all but ≈ 3.5 Gyr old stars, meaning that post-formation heating is less than 95 per cent of their $\sigma_{\text{tot,form}}$. If we assume that stars were heated at the same effective rate as the youngest stars in Fig. 6 for ≈ 12 Myr, our revised $\sigma_{\text{tot,now}}/\sigma_{\text{tot,form}}$ still remains below 1.95 for the same ages. Even if we assume an initial effective heating rate that is 4× larger, $\sigma_{\text{tot,now}}/\sigma_{\text{tot,form}}$ remains less than 2 for the same ages. Thus, our results about the relative importance of post-formation heating remain unchanged.

Lastly, our analysis did not account for uncertainties in stellar age. Stellar ages are difficult to accurately measure, and uncertainties on stellar ages for large populations are generally $\gtrsim 20\text{--}30$ per cent (for example Soderblom 2010), although the synergy of astrometric and spectroscopic surveys with asteroseismological data has driven tremendous progress in the past decade (for example Silva Aguirre et al. 2018; Mackereth et al. 2019b). Previous works showed that the inclusion of observationally motivated age uncertainties of 20–30 per cent in simulations can (1) obscure features from mergers in the age-velocity dispersion relationship (Martig, Minchev & Flynn 2014; Buck et al. 2020), and (2) decrease inferred exponential heating rates via the 'smearing' of age bins (Aumer, Binney & Schönrich 2016). In McCluskey et al. (in preparation) we will examine the kinematic trends of our simulations versus age at $z = 0$, incorporating realistic age uncertainties, which will allow us to compare to observational trends in much more detail.

4.2.2 Disc galaxy formation in FIRE simulations

As we discussed in the Introduction, our work builds upon a series of analyses of FIRE simulations that studied the formation histories of MW-mass galaxies, in particular, the co-evolution of their stellar and gaseous dynamics, morphology, star formation burstiness, metallicity spatial variations, and virialization of the inner CGM (Ma et al. 2017; Garrison-Kimmel et al. 2018; Stern et al. 2021; Yu et al. 2021; Bellardini et al. 2022; Hafen et al. 2022; Trapp et al. 2022; Yu et al. 2022; Gurvich et al. 2023; Hopkins et al. 2023). We further quantified the notion that MW-mass disc galaxy formation occurred in distinct phases, showing that stellar kinematics, both at formation and at present, follow broadly contemporaneous evolutionary phases.

In particular, what we define as the transition from the Early-Disc to the Late-Disc Era coincides with the virialization of the inner CGM (Stern et al. 2021), the transition from bursty to steady star formation (Yu et al. 2021), and a transformation in the energetics and angular momentum of both the accreted gas (Hafen et al. 2022) and the ISM (Gurvich et al. 2023). During this stage of steady star formation in the Late-Disc Era, σ_{form} was low and constant, with the gas turbulence reaching a floor maintained by the thermal and turbulent pressure of the warm neutral/ionized medium (for example Leroy et al. 2008), stellar feedback, and spiral arms.

Yu et al. (2021) quantified the lookback time that star-formation transitioned from bursty to steady – the 'burst time' or t_b – for 12 of the 14 galaxies that we analyse. Table 1 lists each galaxy's t_b from

Yu et al. (2021). The transition to steady star formation typically occurred ≈ 2 Gyr after our ‘time of disc onset’, when stellar $v_\phi/\sigma_{\text{tot}} > 1$. Nominally, this implies that star formation remained bursty for a while during the Early-Disc Era, in agreement with our finding that σ_{form} remained high (but declining) throughout this era. However, our times of disc onset are based on stellar dynamics within the solar annulus today ($R = 6\text{--}10$ kpc and $|Z| < 3$ kpc) and only moderately correlate with t_b , having Spearman coefficients between 0.75–0.78. On the contrary, Yu et al. (2021) included all stars that formed within 20 kpc physical of the galactic centre. If instead we extend our radial range to $R < 12$ kpc, this yields times that (1) are on average 1.2 Gyr later than our original times, because of the hotter kinematics of the inner region, and (2) more strongly correlate with t_b , with Spearman coefficients between 0.94 and 0.96 for our three dynamical thresholds. Thus, these estimates for when a galaxy’s disc began to form/settle can depend on the selection of stars today, with stars in the solar neighbourhood yielding earlier settling times. We will investigate radial trends in these simulations in future work.

