

Lattice-based equation of state with 3D ising critical point

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Abstract.

The BEST Collaboration equation of state combining lattice data with the 3D Ising critical point encounters limitations due to the truncated Taylor expansion up to $\frac{\mu_B}{T} \sim 2.5$. This truncation consequently restricts its applicability at high densities. Through a resummation scheme, the lattice results have been extended to $\frac{\mu_B}{T} = 3.5$. In this article, we amalgamate these ideas with the 3D-Ising model, yielding a family of equations of state valid up to $\mu_B = 700\text{MeV}$ with the correct critical behavior. Our equations of state feature tunable parameters, providing a stable and causal framework-a crucial tool for hydrodynamics simulations.

1 Introduction

The primary objective of the Beam Energy Scan experiment at the Relativistic Heavy-Ion Collider (RHIC) is to map the Quantum Chromodynamics (QCD) phase diagram [1, 2]. Currently, researchers are in the process of analyzing the experimental data [3]. On the theory side, at vanishing chemical potential, high-precision lattice simulations indicate that the transition from a hadron gas to a fluid of strongly interacting quarks and gluons is a smooth crossover [4, 5]. Effective QCD-based models suggest that this crossover eventually transforms into a line of first-order transitions, featuring a critical point at finite density, the precise location of which is still unknown. Unfortunately, lattice simulations at finite baryon density are challenged by the sign problem.

Extrapolation methods that utilize lattice results at zero chemical potential are thus still needed for the equation of state at finite density. Taylor expansion is widely used and is

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effective for low-density physics, approximately $\mu_B/T < 3$; [6–8]. However, its applicability is limited, making it difficult to model high-density critical phenomena.

The BEST Collaboration constructed a parametric family of possible equations of state which, at zero μ_B , agree with the lattice Taylor expansion data, while featuring a critical point at chosen T_C and μ_{BC} [9]. However, the construction appears to be limited to $\mu_B \lesssim 450$ MeV due to unphysical features ("wiggles") induced by extrapolation of Taylor expansion to higher μ_B . On the other hand, the range of $\mu_B > 450$ MeV is needed to describe hydrodynamics in the regime relevant for RHIC Beam Energy Scan.

The Wuppertal-Budapest collaboration developed a novel re-summation expansion scheme [10] that exhibits smooth behavior at high density, producing no visible artifacts up to $\mu_B/T \sim 3.5$, by introducing a chemical potential-dependent transition temperature. This scheme simply relates the baryon density $\chi_1^B(T, \mu_B) \equiv n_B(T, \mu_B)$ with the baryon susceptibility $\chi_2^B(T, \mu_B = 0)$ at zero chemical potential:

$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T', 0), \quad \text{where} \quad \chi_n^B(T) = \frac{\partial^n P/T^4}{\partial (\mu_B/T)^n}, \quad (1)$$

via an effective temperature

$$T'(T, \mu_B) = T \left(1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T} \right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T} \right)^4 + \dots \right). \quad (2)$$

This scheme has inherent advantages, because $\kappa_2^{BB}(T)$ does not introduce wiggles since it is constant near the transition, and κ_4^{BB} is consistent with zero. The remaining task is to introduce a critical point in this equation of state.

1.1 Methodology

1.1.1 Mapping Ising to QCD

Near the critical point, the divergence of the correlation length results in a system whose behavior is determined largely by (global) symmetries, and not by microscopic degrees of freedom. Systems with shared symmetries exhibit similar critical behavior. This fact has been used in Ref. [9] to map the universal critical behavior from the 3D Ising model to the vicinity of the QCD critical point. To combine the approach of Ref. [9] with the Taylor resummation approach of Ref. [10], in this work we employ a two-step mapping with the intermediate step using (T', μ_B) coordinates, where the line of first-order phase transition remains constant at $T' = T_0$. Upon substitution of T'_{lattice} for T' in the subsequent step, the critical line follows the curvature in the QCD coordinates, as illustrated in Fig. 1.

The mapping depicted in Fig. 1 involves free parameters: w' , ρ' , and α'_{12} which can easily be related to the BEST parameters w , ρ , and α_1 and α_2 in [9], as detailed in [11]. Being inherently even in μ_B , this mapping ensures baryon-antibaryon symmetry. In Section 1.1.2, we explain how to merge the critical behavior with the lattice data.

1.1.2 Embedding Ising Critical point into lattice equation of state

It is crucial to interpret Eq. (1) as a definition of T' . The physics at finite density should be incorporated into T' , including the singularity at the critical point. To ensure smooth behavior, we initially merge the lattice data $\kappa_{2,\text{lattice}}^{BB}$ and $\chi_{2,\text{lattice}}^{BB}$ with HRG data (based on the PDG 2021+ particles list [12]), and parameterize in [11] to cover the temperature range relevant for hydrodynamics. Using Eq. (1), we construct n_B^{full} as follows:

