

Composition of Earth's initial atmosphere and fate of accreted volatiles set by core formation and magma ocean redox evolution

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Abstract

The origins of volatile depletion in the Bulk Silicate Earth (BSE) remain controversial due to the numerous processes during accretion and differentiation that can alter volatile abundances. Here, we integrate realistic impact histories from *N*-body simulations to determine the distribution of H, C, and N between Earth's core, mantle, and atmosphere, including loss to space, during the giant impact stage of accretion. Estimates of the BSE's volatile abundances can be matched if accreting materials were depleted ~1–2 orders of magnitude relative to CI chondrite volatile abundances, indicating early volatile loss on planetesimals and embryos prior

to terrestrial accretion. We find the proto-Earth's atmosphere to be dominated by CO (>97%), which leads to the preferential erosion of C relative to H and N by subsequent impacts. Erosion can remove significant C (~18–45% of accreted C), and potentially substantial N as well (~12–23%), if N is not extremely siderophile at high pressures and temperatures. Out of the accreted volatiles, the majority are sequestered in the core, which is the largest reservoir of H (~57–75%), C (~54%–98%), and N (~69%–99.9%) in the bulk Earth after accretion. Our results show that the BSE's volatile depletion relative to chondrites is dominated by loss on precursor bodies (>92%), followed by core formation (~77–99% of remaining <8%), while atmosphere erosion accounts for the rest. Formation scenarios that source more material from the outer Solar System, such as Classical and Grand Tack, result in Earth analogs that have higher total volatile concentrations, lower proportions of volatiles stored in Earth analog cores, and have lost a larger fraction of their accreted C due to atmosphere erosion. The BSE's superchondritic C/N ratio can be matched or exceeded, even if Earth's building blocks have CI chondritic C/N, if N is more strongly siderophile at high pressures and temperatures. If N does not become sufficiently siderophile at high pressures and temperatures, however, Earth's building blocks must have accreted with superchondritic C/N ratios.

1. Introduction

Volatile elements, such as H, C, and N, are essential for Earth's habitability. The distribution of these elements during Earth's formation sets their availability at the surface at the beginning of Earth's evolution. Therefore, understanding the processes that affect the abundances of volatile elements at Earth's surface during its formation are crucial to understanding how Earth came to be habitable. The Bulk Silicate Earth (BSE) is depleted in

volatile elements relative to chondrites (Broadley et al., 2022; Halliday, 2013; Marty and Zimmermann, 1999), but the source of this volatile depletion remains unclear, as multiple processes can alter volatile abundances during Earth's formation. Prior to the giant impact stage of Earth's accretion, condensation from the protoplanetary disk and differentiation and growth of planetesimals and embryos shape the volatile inventories of Earth's building blocks (Hirschmann et al., 2021; Li et al., 2021; Lichtenberg et al., 2019; Newcombe et al., 2023; Young et al., 2023). Subsequently, core formation, magma ocean degassing, and impact erosion by planetesimal and embryo impacts could all have altered the abundances of volatiles in the silicate and surface inventories of rocky planets (Abe and Matsui, 1986; Chen and Jacobson, 2022; Genda and Abe, 2005; Hirschmann, 2016; Sossi et al., 2022; Tucker and Mukhopadhyay, 2014). Even in some pebble accretion scenarios, Earth is formed from two large embryos that experienced their own volatile histories before colliding with one another (Johansen et al., 2021).

Estimates of the BSE's H, C, and N budgets indicate that these elements are depleted by several orders of magnitude relative to carbonaceous chondrites (Hirschmann, 2018; Marty, 2012; Marty et al., 2020; Marty and Zimmermann, 1999). In particular, the BSE is characteristically more depleted in N relative to H and C. It is possible that Earth formed from planetesimals and embryos that already had superchondritic C/N ratios (Grewal and Asimow, 2023; Li et al., 2023). The BSE's C/N could also have been modified by core formation, magma ocean outgassing, and/or impact erosion (Bergin et al., 2015; Hirschmann, 2016). The effects of core formation and magma ocean outgassing on the distribution of volatile elements during Earth's accretion remains an open question, however, as previous models neglected the effects of evolving conditions throughout Earth's formation (Chen and Jacobson, 2022; Sakuraba et al., 2021). For example, pressures and temperatures of core formation are expected to increase and

magma oceans are expected to become more oxidized throughout accretion (Rubie et al., 2015). The core may be Earth's largest reservoir of volatiles (Fischer et al., 2020; Okuchi, 1997), so quantifying its volatile budget is key to understanding the effects of core formation on Earth's volatile element ratios. In addition, the redox evolution of magma oceans, which is set by core-forming metal, determines the speciation of H and C in the overlying atmosphere (Deng et al., 2020; Gaillard et al., 2022; Hirschmann, 2012; Zhang et al., 2017). Various atmospheric species have unique solubilities in silicate liquids, resulting in different distributions of volatile elements between magma oceans and atmospheres (Bower et al., 2022; Hirschmann, 2012). As a result, the susceptibility of each volatile element to erosion by subsequent impacts is linked to its distribution between core, mantle, and atmosphere, which is determined by the conditions of core formation and magma ocean redox states.

Recent geochemical models of core formation include evolving pressures, temperatures, and oxygen fugacities (fO_2) of metal–silicate equilibration throughout accretion (e.g., Fischer et al., 2017; Gu et al., 2023; Rubie et al., 2015), but incorporating volatile elements remains a challenge because these models do not include the overlying atmosphere. In addition, the timing and quantity of volatile delivery may be heterogeneous, resulting in selective sequestration of volatile elements in the core under certain conditions (Schönbächler et al., 2010; Sossi et al., 2022). The degree to which impact erosion can modify volatile abundances is a result of the proportion of volatile elements in the atmosphere available to be eroded. Estimates of the initial volatile concentrations in accreting material and resulting core concentrations therefore may be biased (Blanchard et al., 2022; Fischer et al., 2020; Jackson et al., 2021). Understanding the timing of volatile accretion and the distribution of volatile elements between core, mantle,

atmosphere, and loss to space will not only help to constrain the composition of the Earth's core, but also the extent of volatile depletion on Earth's precursor materials.

Here, we model planetary growth, core formation, magma ocean degassing, and impact erosion during the giant impact stage of Earth's accretion (Fig. 1). Our accretion histories are based on *N*-body simulations, combined with a core formation model that incorporates the effects of evolving pressures and temperatures on the partitioning of light elements into Earth's core. This allows $f\text{O}_2$ to evolve over the course of accretion, both at the metal–silicate interface, and throughout the magma ocean, linking the speciation and solubility of each volatile element in the atmosphere to the interior (Armstrong et al., 2019; Deng et al., 2020; Kuwahara et al., 2023). We focus on the distribution of H, C, and N within the proto-Earth to constrain the extents to which core formation and impact erosion contributed to the BSE's present-day volatile depletion, therefore setting bounds on the initial volatile inventory of Earth's building blocks.