Using the same FIRE-2 simulations, Yu et al. (2022) argued that they can subdivide their bursty star-formation era into two eras: an early bursty and chaotic phase, when stars formed with irregular morphology, and a subsequent bursty phase, when stars formed with discy morphology. These bursty star-formation eras generally correspond to our Pre-Disc and Early-Disc eras. Yu et al. (2022) also studied how both the birth and present-day orbits of stars reflect the era when they formed, classifying stars as belonging to the thin-disc, thick-disc, and isotropic spheroid based on their orbital circularity at both formation and at $z = 0$. They found that, in general, a star’s orbital classification at formation reflects the era in which it formed: stars primarily formed on isotropic spheroid orbits in the early bursty star-formation phase, thick-disc orbits in the later bursty star-formation phase, and thin-disc orbits in the steady star-formation phase. However, Yu et al. (2022) noted that this is not a one-to-one mapping: ≈ 34 per cent of stars that formed with thin-disc orbits now have thick-disc orbits, while a similar fraction of stars that formed with thick-disc orbits now have isotropic spheroid orbits. Furthermore, the fraction of stars that were born ‘thin’ but are now ‘thick’ increases with age. Our results agree with theirs, and we further quantify the degree of post-formation heating/perturbations that stars experienced.

4.2.3 Comparison to other cosmological zoom-in simulations

Various works have used different cosmological zoom-in simulations to study the formation histories of MW-mass galaxies (for example Bird, Kazantzidis & Weinberg 2012; Brook et al. 2012; Martig, Minchev & Flynn 2014; Grand et al. 2016; Buck et al. 2020; Agertz et al. 2021; Bird et al. 2021). Across a wide range of different numerical implementations of hydrodynamics, star formation, and feedback, cosmological simulations consistently agree that the dynamical history of a MW-mass disc arises from both cosmological disc settling and post-formation dynamical heating; one cannot neglect one or the other. As we discuss below, most current-generation zoom-in simulations of MW-mass galaxies also broadly agree in both the typical disc-wide velocity dispersion of young stars at $z \approx 0$ and the shape/magnitude of the stellar age dependence, as measured today.

Our results agree with the VINTERGARTEN simulation (Agertz et al. 2021), which used the same initial conditions as m12i but simulated with the adaptive mesh refinement code RAMSES (Teyssier 2002). Their $\sigma_{Z,\text{now}}$ versus age (fig. 14 of Agertz et al. 2021) is

nearly identical to ours for m12i, both in value and shape: both of our simulations show near-exponential increases in $\sigma_{Z,\text{now}}$ with age for stars younger than 6 Gyr and a jump at 6 Gyr. Our results also agree well with the NIHAO-UHD suite (Buck et al. 2020): their five galaxies, which span $M_{\text{star}} = 1.5\text{--}16 \times 10^{10} M_\odot$, have $\sigma_{Z,\text{now}} \approx 20\text{--}35 \text{ km s}^{-1}$ for young stars, and $\approx 60\text{--}120 \text{ km s}^{-1}$ for stars 12 Gyr old, similar to ours. Any difference in the exact age dependence of each galaxy’s dispersion follows from its unique formation history.

Grand et al. (2016) analysed 16 MW-mass galaxies from the Auriga simulations, also finding broadly similar disc-wide $\sigma_{Z,\text{now}}$ and $\sigma_{Z,\text{form}}$ as we do. Similar to FIRE, NIHAO-UHD, and VINTERGARTEN, Auriga simulations have typical disc-wide $\sigma_{Z,\text{now}} \approx 20\text{--}25 \text{ km s}^{-1}$. They also concluded that cosmological disc settling was generally the primary effect that set the relation between the velocity dispersion and the age of stars today. However, Grand et al. (2016) found that bars were the strongest contributor to vertical dynamical heating. This seemingly contradicts our results, because stars in all of our galaxies had significant post-formation increases to their disc-wide σ_Z , including stars in unbarred or weakly-barred galaxies. Furthermore, even our galaxies that do have significant bars near $z = 0$ (such as m12m) did not at early times, and our results indicate that heating is most significant for stars that formed in the Pre-Disc Era. Although we do not study the effect of bars in this work, Ansar et al. (2023) quantifies the incidence of bars in the FIRE-2 simulations, and in future work we will address the drivers of dynamical heating in our simulations, including bars. That said, while it is possible that the bar heated these old stars more recently, bars most likely do not heat only old stars.

Our results may differ from those of Grand et al. (2016) due to numerical differences. First, Auriga uses the subgrid model from Springel & Hernquist (2003), which treats the ISM as a two-phase medium with an effective equation of state that inhibits the formation of dense gas, so that star formation occurs at a much lower density, $n > 0.13 \text{ cm}^{-3}$, compared to $n > 1000 \text{ cm}^{-3}$ in FIRE-2. House et al. (2011) found that low thresholds in density impose a dispersion ‘floor’ in σ_{form} , even in high-resolution simulations (50 pc), because stars form from higher-temperature gas. House et al. (2011) find that this floor effectively disappears if one increases the star-formation threshold from $n = 0.1$ to 1000 cm^{-3} (see also Martig, Minchev & Flynn 2014; Kumamoto, Baba & Saitoh 2017). That said, given that Grand et al. (2016) find that σ_{form} decreases over time, this dispersion ‘floor’ may have little effect on their results.

