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Unraveling the Atmospheric Energy Input and Ionization due to EMIC-Driven Electron Precipitation from ELFIN Observations

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Key Points:

- ELFIN's pitch-angle resolved data provide an ideal input to the BERI model to estimate atmospheric ionization due to electron precipitation

- EMIC-driven precipitation occurs at all longitudes and 50°–70° latitudes; 74% of its input energy flux efficiently ionizes the atmosphere

- Ionization is enhanced in the mesosphere at altitudes ranging from 52 to 74 km, with ionization rates of ~100–200 pairs/cm^3/s

24 **Abstract**

25 Energetic electron precipitation (EEP) from the radiation belts into Earth's atmosphere leads to
26 several profound effects (e.g., enhancement of ionospheric conductivity, possible acceleration of
27 ozone destruction processes). An accurate quantification of the energy input and ionization due
28 to EEP is still lacking due to instrument limitations of low-Earth-orbit satellites capable of
29 detecting EEP. The deployment of the ELFIN (Electron Losses and Fields InvestigatioN)
30 CubeSats marks a new era of observations of EEP with an improved pitch-angle (0°–180°) and
31 energy (50 keV–6 MeV) resolution. Here, we focus on the EEP recorded by ELFIN coincident
32 with electromagnetic ion cyclotron (EMIC) waves, which play a major role in radiation belt
33 electron losses. The EMIC-driven EEP (~200 keV – ~2 MeV) exhibits a pitch-angle distribution
34 (PAD) that flattens with increasing energy, indicating more efficient high-energy precipitation.
35 Leveraging the combination of unique electron measurements from ELFIN and a comprehensive
36 ionization model known as Boulder Electron Radiation to Ionization (BERI), we quantify the
37 energy input of EMIC-driven precipitation (on average, $\sim 3.3 \times 10^{-2}$ erg/cm^2/s), identify its location
38 (any longitude, 50°–70° latitude), and provide the expected range of ion-electron production rate
39 (on average, 100–200 pairs/cm^3/s), peaking in the mesosphere – a region often overlooked. Our
40 findings are crucial for improving our understanding of the magnetosphere-ionosphere-
41 atmosphere system as they accurately specify the contribution of EMIC-driven EEP, which
42 serves as a crucial input to state-of-the-art atmospheric models (e.g., WACCM) to quantify the
43 accurate impact of EMIC waves on both the atmospheric chemistry and dynamics.

44 **Plain Language Summary**

45 Energetic electron precipitation (EEP) from Earth's radiation belts is a source of energy input to
46 the terrestrial atmospheric system and has the potential of impacting its chemistry and possibly
47 dynamics. Available datasets of EEP are incomplete due to instrumental limitations, hindering
48 the accurate quantification of EEP energy input, its properties, and the resulting ionization. Here,
49 we leverage the observations of the ELFIN CubeSats which provide high-resolution data for the
50 first time both in energy and look-direction (pitch-angle). We specifically focus on observations
51 during electromagnetic ion cyclotron (EMIC) waves, known to precipitate the most energetic
52 electrons, thus penetrating the atmosphere at low altitudes. We estimate the EMIC-driven EEP
53 energy input and its resulting ion-electron production rate as a function of altitude using a
54 sophisticated method that takes into account the EEP pitch-angle distribution. We find that 74%
55 of the energy input ionizes the atmosphere, primarily in the mesosphere (peaking between 52 and
56 74 km), a region underestimated by current recommended ionization rates. We provide the
57 region where EMIC-driven EEP is observed and the ionization rates at each location, which can
58 be used as input to comprehensive atmospheric models to ultimately quantify the accurate impact
59 of EMIC waves on Earth's atmosphere.

60 **1 Introduction**

61   Wave-particle interactions (e.g., Li & Hudson, 2019; Thorne, 2010) and field line
62 scattering (e.g., Capannolo et al., 2022a; Sivadas et al., 2019; Sergeev et al., 1983, 1983; Yahnin
63 et al., 2016, 2017) often occur in Earth's magnetosphere and precipitate electrons and ions into
64 the atmosphere. Here, particles interact with the ambient neutral molecules, depositing their
65 energy into the atmospheric system. The higher the particle energy, the deeper the particle will
66 penetrate into the atmosphere, depositing most of its energy and producing a variable number of
67 ion-electron pairs (ionization rate or production rate) along its trajectory (i.e., altitude) (e.g.,

68 Fang et al., 2010; Lazarev, 1967; Roble & Ridley, 1987), peaking at an altitude depending on its
69 initial energy (i.e., primary ionization peak). Figure 1A provides an illustration for this process.
70 Recent studies have demonstrated that the ionization profile also depends on the pitch-angle
71 distribution (PAD) of the impacting electrons (Marshall & Bortnik, 2018; Xu et al., 2018, 2020).
72 At low altitudes (<50 km), a secondary ionization peak is also observed due to the
73 Bremsstrahlung electromagnetic shower induced by the primary ionization (Xu & Marshall,
74 2019; Xu et al., 2021).

75 Ionization due to energetic electron precipitation (EEP) has several effects on the
76 atmospheric system. EEP can enhance the ionospheric density and conductivity (Khazanov et al.,
77 2018, 2021; Robinson et al., 1987; Yu et al., 2018), which not only influences the reliability of
78 communication systems (e.g., Booker & Wells, 1938; Cummer et al., 1998), but also affects the
79 propagation of geomagnetic currents that link the magnetosphere to the atmosphere, critical to
80 accurately model the entire geospace system (e.g., Korth et al., 2014; Peymirat & Fontaine,
81 1994; Xiong et al., 2020). On the other hand, ionization from EEP affects the atmospheric
82 chemistry by increasing the concentration of nitrogen and hydrogen oxides (NO_x, HO_x), which
83 act as catalysts of ozone (O_3) destruction (e.g., Meraner & Shmidt, 2018; Randall et al., 2005;
84 Sinnhuber et al., 2012). This ozone depletion can occur either through local interaction (*direct
85 effect*; Andersson et al., 2014; Zawedde et al., 2018) or via the subsequent transport of NO_x into
86 the stratosphere, with the potential for long-lasting effects (*indirect effect*; Hendry et al., 2021;
87 Maliemi et al., 2020; Randall et al., 2006, 2007). The indirect effect is particularly prevalent
88 during polar winters, as the absence of sunlight effectively removes the primary loss mechanism
89 of NO, extending its lifetime and facilitating downward transport. In turn, ozone loss could
90 accelerate the polar vortex, thus leading to changes in the radiative balance and possibly overall
91 climate (Baumgaertner et al., 2011; Maliemi et al., 2019; Rozanov et al., 2012; Salminen et al.,
92 2020). Finally, secondary ionization due to EEP reaches the stratosphere, possibly even
93 accounting for enhancements of radiation doses at airplane heights (Tobiska et al., 2016, 2018;
94 Xu et al., 2021), together with showers induced by galactic cosmic rays (GCR). All of these
95 effects clearly highlight the importance of quantifying the energy input and impact of EEP.

96 Despite the extensive implications of EEP for the atmosphere, accurate quantification and
97 inclusion of EEP in atmospheric models remain challenging, mainly due to the limitations
98 imposed by instrument capabilities and their spatial coverage. While the input of solar energetic
99 particles (SEP), GCR and auroral electrons has been recognized as important and quantified,
100 electrons of radiation belt origin (>50 keV up to ~multi-MeV) have largely been neglected.
101 Studies show that NO_x and HO_x concentrations are underestimated if atmospheric models only
102 include ionization due to SEP, GCR and auroral electrons (Andersson et al., 2018; Arsenovic et
103 al., 2016; Duderstadt et al., 2021; Pettit et al., 2021; Randall et al., 2015). Only recently, the
104 Coupled Model Intercomparison Project (CMIP6 – a collaboration effort to recommend
105 appropriate inputs to atmospheric models) included medium-energy electrons (MEE; Matthes et
106 al., 2017) obtained from the POES/MetOp constellation of low-Earth-orbit (LEO) satellites
107 orbiting at an altitude of 800–850 km (e.g., Evans & Greer, 2004; Rodger et al., 2010; Sandanger
108 et al., 2015). EEP observations (zenith-oriented telescope, 0°) from three integral channels (E1:
109 >30 keV, E2: >100 keV, E3: >300 keV) are parameterized with the Ap index (APEEP; van de
110 Kamp et al., 2016) and used to calculate ionization rates through the methodology (valid until 1
111 MeV) based on Fang et al. (2010). Simulations show that such MEE contribution is crucial for
112 improving the stratospheric response to ozone (Andersson et al., 2018).

113   Although POES provide extensive data (from 1998) and reasonable spatial coverage, the
114 data suffer from several limitations (e.g., proton contamination, radiation damage, high noise
115 levels, integral energy channels and very coarse energy steps). Nesse Tyssoy et al. (2021)
116 focused on a geomagnetically active period in April 2010 and examined eight different estimates
117 of MEE ionization rates, all derived from POES/MetOp measurements. They showed that
118 different methods of data processing (proton contamination removal, accounting for the detectors
119 degradation, choice of telescopes, electron energy ranges and energy spectra, and L-MLT
120 sampling), assumptions on the background atmospheric density (empirical or full chemical-
121 dynamical atmospheric models; Picone et al., 2002; Schmidt et al., 2006), and ionization rate
122 calculations (Fang et al., 2010; Rees, 1989; Wissing & Kallenrode, 2009), lead to variations in
123 the production rates of about one order of magnitude, with the largest spread found in the
124 recovery phase of the geomagnetic storm. The work by Nesse Tyssoy et al. (2021) particularly
125 exposes that there is still a large uncertainty in quantifying the MEE ionization impact – an
126 essential input to atmospheric models (Sinnhuber et al., 2021). Duderstadt et al. (2021) further
127 investigated the EEP impact on ozone concentration by considering the full energy range of
128 radiation belt electrons as input using data from the NASA Van Allen Probes (RBSP, Mauk et
129 al., 2013; Spence et al., 2013), after scaling them to high-resolution observations from the
130 FIREBIRD-II CubeSat (Johnson et al., 2020) during magnetic conjunctions. They focused on the
131 10-day-long electron loss occurring in March 2013 and showed that despite including MEE, the
132 CMIP6 particle precipitation underestimated ozone loss in the mesosphere by 20%–40%, further
133 motivating the community to account for the entire range of electron energy (>50 keV up to
134 multi-MeV).

135   In this study, we quantify the energy input and ionization rates from EEP driven by
136 EMIC (electromagnetic ion cyclotron) waves (0.1−5 Hz; Erlandson & Ukhorskiy, 2001; Fraser et
137 al., 2010). EMIC waves are typically radially localized, but are more extended in magnetic local
138 time (MLT) (e.g., Blum et al., 2017). In particular, EMIC-driven EEP is an important key loss
139 process of radiation belt electrons (Drozdov et al., 2022; Shprits et al., 2017; Zhang et al. 2016).
140 EMIC waves can indeed interact with electrons (e.g., Thorne, 2010) and precipitate them into the
141 atmosphere, especially from post-noon to post-midnight (Capannolo et al., 2022a; Carson et al.,
142 2012; Gasque et al., 2021). Precipitation has been observed by several LEO satellites or balloons
143 (e.g., Blum et al., 2015; Hendry et al., 2016; Li et al., 2014; Woodger et al., 2018), and occurs
144 frequently, with POES observations indicating an average of ~1−5 EEP events per day. Contrary
145 to auroral precipitation (limited to EEP at $\lesssim$30 keV), EMIC-driven EEP ranges from ~100s keV
146 up to multi-MeV energies, with enhanced precipitation efficiency above several hundred keV
147 (e.g., Angelopoulos et al., 2023; Capannolo et al., 2021, 2023; Hendry et al., 2017). As a result,
148 atmospheric ionization will be triggered over a broad range of altitudes, also extending to the
149 mesosphere and possibly further below it (due to the ~MeV electrons) – regions not accessed by
150 auroral electrons. Additionally, EMIC waves typically drive EEP at lower latitudes than auroral
151 EEP, likely leading to distinct atmospheric effects within these regions.

152   Capannolo et al. (2023) analyzed 144 EMIC-driven EEP events observed by the Electron
153 Losses and Fields InvestigatioN (ELFIN) CubeSats (Angelopoulos et al., 2020). For the first
154 time, ELFIN provided the electron PAD (0°−180°, resolution ~22.5°; Angelopoulos et al., 2023)
155 at several differential energy channels, making these measurements the optimal input to the
156 Boulder Electron Radiation to Ionization (BERI) model (Xu et al., 2020), which calculates the
157 ionization profiles accounting for both electron energy and PAD. We evaluate how the
158 production rates change as a function of altitude, energy, and pitch-angle and provide a map of

159 the locations of precipitation events with associated ionization rates. Finally, we quantify the
160 input energy flux resulting from the EMIC-driven precipitation events observed by ELFIN and
161 the effective ionization production in the atmosphere as predicted by BERI.

2 Methodology

2.1 ELFIN CubeSats and the Dataset of EMIC-Driven Precipitation Events

164 The ELFIN CubeSat pair (3 sec spinning period, ~90 min orbital period at ~350–475 km
165 of altitude) was launched in September 2019 and operated until its re-entry in mid-September
166 2022 (Angelopoulos et al., 2020). The CubeSats measured 50 keV–6 MeV electrons, from pitch-
167 angles of 0° up to 180°. Compared to the commonly used POES/MetOp satellites, ELFIN
168 offered significantly improved energy and pitch-angle resolution. ELFIN was able to observe
169 electrons locally trapped (mirroring), precipitating (inside the bounce loss cone, LC; likely
170 precipitated by some local mechanism like waves or field line scattering) and backscattered
171 (inside the anti-bounce-loss cone, ALC; likely due to atmospheric backscattering). Over the
172 majority of its orbits (i.e., most latitudes and longitudes) at ~450 km altitude, ELFIN measures
173 electrons that are entirely within the atmospheric loss cones (i.e., even electrons at 90° local
174 pitch-angle are in the drift LC), thus they will be eventually lost to the atmosphere. We
175 previously collected 144 EEP events driven by EMIC waves (details and analysis in Capannolo
176 et al., 2023) using 10s–100s keV proton precipitation (observed by POES/MetOp nearby the
177 ~MeV precipitation observed by ELFIN) as a proxy for EMIC waves (e.g., Capannolo et al.,
178 2021; Carson et al., 2012; Lyu et al., 2022; Miyoshi et al., 2008). Capannolo et al. (2023) solely
179 focused on the wave effects on electrons and the resulting local EEP, neglecting the consequent
180 backscattered electrons. Backscattered electrons would either precipitate in the opposite
181 hemisphere (i.e., ALC) or become quasi-trapped (i.e., drift LC) and eventually precipitate, albeit
182 not local to the ELFIN location, but further eastward along the electron drift shells. On average,
183 Capannolo et al. (2023) showed that EMIC-driven precipitation is observed on localized scales
184 ($\Delta L \sim 0.3$) over 15–24 MLT and 5–8 L shells. Note that this dataset of EMIC-driven EEP events
185 is constrained to nearby POES/MetOp observations of proton precipitation (i.e., a proxy for
186 EMIC waves), thus the considered 144 events are not a complete set of EEP from 2019 to 2022.
187 The scattering efficiency rises with increasing energy, and PAD within the loss cone becomes
188 flatter as energy increases, consistent with the understanding that EMIC waves primarily
189 precipitate high-energy electrons ($\gtrsim 1$ MeV). Low-energy electrons (down to ~200 keV) are
190 precipitated as well, though less efficiently (e.g., Angelopoulos et al., 2023; Capannolo et al.,
191 2021; Hendry et al., 2017). ELFIN's energy and pitch-angle resolution provides an
192 unprecedented opportunity to quantify the ionization effects resulting from the realistic PAD
193 observed during EMIC wave activity. Here, we use the PAD as measured by ELFIN for each 144
194 EMIC-driven precipitation events and use it as input to the BERI model, which is described in
195 Section 2.2. Given the field-of-view width of the pitch-angle sectors (22.5°), electrons at pitch-
196 angles close to the local LC angle encompass a mixed population (partly precipitating and partly
197 trapped), making fluxes measured in this transition region not entirely reliable. To address this
198 issue, we exclude fluxes at pitch-angles within ±22.5° from the local LC angle, and extrapolate
199 their values using the flux at the nearest pitch-angle (outside the ±22.5° range from the LC
200 angle).