2. Methods

2.1. Initial conditions

Accretion histories of Earth analogs are taken from 433 published *N*-body simulations spanning Classical, Annulus, Grand Tack, and Early Instability models of Solar System formation (Clement et al., 2021; Jacobson and Morbidelli, 2014; O'Brien et al., 2006; Raymond et al., 2009; Raymond and Izidoro, 2017). Earth analogs are defined as the planet at the end of each simulation closest to 1 Earth mass (in the range 0.8–1.25 Earth masses) and 1 AU (in the range 0.8–1.2 AU), resulting in a total of 85 Earth analogs (Gu et al., 2023). All bodies are assigned initial compositions based on refractory elemental abundances in CI chondrites (Palme and O'Neill, 2014). Volatile abundances and bulk $f\text{O}_2$ of accreting bodies are prescribed based

on their original semi-major axes (Fischer et al., 2017; Rubie et al., 2015). Following previous models, we use different values of fixed $f\text{O}_2$ for bodies from within (IW–2.75) and beyond (IW–1.50) 2 AU, with inner Solar System bodies being more reduced ($f\text{O}_2$ expressed in log units relative to the iron–wüstite buffer). The $f\text{O}_2$ values and $f\text{O}_2$ cutoff at 2 AU are chosen to match Earth’s mantle FeO, as well as be consistent with Mars’ accretion (Brennan et al., 2022; Gu et al., 2023). Likewise, volatile concentrations are assigned in CI chondritic proportions (Vacher et al., 2020; Wasson and Kallemeyn, 1988), but are depleted by a constant factor based on each body’s initial semi-major axis. The model values for each of these parameters are given in Table 1. We assume that material originating from beyond 2 AU is ten times more enriched in volatile elements than material originating from within 2 AU because carbonaceous chondrites have about an order of magnitude more H, C, and N than enstatite chondrites (Table 1 of Sakuraba et al., 2021 and references therein). The sensitivity of our model to these parameters (location of $f\text{O}_2$ cutoff, $f\text{O}_2$ values, chondrite type, and volatile abundances) is explored in Supplementary Text S5. Volatile elements are initially distributed between cores and mantles of planetesimals and embryos (Supplementary Text S1). We do not explicitly model volatile loss on these bodies. Therefore, assigned initial volatile abundances represent bulk volatile abundances at the time of their accretion onto the proto-Earth.

2.2. Atmosphere erosion

The fraction of existing atmosphere eroded by each impact is determined based on the properties of the impactor and atmosphere at the time of the impact. We adopt separate prescriptions for planetesimal and embryo impacts, taking into account the size distribution of planetesimals and differences in erosional properties of aerial bursts and giant, cratering, and

fragmenting impacts (Kegerreis et al., 2020; Schlichting et al., 2015; Shuvalov, 2009; Shuvalov et al., 2014; Sinclair and Wyatt, 2022).

Planetesimals in N -body simulations represent conglomerations of smaller bodies that are assumed to impact simultaneously. A power law size distribution with a differential power law index of $q = 2$ is used to describe the impactor size distribution for each planetesimal such that $\frac{dN}{dD} \propto D^{-q}$, where N is the number of impactors smaller than diameter D (Fig. S1) (Sakuraba et al., 2021). Impactors have minimum (D_{min}) and maximum (D_{max}) diameters of 1 and 1000 km, respectively. Assuming an impactor density of $\rho_i = 3.3 \text{ g/cm}^3$ and that the mass of a planetesimal scales with its diameter as $m_i = \frac{\pi \rho_i D^3}{6}$, the size distribution of planetesimals is converted to a mass distribution. A random sample equivalent to a specified total planetesimal mass, M_{ptsm} , is generated from this distribution:

$$N = \frac{12M_{ptsm}}{\pi \rho_i} (D_{max} D_{min} (D_{max} + D_{min}))^{-1} \quad (1)$$

Atmosphere loss from planetesimal accretion can be induced by cratering if the impactor reaches the surface intact. Otherwise, impactors may fragment or form aerial bursts before reaching the surface. We adopt the prescription of Shuvalov et al. (2014), as modified by Sinclair and Wyatt (2022), to describe fragmenting impacts and aerial bursts and that of Shuvalov (2009) for cratering impacts (Fig. S2). We also include the polar cap limit from Schlichting et al. (2015), which limits the atmosphere loss due to a single impact. Further details of these methods are described in Supplementary Text S2.

Giant impacts associated with embryo collisions erode existing atmospheres by ejecting mass near the impact site, but also by causing a shock wave to propagate through the planet, resulting in global ground motion (Genda and Abe, 2005; Schlichting et al., 2015). Embryo masses are taken directly from N -body simulations. We apply the scaling law of Kegerreis et al.

(2020) to describe the fraction of atmosphere lost (X_{loss}) for each embryo impact, defined as impactors with masses >0.01 Earth masses:

$$X_{loss} = 0.64 \left[\left(\frac{v_i}{v_{esc}} \right)^2 \left(\frac{M_i}{M_i + M_p} \right)^{\frac{1}{2}} \left(\frac{\rho_i}{\rho_p} \right)^{\frac{1}{2}} f_M(b) \right]^{0.65} \quad (2)$$

where $\frac{v_i}{v_{esc}}$ is the impact velocity normalized to the escape velocity, $\frac{M_i}{M_i + M_p}$ is the impactor:total mass ratio, and $\frac{\rho_i}{\rho_p}$ is the impactor:target density ratio. $f_M(b)$ is a fractional interacting mass (Supplementary Text S2). We do not incorporate atmospheric loss induced by thermal expansion of the atmosphere since thermal loss is inefficient for high mean molecular mass envelopes (Biersteker and Schlichting, 2021).

2.3. Volatile distribution between metal, silicate, and atmosphere

Following each impact, it is assumed that a fraction of the impactor's core, the molten portion of the target's mantle, and the entire remaining atmosphere is equilibrated. The fraction of the mantle that is melted by the impact (k_{mantle}) is determined from existing scaling laws (Abramov et al., 2012; Nakajima et al., 2021). For impactors >0.01 Earth masses (embryos), the impactor and target masses, impact velocity, and impact angle of each collision were extracted from the outputs of each N -body simulation and used in the melt-scaling law from Nakajima et al. (2021), which parameterizes outputs from hydrodynamic simulations to determine the volume and geometry of melting in the target's mantle. Impact angles were rounded to the nearest available angle (0° , 30° , 45° , 60° , 90°) with surface entropy set to 1100 J/K/kg, corresponding to surface temperatures of ~ 300 K. Only a fraction of impacting embryo cores participates in equilibration ($k_{core_emb} = 0.3$), since cores of larger bodies are more likely to directly merge with the proto-Earth's core (Marchi et al., 2018; Rubie et al., 2015). Planetesimal cores, on the other

hand, are considered to completely equilibrate (Kendall and Melosh, 2016). These planetesimals were equilibrated with a portion of the mantle equal to five times the impactor’s mantle mass, consistent with estimates of melting using the average impact velocity and averaged over all impact angles (Abramov et al., 2012; Gu et al., 2023). Planetesimal cores are assumed to remain in the mantle and are then re-equilibrated during the next embryo impact (Fleck et al., 2018).

Major elements Fe, Si, Ni, and O are partitioned between the fraction of equilibrating impactor core and the molten portion of the mantle, whereas Mg, Ca, and Al are assumed to be perfectly lithophile. The partitioning parameterization used for each element can be found in Table S1. Equilibration pressures evolve along a predetermined pressure history, corresponding to the evolving average pressure at the base of magma oceans, which is derived from embryo impacts averaged over many *N*-body simulations coupled with a melt-scaling law (Fig. S3) (Gu et al., 2023; Nakajima et al., 2021). This is done to simplify pressure calculations, remove the dependence of pressure on initial embryo masses in different simulations, and avoid the need for an ad hoc parameter describing the fractional depth at which metal–silicate equilibration occurs. At each equilibration pressure, the equilibration temperature is set to that of the chondritic liquidus (Andrault et al., 2011), and temperature is extrapolated upwards through the magma ocean along the liquid adiabat (Thomas and Asimow, 2013). Partitioning of major elements sets the fO_2 of the core–mantle boundary relative to the iron–wüstite buffer, approximated (assuming ideal mixing) as:

$$\Delta IW = 2 \log \left(\frac{X_{FeO}}{X_{Fe}} \right) \quad (3)$$

where X_{FeO} is the mole fraction of FeO in the mantle and X_{Fe} is the mole fraction of Fe in the core. Assuming a constant Fe^{3+}/Fe_{tot} throughout the magma ocean due to vigorous convection, fO_2 is extrapolated upwards along the liquid adiabat (Thomas and Asimow, 2013) of the magma

ocean to determine the fO_2 at the magma ocean–atmosphere interface (Deng et al., 2020; Hirschmann, 2021, 2022). Magma ocean fO_2 evolves with increasing Fe^{3+}/Fe_{tot} as Earth analogs grow and magma ocean depths increase (Deng et al., 2020). Details of these calculations are discussed in Supplementary Text S3. We also test a scenario, based on recent experimental results from Kuwahara et al. (2023), where magma oceans are more oxidized with Fe^{3+}/Fe_{tot} of ~ 0.3 – 0.4 in Supplementary Text S4. However, oxidized magma oceans are inconsistent with the present-day oxidation state of Earth’s mantle (Section 4.1.) (Hirschmann, 2023). Vertically linking fO_2 from the equilibrating impactor core to the atmosphere allows for the self-consistent calculation of volatile distributions of between metal, silicate, and atmosphere.