Secondly, certain mechanisms that induce dynamical heating in our simulations may be different (and less significant) in the Auriga simulations. The simulations in Grand et al. (2016) had baryonic mass resolution of $4 \times 10^4 M_\odot$, 8–10 times larger than ours, and they used lower spatial resolution, with gravitational force softenings of 375 pc physical for stars and gas (which is comparable to the thin-disc scale height of the MW), while our FIRE-2 simulations use 2.7–4 pc for star particles and ≈ 1 pc minimum and ~ 4 pc ISM-averaged for gas cells. Thus, Auriga simulations do not resolve the dense ISM, including important dynamical structures like GMCs that likely heat star particles after their formation in FIRE. Furthermore, the modeling of stellar feedback in Auriga is markedly different: supernovae launch ‘wind particles’ that are temporarily decoupled from the hydrodynamics until the particle encounters sufficiently low-density gas, which does not couple that feedback directly to dense gas in a star-forming region. In turn, our simulations may exhibit more gas turbulence. That said, the Auriga simulations do model certain physics that these FIRE-2 simulations do not include, most importantly, AGN feedback.

Grand et al. (2016) also found that spiral arms did not considerably heat the disc, while we proposed that spiral arms are probably responsible for the significant increase in $\sigma_{\text{tot, now}}$ of young stars, in part because σ_R dominates the overall dispersion (see also Orr et al. 2022). Our two results are not necessarily in disagreement, however, because Grand et al. (2016) only considered vertical heating. Spiral arms predominately increase the in-plane dispersion (directly), and while GMC deflections can theoretically redirect this in-plane heating in the vertical direction, the Auriga simulations do not resolve GMCs. That said, although FIRE-2 can resolve massive GMCs, young stars in our simulations have vertical effective heating rates that are still $\approx 4 \times$ lower than their radial rates, as Fig. 6 showed, indicating that (in-direct) spiral-driven vertical heating is minor in our simulations at present-day.

Many previous works – both observational and theoretical – only examine the vertical velocity dispersion. However, vertical post-formation dynamical heating operates differently than total and in-plane heating. For example, Fig. 3.4 (top) shows that the amount of vertical dispersion added post-formation increases approximately linearly across all ages, while in-plane heating plateaus for Early-Disc stars. Similarly, Fig. 3.4 (bottom) shows that the present-day vertical dispersion has the largest relative impact from post-formation heating: the dispersion added from heating typically begins to equal (or exceed) the dispersion from formation at ≈ 4 Gyr for the vertical component, but does not rival the in-plane dispersions from formation until ≈ 11 Gyr. Furthermore, the vertical dispersion typically exhibits the smallest values. For these reasons, we do not recommend using vertical heating to draw inferences regarding total heating.

As we discussed above, nearly all current cosmological zoom-in simulations form MW-mass galaxies with typical disc-wide $\sigma_{Z, \text{now}} \gtrsim 20 \text{ km s}^{-1}$, not as dynamically cold as the (solar-neighbourhood) measurements of MW today. However, Bird et al. (2021) analysed a single MW-mass galaxy, h277, selected from the g14 suite of simulations (Christensen et al. 2012), generated using the parallel *N*-body + SPH code GASOLINE (Wadsley, Stadel & Quinn 2004). That simulation had disc-wide $\sigma_{Z, \text{now}} < 10 \text{ km s}^{-1}$, measured within an annulus of $R = 8\text{--}10$ kpc, for stars younger than 1 Gyr, similar to (primarily locally measured) values for the MW (for example Nordström et al. 2004; Mackereth et al. 2019b).

The cold disc-wide kinematics of h277 likely reflect its formation history: its disc began to settle early, ≈ 9 Gyr ago, and it had no major mergers since $z \approx 2$. That said, one of our galaxies, Romeo, has a similar time of disc onset and merger history as h277 but has a larger (disc-wide) $\sigma_Z \approx 16 \text{ km s}^{-1}$ for stars 1 Gyr old today. However, we note that Romeo has a local $\sigma_Z \approx 12 \text{ km s}^{-1}$, which is closer, albeit still higher, than MW measurements. Therefore, differences in our simulations' physical and numerical modeling likely contribute to our unequal dispersions. Although the baryonic mass resolution, modeling of dense ISM, and molecular-based star-formation criteria are broadly similar between the FIRE-2 simulations and h277, key numerical differences exist between our simulations: h277 used SPH for its hydrodynamics solver (instead of MFM); used larger gravitational force softening of 173 pc for all particle species; used the 'blastwave' feedback model for core-collapse supernovae in which cooling is turned off for gas within the blast wave radius to mimic the adiabatic expansion of a remnant (Stinson et al. 2006); used only thermal energy injection for white-dwarf (Type Ia) supernovae; and did not include stellar winds or radiative feedback. While it is unclear how any of these individual aspects do or do not affect their results, differences in feedback modeling may contribute to their lower dispersions to some extent. FIRE-2 accounts for energy and momentum injection for both core-collapse and white-dwarf