2.2 BERI Model

The Boulder Electron Radiation to Ionization (BERI) model was developed by Xu et al. (2020) to simulate the electron precipitation effects. It builds on models previously developed and validated at CU Boulder (e.g., Marshall & Bortnik, 2018; Marshall et al., 2014, 2019; Xu et al., 2018; Xu & Marshall, 2019), based on the energetic precipitation Monte Carlo (EPMC) model described in Lehtinen et al. (1999). BERI considers as input an arbitrary background mass density profile and an electron PAD defined for each electron energy at an altitude of 500 km. The magnetic field is assumed along the zenith with a magnitude of ~40,000 nT (as measured at Poker Flat in Alaska at 700 km at night). The LC angle is ~66° at 500 km (defined by a mirror altitude of 100 km). In this study, we use the NRLMSIS2 (Naval Research Laboratory Mass Spectrometer and Incoherent Scatter; Emmert et al, 2020) model for the atmospheric density profile (Python package developed by Lucas, 2023) and the PAD observed by ELFIN for the 144 events associated with EMIC wave activity. Each electron motion is simulated through the atmospheric and magnetic background, accounting for mirroring forces and electron-neutral collisions. Particle scattering includes both elastic and inelastic collisions, during which a fraction of the electron energy is deposited into the atmosphere, while the remaining energy allows the electron to continue its motion until the electron energy is reduced to 2 keV, marking the end of the simulation (Marshall & Bortnik, 2018). As a result, BERI provides the ion-electron pair production rate as a function of atmospheric altitude, which is calculated from the energy deposition, assuming that an energy of 35 eV is needed to produce an ion-electron pair on average (Berland et al., 2023; Rees, 1989, p. 40). BERI also accounts for backscattered electrons, defined as upward-moving particles that are in the local bounce LC, either mirroring without collisions with the atmosphere or electrons that collided with the neutrals but mirrored back with some fraction of the initial energy. This is particularly useful to quantify the bounce loss cone flux for electrons precipitating into the atmosphere at different energies and pitch-angles (Berland et al., 2023; Marshall & Bortnik, 2018). The key advantages of using BERI are: (1) inclusion of the pitch-angle dependence of downward-going electrons and (2) increased computational efficiency of the ionization profiles through the use of a lookup table, which provides the atmospheric ionization response to electrons at different pitch-angles (0°–90°, 1° resolution) and energies (3 keV–3 MeV, uniformly spaced in a logarithmic scale). Extensive validation tests described in Xu et al. (2020) demonstrate that BERI reliably converts satellite observations to ionization production rates, needed as crucial inputs for atmospheric chemistry models.

3 Results

3.1 Quantification of the Atmospheric Ionization

In this section, we compare the ionization rates from ELFIN's high-resolution PAD and POES low-resolution PAD using the dataset analyzed in Capannolo et al. (2023). We illustrate the average ionization rates due to the realistic PAD observed by ELFIN during EMIC wave activity and put these in context with the current CMIP6 recommendations for atmospheric modeling.

3.1.1 Comparison During an ELFIN&MetOp Conjunction

On 2 February 2021, ELFIN-A observed EMIC-driven precipitation at L~4.8 and MLT~19.4 at ~01:59:30 UT, lasting ~30 s. Five minutes later, MetOp-02 observed both proton

244 and electron precipitation (a tell-tale signature of EMIC wave activity; e.g., Capannolo et al.,
245 2019a; Miyoshi et al., 2008) in a very similar L-MLT region (≲0.3 L and ≲2.3 MLT). As the
246 two satellites crossed a similar region within a few minutes of each other (i.e., magnetic
247 conjunction; Figure 1B), they likely observed the same or very similar patch of precipitation
248 driven by EMIC waves since EMIC waves last much longer than a few minutes (e.g., Blum et
249 al., 2017; Clausen et al., 2011). ELFIN PADs at three energy ranges are shown in Figure 1C–1E
250 and POES proton and electron fluxes are illustrated in Figure 1F-1G, with a purple bar indicating
251 the precipitation interval. Further details on this event are described in Capannolo et al. (2023).
252 While ELFIN measured significant fluxes from 63 keV up to 2.9 MeV by observing the full
253 electron pitch-angle distribution (Figure 2A), MetOp only observed fluxes deep inside (pitch-
254 angle of ~27°, mapped to ELFIN's altitude using a dipole approximation) and outside (~90°) the
255 local bounce LC in four integral energy channels (>30 keV, >100 keV, >300 keV, >700 keV).

256   We examine the atmospheric ionization using BERI. To use BERI, MetOp data (integral
257 electron flux) required a conversion to a PAD of differential number flux, calculated with the
258 Peck et al. (2015) method, which considers both the 0° and 90° POES telescopes. Figures 2A-2B
259 display the PADs for ELFIN and MetOp during the EMIC-driven precipitation event,
260 respectively. The LC angle is indicated by a vertical dashed line. Figures 2C-2D show the PADs
261 for ELFIN and MetOp (logarithmically) interpolated to the resolution of the BERI lookup table,
262 respectively. Since proton precipitation in the >30 keV and >100 keV electron channels was
263 strong during this event, the energy range analyzed is from ~300 keV to 1.4 MeV, based on the
264 Peck et al. (2015) routine applied to MetOp observations. For the comparison, we limit our
265 analysis to MetOp energy range (~300 keV up to ~1.4 MeV), thus ruling out ELFIN observations
266 outside this range. Note that while ELFIN observations provide a high pitch-angle resolution,
267 MetOp's PAD is interpolated from only 8 true measurements of integral electron fluxes. We
268 assume the simplest PAD for MetOp (a linear PAD in logarithmic space, Figure 2D), which is
269 likely less accurate than ELFIN's. Using the background atmospheric neutral density (Figure 2E)
270 at the locations of ELFIN (black) and MetOp (blue), we provide the BERI ionization rates in
271 Figure 2F. The profiles exhibit similar peaks (different by ~27 pairs/cm^3/s) and altitudes of peak
272 ionization (~4 km apart). These minor discrepancies are attributed to the different PADs used as
273 inputs to BERI. In particular, MetOp observed a larger contribution of low-energy electrons
274 (≲700 keV) compared to ELFIN, resulting in increased energy deposition at higher altitudes.
275 Conversely, the high-energy flux (≳1 MeV) observed by ELFIN is approximately twice that
276 observed by MetOp, likely due to the high instrumental noise level of POES/MetOp which
277 underestimates the high-energy precipitation (e.g., Nesse Tyssoy et al., 2016). Consequently, the
278 ionization peak is shifted to lower altitudes than that calculated from MetOp observations. In
279 summary, for this case study, the shape and magnitude of the ion pair production profiles in
280 Figure 2F are similar; however, ELFIN's improved resolution within the LC likely provides
281 more accurate profiles.

282   **3.1.2 Statistical Ionization Rates: Dependence on Altitude, Energy, Pitch-Angle, and**
283   **Location**

284   To calculate the expected atmospheric ionization rate due to EMIC-driven EEP, we
285 binned the PAD into 10° pitch-angle bins for each event in the dataset and for each energy, from
286 63 keV up to 2.1 MeV, and calculated the average values (Figure 3A). Although ELFIN can
287 observe electrons up to 6 MeV, the count rates in high energy channels are low, resulting in a
288 limited statistical significance; therefore, we limit the PAD to 2.1 MeV. The average PAD in

289 logarithmic scale was then interpolated to the pitch-angle and energy resolution in the BERI
290 lookup table (Figure 3B). Figure 3C further characterizes the precipitation input by displaying
291 the PAD of differential energy flux. At low pitch-angles, the LC is more filled by high-energy
292 electrons ($\gtrsim$300 keV) and the PAD flattens as energy increases. Both features are a consequence
293 of more efficient scattering for high-energy electrons than low-energy electrons, consistent with
294 EMIC waves as the driver of this precipitation. The atmospheric density profile is averaged from
295 the atmospheric density of each event (Figure 3D).

296   Figure 4 highlights the energy and pitch-angle dependence of the ionization rates, given
297 the average EMIC-driven PAD (Figure 3B). Figure 4A shows the ionization profile for
298 monoenergetic electrons (color-coded in the legend) considering the contribution from all pitch-
299 angles. Figure 4B illustrates the dependence on energy of the ionization peak (black) and its
300 associated altitude (blue). As expected (e.g., Xu et al., 2020), high-energy electrons precipitate
301 much deeper into the atmosphere, while low-energy electrons deposit their energy at higher
302 altitudes. The altitude of the ionization peak decreases by ~22 km per decade in energy, in
303 agreement with Berland et al. (2023). Additionally, the ionization rate increases as a function of
304 energy indicating that, given the PAD in Figure 3B, high-energy electrons contribute more to
305 atmospheric ionization compared to low-energy electrons. Figures 4C-4D show the variability of
306 the ionization profiles given a specific pitch-angle (color-coded in the legend) considering the
307 realistic PAD (Figure 3B) and the contribution from all energies. Electrons with lower pitch-
308 angles tend to penetrate deeper into the atmosphere and ionize it at lower altitudes than electrons
309 with higher pitch-angles, as the latter encounter a longer effective path through the atmosphere
310 (Marshall & Bortnik, 2018). Electrons outside the loss cone ($\gtrsim$66°) very rarely reach below 100
311 km altitude in the atmosphere, thus their ionization rates are negligible. The ionization peak
312 reaches the maximum at the LC edge (~60°), given that the energy input is highest at this
313 boundary (see the shape of the PAD in Figure 3A and 3B: inside the LC, the highest flux is close
314 to the LC angle). The ionization peak variability with pitch-angle in altitude (~7 km) and
315 production rate (from ~2 to ~5 pairs/cm^3/s) is less evident than the dependence on electron
316 energy, suggesting that the ionization rates depend more on the electron energy and less on their
317 pitch-angle. However, Figures 4C-4D highlight that different pitch-angles can affect different
318 altitudes, thus an accurate input is preferred to provide the most realistic estimates.

319   Figure 5A illustrates the variability of the total ionization rates resulting from the 144
320 EMIC-driven EEP events observed by ELFIN. Different statistical calculations of the PAD are
321 presented (mean, quartiles, min, max), with each line representing the associated ionization
322 profile. The total ionization rate from the mean PAD is shown in black, with the dotted and
323 dashed lines indicating its spread considering the standard error ($\sigma/\sqrt{N}$, where σ is the standard
324 deviation in each pitch-angle bin and N is the number of data points in each pitch-angle bin, for a
325 fixed energy). The quartiles of PAD are drawn in green: dashed line for the 25th, solid line for the
326 50th, and dotted line for the 75th. Additionally, we provide the possible range of ionization rates
327 in blue: the minimum (solid) and maximum (dotted) ionization rates are calculated using the
328 minimum and maximum PAD, respectively. The average atmospheric ionization peaks at 185
329 pairs/cm3/s ($\pm$ 42 pairs/cm^3/s) at 59 km, while the median ionization peaks at 70 pairs/cm^3/s at
330 the same altitude, suggesting that EMIC waves produce ionization of an order of 10^2 pairs/cm^3/s.
331 Furthermore, EMIC waves produce average ionization with rates $\gtrsim$100 pairs/cm^3/s over a broad
332 altitude range (from ~52 km up to ~74 km), covering the entire extent of the mesosphere.
333 Although the statistical significance of our dataset prevented the analysis of precipitation at

334 energies above 2.1 MeV, EMIC waves also precipitate electrons at higher energies, potentially
335 extending the altitude range even further, likely to ≲50 km. We also provide an estimate of the
336 energy deposition (energy flux per unit altitude), shown with the additional x-axes at the bottom
337 of Figure 5A (maroon for $eV/cm^2/s/km$ and blue for $erg/cm^2/s/km$). For this calculation, we
338 assume that an energy of 35 eV is needed to produce an ion-electron pair on average (Berland et
339 al., 2023; Rees, 1989, p. 40), and transform the pairs/cm^3/s in energy flux per unit length, which
340 is the ionization energy flux deposited at each km of altitude.

341 Supporting Figure S1 illustrates the distribution of the 144 EMIC-driven EEP events in
342 geographic (A-B) and magnetic (C-D) coordinates, with their associated peak of ionization rate
343 (Figure S1 A-C) and corresponding altitude (Figure S1 B-D). The event locations align with the
344 extent of the outer radiation belt, as also shown in Capannolo et al. (2023), and demonstrate
345 EMIC-driven precipitation occurs at any geographical longitude. Identifying the latitudinal and
346 longitudinal location of EEP is valuable for specifying regional energy deposition and global
347 extent in atmospheric models. It would be interesting to investigate whether the location of
348 precipitation (both in latitude and longitude) impacts atmospheric model outputs (e.g., HO_x and
349 NO_x production, O_3 reduction, neutral wind trends, etc.) and whether results are improved
350 compared to models that rely on zonal (longitudinal) averages (Verronen et al., 2020).

3.1.3 Comparison with CMIP6 Recommendations

352 Here, we compare the ionization rates calculated with BERI due to EMIC-driven
353 precipitation with those recommended as input to atmospheric models by the CMIP6 (Matthes et
354 al., 2017), which for the first time considers electron precipitation as an additional input of
355 climate models. The solar forcing depends on solar variability and includes both radiative forcing
356 (total solar and spectral solar irradiance and the F10.7 solar radio flux) and particle forcing from
357 GCR, SEP, and MEE precipitation, parameterized with Ap and Kp indices (3-hour indices). This
358 method also assumes that the electron energy is completely deposited into the atmosphere rather
359 than being partly deposited and partly backscattered (as assumed in BERI, more details in
360 section 2.2).