Volatile elements are partitioned between the equilibrating impactor core, molten mantle, and remaining atmosphere by mass balance, such that:

$$M_i^{tot} = X_i^{met} M^{met} + X_i^{sil} M^{sil} + \sum \frac{4\pi R^2 m_i^{atm}}{g\mu r_i} P_i \quad (4)$$

where M_i^{tot} is the total mass of volatile element i , P_i is the partial pressure of each species of element i in the atmosphere, R is the current radius of the proto-Earth at the time of the impact, g is its associated gravitational acceleration, μ is the mean molecular mass of the atmosphere, m_i^{atm} is the molecular mass of each species of element i in the atmosphere, and r_i is the mass ratio of the volatile species to the element of interest (e.g., $m_{CO}/m_C = 28/12$) (Bower et al., 2019; Sakuraba et al., 2021). X_i^{met} and X_i^{sil} are the concentrations of element i in the equilibrating metal and silicate reservoirs, respectively, which have total masses M^{met} and M^{sil} . The summation for the atmospheric term in Eq. 4 is for H and C, which can exist as multiple species (H_2O or H_2 , and CO_2 or CO , respectively), whereas N_2 is assumed to be the only species for N. The mass balance can be solved given a metal–silicate partition coefficient ($D_i = \frac{X_i^{met}}{X_i^{sil}}$), a solubility law relating the partial pressure of each species in the atmosphere to its concentration in the silicate melt, and

the relative proportions of H₂O/H₂ or CO₂/CO in the atmosphere. Metal–silicate partitioning parameterizations (Fig. S4) (Blanchard et al., 2022; Jackson et al., 2021; Shi et al., 2022; Tagawa et al., 2021) and solubility laws (Dasgupta et al., 2022; Lichtenberg et al., 2021) are described in Supplementary Text S3. We do not account for the interactions between light elements in the metal, as the magnitudes of these interactions at high pressures and temperatures are uncertain and likely to be small for these X_i^{met} . We use two parameterizations for the metal–silicate partitioning of N to bound its partitioning behavior, where “J21” refers to the more siderophile parameterization of Jackson et al. (2021) and “S22” refers to the less siderophile parameterization of Shi et al. (2022). For C, we use the parameterization from Blanchard et al. (2022), which predicts a higher siderophility of C at high pressures and temperatures than the parameterization from Fischer et al. (2020). C being more siderophile diminishes its abundance in the mantle and atmosphere, so that the amount of C lost to space in our model represents a lower bound (Supplementary Data). At the surface of the magma ocean, chemical equilibrium distributes H and C budgets into H₂/H₂O and CO/CO₂ given the surface fO_2 and temperature. The two reactions involved are:



The equilibrium constants for these two reactions are computed following the methods of Robie and Hemingway (1995), with the thermodynamic quantities from Linstrom (1998).

We define a set of plausible reference parameters that reproduce mantle FeO and trace element compositions (Gu et al., 2023; McDonough and Sun, 1995; Palme and O’Neill, 2014) to constrain $k_{core_emb} = 0.3$, inner $fO_2 = IW-2.75$, outer $fO_2 = IW-1.5$, fO_2 cutoff = 2 AU, $P_{max} = 60$ GPa, and k_{mantle} set by melt-scaling laws, which are the parameters to which core formation

models are most sensitive (Fischer et al., 2017). The sensitivity of our model to individual parameters is discussed in Supplementary Text S5. Initial volatile concentrations are then varied to provide reasonable matches to estimates for the BSE’s volatile concentrations. In our reference case, inner Solar System materials have 2% and outer Solar System materials have 20% CI chondrite concentrations of H, C, and N (Table 1). Our prescription of equilibration corresponds to magma oceans that crystallize between impacts such that we can employ scaling laws for mantle melting and atmosphere erosion that assume solid surfaces. This is because the timescale of magma ocean crystallization for atmospheres formed by our model are fleeting relative to the time between impacts (Gu et al., 2023; Lichtenberg et al., 2021; Nikolaou et al., 2019). We assume that equilibration between core, mantle, and atmosphere is shorter than the magma ocean crystallization timescale (Salvador and Samuel, 2023; Thomas et al., 2023) and that the distribution of volatile elements between the mantle and atmosphere is not altered due to magma ocean crystallization between impacts. We run additional models to test the effects of volatile outgassing between impacts in Section 4.2. Our model links realistic impact histories to the self-consistent distribution of volatile elements between the proto-Earth’s core, mantle, atmosphere, and loss to space.

3. Results

3.1. Fate of accreted volatiles

Earth analog mantles naturally become progressively oxidized over the course of accretion due to the incorporation of more Si in the core at higher P – T (with more corresponding oxidation of Fe) and the tendency of Earth analogs to accrete more oxidized bodies from beyond 2 AU at later times (O’Brien et al., 2006), resulting in higher Fe_{tot} and $\text{Fe}^{3+}/\text{Fe}_{\text{tot}}$ (Fig. 2) (Rubie

et al., 2011). Oxidation of the mantle, combined with the increasing depths of magma oceans, generates increasingly oxidized surface fO_2 values that evolve from $\sim IW-3$ to $\sim IW-1$ (Fig. 3a) (Deng et al., 2020; Hirschmann, 2022, 2021). These fO_2 values correspond to Fe^{3+}/Fe_{tot} values of ~ 0.01 and ~ 0.05 (Fig. 2b). Despite becoming more oxidized relative to equilibrium with Fe alloy, surface fO_2 values remain below the iron–wüstite buffer. We note that using the data of Sossi et al. (2020) to determine fO_2 would result in more oxidized surface fO_2 by ~ 1 log unit. We also consider a more extreme oxidized case in Supplementary Text S4, where magma oceans have high Fe^{3+}/Fe_{tot} (Kuwahara et al., 2023), but note that this case is inconsistent with the mantle's redox budget and long-term evolution (Hirschmann, 2023).

The abundance of H, C, and N stored in the core and lost to space is a function of both their solubilities and their siderophilicities at the conditions of core formation. All three elements are predicted to be siderophile, so most of each ends up in the core, with the core making up $\sim 57\text{--}75\%$ and $\sim 43\text{--}75\%$ of accreted H and C, respectively (Fig. 4). These bounds are taken from the smallest 10th percentile and the largest 90th percentile of the distributions of Earth analogs. Storage of N in Earth analog cores is governed by its metal–silicate partition coefficient, which remains uncertain above ~ 25 GPa, but can account for $\sim 55\text{--}98\%$ of accreted N (Jackson et al., 2021; Shi et al., 2022). If N is siderophile enough at high pressures and temperatures, core formation could elevate the BSE's C/N ratio relative to its chondritic precursors such that the modeled C/N ratio matches or exceeds the BSE's observed superchondritic C/N ratio (Fig. 5). The exact C/N ratio in this scenario depends on the solubility of N in magma oceans (Fig. S10). However, if N is not as siderophile at high pressures and temperatures, our model fails to reproduce the BSE's superchondritic C/N.