supernovae, and includes stellar winds, radiation pressure, and photo-ionization. This feedback launches starburst-induced star-forming outflows at early times [see the top right panel of our Fig. 1 and El-Badry et al. (2018a)], which Bird et al. (2021) notes are absent in their simulation. As such, it is possible that our additional feedback channels contribute to larger gas velocity dispersions and subsequently larger σ_{form} values in our simulations, although it is unclear if this persists to later cosmic times. Similarly, Hopkins, Quataert & Murray (2012) and Agertz et al. (2013) both show that the structure of the ISM can non-trivially vary depending on the feedback model used. Despite any differences, we again emphasize that both the disc-wide $\sigma_{Z, \text{form}}$ and $\sigma_{Z, \text{now}}$ exhibited similar trends with age in h277 and our FIRE-2 simulations.

4.2.4 Connections to the early formation of the MW

As we discussed in Sections 3.6 and 3.8, the cold kinematics of MW stars may indicate that the MW's disc began to settle unusually early in its history. For example, using our predicted result that the disc-wide kinematics of young stars today correlates with when the disc began to settle, and placing the MW's value of $v_\phi/\sigma_{\text{tot}} \approx 8$ along the relation from our simulations in Fig. 10, this implies that the MW's time of disc onset was $\gtrsim 11$ Gyr ago, which agrees with recent estimates from observations (for example Belokurov & Kravtsov 2022; Conroy et al. 2022; Xiang & Rix 2022). Thus, our analysis of disc settling times in cosmological simulations and its relation to present-day kinematics complements and augments these works.

That said, measurements of the MW's kinematics are largely restricted to one local patch: the solar neighbourhood. As we show in Fig. 8, locally measured velocity dispersions in our simulations (at formation) are typically ≈ 2 times smaller than disc-wide measurements. In turn, while all of our disc-wide dispersions are larger than the MW's, three of our galaxies have present-day local dispersions lower than the MW's. Furthermore, the galaxy-to-galaxy scatter in the relation between disc-wide and local kinematics is significant. For example, stars today younger than 100 Myr in one of our most MW-like galaxies, Romeo, have disc-wide σ that is 1.1–1.2 times higher than its local value, while m12i, which is otherwise similar to Romeo, has disc-wide values that are 1.5–1.7 times higher than local values. More work is needed to determine how the MW's disc-wide dispersion compares to its local dispersion, and if these differences affect its placement on our disc onset – young stellar kinematics relation. In McCluskey et al. (in preparation) we will more rigorously compare our simulations to the MW as well as other observed galaxies.

We now discuss connections with some of these recent observational works. Conroy et al. (2022) combined Gaia astrometry and H3 Survey spectroscopy to posit that the MW disc formed in three dynamical eras, broadly consistent with our Pre-Disc, Early-Disc, and Late-Disc Eras. Conroy et al. (2022) proposed that the MW began in a 'simmering phase', during which the star formation efficiency was low and stars formed kinematically hot, but that $\approx 12\text{--}13$ Gyr ago the MW transitioned into a 'boiling phase', during which the SFE strongly increased and stars formed with increasingly 'discy' dynamics, marking the 'birth of the Galactic disc'. This thick disc continued to grow and settle in a 'boiling phase' for $\approx 3\text{--}4$ Gyr (from $z \approx 4$ to $z \approx 1$) until the GSE merger, after which the star formation efficiency decreased and stars formed in a dynamically colder, thin disc until today.

Similarly, Belokurov & Kravtsov (2022) combined Gaia astrometry with spectroscopy from the APOGEE Data Release 17. Although

Belokurov & Kravtsov (2022) lacked empirical age estimates, their general picture agrees with ours: old, metal-poor, *in-situ* stars were born in a ‘messy’ phase, then the MW ‘spun-up’ as stars rapidly became metal-rich and rotational earlier than ≈ 8 Gyr ago. Belokurov & Kravtsov (2022) posited that this rapid formation of the MW’s disc took place over 1–2 Gyr, with the GSE merger occurring after the initial formation of a thick-disc, such that the merger heated disc stars onto halo-like orbits. On the other hand, Xiang & Rix (2022) conclude that the formation of the *in-situ* halo and thick-disc overlapped. In this picture, the GSE merger occurred ≈ 11 Gyr ago, 1–2 Gyr earlier than previous estimates, and did not cause a transition to thin-disc formation but instead enhanced thick-disc formation.

Our results help connect these present-day observations to understanding the MW’s formation history. In particular, the presence of coherent rotation in a stellar population today does not necessarily indicate the presence of coherent rotation at the time of formation in our simulations. Observations show that older, more metal-poor stars in the MW rotate slower than younger stars, but they still display some coherent rotation. In our simulations, this is not primarily the result of these Pre-Disc stars being born with coherent rotation, but rather, this tends to arise from post-formation torquing (see Fig. 1), for example, caused by mergers (see Bellardini et al., in preparation). m12m provides an illustrative example: it experienced several major mergers at early times, broadly similar to estimates of when the GSE merger occurred. As Fig. 3 shows, m12m has old stars with $v_{\phi,\text{now}} \approx 200 \text{ km s}^{-1}$, despite having $v_{\phi,\text{form}} \approx 0 - 50 \text{ km s}^{-1}$. Thus, mergers like the GSE merger could have torqued old stars in the MW onto more coherently-rotating orbits.