361 Note that the CMIP6 MEE input (van de Kamp et al., 2016) might underestimate the
362 contribution of high-energy electrons, since the >700 keV POES electron channel is not directly
363 included, and the Ap and Kp indices might not be a good proxy for MeV EEP (Hendry et al.,
364 2021). Additionally, although measurements of high-energy EEP are frequent (on average, ~1–5
365 times a day for POES), they are impulsive in flux and short-lived (~10s seconds) due to the fast
366 POES LEO orbit through the precipitation region. EEP generally lasts less than a day, probably
367 several hours at most (the most prolonged recorded EMIC wave activity in situ lasted 8–10
368 hours; Blum et al., 2020; Engebretson et al., 2015). Moreover, EEP is detectable only if high-
369 energy electrons are still populating the outer belt and have not been otherwise lost from the
370 radiation belts, due to factors like magnetopause shadowing or persistent and strong EMIC wave
371 scattering. As a result, the contribution of ~MeV EEP might be smoothed out in the daily
372 averages of CMIP6.

373 From the daily resolution of solar and particle forcing, we select the 144 days when
374 EMIC-driven precipitation was observed by ELFIN and average the ionization rates at each
375 altitude. We show these CMIP6 ionization rates in blue in Figure 5B, together with the rates
376 calculated from BERI (in black and green for the average and median ELFIN PAD,
377 respectively). At first glance, the ionization rates from BERI appear to be ~10 times higher than

378 the MEE recommended by CMIP6. However, BERI ionization rates are calculated specifically
379 during the spatially localized and short-lived EMIC-driven precipitation when the precipitating
380 fluxes are high. Since CMIP6 solar forcing is primarily intended for long-term studies, the MEE
381 contribution is instead averaged on a daily basis, likely smoothing out the short-lived high-
382 energy precipitation. Consequently, caution should be exercised when making a one-to-one
383 comparison of the ionization rates derived from BERI and MEE.

384 The MEE contribution is substantial over 80–120 km, thus primarily impacts the D and
385 E-regions of the ionosphere, but diminishes considerably at lower altitudes. In contrast, BERI's
386 estimates highlight that EMIC-driven precipitation primarily impacts the mesosphere, including
387 the D-region ionosphere and possibly extending into the upper stratosphere, where the ozone
388 concentration starts to rise. This discrepancy is primarily attributed to the underestimate of high-
389 energy precipitation in the MEE forcing (limiting the lower altitude boundary of the energy
390 deposition), while in BERI we specifically focus on the precipitation time intervals when fluxes
391 at high energies are significant, thus energy deposition can extend to lower altitudes. SEP also
392 ionizes the mesosphere, but their impact becomes significant only during SEP events, which are
393 less frequent compared to EMIC-driven EEP events.

394 This qualitative comparison demonstrates that accounting for accurate radiation belt
395 precipitation is critical to obtain realistic ionization profiles at each atmospheric altitude,
396 particularly in the mesosphere. Without BERI's estimates, CMIP6 recommendations would
397 suggest that solar forcing is significant only above 80 km and below 30 km of altitude,
398 overlooking the altitude range where high-energy electrons deposit the most energy. Note that
399 the latitudes affected by these ionization rates are also different: MEE mostly precipitate in
400 auroral latitudes, SEP are often observed at the poles and GCR are global. Therefore, direct
401 comparisons of ionization rates should be made with caution, since each latitudinal region will
402 be dominated by different energies of precipitation. Furthermore, EMIC-driven precipitation
403 typically occurs on much shorter timescales than atmospheric circulation and climate, and it is
404 still unclear if the cumulative effects of this short-lived precipitation could affect climate or only
405 have localized and transitory effects during ongoing precipitation. In the future, we plan to
406 explore if EMIC-driven ionization peaking in the mesosphere affects the atmospheric dynamics
407 and chemistry on the short (~hours) or long (~days) timescales.

408 ### 3.1.4 Comparing Ionization Rates based on a Realistic, Isotropic and Sinusoidal
409 ### Energy-Pitch-Angle Distribution

410 Supporting Figures S2 and S3 compare the ionization rates assuming different PADs as
411 input to the BERI model. The isotropic PAD assumption yields the highest ionization rate, as
412 expected, since the LC is filled for all pitch-angles. The sinusoidal PAD provides an intermediate
413 ionization rate, indicating that the realistic PAD falls off more rapidly than the sinusoidal PAD
414 with decreasing pitch-angle within the loss cone. We find that, on average, the sinusoidal PAD
415 produces ionization rates close to those from the more realistic PAD observed by ELFIN. For the
416 median, the isotropic PAD might be preferred in absence of high-resolution PAD from LEO
417 satellite data, with the caveat that ionization rates might be slightly overestimated. Future work
418 will assess whether these differences have a substantial effect on the outputs of atmospheric
419 models, either on short or long timescales.

3.2 Estimate of the Energy Input during EMIC-Driven Precipitation and Corresponding Ionizing Energy Flux and Backscattered Energy Flux

Given the PAD of each of the 144 EMIC-driven events, we calculate the input energy flux (eV/cm^2/s) of downward-going electrons for each precipitation event by multiplying the differential number flux (#/cm^2/s/sr/keV; Figure 3B) by the central energy of each ELFIN energy channel and integrating in pitch-angle and energy. We further calculate the energy flux responsible for ionizing the atmosphere by integrating the ionization profiles obtained with BERI across altitude, and converting the unit of ion-electron pairs into energy units (35 eV for each ion-electron pair on average; Rees, 1989, p. 40). The energy input for each event is shown on a latitude-longitude map in Figure 6 (A: from ELFIN's measurements; B: from BERI's ionization rates) in both eV/cm^2/s (black) and erg/cm^2/s (blue). Since part of the input energy flux measured by ELFIN is backscattered by the atmosphere, it is reasonable that the energy flux from BERI's estimates is lower. On average, EMIC-driven precipitation provides an input energy flux of ~2.06×10^{10} eV/cm^2/s (~3.29×10^{-2} erg/cm^2/s) in the LC (0°–66°). The average ionization rates from BERI provide an energy flux of ~1.53×10^{10} eV/cm^2/s (~2.45×10^{-2} erg/cm^2/s), indicating that ~74% of the input energy flux is ionizing the atmosphere, while the rest is backscattered (~0.53×10^{10} eV/cm^2/s, ~0.84×10^{-2} erg/cm^2/s). Figure 7 shows a cartoon of the energy budget due to local EMIC-driven precipitation: input and output energy fluxes are shown with the orange and pink arrows, respectively, and the effective atmospheric ionization due to such precipitation is shown in white.

Furthermore, since ELFIN is also able to observe upward-going electrons, we can quantify the backscattered (anti-loss cone, ALC) energy flux during the EMIC-driven precipitation events and compare it with the values expected from BERI. On average, ELFIN observes an energy flux of ~5.20×10^{10} eV/cm^2/s (~8.33×10^{-2} erg/cm^2/s) for upward-going electrons (90°–180°) and ~0.74×10^{10} eV/cm^2/s (~1.18×10^{-2} erg/cm^2/s) in the ALC (114°–180°). The estimate of upward-going electrons from BERI's model is ~5.18×10^{10} eV/cm^2/s (~8.29×10^{-2} erg/cm^2/s), in agreement with ELFIN's observations, providing strong evidence that BERI accurately computes electron backscatter rates. We calculate BERI's backscattered energy flux either (1) by adding the energy flux that is not ionizing the atmosphere (~26% of the LC energy flux from ELFIN) to the energy flux outside the LC (over 66°–90°) measured by ELFIN or (2) by subtracting the energy flux ionizing the atmosphere estimated from BERI from the input energy flux from ELFIN. More in depth comparison with backscattering measurements is provided in Berland et al. (2023).

To compare the contribution from EMIC-driven precipitation with the contribution due to auroral precipitation estimated in Newell et al. (2009), we convert the observed input energy flux inside the LC to energy per unit time (power expressed in Watts). Considering a conservative extent of EMIC waves of ~0.5 L at ~6.5 L and ~3 MLT (corresponding to ~1° magnetic latitude and 45° magnetic longitude; Clausen et al., 2011; Blum et al., 2017; Hendry et al., 2020; Mann et al., 2014), we estimate a total hemispheric energy flux of ~6 MW. Only ~74% of is expected to ionize the atmosphere (~4 MW). However, previous studies show that EMIC waves can persist for several hours and can extend over many MLT sectors (e.g., Blum et al., 2020; Engebretson et al., 2015; Kim et al., 2016; Mann et al., 2014), though remain radially localized (Blum et al., 2017; Paulson et al., 2014; Mann et al., 2014), suggesting that EMIC-driven precipitation is indeed possible over wide longitudinal extents though narrower in latitudinal scales. This agrees with the spatial distribution of EMIC-driven precipitation events, as shown in Figure 6 and

465 Figure S1. Thus, we can provide an upper limit of the EMIC-driven hemispheric contribution of
466 ~22 MW (with ionization estimated at ~16 MW). This calculation is derived under the
467 assumption of the same latitudinal extent, but with a broader longitudinal range of ~180°,
468 corresponding to 12h in MLT – the upper MLT limit of EMIC waves observed by Engebretson
469 et al. (2015). The actual EMIC-driven hemispheric input power lies between these two bounds
470 (6–22 MW), and it notably depends on the ongoing EMIC wave activity, including its intensity
471 and spatial extent, as well as the resulting EEP occurring in specific geographical regions.
472 Accurately quantifying the regional extent of EMIC waves and EMIC-driven precipitation poses
473 a significant challenge, primarily due to limitations in equatorial and LEO satellite coverage. It
474 requires a case-by-case analysis, relying on a combination of measurements from a constellation
475 of spacecraft and ground-based networks (e.g., Hendry et al., 2020; Engebretson et al., 2015;
476 Yahnin et al., 2021) and/or imaging of precipitation regions (e.g., Marshall et al., 2020).
477 Furthermore, the specific input of EMIC-driven EEP can significantly vary over space and time
478 and should be quantified case-by-case as it depends not only on the EMIC wave activity, but also
479 on the radiation belt content (i.e., the outer radiation belt should be populated by enough
480 electrons) and the resulting spatial extent of efficient EEP (i.e., where the resonance conditions
481 are met). In summary, the estimated EMIC-driven power is lower than the auroral power
482 estimated by Newell et al. (2009) (1–6 GW), as expected given the broader geographical extent
483 of the aurora. However, while auroral precipitation mainly affects the E-region and partly the D-
484 region of the ionosphere, EMIC waves are capable of scattering much more energetic electrons,
485 which deposit their energy at lower altitudes. Additionally, auroral precipitation maps to much
486 higher latitudes in comparison to EMIC-driven precipitation. Consequently, the atmospheric
487 effects of these two phenomena may exhibit variations due to their distinct characteristics and
488 geographical locations.

4 Discussion & Conclusions

490   In this study, we quantify the atmospheric ionization from EMIC wave-driven electron
491 precipitation, utilizing the observations of 144 EMIC-driven EEP events from the ELFIN
492 CubeSats from 2019 to 2022 (further described in Capannolo et al., 2023) as input to the BERI
493 model (Xu et al., 2020). We provide the statistical PAD of EMIC-driven precipitation, its energy
494 input, and its resulting ionization rates. EMIC-driven EEP exhibits a larger contribution of high-
495 energy electrons compared to low-energy electrons because EMIC waves more efficiently scatter
496 high-energy electrons. Due to the limited statistical significance at high energies, our results are
497 valid for the energy range over 68 keV–1.932 MeV (or 63 keV–2.1 MeV from ELFIN's data).
498 Our key findings are summarized and discussed below.

499   1. ELFIN data provides unprecedented measurements of precipitating electrons with high
500   energy and pitch-angle resolution, which serve as optimal inputs to the BERI model,
501   significantly enhancing the accuracy of electron precipitation input and associated
502   ionization rate estimates.

503   We showed an example of the data quality comparison between POES and ELFIN
504 (Section 3.1.1), highlighting that POES observations require multiple assumptions to remove
505 proton contamination from low-energy electron measurements and to convert the measured
506 integral energy fluxes to the differential electron flux. The example also highlights the
507 limitations of the POES detector in measuring low flux (Nesse Tyssoy et al., 2016), especially at
508 high-energy, which in turns leads to a peak in ionization rates at a slightly higher altitude than

509 what is estimated using the ELFIN data. Although the differences in ionization profiles are
510 moderate, it remains unknown whether a small change in the altitude distribution of ionization
511 substantially impacts the atmospheric chemistry and dynamics. Whenever available, we
512 recommend using ELFIN data to estimate ionization rates since it offers improved resolution in
513 energy and pitch-angle. However, ELFIN data has more constrained spatial and temporal
514 coverage compared to the POES constellation. Therefore, we encourage the community to
515 contemplate launching mutliple ELFIN-like satellites into LEO, as single-point observations are
516 insufficient to capture the spatiotemporal variability of EEP across the globe. In the meantime,
517 careful comparison and cross-calibration between the ELFIN and POES spacecraft (e.g., during
518 the numerous magnetic conjunctions) could allow us to potentially obtain a hybrid dataset that
519 provides improved energy and pitch-angle resolution (from ELFIN's characteristics) and broad
520 spatial and temporal coverage (from POES). Although this is a challenging task, it might provide
521 a solution to accurately quantify atmospheric ionization over an extensive spatial coverage.

522   2. Electrons scattered by EMIC waves primarily deposit their energy in the mesosphere,
523   with average values of ~100–200 pairs/cm^3/s (over a broad altitude range of ~52–74
524   km) and energy deposition of ~10^{-3} erg/cm^2/s/km.

525   EMIC waves primarily scatter high-energy electrons, thus affect the D-region and
526 altitudes below, potentially extending even into the stratosphere through Bremsstrahlung-induced
527 secondary ionization (Berland et al., 2023; Xu et al., 2021). Our analysis provides an estimate of
528 the average ionization profile and the possible range of ionization due to EMIC waves, which
529 could serve as an input to atmospheric models to quantify the effect of EMIC-driven
530 precipitation in the production of HO_x and NO_x and possible subsequent ozone depletion.
531 Previous studies have shown that radiation belt electrons could have a substantial impact on the
532 atmosphere's chemistry; thus, the contribution from EMIC-driven precipitation (i.e., the most
533 energetic population) should not be excluded from consideration in atmospheric models. We
534 have shown that the current CMIP6 recommendations are underestimating the contribution of
535 high-energy EEP in the mesosphere.

536   3. For the events studied here, EMIC-driven precipitation averages: ~0.033 erg/cm^2/s of
537   input energy flux, ~0.025 erg/cm^2/s of energy flux that ionizes the mesosphere, and
538   ~0.008 erg/cm^2/s energy flux of backscattered electrons. The hemispheric input power
539   of EMIC-driven precipitation is in the range of ~6–22 MW, depending on the area of
540   precipitation (1° magnetic latitude extent and ranging from 45° to 180° magnetic
541   longitude).