On the other hand, the availability of each element to be lost by atmospheric erosion depends on their solubilities, where H is more soluble than N and C. The high solubility of H, combined with the reduced surfaces of magma oceans, results in CO-dominated atmospheres throughout accretion, with C making up >97% of atmosphere species (Fig. 3b) (Bower et al., 2022; Sossi et al., 2020). Because of this, C is preferentially eroded in subsequent impacts, with ~18–45% accreted C lost due to atmosphere erosion (Fig. 4). N, on the other hand, can only be eroded to a significant degree (~12–33%) if it is not strongly siderophile at high pressures and temperatures (Fig. 4). The distribution of volatiles and the abundances of volatile elements that end up in Earth analogs are slightly dependent on simulation type (Figs. 4 and 5). Earth analogs comprised of more material originating from beyond 2 AU, such as those from Classical and Grand Tack simulations, have higher total volatile concentrations, but are more oxidized and so have lower fractions of accreted volatiles in their cores and experience more volatile loss due to atmosphere erosion.

3.2. Origin of the BSE's volatile depletion

Matching estimates for the BSE's H and C contents at the end of accretion, including surface reservoirs (Hirschmann, 2018; Marty, 2012; Marty et al., 2020), requires initial building blocks that are volatile-depleted relative to CI chondrites by ~1–2 orders of magnitude, on average. If inner Solar System materials have 2% and outer Solar System materials have 20% CI chondrite values (Table 1), the BSE's H and C concentrations can be matched, on average, using inner Solar System materials with ~0.7–5% CI chondrite concentrations (Fig. S7). This range corresponds to Earth analogs that are made of building blocks, including those originating from beyond 2 AU, containing 0.9%–7.8% CI chondrite concentrations, on average. This means that

despite Earth's core being a major reservoir of H, C, and N, and despite significant C, and potentially N loss due to atmosphere erosion, >92% of the BSE's depletion relative to chondrites must have occurred prior to accretion onto Earth (Hirschmann et al., 2021; Li et al., 2021; Newcombe et al., 2023). Core formation would make up ~99%, ~78%, and ~77–99% of the remaining <8% depletion in H, C, and N relative to Earth's building blocks with the rest being lost due to atmosphere erosion. Even though atmosphere erosion would be insignificant in setting the BSE's volatile depletion relative to chondrites, it can still erode significant portions of accreted C and N, and therefore, still be important in setting the BSE's C/N ratio.

The estimated average abundances of volatiles in Earth's building blocks are not affected by our assumption of the inner-to-outer Solar System ratio of volatile abundances. If we take an extreme case and assume the inner Solar System is completely dry, our model still reproduces the high C/N ratio of the BSE, on average (Supplementary Text S5). In this scenario, volatiles are delivered based on the timing and quantity of outer Solar System material accreted in each simulation, increasing the spread in final H, C, and N abundances of Earth analogs. Matching Earth's H and C abundances still requires initial building blocks to be depleted by ~1–2 orders of magnitude relative to CI chondrites, on average. Even though the timing and quantity of volatile delivery varies between simulations, core formation always elevates the BSE's C/N ratio relative to its chondritic precursors, challenging the delivery of volatiles solely during the late veneer (e.g., Albarède, 2009). In addition, regardless of the assumed inner-to-outer Solar System ratio of volatiles, atmospheres remain dominated by CO, leading to the preferential loss of C due to atmospheric erosion.

3.3. Core composition

The core is generally the largest reservoir of accreted volatiles. When only considering the bulk volatile concentrations remaining in Earth analogs, the core comprises ~57–75%, ~54%–98%, and ~69%–99.9% of the bulk planet’s H, C, and N, respectively. These core concentrations are associated with 0.05–0.09 wt% H, 0.25–0.47 wt% C, and 0.01–0.03 wt% N in Earth analog cores. Allowing initial inner Solar System volatile concentrations to vary within the ranges allowed to match estimates of BSE volatile contents (0.7–5%) corresponds to H, C, and N concentrations in Earth analog cores of 0.02–0.24 wt%, 0.08–1.14 wt%, and 0.01–0.08 wt%, respectively. The upper limits of H and C concentrations in the core can account for ~50% of the core’s density deficit, but using intermediate values results in H and C accounting for only ~15% of the core’s density deficit (Li et al., 2019). Therefore, other light elements, such as Si, O, and S, likely account for the majority of the core’s light element budget and density deficit (Fischer et al., 2017) unless there are mechanisms other than core formation that can store light elements in the core (e.g., Young et al., 2023).

Even though our core compositions are affected by averaging over simulation types, our overall model results regarding the distribution of volatile elements during accretion remain robust. H, C, and N abundances in the core vary by <1 mol% throughout accretion, whereas Si and O increase with increasing pressure and temperature as the planet grows based on their metal–silicate partitioning behaviors (Fig. 6). Our reference model yields core concentrations of $2.4^{+0.6}_{-0.5}$ wt% Si and $0.6^{+0.2}_{-0.1}$ wt% O. These abundances are relatively low compared to the ~5–10% concentration of light elements required to match the core’s density deficit. However, Si and O are much more sensitive than H, C, and N to the pressure, temperature, and fO_2 of metal–silicate equilibration (Fischer et al., 2017; Gu et al., 2023). We also do not include the mutual interactions of Si and O in the metal (Fischer et al., 2015). Since we are averaging pressures of

equilibration and the bulk fO_2 of accreting bodies over different simulation types, the core concentrations of Si and O could vary significantly if these variables are treated independently for each simulation type. For example, simulations that source more material from the outer Solar System would require more initial bodies to be more reduced, which would promote the incorporation of Si in the core. In these simulations, decreasing the inner fO_2 to IW-3.5 increases Si concentrations in the core to $5.4^{+4.7}_{-1.6}$ wt%, but does not significantly affect volatile element concentrations in the core (Supplementary Data), so there are model conditions that can account for whatever portion of the core's density deficit is not explained by H, C, and N through incorporation of other light elements in the core. This initial inner fO_2 is consistent with the mantle's FeO content in some Grand Tack and Classical simulations which source relatively larger fractions of material from beyond 2 AU (Fischer et al., 2017; Gu et al., 2023; Rubie et al., 2015).

4. Discussion

4.1. Magma ocean redox

Magma oceans become more oxidized throughout accretion due to core formation (Rubie et al., 2011). However, magma ocean surfaces remain below the iron-wüstite buffer, resulting in CO-rich atmospheres throughout accretion. Our results rely on the abundance of Fe^{3+} in equilibrium with metal alloy from the molecular dynamics study of Deng et al. (2020), which predicts $Fe^{3+}/Fe_{tot} < 0.1$ (Fig. 3). These results conflict with experimental studies that indicate that Fe^{3+}/Fe_{tot} could increase dramatically with pressure, reaching values >0.35 (Armstrong et al., 2019; Kuwahara et al., 2023). However, subsequent crystallization and evolution of Earth's mantle must be consistent with basalt source regions that match the present-day upper mantle

$\text{Fe}^{3+}/\text{Fe}_{\text{tot}}$ of 0.04 ± 0.02 (Hirschmann, 2022). The effects of magma ocean solidification should be limited due to the relative compatibility of Fe^{3+} in bridgmanite and by homogenization of the mantle by convection (Kuwahara and Nakada, 2023). Iron disproportionation and hydrogen escape could both oxidize the mantle over time (Catling et al., 2001; Frost et al., 2004; Wordsworth et al., 2018; Zahnle et al., 2013). In contrast, late accretion may deliver metallic iron that reduces the mantle (Itcovitz et al., 2022; Zahnle et al., 2020), though it is unlikely that this mechanism is sufficient to reduce the mantle from the very oxidized values proposed by Kuwahara et al. (2023) (Hirschmann, 2023). We therefore favor the reduced case, but note that oxidized magma oceans would promote comparatively greater H and N loss relative (Supplementary Text S4). Future work on the redox evolution of magma oceans and Earth's mantle over time will be instrumental in constraining the distribution of volatile elements during Earth's formation.