Ultimately, the MW’s kinematics reflect its individual formation history. While the MW provides a useful benchmark, the fact that most current cosmological simulations form galaxies with hotter kinematics than the MW may speak to the MW’s early settling time rather than to limitations of the simulations. Indeed, the comparison of the MW to M31 and M33 in Fig. 7 affirms that the MW may be an outlier among galaxies at its mass. Therefore, future works that study the kinematics of stars and their relation to stellar ages in external galaxies will provide excellent opportunities to place the MW in a statistical context and study the combination of cosmological disc settling and dynamical heating across a range of formation histories (Dorman et al. 2015; Leaman et al. 2017; Quirk et al. 2022).

ACKNOWLEDGEMENTS

We thank Jonathan Stern for helpful comments and the participants of the Disk Formation Workshop at UC Irvine for their enlightening discussions. FM and AW received support from: NSF via CAREER award AST-2045928 and grant AST-2107772; NASA ATP grant 80NSSC20K0513; HST grants AR-15809, GO-15902, GO-16273 from STScI; and a Scialog Award from the Heising-Simons Foundation. CAFG was supported by NSF through grants AST-2108230 and CAREER award AST-1652522; by NASA through grants 17-ATP17-0067 and 21-ATP21-0036; by STScI through grant HST-GO-16730.016-A; by CXO through grant TM2-23005X; and by the Research Corporation for Science Advancement through a Cottrell Scholar Award. PFH received support by NSF Research Grants 1911233 and 20009234, NSF CAREER grant 1455342, and NASA grants 80NSSC18K0562 and HST-AR-15800.001-A. We performed some of this work at the Aspen Center for Physics, supported by NSF grant PHY-1607611. We ran simulations using: XSEDE, supported by NSF grant ACI-1548562; Blue Waters, supported by the NSF; Frontera allocations AST21010 and AST20016, supported by the

NSF and TACC; Pleiades, via the NASA HEC program through the NAS Division at Ames Research Center.

DATA AVAILABILITY

All of the PYTHON code that we used to generate these figures are available at <https://fmclluskey.github.io>, which uses the publicly available package https://bitbucket.org/awetzel/gizmo_analysis (Wetzel & Garrison-Kimmel 2020). The FIRE-2 simulations are publicly available (Wetzel et al. 2023) at <http://flatuhub.flatironinstitute.org/fire>. Additional FIRE simulation data is available at <https://fire.northwestern.edu/data>. A public version of the GIZMO code is available at <http://www.tapir.caltech.edu/~phopkins/Site/GIZMO.html>.

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APPENDIX A: EFFECT OF VERTICAL SELECTION

By default, we select stars currently at $|Z| < 3$ kpc, which includes stars in both the thin and thick disc, to sample disc stars at all ages in an unbiased manner and avoid assumptions about thin- versus thick-disc dynamics. Fig. A1 shows the effect of the vertical selection on $\sigma_{Z, \text{now}}$ and $\sigma_{\text{tot}, \text{now}}$ versus stellar age, specifically, the fractional change relative to our fiducial selection of $|Z| < 3$ kpc. We also examined the effect on $\sigma_{R, \text{now}}$ and $\sigma_{\phi, \text{now}}$ and they show trends nearly identical to $\sigma_{\text{tot}, \text{now}}$ so we do not show them separately.

At most ages, our vertical selection has minimal impact on any component of the dispersion, $\lesssim 3$ per cent. This is because most