542   The energy input and the input power due to EMIC-driven precipitation are lower than
543 those estimated for auroral precipitation (Newell et al., 2009). This is reasonable because aurora
544 covers a much broader area than EMIC-driven precipitation and the low-energy electrons
545 producing aurora are more abundant than those scattered by EMIC waves. However, it is
546 important to emphasize that EMIC-driven precipitation affects both lower altitudes and lower
547 magnetic latitudes (~56°–68°) than the auroral precipitation (>70°). Therefore, its effect on the
548 atmospheric chemistry and dynamics is distinct from that of auroral precipitation. Our relatively
549 limited dataset of EMIC-driven precipitation events demarcates a region where EMIC-driven
550 precipitation is possible, however, does not allow us to infer the exact spatial or temporal extent
551 of the precipitation region. A combination of satellite and possibly ground-based measurements
552 (e.g., Hendry et al., 2020; Engebretson et al., 2015; Yahnin et al., 2021) or imaging of

553 precipitation (Marshall et al., 2020) would overcome this limitation and improve the estimate on
554 the evolution in time and space of the precipitation region.

555   Furthermore, our analysis reveals that, in the absence of a realistic PAD (as provided by
556 ELFIN), a sinusoidal PAD for EMIC-driven precipitation provides a reasonable assumption for
557 average ionization rates. We also showed that, as expected, higher energy electrons and field-
558 aligned electrons deposit their energy at 50–60 km, while electrons with lower energy and with
559 pitch-angles closer to the loss cone impact higher altitudes. Moreover, while EMIC-driven
560 precipitation is spatially localized in latitude and typically occurs over 50°–70° in latitude, it is
561 observed over a much broader geographical longitude, due to EMIC waves being more extended
562 in local time than radially. As an additional comparison, we demonstrated that BERI estimates an
563 energy flux of backscattered electrons comparable to ELFIN's observations, opening a new line
564 of research in accurate modeling and comparison of the PAD of backscattered electrons. Some
565 progress has been made recently by Berland et al. (2023).

566   It is noteworthy that our estimates of energy flux and ionization can only account for the
567 local precipitation driven by EMIC waves and disregard any other contributions due to drifting
568 and backscattered electrons. In fact, electrons backscattered by the atmosphere, as predicted by
569 BERI and previous models (e.g., Cotts et al., 2011; Marshall & Bortnik, 2018), with pitch-angles
570 inside the ALC (symmetric to the LC) are expected to precipitate in the opposite hemisphere,
571 thus providing some energy input there as well. From ELFIN's observations, the ALC (114°–
572 180°) energy flux during EMIC-driven precipitation is ~36% of that in the LC, providing a total
573 energy flux of ~0.74×10^{10} eV/cm^2/s (~1.18×10^{-2} erg/cm^2/s). Additionally, at ELFIN's altitude,
574 the electrons outside the LC and the ALC are likely only locally trapped or quasi-trapped in the
575 drift LC (Tu et al., 2009), indicating that they will also eventually deposit energy into the
576 atmosphere over an electron drift period. These estimates require further modeling and
577 understanding that we leave as future work.

578   Pilot studies by Hendry et al. (2021) and Ozaki et al. (2022) evaluated the ozone
579 depletion specifically due to EMIC-driven precipitation using modeling, such as 1D Sodankyla
580 Ion and neutral Chemistry model (Turunen et al., 2009; Verronen et al., 2005, 2016) and a
581 combination of observations, such as the SABER experiment board the TIMED satellite (Rong et
582 al., 2009; Russel et al., 1999). While Hendry et al. (2021) estimated ozone depletions of up to
583 ~10% (short-lived during Summer EEP and longer-lived during Winter), Ozaki et al. (2022)
584 observed changes up to 60%, indicating that the effect of EMIC-driven precipitation on ozone
585 depletion can be quite significant. Although these results are based on case studies, and a
586 comprehensive analysis of the overall effects of EEP on the atmosphere as a whole, especially
587 over long timescales, is still unavailable, they provide further evidence of the critical importance
588 of incorporating and exploring the impact of EMIC-driven precipitation into the atmospheric
589 system.

590   In summary, the analysis we presented in this paper provides a database of accurate
591 ionization rates associated with each EEP, together with an average ion-electron pair production
592 due to EMIC waves. These outcomes can serve as inputs to advanced models, such as WACCM,
593 to assess the impact of EMIC-driven precipitation on the entire atmosphere. This assessment can
594 encompass various aspects, including the concentration of HO_x, NO_x and ozone, as well as the
595 resulting effects on its dynamics, both on short and long timescales. Efforts to extend our limited
596 knowledge on the temporal and spatial variation of the precipitation region are also worth
597 pursuing to improve estimates on the EEP extent in latitude, longitude and duration. Through

598 these efforts, we can strive towards a more precise and complete representation of the
599 magnetosphere-ionosphere-atmosphere coupling and the Earth's atmospheric system itself,
600 advancing our understanding of its underlying physical processes and ultimately enhancing the
601 accuracy of atmospheric models.

Open Research

618 ELFIN data are available at https://data.elfin.ucla.edu/ and processed using SPEDAS
619 routines specifically written for processing ELFIN data by the ELFIN UCLA team
620 (Angelopoulos et al., 2019). POES/MetOp data are available at:
621 https://www.ncei.noaa.gov/data/poes-metop-space-environment-monitor/access/l1b/v01r00/. The
622 dataset of the 144 EMIC-driven events is available at https://doi.org/10.5281/zenodo.7697272.
623 The BERI code used to perform the calculations is at https://doi.org/10.5281/zenodo.3945306.
624 We used the Python package at https://doi.org/10.5281/zenodo.10922485 (release v0.7.0) to
625 obtain the MSIS background atmosphere. CMIP6 recommendations are available at
626 https://solarisheppa.geomar.de/cmip6. Data used to produce the figures is at
627 10.5281/zenodo.8222649. Data analysis was conducted using MATLAB. Figures 1A-1B and
628 Figure 7 are created with Keynote for Mac OS.

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1025

Figure 1: A: Illustration of the different altitudes where energetic particles ionize the atmosphere, depending on their energy; the radiation belt precipitation is marked in purple. Plasmasheet and auroral precipitation is for electrons of ≲30 keV, radiation belt precipitation is for electrons at energies of 10s keV up to ~10 MeV, SEP are solar protons of energies of 10s to 100s MeV, and GCR (protons and α-particles) have typical energies of 100s MeV up to 100s GeV. B: overview of the geometry of the ELFIN&MetOp magnetic conjunction and diagram of electron precipitation (purple arrow) resulting from wave-particle interactions in the outer radiation belt; LEO satellite orbits are shown in blue (ELFIN) and green (MetOp); their intersection determines the magnetic conjunction, which occurs nearby the electron precipitation. C–G: figure adapted from Capannolo et al. (2023), which depicts the observations during the ELFIN&MetOp conjunction. C–E: PAD observed by ELFIN at 3 different energy intervals; solid (dashed) horizonal lines indicate the local loss (anti-loss) cone; the precipitation is highlighted with a purple bar. F-G: proton and electron flux observed by MetOp in conjunction with ELFIN; solid (dotted) lines indicate the precipitating (trapped) particles; the precipitation of both protons and electrons is highlighted with a purple bar.