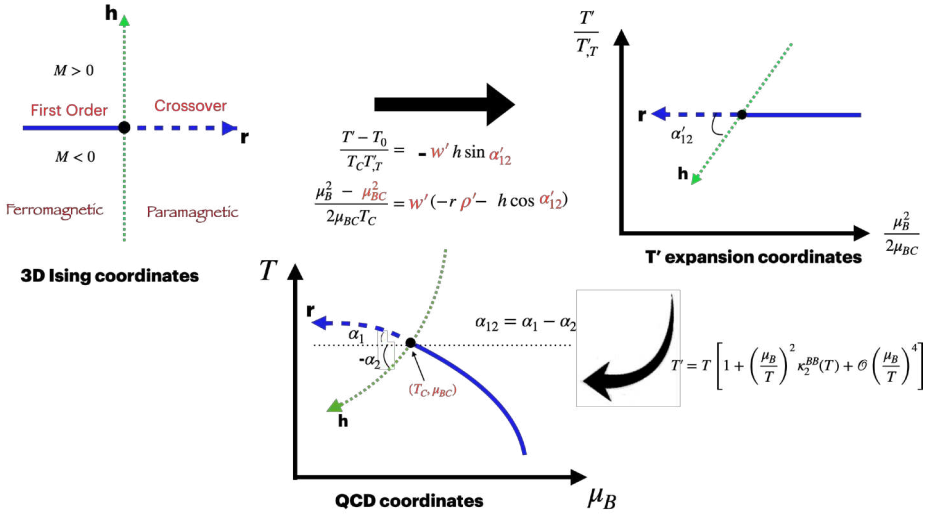


Figure 1. The top left plot represents the 3D Ising model with a critical point located at $(r = 0, h = 0)$. The top right plot displays the alternative T' -expansion scheme coordinates, with a critical point at $(T' = T_0, \mu_B = \mu_{BC})$. Finally, the bottom plot corresponds to the QCD coordinates, featuring a critical point located at (T_C, μ_{BC}) . The tunable (yet unknown) parameters are in red.

$$\frac{n_B(T, \mu_B)}{T^3} = \frac{\chi_1^B(T, \mu_B)}{T^3} = \left(\frac{\mu_B}{T}\right) \chi_{2, \text{lattice}}^B(T'_{\text{full}}(T, \mu_B), 0) \quad (3)$$

$$T'_{\text{full}}(T, \mu_B) = \underbrace{T'_{\text{lattice}}(T, \mu_B)}_{\text{lowest orders in } (\mu_B/T)} + \underbrace{T'_{\text{crit}}(T, \mu_B) - \text{Taylor}[T'_{\text{crit}}(T, \mu_B)]}_{\text{higher order in } (\mu_B/T)}. \quad (4)$$

T'_{lattice} is kept to $O((\mu_B/T)^2)$, and the critical part contributes to higher orders. T'_{crit} is:

$$T'_{\text{crit}}(T, \mu_B) \approx \left(\frac{\partial \chi_{2, \text{lattice}}^B(T, 0)}{\partial T} \Big|_{T_0} \right)^{-1} \frac{n_B^{\text{crit}}(T, \mu_B)}{(\mu_B/T)} \quad (5)$$

where n_B^{crit} is the singular contribution to n_B obtained by mapping the 3D Ising model to QCD coordinates. By substituting Eq. (5) into Eq. (4), we reconstruct the full baryon density using Eq. (3).

1.2 Results

In this contribution we present an enhancement in QCD thermodynamics with respect to the BEST collaboration result [9], as illustrated in Fig. 2. This improvement extends the analysis to a higher chemical potential, reaching up to $\mu_B = 700$ MeV. The transition line in the figure exhibits the expected curvature in QCD, thereby reducing the number of free parameters. Specifically, we fix α_1 to be the slope of the curve in Fig. 1. Furthermore, our parametric equations of state remain valid for smaller mapping angles α_{12} , corresponding to physical (small) quark masses, according to [13].

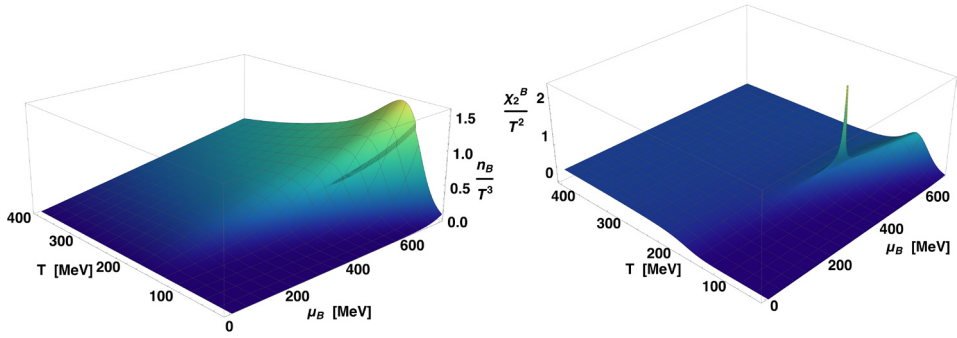


Figure 2. Left panel: baryon density, with a discontinuity marking a first-order transition. Right panel: 2nd order baryon susceptibility with a divergence indicating the critical point. Both have been obtained using parameters $\mu_{BC} = 500\text{MeV}$, $T_C = 116.49\text{MeV}$, $w = 15.0$, $\rho = 0.3$, $\alpha_1 = 11.19^\circ$ and $\alpha_{12} = \alpha_1$.

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