4.2. Magma ocean crystallization

The distribution of volatiles between mantle and atmosphere after magma ocean crystallization determines the subsequent evolution and cycling between the two reservoirs. At the end of accretion, but before magma ocean crystallization, $<<1\%$ H, $\sim 92\%$ C, and $37\text{--}53\%$ N of BSE (mantle + atmosphere) H, C, and N are in the atmosphere. These values contrast with the $\sim 57\%$ H, $\sim 19\%$ C, and $\sim 58\%$ N estimated to be on the surface and in the atmosphere today (Hirschmann, 2018). Compared to C, H and N are inefficiently recycled between the mantle and the surface over geologic timescales, suggesting that Earth's present-day surface H and N reservoirs were established early, shortly after magma ocean solidification (Hirschmann, 2018).

During crystallization of magma oceans, oxidation of the residual liquid could promote oxidation of atmospheres in equilibrium with the magma ocean, although the extent to which the residual liquid is oxidized depends on the partitioning of Fe^{3+} into bridgmanite (Boujibar et al., 2016; Kuwahara and Nakada, 2023). Oxidation and outgassing of some of the magma ocean's N could provide N from the mantle to the surface without requiring significant subsequent interior-atmosphere N exchange. For H, if magma ocean crystallization releases a maximum of 30% H (Bower et al., 2022; Miyazaki and Korenaga, 2022), outgassing of 1 ocean mass would require a total BSE H budget of >3.3 ocean masses. Smaller total H would require additional mantle degassing early in Earth's history (Médard and Grove, 2006) or late delivery of H that contributes to Earth's surface H (e.g., Wang and Becker, 2013). The strong temperature dependence of C solubility in silicate melts would lead to precipitation of diamond or graphite within the cooling magma ocean, which could incorporate C that was originally in the atmosphere into the mantle (Armstrong et al., 2019; Hirschmann, 2012; Keppler and Golabek, 2019; Sossi et al., 2020). Future work on the distribution of volatile elements during magma ocean crystallization will clarify the distribution of H, C, and N between Earth's mantle and atmosphere in the early Earth.

4.3. Additional model limitations

Our model provides constraints on the distributions of H, C, and N during Earth's formation by combining astrophysical *N*-body simulations with interior-atmosphere chemistry and atmosphere erosion. However, there are limitations within the scaling laws and parameterizations we apply that link these different processes together. Below, we outline

additional assumptions that could affect our model results. Outputs of all model runs are provided in Supplementary Data.

We consider variations in individual parameters, including k_{core_emb} , impactor fO_2 , the pressures and temperatures of metal–silicate equilibration, inner:outer Solar System volatile concentration ratio, and initial H:C:N volatile ratios (Supplementary Text S5). Of the free parameters in our model, k_{core_emb} is the most influential. Even though the abundances of H, C, and N in Earth analogs remain consistent with present-day BSE estimates for all values of k_{core_emb} (Fig. S8), complete equilibration ($k_{core_emb} = 1$) decreases the fraction of C in Earth analog cores from ~63% to ~34% and increases C loss due to atmosphere erosion from ~26% to ~47% (Fig. S8). This is because increasing k_{core_emb} prevents the erosion of C by sequestering more C in the core. Increasing k_{core_emb} also increases the oxidizing effect of dissolving Si into the core, therefore oxidizing Earth analog mantles. Altering k_{core_emb} independently is consequently inconsistent with Earth’s mantle FeO and trace element (Ni, Co, V, Cr, etc.) concentrations because it increases or decreases the abundance of metal that can sequester siderophile elements (Fischer et al., 2017). While the set of parameters we choose is non-unique, we conclude that adjusting parameters to match Earth’s mantle composition will likely yield similar distributions of H, C, and N. For example, for a higher value of k_{core_emb} , decreasing the fO_2 of initial bodies will decrease C loss, negating the effects of changing k_{core_emb} .

We also vary the initial concentrations of H, C, and N at the beginning of N -body simulations. Changing only initial concentrations while keeping the inner:outer Solar System ratio of initial H, C, and N constant has very little effect on the proportion of H, C, and N in each reservoir, but affects the BSE abundances of each element (Fig. S8). Only C loss is noticeably affected, where using the smallest initial concentration consistent with the BSE’s C content

increases the proportion of C lost due to atmosphere erosion by ~5%. The effects of varying other parameters individually are discussed in Supplementary Text S5.

Magma oceans are assumed to crystallize before subsequent embryo impacts and are assumed to do so in our model (e.g., Nikolaou et al., 2019). However, we do not explicitly account for compositional effects of magma ocean crystallization, instead, holding mantle and atmosphere abundances of H, C, and N constant between impacts. Magma ocean crystallization could influence the interior–atmosphere distribution of H, C, and N by outgassing volatiles from the mantle to the atmosphere, potentially increasing the availability of each element to be eroded by subsequent impacts (Fig. S9). We have run additional models to assess the effects of magma ocean outgassing, in which a fixed fraction of the mantle’s volatiles is outgassed and added to the atmosphere between embryo impacts. The distribution of H is most affected by outgassing between impacts because of the high solubility of H in magma oceans, increasing the fraction of H in the core from ~69% to ~91% when all volatiles are outgassed. However, recent studies of magma ocean crystallization that incorporate the trapping of volatiles in interstitial melts suggest that H could be efficiently retained in crystallized mantles (<30% outgassed), limiting the effects of outgassing on the distribution of H (Bower et al., 2022; Hier-Majumder and Hirschmann, 2017; Miyazaki and Korenaga, 2022).

It is possible that H₂/He atmospheres existed on initial planetary embryos, but these low mean molecular mass atmospheres are efficiently lost following impacts (Biersteker and Schlichting, 2021). The reaction of nebular H₂ with silicate or iron in embryos could increase the abundance of H relative to C and N and would result in non-chondritic H abundances (Horn et al., 2023; Young et al., 2023). However, the efficient storage of H in Earth analog interiors means that this additional H would most likely be stored in Earth’s core, increasing the

abundances of H in the core proposed in Section 3.3. H escape could alter the redox evolution of Earth analogs and further oxidize magma oceans (Catling et al., 2001; Wordsworth et al., 2018; Zahnle et al., 2013), but we do not consider hydrodynamic escape of volatile elements, as H escape is inefficient in high mean molecular mass and H-poor atmospheres (Biersteker and Schlichting, 2019).

There are also assumptions involved in our prescription of atmosphere erosion. Altering the minimum and maximum impactor size or increasing the differential power law index will change the impactor size distribution and therefore slightly modify the efficiency of planetesimals in eroding existing atmospheres (Supplementary Data). The scaling law for embryo erosion from Kegerreis et al. (2020) is based on loss of hydrogen–helium atmospheres. Erosion of higher mean molecular weight atmospheres would likely be less efficient, making our loss estimates upper bounds. Furthermore, atmosphere loss may be more efficient early in Earth’s accretion when the proto-Earth had a smaller escape parameter (Sossi et al., 2022). The BSE’s present-day volatile budget is therefore likely more representative of later-accreted material, as volatile loss was less efficient later in Earth’s accretion, in addition to the more efficient sequestration of earlier-accreted volatiles in Earth’s core (Dauphas, 2017).