Figure 2: Ionization rates during the EMIC-driven EEP event observed during the ELFIN&MetOp conjunction (Figure 1C–G). A-B: PADs from (A) ELFIN-A data and (B) MetOp-02 data processed with the Peck et al. (2015) routine, color-coded by energy; C-D: PADs interpolated to the BERI lookup table resolution; E: background atmospheric density at the satellite locations (ELFIN: black, MetOp: blue); F: ionization rates from BERI for each satellite (ELFIN: black, MetOp: blue). The loss cone at 500 km is indicated by the vertical dashed line in A, B, C, D.
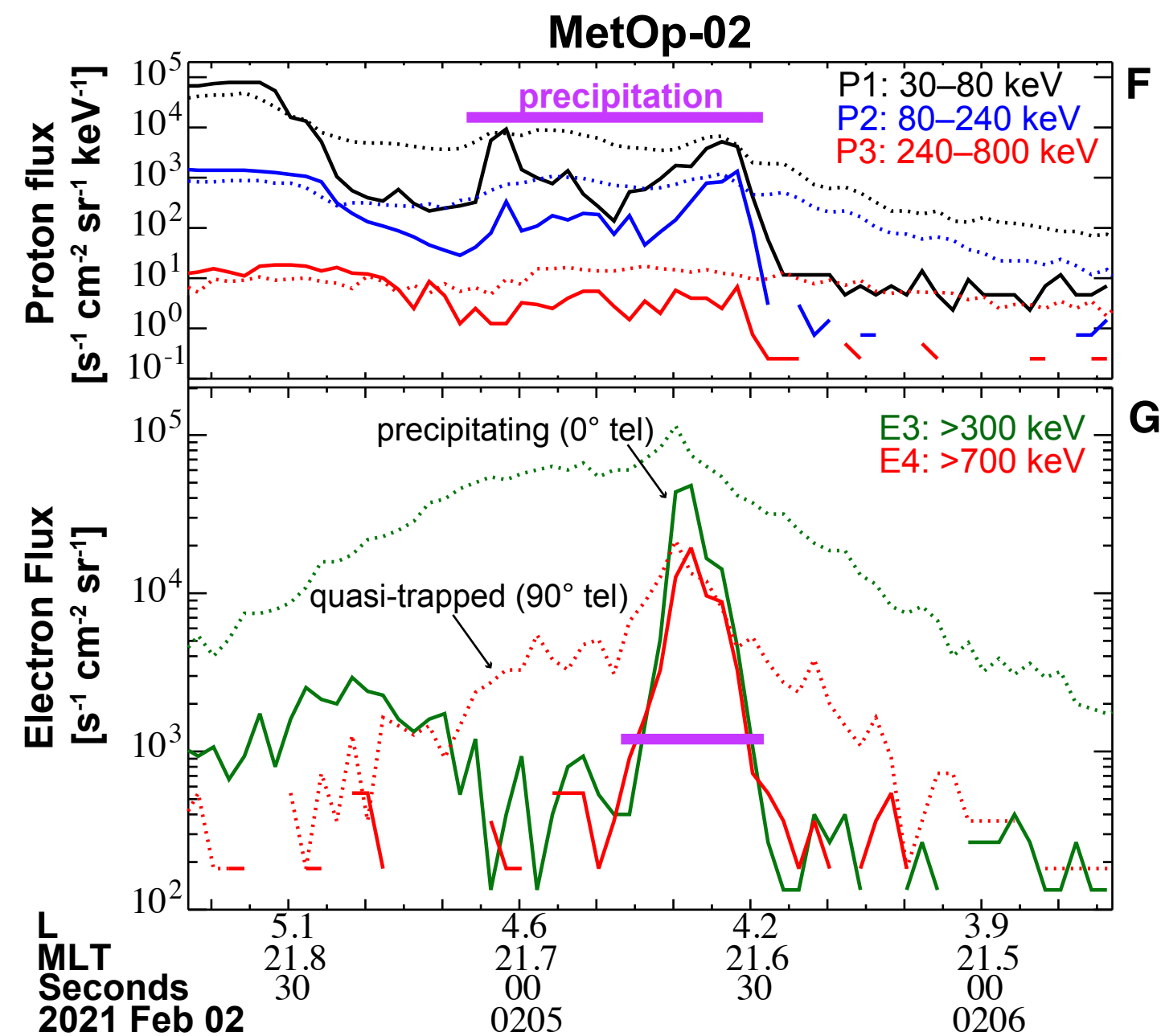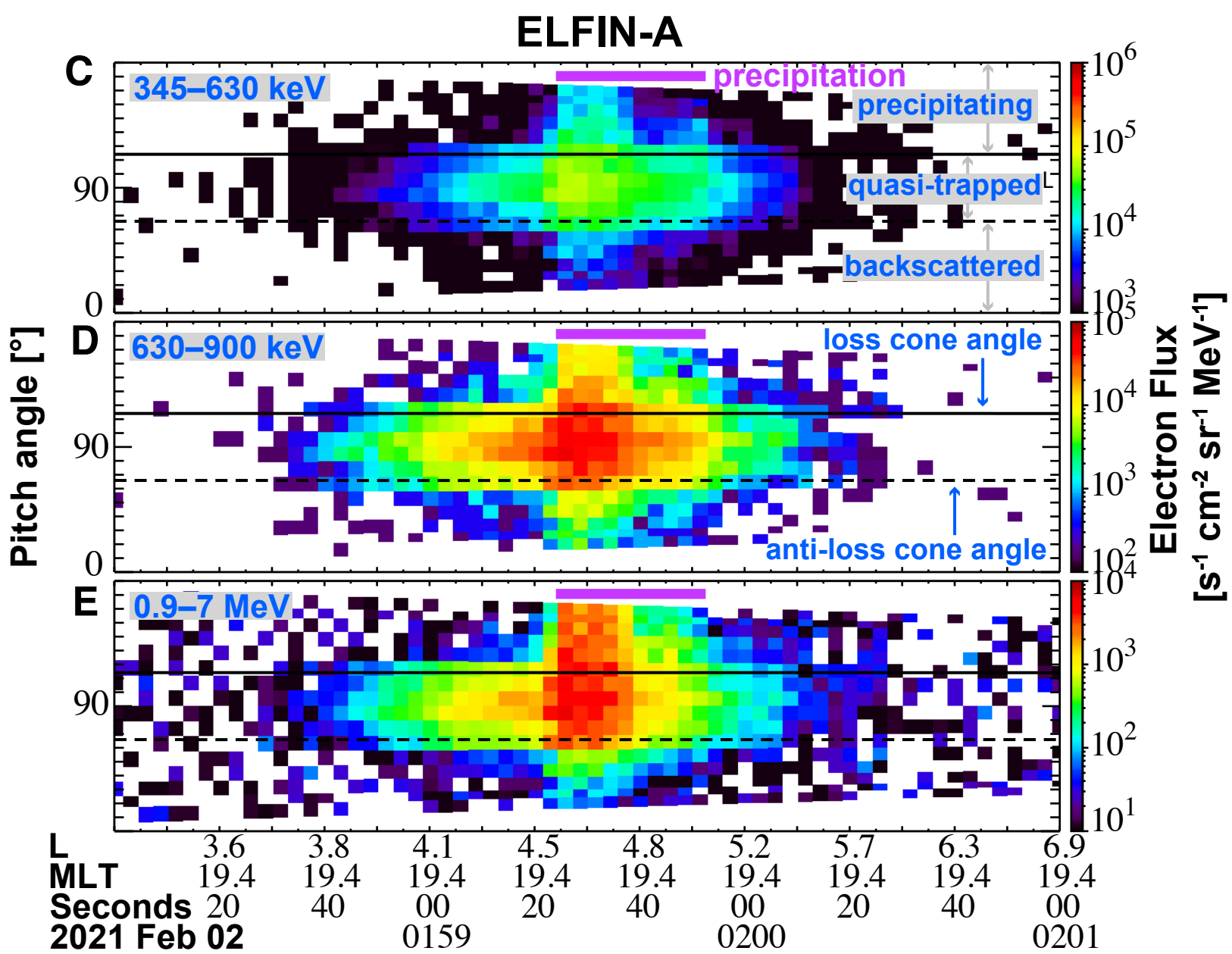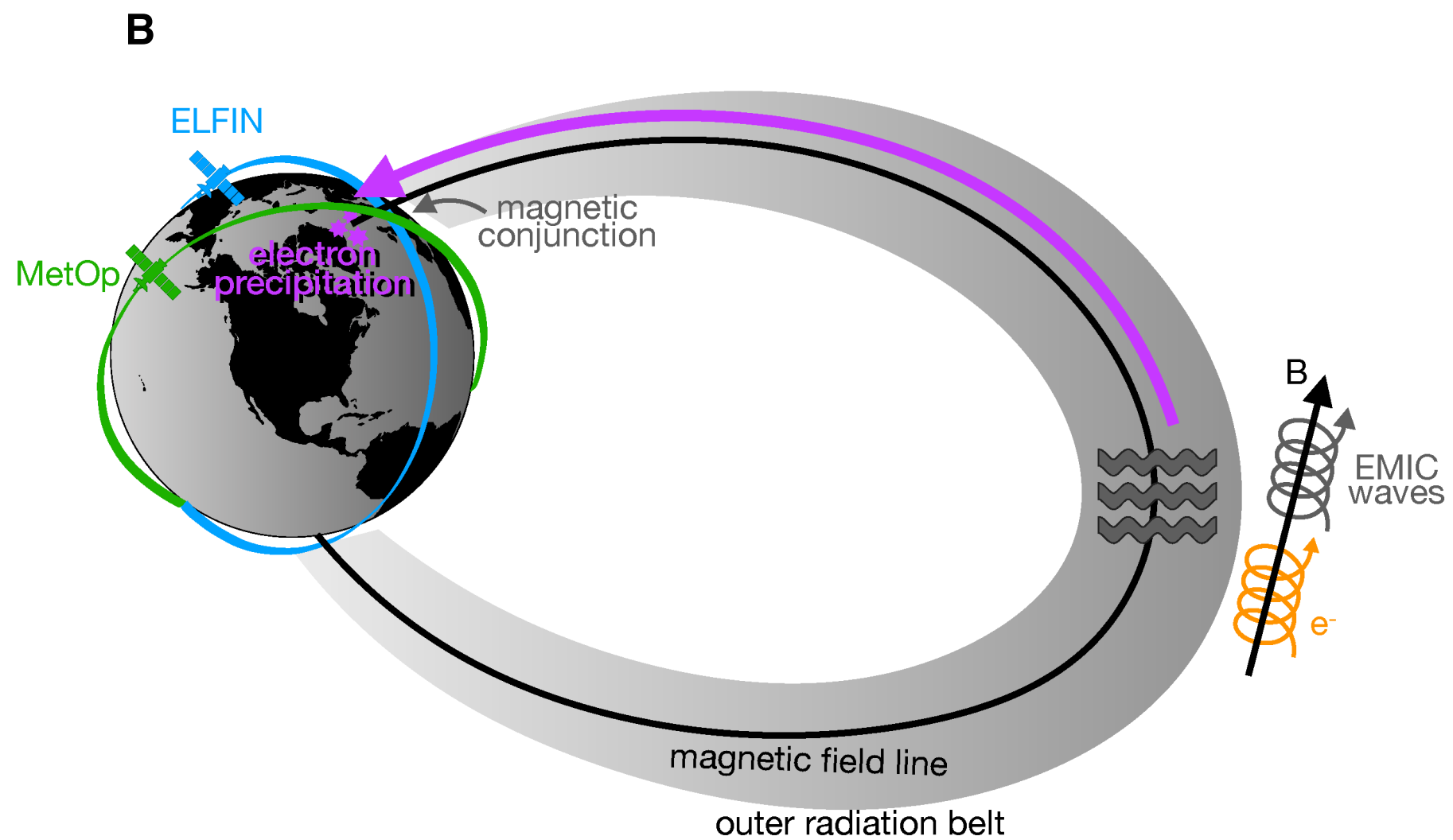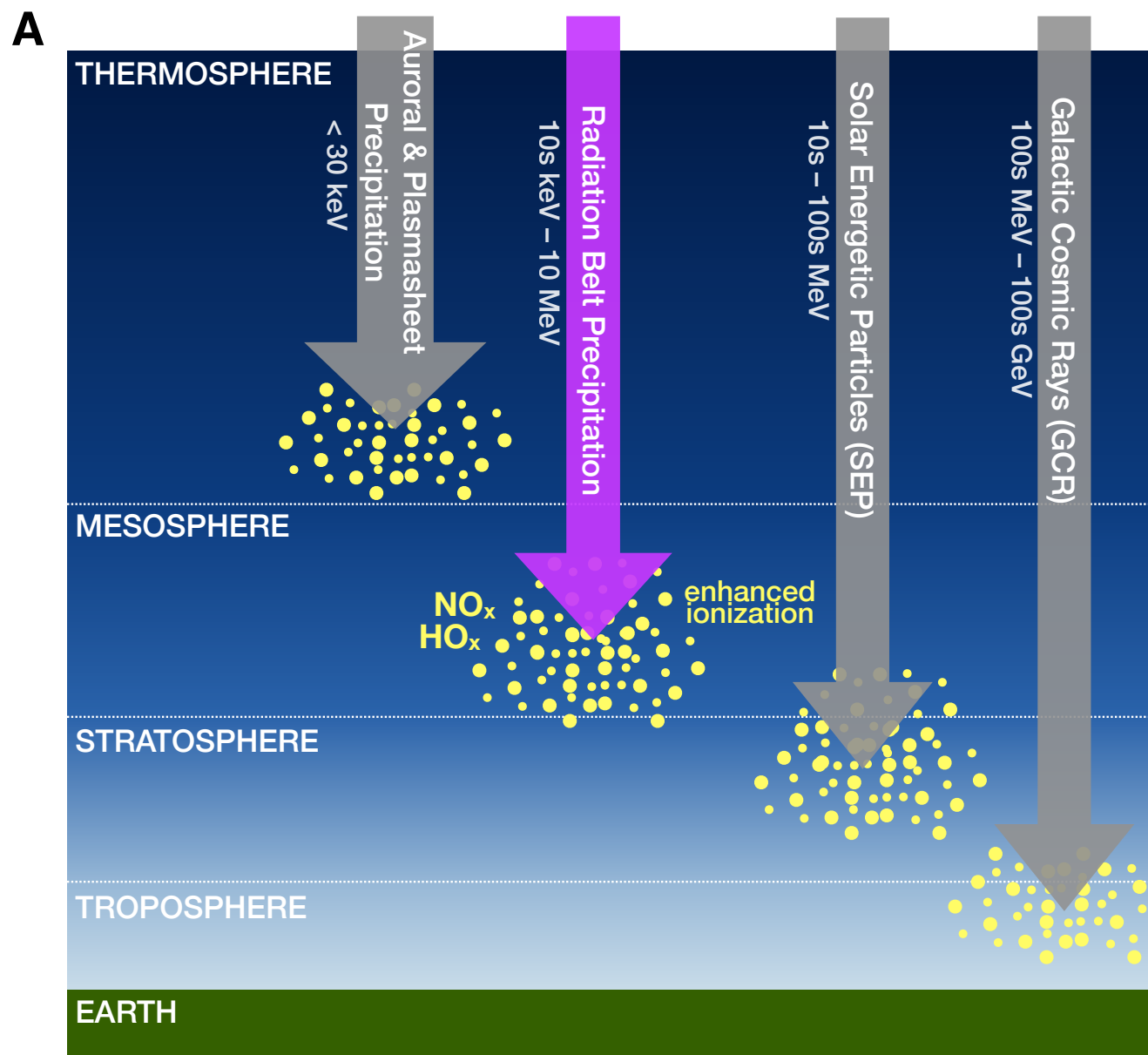
Figure 3: Inputs to BERI model. A: PAD from statistical EMIC-driven EEP data (from ELFIN), color-coded by the center energies in each energy channel; B: PAD interpolated to the lookup