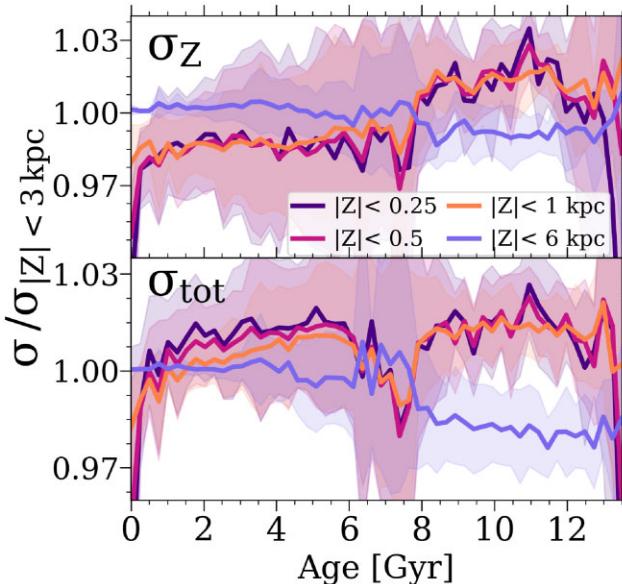


Figure A1. Effect of vertical selection on σ_Z (top) and σ_{tot} (bottom): present-day velocity dispersion for stars selected within four vertical ranges, divided by the dispersion for our fiducial range of $|Z| < 3$ kpc, versus stellar age. We select all stars at $R = 6\text{--}10$ kpc. Thinner vertical selections lead to slightly higher σ_Z and σ_{tot} for stars older than ≈ 8 Gyr, but slightly lower σ_Z and higher σ_{tot} for stars younger. However, different vertical selections change the dispersion by $\lesssim 3$ per cent, with the exception of the oldest and youngest stars, which have ≈ 5 per cent lower dispersions for thinner vertical ranges.

disc stars lie well within $|Z| < 1$ kpc, so they dominate regardless of vertical selection. Thinner vertical selection does yield slightly higher $\sigma_{\text{tot}, \text{now}}$, although with significant scatter. Interestingly, thinner selection only yields higher $\sigma_{Z, \text{now}}$ for stars older than ≈ 8 Gyr, and instead yield slightly lower values of $\sigma_{Z, \text{now}}$ for younger stars. The youngest population formed near the mid-plane and experienced ≈ 125 Myr of post-formation heating, on average, so most are still near the mid-plane and have low dispersions. Any that are above $|Z| = 0.25\text{--}1$ kpc likely experienced an atypical amount of heating. Again, all of these effects are at the level of $\lesssim 3$ per cent, where non-equilibrium effects like disc warping and asymmetries may become important.

If instead we measure the dispersion for the galaxy overall (not in age bins) as in Section 7, larger vertical selections non-trivially increase the dispersion. Vertical dispersions within $|Z| < 3$ kpc are 10–35 per cent larger than within $|Z| < 1$ kpc, because larger vertical selections include a higher fraction of older, higher-dispersion stars.

APPENDIX B: EFFECT OF IN-SITU SELECTION

By default, we select only stars that are at $R = 6\text{--}10$ kpc and $|Z| < 3$ kpc at present. We include all of these stars when measuring present-day kinematics, but when measuring kinematics at formation, we select only *in-situ* stars that formed within 30 kpc comoving of the galactic centre. Fig. B1 shows the effect of also including this *in-situ* selection on the current velocity dispersion versus stellar age. As above, we only show $\sigma_{Z, \text{now}}$ and $\sigma_{\text{tot}, \text{now}}$, because $\sigma_{R, \text{now}}$ and $\sigma_{\phi, \text{now}}$ exhibit the same trends as $\sigma_{\text{tot}, \text{now}}$. Overall, our *in-situ* selection has no effect on stars younger than ≈ 6 Gyr, because almost all of them formed in situ. However, the effects of *in-situ* selection become increasingly significant for older stars, when a more significant fraction formed ex situ. Not surprisingly, including those older *ex-situ* stars increases the dispersion. That said, this effect is $\lesssim 10$ per cent for all but the oldest stars [see El-Badry et al. (2018b) for analysis of these ancient stars].

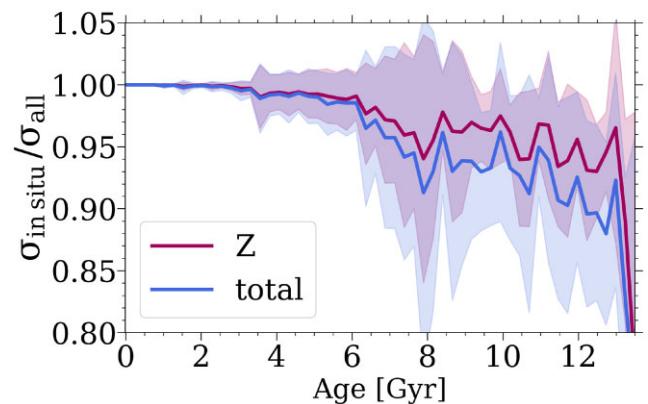


Figure B1. Effect of excluding *ex-situ* stars on the present-day velocity dispersion: $\sigma_{Z, \text{now}}$ (top) and $\sigma_{\text{tot}, \text{now}}$ (bottom) of only stars that formed *in situ* divided by that of all (*in-situ* + *ex-situ*) stars, versus age, averaged across 11 galaxies. Stars with ages $\lesssim 6$ Gyr show no significant difference, because almost all young disc stars formed *in situ*. Older than 6 Gyr, the dispersion of *in-situ* stars relative to all stars decreases with age, because the fraction of *ex-situ* stars increases with age, and *ex-situ* stars have larger present-day dispersions. However, at almost all ages that we investigate, the difference in velocity dispersion is $\lesssim 10$ per cent.