5. Conclusions

We explore the effects of core formation and magma ocean redox on the distribution of H, C, and N during Earth’s formation to determine the evolution of Earth’s earliest atmosphere. Our model predicts that Earth had a CO-dominated atmosphere throughout accretion. This C-dominated atmosphere results in the preferential erosion of C (~18–45% is lost) relative to H and N, but significant N (~12–23%) can be lost if N is not strongly siderophile at high pressures and

temperatures. Matching the BSE's C and H budgets requires volatile abundances in accreting materials to be depleted by ~ 1 – 2 orders of magnitude relative CI chondrite concentrations, on average, consistent with early volatile depletion on planetesimals and embryos. Volatile depletion on precursors is responsible for $>92\%$ of the BSE's volatile depletion relative to chondrites. Core formation and atmosphere erosion make up the remaining $<8\%$ of the BSE's volatile depletion but are both still crucial in setting the BSE's C/N ratio. The core is the major reservoir of H, C, and N in the Earth, making up ~ 57 – 75% , $\sim 54\%$ – 98% , and $\sim 69\%$ – 99.9% of the bulk Earth's budgets of these elements, respectively. Simulations that form Earth analogs with larger contributions of outer Solar System material, such as Classical and Grand Tack simulations, result in Earth analogs with higher total volatile concentrations and that lose a larger fraction of their accreted volatiles due to atmosphere erosion. Core formation alone can match or exceed the BSE's superchondritic C/N ratio if N is sufficiently siderophile at higher pressures and temperatures. If N is not sufficiently siderophile at high pressures and temperatures, however, the BSE's processes on precursor bodies must be responsible for elevating the BSE's C/N ratio. Future work on the redox states of magma oceans and the metal–silicate partitioning behavior of N will help to elucidate the fate of Earth's accreted volatile elements.

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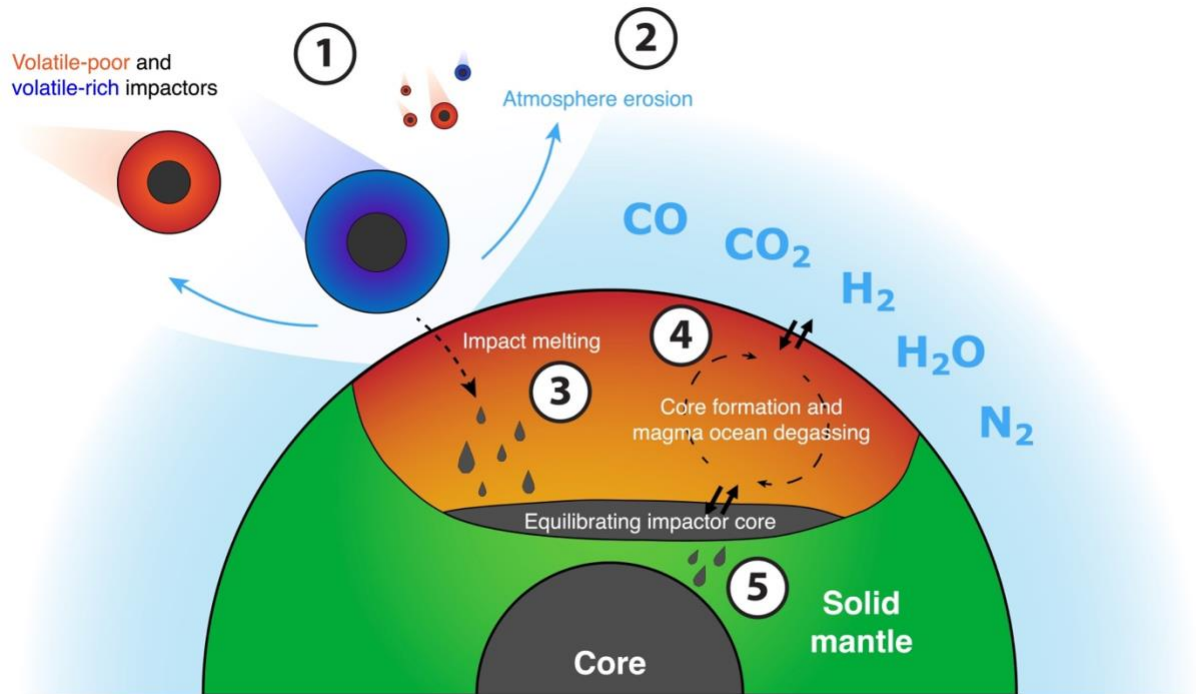


Figure 1. Schematic representation of model steps for each impactor. 1) Volatile elements are assigned to impactors based on their initial semi-major axes, which are taken from N -body simulations. 2) Each impactor erodes a portion of the existing atmosphere and 3) melts a fraction of the proto-Earth's mantle. 4) Volatile elements are distributed between a portion of the impactor's core, the molten mantle, and the remaining atmosphere based on metal–silicate partitioning coefficients, solubilities of atmosphere species in the magma, and the fO_2 at the surface of the magma ocean. 5) The equilibrated core material sinks and merges with the proto-Earth's core.

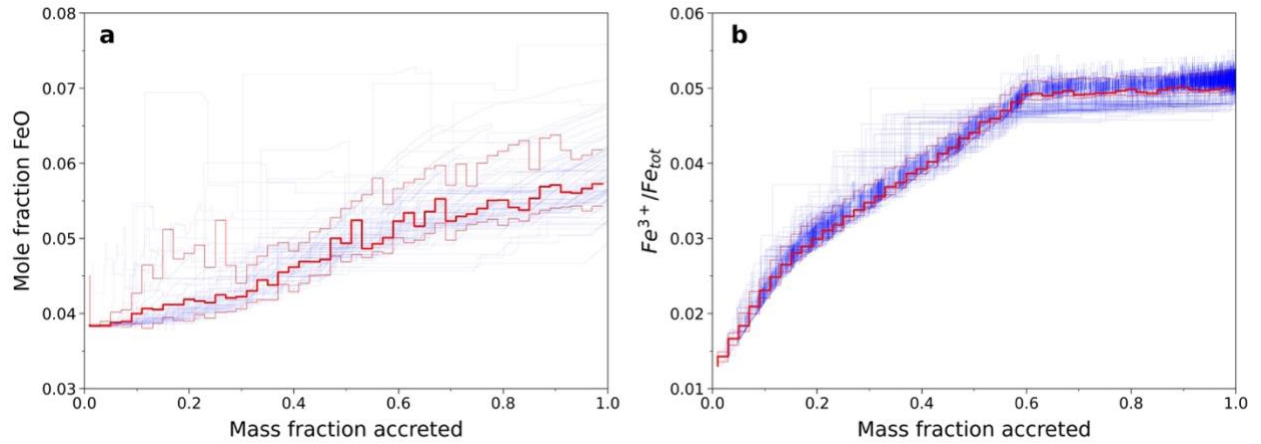


Figure 2. (a) Mantle FeO contents and (b) Fe^{3+}/Fe_{tot} ratios in Earth analogs throughout accretion. Bold red lines are the medians and semi-bold red lines are the 10th and 90th percentiles of all Earth analogs. Thin blue lines correspond to the evolution of individual Earth analogs. Bulk fO_2 of accreting bodies are adjusted so that Earth analogs match Earth's ~8 wt% FeO, on average, at the end of accretion. FeO contents in Earth analog mantles progressively increase due to core formation. Higher FeO and deeper magma oceans promote oxidation of the mantle.