table resolution, color-coded by energy; C: PAD for the energy flux from the PAD in B; D: background density for each event (black) and averaged at each altitude (red).
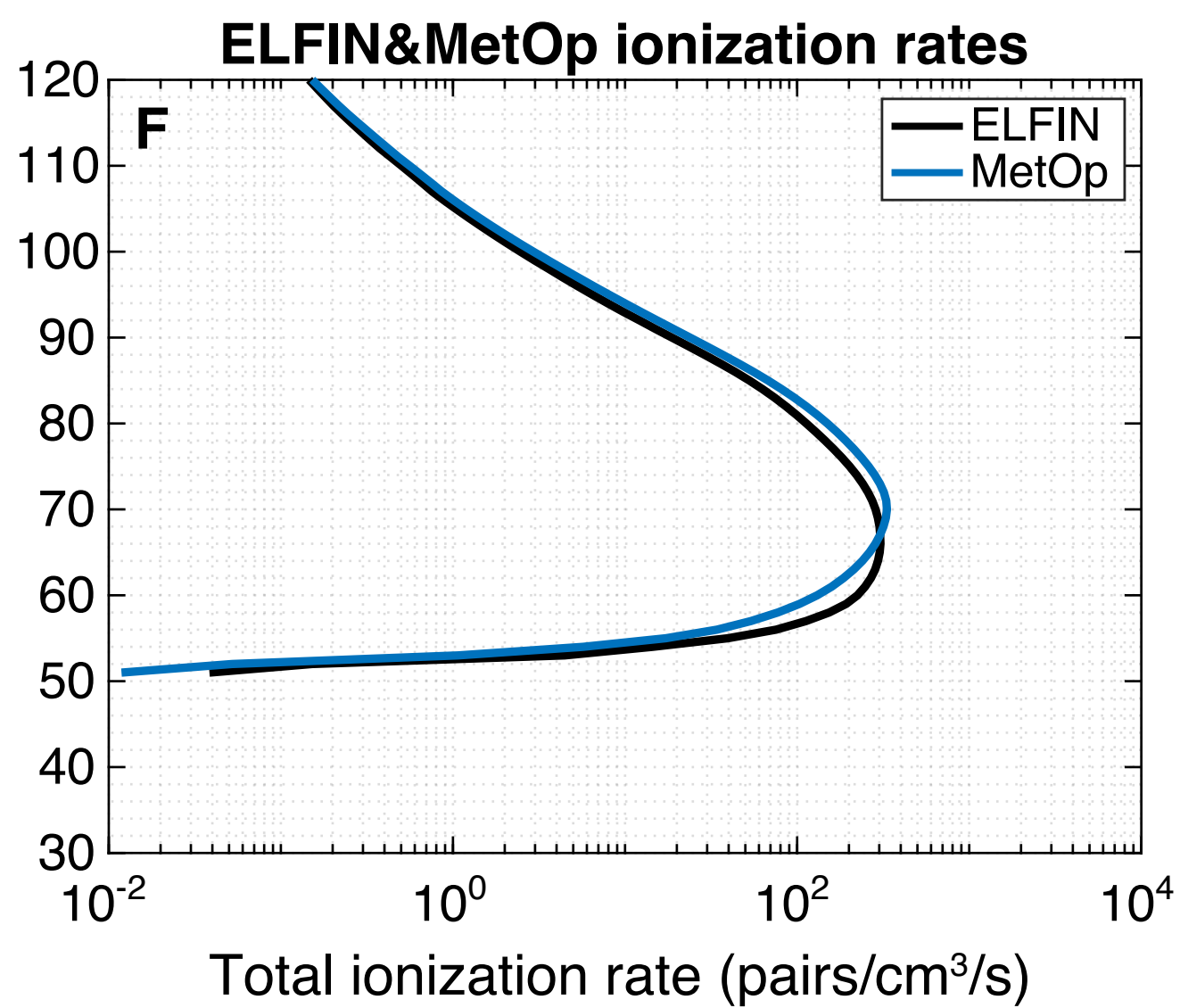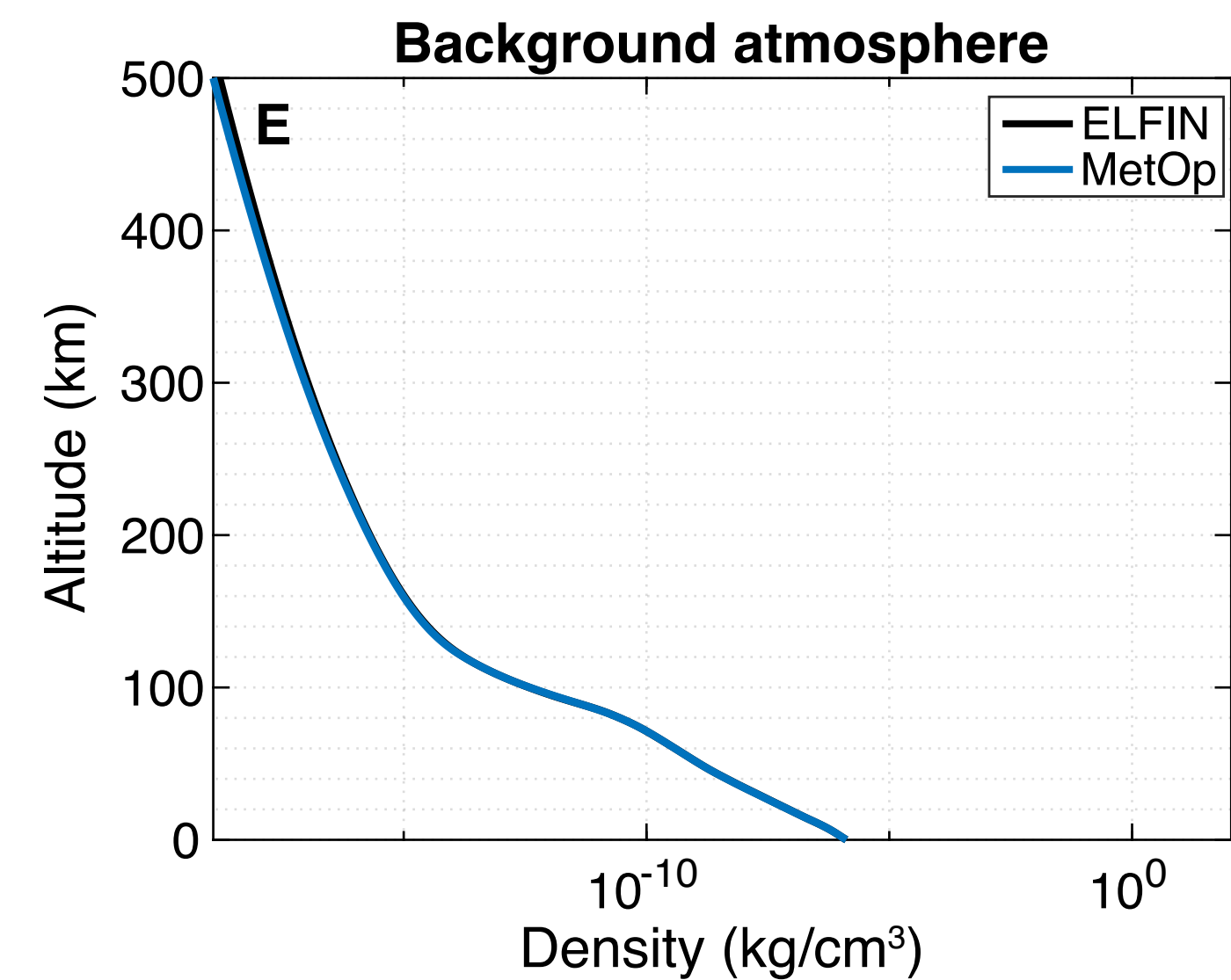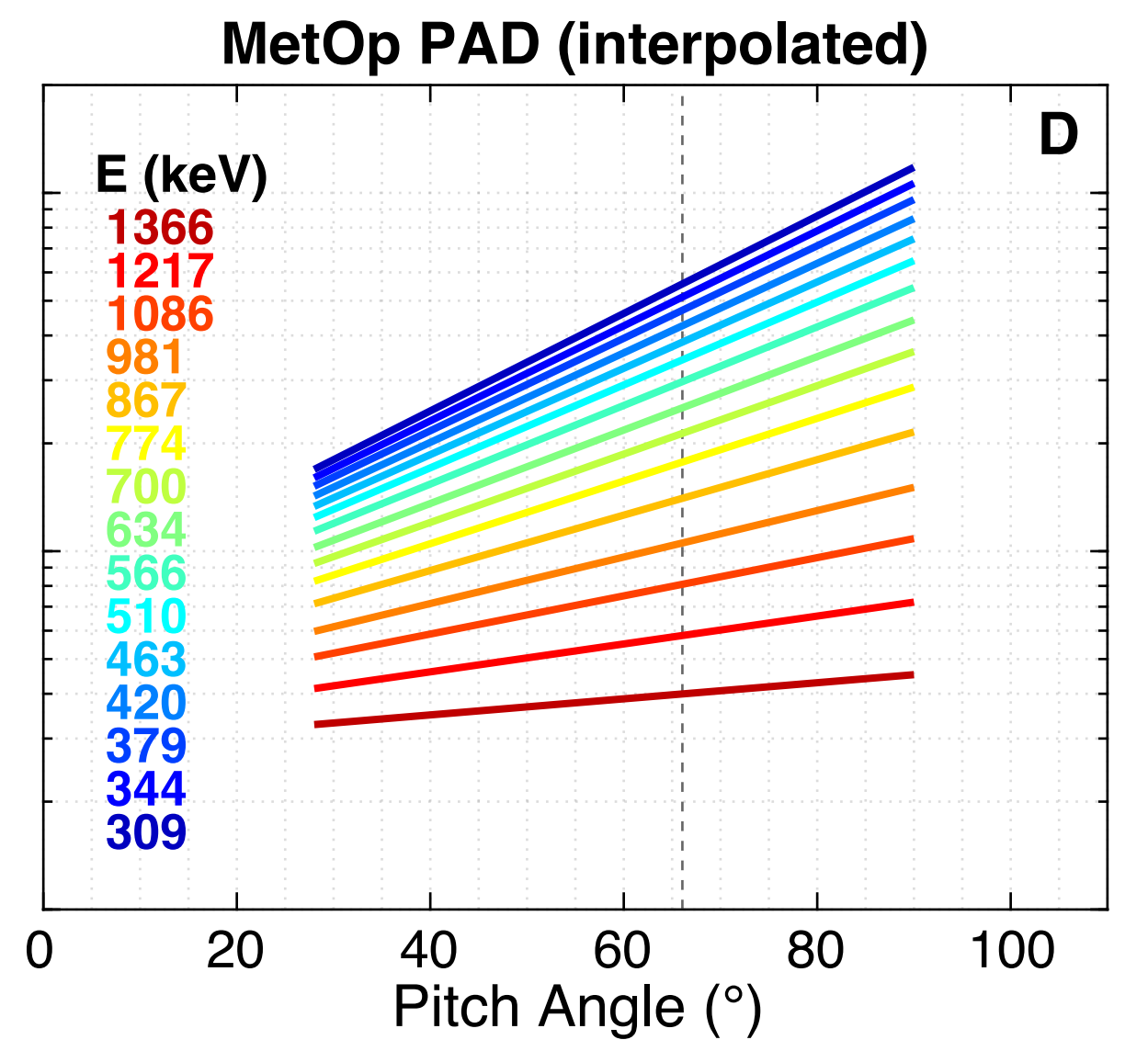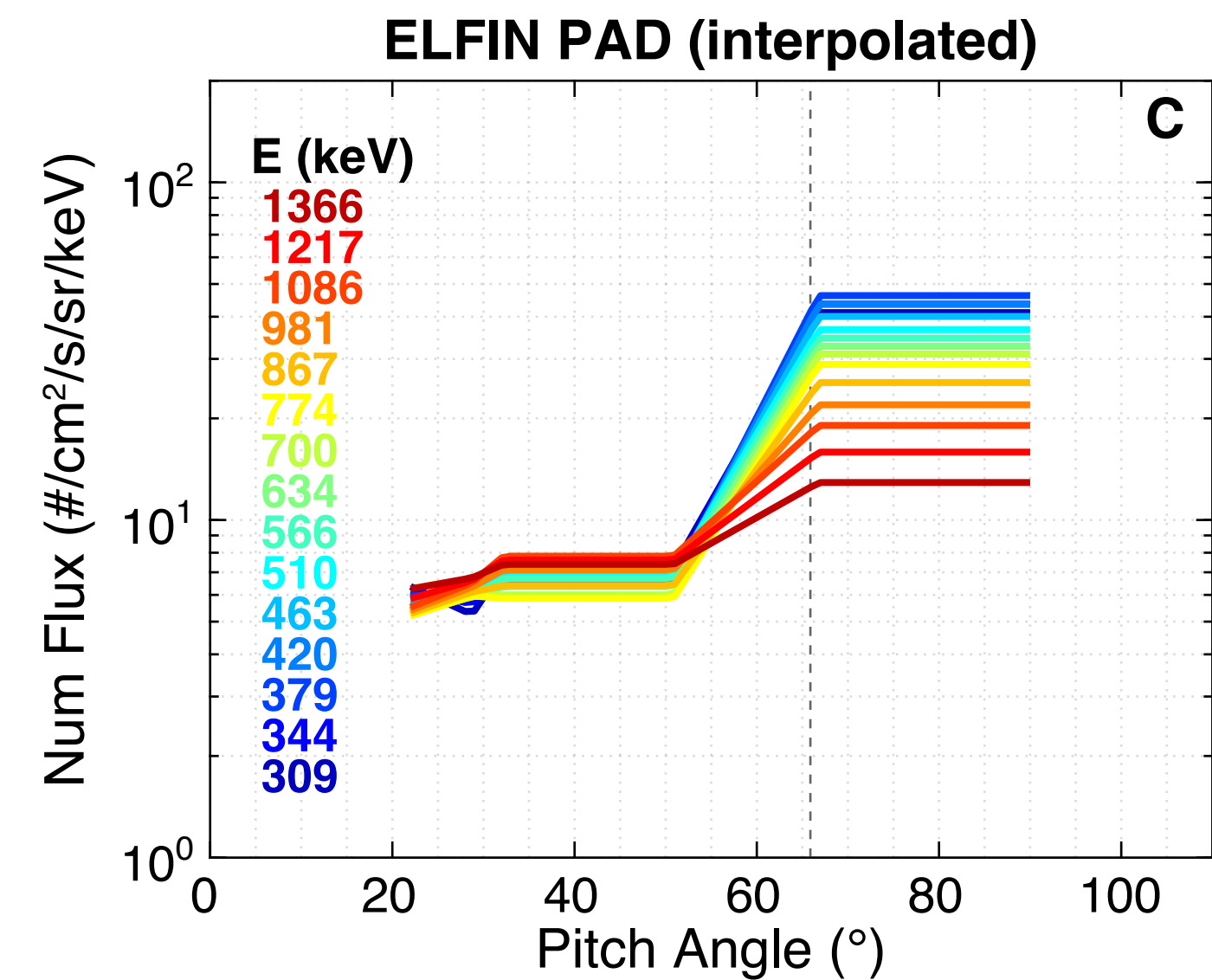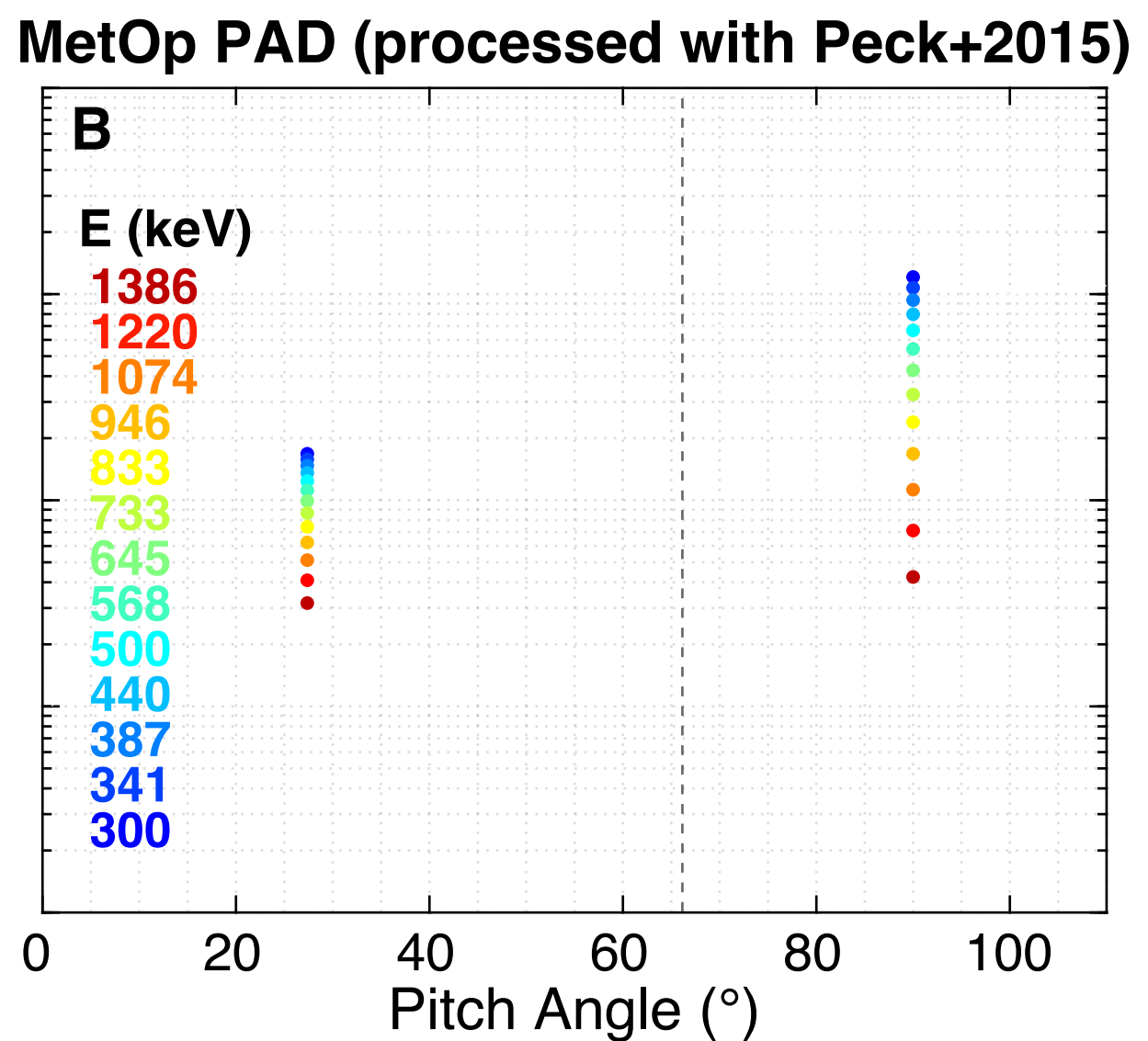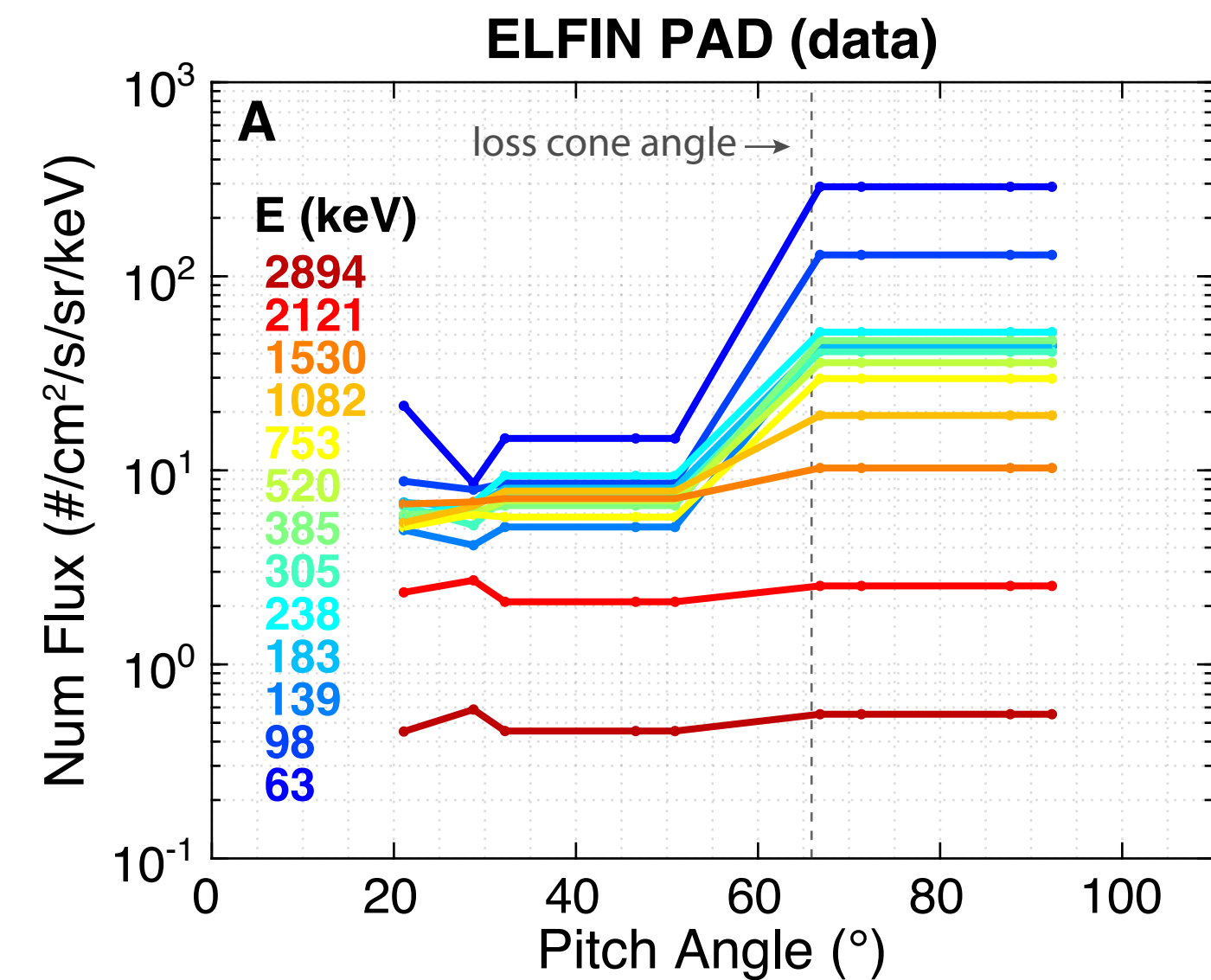
Figure 4: Dependence of the BERI ionization rates in energy (A-B) and pitch-angle (C-D). Ionization rates for a fixed electron energy (A) and pitch-angle (C). Peak of ionization rate (black) and corresponding altitude (blue) as a function of energy (B) and pitch-angle (D).
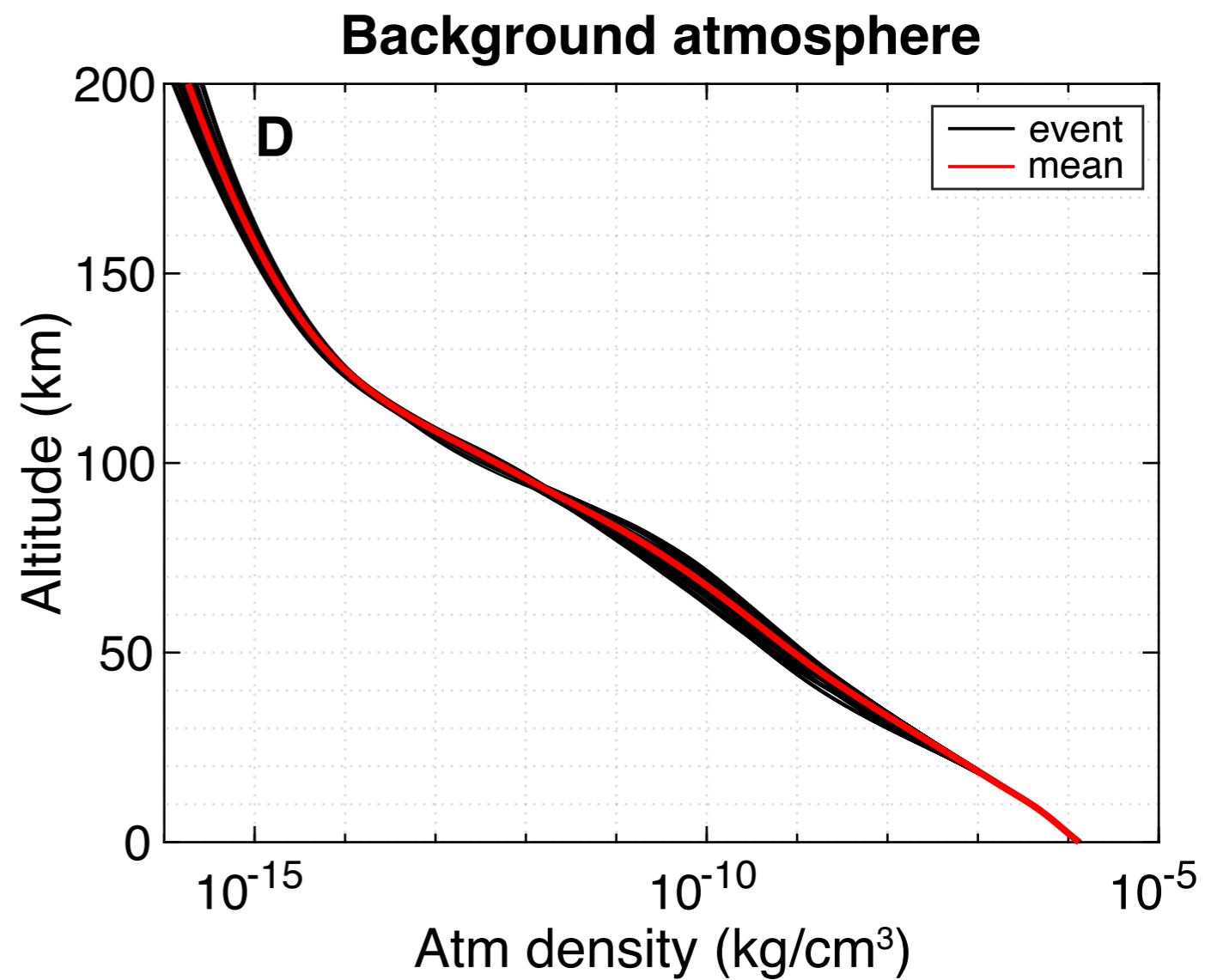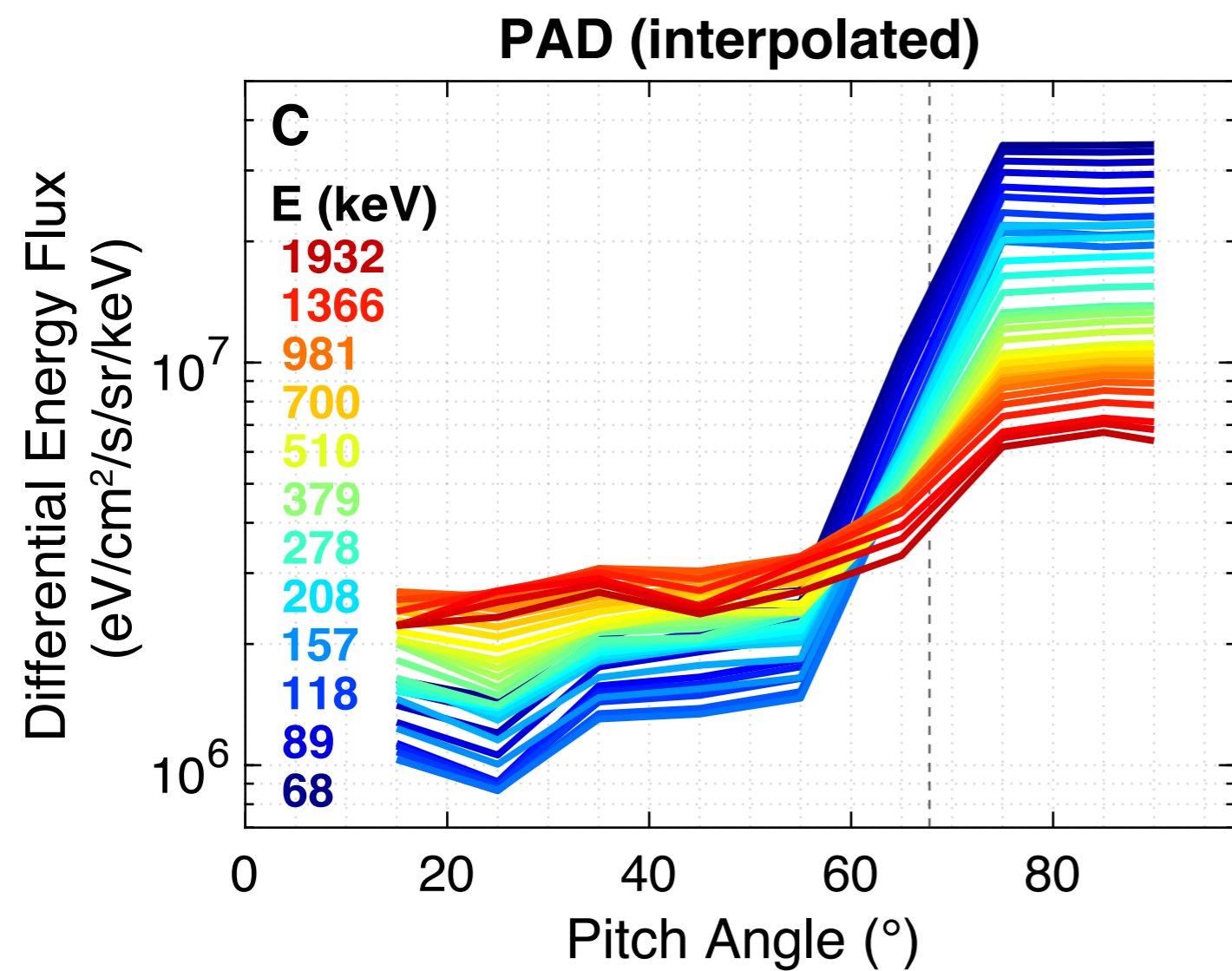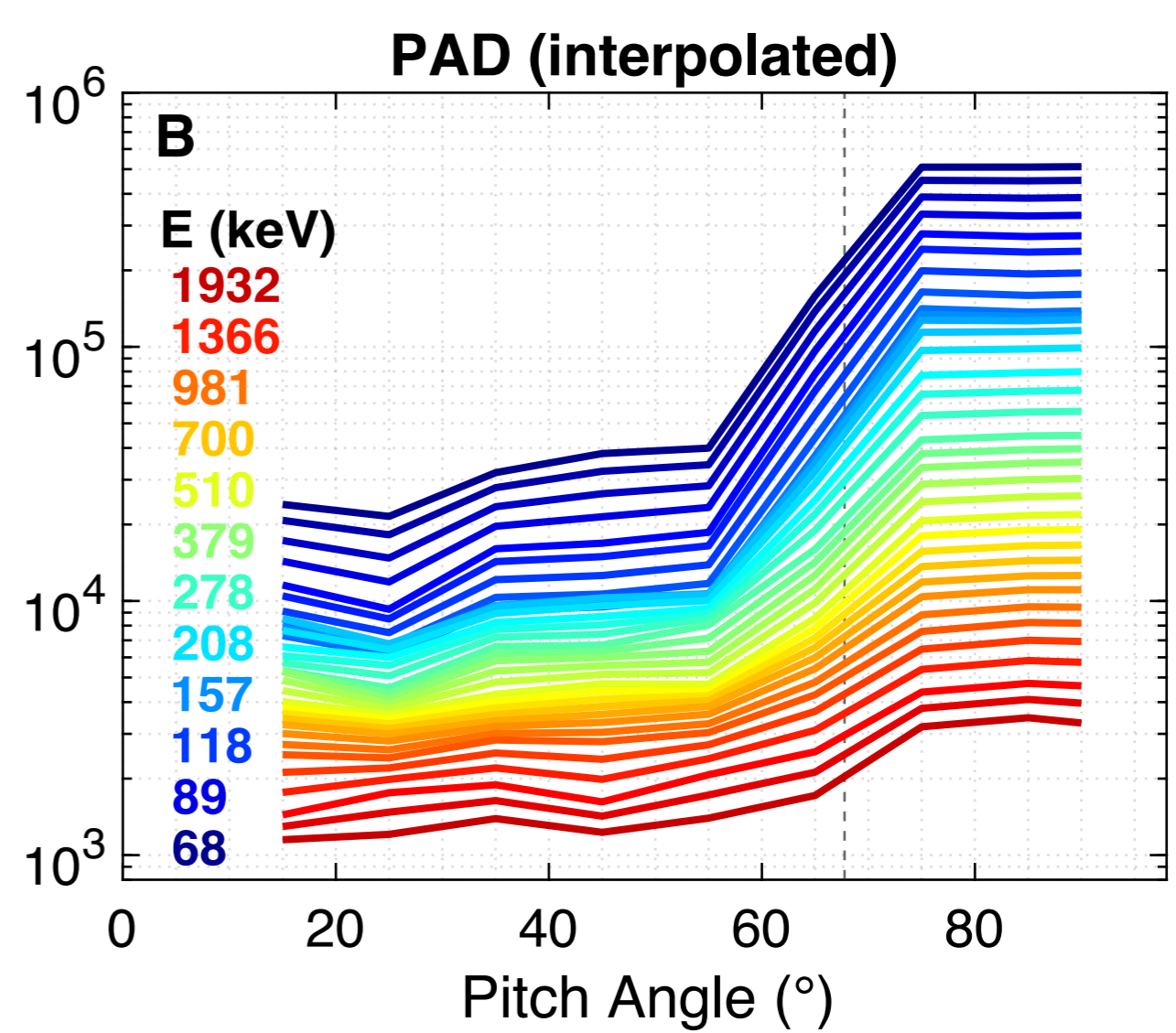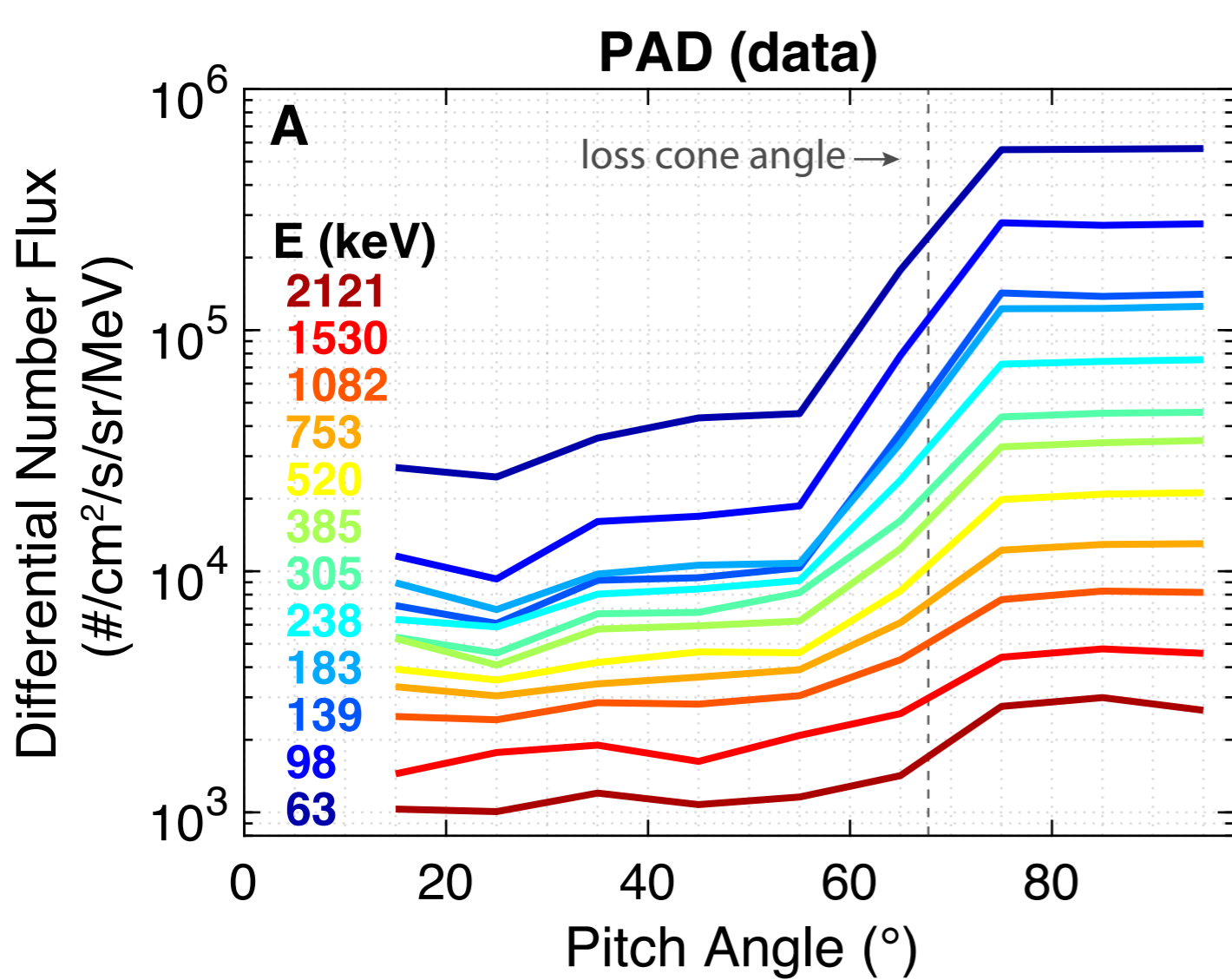
Figure 5: Overview of the ionization rates from BERI due to EMIC-driven EEP. A: range of ionization rates using a PAD as input as described in the legend (e.g., black solid line corresponds to the average PAD shown in Figure 3B; blue solid line corresponds to the result using the minimum PAD from the statistics). Energy deposition is indicated with the additional x-axes at the bottom of panel A (blue for $eV/cm^2/s/km$ units and maroon for $erg/cm^2/s/km$ units). B: daily averaged ionization rates recommended by CMIP6 compared to the EMIC-driven BERI ionization rates (black for average PAD and green for median PAD), where MEE: medium-energy electrons, SEP: solar energetic particles, GCR: galactic cosmic rays. The ionospheric E and D-regions are shaded in orange and blue, respectively. The ozone layer is shaded in green with a darker green indicating the approximate altitude of the highest mixing ratio. Atmospheric layers are also indicated.