APPENDIX C: EFFECT OF COSMIC-RAY FEEDBACK

The fiducial FIRE-2 simulations that we examine do not include self-consistent cosmic ray (CR) injection and transport. However, CRs may be important in regulating the dynamical state of the ISM, and therefore the resultant stellar populations, because in the solar neighbourhood the CR energy density is comparable to the gas thermal and turbulent energy densities (for example Boulares & Cox 1990).

Here, we explore the effects of CR feedback from stars on the dynamical state of our MW-mass discs. One proposed way to form kinematically colder discs in cosmological simulations is the inclusion of previously neglected feedback channels, including CRs, because they can provide a spatially and temporally smoother form of feedback and potentially launch cooler galactic winds. Previous work using FIRE-2 simulations found that the inclusion of magnetic fields, anisotropic conduction, and viscosity has no appreciable effects on any bulk galaxy properties, but that the inclusion of CRs, within reasonable uncertainty for the treatment of the CR diffusion coefficient, significantly can alter galaxy properties (Su et al. 2017; Hopkins et al. 2020). FIRE-2 simulations with CRs also include anisotropic CR transport with stream and advection/diffusion (with constant parallel diffusivity, $\kappa = 3 \times 10^{29} \text{ cm}^2 \text{ s}^{-1}$), CR cooling (hadronic and Compton, adiabatic, and streaming losses), CR injection in supernova shocks, and CR-gas coupling (see Chan et al. 2019, for details of the numerical implementation of CRs in these simulations).

Hopkins et al. (2020) showed in the FIRE-2 simulations that the primary effect of CR feedback is to prevent the accretion of halo gas into the galaxy, leading to lower star formation and stellar mass by $\approx 2-4 \times$ in MW-mass galaxies at $z \lesssim 2$. Furthermore, Chan et al. (2021) found that the inclusion of CRs in FIRE-2 leads to lower velocity dispersions in gas, which might suggest dynamically cooler discs, although they did not examine v_ϕ/σ . We show here that the opposite is true for young stars when considering v_ϕ/σ , that is, the inclusion of CRs leads to dynamically hotter stellar discs.

Fig. C1 shows the effect of CRs on present-day stellar kinematics versus age. We include only the isolated galaxies from the Latte suite, which have versions with CR feedback. We also include the lower-mass galaxies m12r and m12z, because our goal is only to examine the differential effects of including CRs. Fig. C1 (top) shows the current median azimuthal velocity, v_ϕ , now, versus stellar age, averaged across the same galaxies in the CR simulations and the fiducial simulations. At all ages, galaxies with CR feedback have lower v_ϕ , with the largest absolute difference for the youngest stars, which arises primarily because of their lower stellar masses.

Fig. C1 (middle) shows the present vertical velocity dispersion, σ_z , versus age. The CR simulations have lower σ_z at all ages, except for the youngest stars, which may arise from the effects of CRs on the ISM and halo gas occurring over long time-scales. We also examined σ_R , σ_ϕ , and σ_{tot} and found qualitatively similar results. Overall, the inclusion of CR feedback leads to lower velocity dispersions in stars, consistent with analyses of gas in these simulations (Hopkins et al. 2020; Chan et al. 2021).

Fig. C1 (bottom) shows the sample-averaged v_ϕ/σ_z versus age. For most ages, stars in the CR simulations have slightly less rotational support than in the simulations without CR feedback; the youngest stars show the most significant reduction at ≈ 30 percent. Thus, including CR feedback leads to fractionally dynamically hotter galaxies, despite their lower absolute velocity dispersions. However, Fig. C2 shows that CR feedback leads to galaxies with $\approx 3 \times$ lower stellar mass at present, as Hopkins et al. (2020) showed.

Fig. C2 (top row) shows v_ϕ/σ_z for all stars versus stellar mass, while Fig. C2 (bottom row) shows the same but for young stars