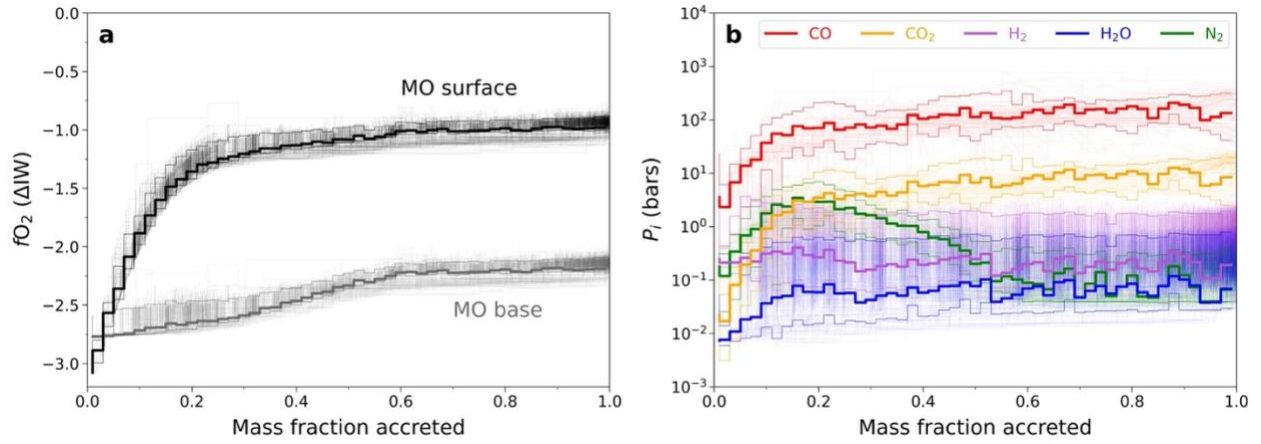


Figure 3. Evolution of (a) fO_2 at the base and surface of the magma ocean and (b) atmosphere composition of Earth analogs. Bold lines are medians and semi-bold lines are 10th and 90th percentiles of all Earth analogs. Thin lines are evolutions of individual Earth analogs. Magma oceans remain below the iron–wüstite buffer, resulting in CO-dominated atmospheres throughout accretion.

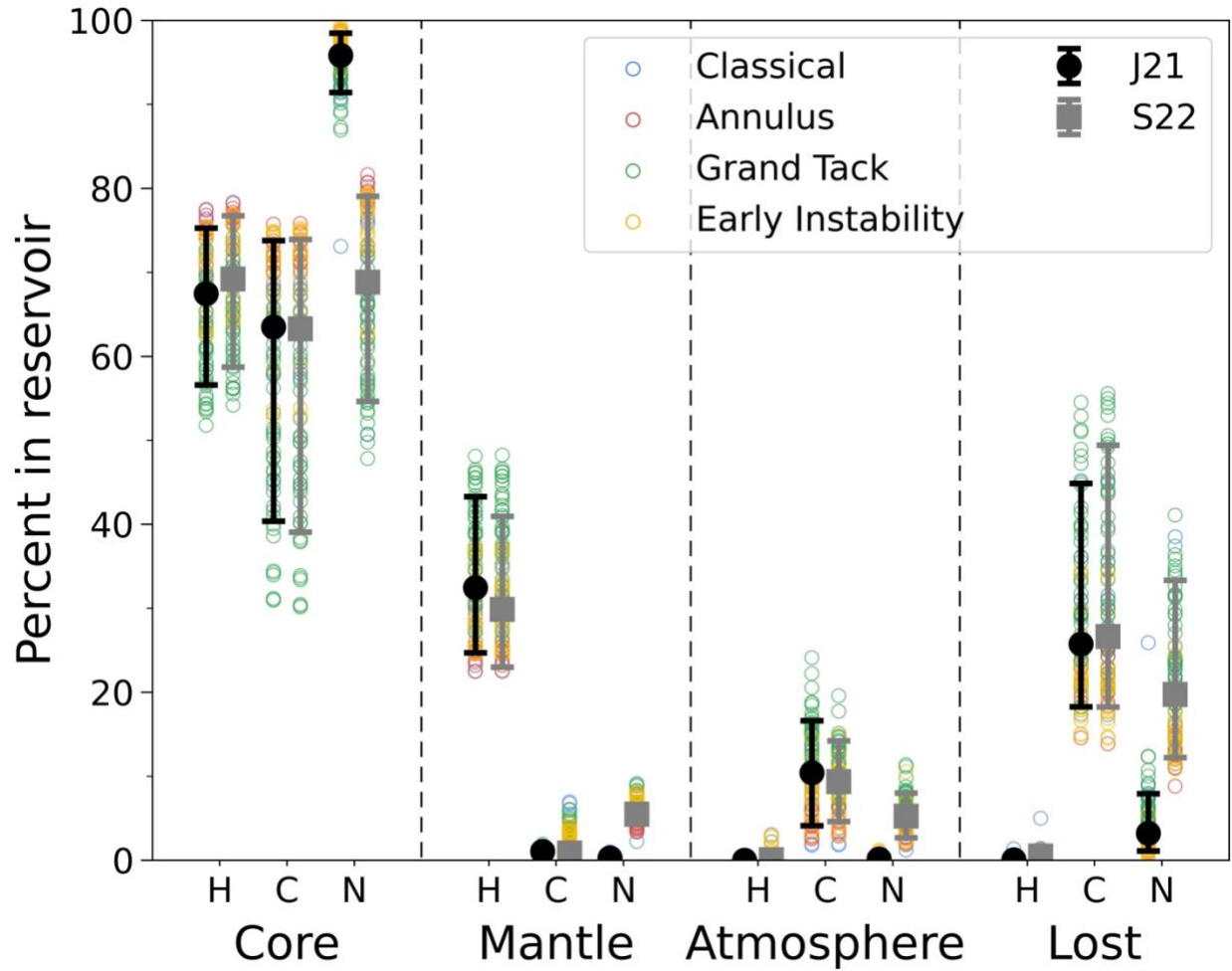


Figure 4. Distribution of H, C, and N between core, mantle, atmosphere, and loss to space at the end of accretion before magma ocean crystallization. “J21” and “S22” refer to metal–silicate partitioning parameterizations of N from Jackson et al. (2021) and Shi et al. (2022), respectively. Open symbols represent individual Earth analogs and solid black symbols are medians with error bars corresponding to 10th and 90th percentiles of all Earth analogs. The majority of Earth’s present-day H, C, and N lies in the core. Atmosphere erosion preferentially erodes C, but can also remove significant N if N is less siderophile.