Figure 6: Spatial distribution of the 144 EMIC-driven events in geographical coordinates. Panel A shows the input energy flux calculated for each event from ELFIN data and Panel B shows the ionization energy flux calculated from BERI's ionization rate for each event. The colorbar is in $eV/cm^2/s$ units (black) and $erg/cm^2/s$ units (blue).

1070   **Figure 7**: Cartoon of the local energy budget: orange arrow indicates the energy input due to
1071   EMIC-driven precipitation (calculated from ELFIN data in the loss cone, LC), white dots
1072   indicate the ion-electron pairs produced in the mesosphere due to ionization, with the
1073   corresponding atmospheric ionization energy flux (calculated from the BERI's ionization rate),
1074   and the pink arrow shows the backscattered energy flux (from BERI's estimate).

A

THERMOSPHERE

Auroral & Plasmasheet Precipitation
< 30 keV

Radiation Belt Precipitation
10s keV – 10 MeV

Solar Energetic Particles (SEP)
10s – 100s MeV

Galactic Cosmic Rays (GCR)
100s MeV – 100s GeV

MESOSPHERE

NO$_x$
HO$_x$

enhanced ionization

STRATOSPHERE

TROPOSPHERE

EARTH

B

ELFIN

MetOp

electron precipitation

magnetic conjunction

B

EMIC waves

e⁻
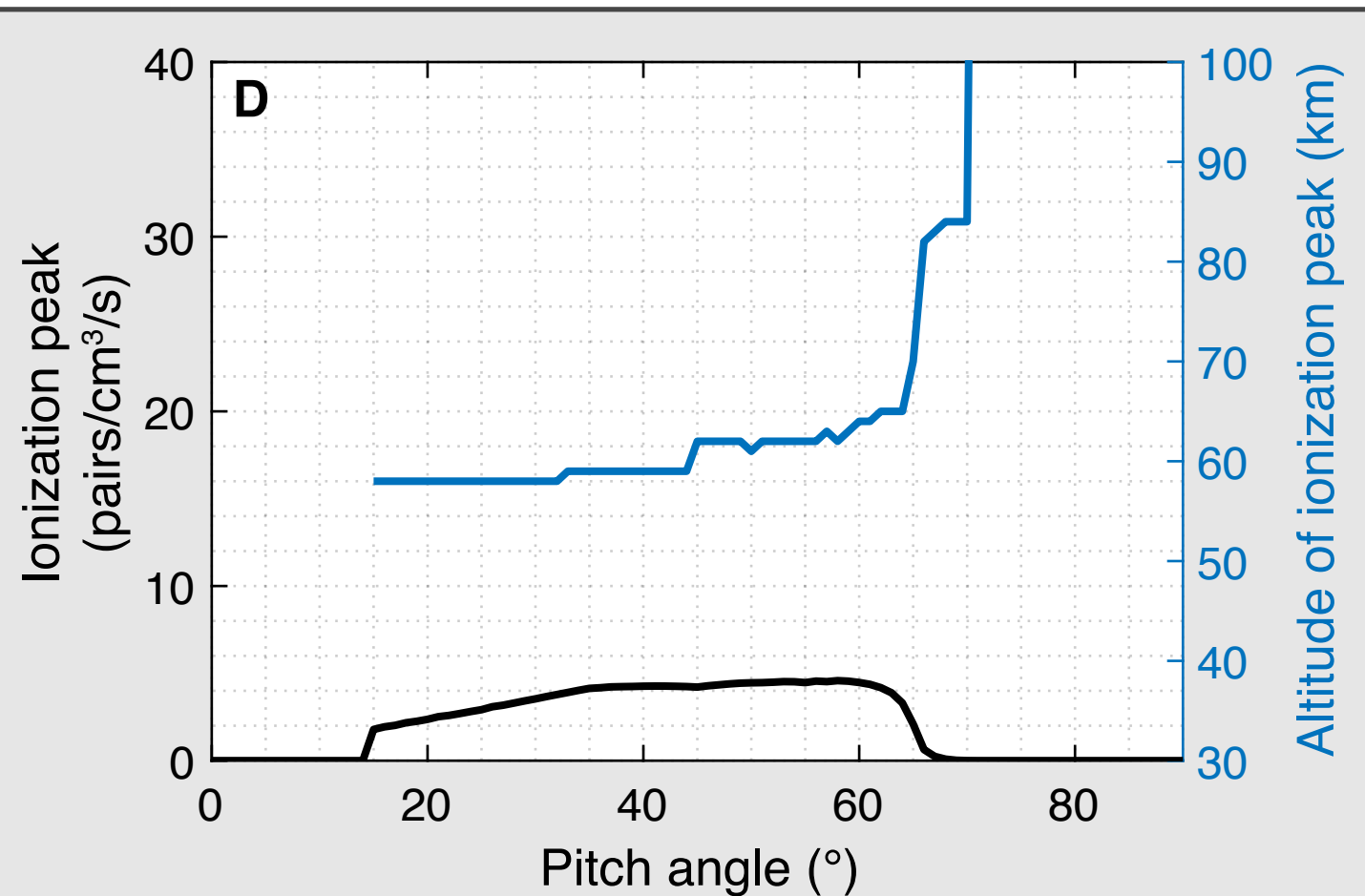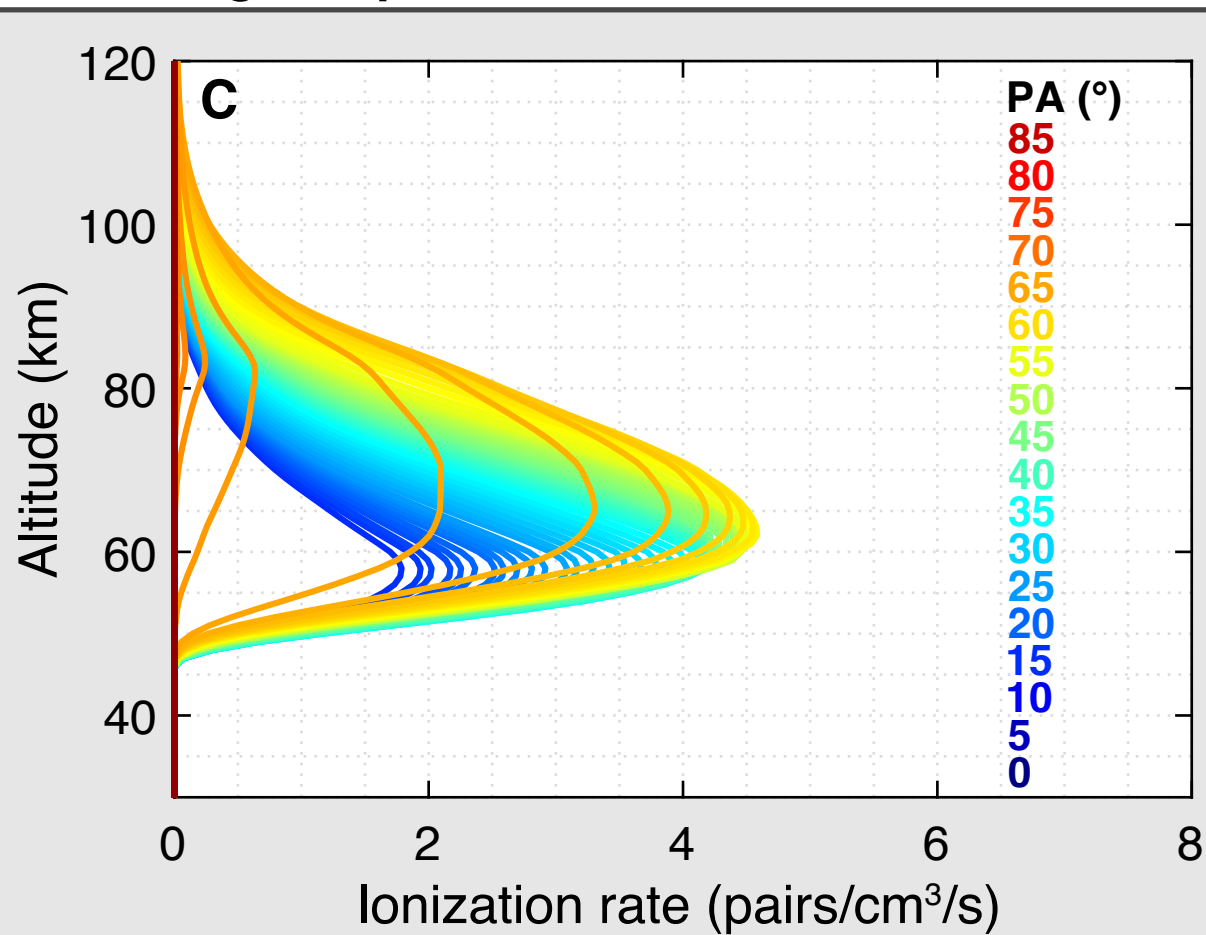
magnetic field line

outer radiation belt

ELFIN-A

C 345–630 keV

precipitation

precipitating

quasi-trapped

backscattered

D 630–900 keV

loss cone angle

anti-loss cone angle

Pitch angle [°]

Electron Flux [s⁻¹ cm⁻² sr⁻¹ MeV⁻¹]

E 0.9–7 MeV

L	3.6	3.8	4.1	4.5	4.8	5.2	5.7	6.3	6.9
MLT	19.4	19.4	19.4	19.4	19.4	19.4	19.4	19.4	19.4
Seconds	20	40	00	20	40	00	20	40	00
2021 Feb 02			0159			0200			0201

MetOp-02

F

precipitation

P1: 30–80 keV
P2: 80–240 keV
P3: 240–800 keV

Proton flux [s⁻¹ cm⁻² sr⁻¹ keV⁻¹]

G

precipitating (0° tel)

quasi-trapped (90° tel)

E3: >300 keV
E4: >700 keV

Electron Flux [s⁻¹ cm⁻² sr⁻¹]

L	5.1	4.6	4.2	3.9
MLT	21.8	21.7	21.6	21.5
Seconds	30	00	30	00
2021 Feb 02		0205		0206

A ELFIN PAD (data)

loss cone angle →

E (keV)
2894
2121
1530
1082
753
520
385
305
238
183
139
98
63

B MetOp PAD (processed with Peck+2015)

E (keV)
1386
1220
1074
946
833
733
645
568
500
440
387
341
300

C ELFIN PAD (interpolated)

E (keV)
1366
1217
1086
981
867
774
700
634
566
510
463
420
379
344
309

D MetOp PAD (interpolated)

E (keV)
1366
1217
1086
981
867
774
700
634
566
510
463
420
379
344
309

E Background atmosphere

— ELFIN
— MetOp

F ELFIN&MetOp ionization rates

— ELFIN
— MetOp

A PAD (data)

loss cone angle →

E (keV)
2121
1530
1082
753
520
385
305
238
183
139
98
63

Differential Number Flux (#/cm²/s/sr/MeV)

Pitch Angle (°)

B PAD (interpolated)

E (keV)
1932
1366
981
700
510
379
278
208
157
118
89
68

Pitch Angle (°)

C PAD (interpolated)

E (keV)
1932
1366
981
700
510
379
278
208
157
118
89
68

Differential Energy Flux (eV/cm²/s/sr/keV)

Pitch Angle (°)

D Background atmosphere

— event
— mean

Altitude (km)

Atm density (kg/cm³)

Energy dependence

Pitch-angle dependence

Panel A — Range of EMIC-driven ionization rates

Legend:
- avg
- median
- 25th
- 75th
- min
- max
- avg+SE
- avg−SE

Ionization rate (pairs/cm³/s)

Energy deposition

(eV/cm²/s/km)

(erg/cm²/s/km)

Panel B — Comparison with CMIP6

MEE, SEP, BERI avg, BERI median, GCR

E region, D region, Ozone layer

THERMOSPHERE, MESOSPHERE, STRATOSPHERE, TROPOSPHERE

Altitude (km)

Pressure (hPa)

Ionization rate (pairs/cm³/s)

EMIC-driven (LC) INPUT
~0.033 $\frac{erg}{cm^2 \, s}$
||from ELFIN||

Backscattered OUTPUT
~0.008 $\frac{erg}{cm^2 \, s}$
||from BERI||

i+   e−

~0.025 $\frac{erg}{cm^2 \, s}$ ionizes mesosphere
(~74% of input energy flux)
||from BERI||

STRATOSPHERE

MESOSPHERE