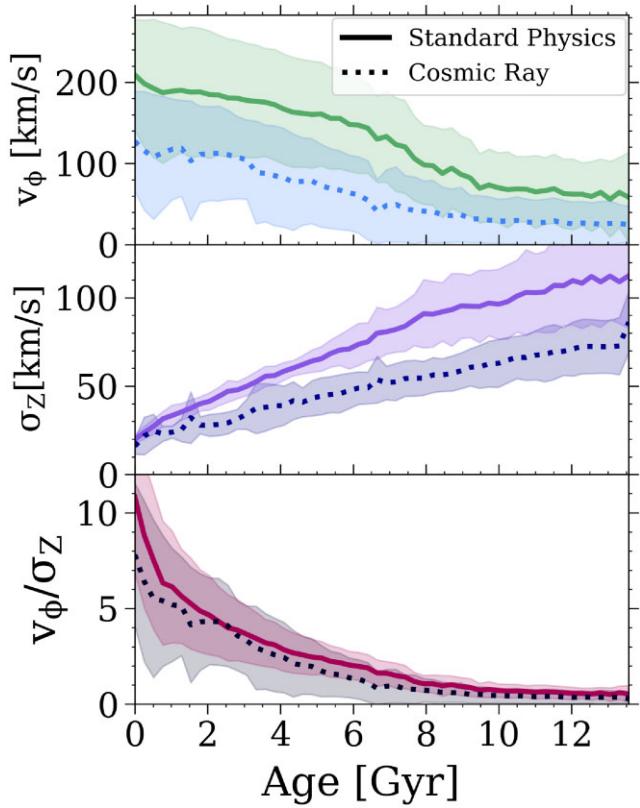


Figure C1. Effects of CR feedback on disc dynamical history: current stellar kinematics versus age, averaged across 8 isolated MW-mass galaxies simulated with fiducial FIRE-2 physics (solid) and additionally including CR injection and transport from stars (dotted), assuming a constant effective diffusion coefficient $\kappa = 3 \times 10^{29} \text{ cm}^2 \text{ s}^{-1}$. The shaded regions show the galaxy-to-galaxy standard deviation. *Top:* Median azimuthal velocity, v_ϕ . The addition of CR feedback reduces overall star formation, leading to lower-mass galaxies with lower velocities at all ages. *Middle:* Vertical velocity dispersion, σ_z . Given their lower stellar mass, the CR simulations have lower velocity dispersions, though the CR simulations converge to the fiducial simulations for the youngest stars. *Bottom:* Median v_ϕ/σ_z , as a measure of rotational support and ‘disciness’. Given the reduction of both v_ϕ and σ_z in the CR simulations, v_ϕ/σ_z is similar at nearly all ages for the fiducial and CR simulations. For the youngest stars, the CR simulations are fractionally dynamically hotter and less rotationally supported. We find similar trends for the other velocity components and the total dispersion.

(age < 250 Myr). We show our six fiducial LG-like galaxies (as black stars) to clarify our fiducial sample’s mass-relation across this entire mass range, given that none of our fiducial isolated simulations have masses within $2 - 4 \times 10^{10} M_\odot$, the range containing half of their CR counterparts. That said, we advise against making direct comparisons between any individual fiducial LG-like simulations and CR isolated simulations, because each galaxy has its own unique history and our fiducial LG-like simulations generally have slightly colder kinematics than our isolated simulations. We emphasize that Fig. C2 shows the same trends with stellar mass as for our fiducial simulations without CR feedback in Fig. 7. Thus, the inclusion of CR feedback primarily acts to reduce a galaxy’s stellar mass and correspondingly shifts its v_ϕ and v_ϕ/σ_z along the existing relation. This leads to galaxies that are slightly less discy, with slightly lower v_ϕ/σ_z . We conclude that the inclusion of CR feedback, at least as implemented in FIRE-2 and assuming a constant diffusion coefficient, does not lead to dynamically colder discs.

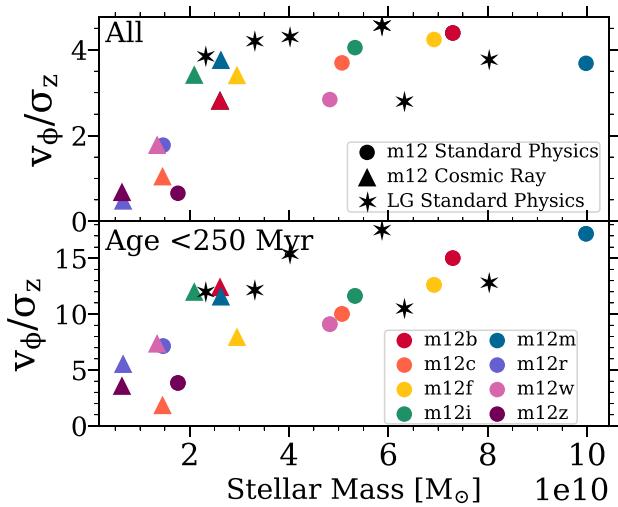


Figure C2. Effects of CR feedback on current stellar v_ϕ/σ_Z versus stellar mass, across eight isolated MW-mass galaxies simulated with fiducial FIRE-2 physics (circles) and additionally with CR injection and transport (triangles). Stars show LG-like galaxies simulated with fiducial physics; although these galaxies lack CR counterparts, we include them in order to illustrate our fiducial sample's v_ϕ/σ_Z -stellar mass relation. CR feedback leads to galaxies with ≈ 3 times lower stellar mass. *Top row:* v_ϕ/σ_Z for stars of all ages. The CR simulations have lower σ_Z but also lower stellar masses, leading them to land on a similar relation in v_ϕ/σ_Z . *Bottom row:* Same but for young stars with ages < 250 Myr. CR simulations have slightly lower v_ϕ/σ_Z for seven of the eight galaxies, though again, they remain on a similar scaling relation. Adding CR feedback does not lead to dynamically colder discs; it primarily reduces the stellar mass formed and moves galaxies along a similar relation.

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