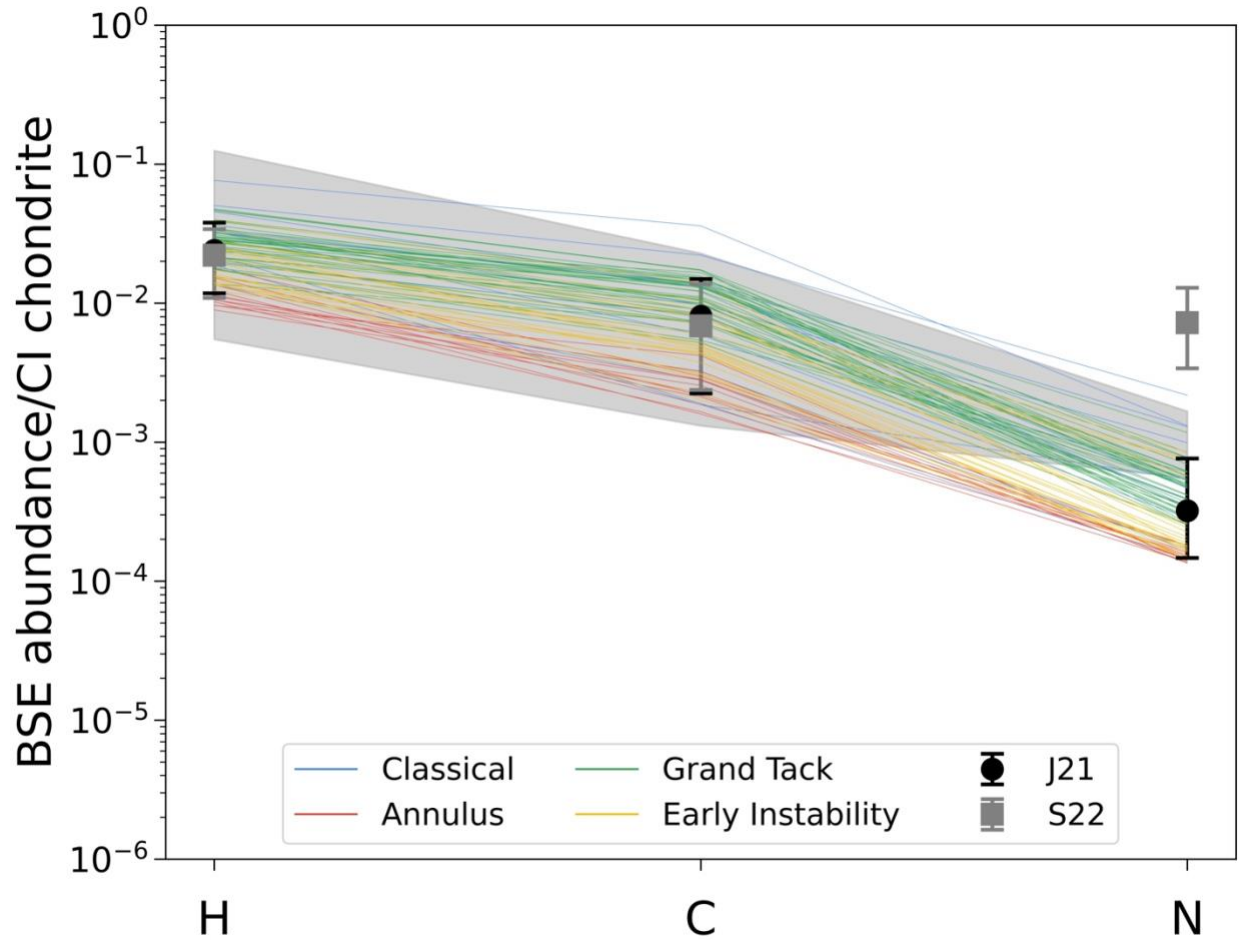


Figure 5. Abundances of volatile elements in the mantle and atmosphere relative to CI chondrites at the end of accretion. “J21” and “S22” refer to metal–silicate partitioning parameterizations of N from Jackson et al. (2021) and Shi et al. (2022), respectively. The shaded gray region represents estimated bulk Earth abundances (Hirschmann, 2018; Marty, 2012; Marty et al., 2020); lines represent individual Earth analogs from the “J21” scenario. Symbols are medians with error bars corresponding to 10th and 90th percentiles of all Earth analogs. Core formation alone could account for Earth’s superchondritic C/N ratio if N becomes siderophile enough at high pressures and temperatures.

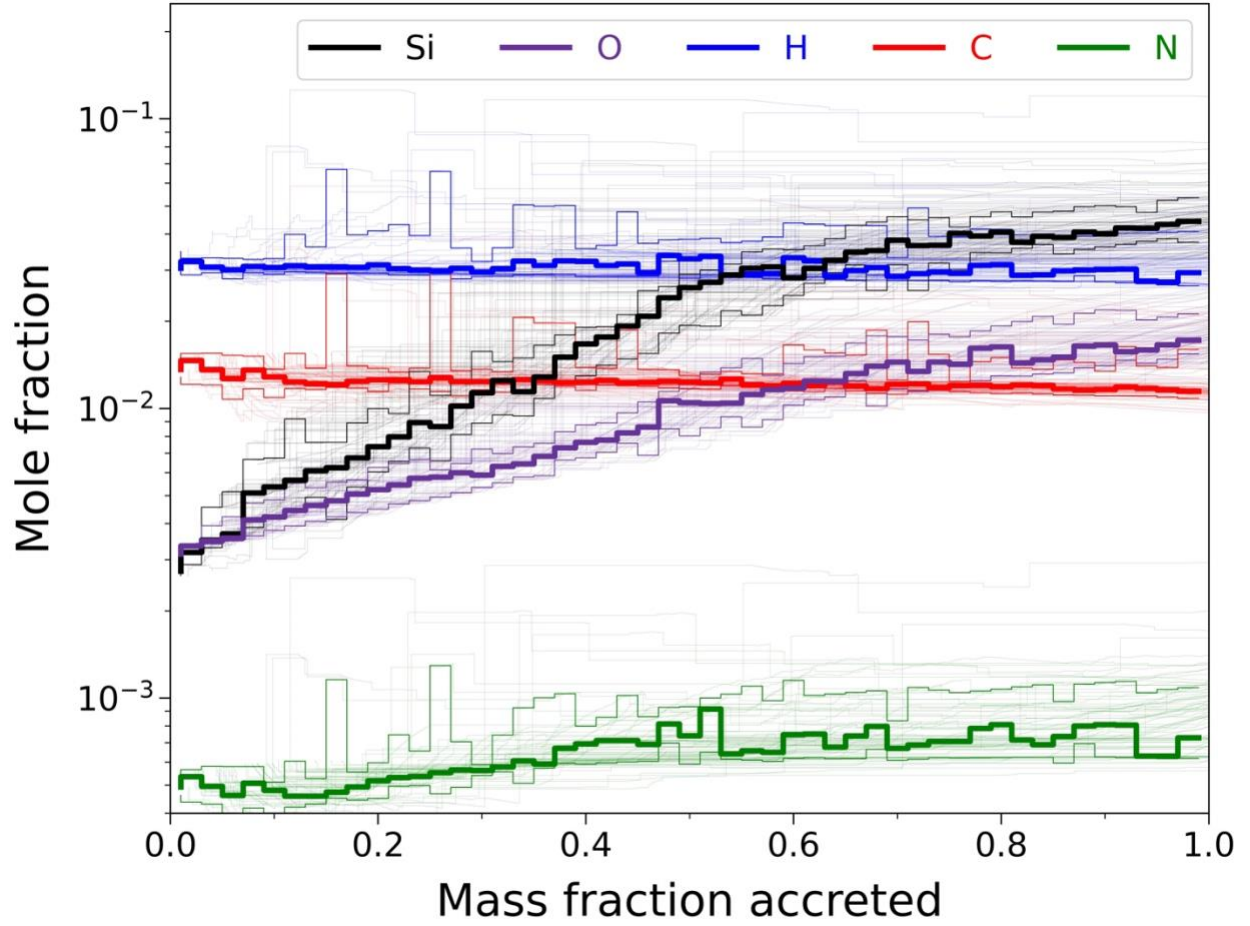


Figure 6. Concentration of light elements in the core throughout accretion of Earth analogs using the metal–silicate partitioning of N from Jackson et al. (2021). Bold lines are medians and semi-bold lines are 10th and 90th percentiles of all Earth analogs. Thin lines are evolutions of individual Earth analogs. At the end of accretion, Earth analog cores contain $2.4^{+0.6}_{-0.5}$ wt% Si, $0.5^{+0.2}_{-0.1}$ wt% O, $0.06^{+0.03}_{-0.01}$ wt% H, $0.25^{+0.20}_{-0.04}$ wt% C, and $0.03^{+0.01}_{-0.01}$ wt% N. Si and O concentrations in the core increase as pressures and temperatures of equilibration increase, whereas concentrations of light elements remain relatively constant due to the siderophile nature of H, C, and N.

Table 1. Volatile concentrations (ppmw) in the BSE, CI chondrites, and prescribed to accreting bodies. Inner and outer refer to bodies originating from within and beyond 2 AU, respectively. The effects of varying volatile concentrations, H:C:N ratios, and inner:outer Solar System abundance ratios are explored in Supplementary Text S5 and the Supplementary Data.

	H	C	N
BSE ¹	33–750	42–730	0.83–2.5
CI chondrite ²	6000	32000	1500
Inner (2% CI)	120	640	30
Outer (20% CI)	1200	6400	300

¹BSE constraints are from Marty (2012), Hirschmann (2018), and Marty et al. (2020).

²CI chondrite values are from Wasson and Kallemeyn (1988) and Vacher et al. (2020).